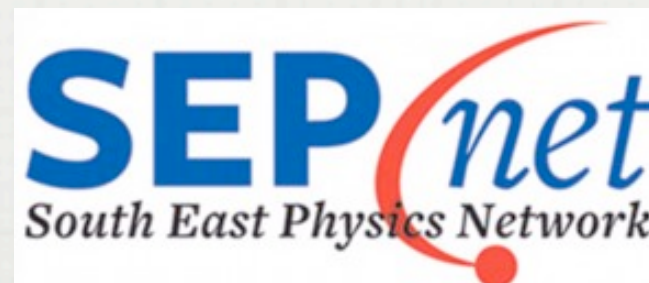


BSM PARTICLE AND ASTRO-PARTICLE THEORY

STEPHEN WEST

ROYAL HOLLOWAY, UNIVERSITY OF LONDON



IOP JOINT HEPP AND APP

APRIL 2012

IOP joint HEPP and APP 2012

INTERESTING TIME IN ASTRO-PARTICLE PHYSICS

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□ MY FOCUS IS - DARK MATTER PHYSICS

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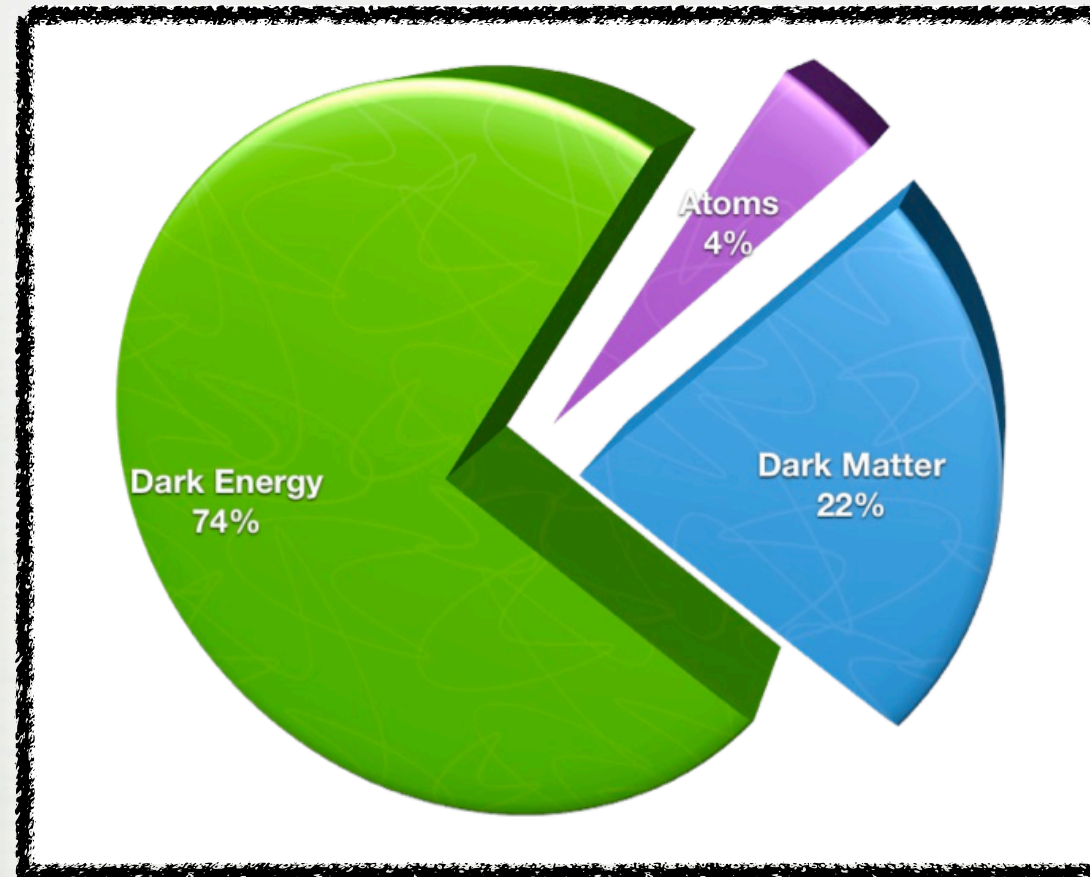
□ STRUCTURE FORMATION

□ EDM MEASUREMENTS

□ ...

INTRODUCTION

□ COMPOSITION OF THE UNIVERSE:

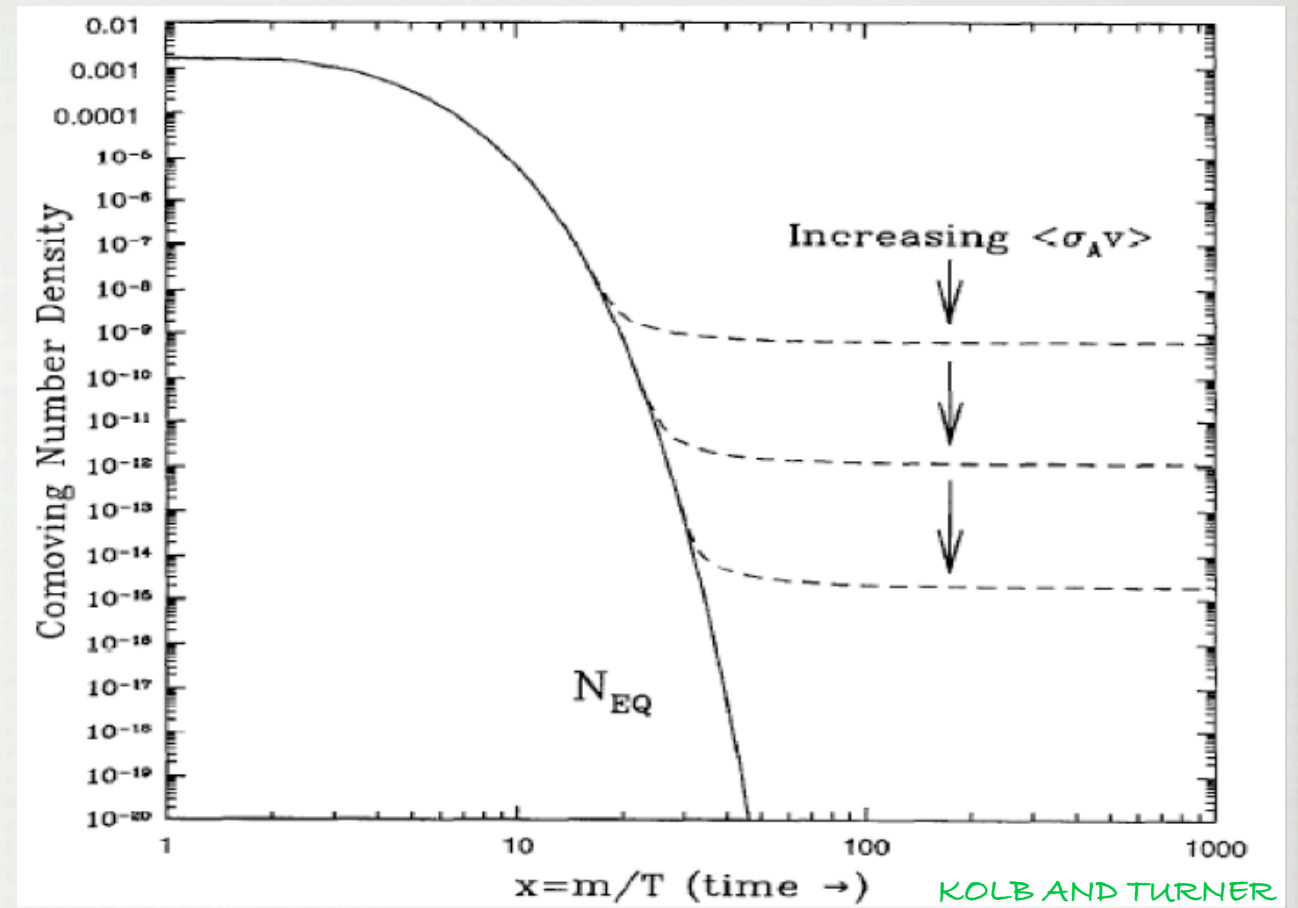
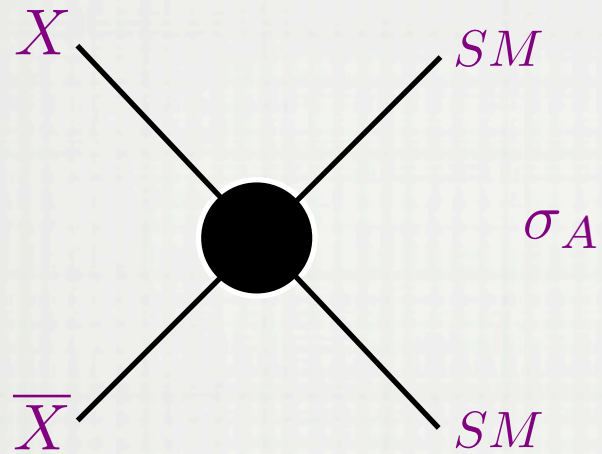


□ DARK MATTER IS A MAJOR COMPONENT - DEMANDS INVESTIGATION

□ MUCH WORK ALREADY DONE...MUCH DONE IN A FAMILIAR FRAMEWORK

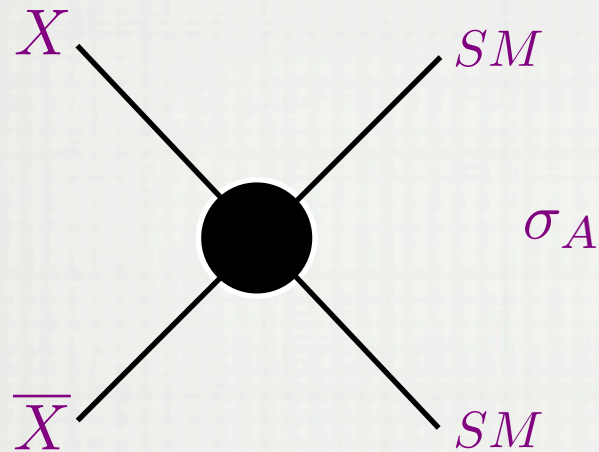
THE "STANDARD PICTURE" - FREEZE-OUT

□ DARK MATTER X INITIALLY IN THERMAL EQUILIBRIUM $T > m_X$

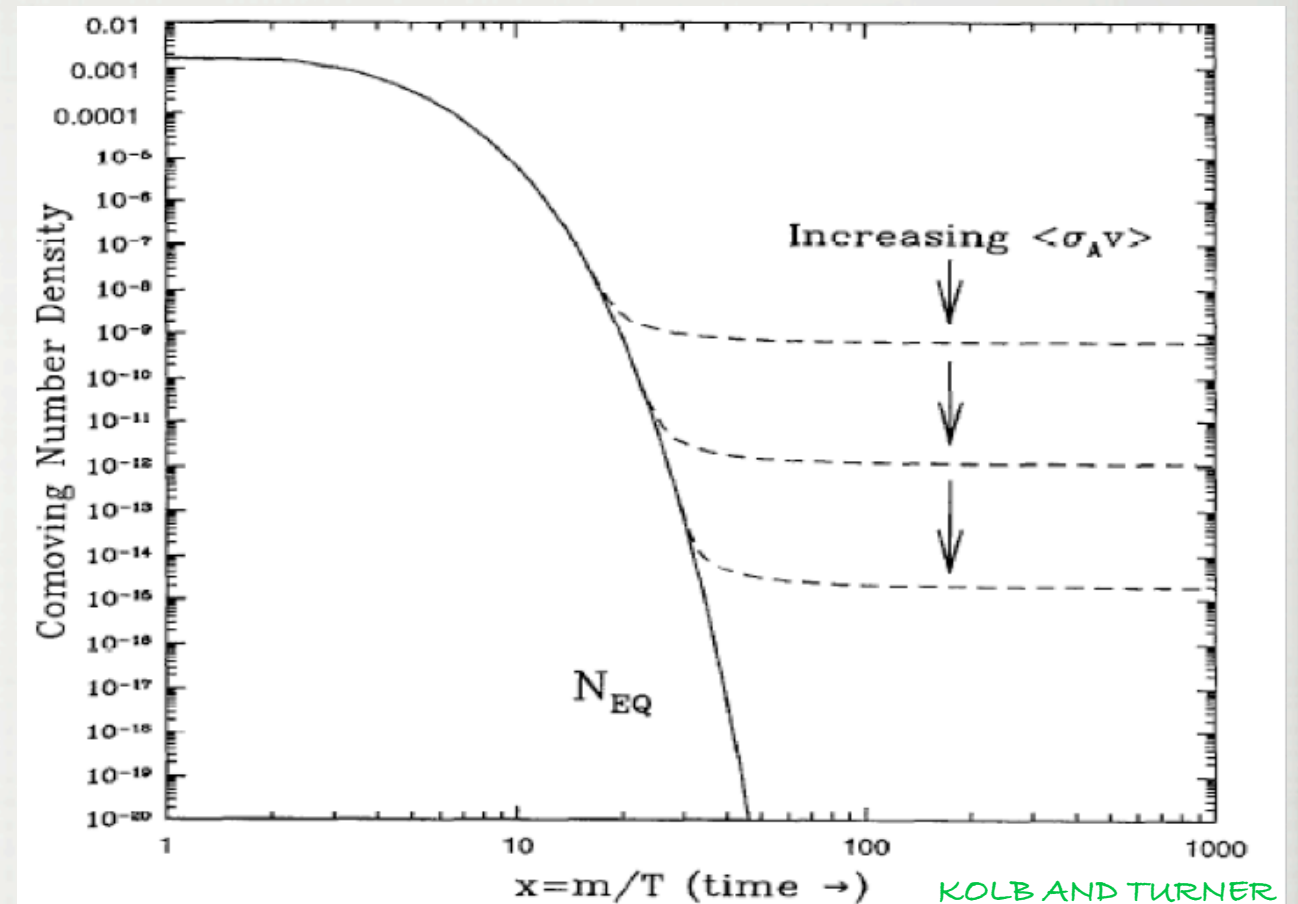


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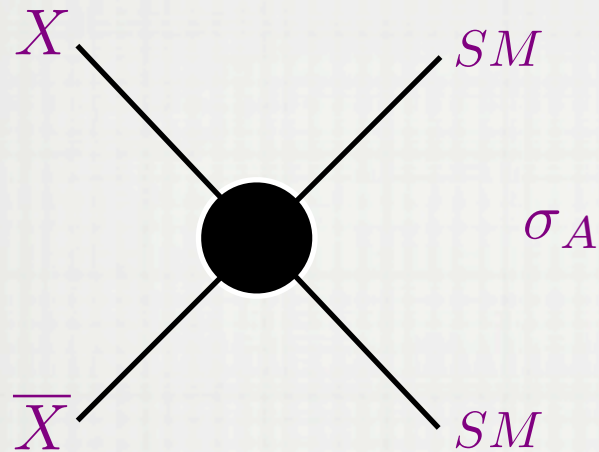


□ AS THE TEMP DECREASES $T < m_X$ CREATION OF X BECOMES EXPONENTIALLY SUPPRESSED

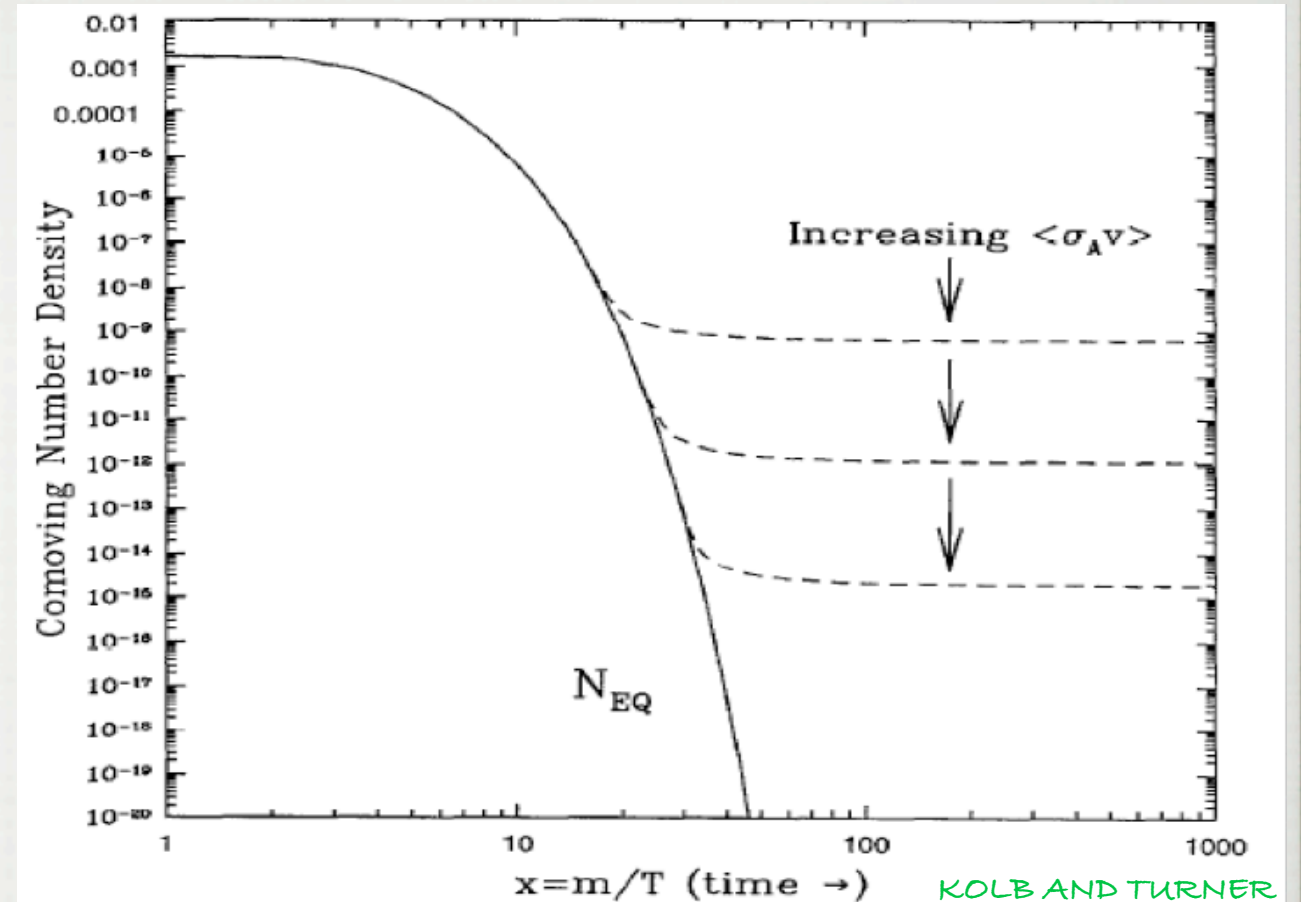


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□ ANNIHILATION OF X STILL PROCEEDS, NUMBER DENSITY OF X GIVEN BY

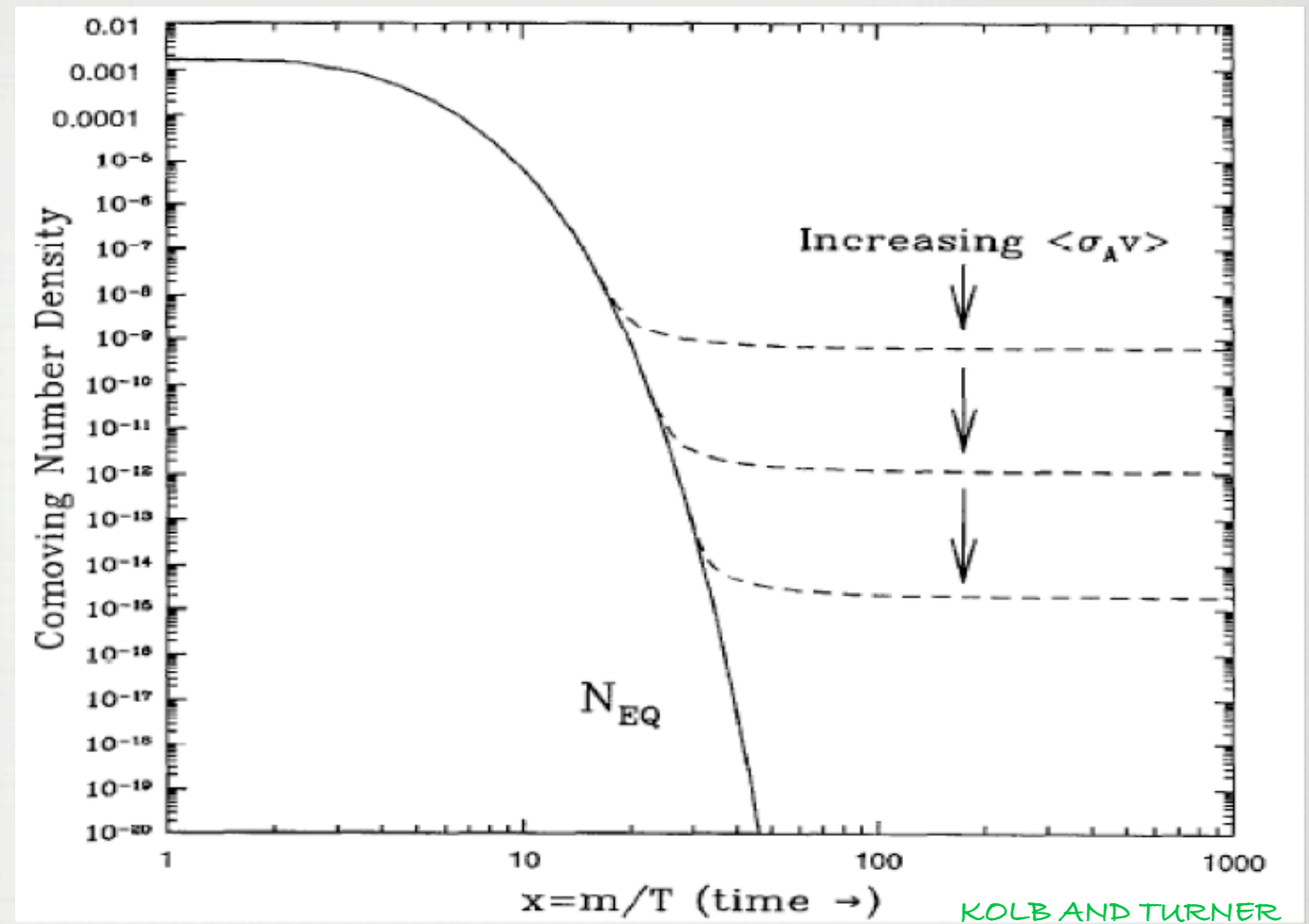
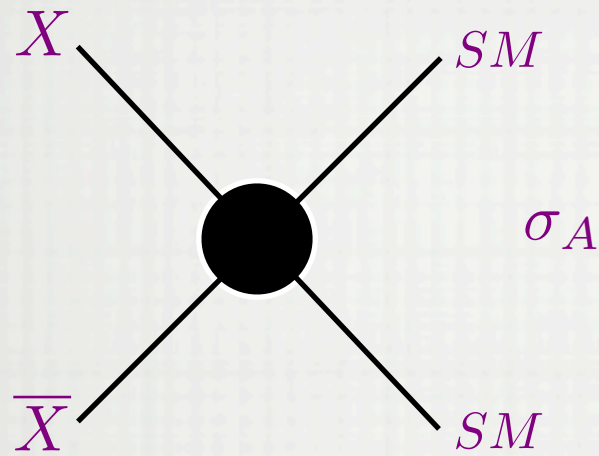
$$n_{X,eq} \approx g_X \left(\frac{m_X T}{2\pi} \right)^{3/2} e^{-m_X/T}$$

$$n_{X,eq} \rightarrow 0$$

$$\text{as } T \rightarrow 0$$

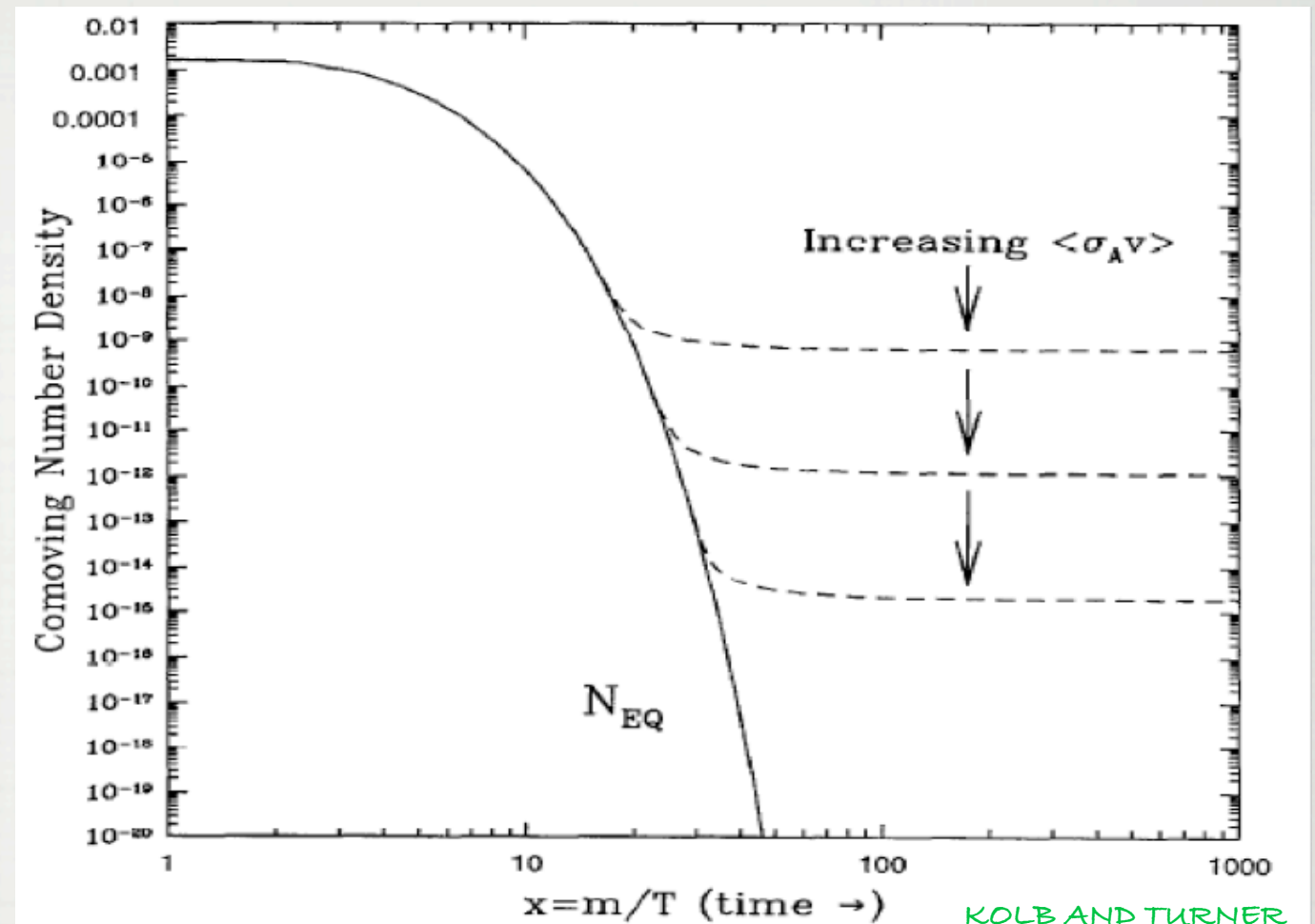
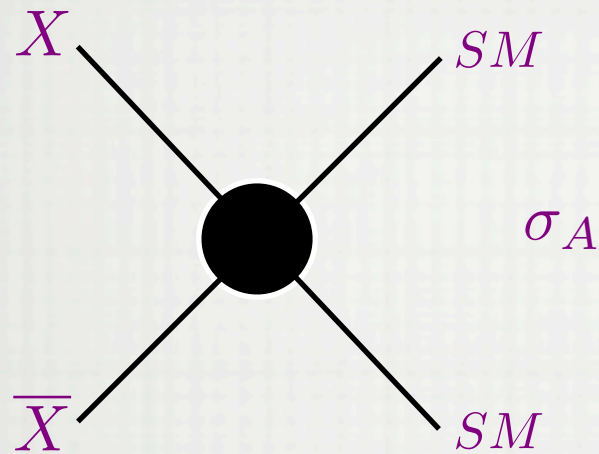
□ DUE TO EXPANSION, DARK MATTER
NUMBER DENSITY FREEZES-OUT

WHEN: $\Gamma = n_X \langle \sigma_A v \rangle < H$



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WHEN: $\Gamma = n_X \langle \sigma_{AV} \rangle < H$

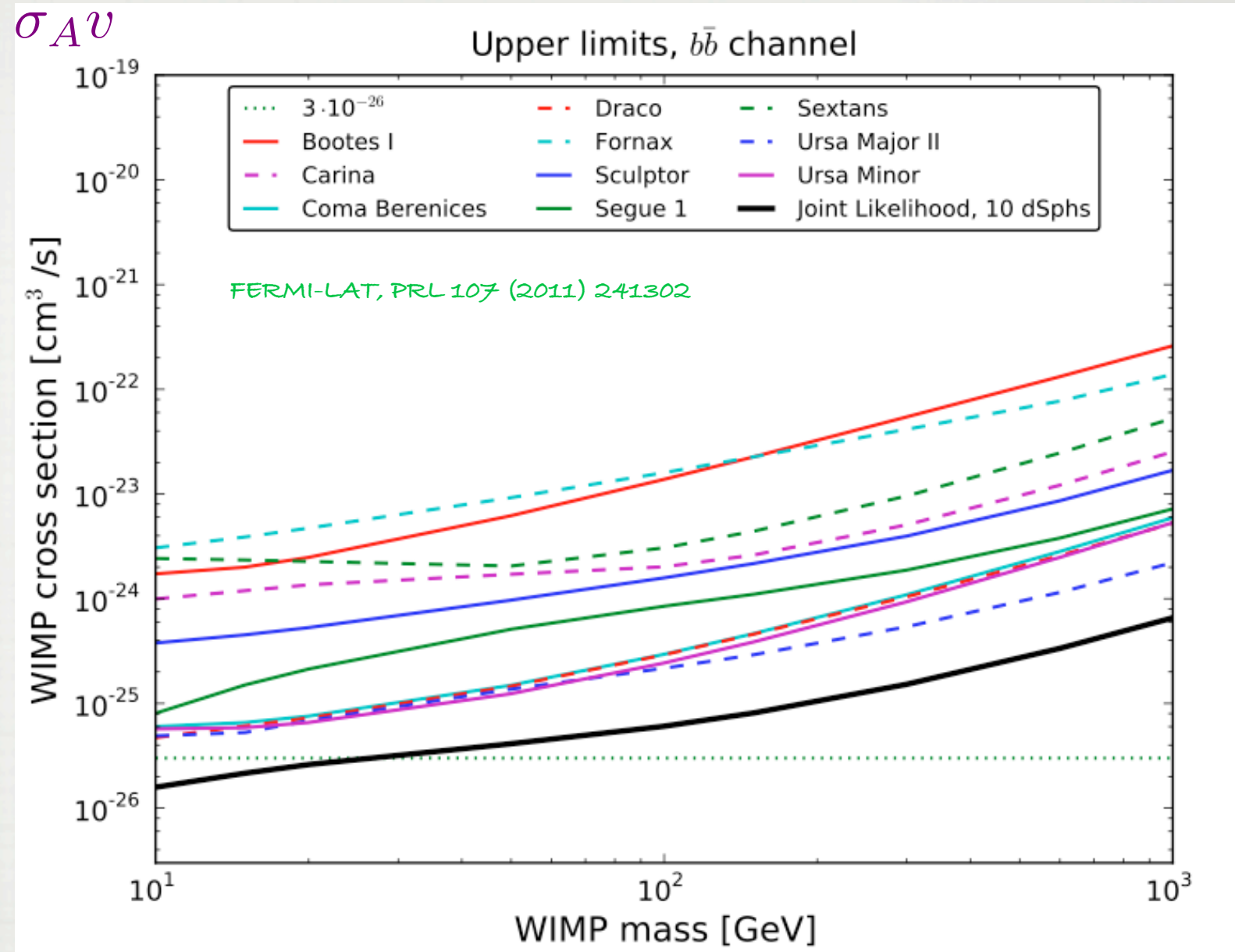


□ FINAL ABUNDANCE:

$$\Omega h^2 \sim 0.1 \frac{3 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}}{\langle \sigma_{AV} \rangle}$$

APPROX. WEAK SCALE CROSS SECTION - WIMPS

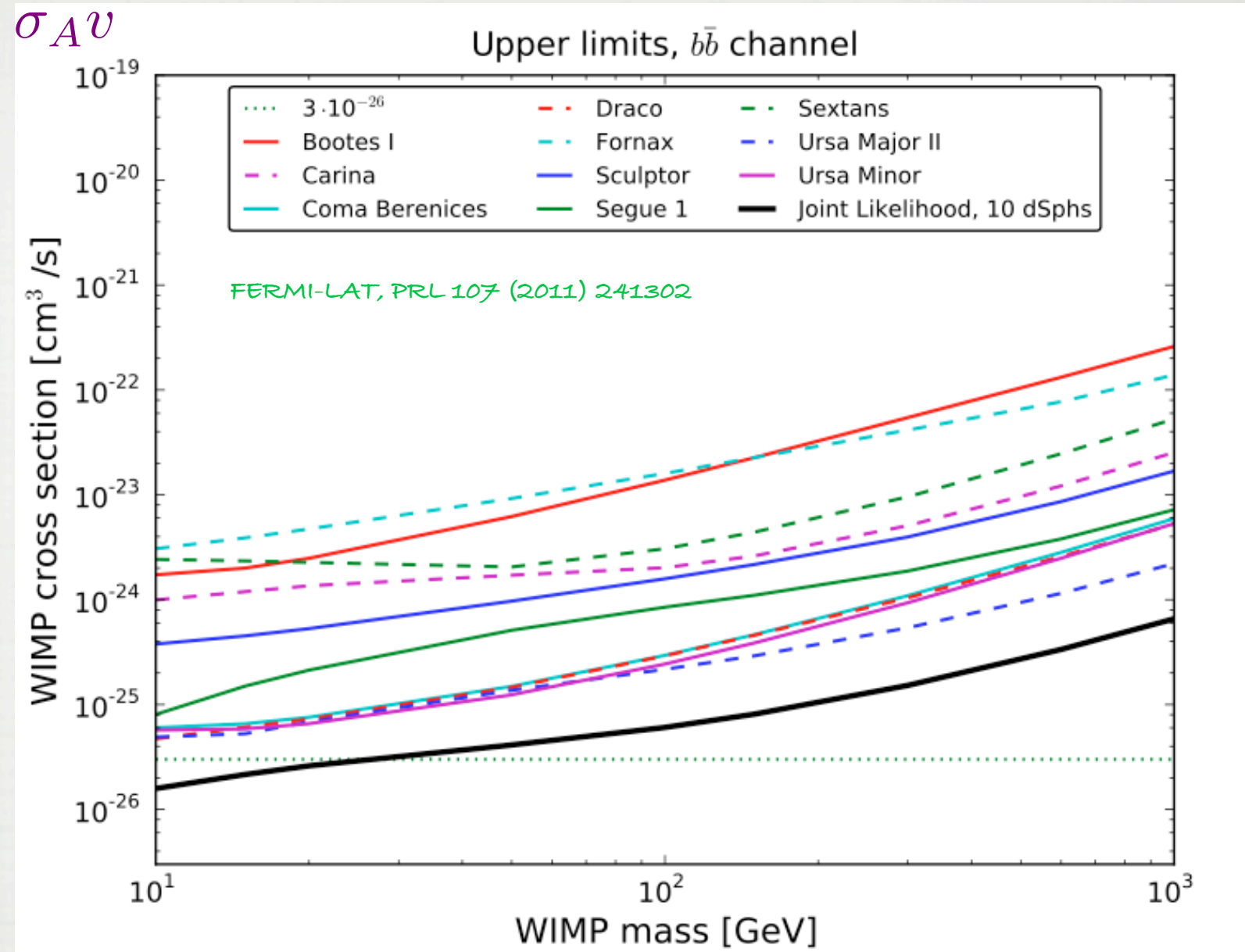
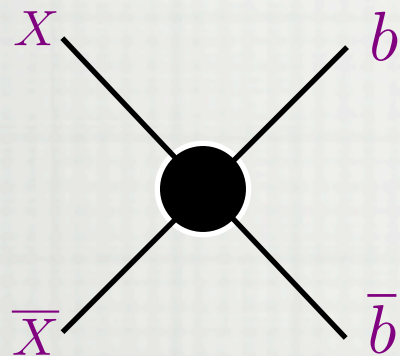
INDIRECT DETECTION IN STANDARD PICTURE



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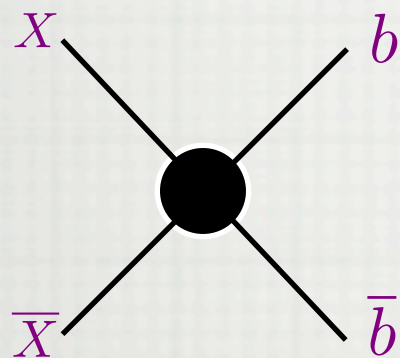
□ PUTS LIMITS ON THE **WIMP ANNIHILATION CROSS SECTION**



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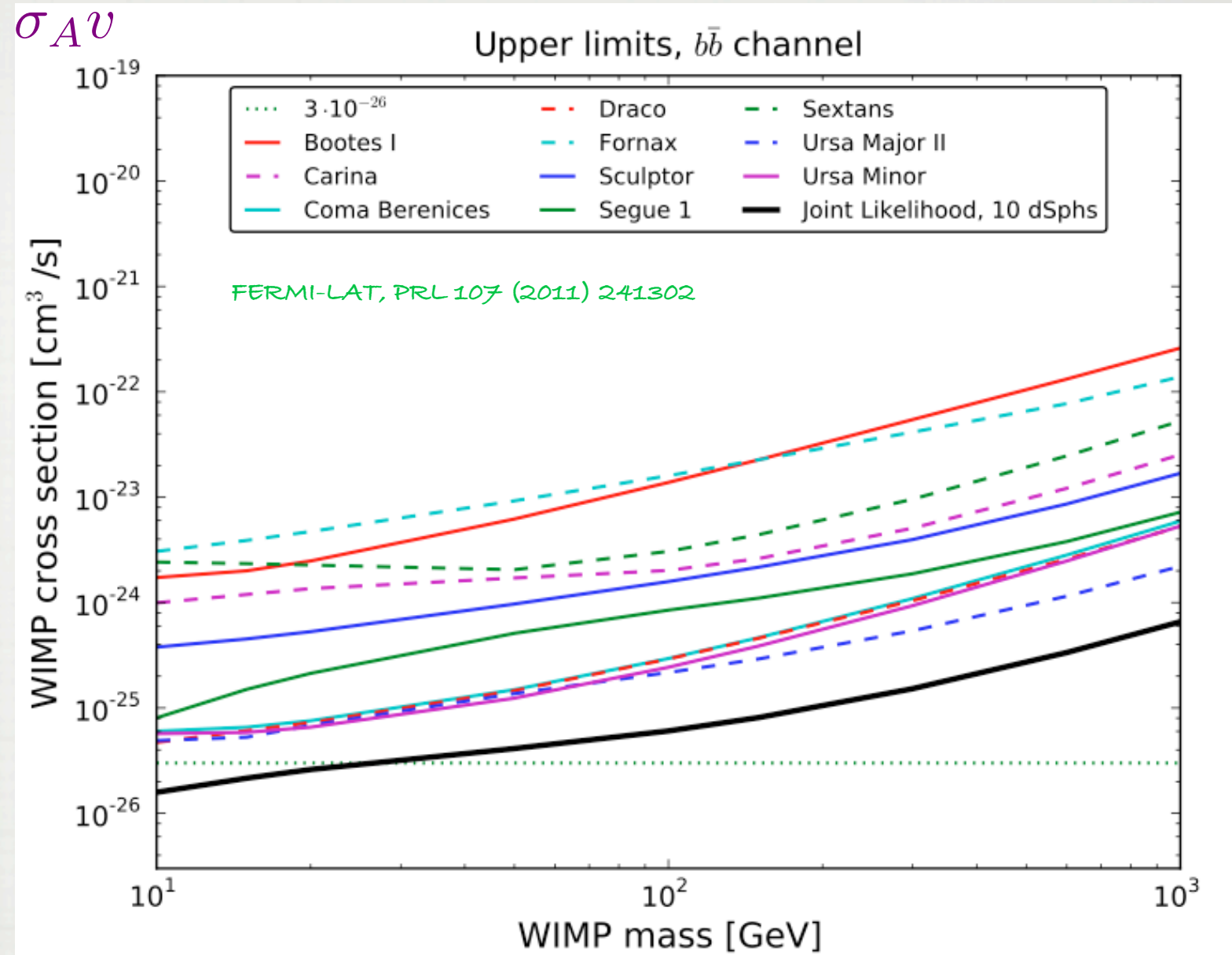
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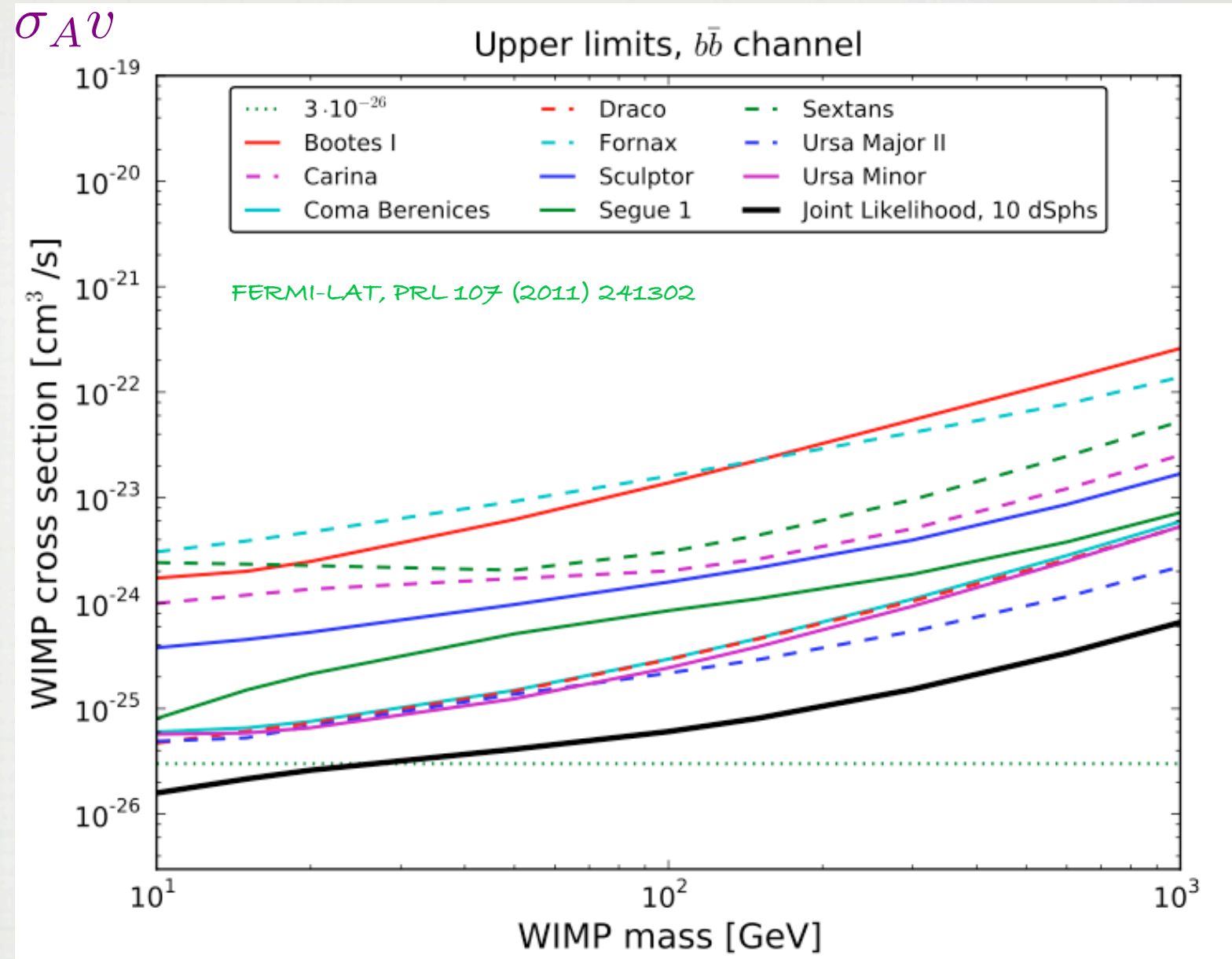
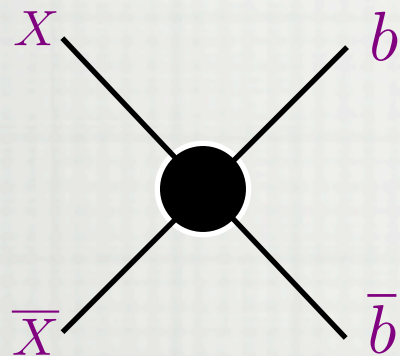
$$\sigma_{AV} \approx a + bv^2 + \dots$$



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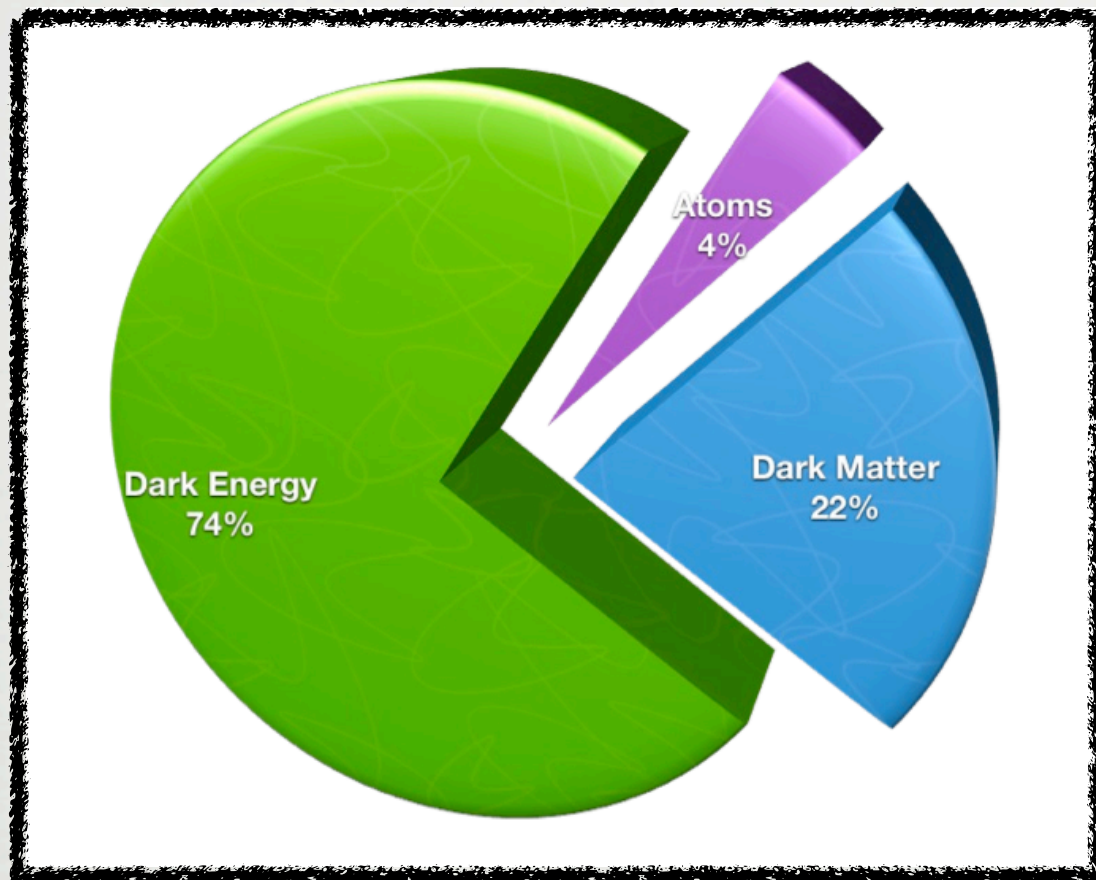


□ PROVIDED THE ANNIHILATION IS S-WAVE $\sigma_{AV} \approx a + bv^2 + \dots$

□ IMPROVING ALL THE TIME, STARTING TO IMPINGE ON DM MODELS...SEE PREVIOUS TALK!

ASYMMETRIC DM MOTIVATION

□ BACK TO THE COMPOSITION OF THE UNIVERSE:



□ USUALLY THESE TWO NUMBERS ARE DETERMINED BY INDEPENDENT DYNAMICS

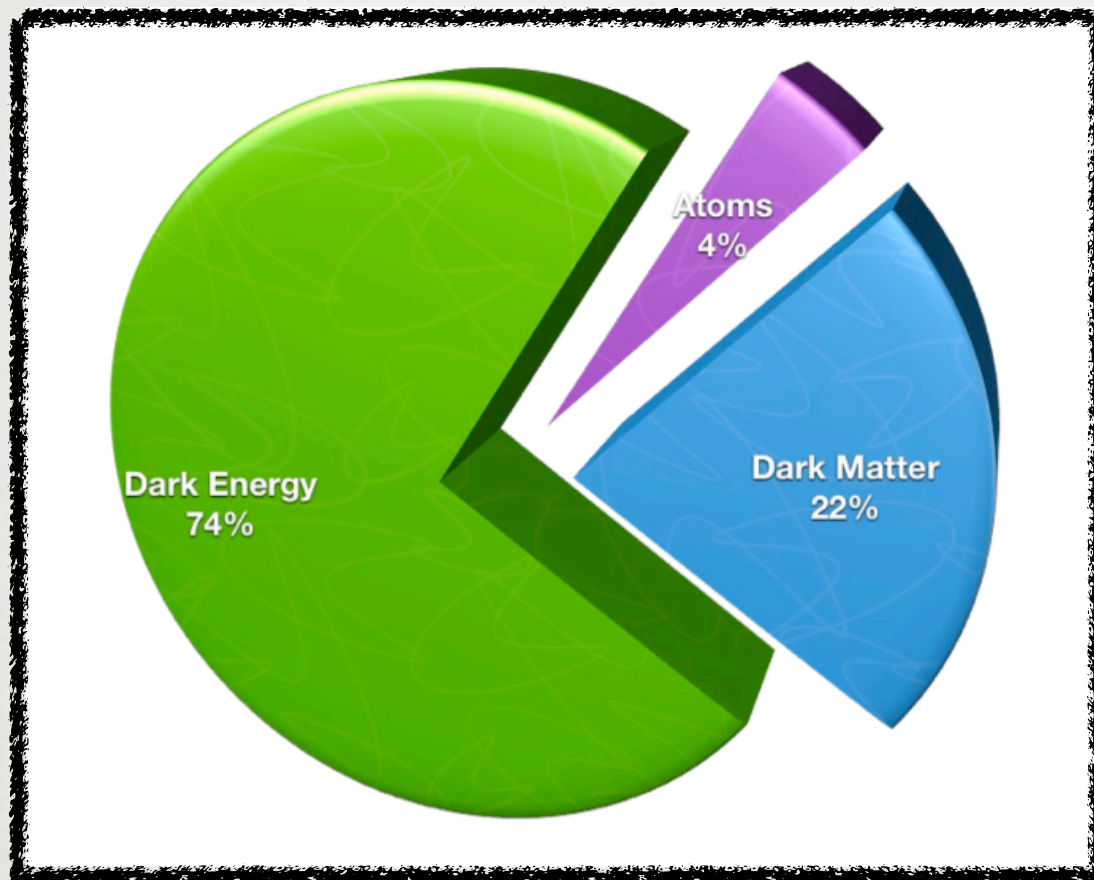
Ω_{dm} BY FREEZE-OUT OF WIMPS

Ω_b BY BARYOGENESIS/LEPTOGENESIS

$$\frac{\Omega_{dm}}{\Omega_b} \sim 5$$

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□ TAKE SERIOUSLY THE CLOSENESS OF THESE VALUES - INVESTIGATE DYNAMICS THAT LINK THE TWO...

...LEADS TO IDEAS OF ASYMMETRIC DM

ASYMMETRIC DARK MATTER BASICS

□ INTRODUCE AN ASYMMETRY IN DM NUMBER DENSITY (OR THE BARYON SECTOR)

$$\eta_{dm} = n_{dm} - n_{\overline{dm}} \neq 0$$

OR

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- THE VALUE OF C DEPENDS ON THE DETAILS OF THE DYNAMICS CONNECTING DM AND BARYONS...A SHARED SYMMETRY

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- CANDIDATES: COMPLEX SCALARS AND DIRAC FERMIONS (+USUAL REQUIREMENTS FOR DM, NO EM OR COLOUR CHARGE ETC)
- CANNOT USE MAJORANA

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- NEED A SHARED QUANTUM NUMBER, E.G. A CHARGE ASSOCIATED WITH A GLOBAL $U(1)$, RELATED TO B-L NUMBER

ASYMMETRIC DARK MATTER

A (PARTIAL) HISTORY

□ 80'S AND 90'S

COSMIONS AS ~ 5 GEV ADM - SOLUTION TO SOLAR NEUTRINO PROBLEM:

GELMINI, HALL, LIN (1987); GIUDICE, RABY (1990)

WEAK SCALE ADM: NUSSINOV (1985); BARR, CHIVUKULA, FARHI (1990), BARR (1991); DB KAPLAN (1992); THOMAS (1995);

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□ 00'S

WEAK SCALE ADM: FUJII, YANAGIDA (2002); FARRAR, ZAHARIJAS (2004), HOOPER, MARCH-RUSSELL, SMW (2004); KITANO, LOW (2004); AGASHE, SERVANT (2004); TYTGAT (2006).

□ MANY RECENT DEVELOPMENTS - LOTS OF OTHERS:

~ 5 GEV OR TEV ADM

MURAYAMA, RATZ, KAPLAN (DE), LUTY, ZUREK, COHEN, CAI, FRANDSEN, SARKAR, SCHMIDT-HOBERG, PHALEN, SANNINO, DAVOUDIASHVILI, MORRISSEY, SIGURDSEN, TULIN, HABA, MATSUMOTO, BUCKLEY, RANDALL, CHUN, GU, LINDNER, SARKAR, ZHANG, BLENNOW, DASGUPTA, FERNANDEZ-MARTINEZ, MCDONALD, GRAESSER, SHOEMAKER, VECCHI, IMINNIYAZ, DREEZE, CHEN, HALL, MARCH-RUSSELL, SMW...MANY MORE

GENERATING THE ASYMMETRY: CO-GENESIS VS SHARING

□ CO-GENESIS

- ASYMMETRIES IN DM AND BARYONS GENERATED SIMULTANEOUSLY
- DM GENESIS/BARYOGENESIS ALL WRAPPED UP IN ONE MECHANISM
- POTENTIAL TO TEST BOTH DM GENESIS AND BARYOGENESIS

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□ SHARING

- ASSUME PRE-EXISTING ASYMMETRY (EITHER IN BARYONS OR DM)
- ASYMMETRY TRANSFERRED AND SHARED BETWEEN SECTORS
- OPERATORS FOR TRANSFER COULD BE TESTABLE
- GENERALLY HARD TO TEST GENERATION OF INITIAL ASYMMETRY

POSSIBLE MASSES

SIMPLEST CASES, THERE ARE TWO MASS REGIONS

$$m_{dm} \sim 5 \text{ GeV} \text{ AND } m_{dm} \sim 1 \text{ TeV}$$

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IN EQUILIBRIUM AS DM FREEZES-OUT

$$n_{dm} = C n_b \text{ WITH } C \sim \mathcal{O}(1)$$

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WITH $x = \frac{m_{dm}}{T_d}$ CORRECT RATIO FOR

$$m_{dm} \sim 1 \text{ TeV}$$

(ACTUALLY ONLY REALLY CORRECT FOR $m_{dm} \gg T_d$)

SHARING EXAMPLE

KAPLAN, LUTY, ZUREK (2009)

- AT HIGH T, A B-L ASYMMETRY IS GENERATED
- TRANSFER OPERATORS PRESERVE A GLOBAL CHARGE (COMBINATION OF DM AND LEPTON NUMBER)

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- WHEN IN EQUILIBRIUM, THIS OPERATOR TRANSFERS AN L ASYMMETRY INTO THE DM X SECTOR
- NEED TO FIND RELATIONSHIP BETWEEN X ASYMMETRY AND B - NEED TO SOLVE USUAL EQUILIBRATION CONDITIONS

SEE E.G. J. A. Harvey and M. S. Turner, Phys. Rev. D 42, 3344 (1990); T. Inui, T. Ichihara, Y. Mimura and N. Sakai, Phys. Lett. B 325, 392 (1994) [arXiv:hep-ph/9310268].

□ ASSUMING TRANSFER PROCESS DROPS OUT OF THERMAL EQUILIBRIUM ABOVE E-WEAK PHASE TRANSITION

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□ X ASYMMETRY CAN BE CALCULATED IN TERMS OF B-L

$$\eta_X = -\frac{11}{79}(\eta_B - \eta_L)$$

□ THROUGH THE E-WEAK ANOMALY B-L IS TRANSFERRED INTO B

$$\eta_B \approx 0.31(\eta_B - \eta_L)$$

□ FINALLY BY INVERTING $\frac{\Omega_X}{\Omega_b} \sim \frac{\eta_X m_X}{\eta_B m_b}$ A PREDICTION FOR

$$m_X \approx \frac{\eta_B}{\eta_X} \frac{\Omega_X}{\Omega_b} \approx 11 \text{ GeV}$$

CONSTRAINING/TESTING ADM

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□ TWO MAIN AVENUES:

REMOVING SYMMETRIC
DM COMPONENT

LHC LIMITS - MONOJETS, MONOPHOTONS

DIRECT DM DETECTION

HEAVY QUARKONIUM DECAYS

BBN, CMB PERTURBATIONS?

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BBN, CMB PERTURBATIONS?

PROBING ASYMMETRY
SHARING OPERATORS

LHC LIMITS - LONG LIVED STATES

REMOVING THE SYMMETRIC CPT

BUCKLEY; MARCH-RUSSELL, UNWIN, SMW

$$\frac{\Omega_X}{\Omega_b} \sim \frac{\eta_X m_X}{\eta_B m_b}$$

THIS IS TRUE ONLY IF THE X DENSITY IS DETERMINED BY THE ASYMMETRY

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$$\frac{\Omega_X}{\Omega_b} \sim \frac{\eta_X m_X}{\eta_B m_b}$$

THIS IS TRUE ONLY IF THE X DENSITY IS DETERMINED BY THE ASYMMETRY

OTHERWISE:

$$\frac{\Omega_X}{\Omega_b} \sim \frac{n_X + n_{\bar{X}}}{n_B} \frac{m_X}{m_b}$$

LOOSE RELATIONSHIP BETWEEN ABUNDANCES

NEED:

$$n_X + n_{\bar{X}} \approx n_X - n_{\bar{X}}$$

□ NEED TO ANNIHILATE AWAY THE SYMMETRIC COMPONENT...

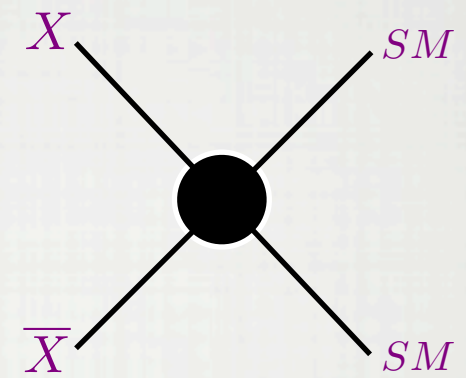
REMOVING THE SYMMETRIC CPT

REMOVING THE SYMMETRIC CPT

□ THREE OPTIONS...

1) ANNIHILATE DIRECTLY TO SM STATES

INTERESTING EXISTING LIMITS ...SEE LATER



REMOVING THE SYMMETRIC CPT

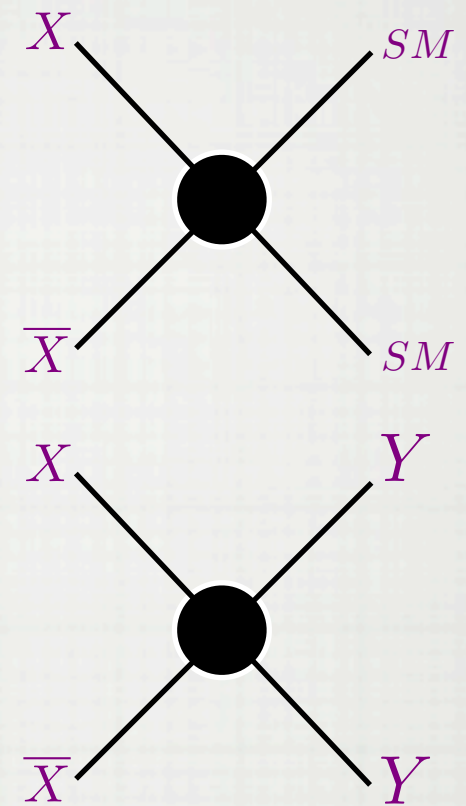
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2) ANNIHILATE DIRECTLY TO LIGHT HIDDEN
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POSSIBLE LONG RANGE DM FORCES



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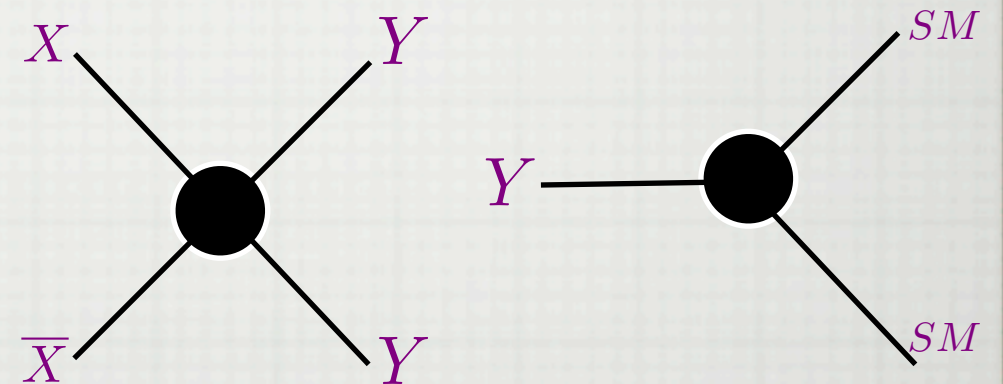
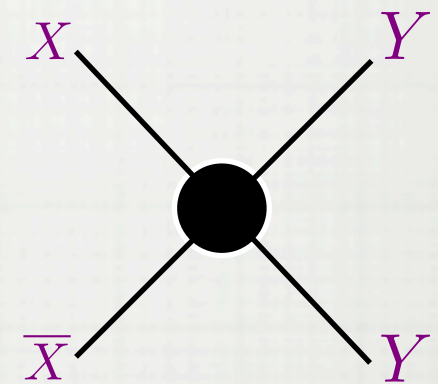
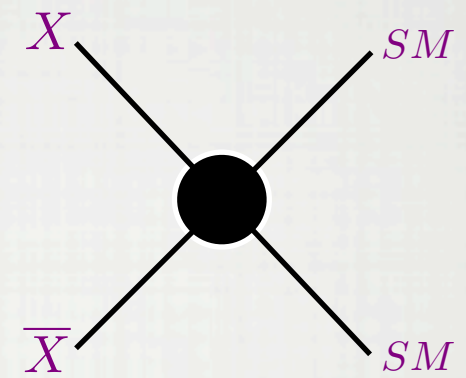
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LATE TIME ENERGY INJECTION IN EARLY UNIVERSE



REMOVING THE SYMMETRIC CPT

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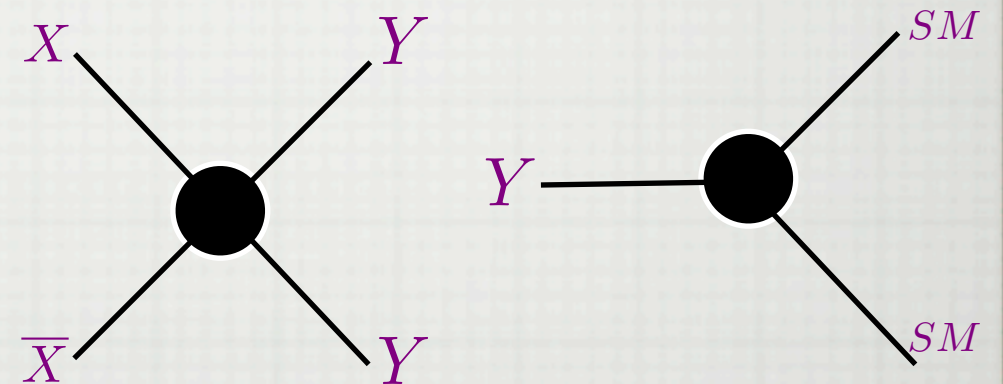
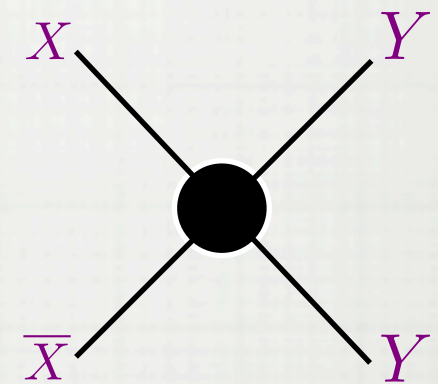
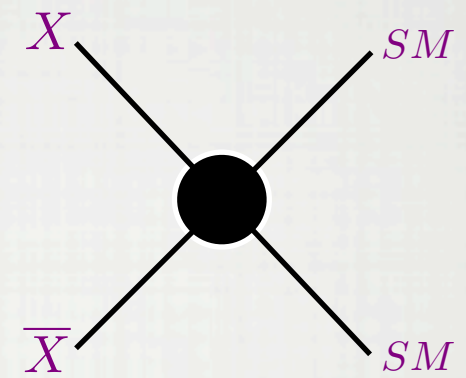
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FOR MORE DETAILS SEE IMINNIYAZ, DREES, CHEN
(1104.5548); GRAESSER, SHOEMAKER, VECCHI, (1103.2771)

IOP joint HEPP and APP 2012

ASYMMETRIC FREEZE-OUT

FOR MORE DETAILS SEE IMINNIYAZ, DREES, CHEN
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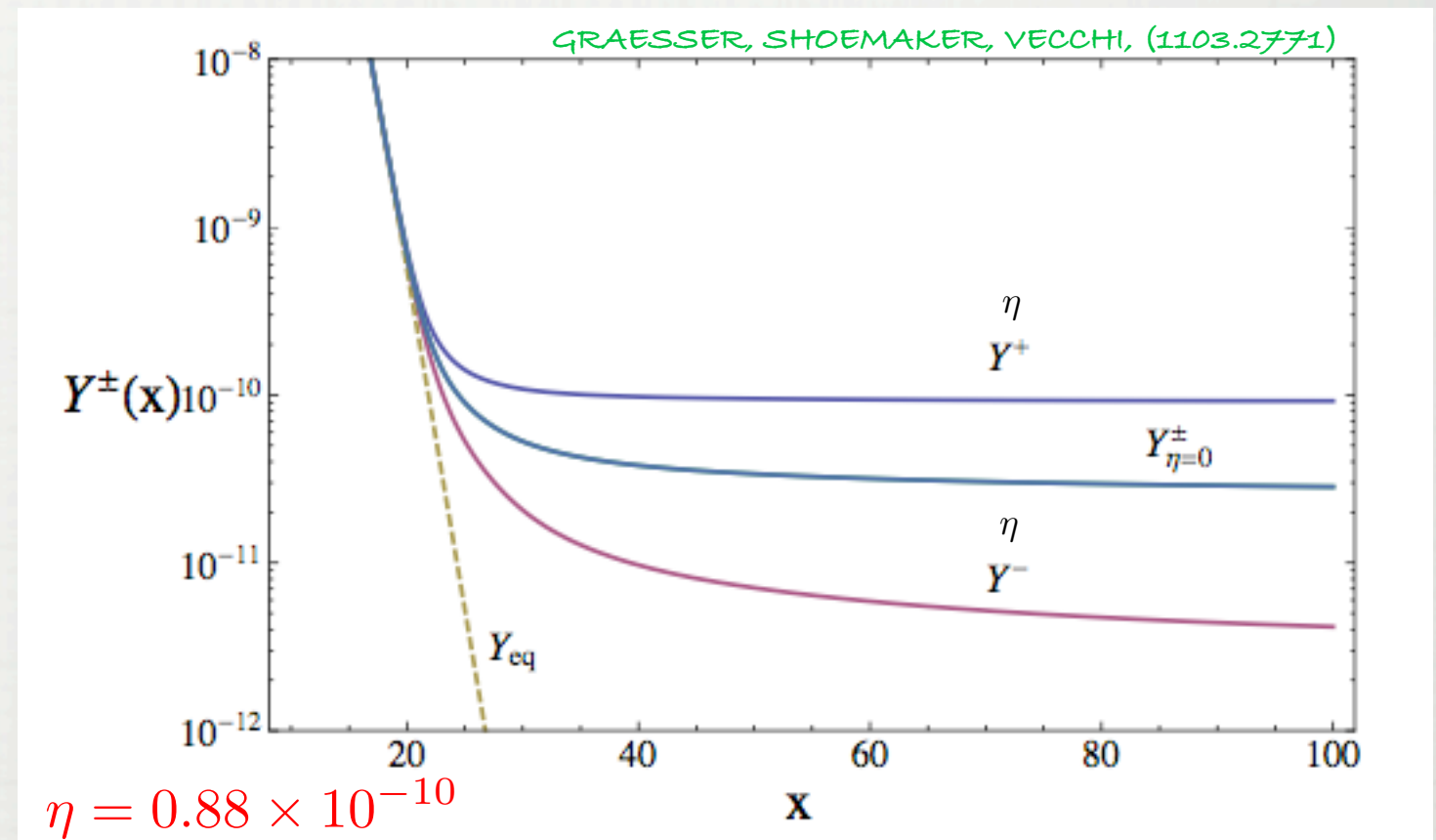
IOP joint HEPP and APP 2012

ASYMMETRIC FREEZE-OUT

□ FREEZE-OUT OPERATES AS USUAL VIA ANNIHILATIONS BUT NOW THE DM HAS AN ASYMMETRY - THIS CHANGES THE FREEZE-OUT DETAILS

ASYMMETRIC AND SYMMETRIC DM FREEZE-OUT, WITH THE SAME ANNIHILATION RATE AND MASS

NEED LARGER ANNIHILATION RATE
NEEDS TO BE APPROX FACTOR OF 2-3 LARGER



REMOVING THE SYMMETRIC COMPONENT

SEE BUCKLEY FOR A FIRST ATTEMPT

REMOVING THE SYMMETRIC COMPONENT

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□ FOR **OPTION 1**), ANNIHILATING DIRECTLY INTO THE SM, WE CAN PARAMETERISE THE INTERACTIONS IN TERMS OF **EFFECTIVE OPERATORS** ASSUMING SOME HEAVY MEDIATOR

□ E.G. DIRAC FERMION DM ψ

$$\frac{m_q}{\Lambda^3} \bar{\psi} \psi \bar{q} q$$

$$\bar{\psi} \psi \rightarrow \bar{q} q$$

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MONOJETS/MONOPHOTONS AT THE LHC $\bar{q} q \rightarrow g \bar{\psi} \psi$ $\bar{q} q \rightarrow \gamma \bar{\psi} \psi$

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$$\psi q \rightarrow \psi q$$

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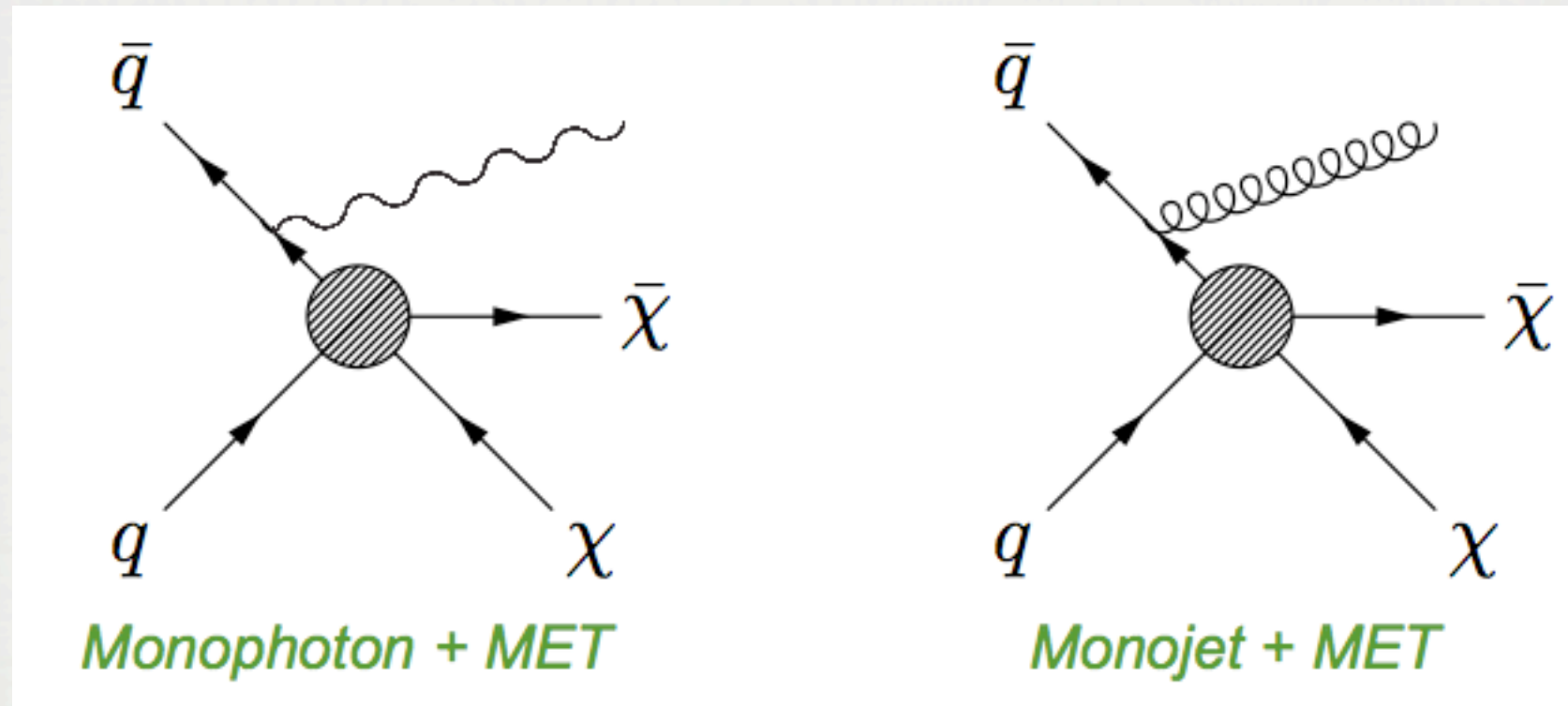
DIRECT DETECTION $\psi q \rightarrow \psi q$

HEAVY QUARKONIUM DECAY $\bar{q} q \rightarrow \bar{\psi} \psi$ $\bar{q} q \rightarrow \gamma \bar{\psi} \psi$

(NOT TOO RESTRICTIVE IN MOST CASES)

MONOJET/MONOPHOTON LIMITS: MODEL INDEPENDENT LIMITS ON DM

STEVE WORM TALK AT MORIOND 2012



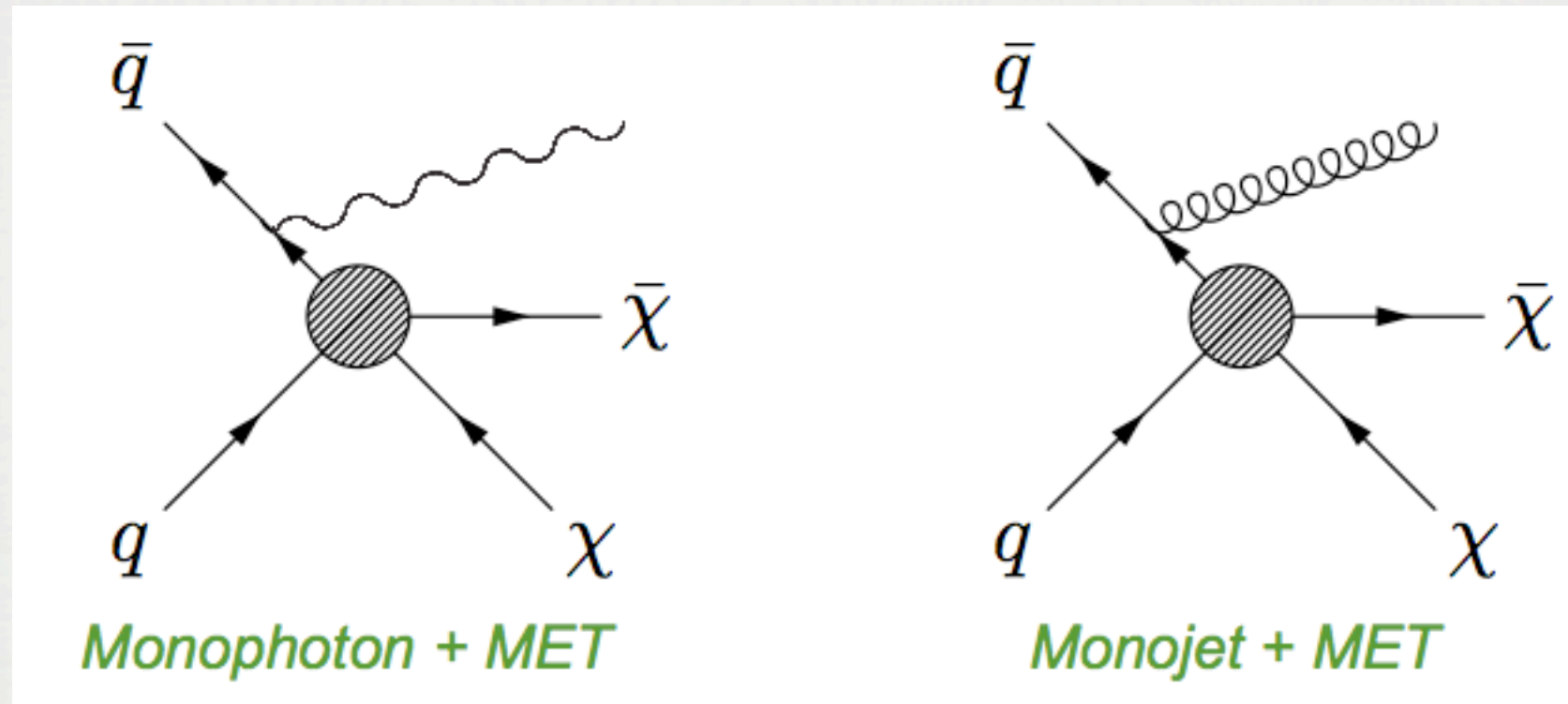
E.G. BAI, FOX AND HARNIK, JHEP 1012:048 (2010); GOODMAN, IBE, RAJARAMAN, SHEPHERD, TAIT, YU, PHYS.REV.D82:116010

ATLAS: ATLAS-CONF-2011-096 (2011)

LATEST CMS: EXO-11-059, EXO-11-096, [HTTPS://TWIKI.CERN.CH/TWIKI/BIN/VIEW/CMSPUBLIC/PHYSICSRESULTSEXO](https://twiki.cern.ch/twiki/bin/view/cmSPublic/PhysicsResultSEXO)

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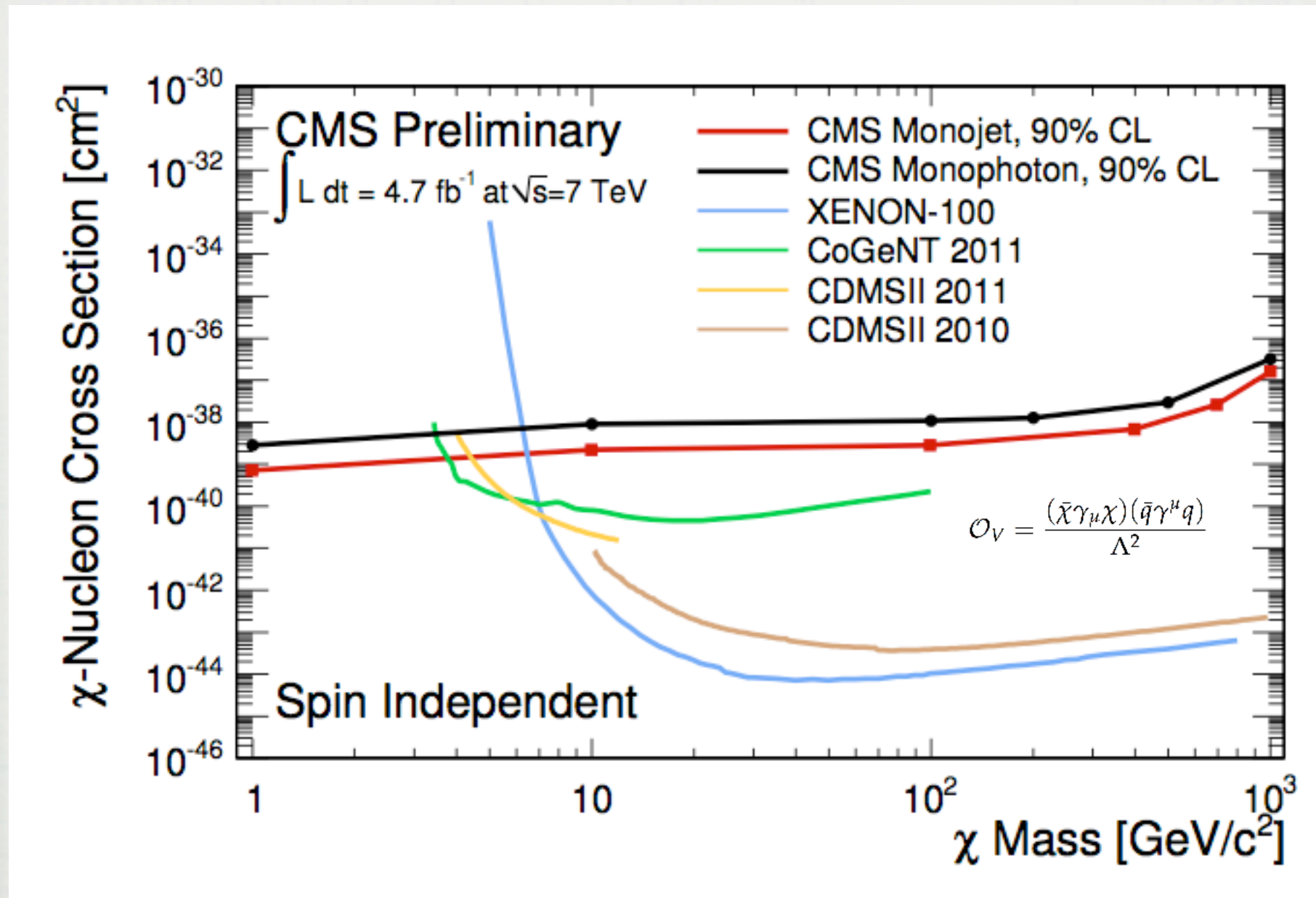
□ ATLAS AND CMS HAVE UPPER LIMITS ON THESE PROCESSES - RECENTLY UPDATED

E.G. BAI, FOX AND HARNIK, JHEP 1012:048 (2010); GOODMAN, IBE, RAJARAMAN, SHEPHERD, TAIT, YU, PHYS.REV.D82:116010

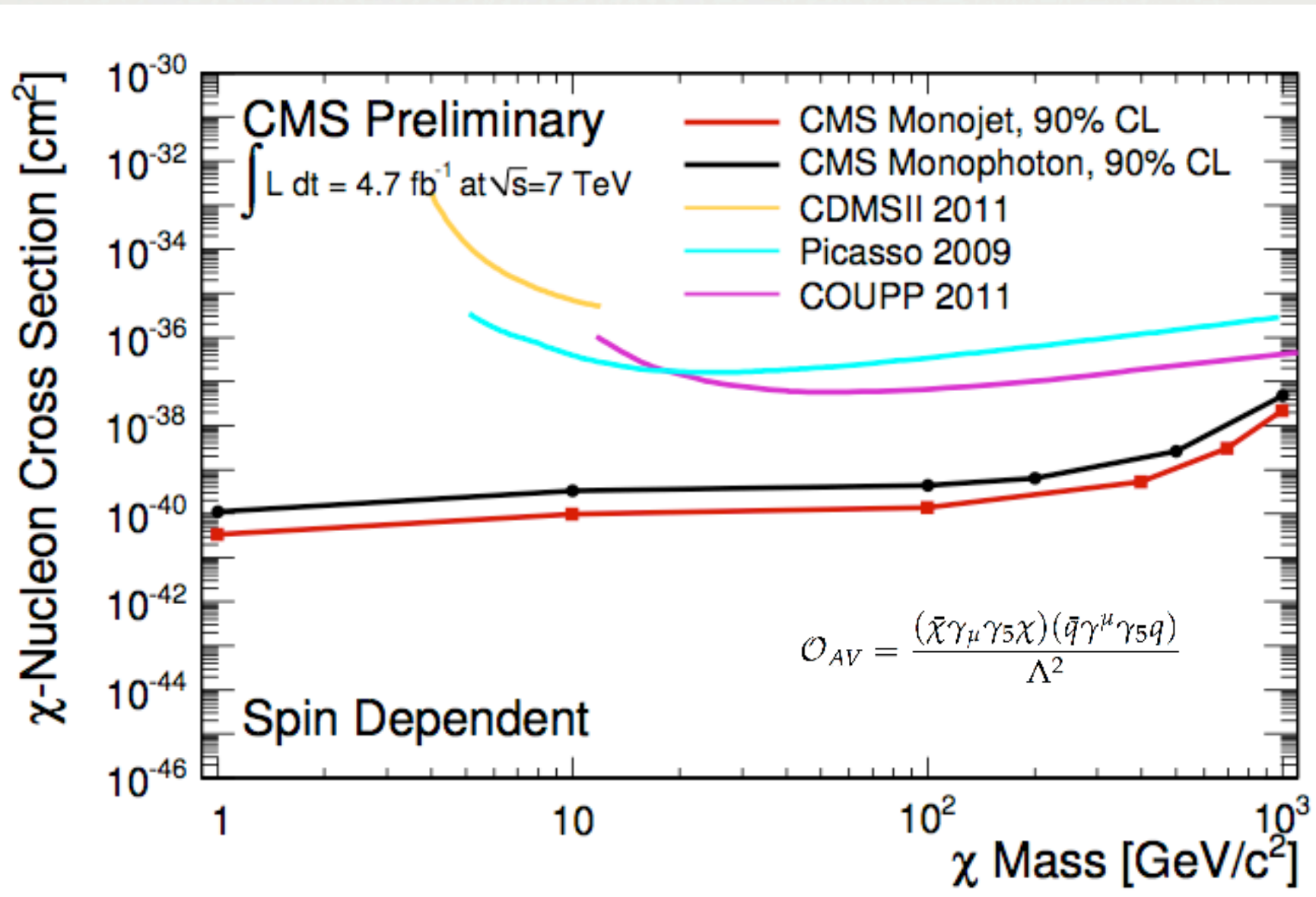
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MONOJET/MONOPHOTON LIMITS

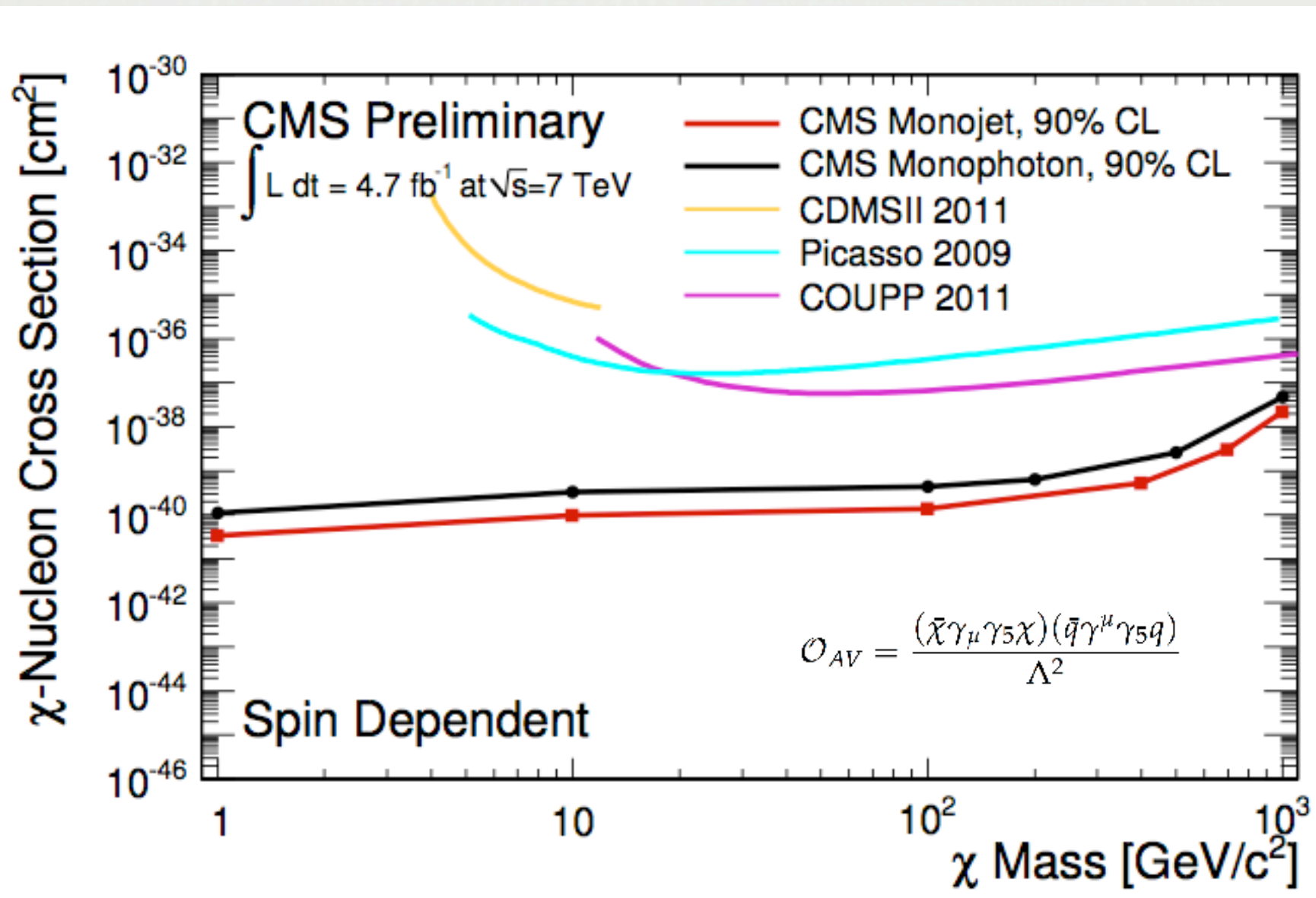


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MONOJET/MONOPHOTON LIMITS



☐ MUST BE CAREFUL WITH THESE LIMITS - NOT ALWAYS CORRECT TO USE EFFECTIVE THEORY DESCRIPTION

LATEST CMS: EXO-11-059, EXO-11-096, [HTTPS://TWIKI.CERN.CH/TWIKI/BIN/VIEW/CMSPUBLIC/PHYSICSRESULTSEXO](https://twiki.cern.ch/twiki/bin/view/CMSPUBLIC/PHYSICSRESULTSEXO)

THE CONCERN WITH EFFECTIVE OPERATOR APPROACH

M_χ (GeV/ c^2)	Spin-dependent		Spin-independent	
	$\sigma(\text{cm}^2)$	Λ (GeV)	$\sigma(\text{cm}^2)$	Λ (GeV)
1	3.37×10^{-41}	730	7.20×10^{-40}	776
10	9.83×10^{-41}	744	2.12×10^{-39}	789
100	1.33×10^{-40}	718	2.65×10^{-39}	776
400	5.14×10^{-40}	514	6.66×10^{-39}	619
700	2.95×10^{-39}	332	2.62×10^{-38}	440
1000	2.15×10^{-38}	202	1.57×10^{-37}	281

□ ANALYSIS BECOMES MODEL DEPENDENT - MUST BE CAREFUL

LATEST CMS: EXO-11-059, EXO-11-096, [HTTPS://TWIKI.CERN.CH/TWIKI/BIN/VIEW/CMSPUBLIC/PHYSICSRESULTSEXO](https://twiki.cern.ch/twiki/bin/view/cmSPublic/PhysicsResultSEXO)

MONOJET + DIRECT DETECTION VS ADM

□ PLACES A LOWER LIMIT ON Λ FOR LARGE CUT OFF ONLY

$$\sigma_{\text{mono}} \propto \frac{1}{\Lambda^6}$$

$$\frac{m_q}{\Lambda^3} \bar{\psi}\psi\bar{q}q$$

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□ CAN ALSO CONVERT LIMITS ON THE SPIN INDEPENDENT ELASTIC SCATTERING CROSS SECTION TO A LOWER LIMIT ON Λ

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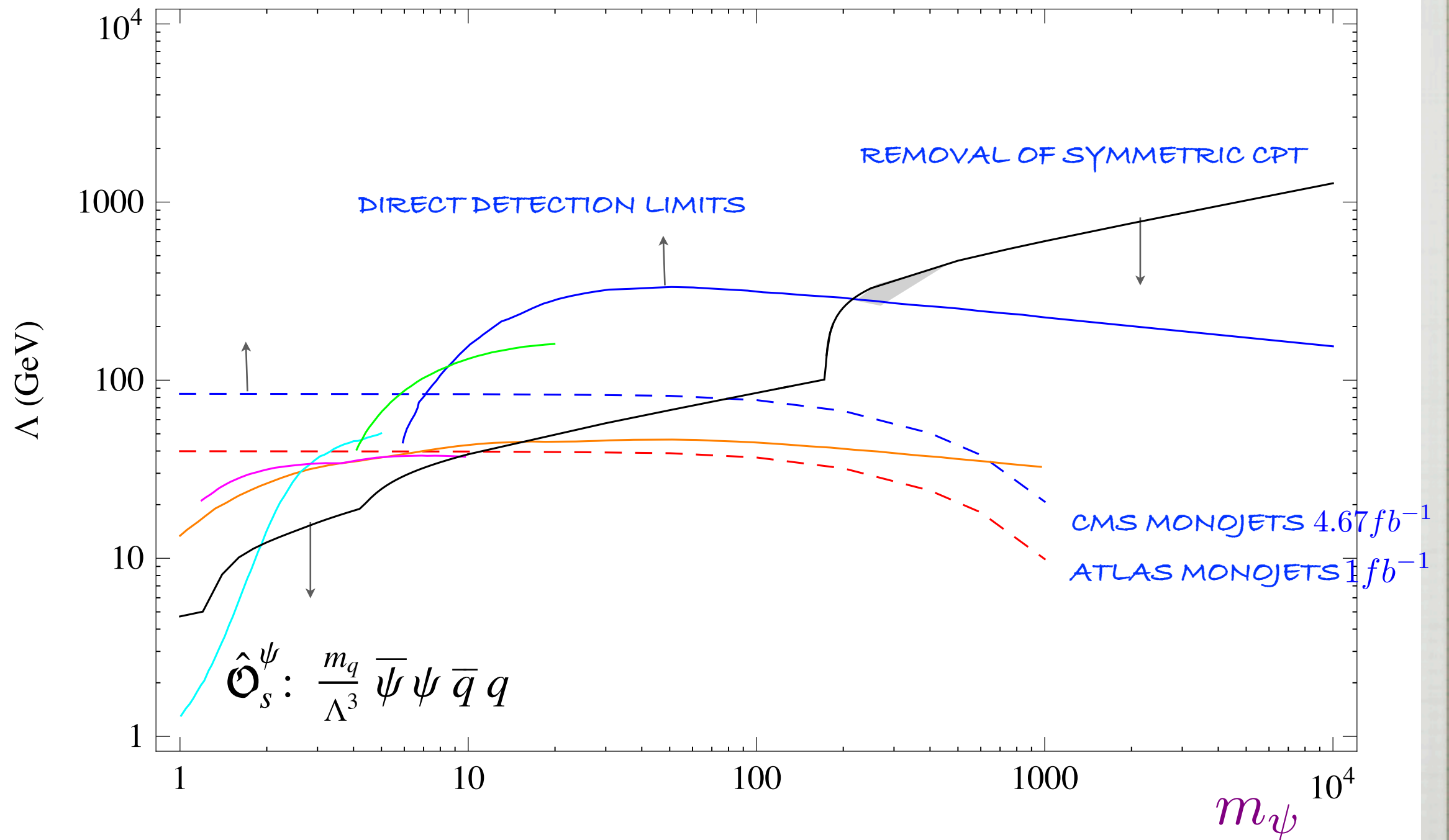
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□ NEEDING TO ANNIHILATE AWAY SYMMETRIC CPT GIVES A MAXIMUM Λ

$$\sigma_{\text{ann}} \propto \frac{1}{\Lambda^6}$$

MONOJET + DIRECT DETECTION VS ADM



LIGHT MEDIATORS

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- FOR LOW CUT OFF/MEDIATOR, EFFECTIVE THEORY NO LONGER APPROPRIATE
- NEED TO TAKE ACCOUNT OF A REAL MEDIATOR BETWEEN DM AND SM SECTOR

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□ SIMPLE EXAMPLE FERMION DM: $\mathcal{L} = \lambda_X \eta \bar{X} X + \lambda'_q \bar{q} q \eta$

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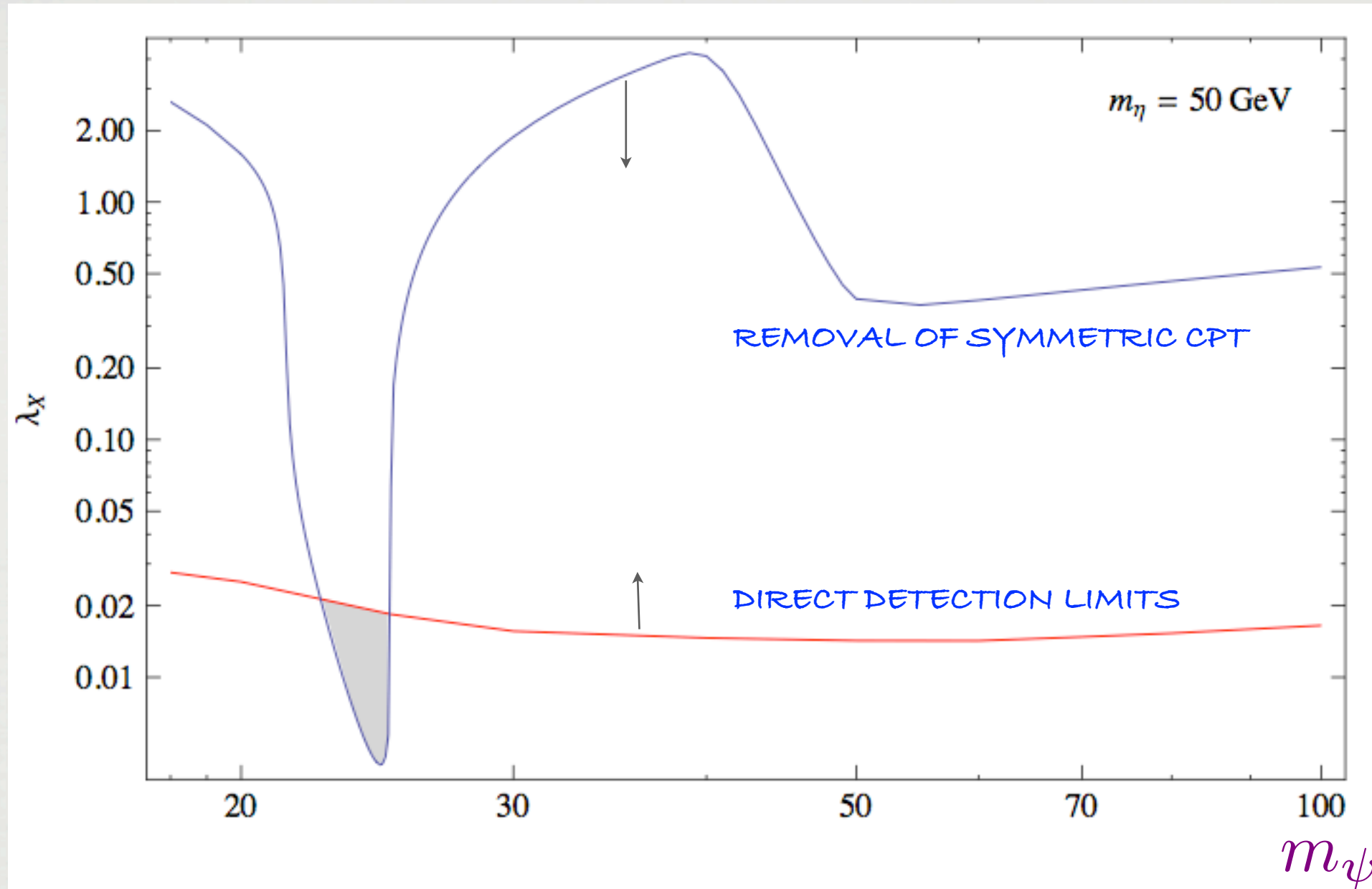
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□ SIMPLE EXAMPLE FERMION DM: $\mathcal{L} = \lambda_X \eta \bar{X} X + \lambda'_q \bar{q} q \eta$

□ FIXING THE SIZE OF THE QUARK COUPLING TO THE MEDIATOR WE FIND THE VALUE OF λ_X THAT GIVES A LARGE ENOUGH ANNIHILATION RATE TO REMOVE SYMMETRIC PART

□ WE CAN APPLY DIRECT DETECTION CONSTRAINTS AS BEFORE

LIGHT MEDIATORS



- MOST MEDIATOR MASSES, ONLY IN THE RESONANT REGION IS IT POSSIBLE
- POSSIBILITIES FOR PSEUDO-SCALAR MEDIATOR...MORE TO COME

BOTTOM LINE FOR DIRECT ANNIHILATION TO SM STATES

□ IT IS VERY CONSTRAINED - POSSIBILITIES FOR LIGHT PSEUDO-SCALAR MEDIATOR, OTHERWISE UNLIKELY

⇒ NEED A MORE COMPLICATED HIDDEN SECTOR

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POSSIBLE LONG RANGE DM FORCES

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LATE TIME ENERGY INJECTION IN EARLY UNIVERSE

SHARING OPERATORS - LHC SIGNALS

- MANY ASYMMETRY TRANSFER OPERATORS CAN LEAD TO LONG LIVED PARTICLES AT THE LHC
- FOR EXAMPLE, IN SUSY MODELS THE $L\bar{O}SP$ CAN BE LONG LIVED IF IT HAS A SMALL DECAY WIDTH TO THE DM STATE THROUGH A CONNECTOR OPERATOR

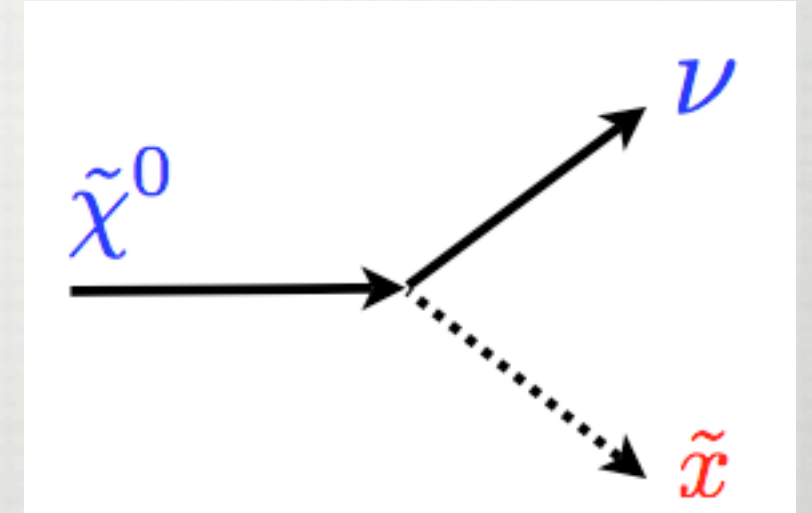
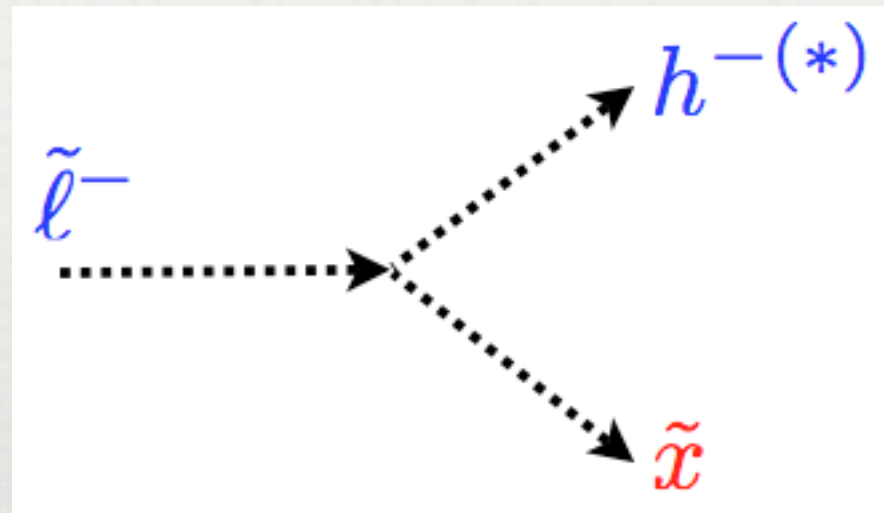
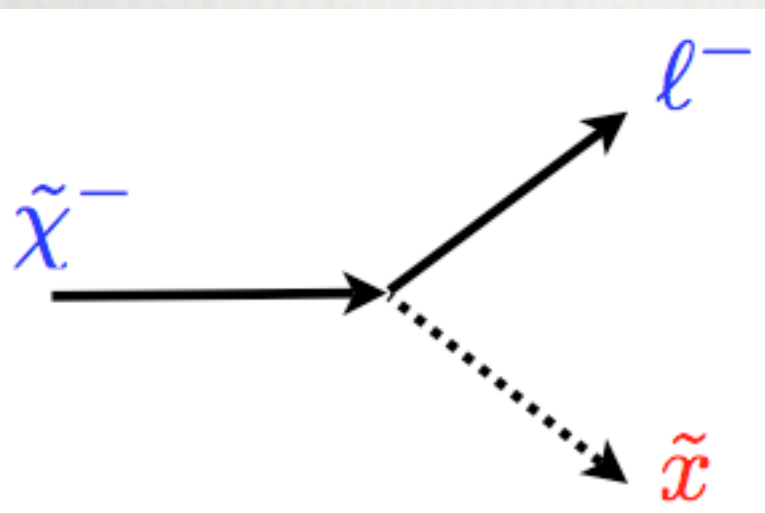
$$\Delta W = \lambda L H_u X$$

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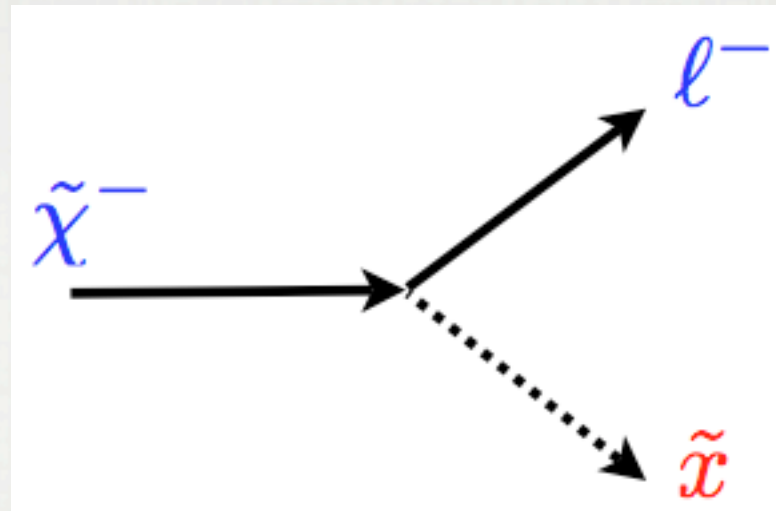
$$\Delta W = \lambda L H_u X$$

- INTRODUCES THE FOLLOWING INTERACTIONS:



□ DEPENDING ON WHAT IS THE LIGHTEST ORDINARY SUSY PARTICLE WE WILL GET THE FOLLOWING AT THE END OF SUSY DECAY CHAIN:

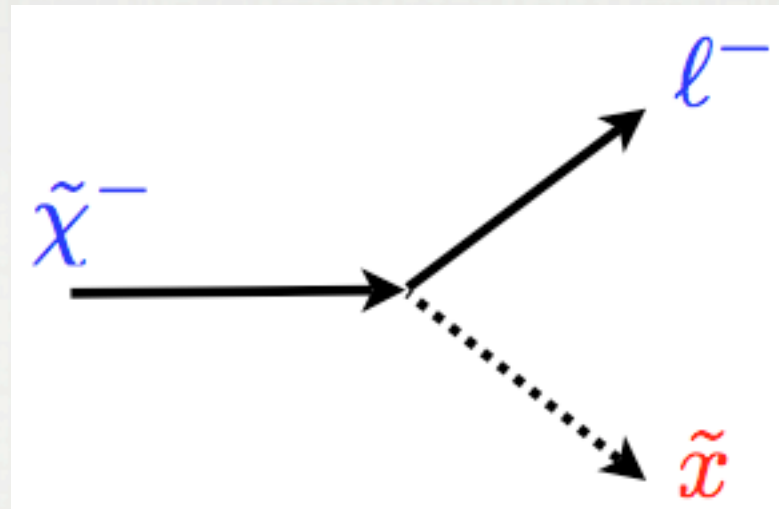
CHARGINO LOSS



CHARGED TRACK, WITH A KINK

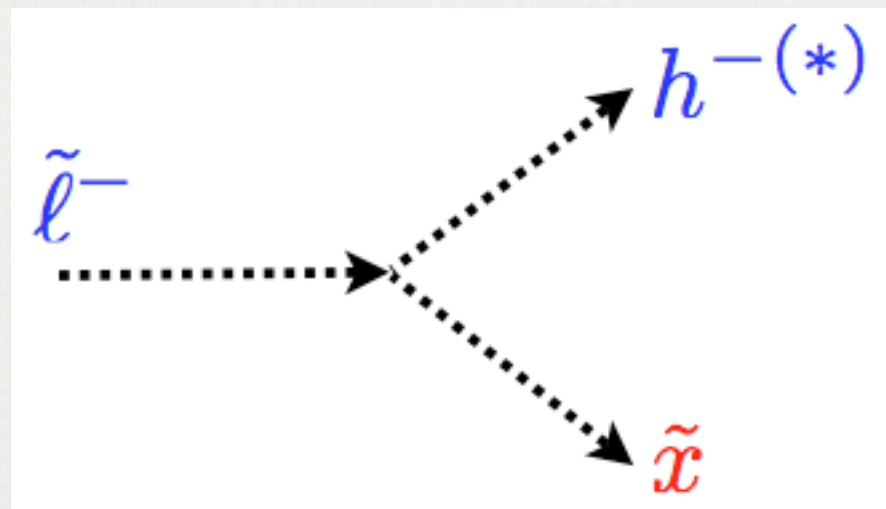
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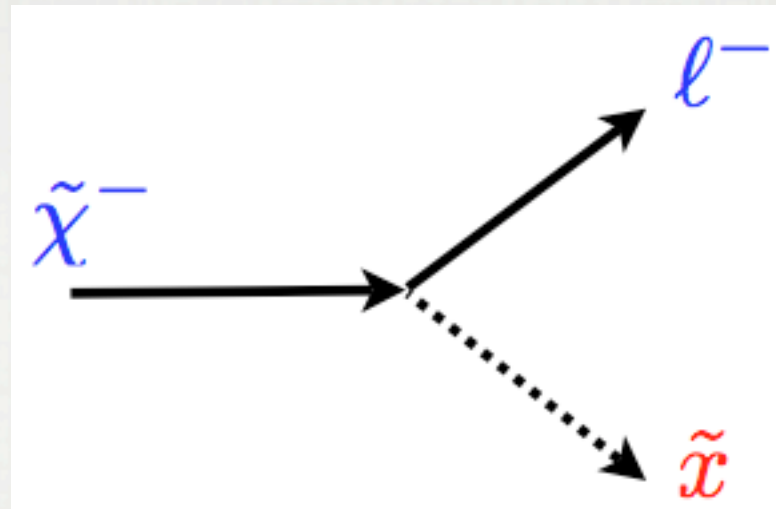
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SLEPTON LOSS



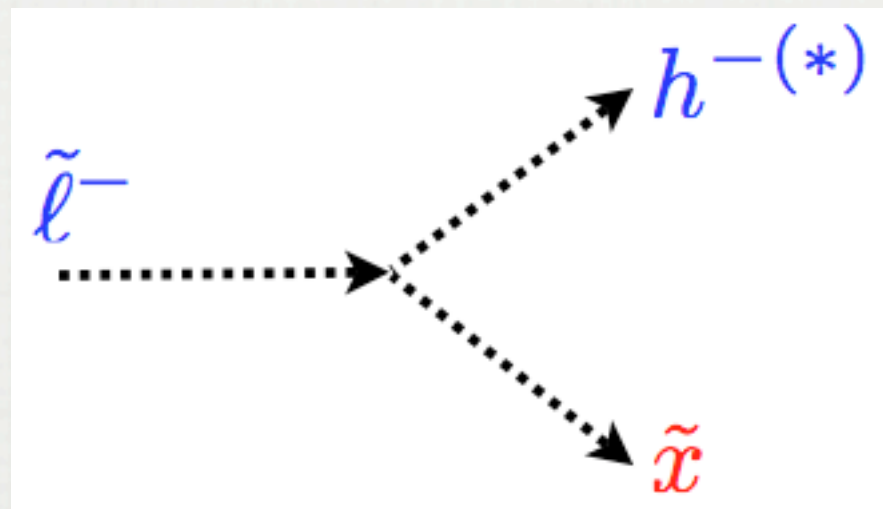
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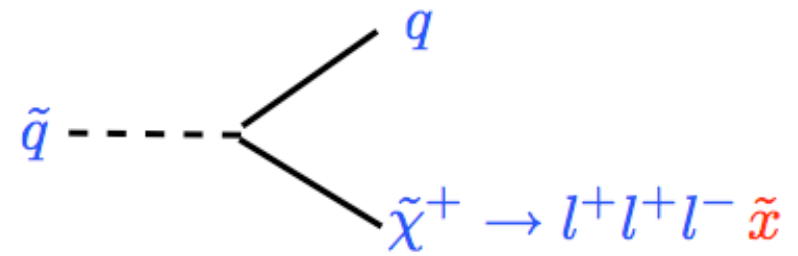
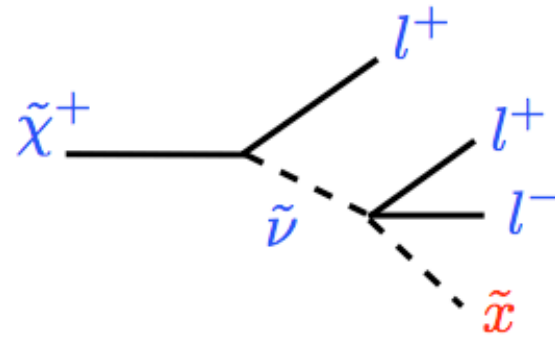
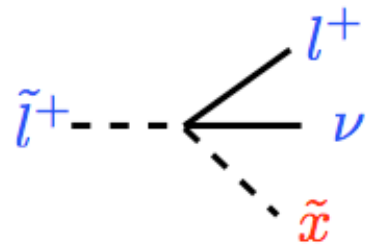
SLEPTON LOSS



□ REMEMBERING THAT THE SHARING OPERATORS MUST DROP OUT OF THERMAL EQUILIBRIUM BEFORE FREEZE-OUT - COUPLING MUST BE SMALL

$CT \sim$ primary vertex - many meters

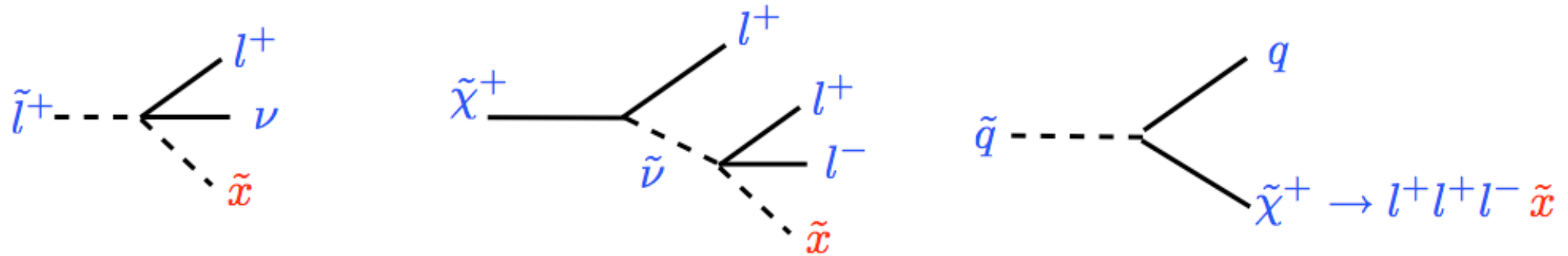
eg, $W = \frac{1}{M} L L E^c X$



□ ANOTHER EXAMPLE

eg, $W = \frac{1}{M} L L E^c X$

MARCH-RUSSELL TALK, HEIDELBERG 2011



□ NOTE: EACH SUSY DECAY CHAIN WILL END IN THIS DECAY

□ ALL DECAYS FORM SHARING OPERATORS VIOLATE **B OR L**

□ NEED A COMBINATION OF LIMITS ON PROMPT EVENT AND DISPLACED VERTEX

CONSTRAINTS FROM THE SUN

- IF DM HAS LARGE SPIN-DEPENDENT SCATTERING CROSS SECTION OR SELF INTERACTING, DM CAN ACCUMULATE IN THE SUN
- OLD IDEA TO SOLVE SOLAR NEUTRINO PROBLEM - COSMIONS/LOW MASS DM IN THE SUN TRANSPORTS ENERGY AWAY FROM CORE
- DM WITH AN ASYMMETRY NEEDED SO THAT ABUNDANCE BUILT UP
- CHANGES TEMP PROFILE, WHICH AFFECTS THE NEUTRINO FLUXES -- OF COURSE NOW SOLVED BY OSCILLATIONS

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- CHANGES TEMP PROFILE, WHICH AFFECTS THE NEUTRINO FLUXES -- OF COURSE NOW SOLVED BY OSCILLATIONS
- IN NEW MODELS OF ADM, THE COSMION CONDITIONS COULD BE REPRODUCED
- CAPTURE OF ADM BY THE SUN, COULD THEN BE CONSTRAINED BY THE PROPERTIES OF THE SUN OR MAY EVEN ALLEVIATE POTENTIAL ISSUES WITH THE STANDARD SOLAR MODEL

SERENLLI, BASU, FERGUSON (2009),
ASPLUND, GREVESSE, SAUVAL (2004, 2009)

ADM/COSMION PAPERS: FAULKNER, GILLILAND (1985); SPERGEL, PRESS (1985); GILLILAND, FAULKNER, PRESS, SPERGEL (1986); GELMINI, HALL, LIN (1987); GIUDICE, RABY (1990); LOPES, SILK, HANSEN, BERTONE (2002) FRANDSEN, SARKAR (2010); CUMBERBATCH, GUZIK, SILK, WATSON, SMW (2010); TAOSO, IOCCO, MEYNET, BERTONE, EGGENBERGER (2010)

CONCLUSIONS

- ADM IS AN INTERESTING AND WELL MOTIVATED DM SCENARIO TO EXPLAIN

$$\frac{\Omega_{dm}}{\Omega_b} \sim 5$$

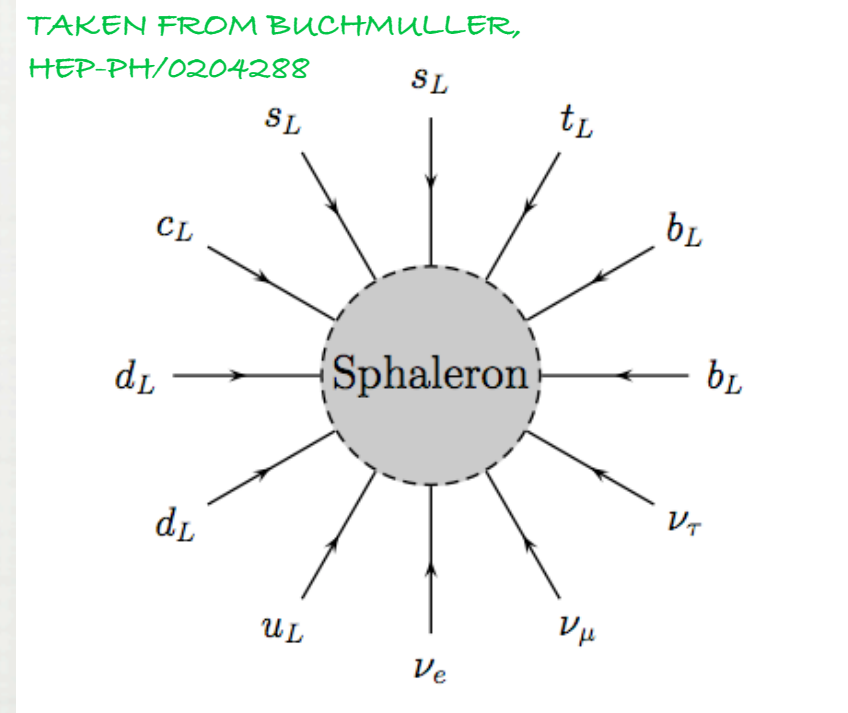
- REQUIRE A **SHARED (GLOBAL) QUANTUM NUMBER** BETWEEN DM AND SM
- TWO MAIN SCENARIOS, **CO-GENESIS** (DM AND B ASYMMETRY GENERATED SIMULTANEOUSLY) AND **SHARING** WHERE A PRE-EXISTING ASYMMETRY IS TRANSFERRED BETWEEN DM AND SM SECTORS
- MANY CONSTRAINTS ON ADM, ESPECIALLY REMOVAL OF SYMMETRIC COMPONENT
- LHC PHENOMENOLOGY OF SHARING OPERATORS- LONG LIVED EXOTICS
- LOTS MORE TO INVESTIGATE...

BACK UPS AND OLD SLIDES

IMPORTANT ASIDE ON THE ELECTROWEAK ANOMALY/SPHALERONS

- HAS AN IMPORTANT INFLUENCE OVER THE DYNAMICS OF ASYMMETRIES IN ANY CHIRAL FERMION CHARGED UNDER $SU(2)_L$
- **B+L VIOLATING** PROCESS, **CONSERVES B-L** EFFICIENTLY OPERATE $10^{12} \text{ GeV} > T \gtrsim 100 \text{ GeV}$ (BELOW EXPONENTIALLY SUPPRESSED)

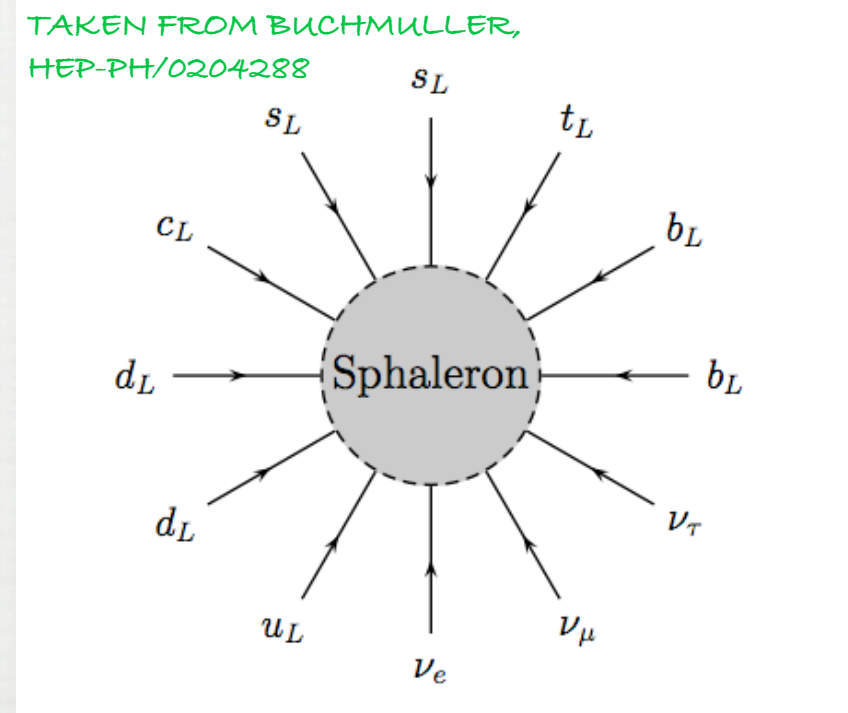
□ CAN EFFECTIVELY BE THOUGHT OF AS MULTI-PARTICLE VERTEX INVOLVING $SU(2)_L$ CHARGED STATES



IMPORTANT ASIDE ON THE ELECTROWEAK ANOMALY/SPHALERONS

- HAS AN IMPORTANT INFLUENCE OVER THE DYNAMICS OF ASYMMETRIES IN ANY CHIRAL FERMION CHARGED UNDER $SU(2)_L$
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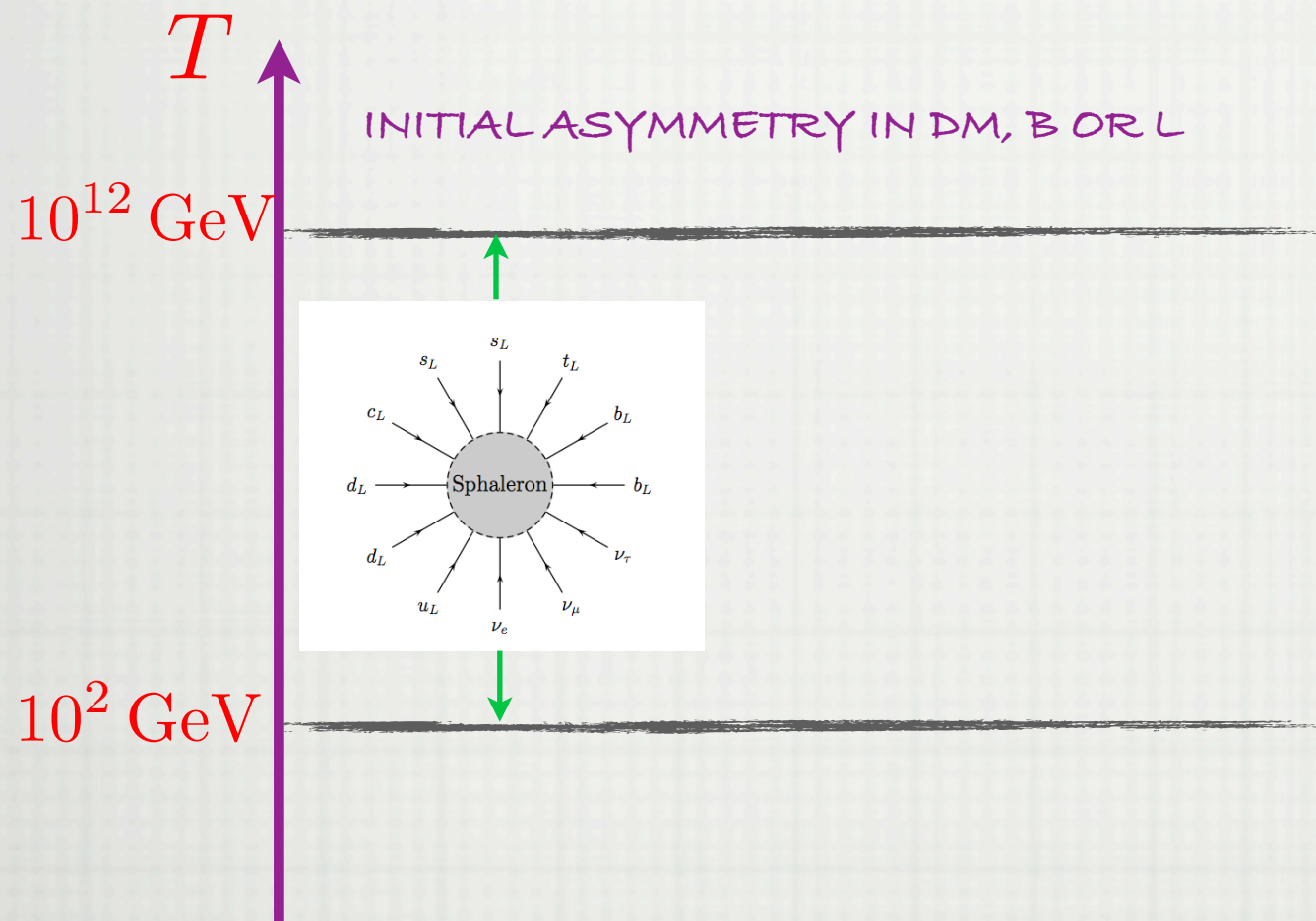
- IF $L \neq 0$, $B=0$ SPHALERONS WILL **REPROCESS L ASYMMETRY INTO B NUMBER**
- IF $B \neq 0$, $L \neq 0$ BUT $B-L=0$, E-WEAK ANOMALY WILL WASH OUT THE ASYMMETRY

SHARING:

- ASSUME THE PRESENCE OF A NON ZERO PRE-EXISTING HIGH SCALE ASYMMETRY IN EITHER DM , B OR L .

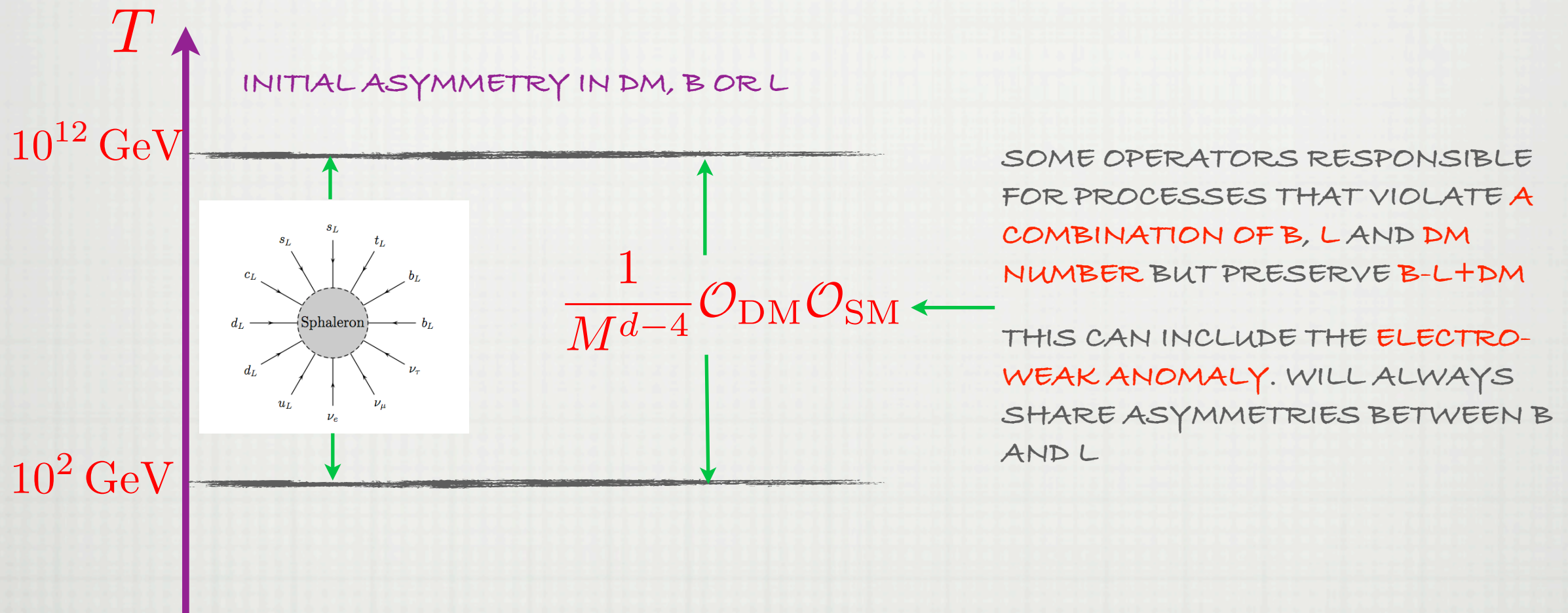
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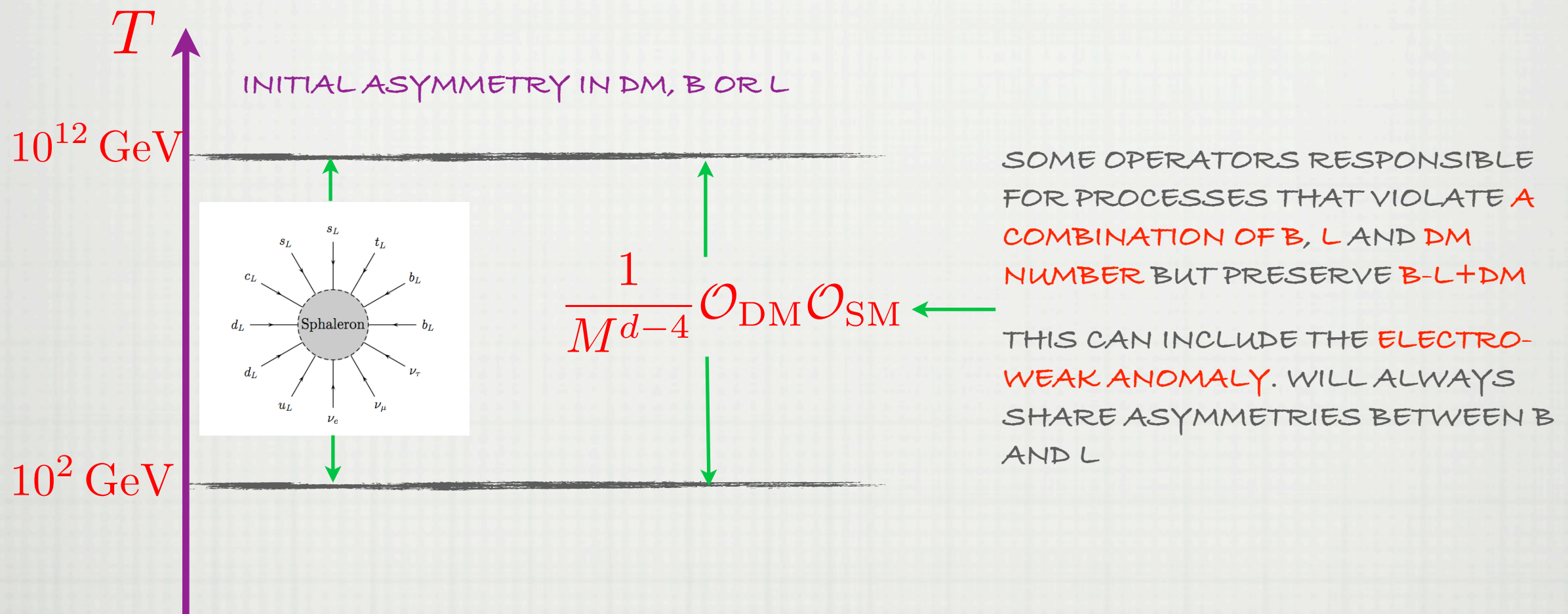
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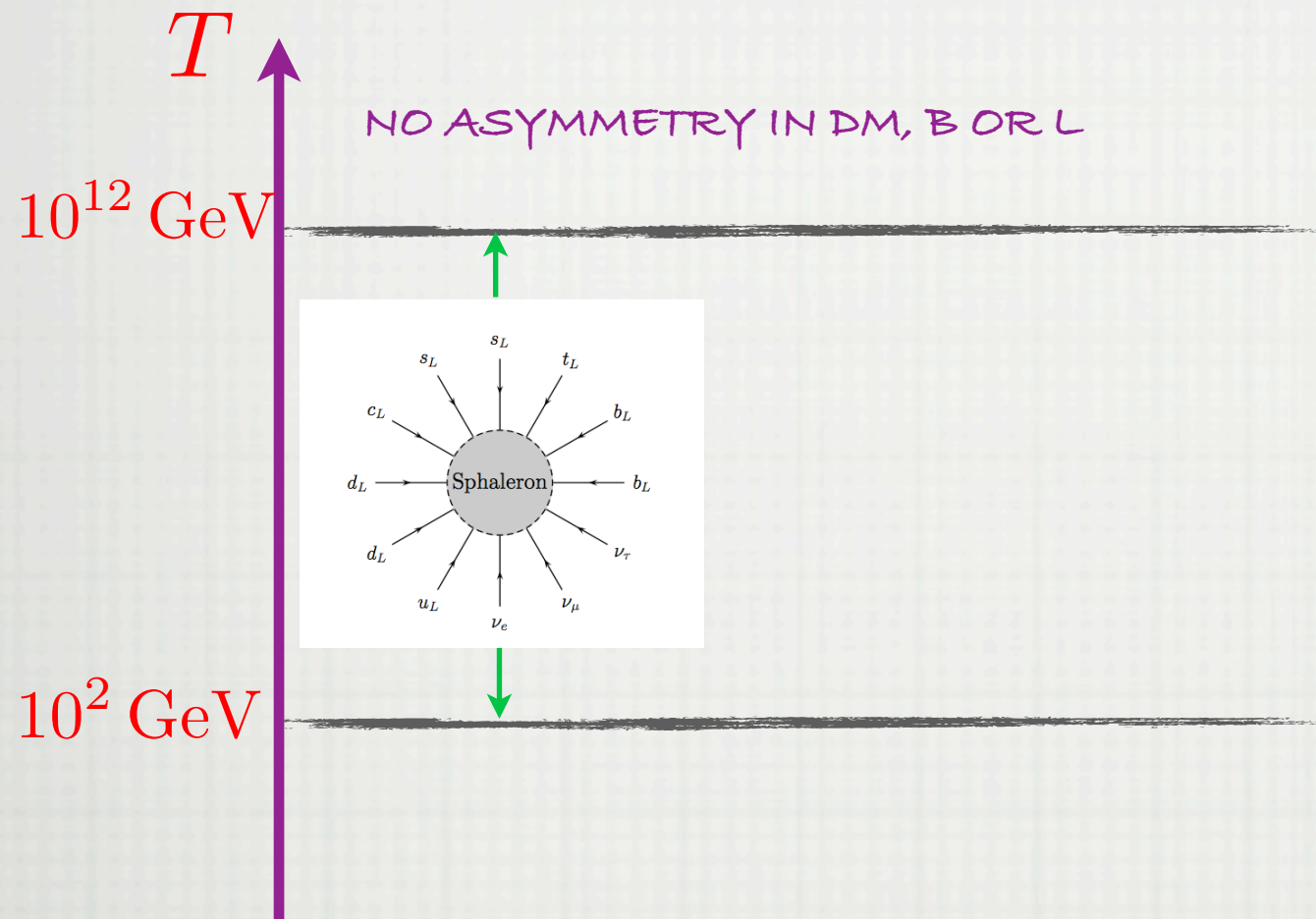
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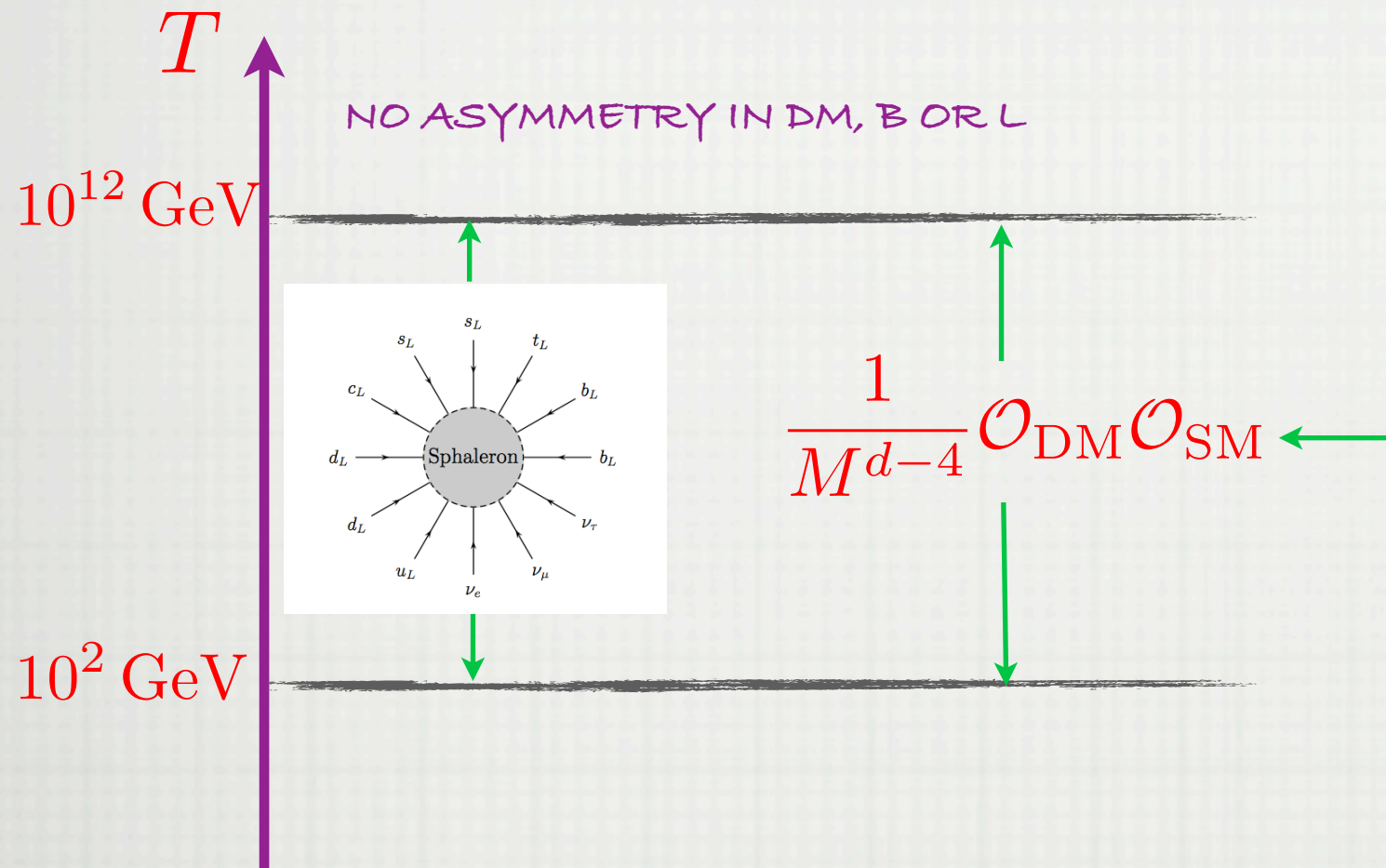


□ THE RESULT IS THAT ANY ASYMMETRIES IN **B**, **L** OR **DM** ARE SHARED AND RELATED BY **B-L+DM** NUMBER - EXAMPLE LATER

CO-GENESIS EXAMPLES - GENERATING AN ASYMMETRY IN B AND DM SIMULTANEOUSLY



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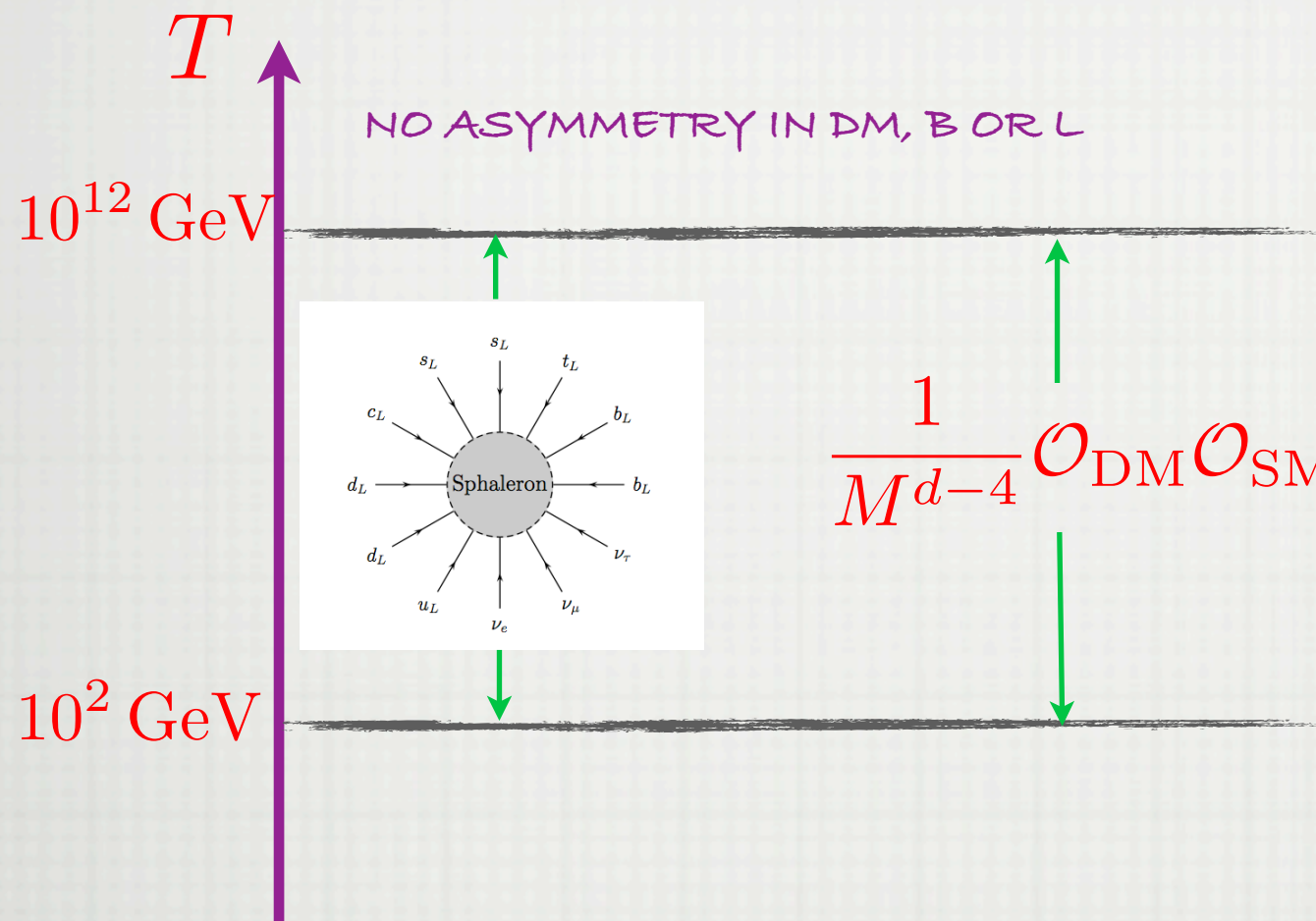


NO ASYMMETRY IN DM, B OR L

AGAIN COMBINATION OF **B**, **L** AND **DM** NUMBER VIOLATED AND **B-L** + **DM** PRESERVED.

BUT NOW, THESE INTERACTIONS ARE RESPONSIBLE FOR GENERATING THE ASYMMETRY.

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BUT NOW, THESE INTERACTIONS ARE RESPONSIBLE FOR GENERATING THE ASYMMETRY.

□ MUCH MORE ELEGANT AND TESTABLE - BUT HARD TO BUILD A WORKING MODEL

CO-GENESIS IS HARD: SOME EXAMPLES

□ E-WEAK BARYOGENESIS (EWB) KAPLAN DB (1992)

▶ EXTRA $U(1)_{DM}$ SYMMETRY WITH WEAK ANOMALY

▶ STABLE PARTICLES CHARGED UNDER $U(1)_{DM}$ WILL BE PRODUCED IN EWB WITH BARYONS

▶ DM STATES CHARGED UNDER $SU(2)_L$

▶ MUST ALSO HAVE LIGHT MASSES (SUB 45GEV)

⇒ SIMPLE MODEL RULED OUT BY COUPLINGS TO Z
(DIRECT DETECTION AND INVISIBLE Z-WIDTH)

⇒ GENERALLY DIFFICULT TO TEST, HIGH SCALE DYNAMICS

SUBSET OF RELATED: THOMAS, DAVOUDIASH, MORRISSEY, SIGURDSON, TULIN, HALL, MARCH-RUSSELL, SMW, CHUN, BLENNOW, ALLAHVERDI, FALKOWSKI, RUDERMAN, VOLANSKY, ZUREK, CHEUNG, MCCULLOUGH.

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□ ASYMMETRIC FREEZE-IN...MORE LATER

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E.G.

$$\mathcal{L} \sim \frac{1}{M^{d-4}} \mathcal{O}_{dm} \mathcal{O}_{sm}$$

d = DIMENSION OF COMBINED OPERATOR

\mathcal{O}_{sm} AND \mathcal{O}_{dm} INDIVIDUALLY CHARGED UNDER GLOBAL $U(1)$, BUT COMBINED OPERATOR IS INVARIANT UNDER $U(1)$

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- IF ASYMMETRY EXISTS IN EITHER SM OR DM SECTOR, THESE OPERATORS WILL SHARE THIS WITH THE OTHER SECTOR
- OPERATORS MUST BE IN THERMAL EQUILIBRIUM ABOVE $T = m_{dm}$
- HOWEVER, THEY MUST DROP OUT OF THERMAL EQUILIBRIUM ABOVE DM FREEZE-OUT OTHERWISE THEY WILL HEAVILY SUPPRESS THE ASYMMETRY - ACTUALLY LEADS TO TEV SCALE POSSIBILITY - SEE LATER

BARYON ASYMMETRY

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$$\xi = \frac{n_B}{n_\gamma} \approx (1.5 - 6.3) \times 10^{-10} \quad n_\gamma \sim 400/cm^3$$

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□ NO SIGNIFICANT ANTI-MATTER IN THE UNIVERSE

$$\eta_B = Y_B - Y_{\bar{B}} = \frac{n_B - n_{\bar{B}}}{s} = \frac{\xi}{7} \quad s \approx 3 \times 10^3/\text{cm}^3$$

$$\Rightarrow \Omega_b h^2 \approx 0.02$$

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□ WHERE DOES THE ASYMMETRY COME FROM? LEPTOGENESIS?
BARYOGENESIS?

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□ MOREOVER WE NOTE:

$$\frac{\Omega_{dm}}{\Omega_b} \sim 5$$

□ TAKE SERIOUSLY THE CLOSENESS OF THESE VALUES -
INVESTIGATE DYNAMICS THAT LINK THE TWO...

...LEADS TO IDEAS OF ASYMMETRIC DM

WHAT ABOUT CO-GENESIS? ASYMMETRIC FREEZE-IN.

FIRST, WHAT IS FREEZE-IN?

FREEZE-IN OVERVIEW

HALL, JEDAMZIK, MARCH-
RUSSELL, SMW, ARXIV:0911.1120

- FREEZE-IN IS RELEVANT FOR PARTICLES THAT ARE FEEBLY COUPLED
(VIA RENORMALISABLE COUPLINGS) - λ
FEEBLY INTERACTING MASSIVE PARTICLES - FIMPS X

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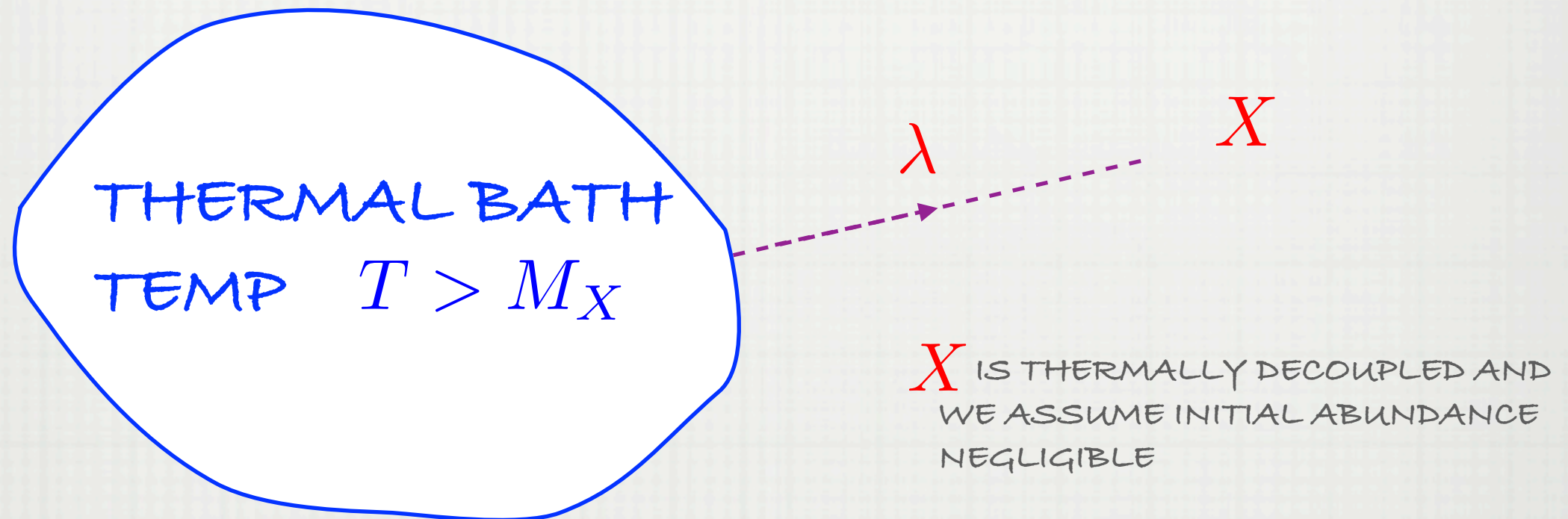
THERMAL BATH
TEMP $T > M_X$

X IS THERMALLY DECOUPLED AND
WE ASSUME INITIAL ABUNDANCE
NEGLIGIBLE

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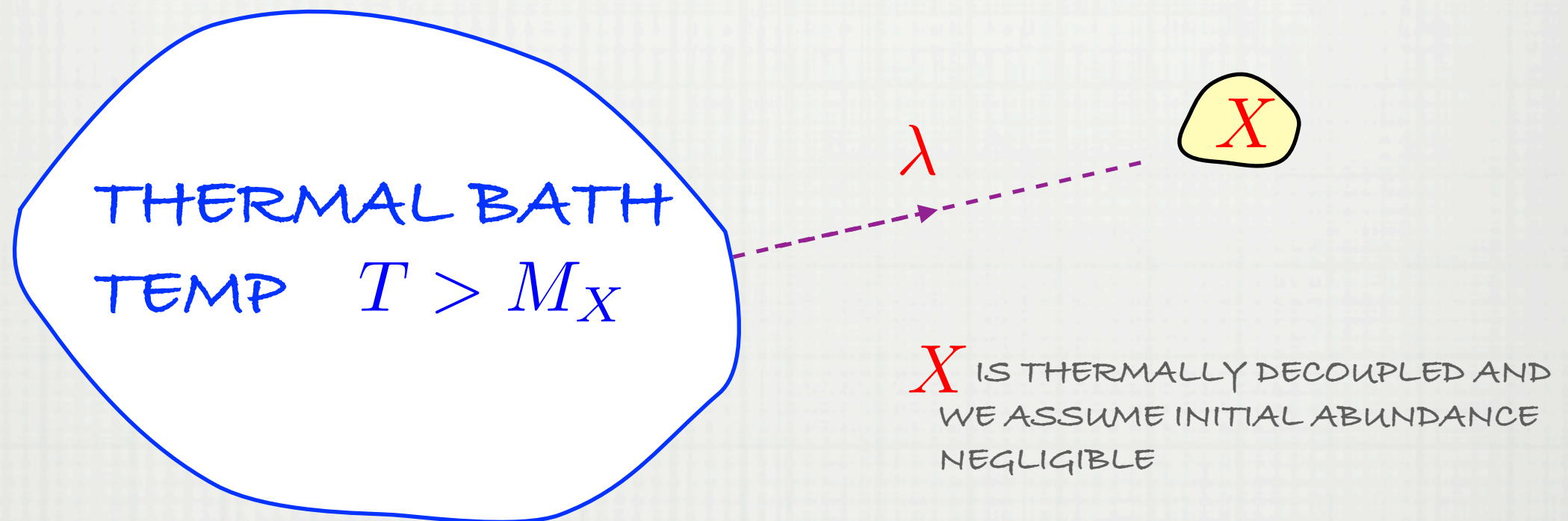


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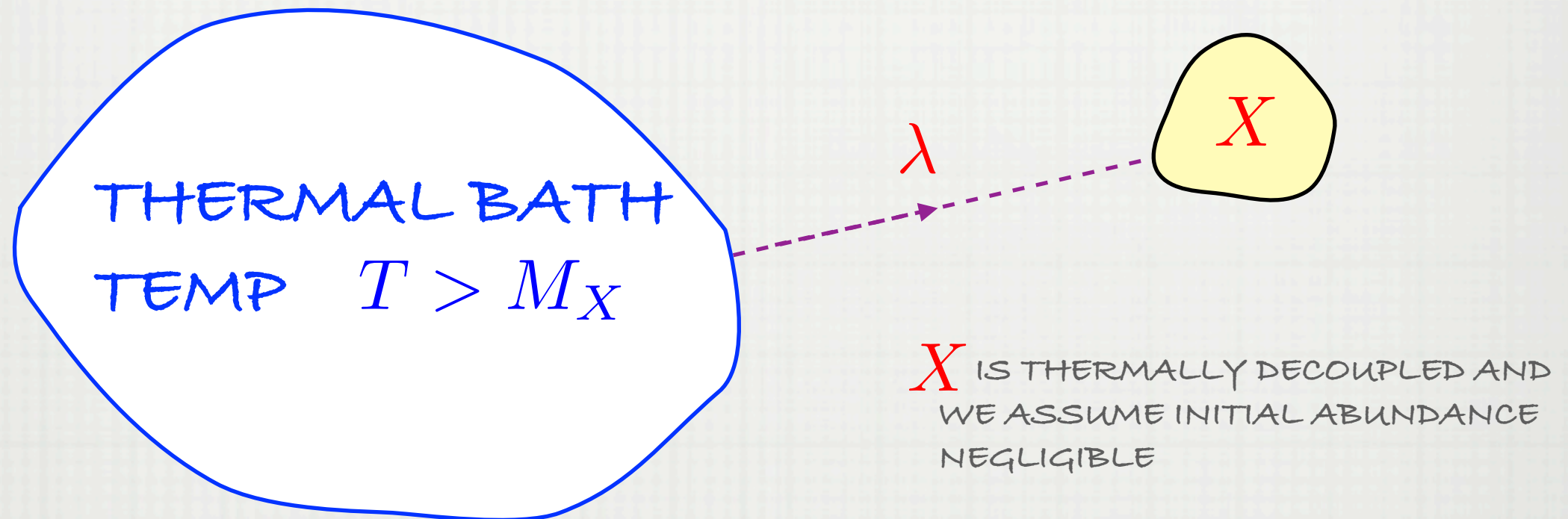


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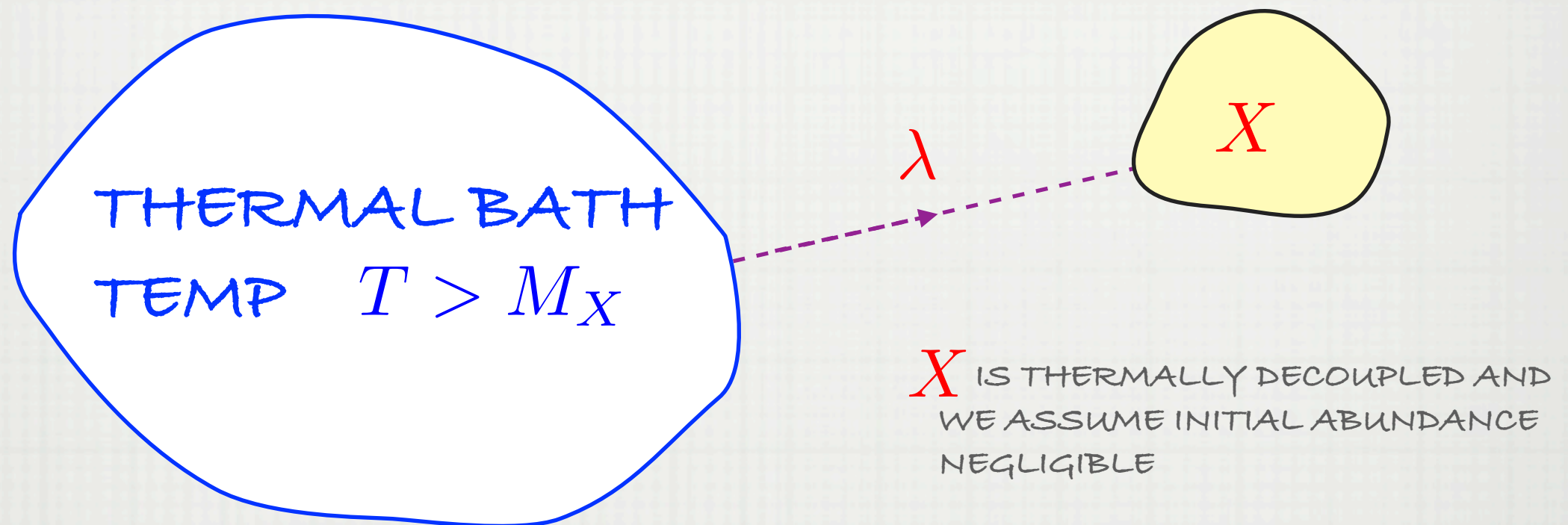


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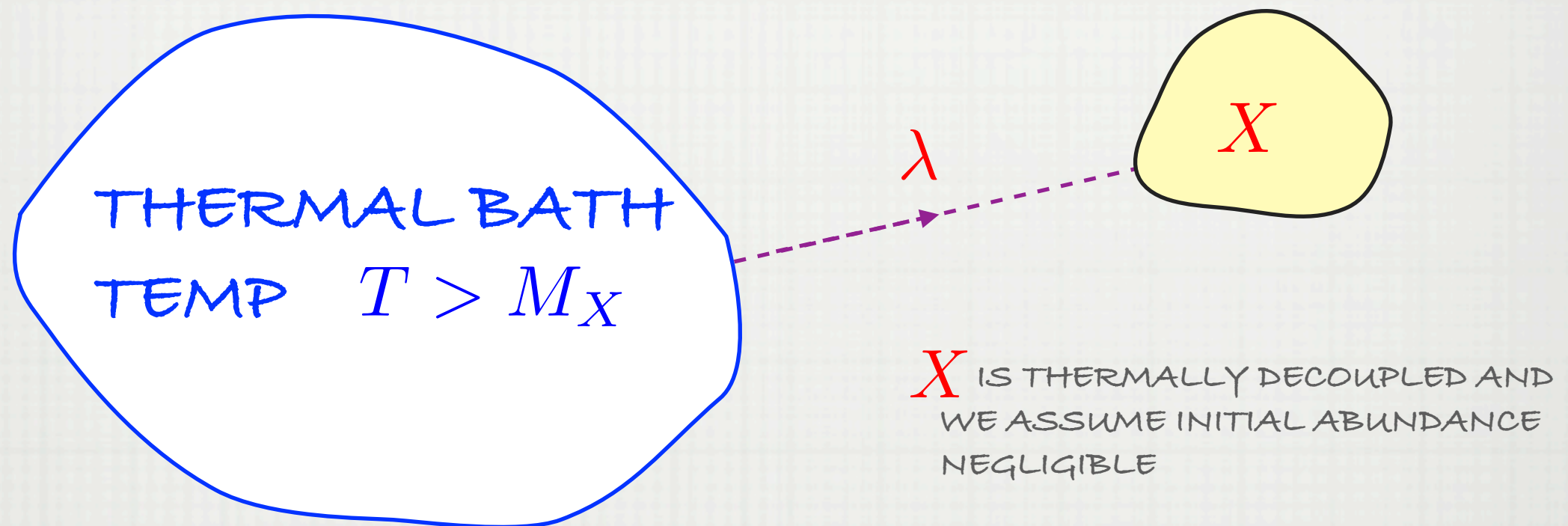


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- ALTHOUGH INTERACTIONS ARE FEEBLE THEY LEAD TO SOME X PRODUCTION
- DOMINANT PRODUCTION OF X OCCURS AT $T \sim M_X$ - IR DOMINANT
- INCREASING THE INTERACTION STRENGTH INCREASES THE YIELD

OPPOSITE TO FREEZE-OUT

FREEZE-OUT VS FREEZE-IN

$$Y_{FO} \sim \frac{1}{\langle \sigma v \rangle M_{Pl} m'}$$

USING $\langle \sigma v \rangle \sim \lambda'^2 / m'^2$

$$Y_{FO} \sim \frac{1}{\lambda'^2} \left(\frac{m'}{M_{Pl}} \right)$$

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FREEZE-IN VIA 2-2 SCATTERING,
DECAYS OR INVERSE DECAYS

COUPLING STRENGTH λ

m MASS OF HEAVIEST
PARTICLE IN INTERACTION

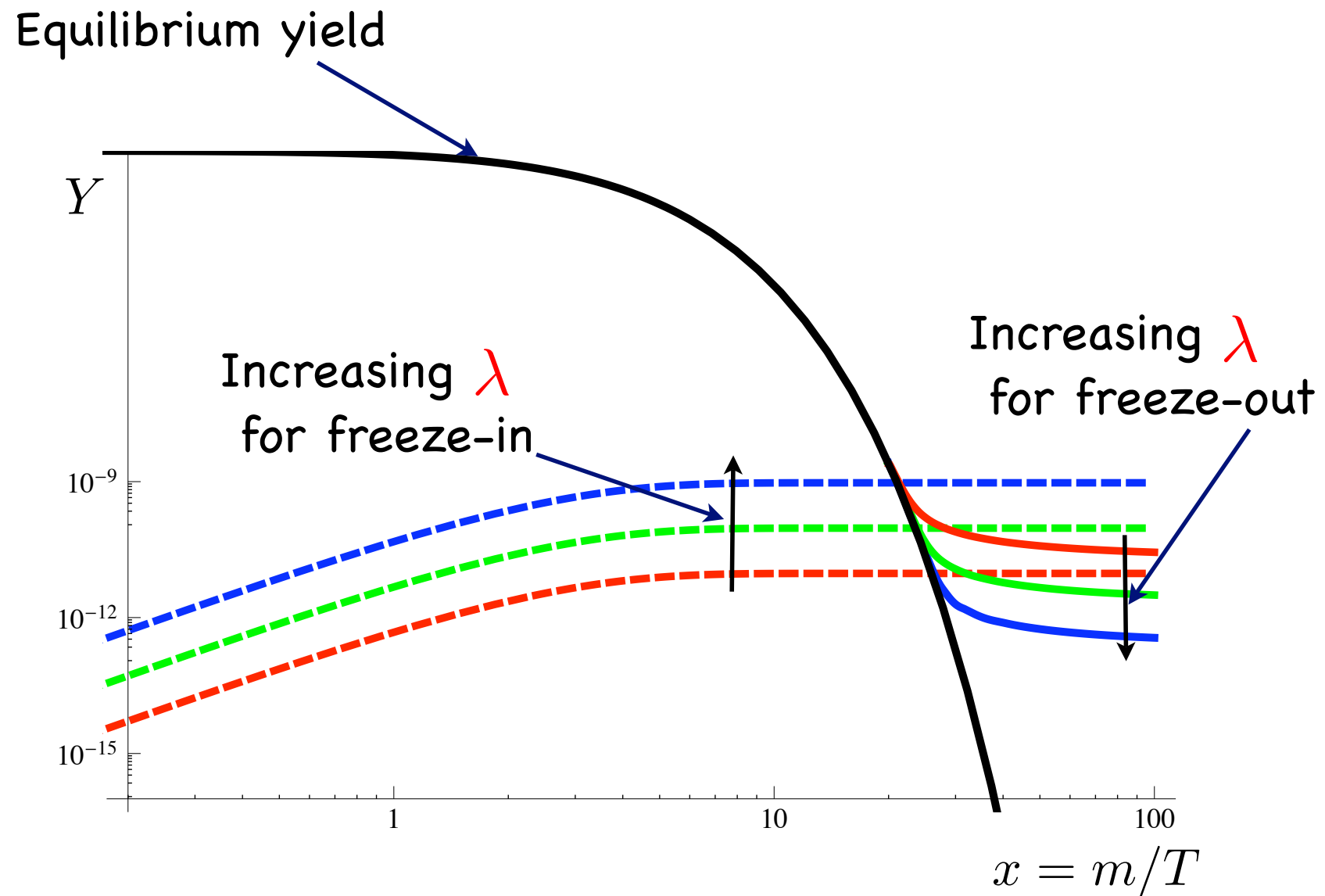
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FREEZE-OUT VS FREEZE-IN

□ AS TEMP DROPS BELOW MASS OF RELEVANT PARTICLE, DM ABUNDANCE IS HEADING **TOWARDS (FREEZE-IN)** OR **AWAY FROM (FREEZE-OUT)** THERMAL EQUILIBRIUM

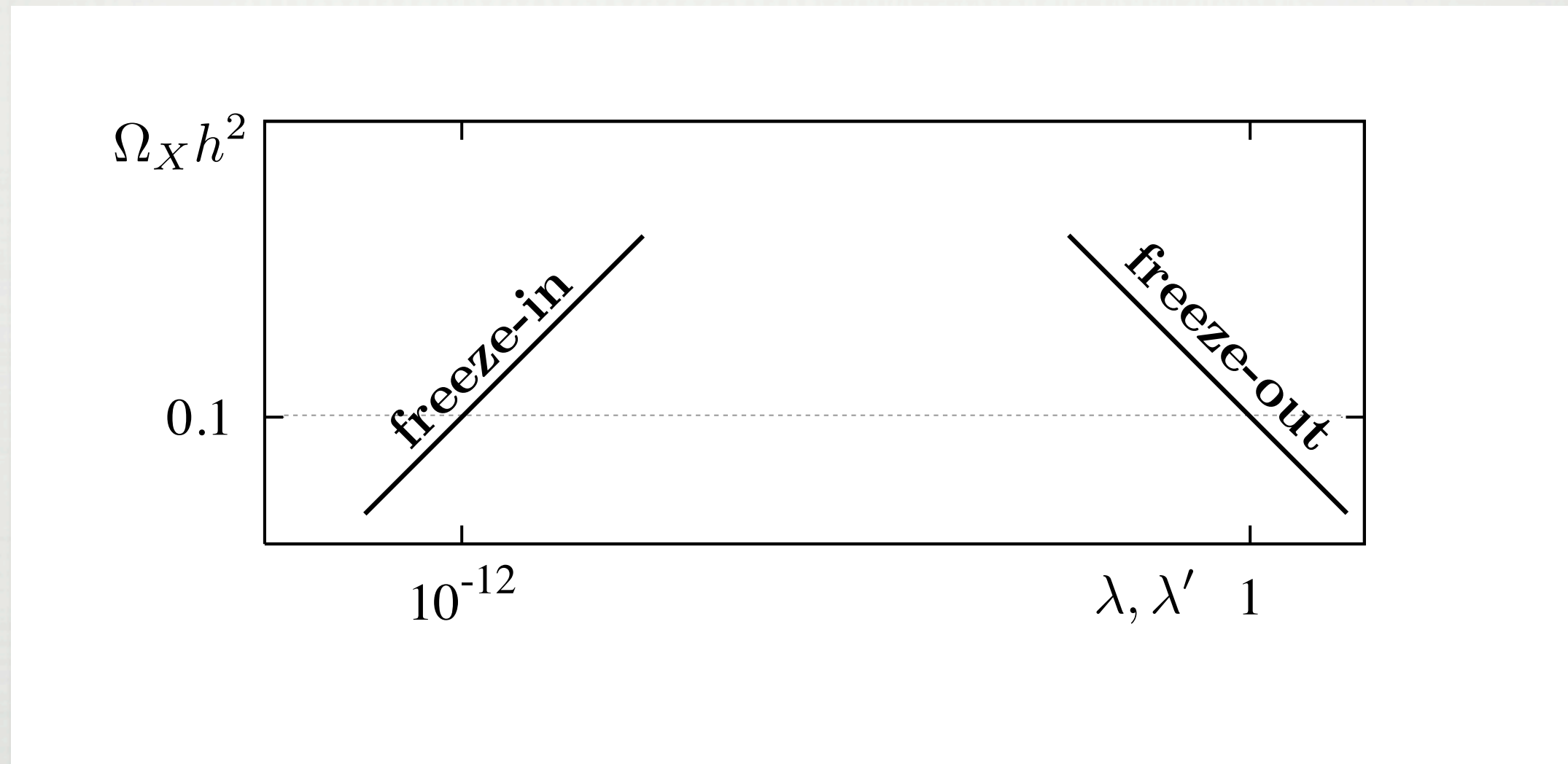
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FREEZE-OUT VS FREEZE-IN

□ FOR A TEV SCALE PARTICLE WE HAVE THE FOLLOWING PICTURE



FIMP MIRACLE VS WIMP MIRACLE

□ WIMP MIRACLE IS THAT FOR $m' \sim v$ $\lambda' \sim 1$

$$Y_{FO} \sim \frac{1}{\lambda'^2} \left(\frac{m'}{M_{Pl}} \right) \sim \frac{v}{M_{Pl}}$$

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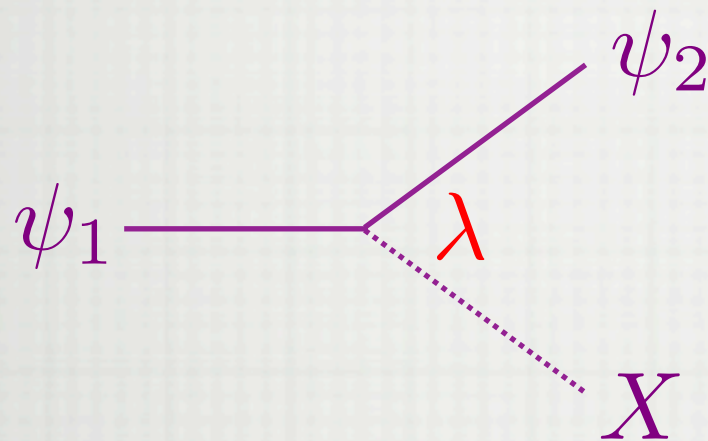
□ CONSIDER FIMP X COUPLED TO TWO BATH FERMIONS, FIMP IS LIGHTEST STATES CARRYING SOME STABILISING SYMMETRY - FIMP IS DM

$$L_Y = \lambda \psi_1 \psi_2 X \quad m_{\psi_1} > m_X + m_{\psi_2}$$

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$$\Omega_X h^2 \sim 10^{24} \frac{m_X \Gamma_{\psi_1}}{m_{\psi_1}^2}$$

ABUNDANCE GOES AS λ^2

CORRECT
ABUNDANCE

FOR $m_X \sim m_{\psi_1}$
 $\Rightarrow \lambda \sim 10^{-11}$

□ GIVES LONG LIVED DECAYS AT LHC, IMPLICATIONS FOR BBN

□ MASS OF FIMP DOES NOT HAVE TO BE SAME SCALE AS BATH PARTICLES, COULD HAVE MUCH SMALLER MASS

EXAMPLE MODEL II

- CONSIDER FIMP X COUPLED TO TWO BATH FERMIONS $B_1 B_2$
- AGAIN ASSUME FIMP IS LIGHTEST PARTICLE UNDER SOME STABILISING SYMMETRY - FIMP IS DM
- CONSIDER SOME QUARTIC INTERACTION OF FIMP WITH TWO BATH SCALARS

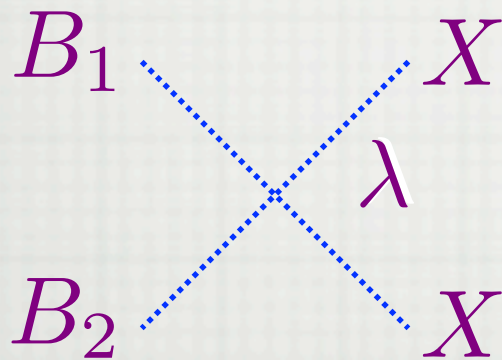
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$$\mathcal{L}_Q = \lambda X^2 B_1 B_2$$

ASSUMING

$$m_X \gg m_{B_1}, m_{B_2}$$



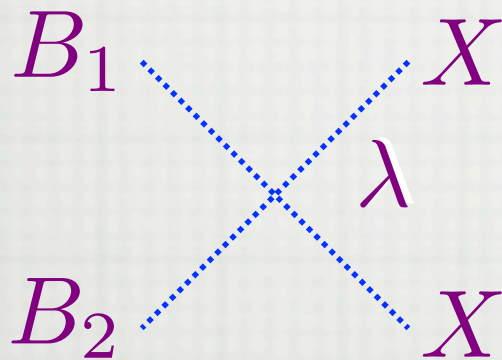
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$$\Omega h_X^2 \approx 10^{21} \lambda^2$$

FOR CORRECT ABUNDANCE $\Rightarrow \lambda \sim 10^{-11}$

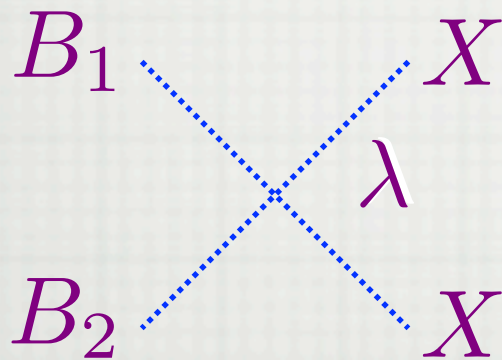
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- NOTE: ABUNDANCE IN THIS CASE IS INDEPENDENT OF THE FIMP MASS, FIMPZILLA?

EXPERIMENTAL IMPLICATIONS

- LONG LIVED "LOSPS" AT THE LHC: FIMPS FROZEN-IN BY DECAY OF LOSP
 - LOSP PRODUCED AT LHC WILL BE LONG LIVED
- LOSP COULD BE CHARGED

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$$\tau_{\text{LOSP}} = 7.7 \times 10^{-3} \text{sec} \left(\frac{m_X}{100 \text{ GeV}} \right) \left(\frac{300 \text{ GeV}}{m_{\text{LOSP}}} \right)^2 \left(\frac{10^2}{g_*(m_{\text{LOSP}})} \right)^{3/2}$$

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- SIGNALS FOR BBN: FIMPS AND LOSPS DECAYING LATE
- ENHANCED INDIRECT AND DIRECT DETECTION: RELIC ABUNDANCE NO LONGER LINKED TO DM ANNIHILATION RATE

ASYMMETRIC FREEZE-IN

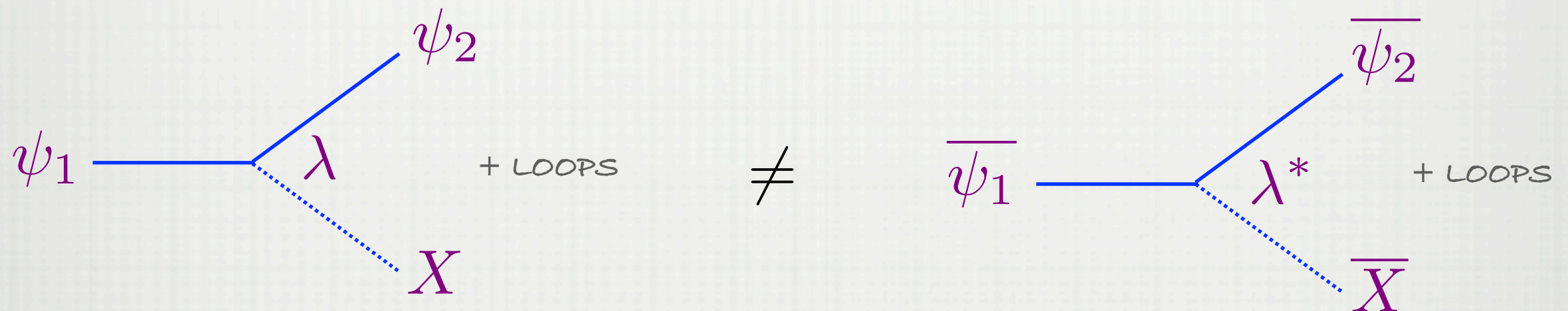
HALL, MARCH-RUSSELL, SMW
ARXIV: 1010.0245 [HEP-PH]

- WE CAN INTRODUCE CP AND B-L VIOLATION IN THE DECAYS THAT FREEZE-IN OUR DARK MATTER

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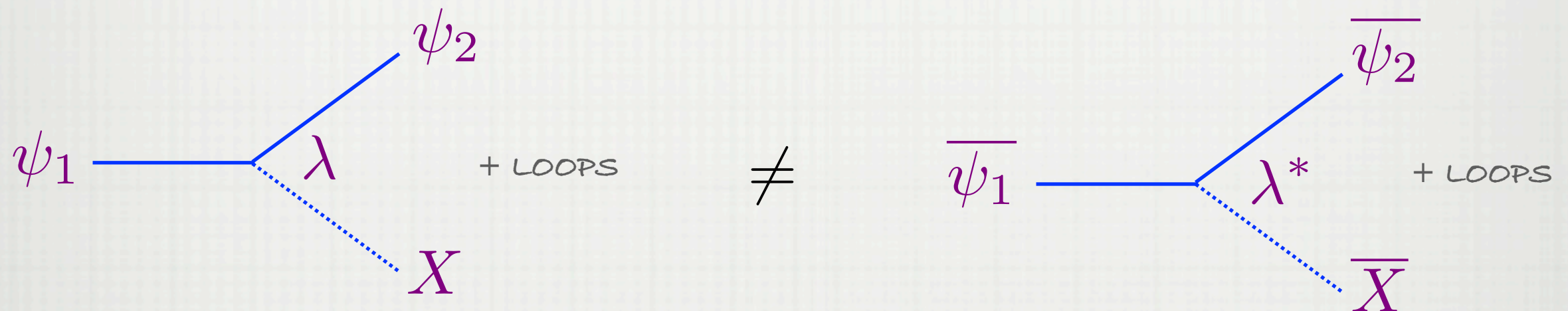
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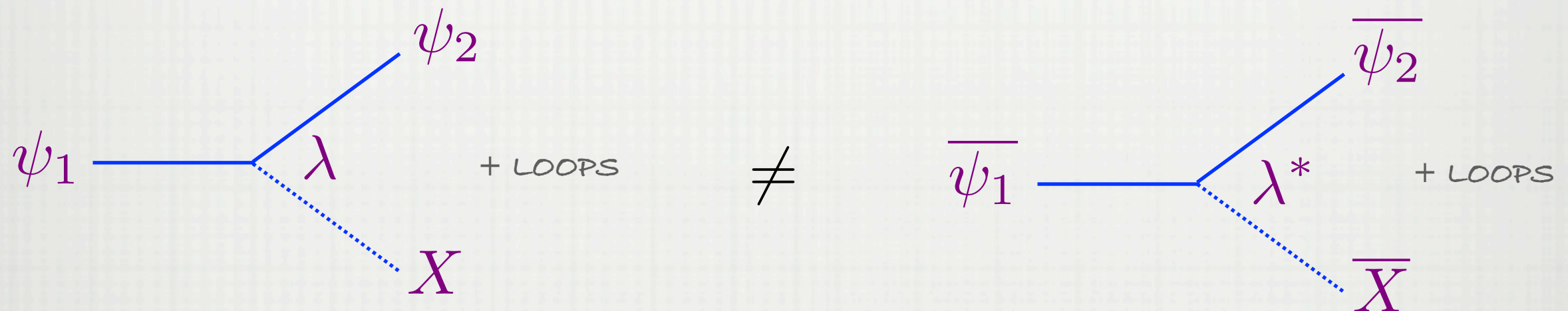


$$\Gamma(\psi_1 \rightarrow \psi_2 X) \neq \Gamma(\bar{\psi}_1 \rightarrow \bar{\psi}_2 \bar{X})$$

ASYMMETRIC FREEZE-IN

HALL, MARCH-RUSSELL, SMW
ARXIV: 1010.0245 [HEP-PH]

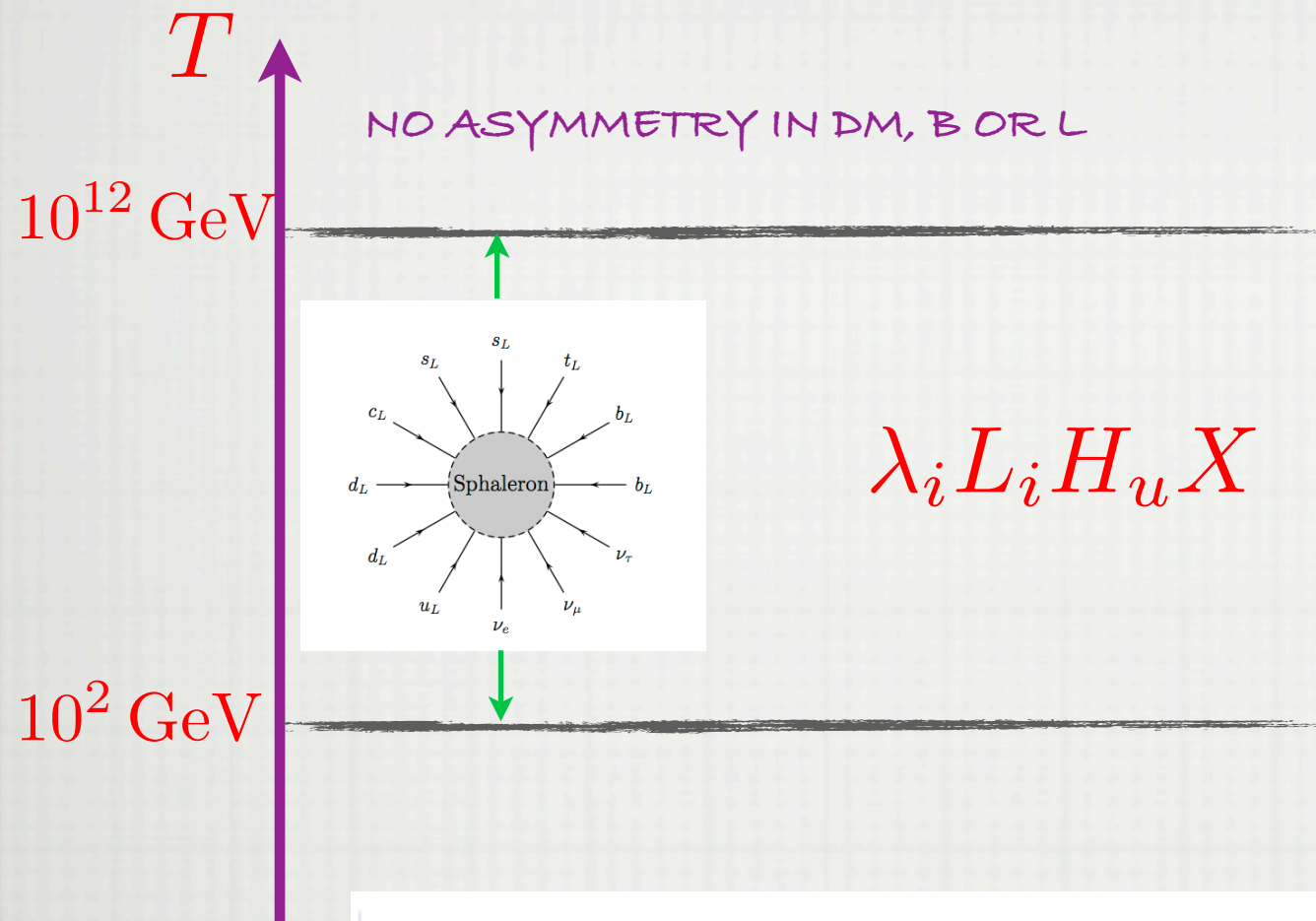
- WE CAN INTRODUCE CP AND B-L VIOLATION IN THE DECAYS THAT FREEZE-IN IN OUR DARK MATTER



$$\Gamma(\psi_1 \rightarrow \psi_2 X) \neq \Gamma(\bar{\psi}_1 \rightarrow \bar{\psi}_2 \bar{X})$$

- WE NEED CP VIOLATION (AND LOOP DIAGRAMS TO INTERFERE WITH THE TREE LEVEL DIAGRAMS)

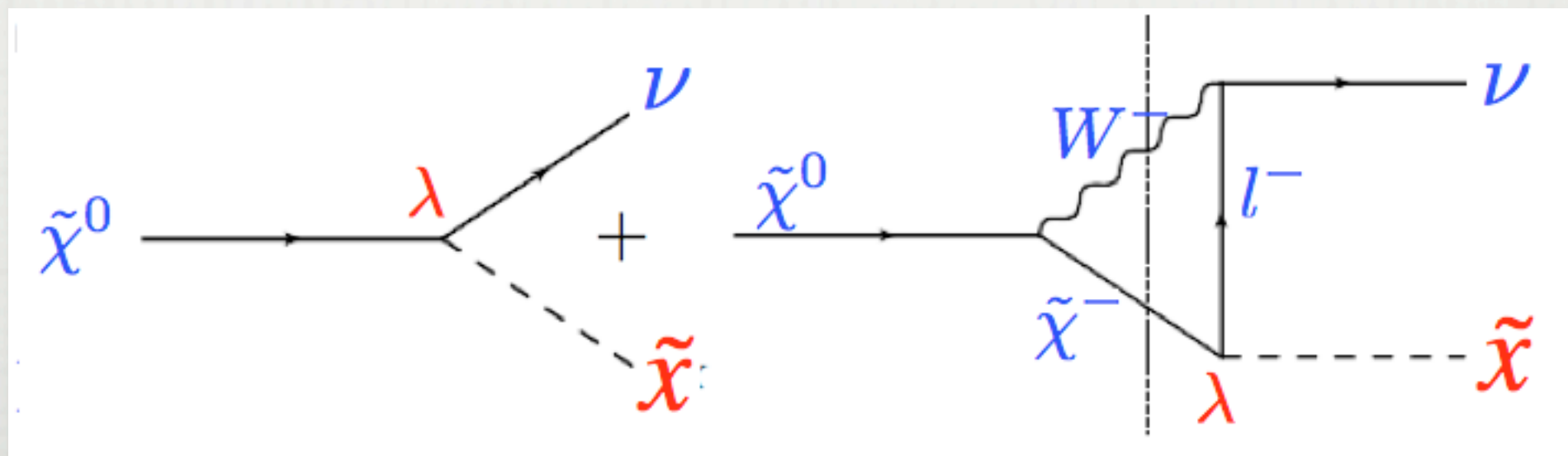
ASYMMETRIC FREEZE-IN EXAMPLE



$$\lambda_i L_i H_u X$$

THIS OPERATOR NOW HAS A SMALL COUPLING AND IS RESPONSIBLE FOR THE ASYMMETRY

HAS A SYMMETRY $U(1)_{B-L+X}$



ASYMMETRIC FREEZE-IN EXAMPLE

- THESE PROCESSES ALREADY CONTAIN OUT-OF-EQUILIBRIUM PROCESSES - FIMP IS NOT IN THERMAL EQUILIBRIUM, IN FACT ALL YOU NEED IS A DIFFERENCE IN TEMPERATURE BETWEEN FIMP AND SM SECTOR

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$$\epsilon = (\text{loop factor}) \sin \phi \quad \Gamma_{\chi^0} \sim \frac{\lambda^2 m_{\chi^0}}{8\pi}$$

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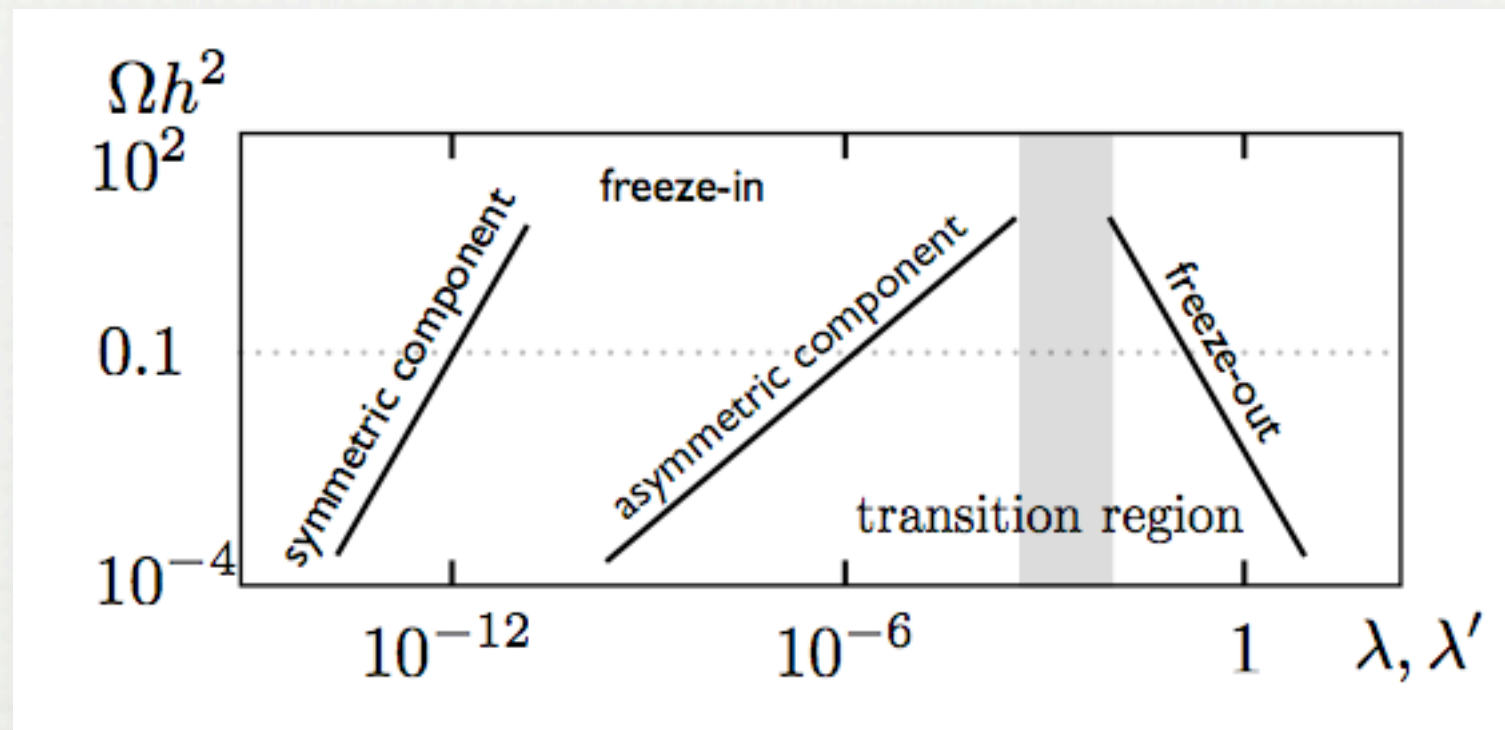
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- ASYMMETRY APPEARS AT λ^2
- CP VIOLATION COULD COME FROM GAUGINO - HIGGSINO SECTOR

ASYMMETRIC FREEZE-IN

ASYMMETRIC FREEZE-IN

□ TURNS OUT, THROUGH NON-TRIVIAL CANCELLATIONS IN THE BOLTZMANN EQUATIONS THE ASYMMETRY APPEARS AT λ^3 HOOK, ARXIV:1105.3728



□ MAKES THE MODEL VERY PREDICTIVE - NOT MUCH PARAMETER SPACE

DEPENDING ON THE MODEL, ASYMMETRIC FREEZE-IN MAY ALLOW "FULL" PROBE OF BARYOGENESIS - DM CONNECTION