

CKM GAMMA FROM TREE DECAYS

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TAKE HOME MESSAGE

- γ from $B \rightarrow DK$ is theoretically extremely clean
- it is an interesting observable to measure precisely

OUTLINE

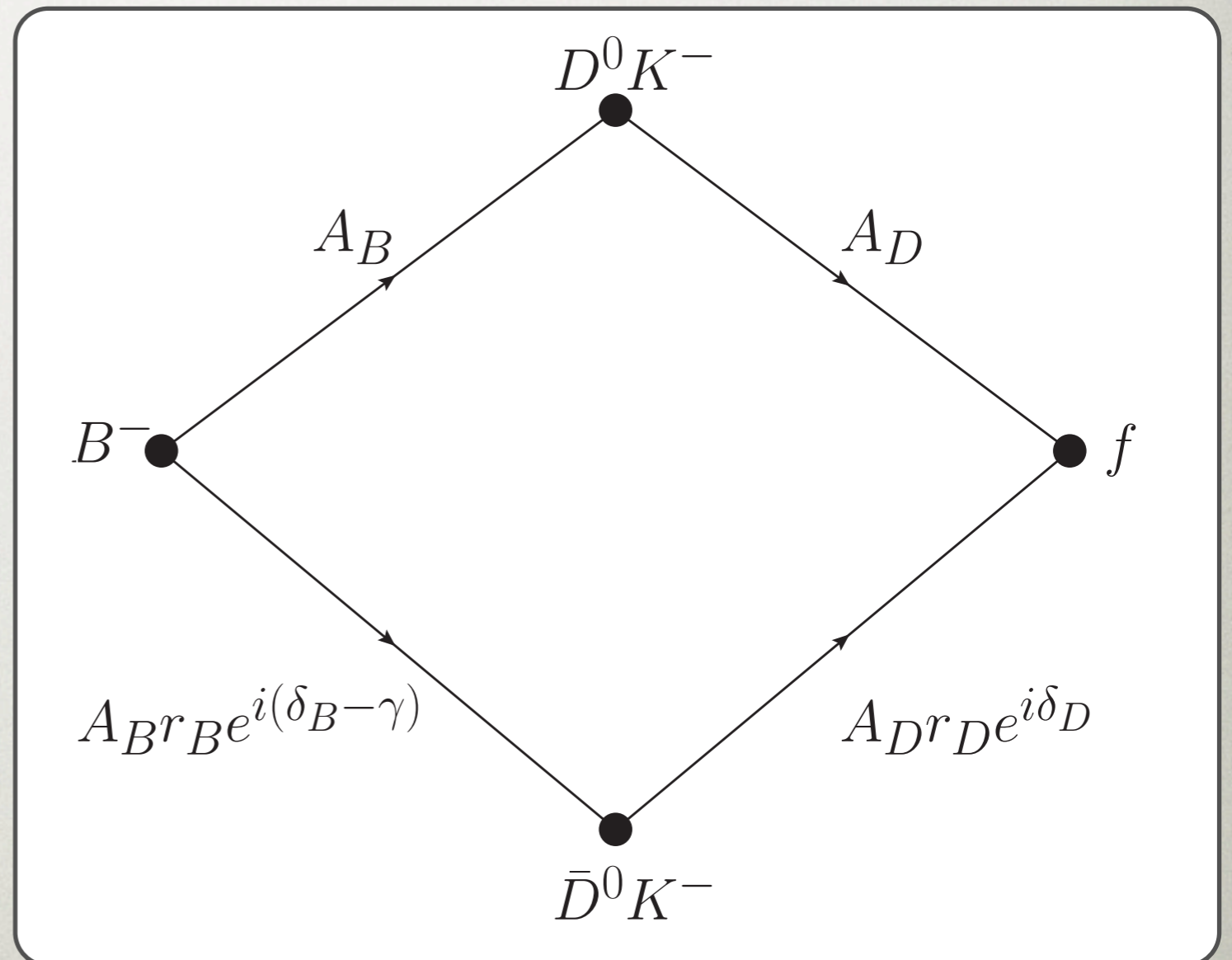
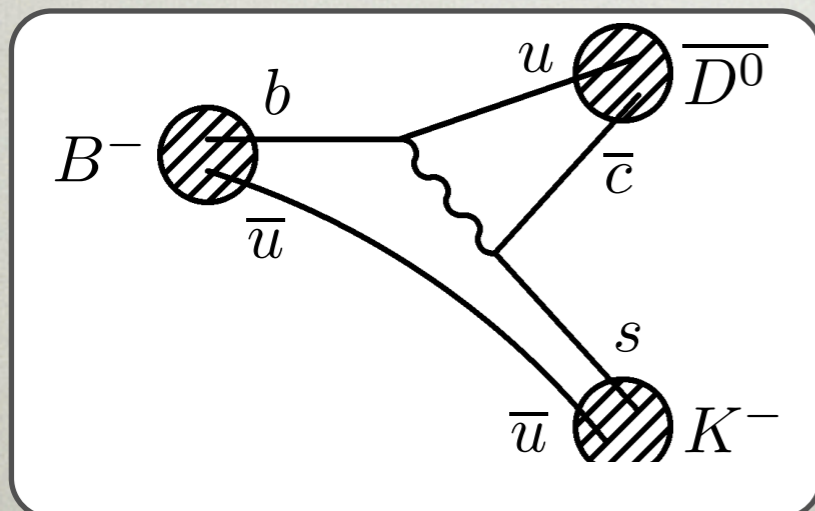
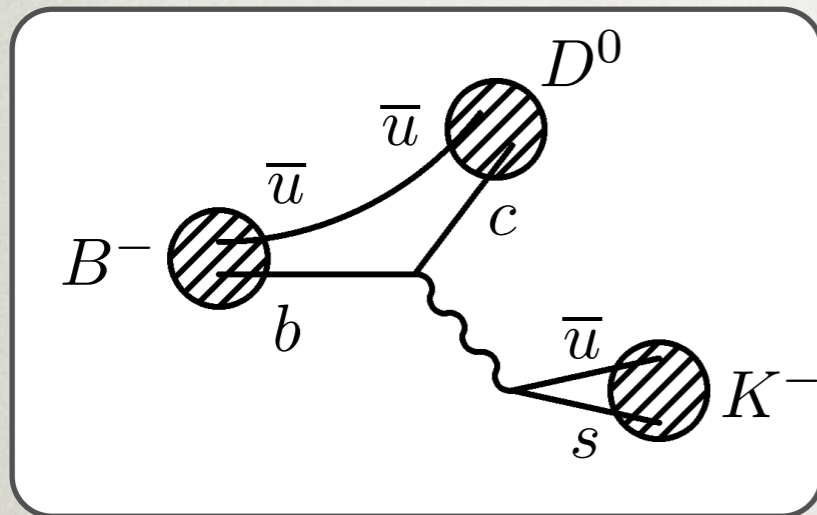
- brief outline of the $B \rightarrow DK$ methods
- what is the ultimate theory error?
 - electroweak corrections
- “the ultimate test of (MFV) NP”

Pirjol, JZ, to appear

OBTAINING GAMMA

- use interference between $b \rightarrow c\bar{u}s$ and $b \rightarrow u\bar{c}s$

Gronau, Wyler, 1991; Gronau, London, 1990



MANY FINAL STATE CHOICES

see also talk by Stefania Ricciardi

- possible choices for final state f in D decay
 - CP- eigenstate (e.g. $K_S\pi^0$) [Gronau, London, Wyler \(1990,1991\)](#)
 - flavor state (e.g. $K^+\pi^-$) [Atwood, Dunietz, Soni \(1997\)](#)
 - singly Cabibbo suppressed (e.g. $K^{*+}K^-$) [Grossman, Ligeti, Soffer \(2002\)](#)
 - many-body final state (e.g. $K_S \pi^+\pi^-$) [Giri, Grossman, Soffer, JZ \(2003\)](#)
[Poluektov et al. \[Belle\] \(2004\)](#)
- other extensions:
 - many body B final states: $B^+ \rightarrow DK^+\pi^0$, $B^0 \rightarrow D\pi^-K^+$
[Aleksan, Petersen, Soffer \(2002\)](#), [Gershon \(2008\)](#), [Gershon, Poluektov \(2009\)](#)
 - use D^{0*} in addition to D^0 [Bondar, Gershon \(2004\)](#)
 - use self tagging D^{0**} , D_2^{*-}
[Sinha \(2004\)](#) [Gershon \(2008\)](#)
 - neutral B_d, B_s decays (time dep., time-integr., self-tag)
[Aleksan, Dunietz, Kayser \(1992\)](#), [Kayser, London \(2000\)](#), [Atwood Soni \(2003\)](#), [Fleischer \(2003\)](#),
[Gronau et al \(2004\)](#)

WHY THE METHODS WORK

- many methods: GLW, ADS, Dalitz,...
- one has $\sim N_D N_B$ measurements, but $\sim N_D + N_B$ unknowns
 - \Rightarrow can determine γ
- does it make sense to split into “methods”?
 - we are really interested in γ
 - combined analysis wins
- “only” benefit of splitting: compare diff. γ
 - check for NP or systematics

WHY MEASURE GAMMA?

WHY MEASURE GAMMA?

- now (= in the age of LHCb and upcoming SFF)
 - theoretically clean extraction of the CKM weak phase
 - standard candle of the CKM
- unlikely there is NP in γ (at present precision)
 - search for NP by comparing to other observables
 - can test for consistency in γ extraction in $B \rightarrow DK$ itself

TEST FOR NP IN DECAY AMPLITUDES

- extraction of γ has a built in test for presence of extra NP in decay ampl.

$$A(B^- \rightarrow f_D K^-) \propto r_D e^{i\delta_D} + r_B e^{i\delta_B - \gamma} + r'_B e^{i\delta'_B - \gamma'}$$

$$A(B^+ \rightarrow f_D K^+) \propto r_D e^{i\delta_D} + r_B e^{i\delta_B + \gamma} + r'_B e^{i\delta'_B + \gamma'}$$

- thus for B^+ and B^- different r_B

$$r_{B^+} \rightarrow |r_B e^{i\delta_B + \gamma} + r'_B e^{i\delta'_B + \gamma'}|; \quad r_{B^-} \rightarrow |r_B e^{i\delta_B - \gamma} + r'_B e^{i\delta'_B - \gamma'}|$$

TEST OF DIRECT CPV NP IN $B \rightarrow DK$

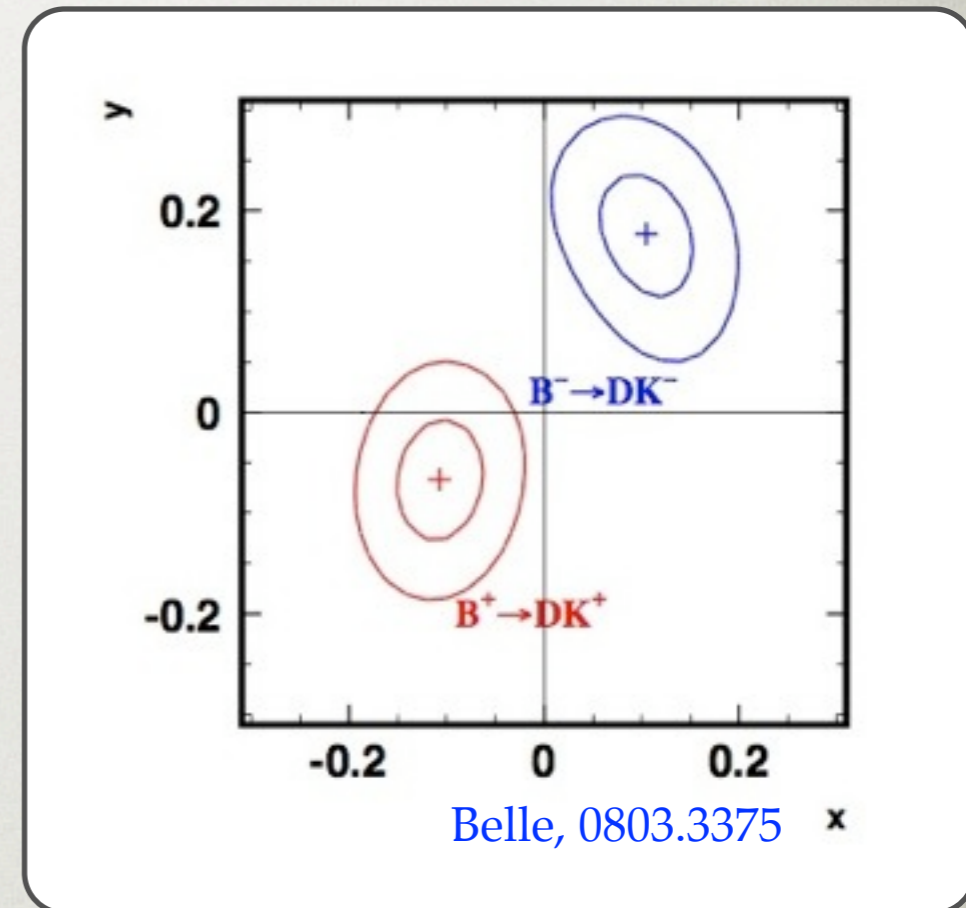
- there is NP in $B \rightarrow DK$ amplitude if

$$r_{B^-} \neq r_{B^+}$$

- Belle and Babar already measure this

$$\begin{aligned}x_{\pm} &= r_B \cos(\gamma \pm \delta_B) \\y_{\pm} &= \pm r_B \sin(\gamma \pm \delta_B)\end{aligned}$$

- even, if $x_+^2 + y_+^2 = x_-^2 + y_-^2$ still possible that γ is shifted



IN THE (FAR) FUTURE

- another test: γ from $B^\pm \rightarrow DK^\pm, B^\pm \rightarrow DK^{*\pm}, B^\pm \rightarrow D^*K^\pm, B^0 \rightarrow DK^0, \dots$ all need to coincide!
- NP with contributions of different chirality could for instance give different shifts in γ
- since the extraction of gamma theoretically clean
 - can be used to search for high scale NP when a lot of statistics
- next: what is the scale we could in principle probe?
 - how clean? what is the theory error?

THEORY ERRORS ON EXTRACTING GAMMA

THEORY ERRORS

- assume SM
- several sources that can induce a shift in γ
- most can be avoided
 - with more statistics
 - example: Dalitz plot theory error
 - by modifying equations
 - example: errors from $D - \bar{D}$ mixing
- remain: errors from electroweak corrections

$D - \bar{D}$ MIXING

- in SM $D - \bar{D}$ mixing is CP conserving \Rightarrow the effect is small
- if D decay info. is from flavor tagged D (i.e. from $D^* \rightarrow D \pi$)
 - then only important changes are in the interf. term
 - change in relative phase: $\delta_f \rightarrow \langle \delta_f \rangle$ Grossman, Soffer, JZ, 2005
 - dilutes the interference: $\dots \rightarrow \dots \times e^{-\varepsilon}$
- the effect on γ is $\varepsilon \sim \mathcal{O}(x_D^2/r_f^2, y_D^2/r_f^2)$
 - applies e.g. to GLW, ADS
 - even for doubly Cabibbo supp. D decays the shift $\Delta\gamma < 1^\circ$
- in model indep. Dalitz analysis no changes needed, if everything from $B \rightarrow DK$
 - one already fits for both $\langle \delta_f \rangle$ and ε by fitting for c_i, s_i

$D - \bar{D}$ MIXING

Bondar, Poluektov, Vorobiev, 1004.2350

- the effect potentially larger, if D decay info from CLEO ($\psi(3770) \rightarrow D\bar{D}$)
- the change since time integr. interv.: $t \in (-\infty, \infty)$
 - the shift in γ is now linear in x_D, y_D
 - but still small: $\Delta\gamma \leq 2.9^\circ$ ($\leq 0.2^\circ$, if $|A_D|^2$ info. comes from $D^* \rightarrow D\pi$)
- most importantly: $D - \bar{D}$ mixing effects can be incl. exactly if x_D, y_D precisely measured

OTHER ERRORS

- for γ from (untagged) $B_s \rightarrow D\phi$ the inclusion of $\Delta\Gamma_s$ depen. important
 - $\Delta\Gamma_s$ needs to be well measured
Gronau, Grossman, Soffer, Surujon, JZ, 2007
- QED radiative corrections
 - CP conserv., in principle no effect on γ
- the remaining (SM) theory error from
 - higher electroweak corrections

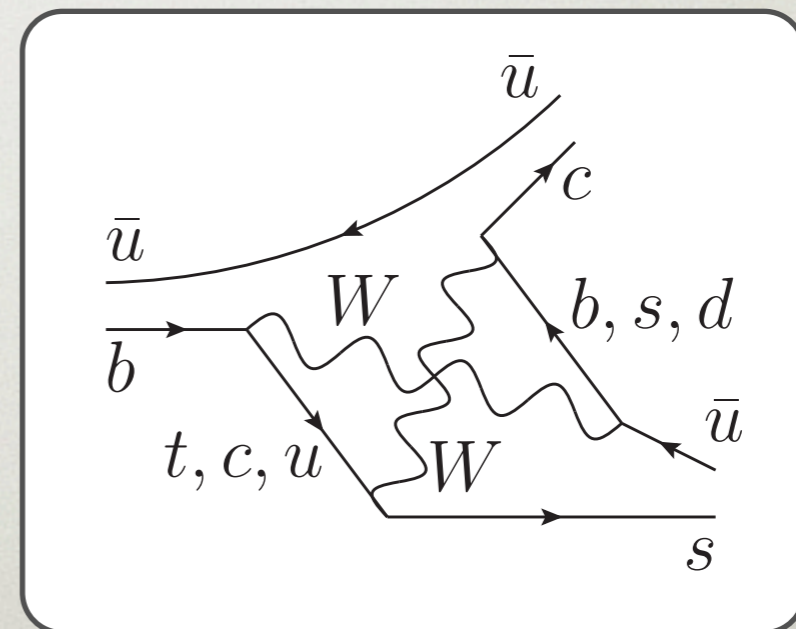
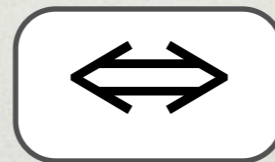
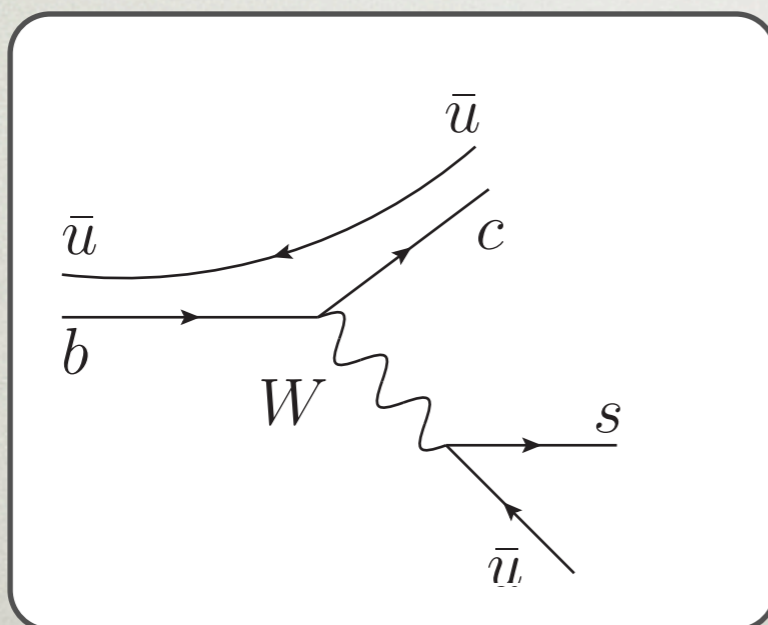
IRREDUCIBLE THEORY

ERROR ON GAMMA

- irreduc. theory error in SM introduced by ew. corrections that change CKM structure

Pirjol, JZ, to appear

- if only vertex corrections no effect on γ extr.
- no effect from Z exchange
- there is effect from box diagrams



IRREDUCIBLE THEORY

ERROR ON GAMMA

- the shift on γ due to the box diagram
- dominant contrib. effectively due to t and b in the loop

- $b \rightarrow us\bar{c}$:

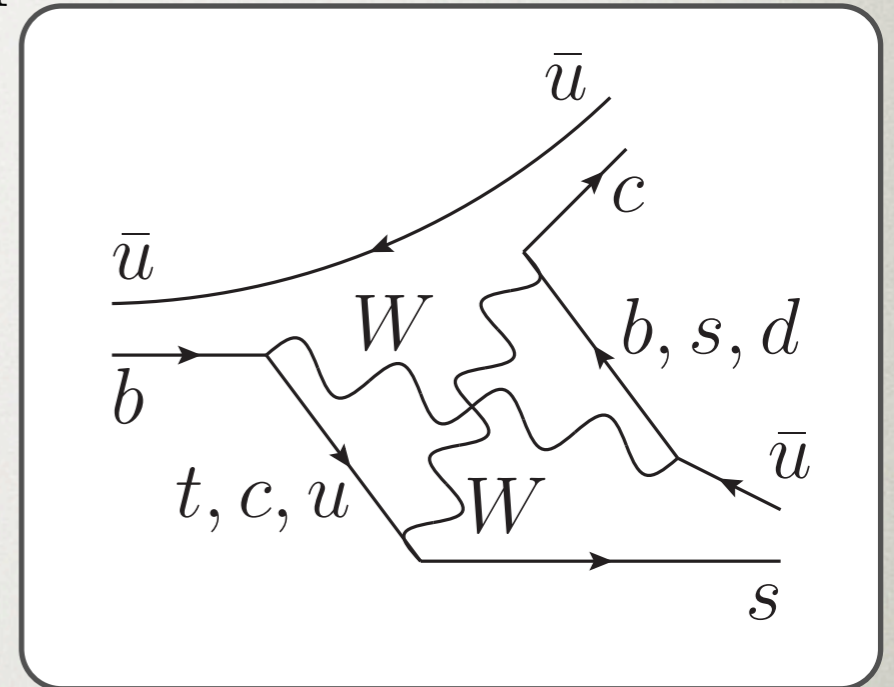
- tree level $\sim V_{ub}V_{cs}^*$
- box diagram $\sim (V_{tb}V_{ts}^*)(V_{ub}V_{cb}^*)$

- same weak phase, does not induce $\delta\gamma$

- $b \rightarrow cs\bar{u}$:

- tree level $\sim V_{cb}V_{us}^*$
- box diagram $\sim (V_{tb}V_{ts}^*)(V_{cb}V_{ub}^*)$

- different weak phase $\Rightarrow \delta\gamma \neq 0$



IRREDUCIBLE THEORY ERROR ON GAMMA

Pirjol, JZ, to appear

- we estimate it in two ways
- integrate both t and b at the same time
 - local operator, but large logs
- resum $\log(m_b/m_W)$ but also nonlocal contriibs
 - keep only local ones
- the precision suffices for our purposes
- irreducible theory error on γ is
$$\delta\gamma/\gamma < O(10^{-6})$$
 - most likely even $\delta\gamma/\gamma \lesssim O(10^{-7})$

ULTIMATE TEST OF MFV

- how high NP scales can we probe using γ from $B \rightarrow DK$?
- assuming MFV: can probe $\Lambda \sim 10^2 \text{TeV}$
- assume gen. FV: can probe $\Lambda \sim 10^3 \text{TeV}$
- this is far future of course
 - $O(10^{18})$ $B\bar{B}$ pairs needed

SOME NUMBERS FOR FUN

JZ, 1101.0134

PROBE	Λ_{NP} for (N)MFV NP	Λ_{NP} for gen. FV NP	No. of BB pairs
γ from $B \rightarrow DK$	$\Lambda \sim O(10^2 \text{ TeV})$	$\Lambda \sim O(10^3 \text{ TeV})$	$\sim 10^{18}$
$B \rightarrow \tau \nu$ ¹⁾	$\Lambda \sim O(\text{TeV})$	$\Lambda \sim O(30 \text{ TeV})$	$\sim 10^{13}$
$b \rightarrow s s \bar{d}$	$\Lambda \sim O(\text{TeV})$	$\Lambda \sim O(10^3 \text{ TeV})$	$\sim 10^{15}$
β from $B \rightarrow J/\psi K_S$ ²⁾	$\Lambda \sim O(50 \text{ TeV})$	$\Lambda \sim O(200 \text{ TeV})$	$\sim 10^{12}$
K - K mixing ³⁾	$\Lambda > 0.4 \text{ TeV} (6 \text{ TeV})$	$\Lambda > 10^3 \text{ TeV} (10^4 \text{ TeV})$	now

1) assuming no err. on f_B , so that ultimate th. error just from ew. corr.

2) assuming pert. error estimates $\delta\beta/\beta \sim 0.1\%$

3) bounds for $\text{Re}C_1$ ($\text{Im} C_1$) from UTfitter 0707.0636

CONCLUSIONS

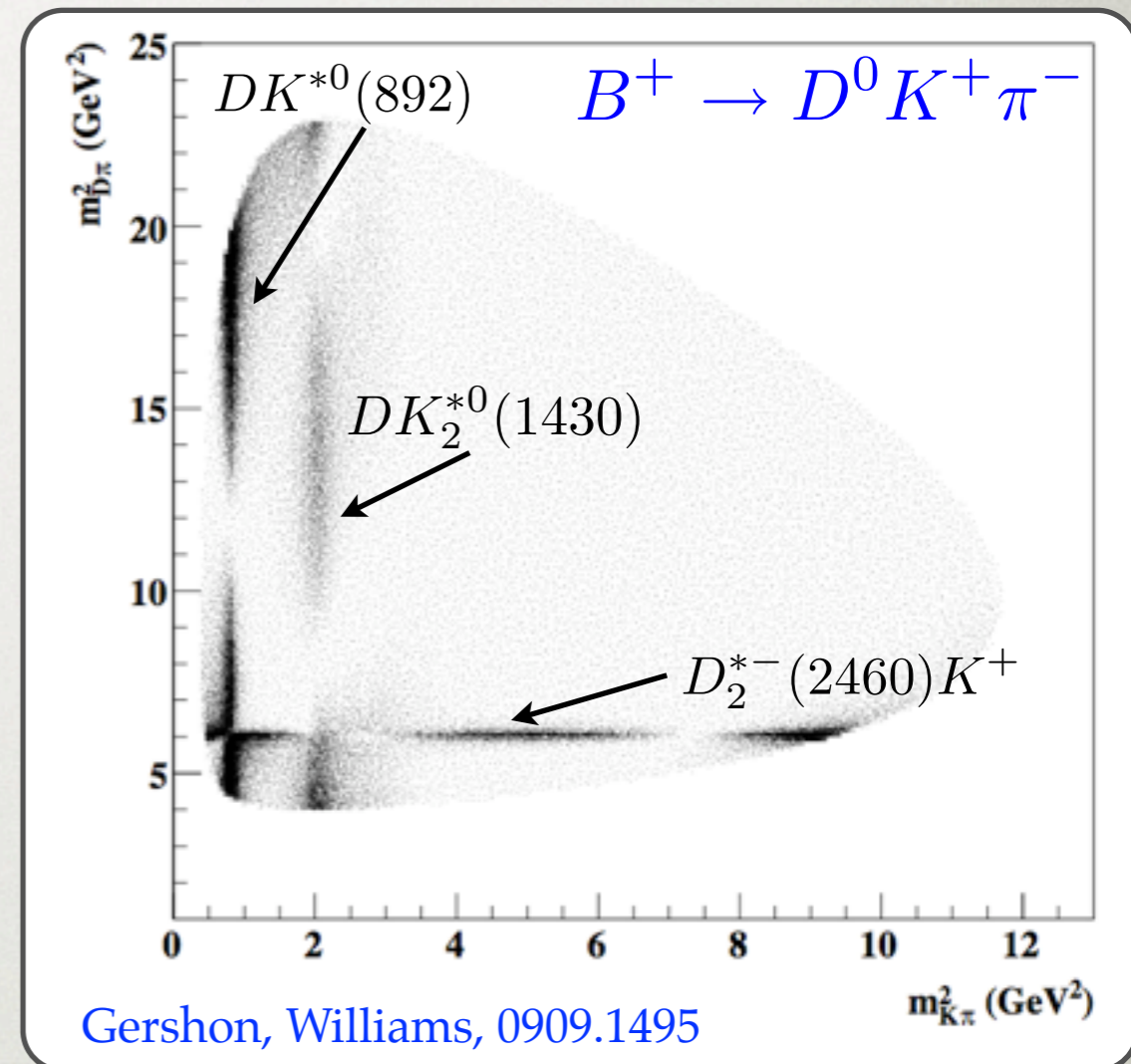
- γ extraction from $B \rightarrow DK$ is theoretically clean
 - irreducible theory error on γ is below $\delta\gamma/\gamma < 10^{-6}$
- measuring γ is important
 - standard candle of the SM
 - search for NP

BACKUP SLIDES

SOME SPECIFIC IDEAS FOR LHCb

- a “method”: a subset of final states allowing for extr. of γ
- multibody $B^0 \rightarrow DK^+\pi^-$
Gershon (2008) Gershon, Williams (2009)

- contains flavor specific
 $D_2^{*-}(2460) \rightarrow \bar{D}^0\pi^-$
- interf. with other resonances
 (e.g. $B^0 \rightarrow DK^{*0}$) gives γ
- many choices for $D \rightarrow f$ still
- equivalent of GLW does not need CP-odd $D \rightarrow K_S\pi^0$ decays (that is difficult for LHCb)



MORE ON

$$B^0 \rightarrow DK^+ \pi^-$$

- compared to quasi-two-body $B^0 \rightarrow DK^{*0}$
 - at least 50% better sensitivity to γ Gershon, Williams, 0909.1495
- extension of model indep. method possible
 - double Dalitz plot analysis $B^0 \rightarrow DK^+ \pi^-$
 $\rightarrow (K_S \pi^+ \pi^-) DK^+ \pi^-$ Gershon, Poluektov, 0910.5437
 - $B^0 \rightarrow DK^+ \pi^-$ Dalitz still poorly known
 - estimates using reasonable models: 20 annual yields of LHCb $\Rightarrow O(1^\circ)$ error