

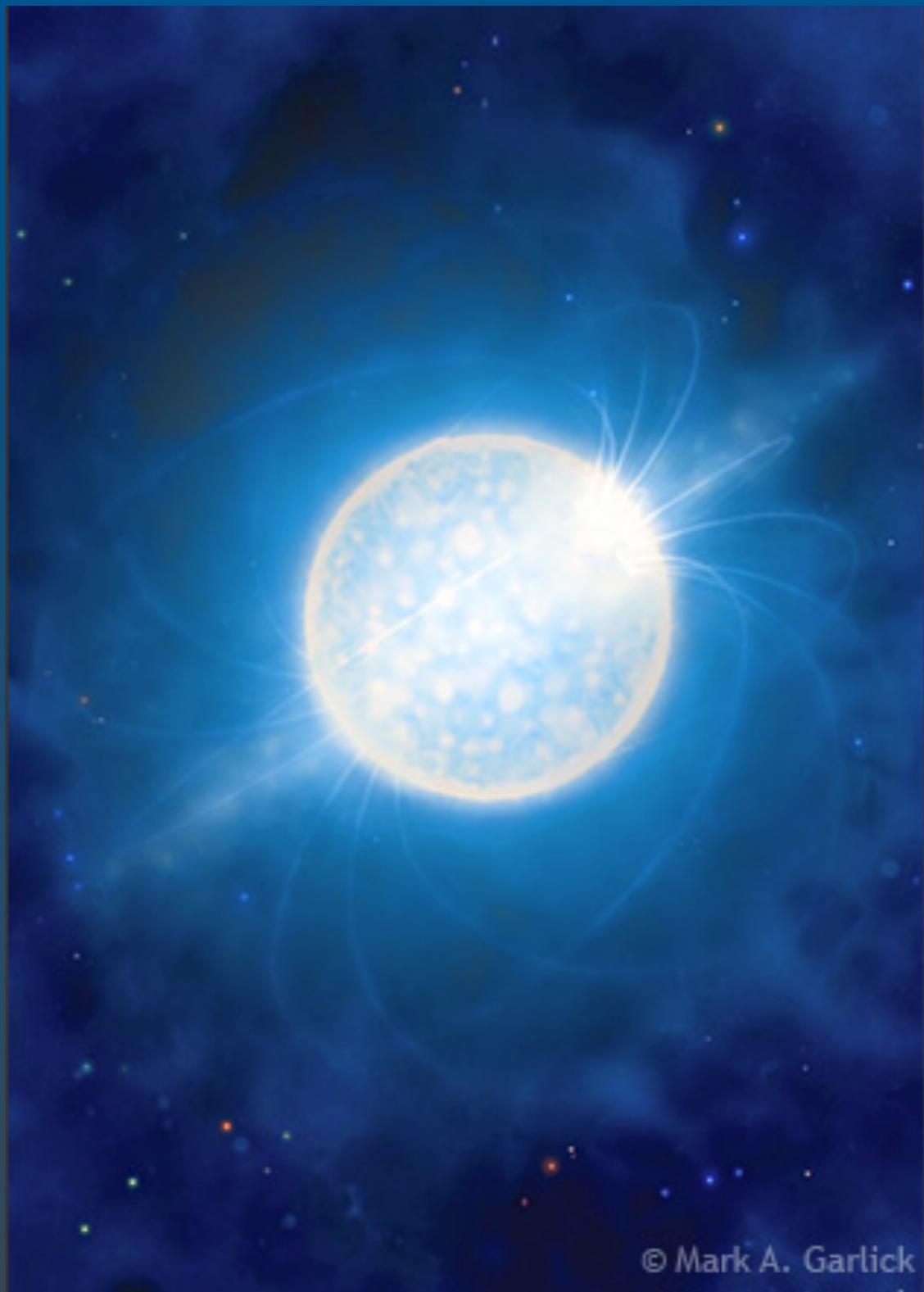
IMPLICATION OF RESULTS FROM HEAVY-ION EXPERIMENTS FOR COMPACT STARS

Debaratí Chatterjee

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Universität Heidelberg, Germany*

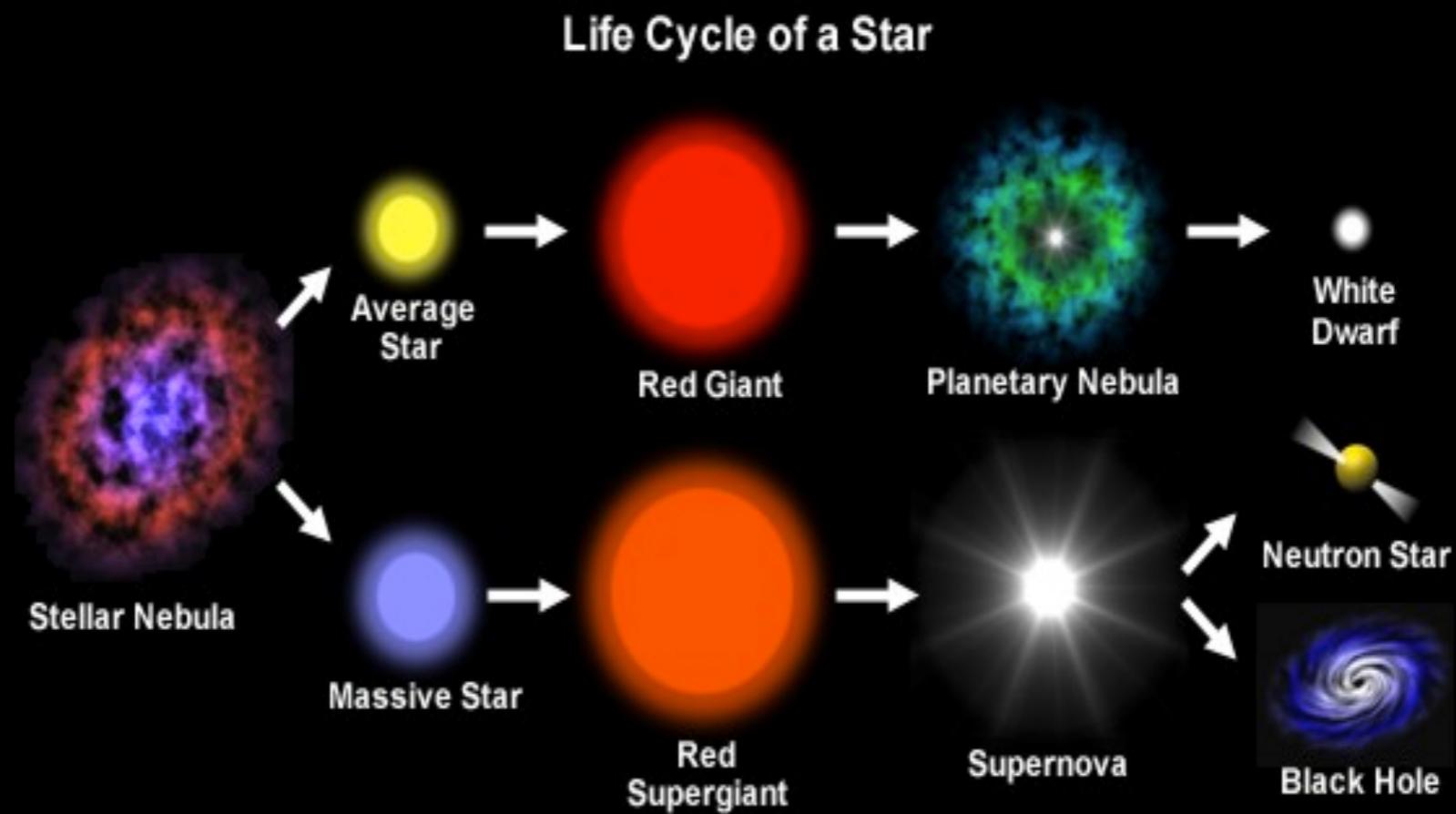
Collaborators: Jürgen Schaffner-Bielich
Simon Weissenborn

Irina Sagert, MSU
Laura Tolos, IEEC, Barcelona
Cristián Sturm, GSI



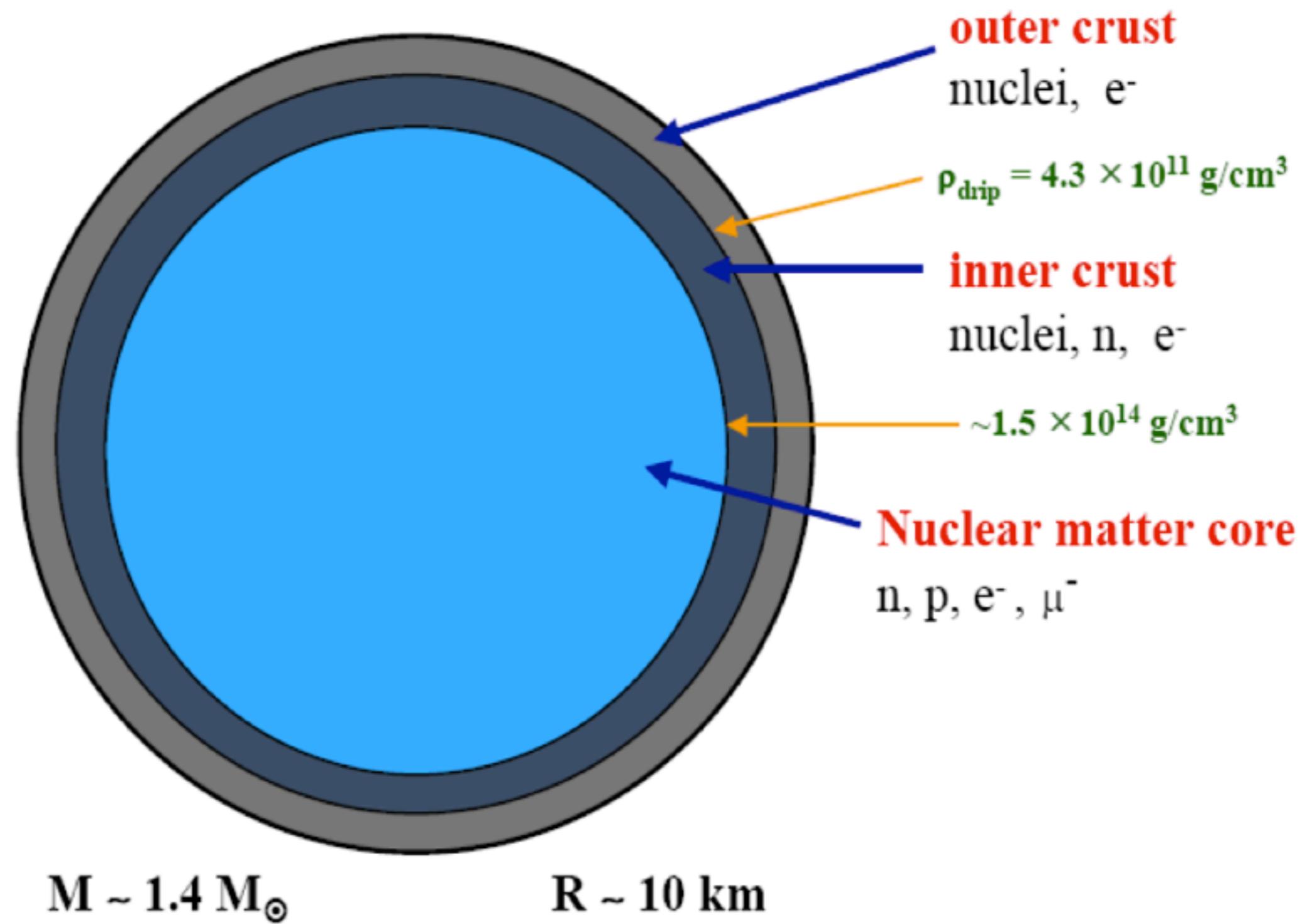
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NEUTRON STARS

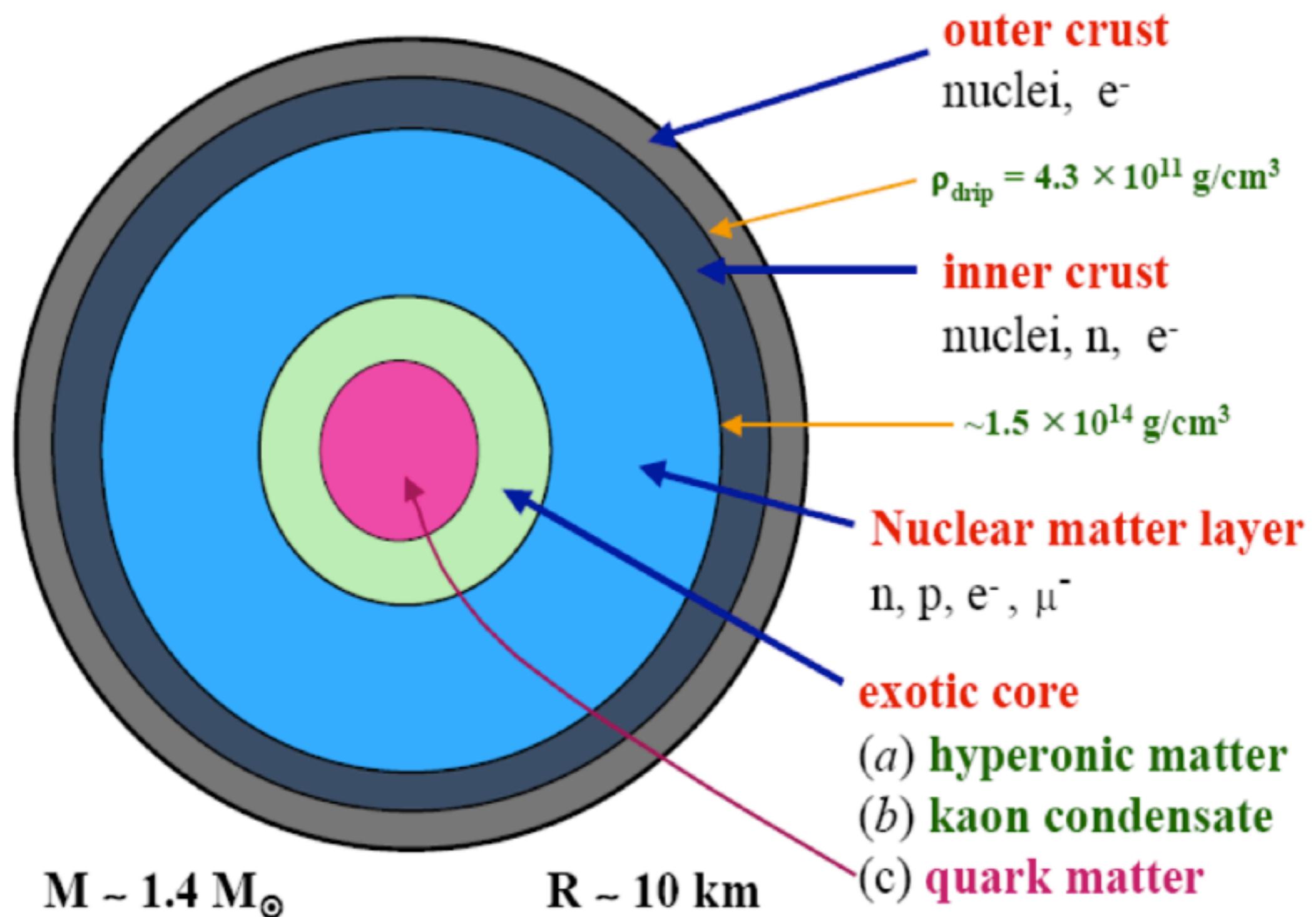


- *Produced in supernova explosions (Type II)*
- *Compact massive objects, $M \sim 1-2 M_{\text{solar}}$, $R \sim 10 \text{ km}$*

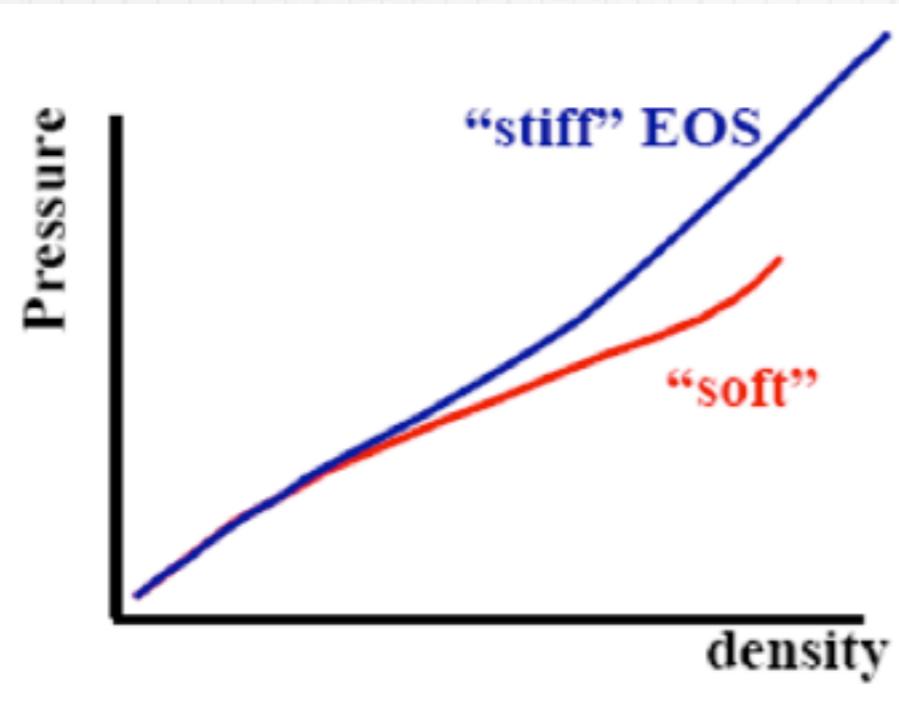
NEUTRON STAR STRUCTURE



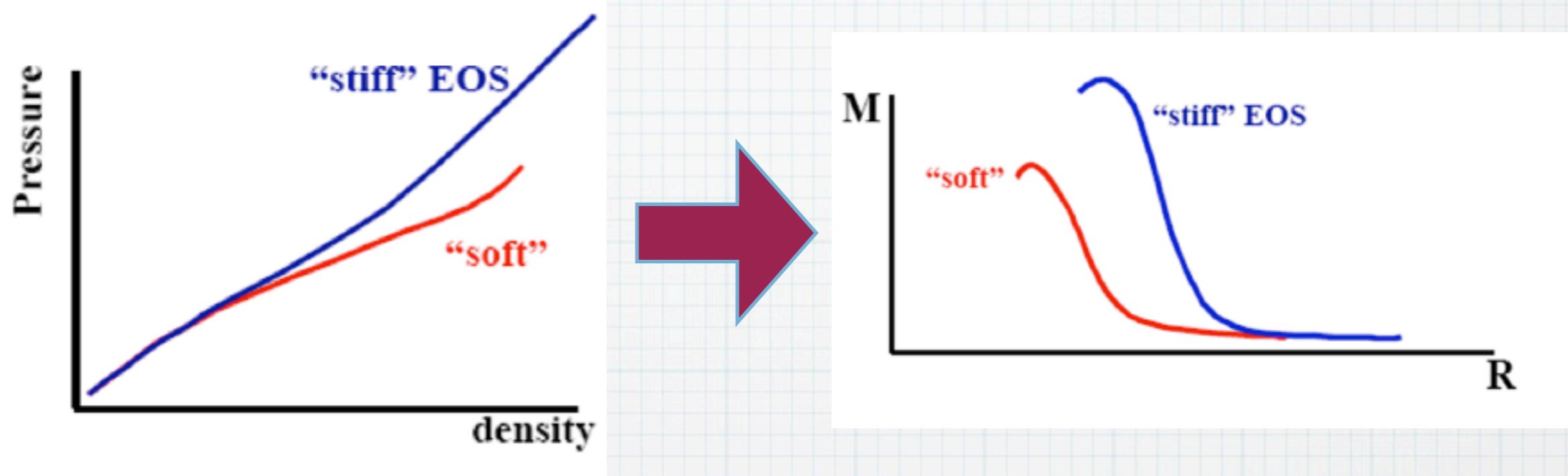
NEUTRON STAR STRUCTURE



Equation of state (EoS)



Equation of state (EoS)



Tolman-Oppenheimer-Volkov equations of relativistic hydrostatic equilibrium:

$$\frac{dp}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi pr^3)(\epsilon + p)}{r(r - 2Gm/c^2)}$$
$$\frac{dm}{dr} = 4\pi \frac{\epsilon}{c^2} r^2$$

Measurement of neutron star masses : Relativistic binaries

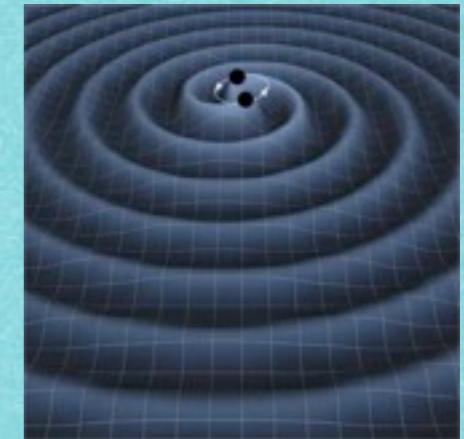
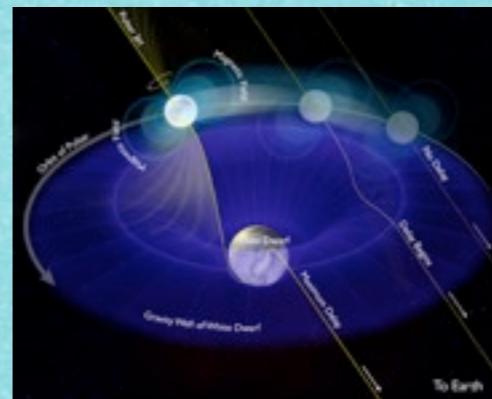
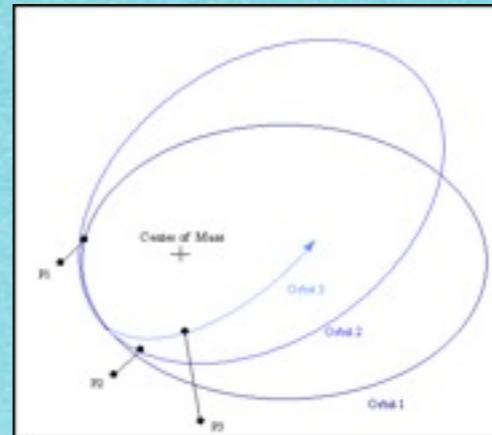
Keplerian parameters

- Orbital period P_b
- Projected semi-major axis $x = (a_p \sin i) / c$
- Orbital eccentricity e
- Longitude of periastron ω
- Epoch of periastron passage T_o

Measurement of neutron star masses : Relativistic binaries

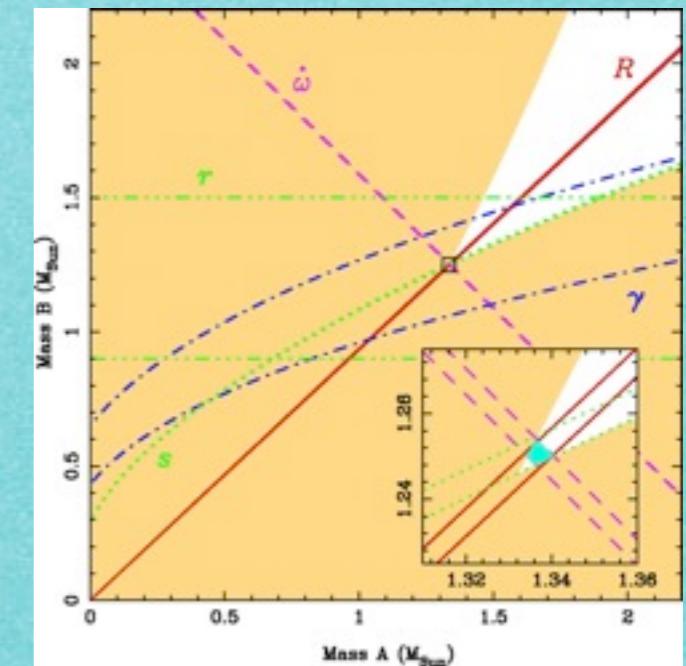
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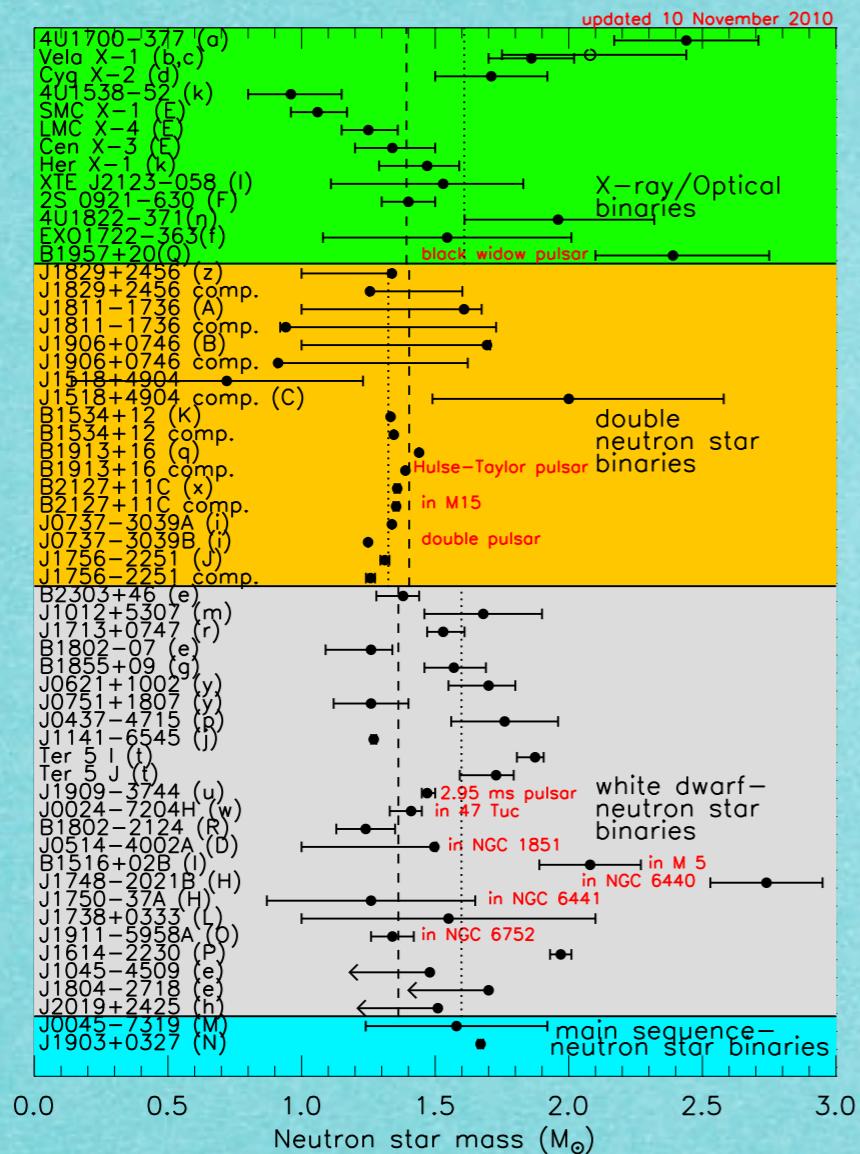


Post-Keplerian Parameters

- Relativistic advance of periastron $\dot{\omega}$
- Gravitational redshift and time dilation γ
- Orbital decay change in period \dot{P}_b
- Shapiro delay range r and shape s

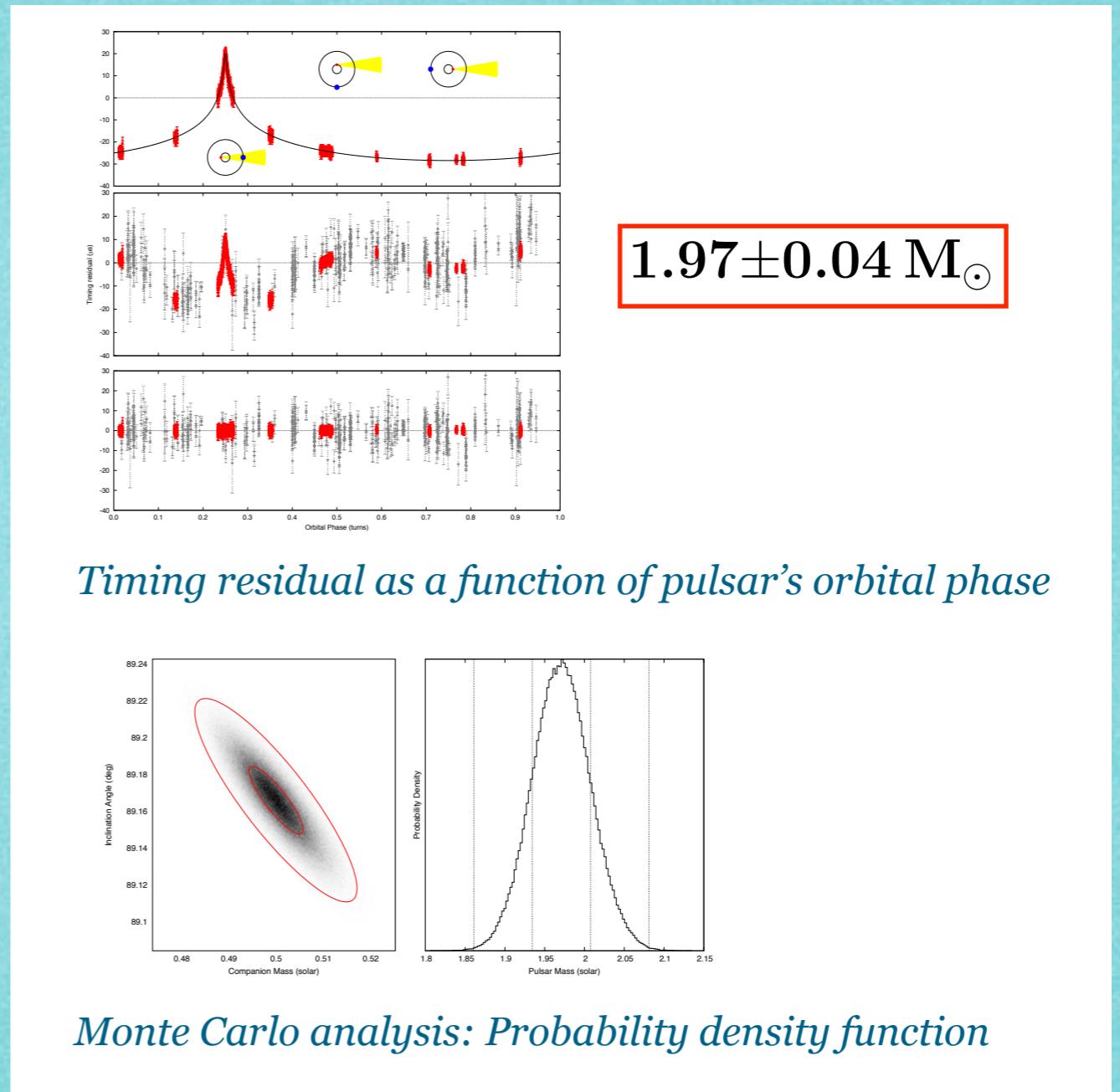
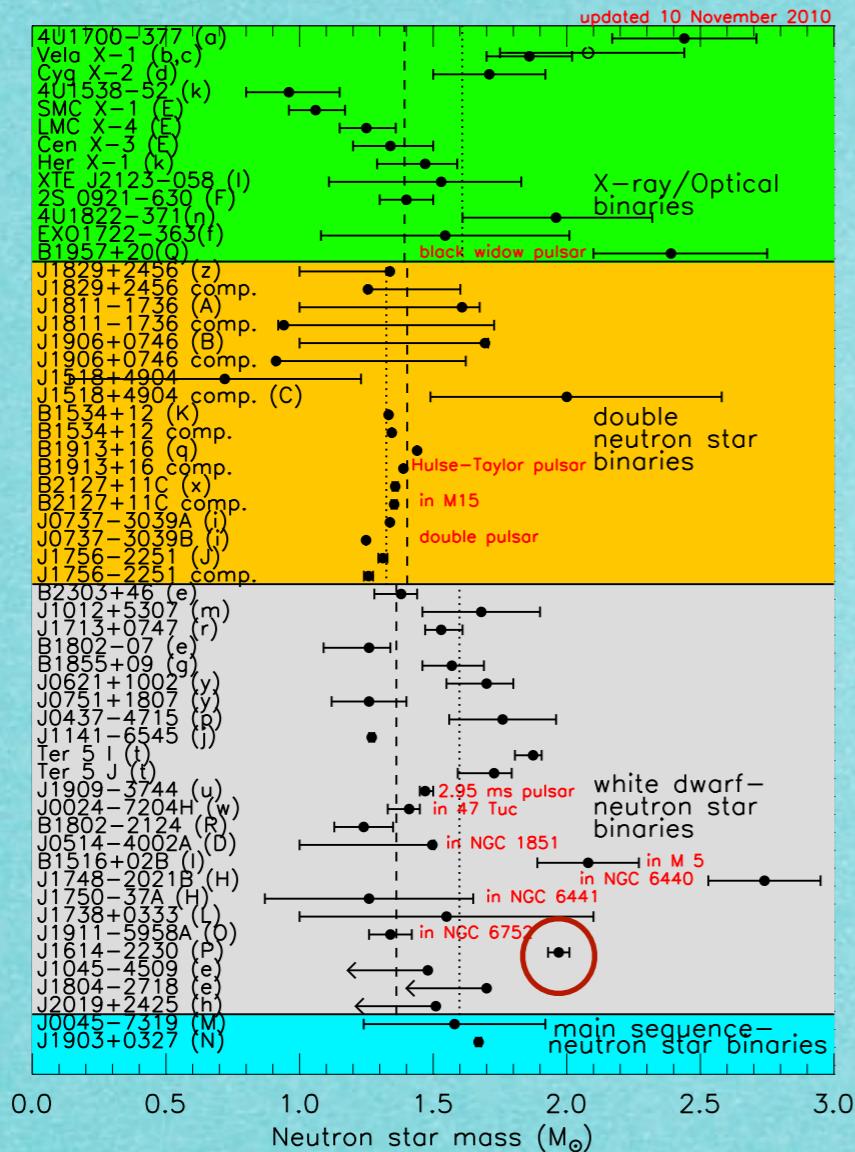


Mass measurements



Lattimer and Prakash, arXiv1012.3208

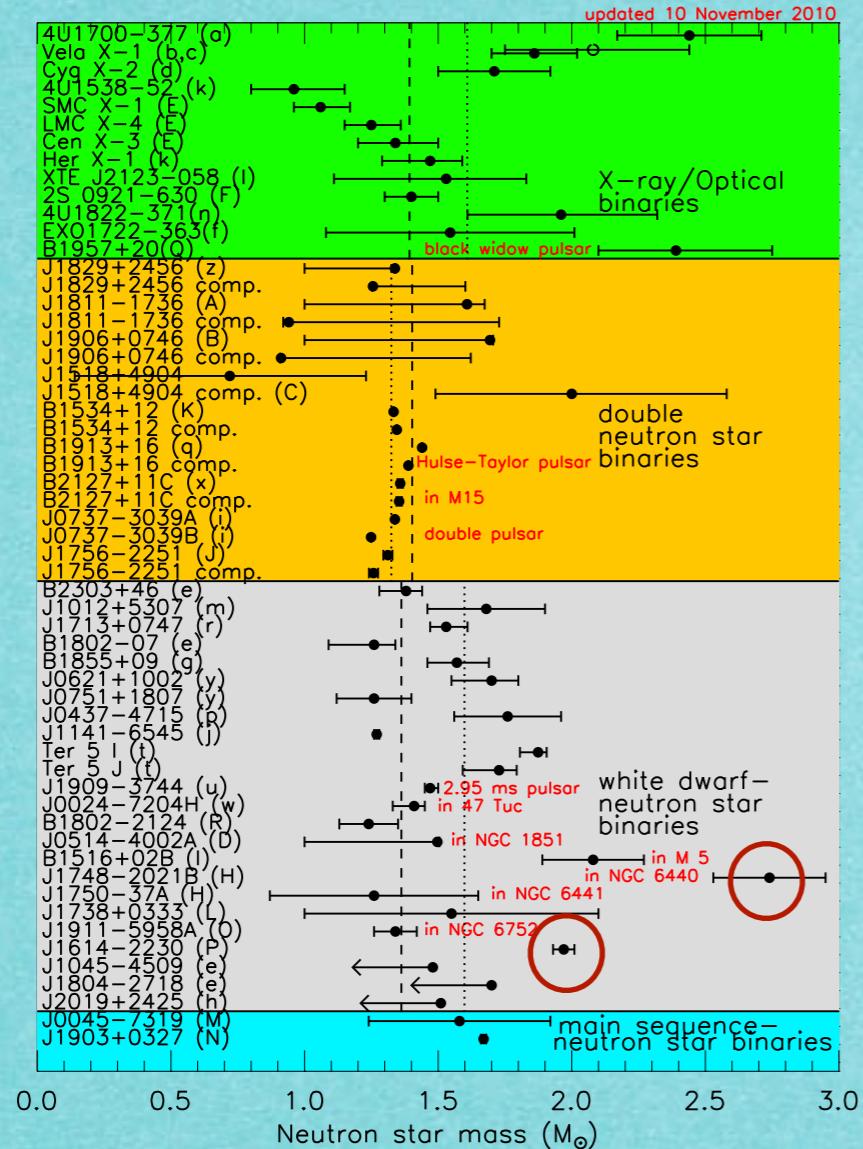
Highest mass measurement : J1614-2230



Lattimer and Prakash, arXiv1012.3208

Demorest et al (Nature 2010)

Mass measurements



Freire et al (Ap. J. 2008)

$$2.74 \pm 0.2 M_{\odot}$$

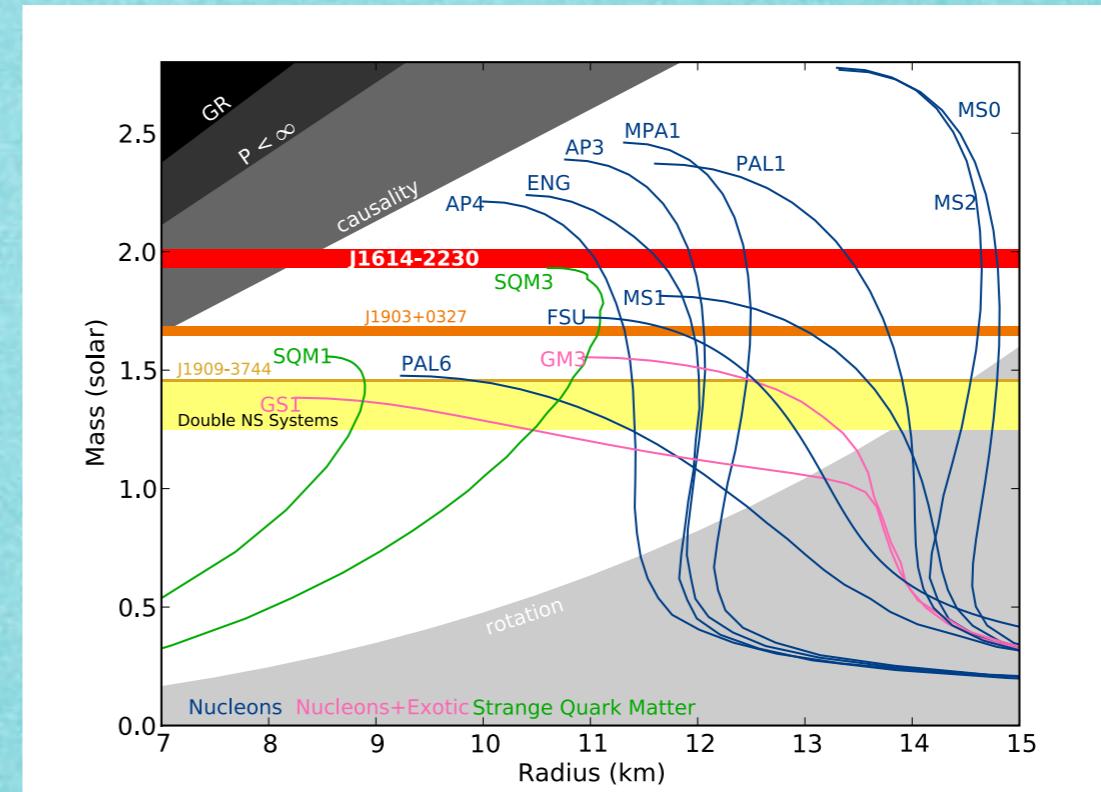
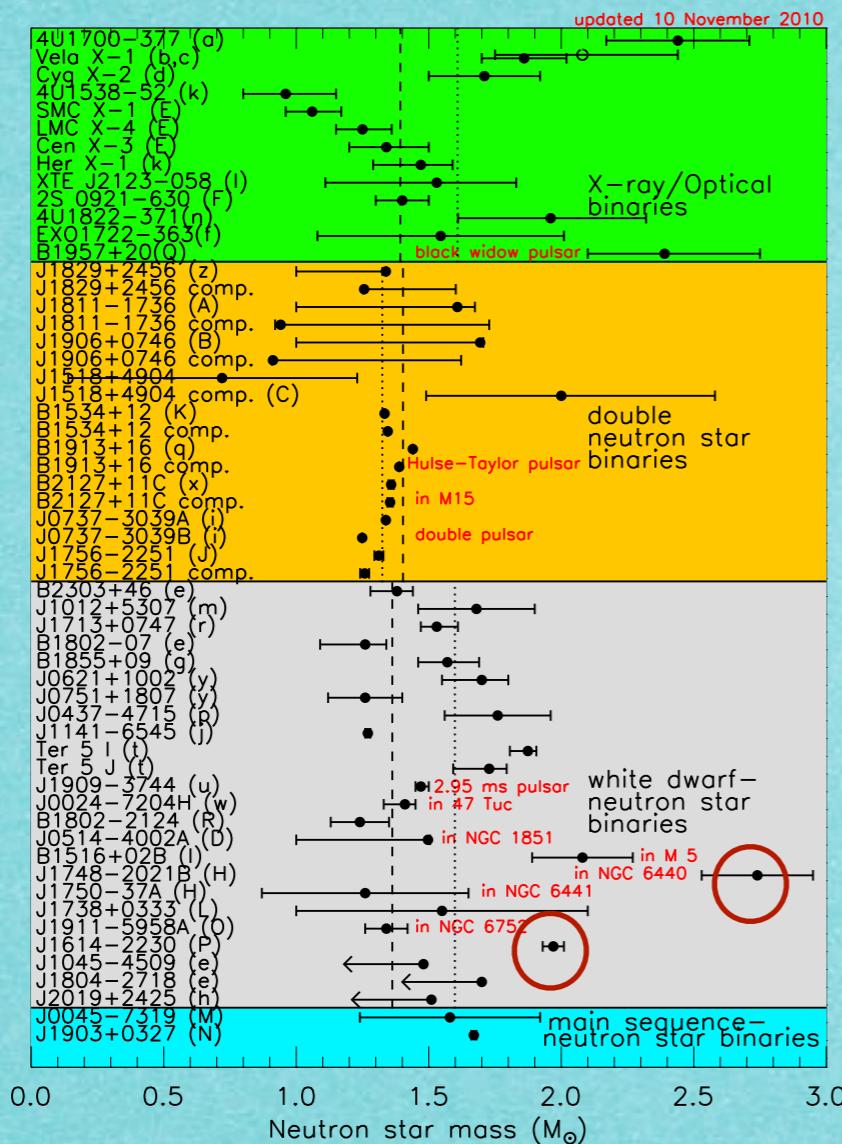
$$1.97 \pm 0.04 M_{\odot}$$

Demorest et al (Nature 2010)

Lattimer and Prakash, arXiv1012.3208

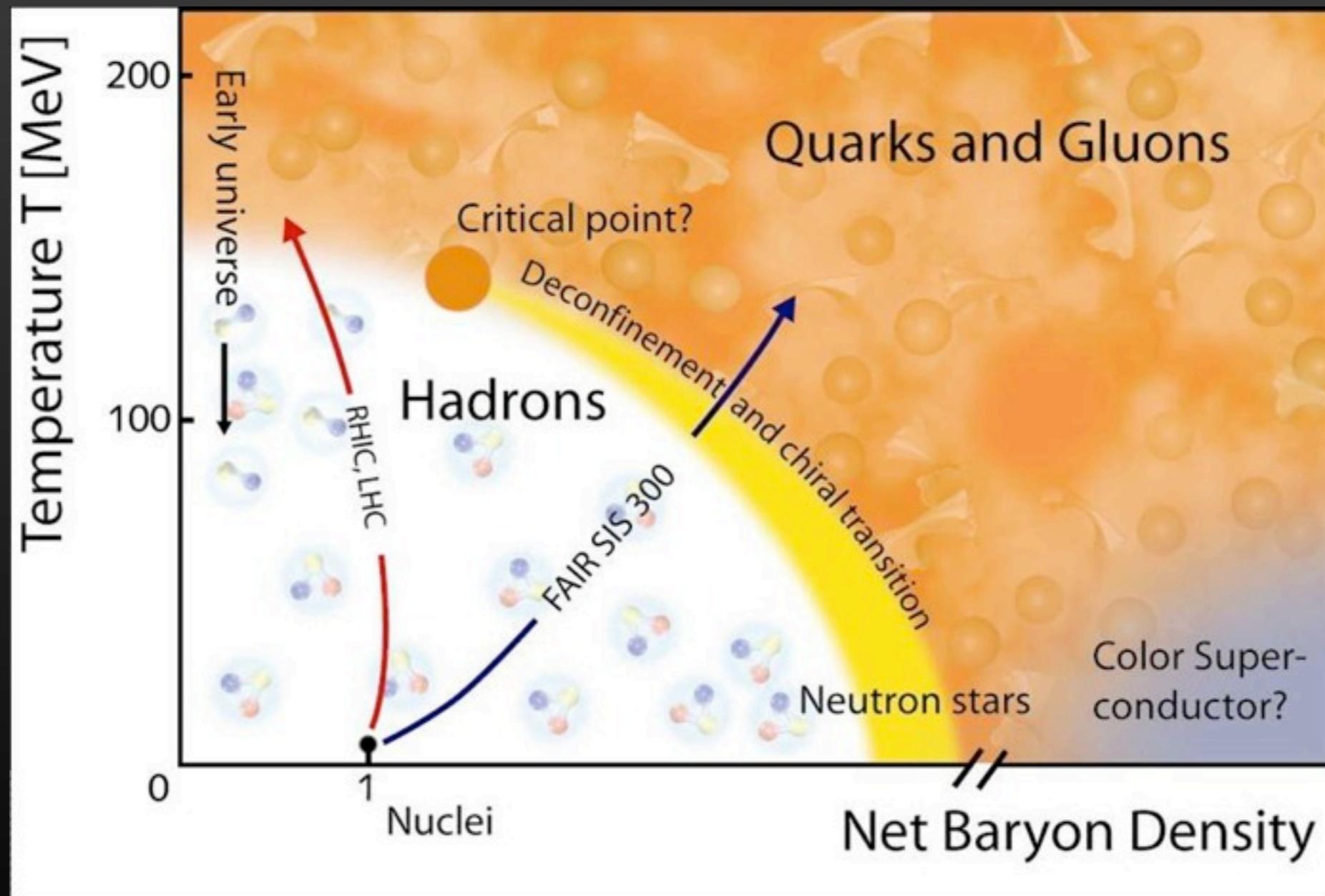
Constraining the EoS

$$M^{\max}(\text{theo}) > M^{\max}(\text{obs})$$

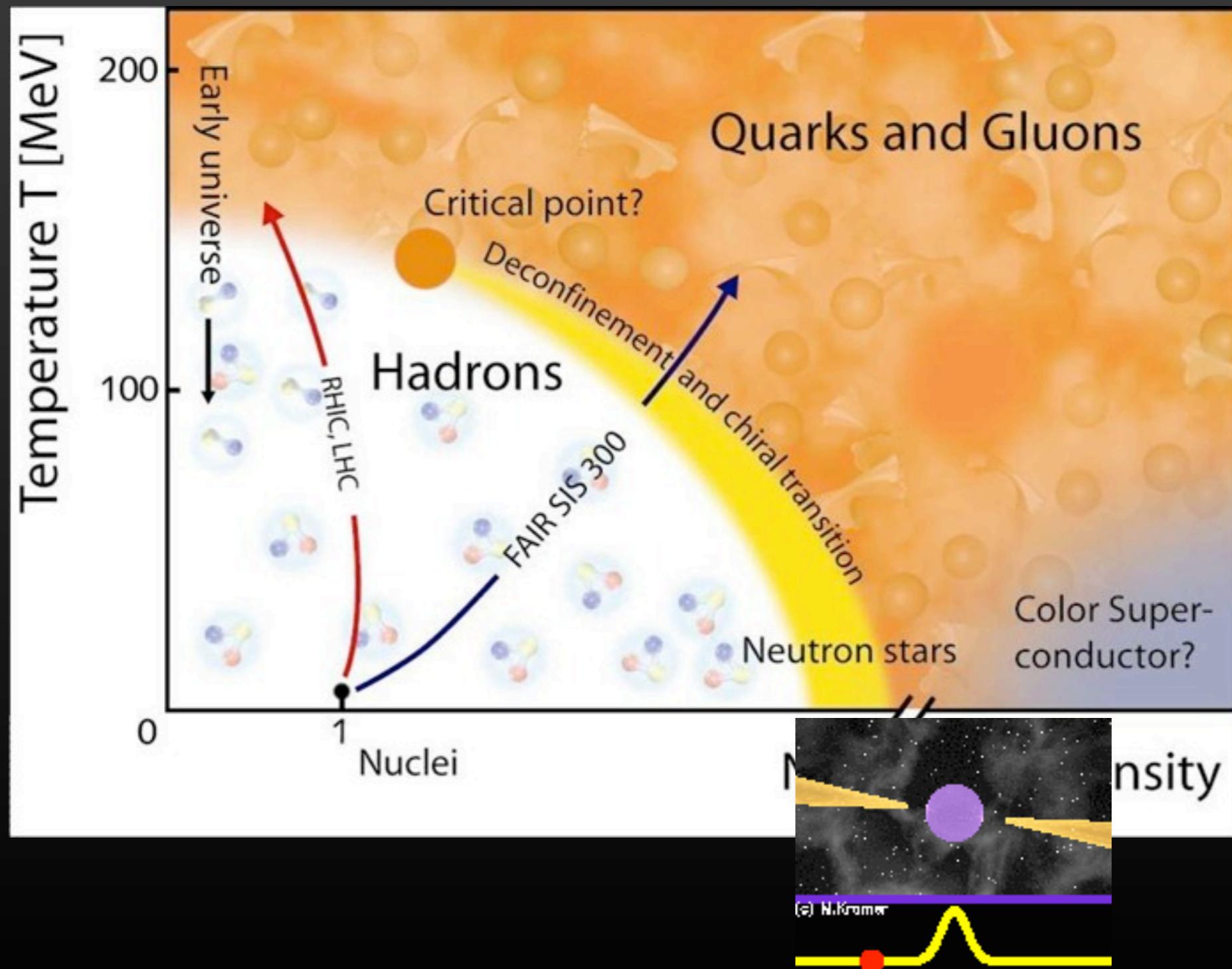


Lattimer and Prakash, arXiv1012.3208

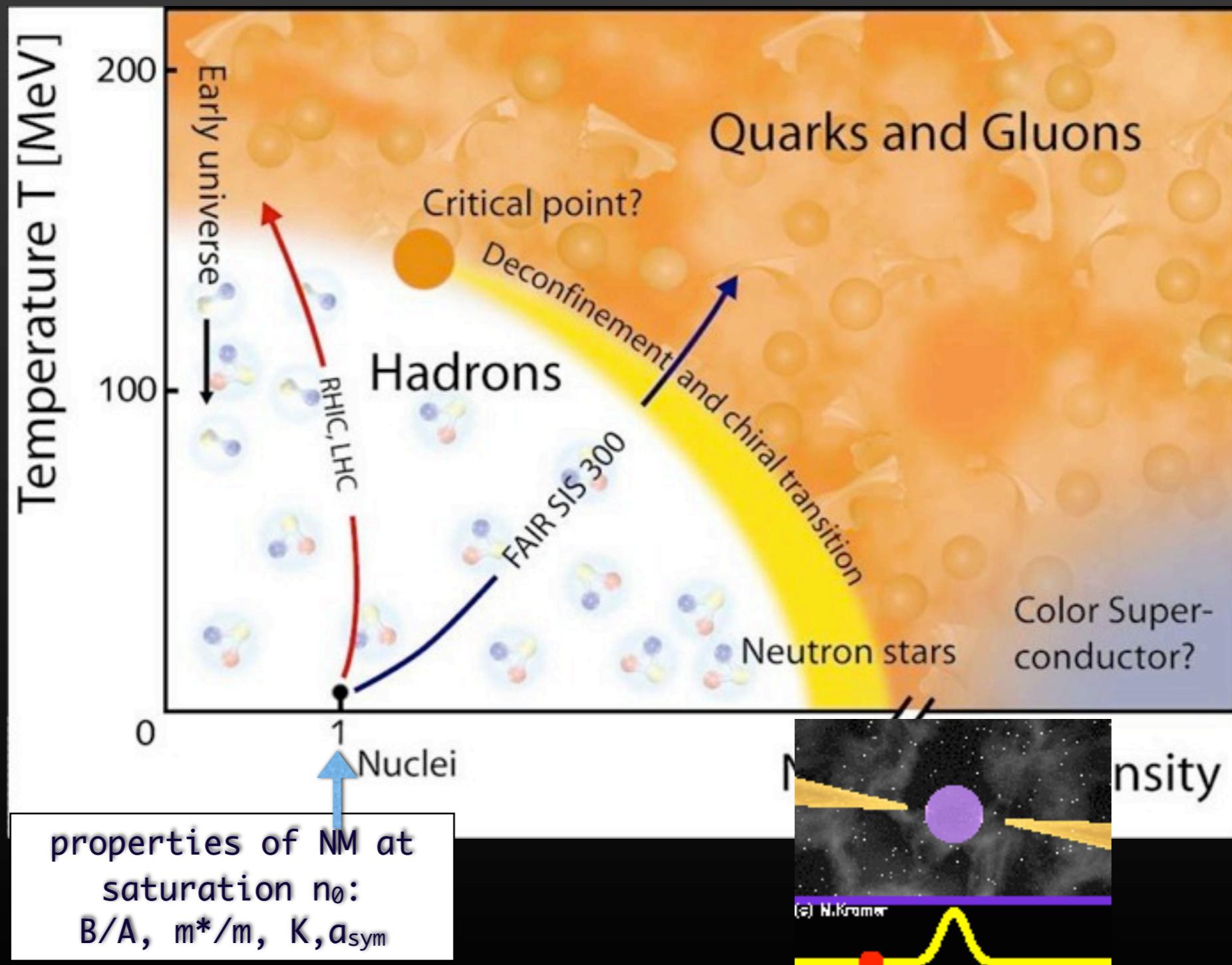
CONSTRAINTS FOR NUCLEAR EOS



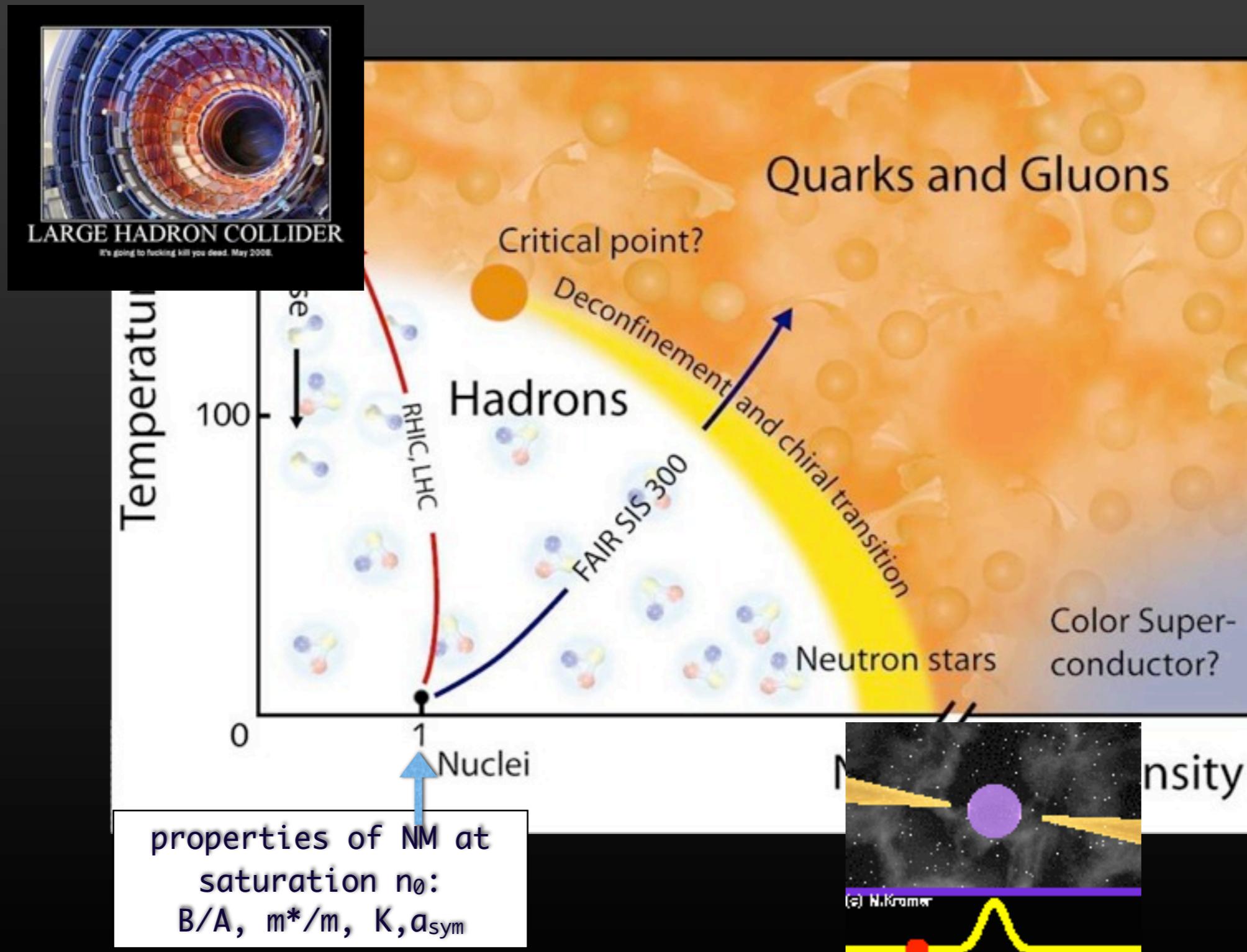
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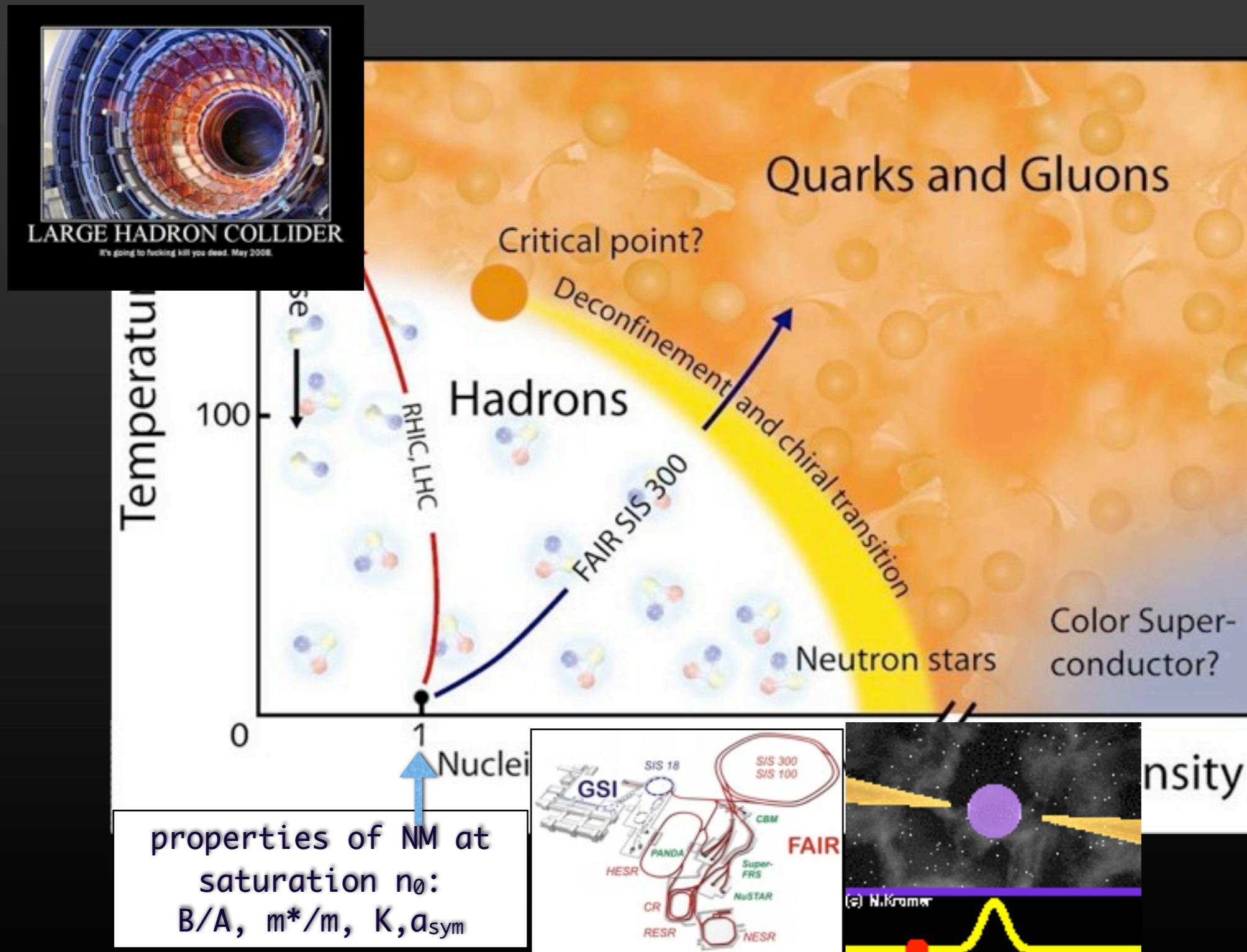
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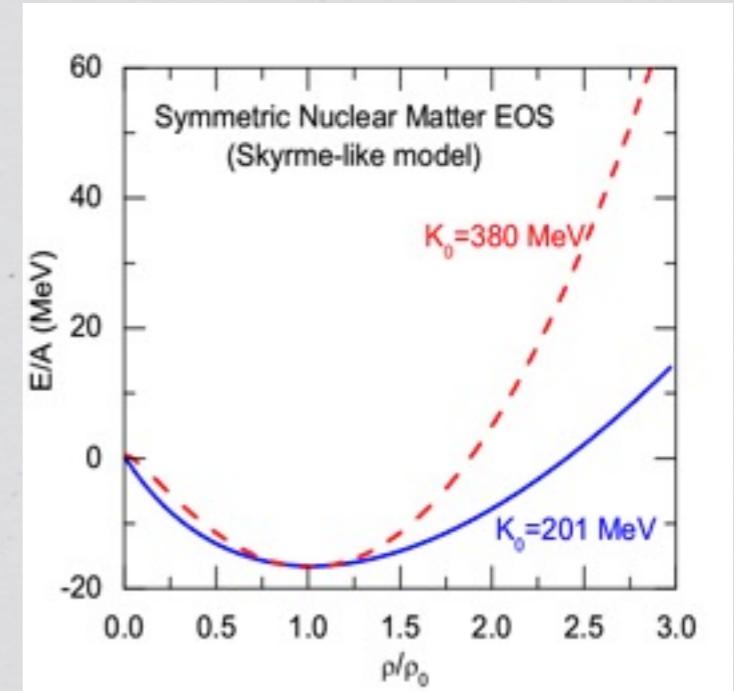
CONSTRAINTS FOR NUCLEAR EOS



Properties of dense nuclear matter

Symmetric nuclear matter at saturation

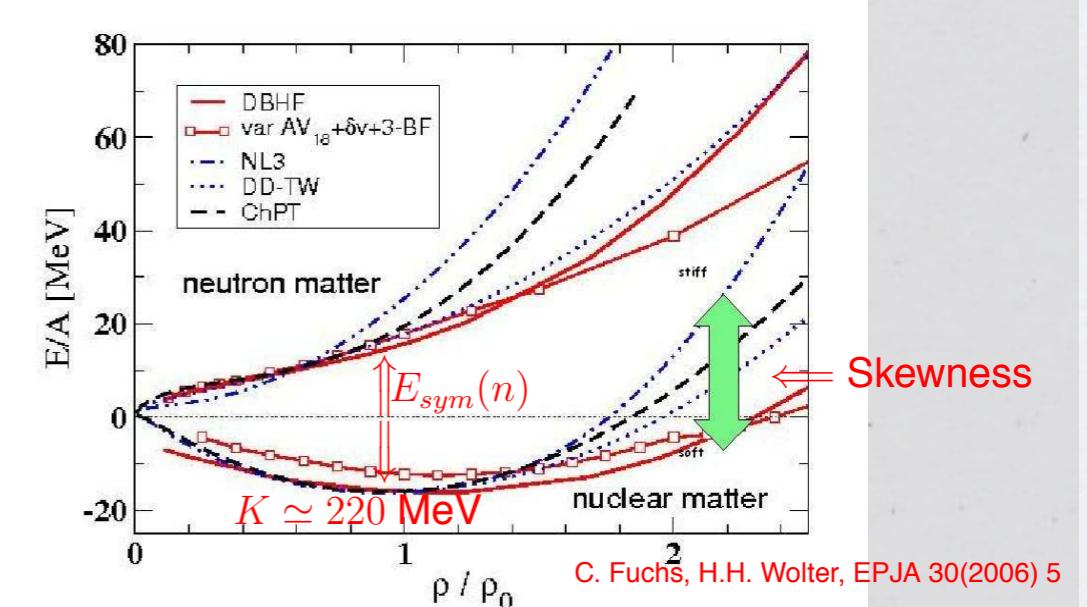
- * *saturation density* $n_0 = 0.17 \text{ fm}^{-3}$
- * *binding energy per nucleon* $B/A = -16.3 \text{ MeV}$
- * *effective nucleon mass* $m^*/m = 0.55-0.8$
- * *incompressibility* $K_o = 235 \pm 14 \text{ MeV}$



Properties of dense nuclear matter

Asymmetric nuclear matter at saturation

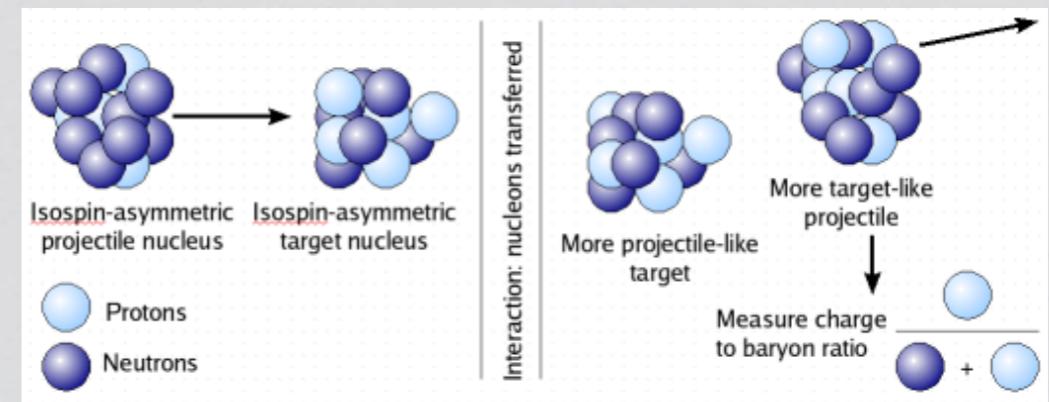
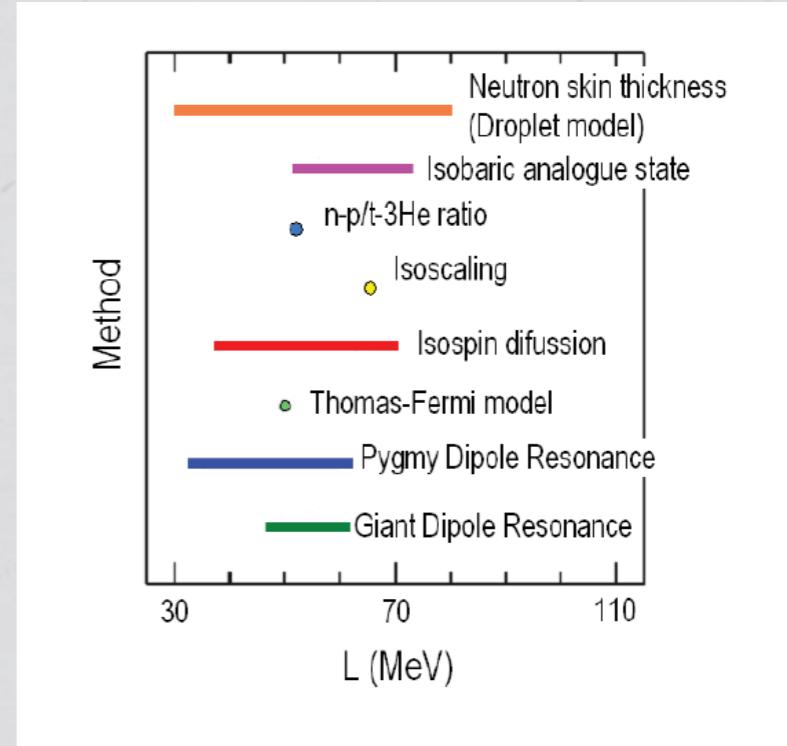
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- * *incompressibility* $K_o = 235 \pm 14 \text{ MeV}$
- * *symmetry energy* $E_{\text{sym}} = 28-32 \text{ MeV}$



Density dependence of Symmetry Energy L

Asymmetric nuclear matter at saturation

- * the density dependence of symmetry energy is a crucial quantity in nuclear physics
- * nuclei $\Rightarrow n < n_0$
- * Isospin diffusion data from intermediate energy HIC provide constraint on L only around n_0
- * neutron skin thickness of heavy nuclei
- * Giant dipole resonance in ^{208}Pb
- * Pygmy dipole resonance in ^{208}Pb

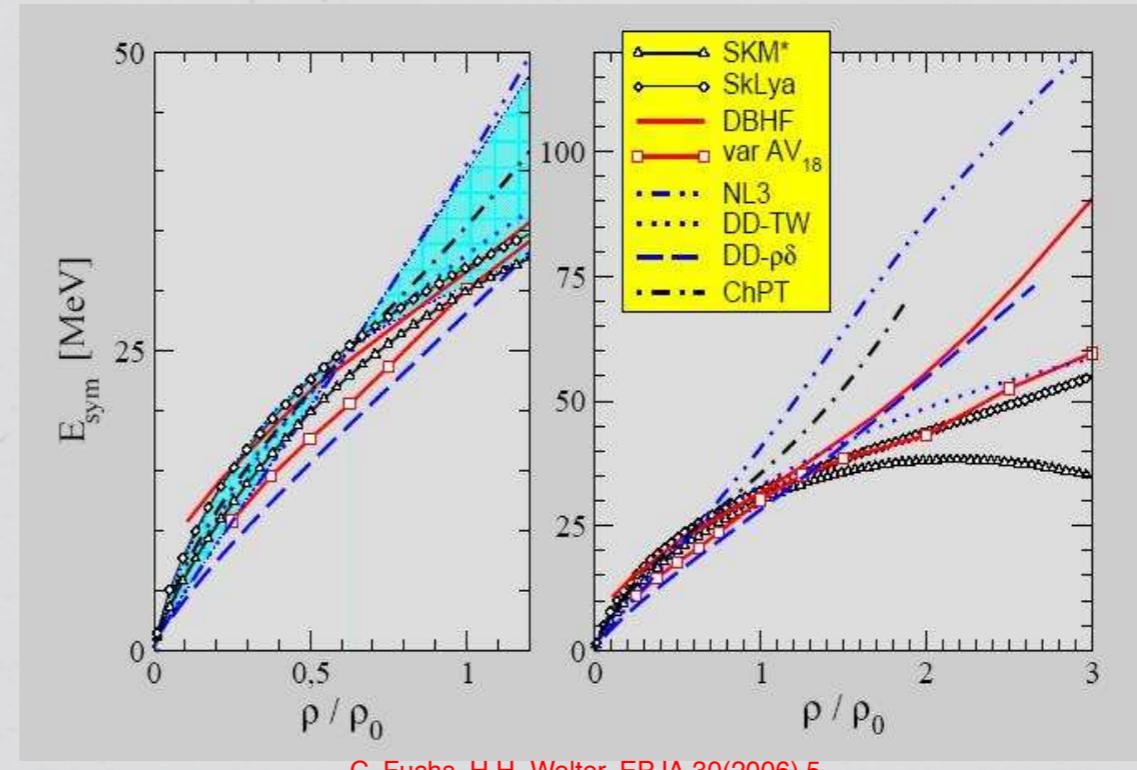


Density dependence of Symmetry Energy L

Nuclear matter beyond saturation

- * Density dependence of symmetry energy “ L ” becomes highly uncertain at $n \gg n_0$

The Uncertain $E_{sym}(n)$

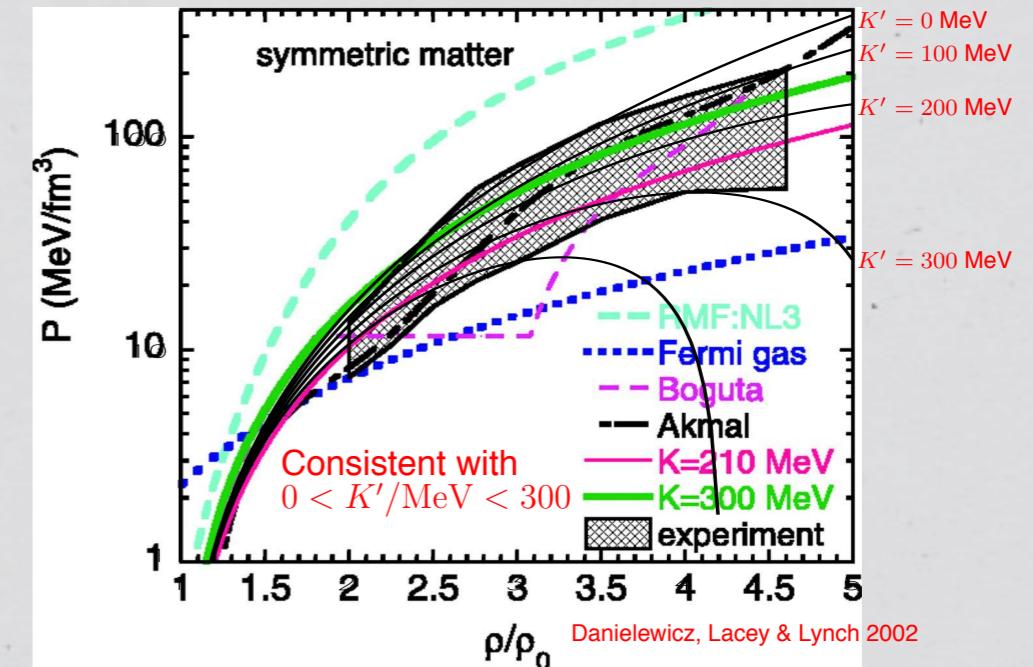


Density dependence of Symmetry Energy L

Nuclear matter beyond saturation

- * Density dependence of symmetry energy “ L ” becomes highly uncertain at $n \gg n_0$
- * Elliptic flow of nucleons in non-central nucleus-nucleus collisions
 - (not conclusive, new degrees of freedom at high energies, momentum dependence of interaction, model dependence of analysis)

Flow Constraint From Heavy Ions



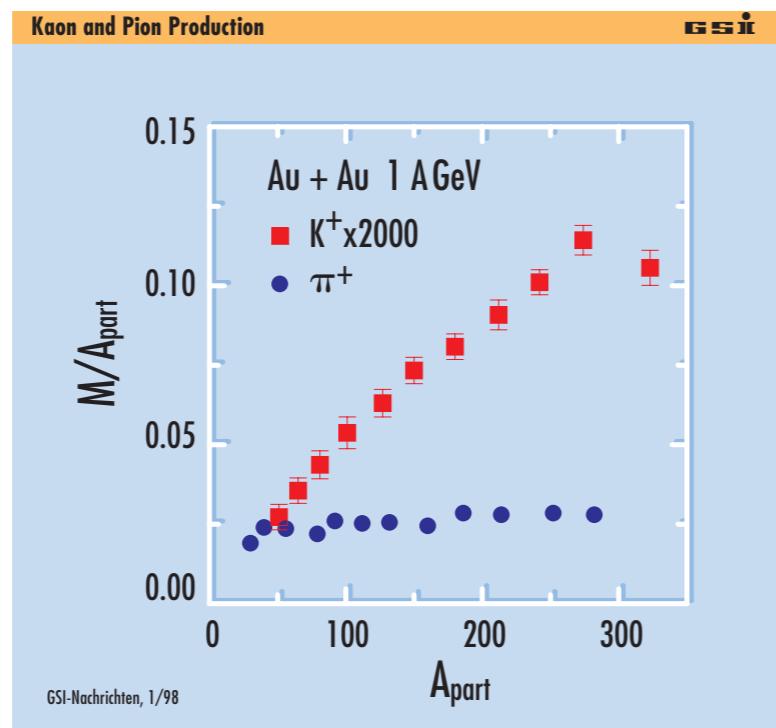
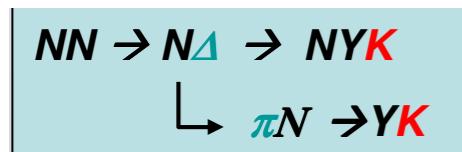
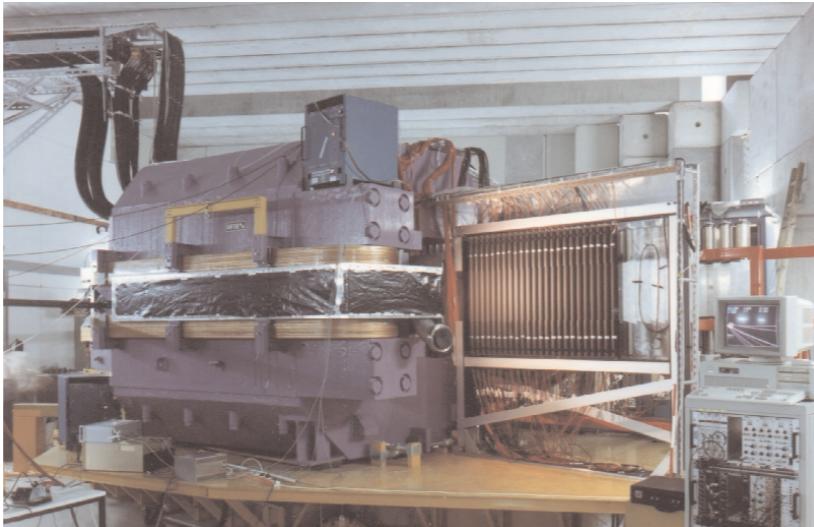
Density dependence of Symmetry Energy L

Nuclear matter beyond saturation

- * *Density dependence of symmetry energy “ L ” becomes highly uncertain at $n >> n_0$*
- * *Elliptic flow of nucleons in non-central nucleus-nucleus collisions*
- (not conclusive, new degrees of freedom at high energies, momentum dependence of interaction, model dependence of analysis)
- * *K^+ meson production in nuclear collisions* ✓

K^+ meson production in heavy-ion collisions

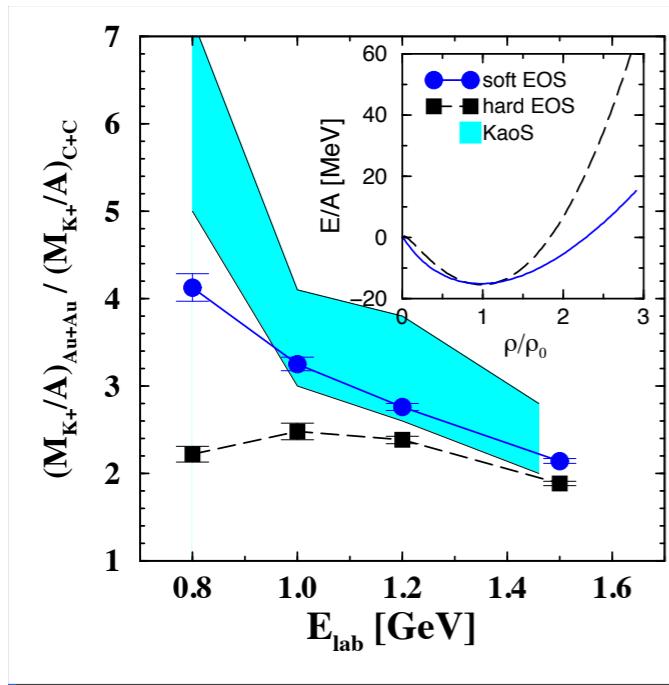
KaoS experiment,
GSI Darmstadt



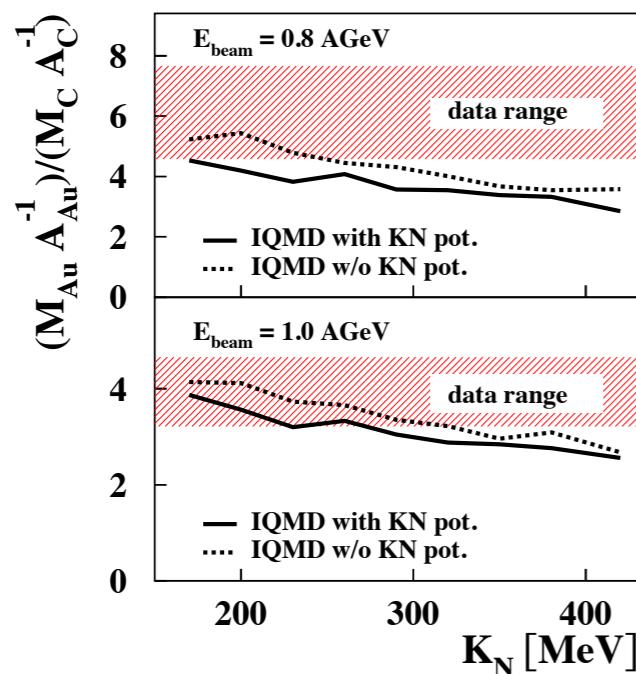
Subthreshold production of K^+ particles

- * K^+ particles produced by multiple NN collisions ($NN \rightarrow N\Lambda K, NN \rightarrow NN\bar{K}\bar{K}$) or secondary collisions ($\pi N \rightarrow \Lambda K, \pi\Lambda \rightarrow N\bar{K}$)
- * Nuclear matter compressed up to $\sim 2\text{-}3 n_0$
- * Production of K^+ particles sensitive to the nuclear EoS
 \Rightarrow tool to probe compressibility of nuclear matter at $\sim 2\text{-}3 n_0$

Soft equation of state from heavy-ion data



Sturm et al. (KaoS collaboration), PRL 2001



Hartnack, Oeschler, Aichelin, PRL 2006

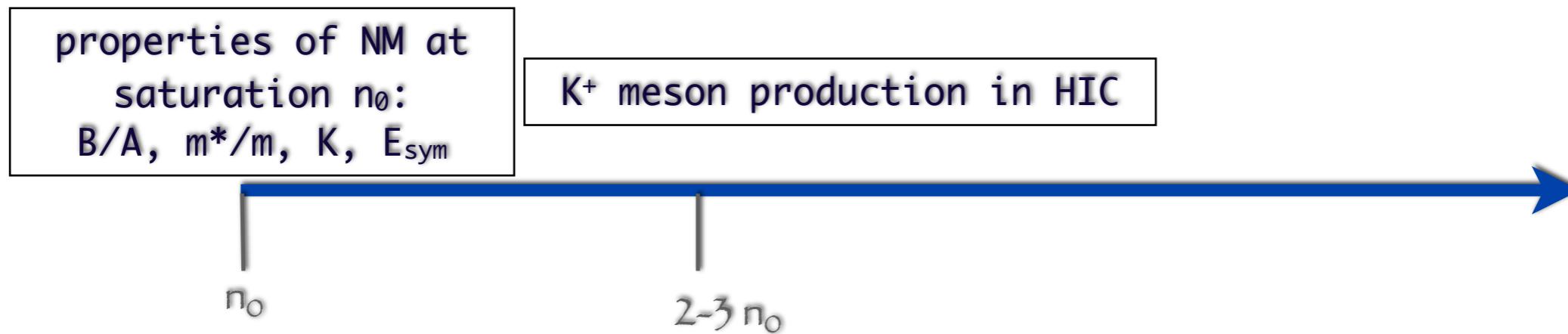
- * K^+ multiplicity ratio in $Au+Au$ and $C+C$ collisions at 0.8 AGeV and 1.0 AGeV is sensitive to the compression modulus of matter
 - * transport model calculations performed: Skyrme-type nucleon potential with $2BF$, $3BF$ were applied, with parameters to reproduce a soft EoS (with $K = 200$ MeV) and a stiff one (with $K = 380$ MeV).
 - * transport models agree, confirm that matter in the collision zone reaches densities up to $2-3 n_0$
 - * only $K \sim 200$ MeV can describe the data (KaoS collaboration, 2007)
- ⇒ *the nuclear EoS is soft*

Testing soft EoS with Neutron Stars

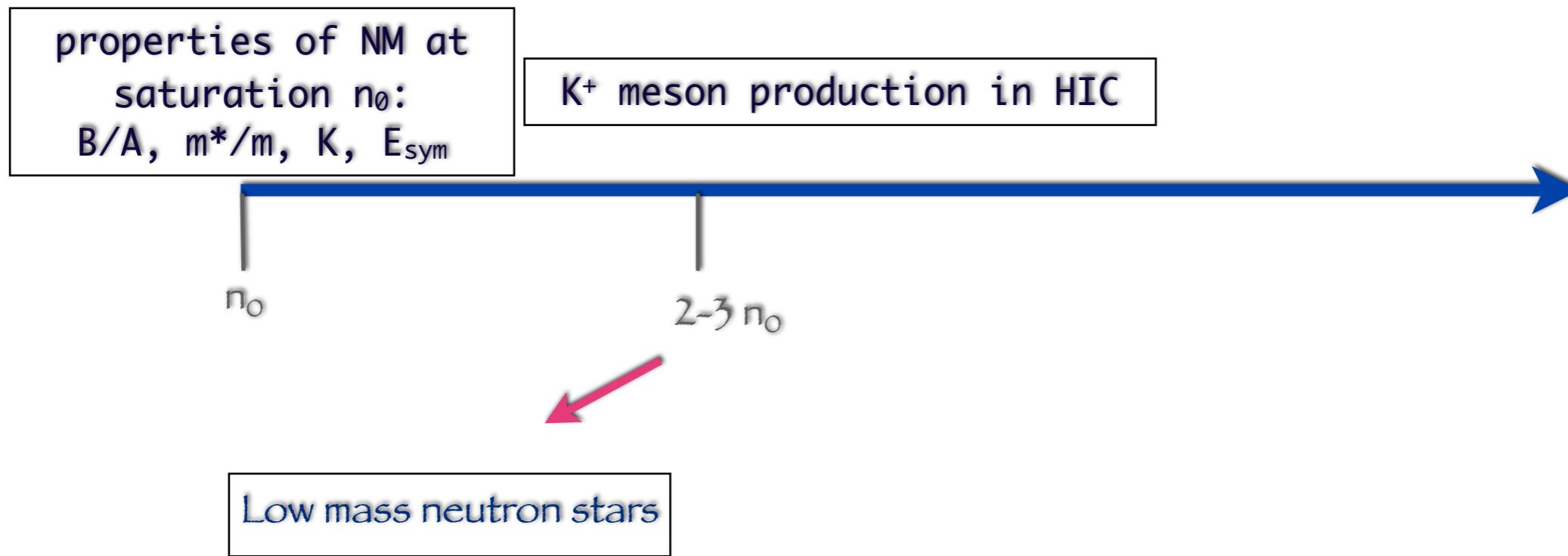
properties of NM at
saturation n_0 :
 B/A , m^*/m , K , E_{sym}



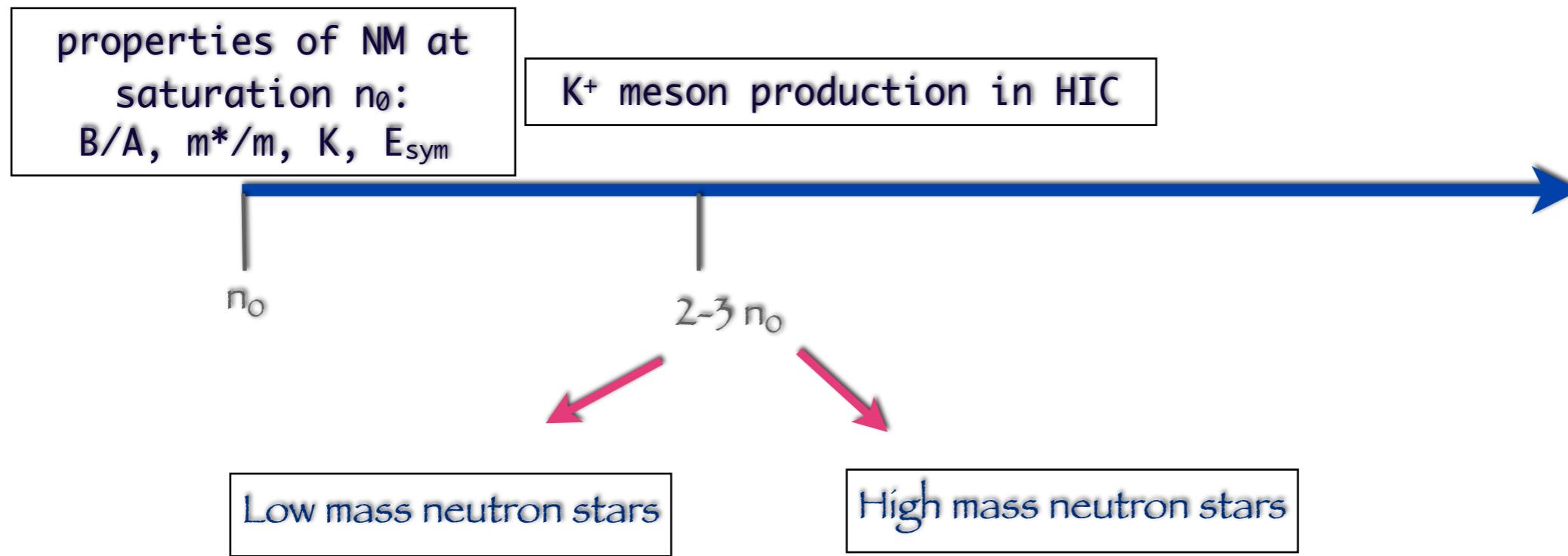
Testing soft EoS with Neutron Stars



Testing soft EoS with Neutron Stars



Testing soft EoS with Neutron Stars



Phenomenological EoS for NS core

$$\begin{aligned} \frac{E}{A} = & m_n (1 - Y_p) + m_p Y_p + E_0 u^{\frac{2}{3}} + B \frac{u}{2} + D \frac{u^\sigma}{(\sigma + 1)} \\ & + (1 - 2Y_p)^2 \left[\left(2^{\frac{2}{3}} - 1 \right) E_0 \left(u^{\frac{2}{3}} - F(u) \right) + S_0 u^\gamma \right], \quad (1) \end{aligned}$$

Skyrme EoS

- * E_0 = binding energy of SNM at n_0
- * baryon number density $u = n/n_0$
- * Y_p = proton fraction
- * density dependence of symmetry energy chosen as a power law with u^γ
- * parameters σ, B, D (2BF, 3BF)

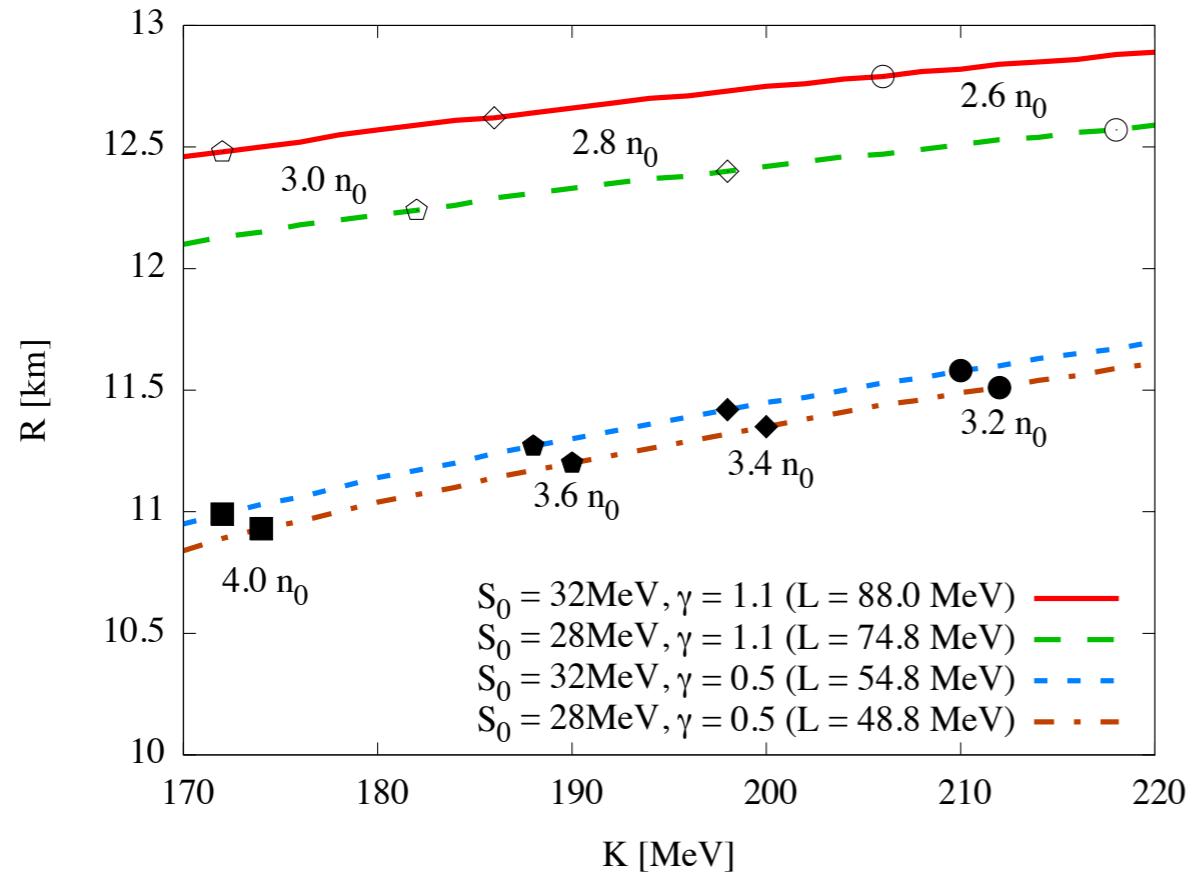
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- * Y_p = proton fraction
- * density dependence of symmetry energy chosen as a power law with u^γ
- * parameters σ, B, D (2BF, 3BF)
- * Parameters fitted to reproduce saturation density, binding energy, stiffness parameter
- * Variation of values:
 - $K = 170 - 220$ MeV
 - $S_0 = 28 - 32$ MeV
 - $\gamma = 0.5 - 1.1$ (motivated by heavy-ion experiments)
- * $M = 1.25 M_{sol}$: lightest pulsar mass deduced from observations

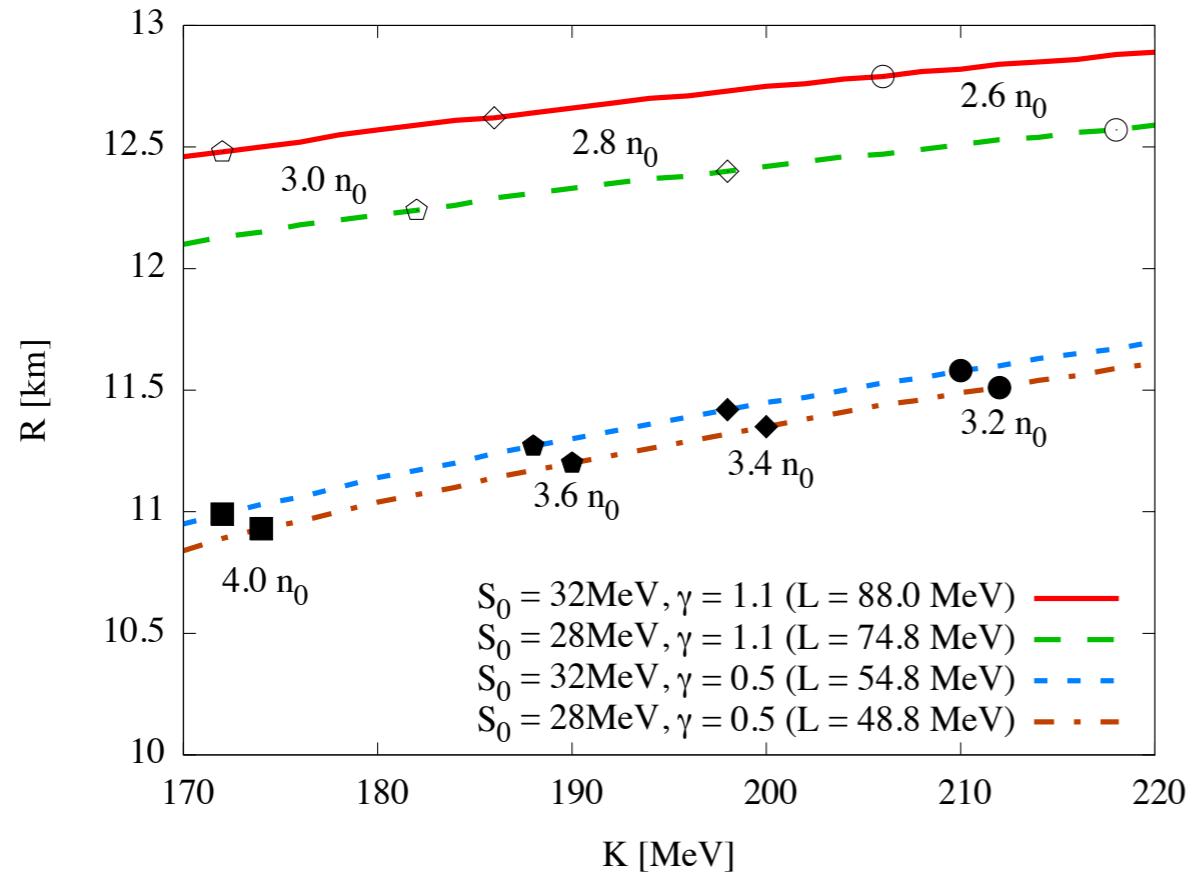
Radius of light neutron stars



*Radii and central densities of
1.25 M_{sol} neutron stars with K ,
for different values of S_0 and γ*

*I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich
and C. Sturm,, arXiv :1112.0234*

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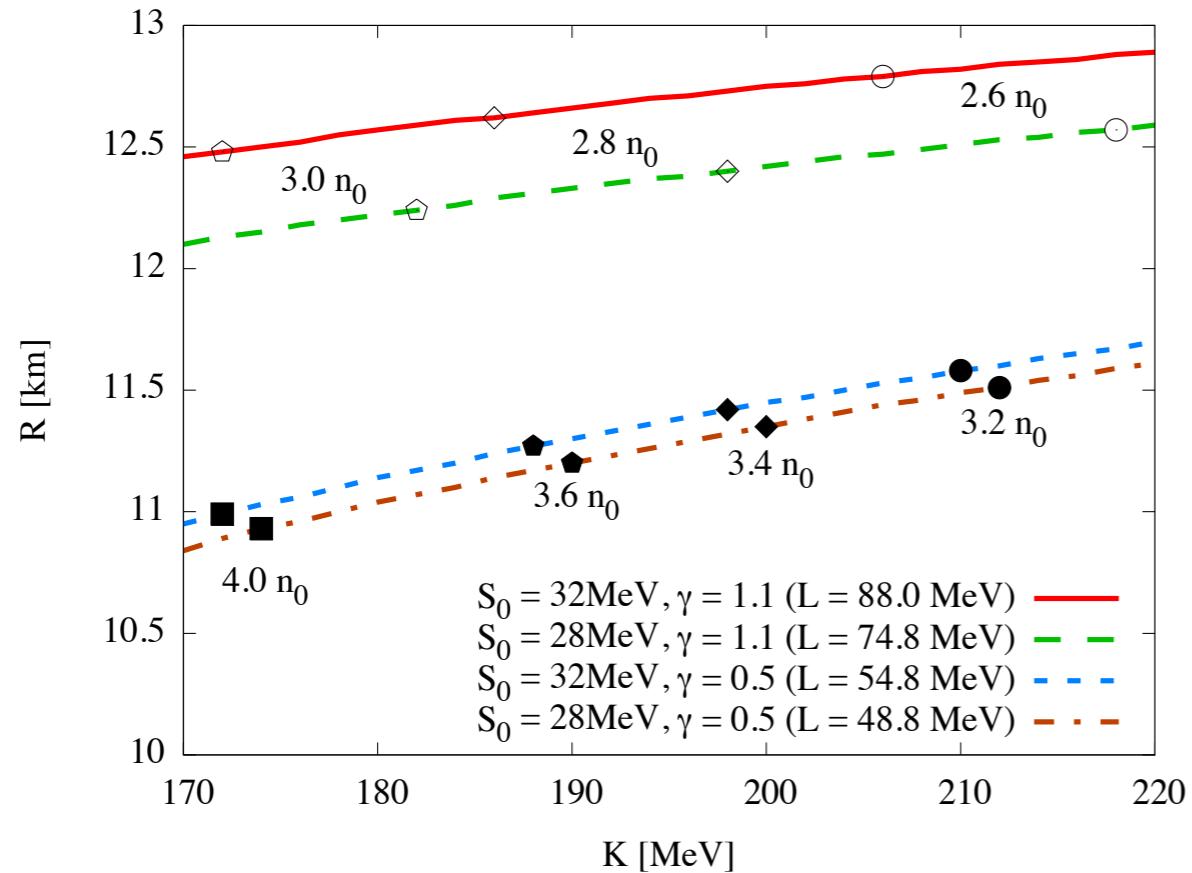


- * The central densities of the corresponding stars are in the range of the density region explored by KaoS.

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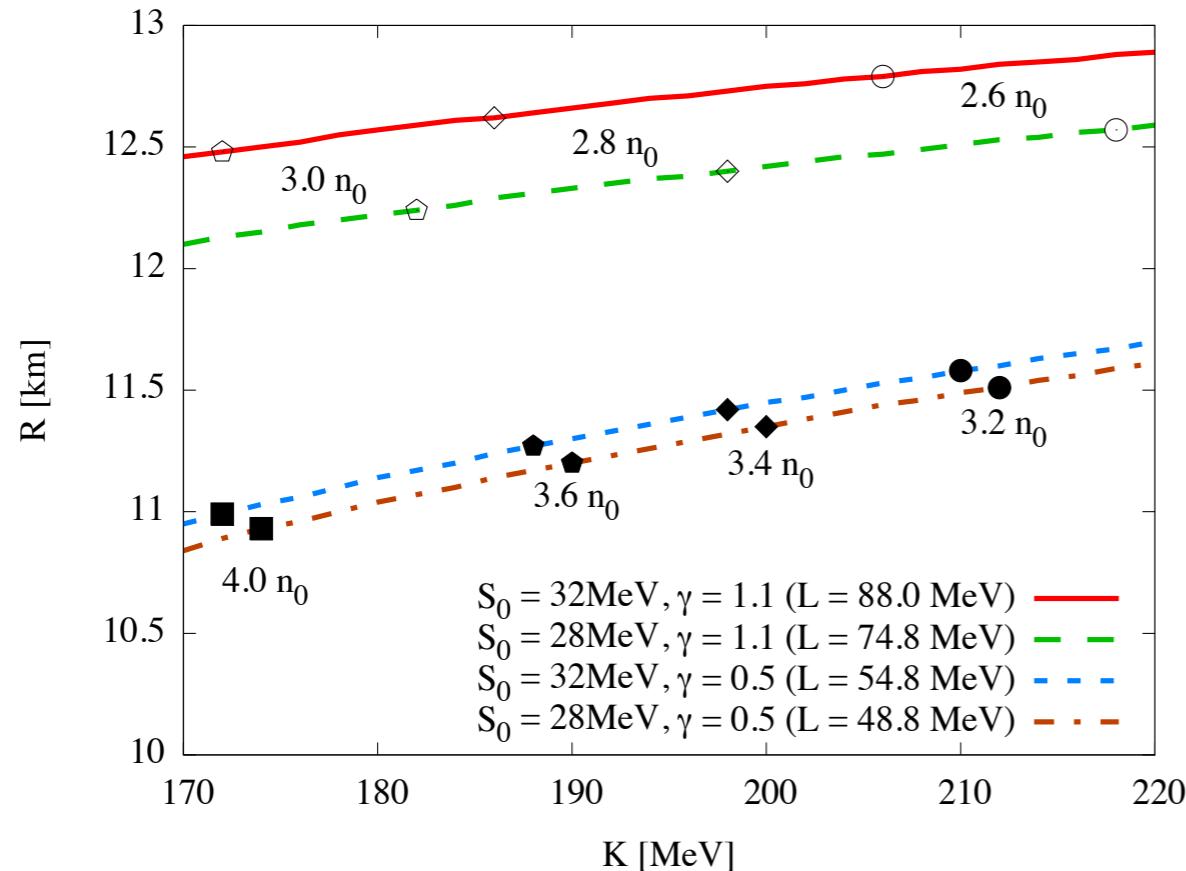


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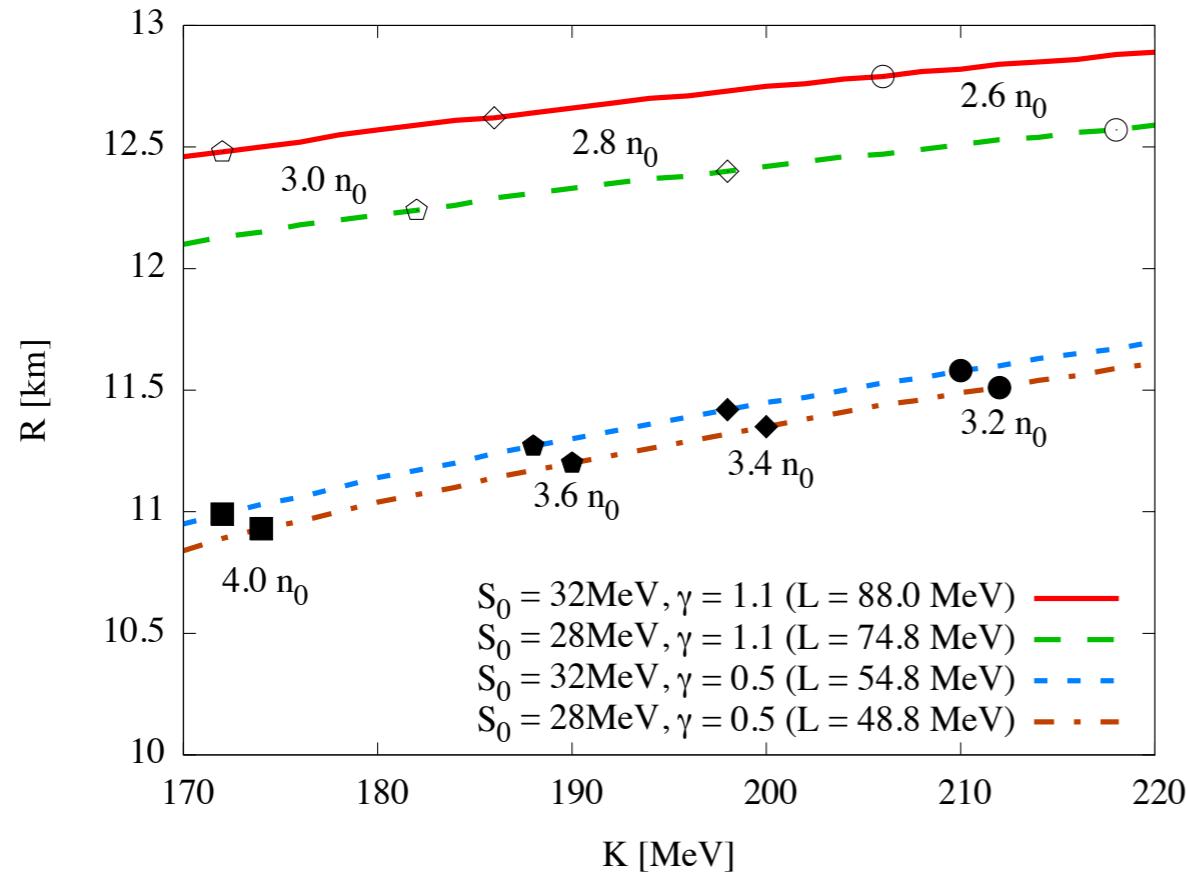


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- * At $K \sim 200 \text{ MeV}$, stiff and soft symmetry energy configurations lead to a difference in the neutron star radius of around $\Delta R \sim (1 - 1.5) \text{ km}$.

Radii and central densities of $1.25 M_{\text{sol}}$ neutron stars with K , for different values of S_0 and γ

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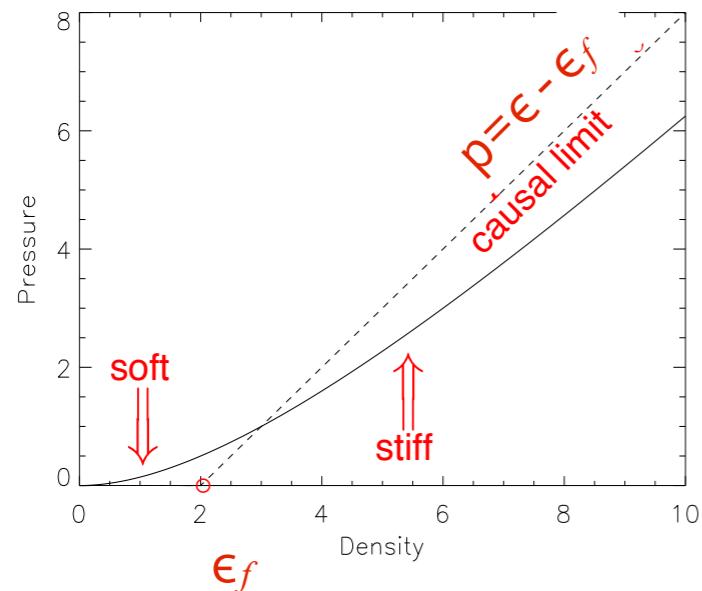


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- * radii of light neutron stars with $M \sim 1.25 M_{\text{sol}}$ are strong candidates for a direct cross check between heavy-ion experiments and astrophysical observations

I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich
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Massive neutron stars



$$M_{\max} = 4.2 M_{\odot} (\epsilon_0 / \epsilon_f)^{1/2}$$

Rhoades & Ruffini (1974)

Hartle (1978)

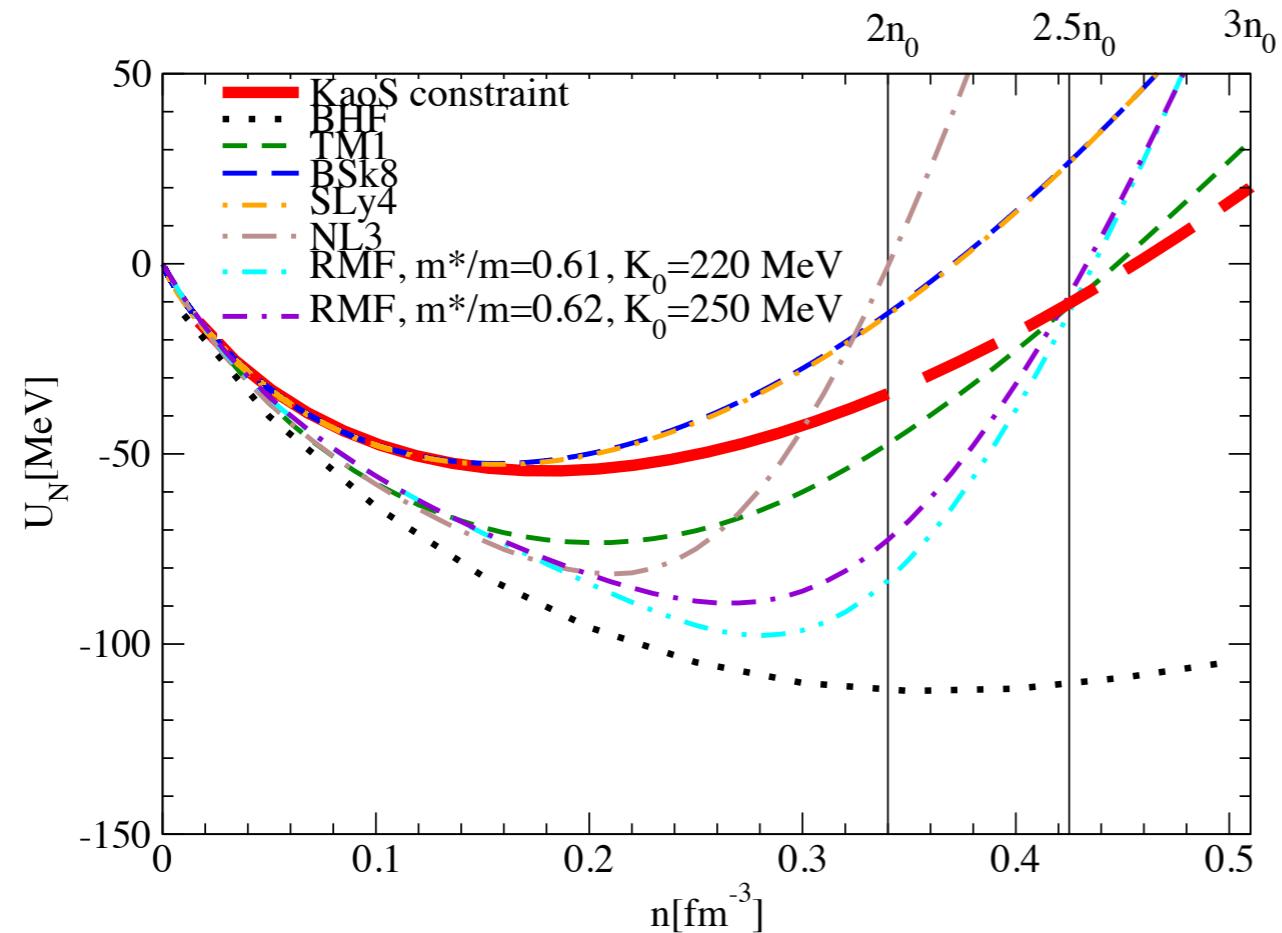
- * *Stiffest causal EoS:*
 $p = \epsilon - \epsilon_f$ above the fiducial density ϵ_f
- * *at high densities, smooth transition to the stiffest EoS*
- * *gives the highest possible mass of a compact star*
- * *At low densities, EoS should satisfy KaoS constraint*
⇒ *new upper mass limit of $3 M_{\odot}$ from heavy-ion data*

Equations of State

Densities around and above saturation

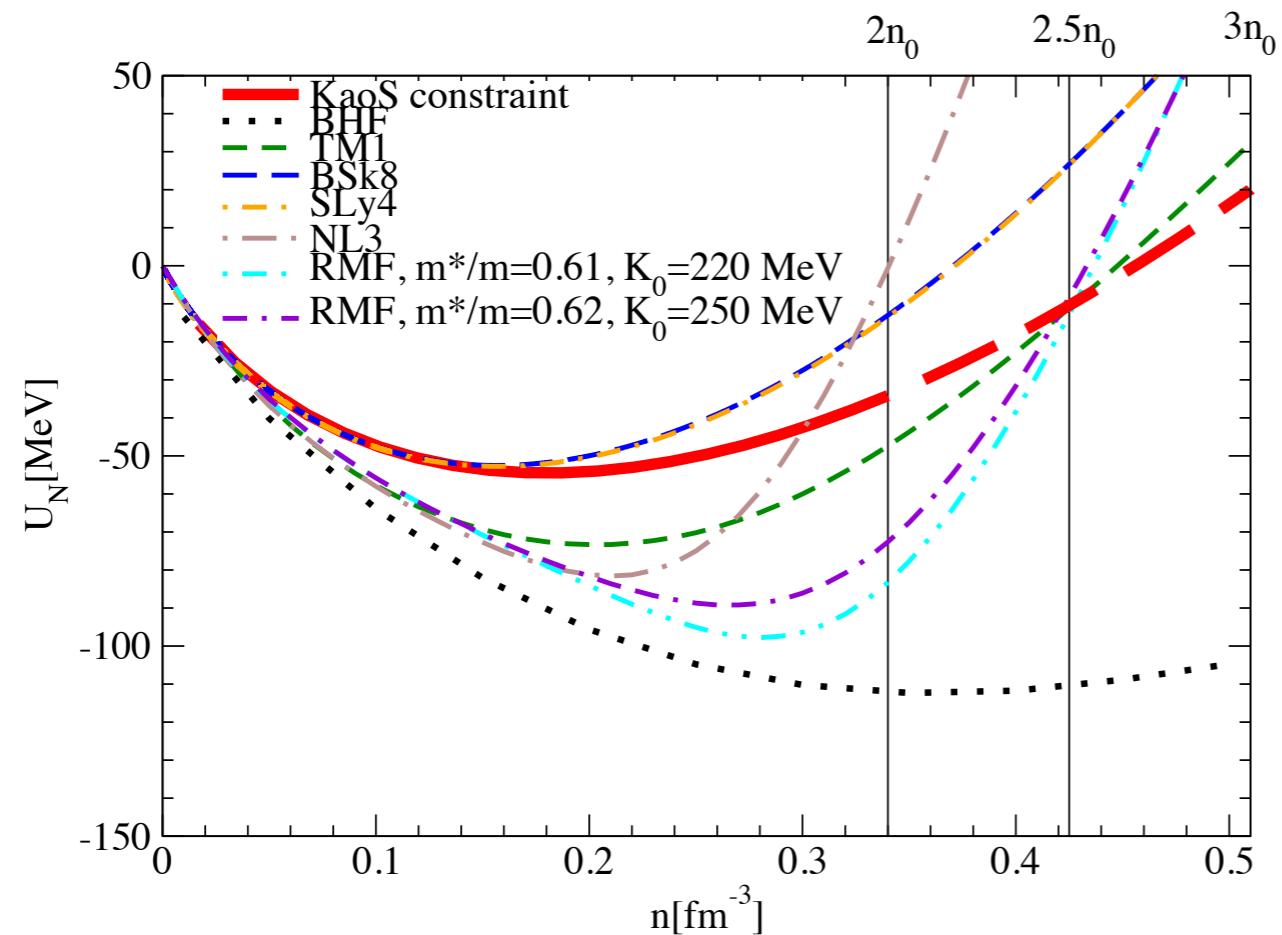
- * *Relativistic Mean Field Models (RMF)*
 - *with non-linear interaction of mesons, fitted to bulk nuclear matter (GL, TM1)*
 - *fitted to properties of nuclei (NL3)*
- * *Brueckner Hartree Fock models (BHF)*
 - *realistic N-N interactions*
- * *Phenomenological models*
 - *Skyrme interactions (Bsk8, SLy4)*

the KaoS constraint



*I. Sagert, L.Tolos, D.C., J. Schaffner-Bielich and C. Sturm,
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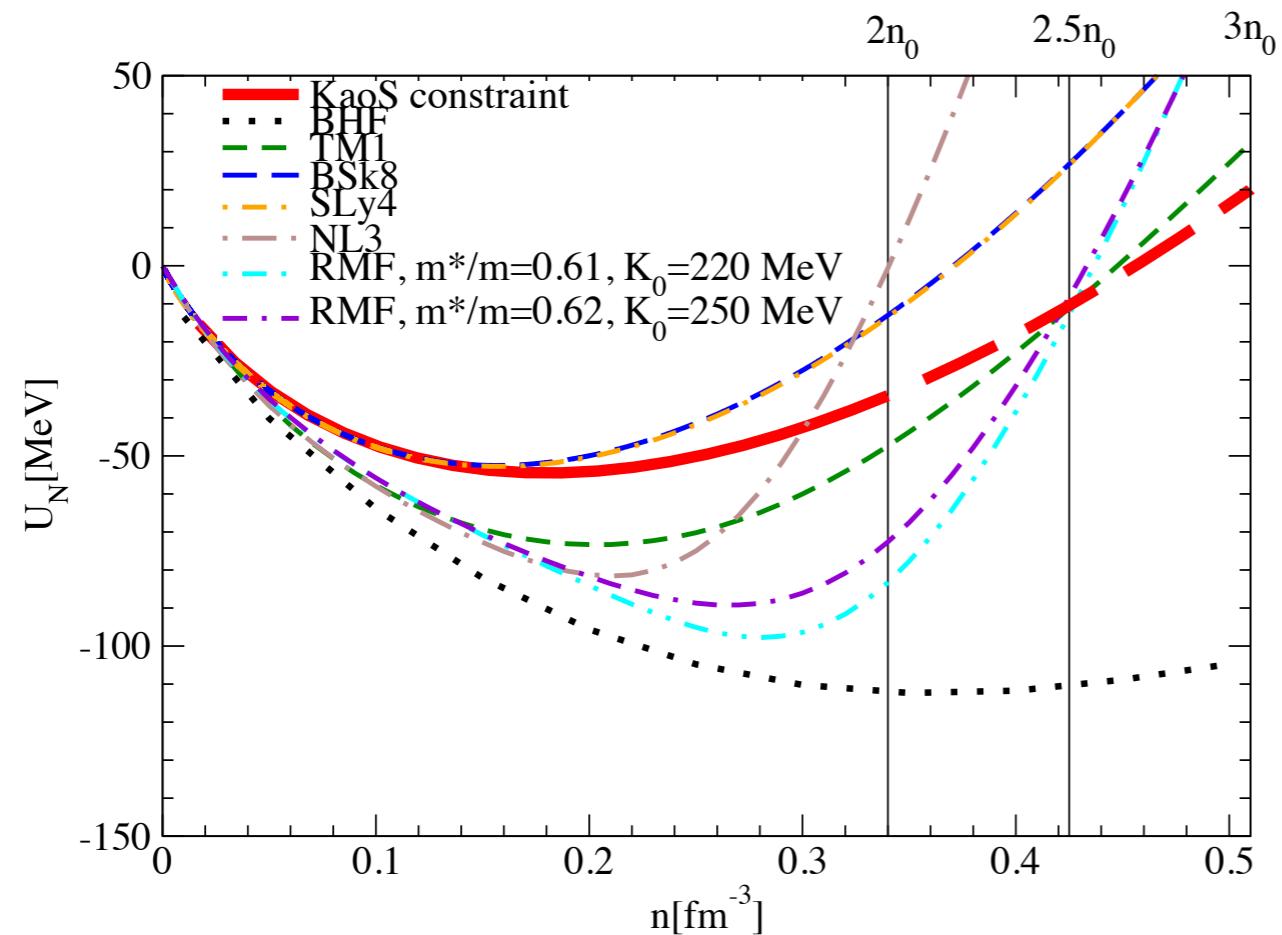
the KaoS constraint



- the EoS is chosen so as to obtain a nucleon potential similar to or more attractive than the KaoS constraint within the density limits

I. Sagert, L. Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
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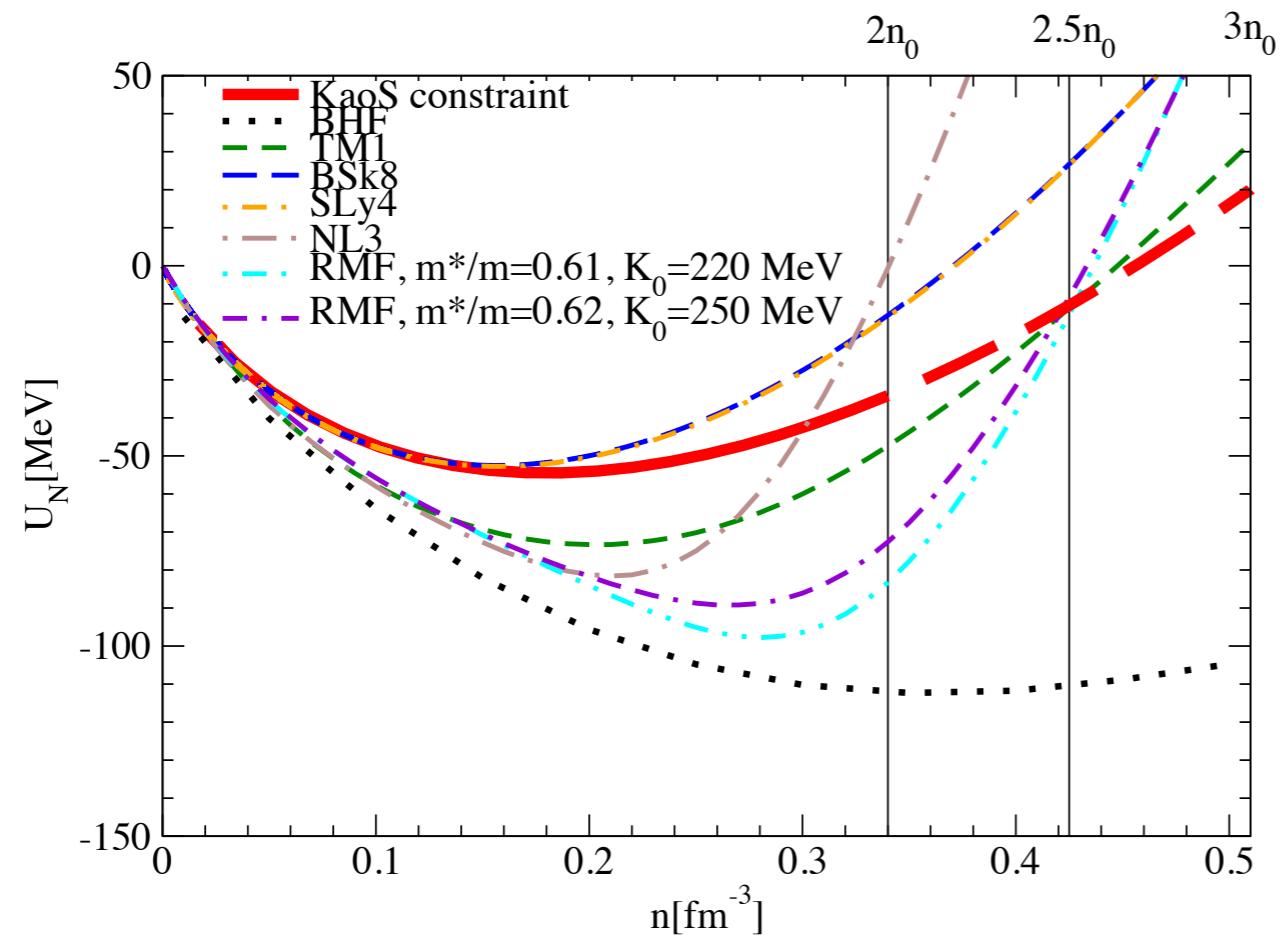


- the EoS is chosen so as to obtain a nucleon potential similar to or more attractive than the KaoS constraint within the density limits

- A more attractive U_N allows a higher compression of matter for the same bombarding energy, enhances multiple scattering processes in subthreshold kaon production.

I. Sagert, L.Tolos, D.C., J. Schaffner-Bielich and C. Sturm,
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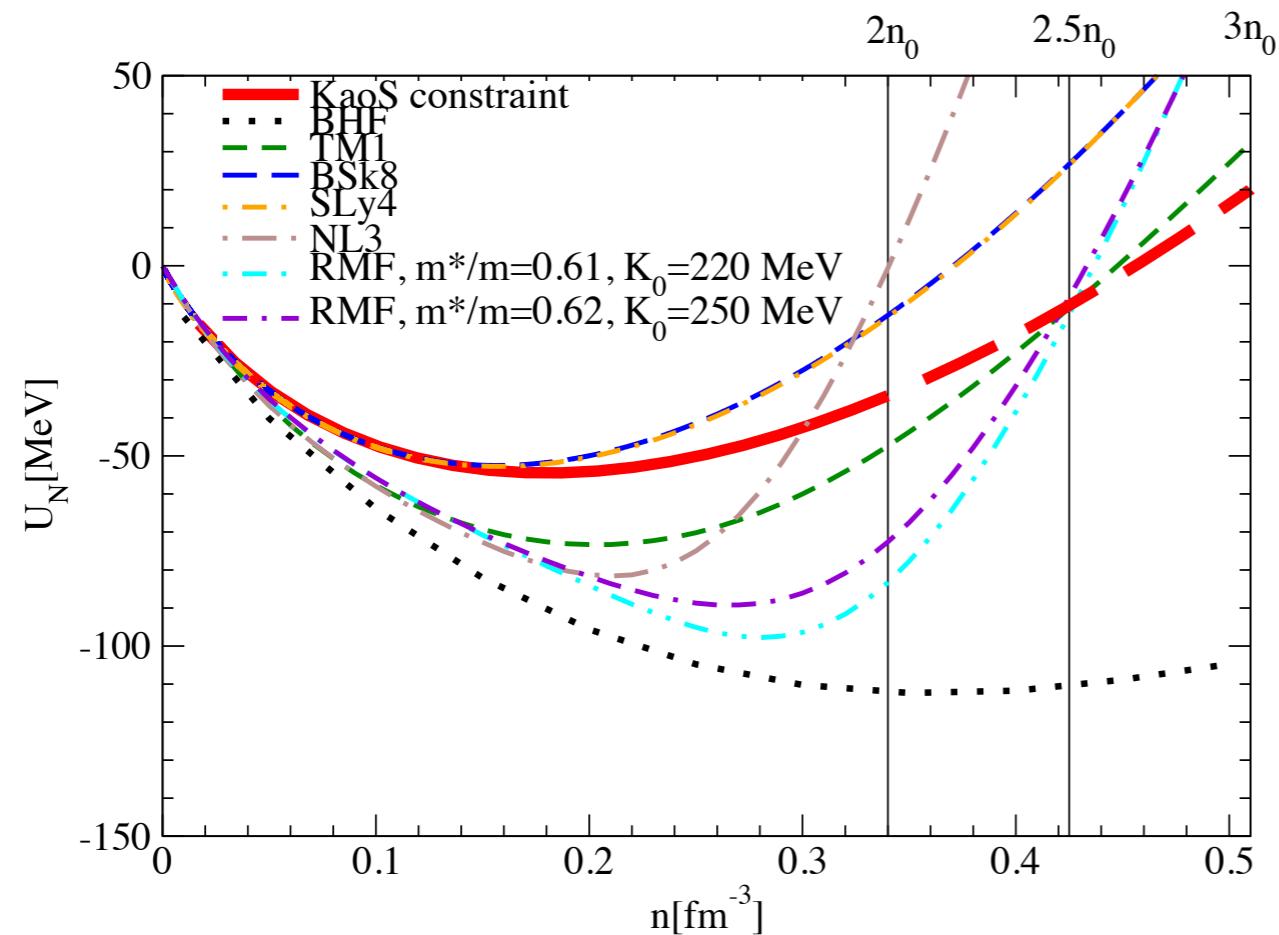
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- For Skyrme type models, $K \leq 200$ MeV

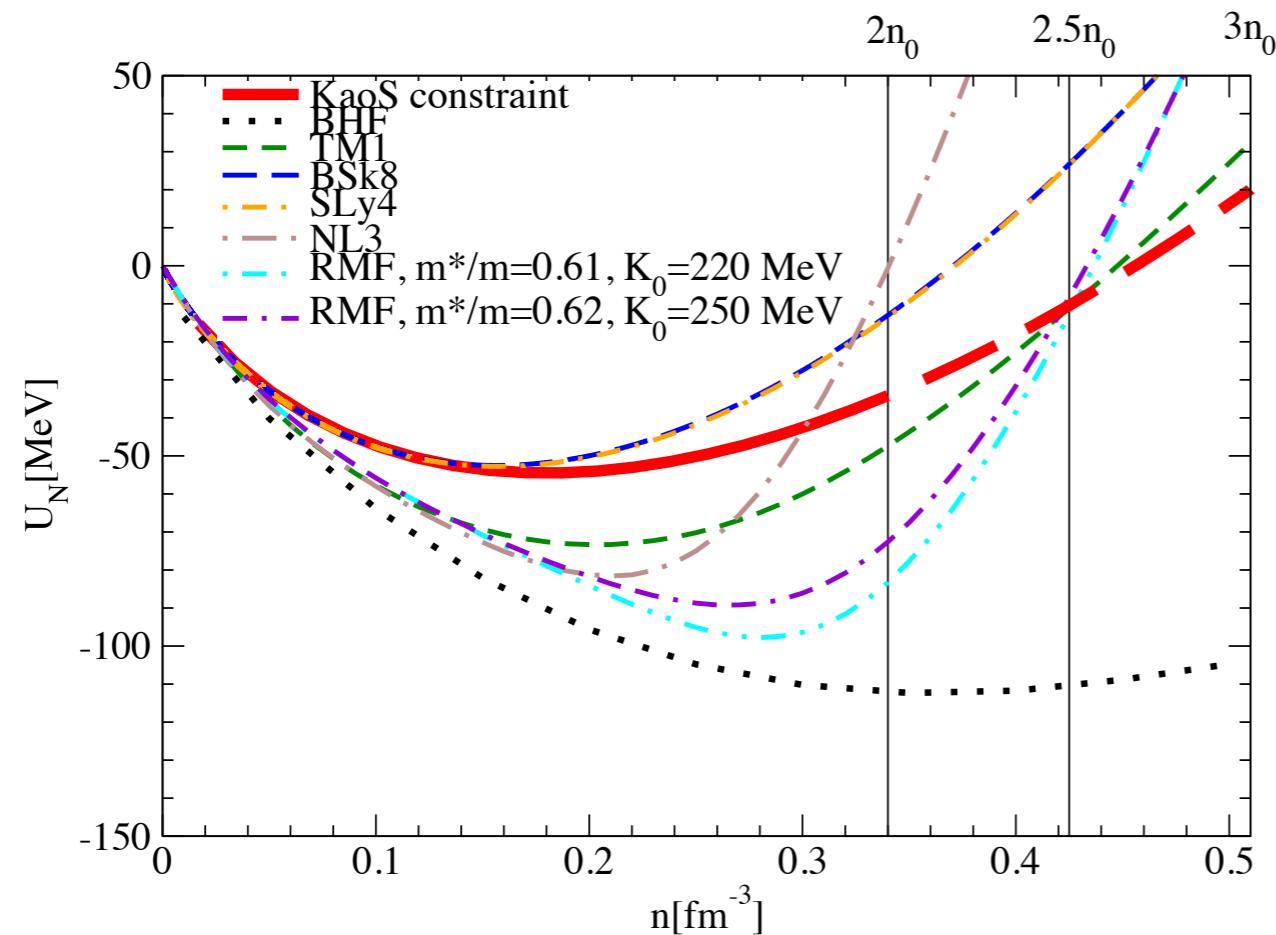
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- For RMF models, stiffness of high density nuclear matter determined by m_* (n_0)
 $\Rightarrow U_N$ is chosen by varying m_* for given K_0

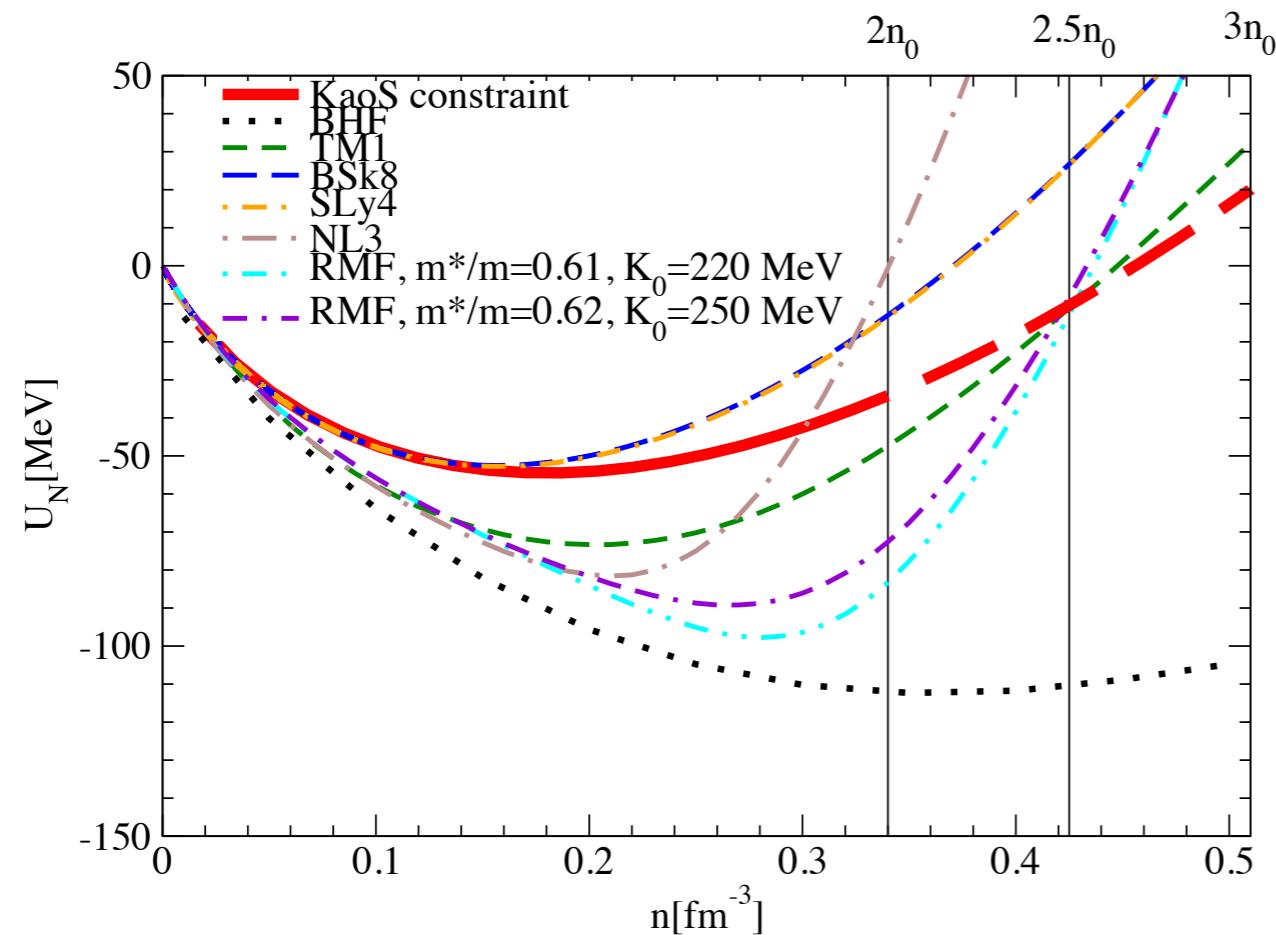
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- TM1, BHF satisfy KaoS constraint

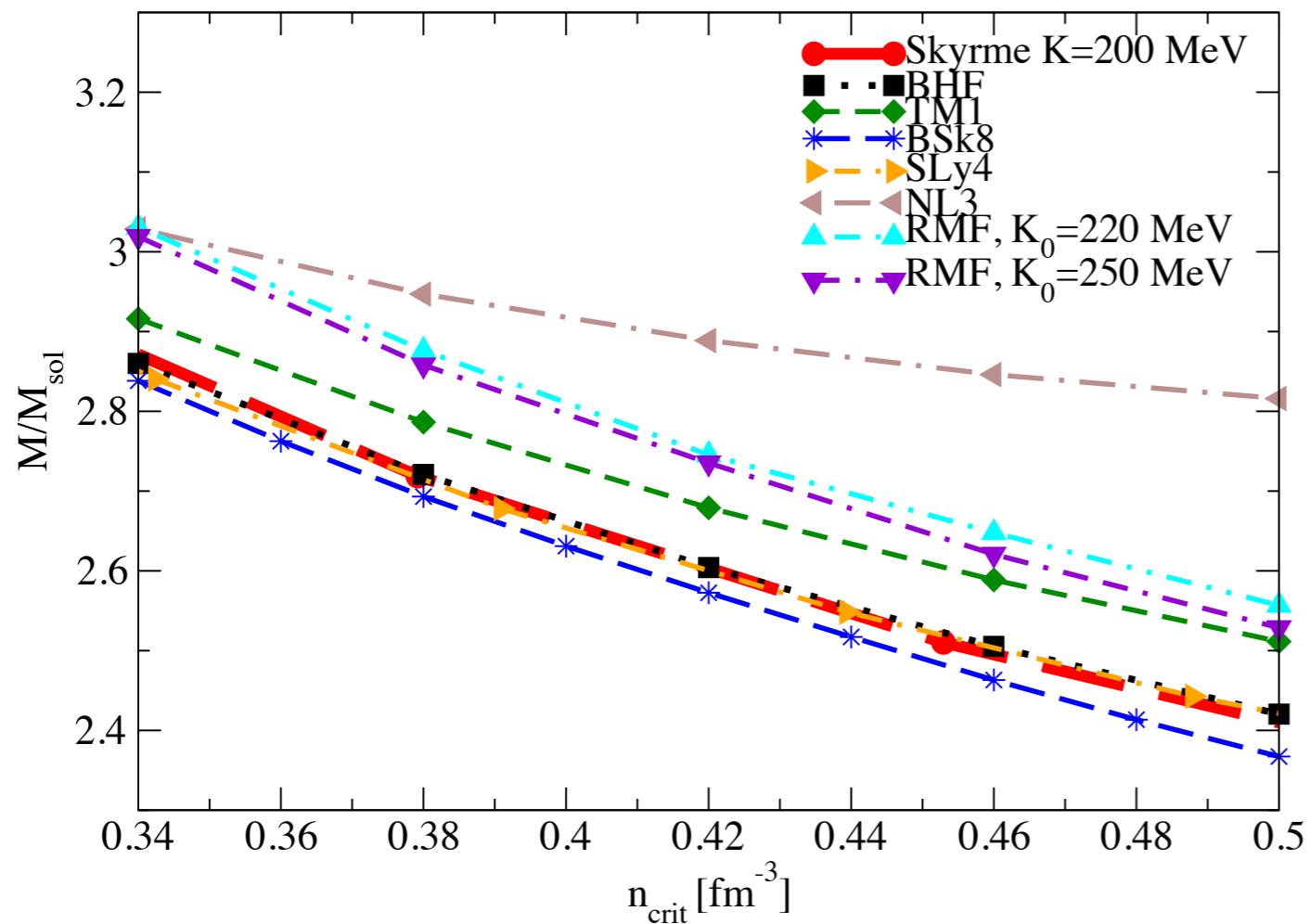
the KaoS constraint



I. Sagert, L.Tolos, D.C., J. Schaffner-Bielich and C. Sturm,
arXiv :1112.0234

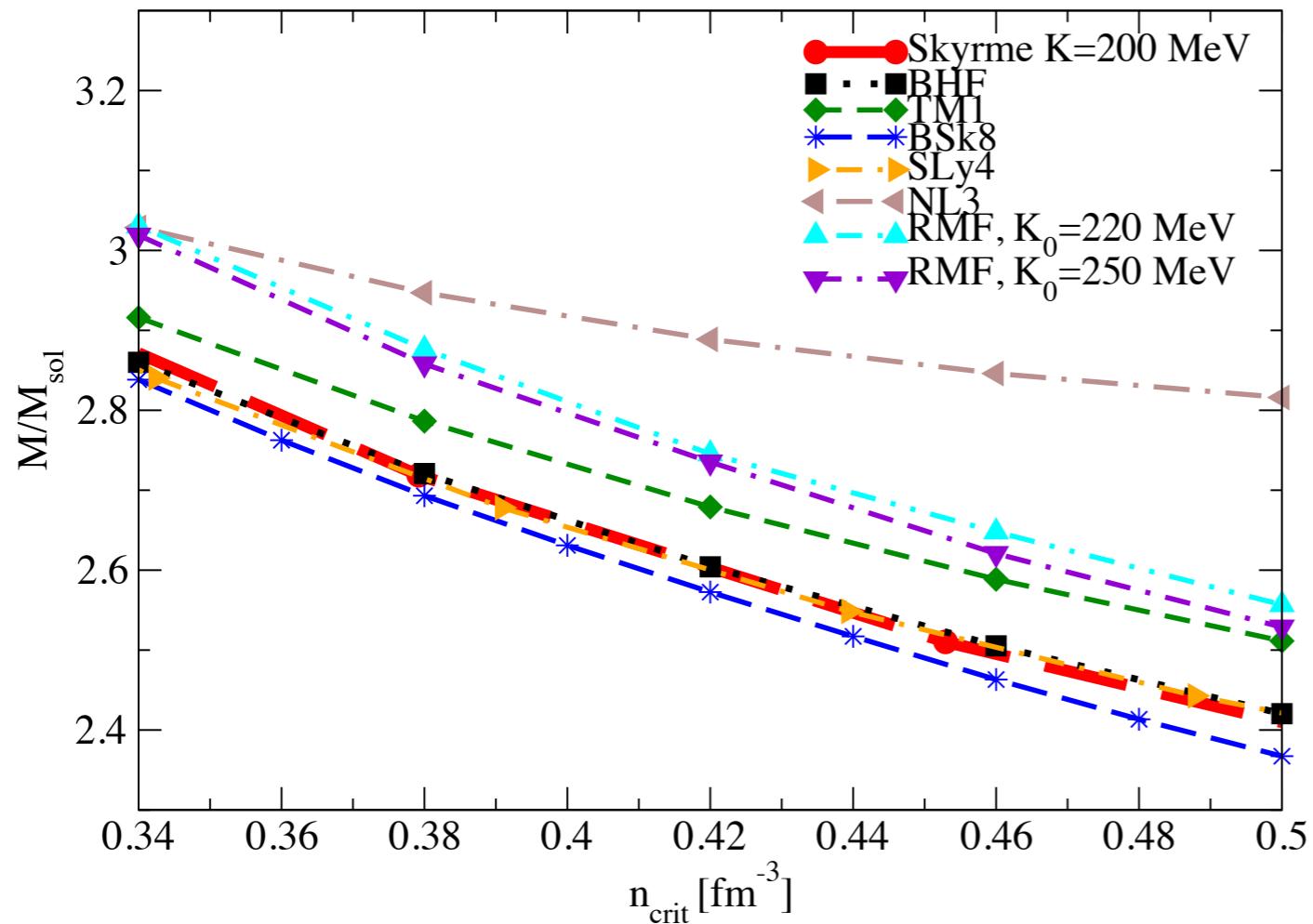
- the EoS is chosen so as to obtain a nucleon potential similar to or more attractive than the KaoS constraint within the density limits
- A more attractive U_N allows a higher compression of matter for the same bombarding energy, enhances multiple scattering processes in subthreshold kaon production.
- For Skyrme type models, $K \leq 200 \text{ MeV}$
- For RMF models, stiffness of high density nuclear matter determined by m_* (n_0)
 $\Rightarrow U_N$ is chosen by varying m_* for given K_0
- TM1, BHF satisfy KaoS constraint
- BSk8, SLy4 or NL3 have more repulsive potentials and are ruled out

Maximally allowed gravitational mass



I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
arXiv :1112.0234

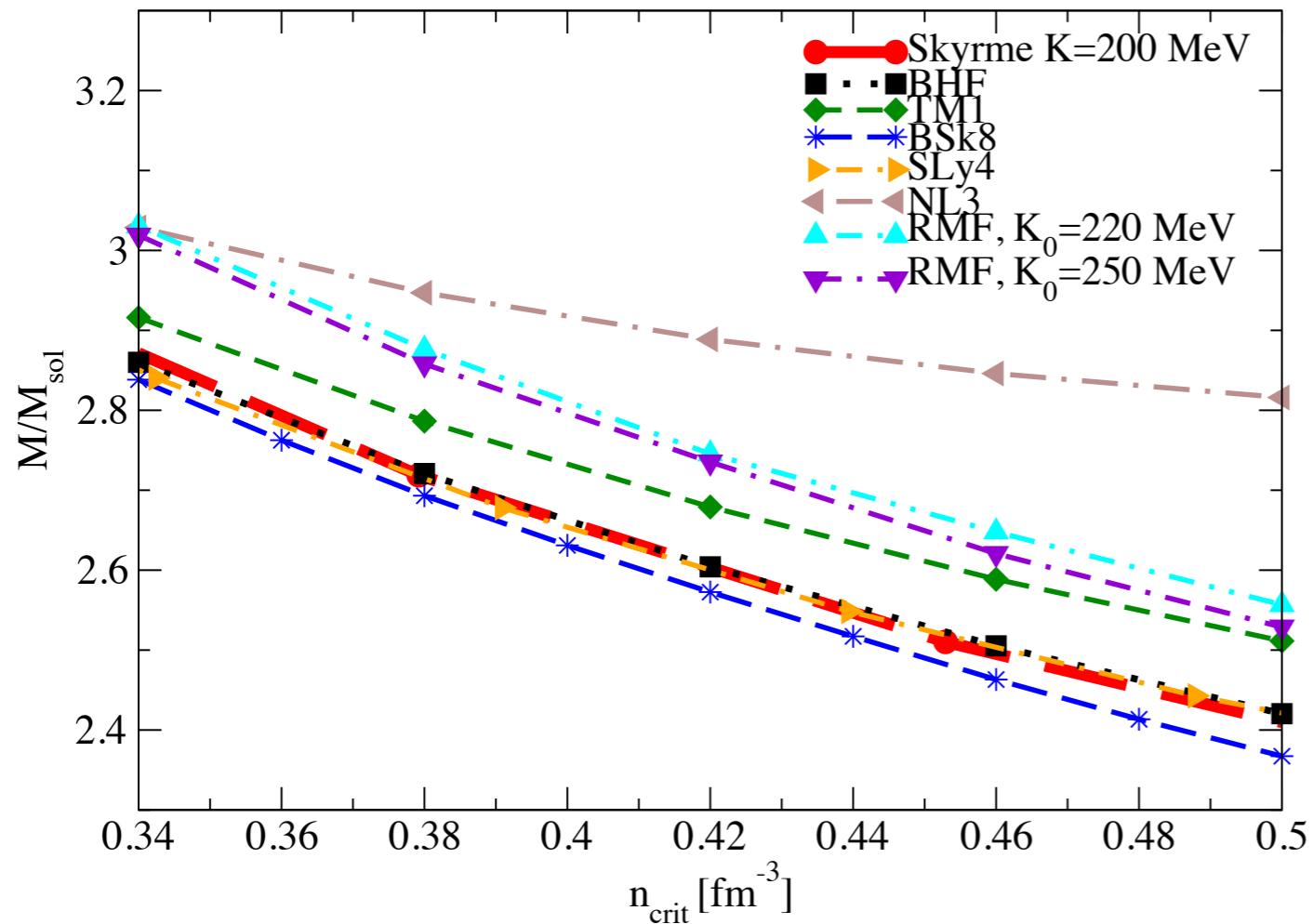
Maximally allowed gravitational mass



- highest allowed neutron star mass $3M_{\text{sol}}$ at $n_{\text{crit}} \sim 2 n_0$

I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
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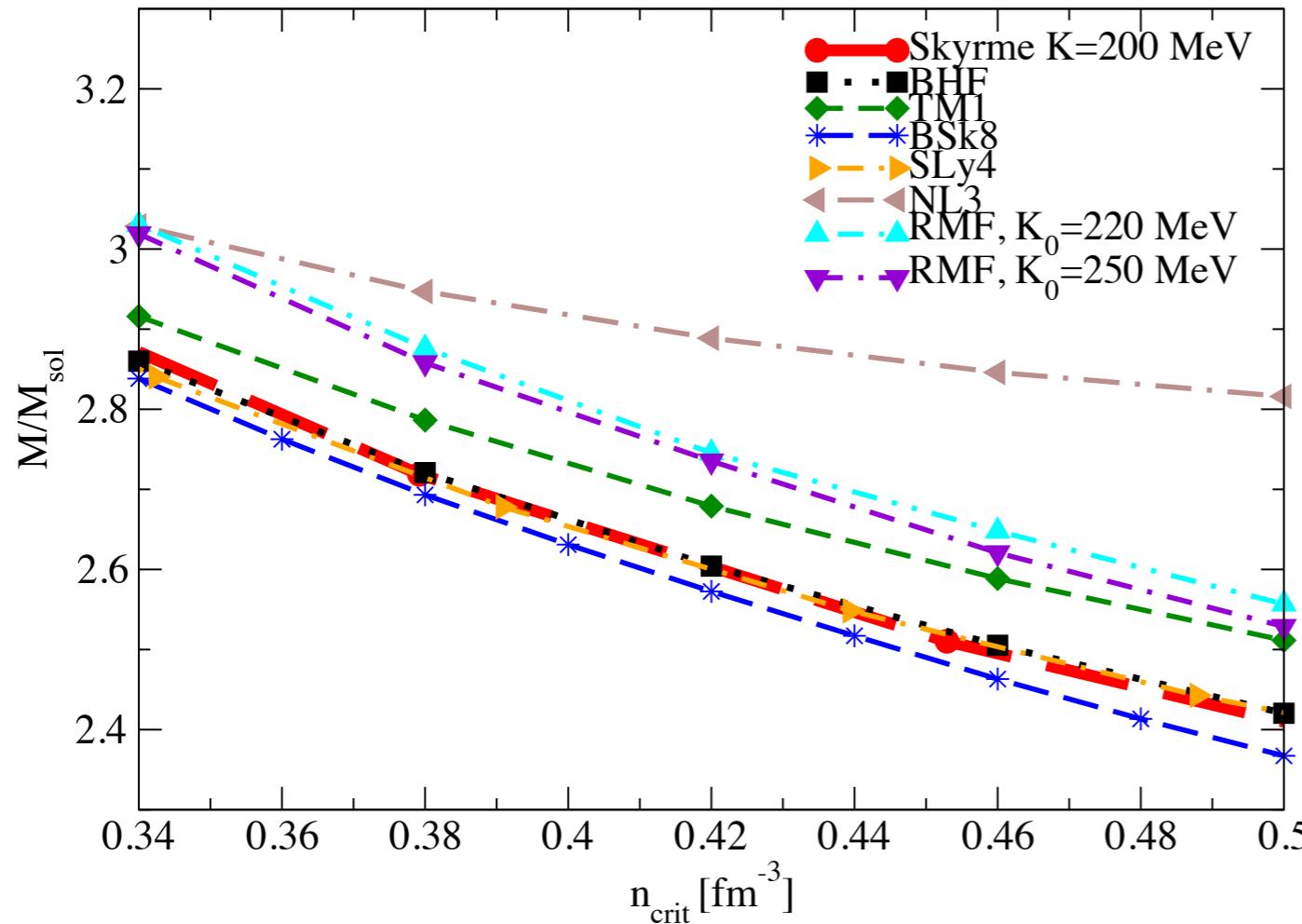
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- highest allowed neutron star mass $3M_{\text{sol}}$ at $n_{\text{crit}} \sim 2 n_0$
- Smaller maximum masses are obtained for $n_{\text{crit}} \sim 3 n_0$

I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
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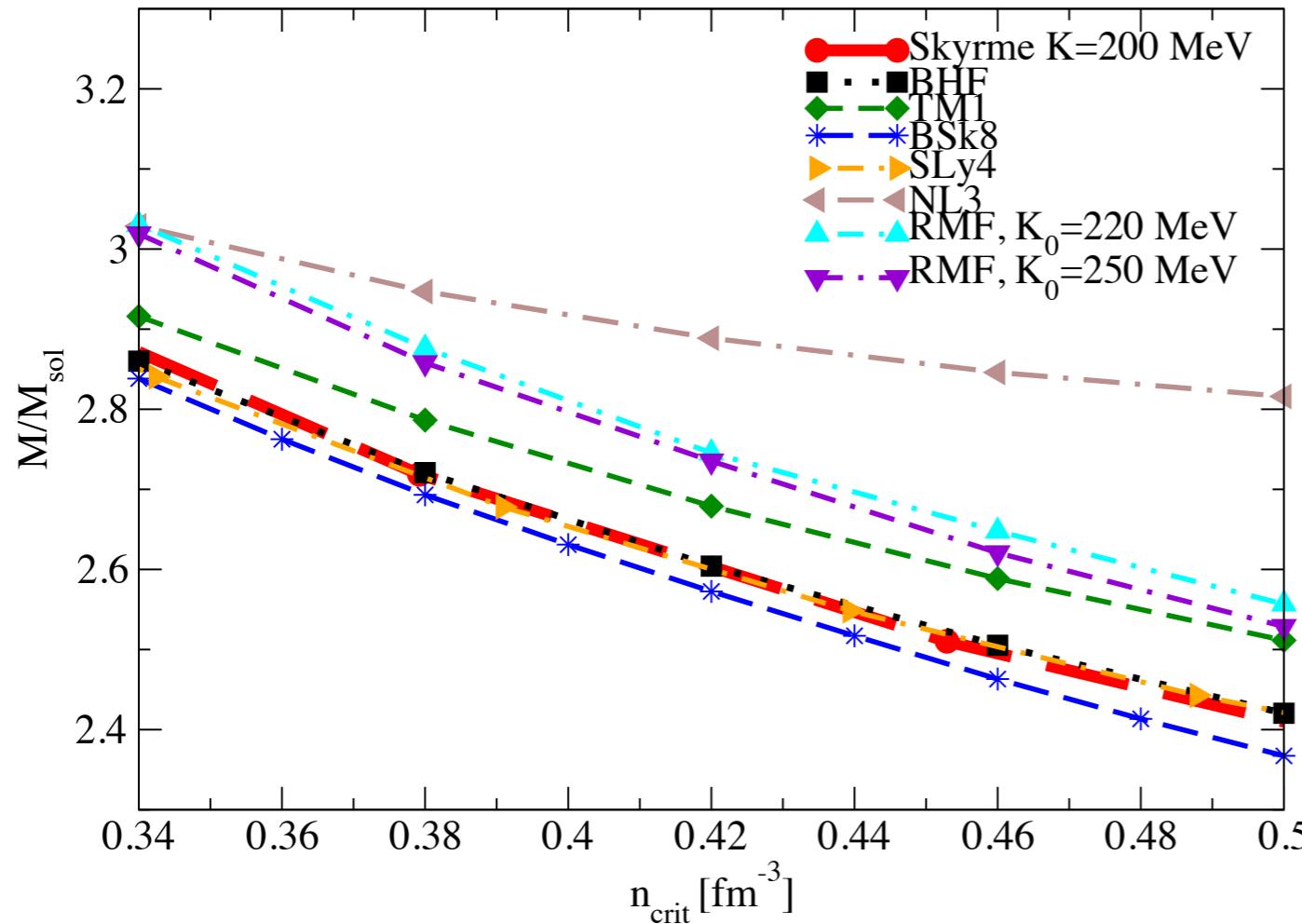
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I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
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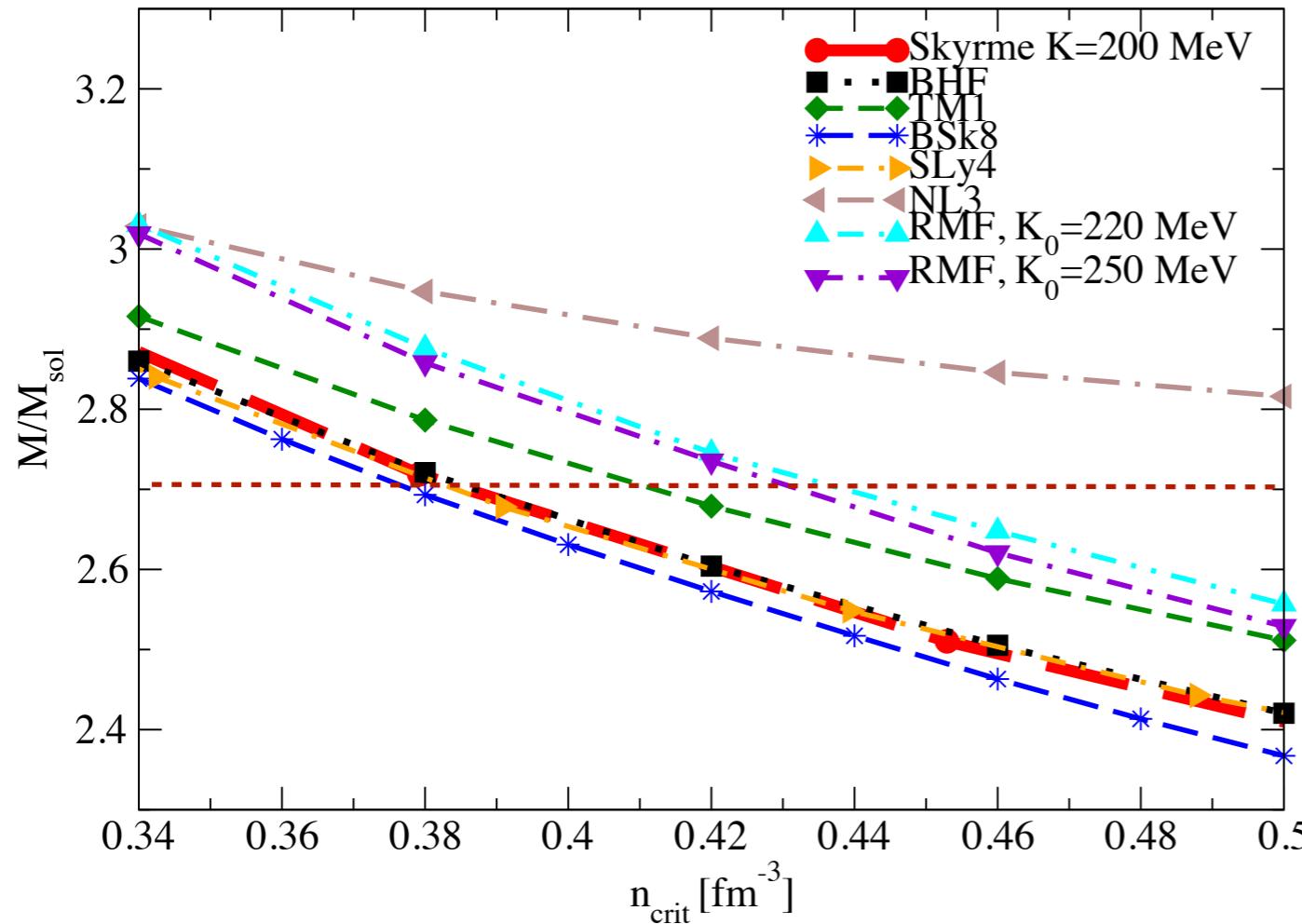
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I. Sagert, L.Tolos, D. C., J. Schaffner-Bielich and C. Sturm,
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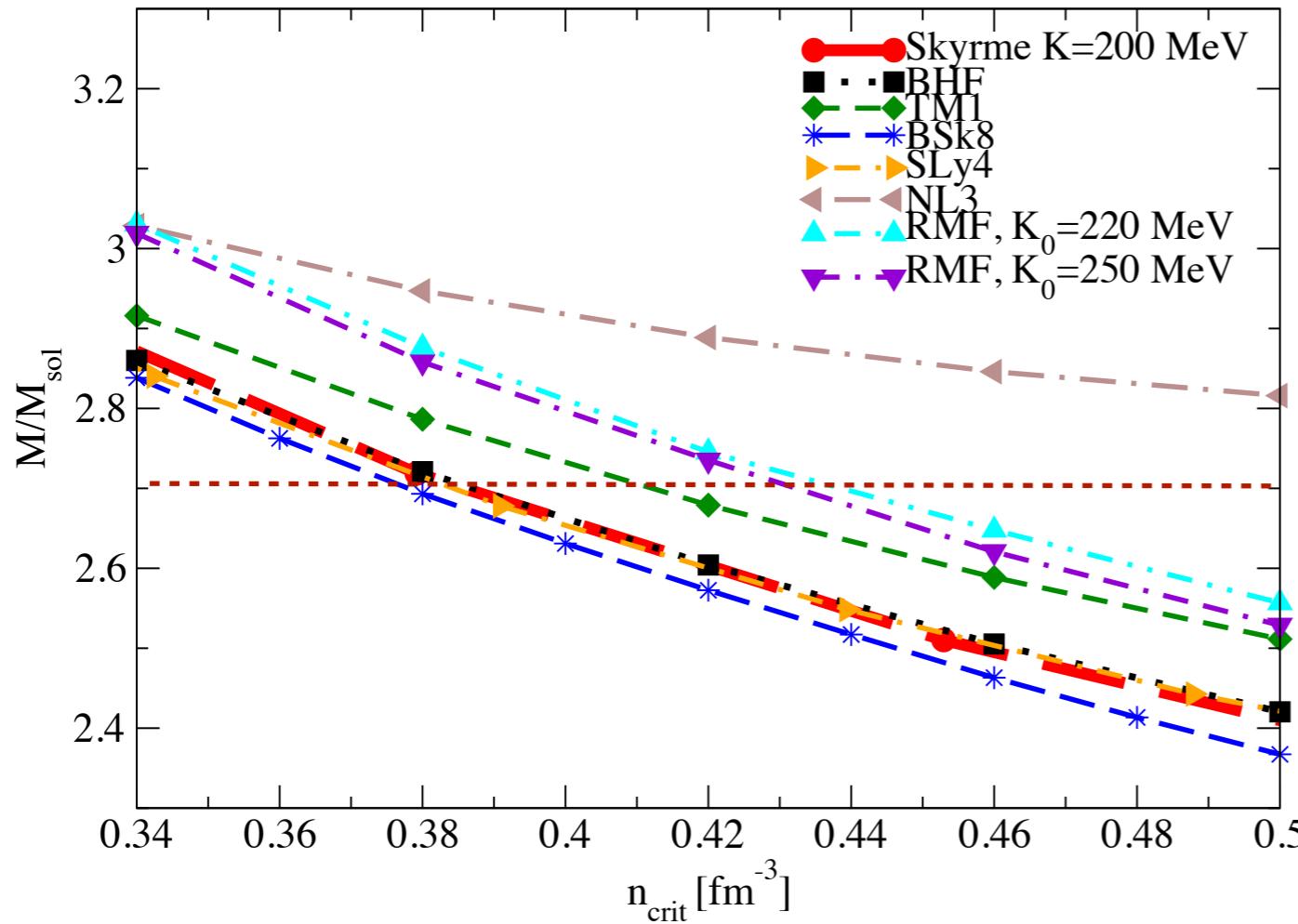
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Maximally allowed gravitational mass



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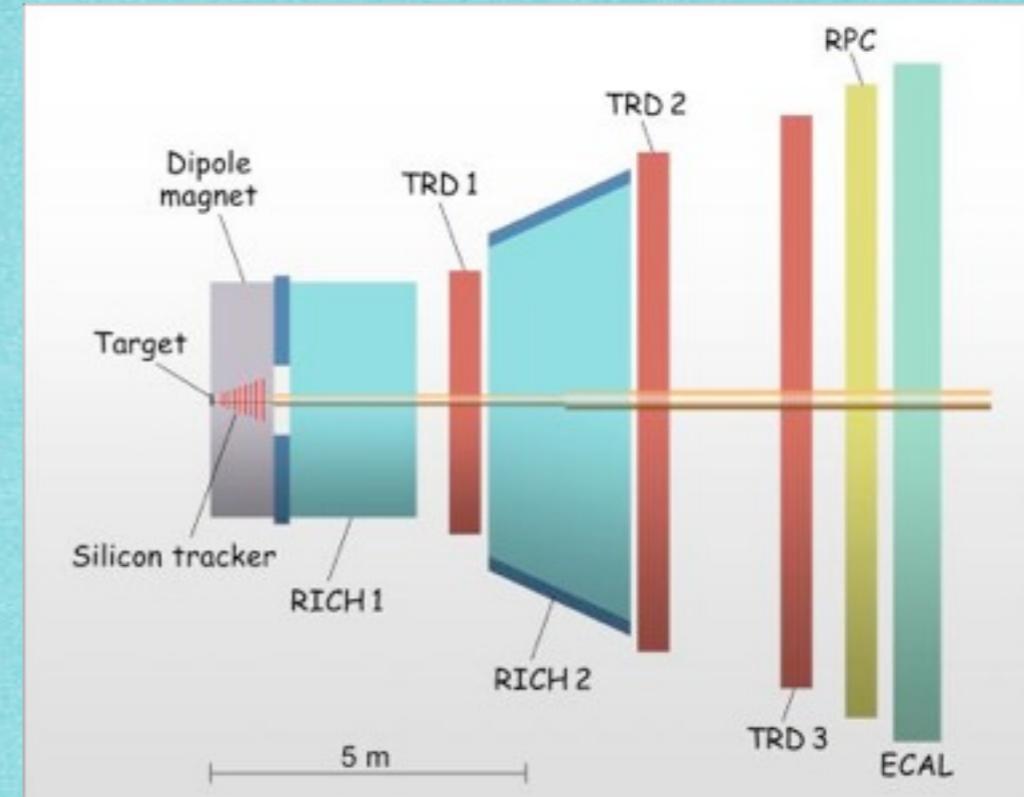
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- a pulsar of $2.7M_{\text{sol}}$ not ruled out by KaoS data, but requires a fiducial density of $\sim 2.2 - 2.5 n_0$
- Since the maximum mass configuration is dominated by the causal high density EoS, the symmetry energy has very little influence on M_{max} : $\Delta M \approx 0.02M_{\text{sol}}$

Summary

- K^+ multiplicities from heavy-ion collisions indicate a soft nuclear EoS for densities of 2-3 n_0
- We test the implications of results on neutron stars
- Light neutron stars with $M \sim 1.25 M_{\text{sol}}$ have central densities $\sim 2-3 n_0$
- Measurement of radii of low mass neutron stars can test KaoS results
- To test if soft nuclear EoS is compatible with massive neutron stars, we apply KaoS results at densities up to 2-3 n_0 , and then introduce the stiffest possible causal EoS to calculate the highest allowed maximum neutron star mass
- KaoS results indicate highest possible neutron star mass of $3 M_{\text{sol}}$
- The massive pulsar of $2.7 M_{\text{sol}}$ requires an onset of the stiffest possible EoS at a fiducial density of $\sim 2.2 - 2.5 n_0$.

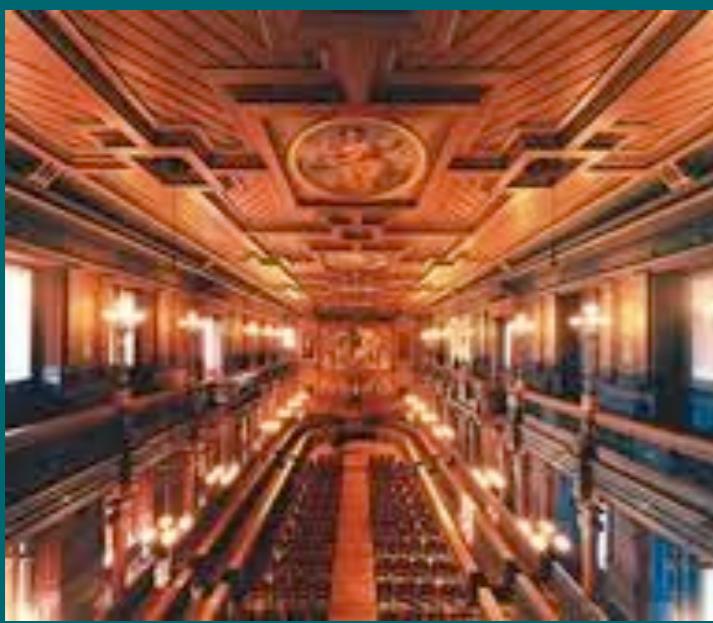
“Soft equation of state from heavy-ion data and implications for compact stars”
I. Sagert, L. Tolos, D.C., J. Schaffner-Bielich and C. Sturm, arXiv:1112.0234

Outlook



The *CBM* (*Condensed Baryonic Matter*) experiment at FAIR will probe densities beyond $3 n_0$, using rare probes such as D-meson and provide better constraints on the maximum neutron star mass.

THANK YOU FOR YOUR ATTENTION!

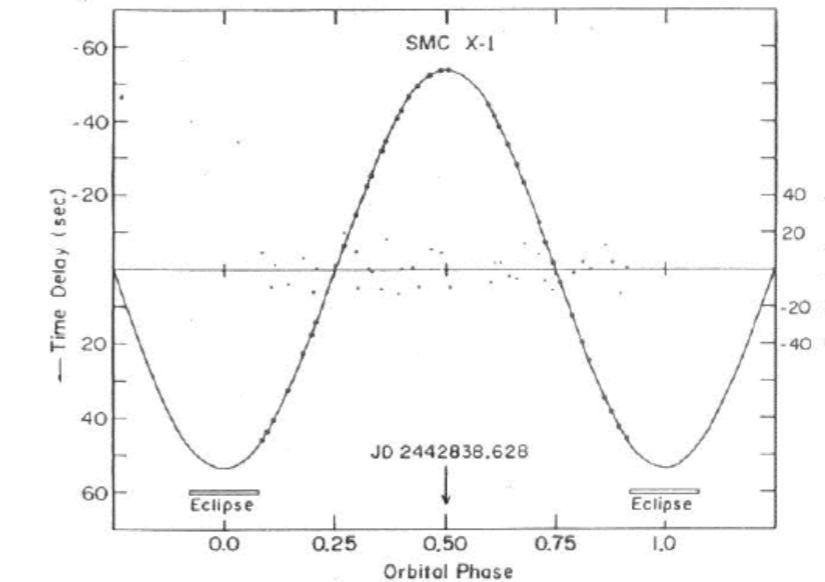


Mass measurements: *in eclipsing X-ray Binaries*

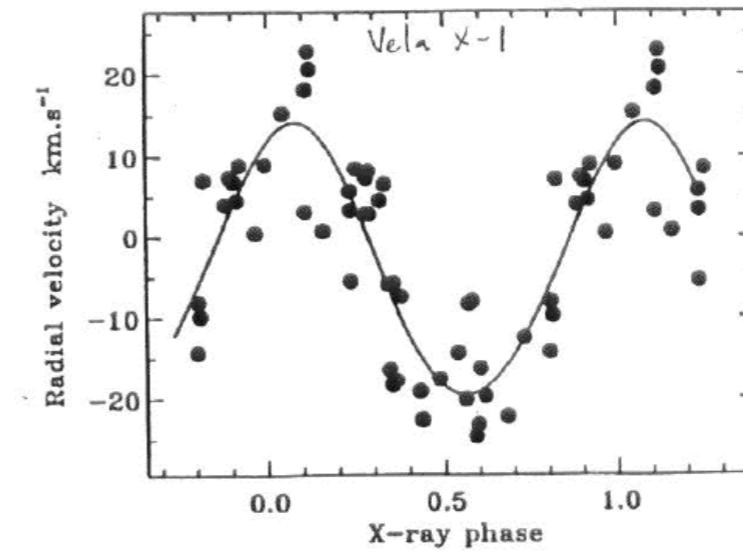
Mass function:

$$f(M_D) = \frac{P(V_x \sin i)^3}{2\pi G} = \frac{(M_D \sin i)^3}{(M_x + M_D)^2}$$

- Doppler shifts of X-ray pulse period
⇒ $V_x \sin i$
- Doppler shifts of Companion's spectral features ⇒ $V_D \sin i$
⇒ $f(M_X)/f(M_D)$
- eclipse ⇒ $\sin i \sim 1$

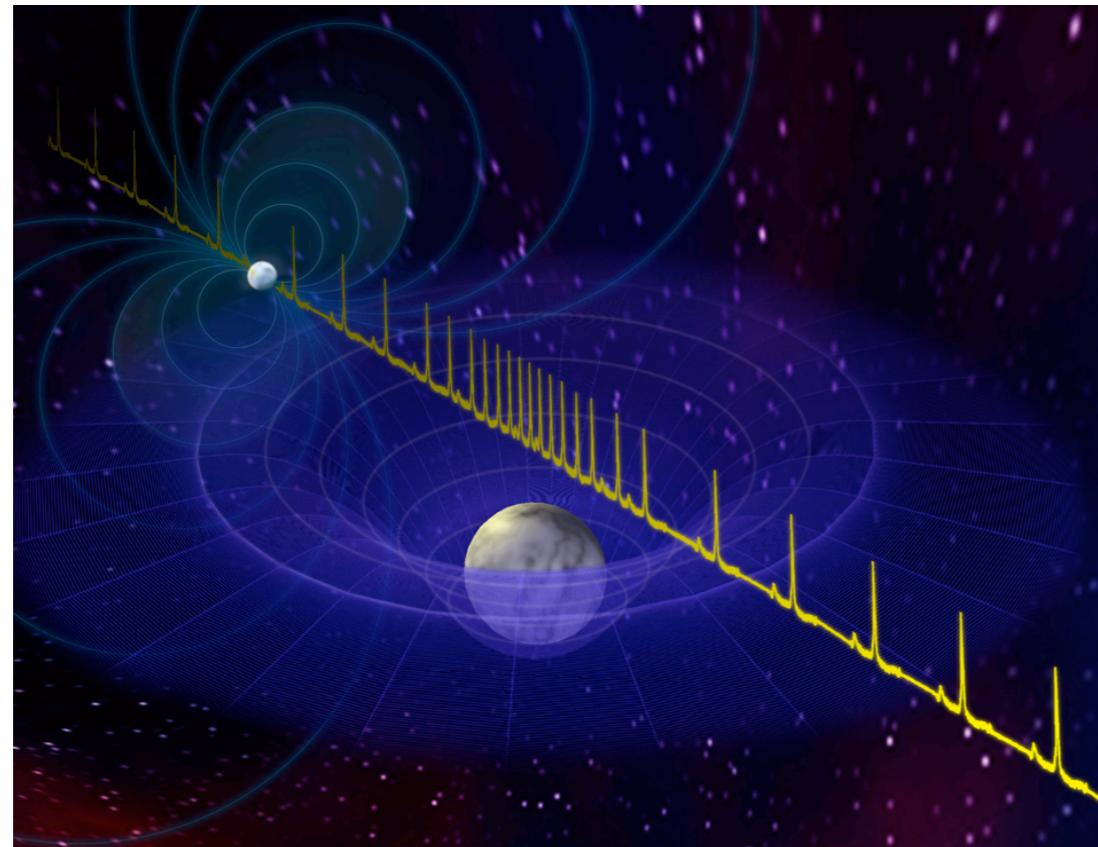


Radial velocity of X-ray pulsar



Radial velocity of optical companion

Shapiro Delay



Credit: Bill Saxton/NRAO

- If there is just one pulsar in the system, it is only possible to measure the mass ratio of a pulsar to its companion M_p/M_c
- If the system is nearly edge on, as the pulse train passes close to the companion, it experiences *Shapiro* delay in the pulses.
- The magnitude and duration of the delay episode is related to the inclination of the binary orbit to the line of sight, and the mass of the companion.
- This completely determines the mass of the pulsar.

Relativistic Mean Field Model

*Properties of asymmetric nuclear matter
at saturation*

$$\begin{aligned}\mathcal{L} = & \sum_B \bar{\psi}_B (i\gamma_\mu \partial^\mu - m_B + g_{\sigma B} \sigma - g_{\omega B} \gamma_\mu \omega^\mu) \\ & - \frac{1}{2} g_{\rho B} \gamma_\mu \vec{\tau}_B \cdot \vec{\rho}^\mu \psi_B \\ & + \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma - m_\sigma^2 \sigma^2) - U(\sigma) \\ & - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu \\ & - \frac{1}{4} \rho_{\mu\nu} \cdot \rho^{\mu\nu} + \frac{1}{2} m_\rho^2 \rho_\mu \cdot \rho^\mu + \mathcal{L}_{YY} \\ & + \sum_{e^-, \mu^-} \bar{\psi}_\lambda (i\gamma_\mu \partial^\mu - m) \psi_\lambda.\end{aligned}$$

where, $U(\sigma) = \frac{1}{3} b m_N (g_{\sigma N} \sigma)^3 + \frac{1}{4} c (g_{\sigma N} \sigma)^4$.

- * saturation density $n_0 = 0.17 \text{ fm}^{-3}$
- * binding energy per nucleon $B/A = -16.3 \text{ MeV}$
- * incompressibility $K_0 = 200-250 \text{ MeV}$
- * symmetry energy $E_{\text{sym}} = 32.5 \text{ MeV}$
- * effective nucleon mass $m^*/m = 0.55-0.8$

TM1 Model

$$\begin{aligned} \mathcal{L} = & \sum_B \bar{\Psi}_B (i\gamma_\mu \partial^\mu - m_B + g_{\sigma B} \sigma - g_{\omega B} \gamma_\mu \omega^\mu - g_{\rho B} \gamma_\mu t_B \cdot \boldsymbol{\rho}^\mu) \Psi_B \\ & + \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma - m_\sigma^2 \sigma^2) - U(\sigma) + U(\omega) \\ & - \frac{1}{4} \omega_{\mu\nu} \omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - \frac{1}{4} \boldsymbol{\rho}_{\mu\nu} \cdot \boldsymbol{\rho}^{\mu\nu} + \frac{1}{2} m_\rho^2 \boldsymbol{\rho}_\mu \cdot \boldsymbol{\rho}^\mu. \end{aligned}$$

$$U(\sigma) = \frac{1}{3} b \sigma^3 + \frac{1}{4} c \sigma^4,$$

$$U(\omega) = \frac{1}{4} d (\omega_\mu \omega^\mu)^2,$$

Nuclear equations of state

Phenomenological:

$$\begin{aligned} U(n_b) &= \frac{A}{2} \left(\frac{n_b}{n_0} \right) + \frac{B}{\sigma+1} \left(\frac{n_b}{n_0} \right)^\sigma \\ E_{sym}(n_b) &\sim S_0 \left(\frac{n_b}{n_0} \right)^\alpha \end{aligned}$$

- BE = -16 MeV, $n_0 \sim 0.16 \text{ fm}^{-3}$
- $K_0 = (160 - 240) \text{ MeV}$
- $S_0 = (28 - 32) \text{ MeV}$, $\alpha = 0.7 - 1.1$

Skyrme and rel. mean field:

	S_0 [MeV]	K_0 [MeV]
Bsk8 (NPA 750 (2005))	28.0	230.2
Sly4 (NPA 635 (1998))	32.0	229.9
TM1 (NPA 579 (1994))	36.9	281