



The 4th **A**sian **T**riangle **H**eavy **I**on **C**onference  
in Pusan, South Korea on November 14-17, 2012

# Perspectives of search for ultra-strong magnetic field via direct virtual photon measurement with ALICE at the LHC

**Asako Tsuji**

for the ALICE collaboration

Hiroshima Univ.



# Outline

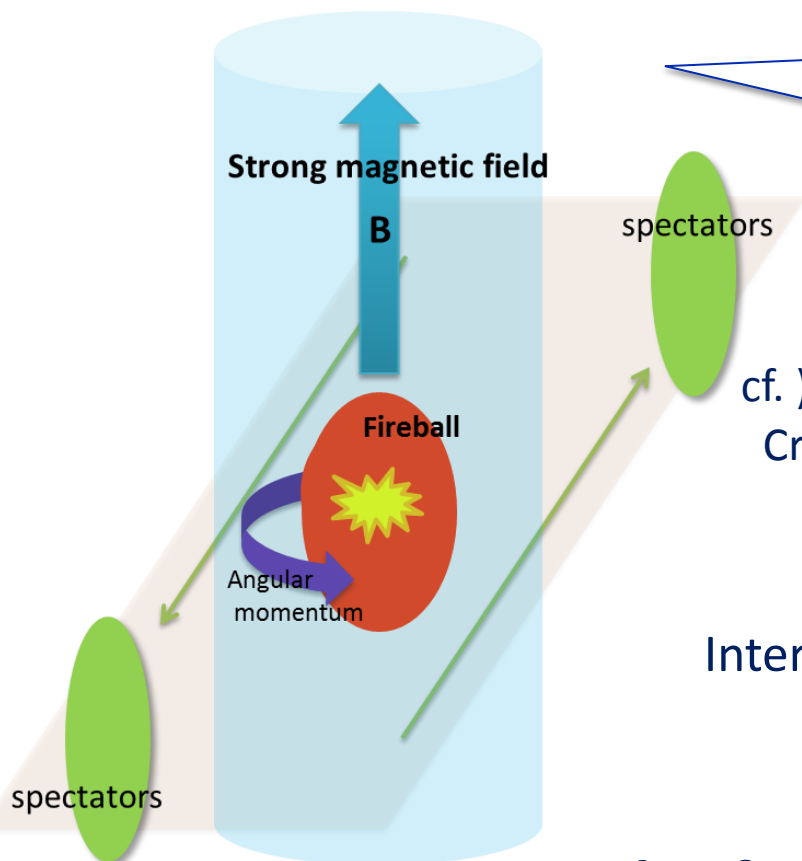
- **Physics Motivation**
  - Ultra-strong magnetic field in HIC
- **Theoretical study**
  - Calculation of vacuum polarization tensor
  - Virtual photon Anisotropy and Polarization estimations
- **Experimental approach**
  - $e^+e^-$  measurement and E.P. determination at ALICE
  - Feasibility of the field detection
- **Ongoing efforts of real data analysis**
- **Summary & Conclusion**

# Physics Motivation

# Ultra-strong magnetic field created in HIC

**B**  $\sim 10^{15}$  Tesla @LHC  
 $\sim 10^{14}$  Tesla @RHIC

**Very Strong!!**



cf.) LHC superconducting magnets : 8.3 Tesla  
Critical magnetic field of electrons  $B_c$  :  $4 \times 10^9$  Tesla  
Magnetars :  $\sim 10^{11}$  Tesla

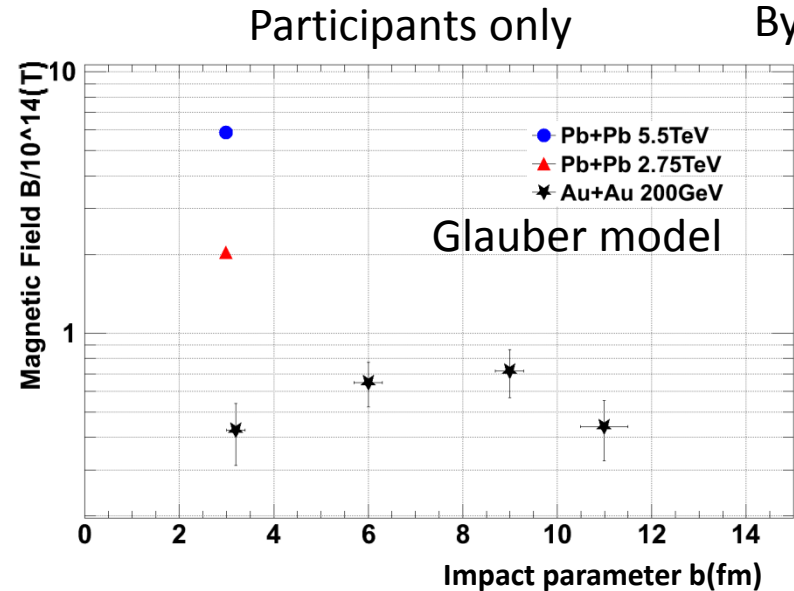
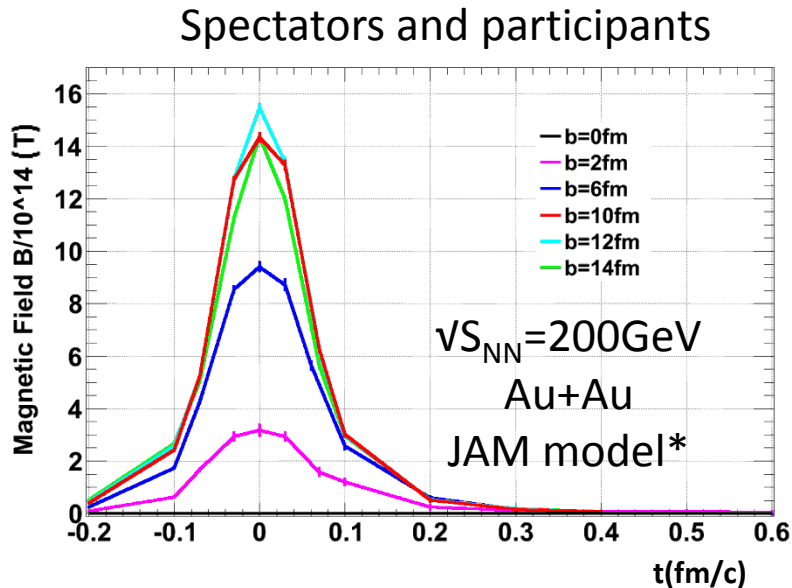
Interesting phenomena : Chiral magnetic effect\*  
Synchrotron radiation  
Non-linear QED effects

**The field is not yet directly detected!!**

\*K.Fukushima, D.E.Kharzeev, H.J.Warringa, PRD78 (2008) 074033

# The field estimation

- Impact parameter, energy, and time dependence

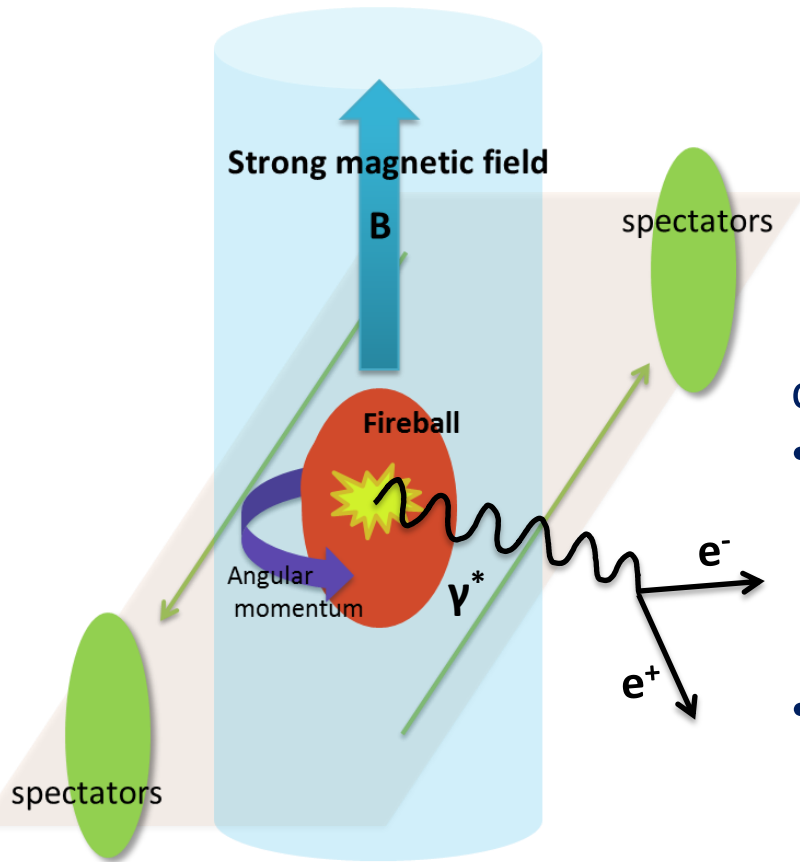


- The field intensity reaches maximum in peripheral collisions ( $B_{\max} \sim 10^{15}$  Tesla) and grows with the beam energy.
- The field rapidly damps, but is still above  $B_c$  for a few fm/c.

\* JAM (Y.Nara, N.Otuka, A.Ohnishi, K.Niita, S.Chiba, PRC61 (2000) 024901)

# A new approach to detect the field

## Via virtual photon (low mass $e^+e^-$ ) measurement



### (1) Anisotropy ( $v_2$ )

### (2) Polarization

cf.)

- Real photon measurement
  - challenging to subtract a large background from hadron decays.
  - no detector to measure polarization.
- Virtual photon (low mass  $e^+e^-$ ) measurement
  - signal to background ratio is dramatically improved by selecting a mass region above  $\pi^0$  mass.
  - achievable with present detectors.

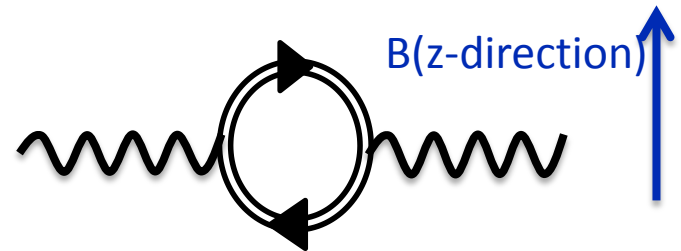
# Theoretical Study

# Calculation of vacuum polarization tensor

- In strong magnetic field

- Photon propagator

$$D_{\mu\nu}(q) = \frac{-i}{q^2} \left[ g^{\mu\nu} - \frac{1}{q^2} \Pi^{\mu\nu}(q, B) \right]^{-1}$$



$q$  : 4D momentum of virtual photon

- Vacuum polarization tensor

$$\Pi^{\mu\nu}(q, B) = \left[ \left( g^{\mu\nu} q^2 - q^\mu q^\nu \right) - \left( g_{\parallel}^{\mu\nu} q_{\parallel}^2 - q_{\parallel}^\mu q_{\parallel}^\nu \right) - \left( g_{\perp}^{\mu\nu} q_{\perp}^2 - q_{\perp}^\mu q_{\perp}^\nu \right) \right] \bar{N}_0 + \left( g_{\parallel}^{\mu\nu} q_{\parallel}^2 - q_{\parallel}^\mu q_{\parallel}^\nu \right) \bar{N}_1 + \left( g_{\perp}^{\mu\nu} q_{\perp}^2 - q_{\perp}^\mu q_{\perp}^\nu \right) \bar{N}_2$$

- When  $q_{\perp}^2 = 0$ , it is integrable for all  $q_{\parallel}^2$  region.

[in terms of DiGamma, Kerbstein et al.

Intl.J.Mod.Phys.[arXiv:1111.5984[hep-ph]].

- For  $q_{\perp}^2 \neq 0$  case, Hattori-Itakura has obtained Landau level summation form recently (Hattori,Itakura [arXiv:1209.2663]).
- One-loop calculation of electron and muon.
- In our calculation,  $\bar{N}_2$  is not included yet.



# Anisotropy of virtual photons

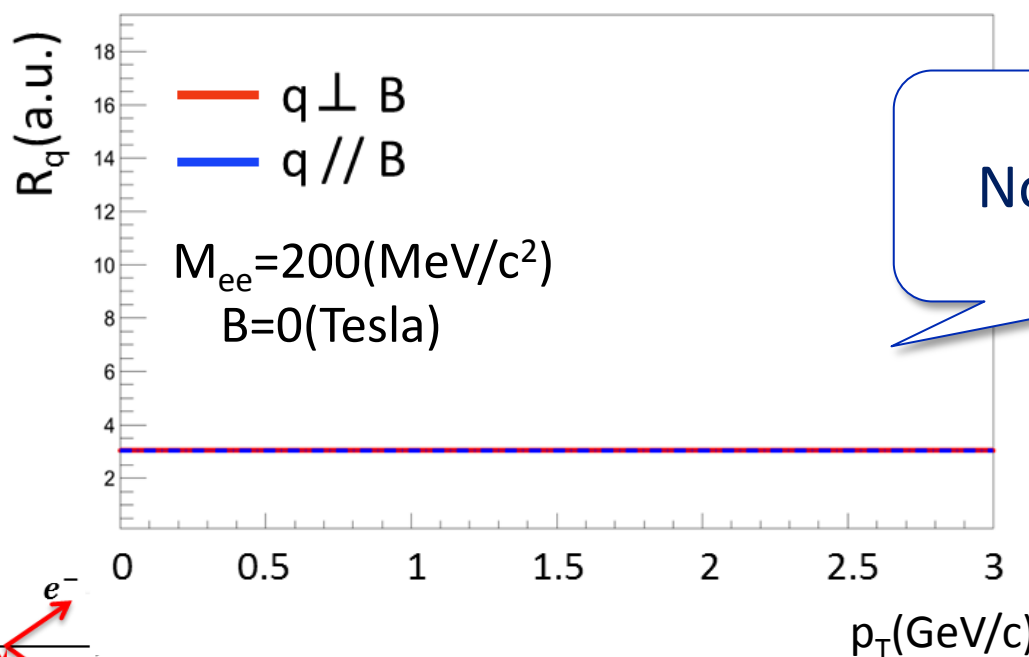
- **Di-electron production rate**

$$R_q \equiv \frac{R_{e^+e^-}}{d^4q} = -\frac{\alpha^2}{3\pi^3} q^2 g^{\mu\alpha} D_{\mu\nu}(q, eB) D_{\alpha\beta}(q, eB)^* \frac{\text{Im} G_R^{\nu\beta}(q, T, eB)}{e^{q^0/T} - 1}$$

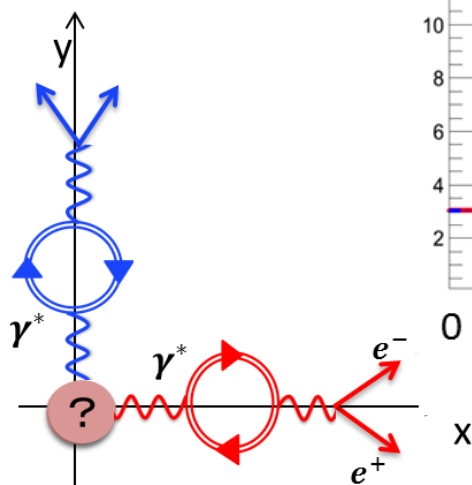
q : 4D momentum of electron pair

G : source of virtual photons

- Last factor is approximated as a constant term.
- **When B=0, the rate is isotropic.**



No anisotropy.



by K.Ishikawa and A.Tsuji

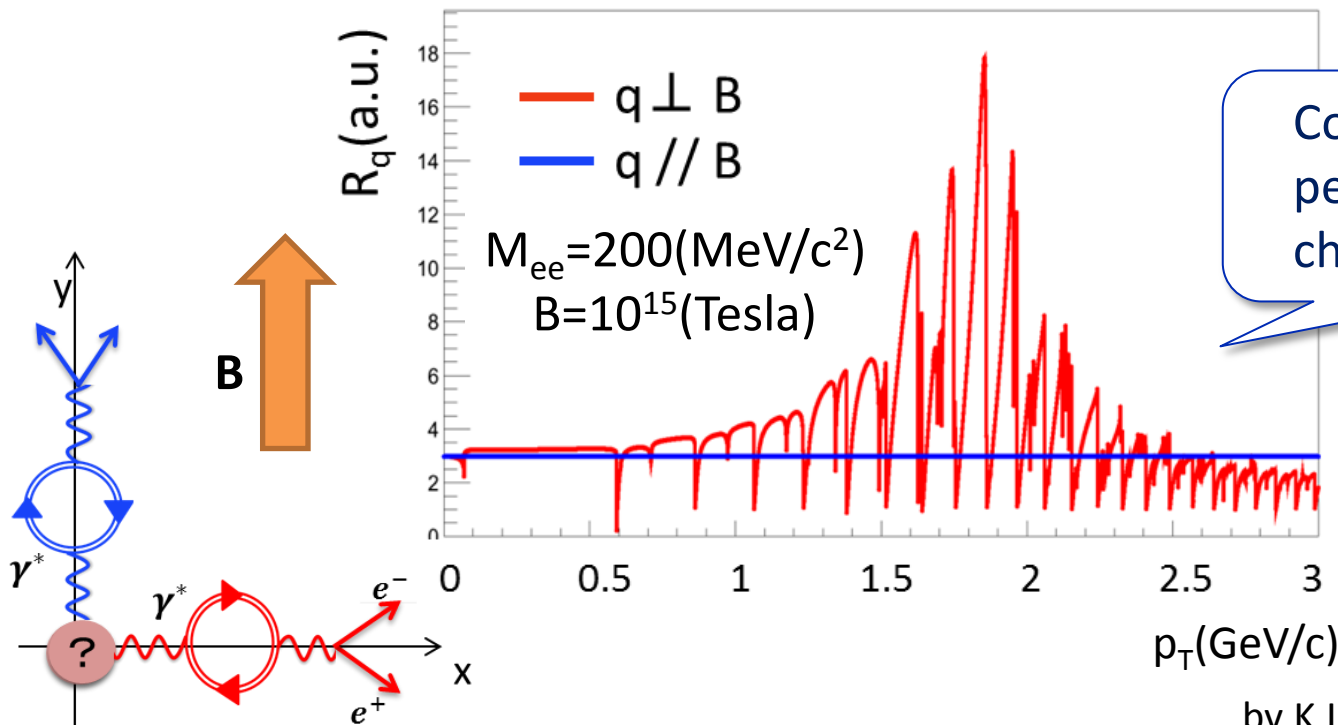
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q : 4D momentum of electron pair  
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- Last factor is approximated as a constant term.
- **For B≠0 case, the rate is anisotropic.**



by K.Ishikawa and A.Tsuji

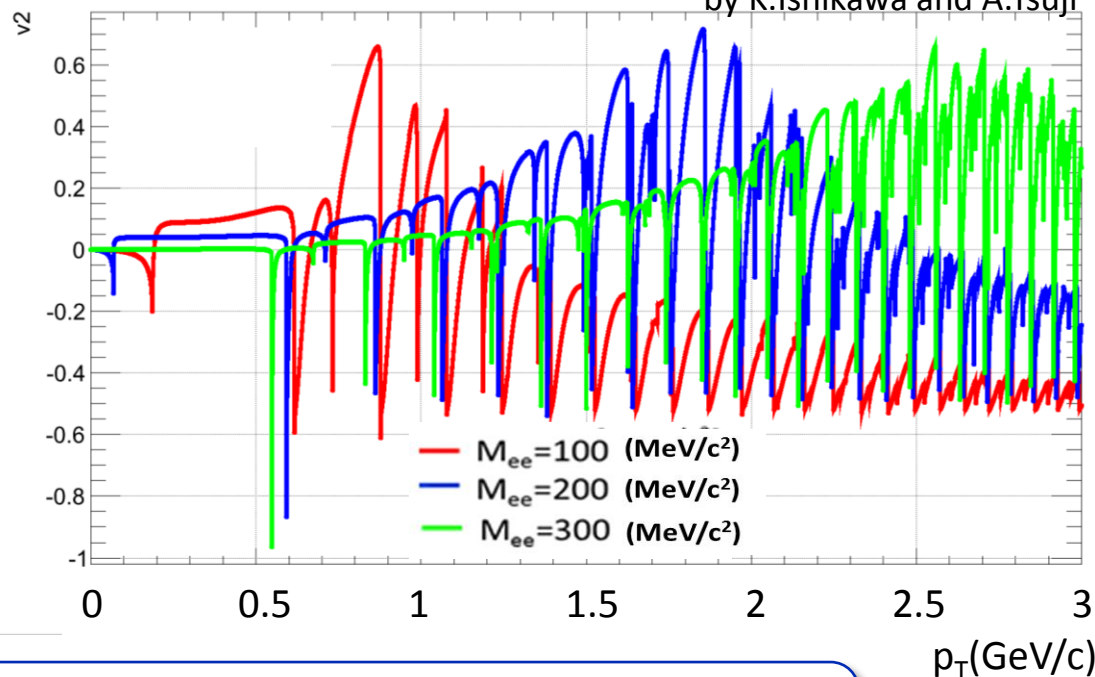
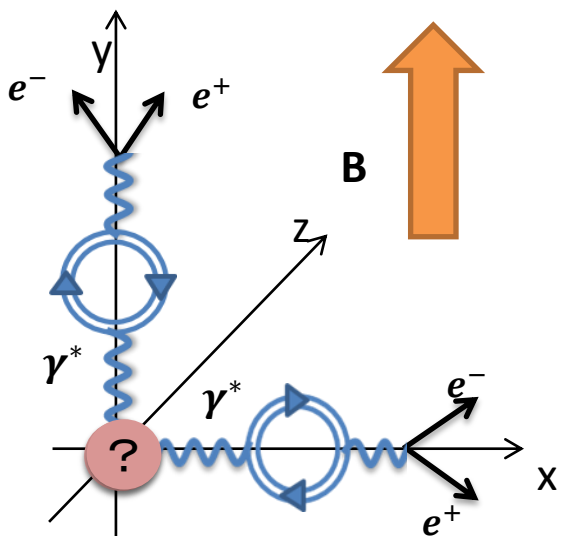
# Anisotropy of virtual photons

$$Anisotropy(v_2) = \frac{R_{\perp} - R_{\parallel}}{R_{\perp} + R_{\parallel}}$$

R: di-electron production rate

Anisotropy of Virtual photon (B=10<sup>15</sup>Tesla)

by K.Ishikawa and A.Tsuji



**The Anisotropy is in the order of 10<sup>-1</sup>!!**

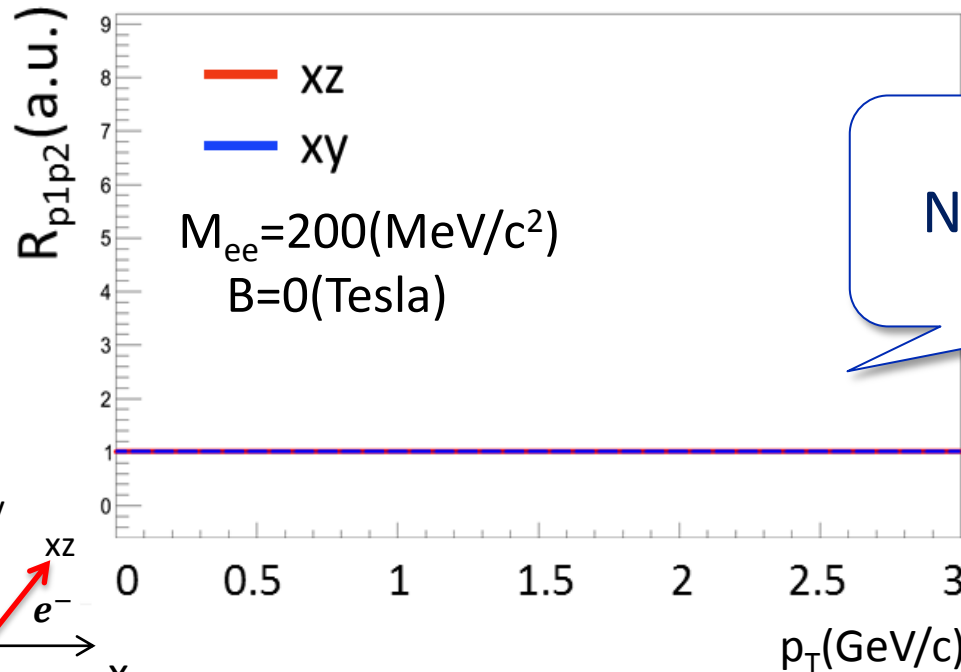
# Polarization of virtual photons

- Di-electron production rate

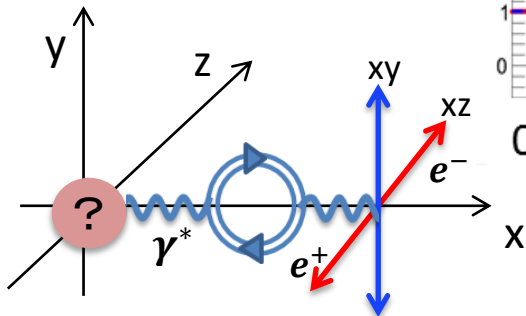
$$R_{p_1 p_2} \equiv \frac{E_1 E_2 dN_{e^+ e^-}}{d^3 p_1 d^3 p_2 d^4 x} = \frac{\alpha^2}{2\pi^4} L^{\mu\nu}(p_1, p_2) D_{\mu\alpha}(q, eB) D_{\nu\beta}(q, eB)^* \frac{\text{Im} G_R^{\alpha\beta}(q, T, eB)}{e^{q^0/T} - 1}$$

- Last factor is approximated as a constant term.
- **When B=0, virtual photon decay is isotropic.**

$p_{1,2}$  : 4D momentum of leg-electron  
 $L$  : kinematics of electron pair



No polarization.



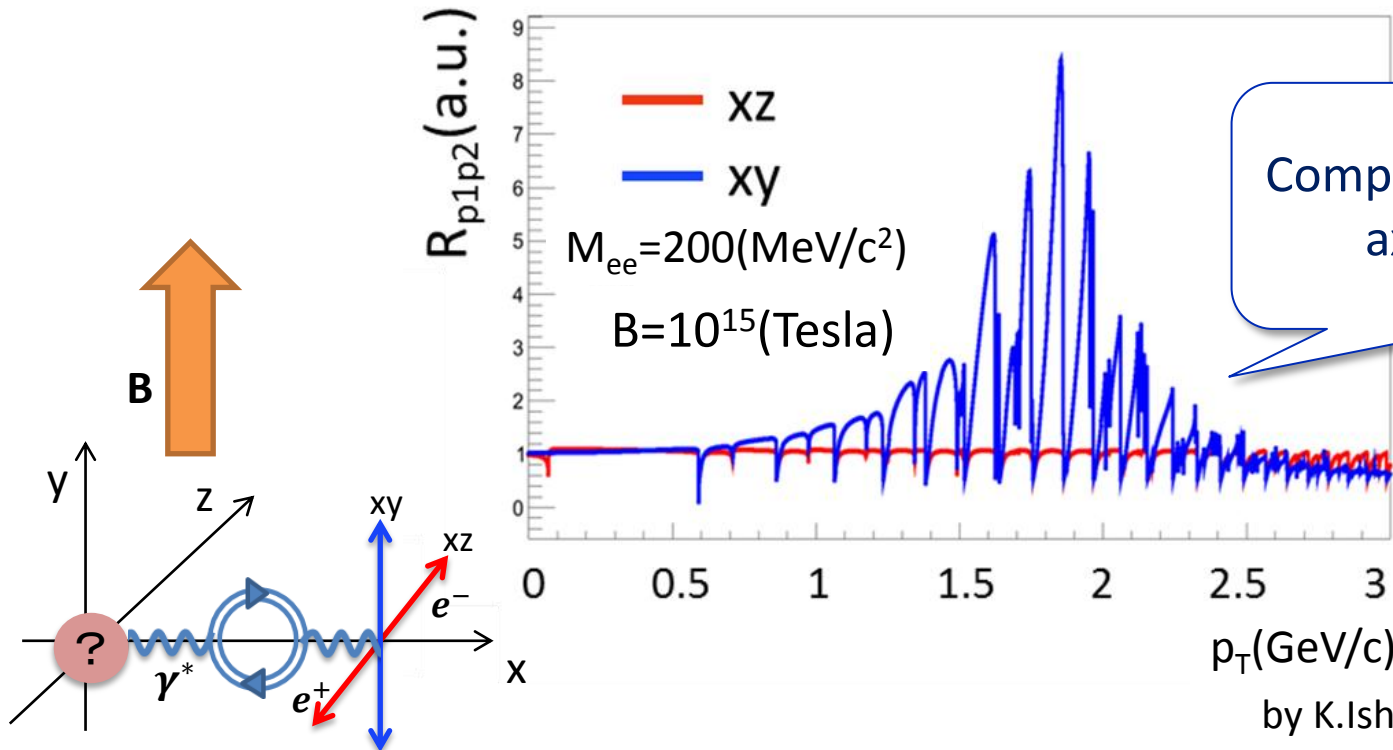
by K.Ishikawa and A.Tsuji

# Polarization of virtual photons

- Di-electron production rate

$$R_{p_1 p_2} \equiv \frac{E_1 E_2 dN_{e^+ e^-}}{d^3 p_1 d^3 p_2 d^4 x} = \frac{\alpha^2}{2\pi^4} L^{\mu\nu}(p_1, p_2) D_{\mu\alpha}(q, eB) D_{\nu\beta}(q, eB)^* \frac{\text{Im} G_R^{\alpha\beta}(q, T, eB)}{e^{q^0/T} - 1}$$

- Last factor is approximated as a constant term.  $p_{1,2}$ : 4D momentum of leg-electron
- For  $B \neq 0$  case, virtual photon decay is anisotropic.  $L$ : kinematics of electron pair



Component along the B axis increases!

by K.Ishikawa and A.Tsuji

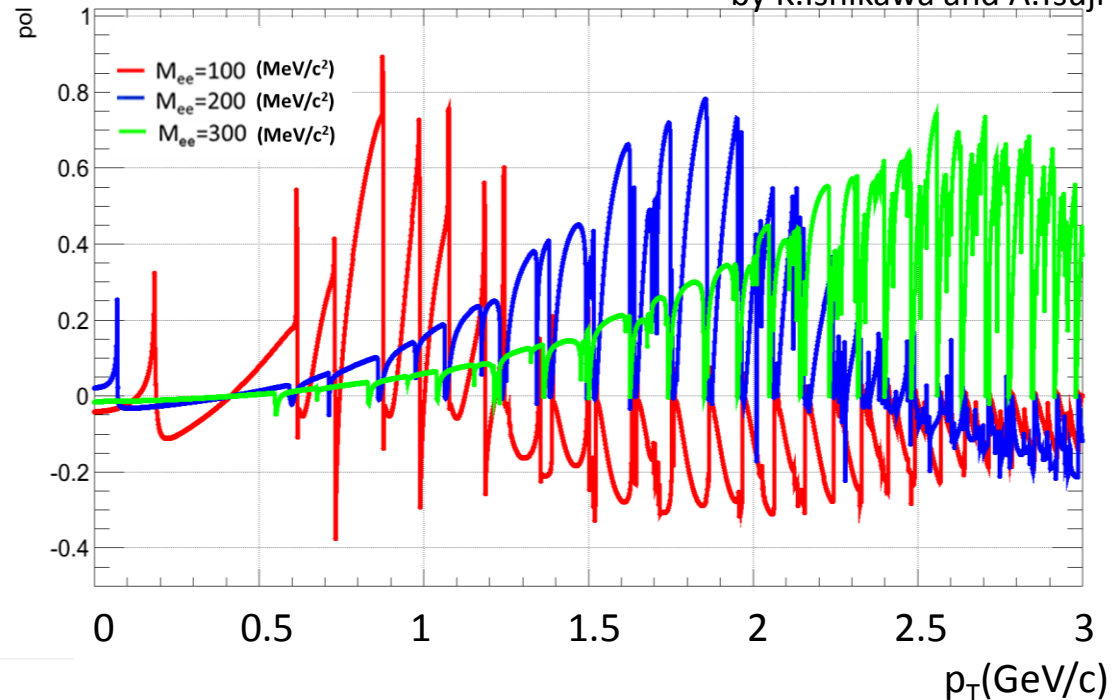
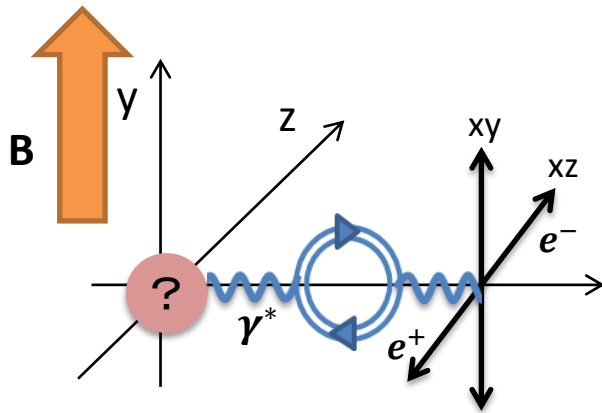
# Polarization of virtual photons

$$\text{polarization} = \frac{R_{xz} - R_{xy}}{R_{xz} + R_{xy}}$$

R: di-electron production rate including leg-electron information

Polarization of virtual photons ( $B=10^{15}$  Tesla)

by K. Ishikawa and A. Tsuji

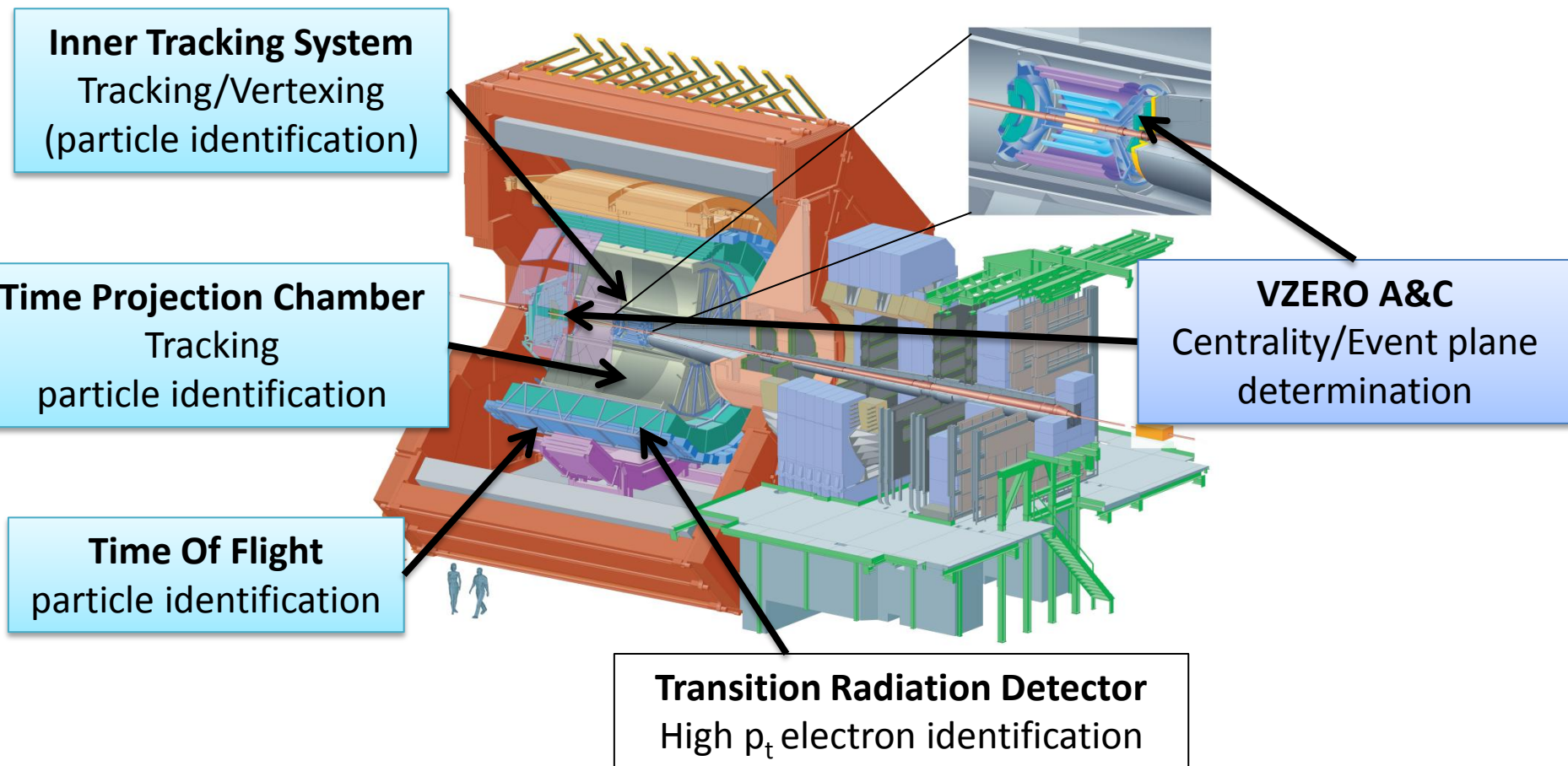


**Virtual photons are polarized in the order of  $10^{-1}$ !!**

# Experimental approach

# The ALICE Detector

## Low mass di-electron analysis with event plane(E.P.)





# Feasibility of field detection @ ALICE

Example: 2011Pb-Pb  $\sqrt{s_{NN}}=2.76\text{TeV}$  0-10% central

(Real analysis will utilize non-central collisions.)

$0.1\text{GeV}/c^2 \leq M_{e\bar{e}} \leq 0.3\text{GeV}/c^2$ ,  $1.0\text{GeV}/c \leq p_t^{\text{pair}} \leq 2.0\text{GeV}/c$

$$\text{statistical significance} = \frac{\text{signal}}{\sqrt{\text{signal} + \text{background}}}$$
$$\approx \frac{N_{e^+e^-}^{\text{all}} \times S/B_{cb} \times (1 - f_{\text{hadronic}}) \times R_{E.P.} \times A}{\sqrt{N_{e^+e^-}^{\text{all}}}}, \quad \frac{N_{e^+e^-}^{\text{all}} / 2 \times S/B_{cb} \times (1 - f_{\text{hadronic}}) \times R_{E.P.} \times P}{\sqrt{N_{e^+e^-}^{\text{all}} / 2}}$$



# Feasibility of field detection @ ALICE

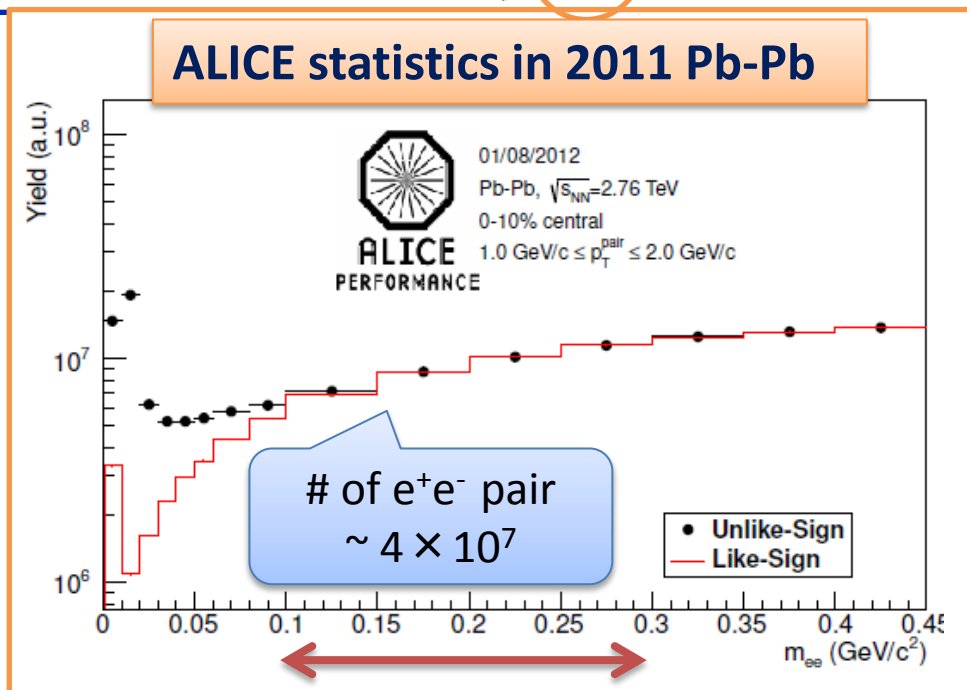
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✓ # of all  $e^+e^-$  pair  $\sim 4 \times 10^7$



# Feasibility of field detection @ ALICE

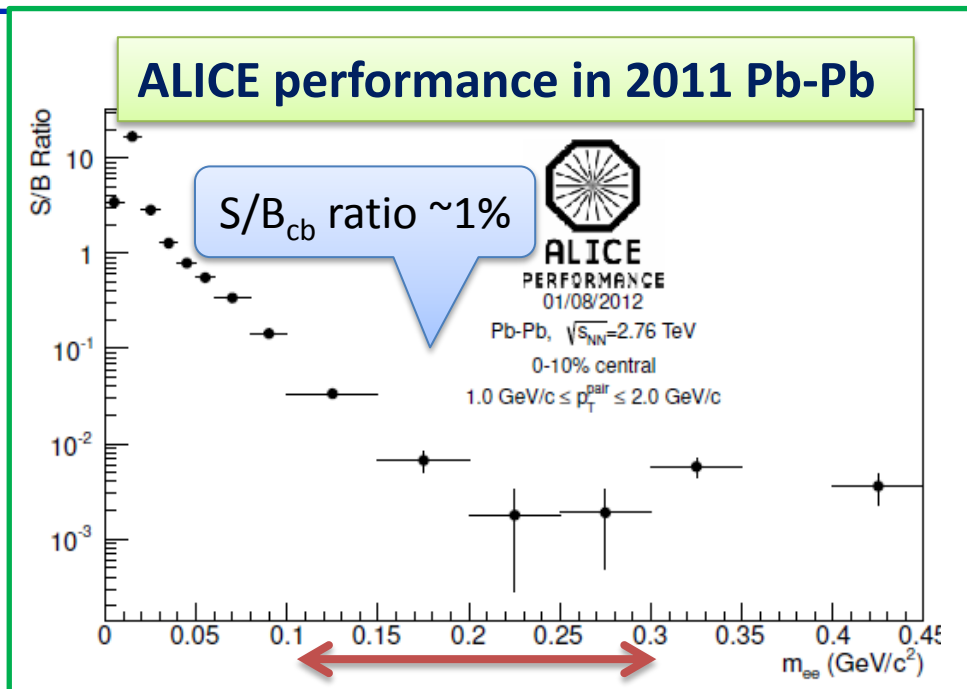
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- ✓ # of all  $e^+e^-$  pair  $\sim 4 \times 10^7$
- ✓  $S/B_{cb}$  ratio  $\sim 1\%$



# Feasibility of field detection @ ALICE

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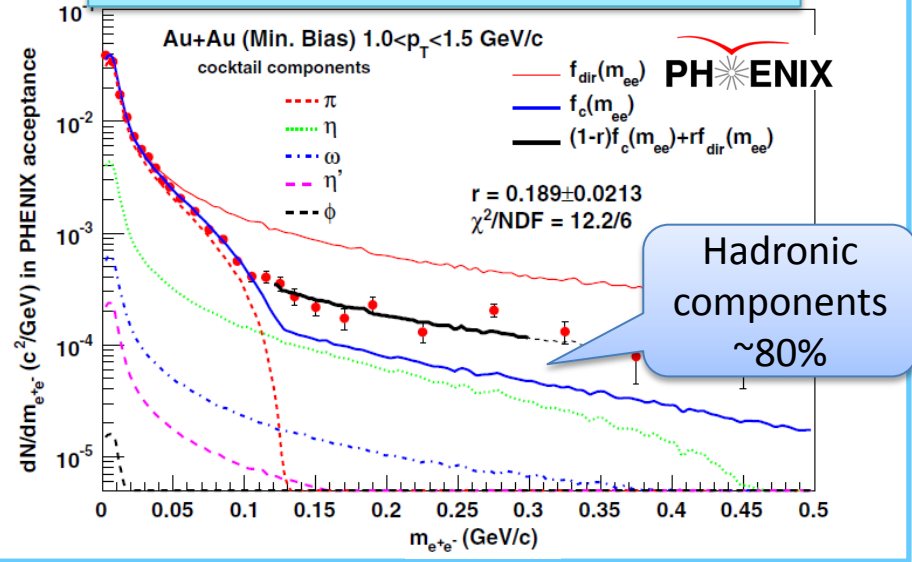
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- ✓ # of  $e^+e^-$  pair  $\sim 4 \times 10^7$
- ✓  $S/B_{cb}$  ratio  $\sim 1\%$
- ✓ Direct components  $\sim 20\%$

borrowed from PHENIX  
 (PRL 104, 132301 (2010))



# Feasibility of field detection @ ALICE

Example: 2011Pb-Pb  $\sqrt{s_{NN}}=2.76\text{TeV}$  0-10% central  
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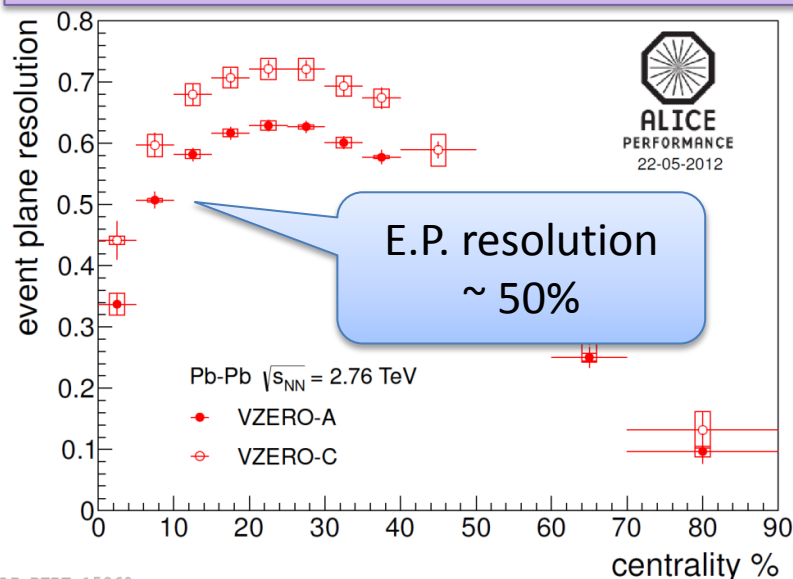
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- ✓ # of all  $e^+e^-$  pair  $\sim 4 \times 10^7$
- ✓  $S/B_{cb}$  ratio  $\sim 1\%$
- ✓ Direct components  $\sim 20\%$
- ✓ E.P. resolution  $\sim 50\%$

## ALICE performance in 2011 Pb-Pb



# Feasibility of field detection @ ALICE

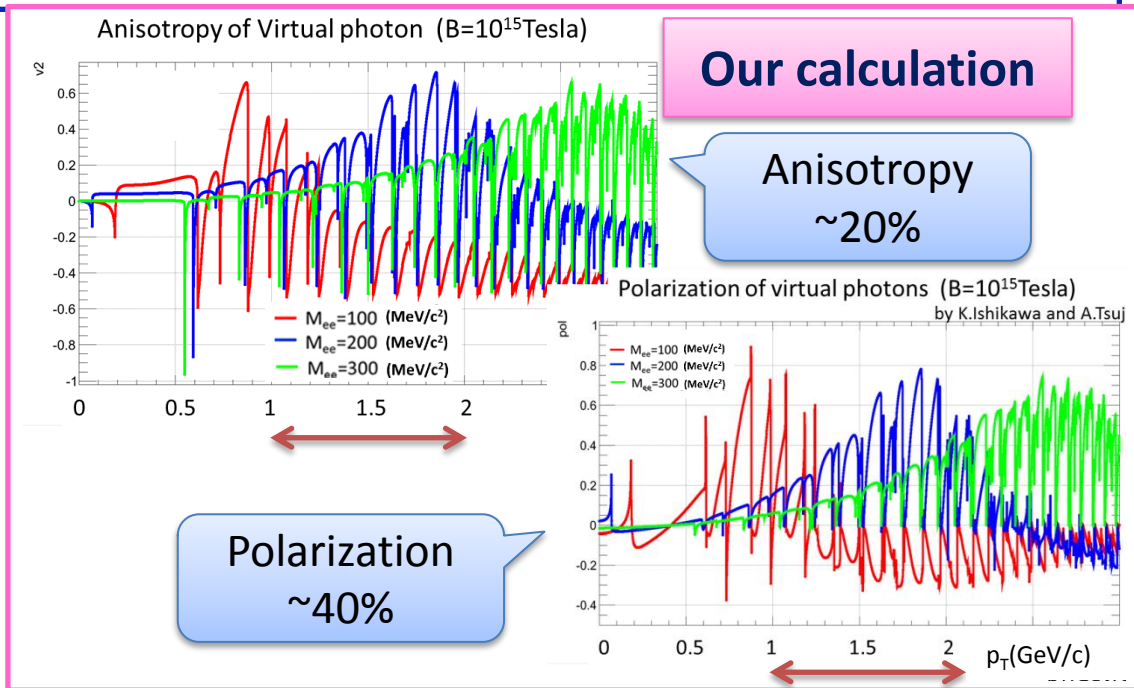
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- ✓ # of all  $e^+e^-$  pair  $\sim 4 \times 10^7$
- ✓  $S/B_{cb}$  ratio  $\sim 1\%$
- ✓ Direct components  $\sim 20\%$
- ✓ E.P. resolution  $\sim 50\%$
- ✓ Anisotropy  $\sim 20\%$
- ✓ Polarization  $\sim 40\%$



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- ✓ Polarization  $\sim 40\%$

Expected statistical significances

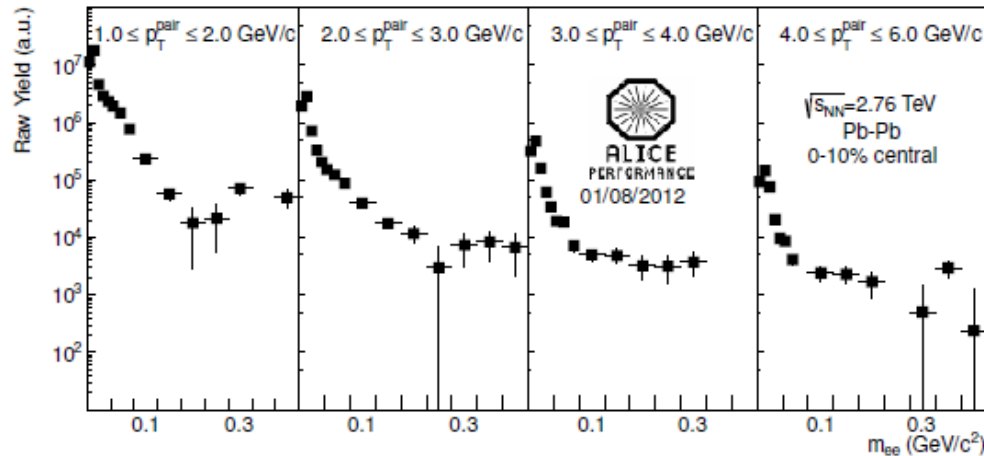
Anisotropy  $\sim 1\sigma$

Polarization  $\sim 2\sigma$

having possibility that we detect the field !!!

# Ongoing efforts of real data analysis

- ALICE low mass  $e^+e^-$  analysis framework already exists.



ALI-PERF-35221

Background subtracted di-electron signal

- **Under development**
  1. combine with event plane information
    - to see the magnetic field axis
  2. add leg-electron three dimensional momentum information
    - to see direction of decay angle with respect to E.P. (B field)



# Summary & Conclusion

- **Ultra-strong magnetic field in HIC**
  - $B \sim 10^{15}$  Tesla @LHC,  $B \sim 10^{14}$  Tesla @RHIC
  - The field is above  $B_c$  even after a few fm/c
  - Various interesting phenomena are under discussion
- **Studying to detect the field via virtual photon measurements at ALICE**
  - Estimation of effects of the field
    - Anisotropy of virtual photon :  $o(10^{-1})$
    - Polarization of virtual photon :  $o(10^{-1})$
  - Feasibility of detection of the field
    - Anisotropy of virtual photon :  $o(1\sigma)$
    - Polarization of virtual photon :  $o(1\sigma)$

**I'm analyzing real data and**

**shall show you results in the near future!!**

**Back up !!**

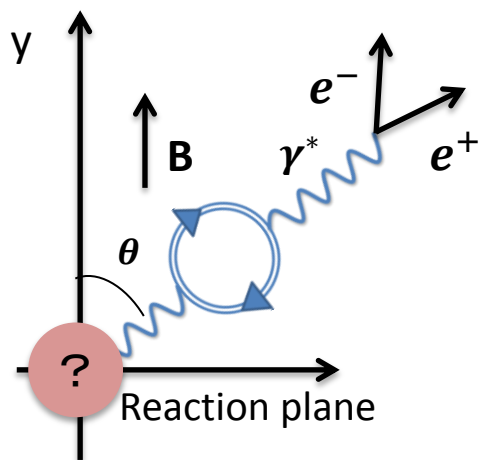
# Di-electron production rate

$$R_q \equiv \frac{R_{e^+e^-}}{d^4q} = -\frac{\alpha^2}{3\pi^3} q^2 g^{\mu\alpha} D_{\mu\nu}(q, eB) D_{\alpha\beta}(q, eB)^* \frac{\text{Im} G_R^{\nu\beta}(q, T, eB)}{e^{q^0/T} - 1}$$

$q$  : 4D momentum of electron pair

$$R_{p_1 p_2} \equiv \frac{E_1 E_2 dN_{e^+e^-}}{d^3p_1 d^3p_2 d^4x} = \frac{\alpha^2}{2\pi^4} L^{\mu\nu}(p_1, p_2) D_{\mu\alpha}(q, eB) D_{\nu\beta}(q, eB)^* \frac{\text{Im} G_R^{\alpha\beta}(q, T, eB)}{e^{q^0/T} - 1}$$

$p_{1,2}$  : 4D momentum of leg-electron



$$\frac{\text{Im} G_R^{\nu\beta}(q, T, eB)}{e^{q^0/T} - 1} = (-g^{\alpha\beta} q^2 + q^\alpha q^\beta) C$$

$$D_{\mu\nu}(q) = \frac{-i}{q^2} \left( g^{\mu\nu} - \frac{1}{q^2} \Pi^{\mu\nu}(q) \right)^{-1}$$

$$L^{\mu\nu} = p_1^\mu p_2^\nu + p_1^\nu p_2^\mu - (p_1 \cdot p_2 + m_e^2) g^{\mu\nu}.$$

- $G$  is the retarded Green function of electric current in the thermal and magnetic environment and source of the virtual photons.
- It is still difficult to argue for the detail of the virtual photon source in the heavy ion collision and here we do not consider them, for which we employ a simple constant form with current conservation projection;

$$\frac{\text{Im} G_R^{v\beta}(q, T, eB)}{e^{q^0/T} - 1} = (-g^{\alpha\beta} q^2 + q^\alpha q^\beta) C$$

# Photon vacuum polarization tensor in strong B

$$\Pi^{\mu\nu}(q, B) = \left[ \left( g^{\mu\nu} q^2 - q^\mu q^\nu \right) - \left( g_{\parallel}^{\mu\nu} q_{\parallel}^2 - q_{\parallel}^\mu q_{\parallel}^\nu \right) - \left( g_{\perp}^{\mu\nu} q_{\perp}^2 - q_{\perp}^\mu q_{\perp}^\nu \right) \right] \bar{N}_0 \\ + \left( g_{\parallel}^{\mu\nu} q_{\parallel}^2 - q_{\parallel}^\mu q_{\parallel}^\nu \right) \bar{N}_1 + \left( g_{\perp}^{\mu\nu} q_{\perp}^2 - q_{\perp}^\mu q_{\perp}^\nu \right) \bar{N}_2$$

$$\bar{N}_j = -\frac{\alpha}{4\pi} \int_{-1}^1 dv \int_{0-i\varepsilon}^{\infty-i\varepsilon} \frac{dz}{z} \left[ e^{-i\Phi} \frac{z}{\sin z} F_j(v, z) - [\text{c.t.}] \right]$$

$$F_0(v, z) = \cos(vz) - v \frac{\cos(z) \sin(vz)}{\sin z}, F_1(v, z) = (1-v^2) \cos z, F_2(v, z) = 2 \frac{\cos(vz) - \cos z}{\sin^2 z}$$

$$[\text{c.t.}] = (1-v^2) e^{-\frac{z}{\mu}}, \Phi \equiv \frac{1-(1-v^2)r}{\mu} z - 2 \frac{\cos(vz) - \cos z}{\sin z} \frac{q}{\mu},$$

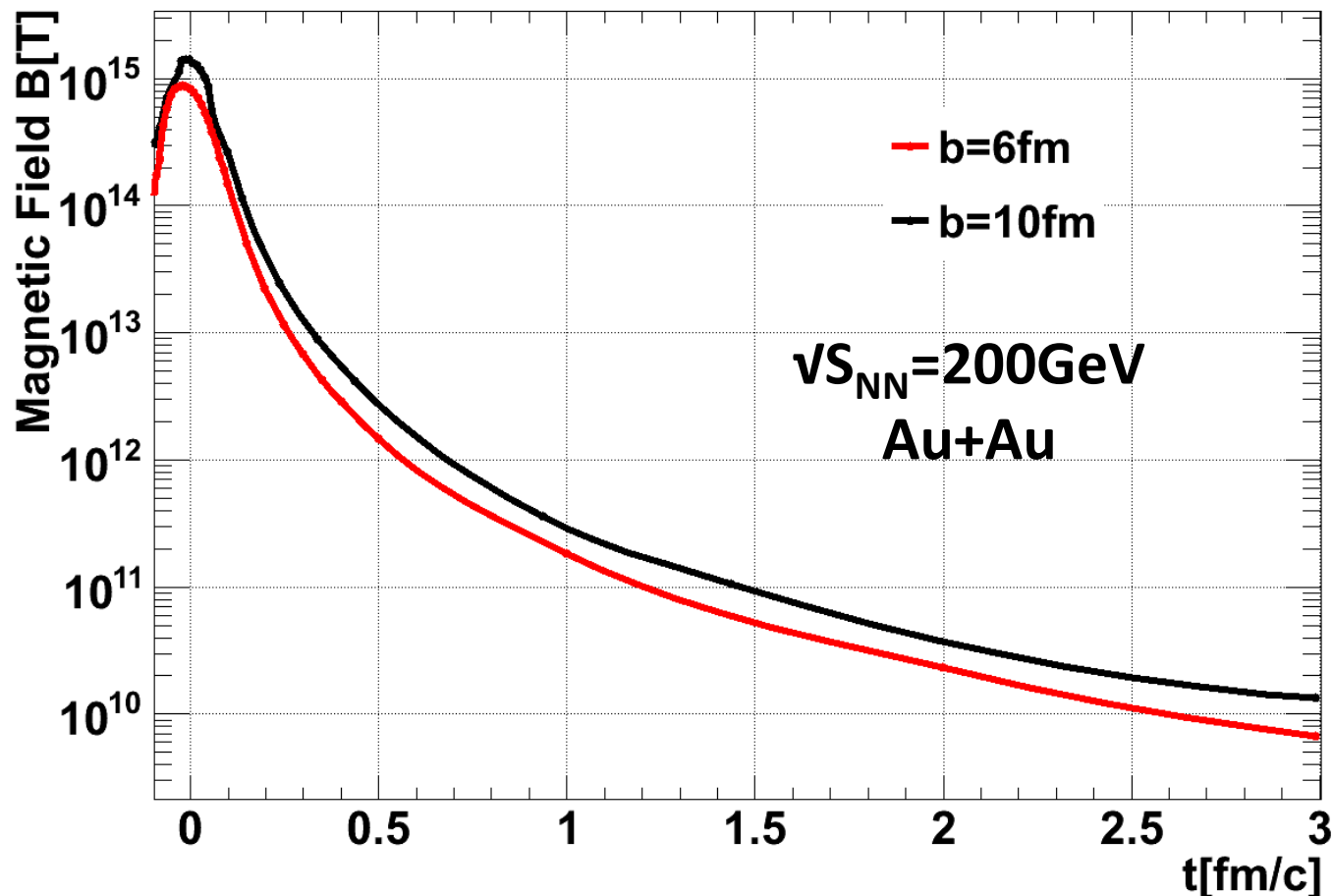
$$\mu \equiv \frac{m^2}{eB}, \quad r \equiv \frac{q_{\parallel}^2}{4m^2}, \quad q \equiv \frac{q_{\perp}^2}{4m^2}, \quad q_{\parallel}^2 = (q_t)^2 - (q_z)^2, \quad q_{\perp}^2 = -(q_x)^2 - (q_y)^2$$

Tsai, Phys. Rev. D10(1974); Baier, Katkov, Strankhovenko, Sov. Phys. JETP41(1975);  
Melrose, Stoneham, J. Phys. A. Math. Gen. 10(1977).....



# Time dependence

- Impact parameter = 6,10fm, spectators only



By using JAM model