

# Freeze-In of FIMP Dark Matter

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# Outline of Talk

## I. Freeze-out of Weakly Interacting Massive Particles

## II. Freeze-In of Feebly Interacting Massive Particles

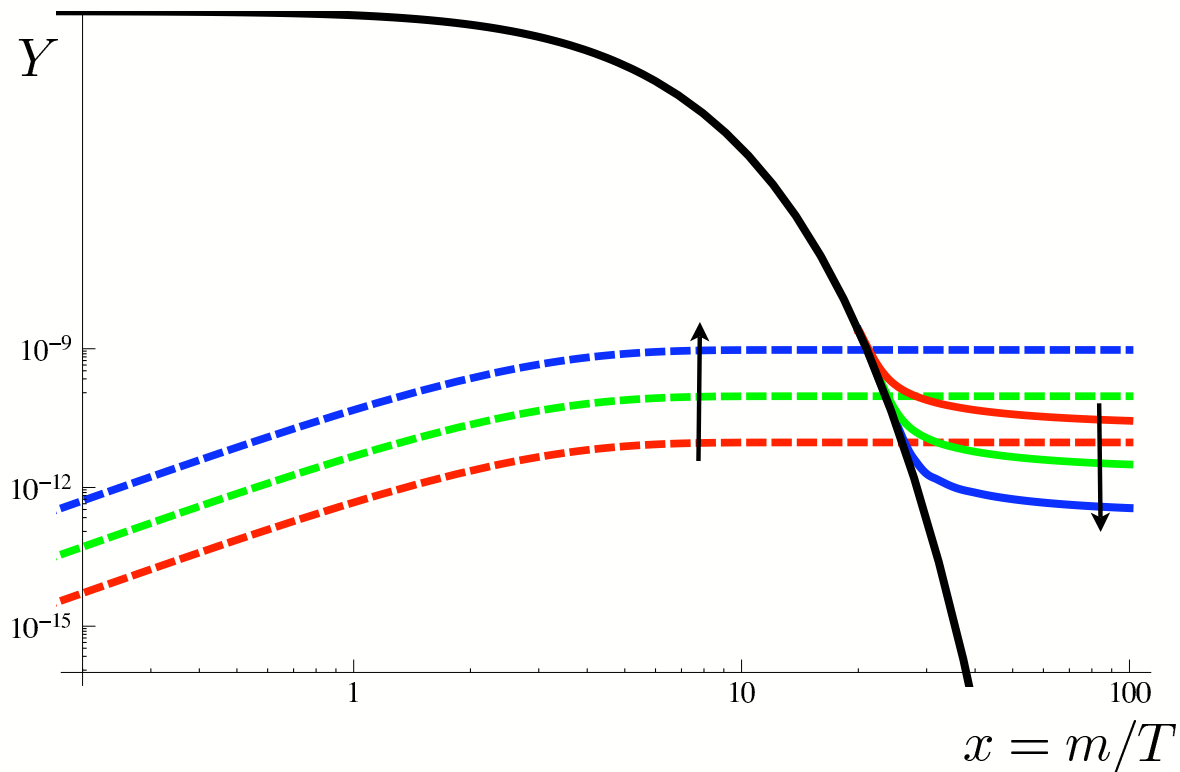
Hall, K.J., March-Russell, West

- The Freeze-In Process
- Comparison to super-WIMPs
- A Unified View of Freeze-In and Freeze-Out
- Detectability
- Candidate Particles

## III. Conclusions

# Freeze-Out of Dark Matter

- need some dark matter particle  $X$  stabilizing symmetry (parity)
- annihilation reactions at  $X + \bar{X} \rightarrow \text{standard model particles}$  freeze out at some  $T \lesssim m_X$  and  $n_X \ll T^3$



# Virtues of Freeze-Out Production of Dark Matter

minimalistic assumptions as well as accelerator testability

- thermodynamic and chemical equilibrium shortly before freeze-out

seemingly reasonable assumption since typically  $t_{equ}/t_{Hubble} \ll 1$

- $\Omega h^2 \approx 0.1 \left( \frac{\sigma v}{3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}} \right)^{-1}$  - required interactions in principle accelerator testable

- WIMP-"miracle": the required interaction strength is reached roughly at the electroweak scale where new physics is expected

reminiscent to conditions which led to the standard Big Bang nucleosynthesis model

# Question:

*Is freeze-out of dark matter the ONLY accelerator testable dark matter production mechanism in thermodynamic equilibrium conditions ?*

**No !**

# FIMP Dark Matter

production per Hubble time

imagine a particle  $X$  which is so feebly interacting with the plasma (in TE) that it will never reach equilibrium abundance

call it FIMP  $\equiv$

*"Feebly Interacting Massive Particle"*

take interaction  $\mathcal{L} \sim \lambda X B_1 B_2$  with  $\lambda \ll 1$   
where  $B_1$  and  $B_2$  are bath particles

the plasma produces it in attempting to attain equilibrium via  $B_1 \rightarrow B_2 + X$  decay production

$$\begin{aligned}\Delta n_X / s &\sim \frac{n_{B_1} \Gamma_{B_1 \rightarrow B_2 + X} t_H}{s} \\ &\sim \frac{g_{B_1} T^3 \lambda^2 m_{B_1} M_{pl} / T^2}{g T^3} \\ &\sim \frac{g_{B_1} \lambda^2 m_{B_1} M_{pl}}{g T^2}\end{aligned}$$

prod. infrared dominated !!!

$$\rightarrow \Omega_X \sim \frac{g_{B_1}}{g} \lambda^2 M_{pl} \frac{m_X}{m_{B_1}}$$

# Difference to super-WIMPs

- super-WIMPs as **gravitinos** or **axinos** are also very weakly interacting
- $\Delta n_G/s \sim n^2 \sigma v t_H/s \sim g^2 M_{pl} T \sigma v$  with  $\sigma \sim 1/M_{pl}^2$  for weak mass scale gravitino, for example
- → their production is ultraviolet dominated and reheat temperature  $T$  dependent

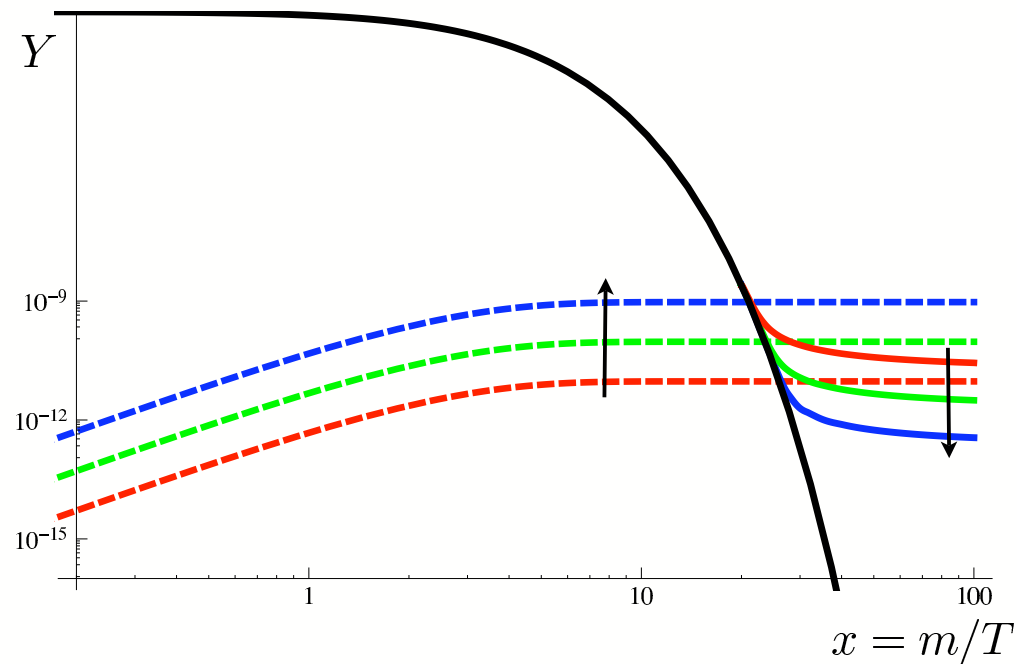
*reheat temperature essentially non-testable in accelerators –*

*requires detailed information of the inflaton sector*

*difference between super-WIMPs and FIMPs is renormalizability of interaction*

# Freeze-In of Dark Matter

- production reactions  $B_1 \rightarrow X + B_2$  become inefficient at  $T \lesssim m_{B_1}$  freezing-in (thawing-in) the dark matter abundance at  $n_X \ll T^3$



- production goes up with interaction strength



## Required Interaction Strength

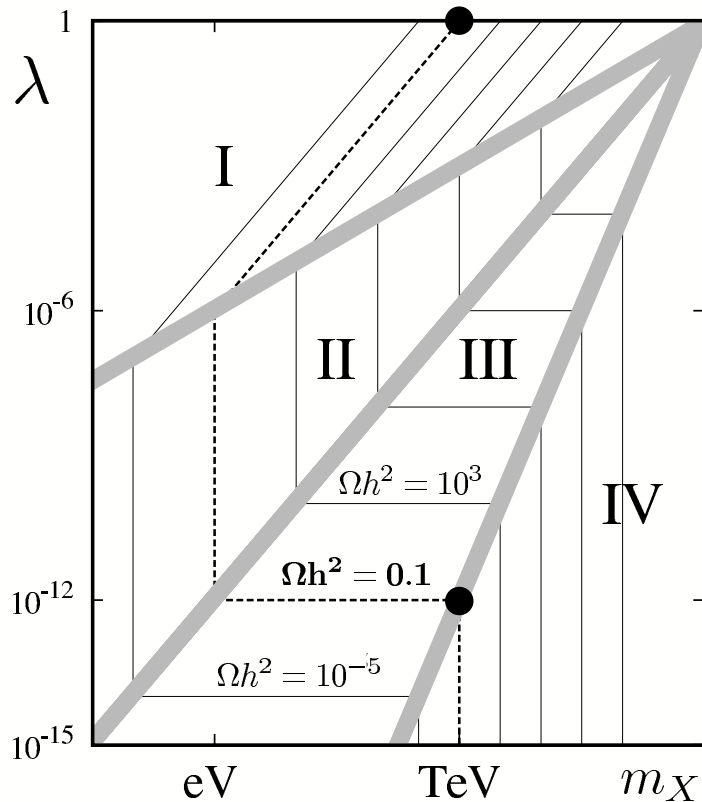
$$\lambda \simeq 1.5 \times 10^{-12} \left( \frac{m_X}{m_{B_1}} \right)^{1/2} \left( \frac{g_*(m_X)}{10^2} \right)^{3/4} \left( \frac{1}{g_{bath}} \right)^{1/2}$$

this is close to  $M_{EW}/M_{GUT} \sim 10^{-13}$

$g_{bath} \gg 1$  possible

# A Unified View of Freeze-In and Freeze-Out

$$\mathcal{L} \sim \lambda X B_1 B_2 \text{ and } M_x \lesssim M_{B_1}$$

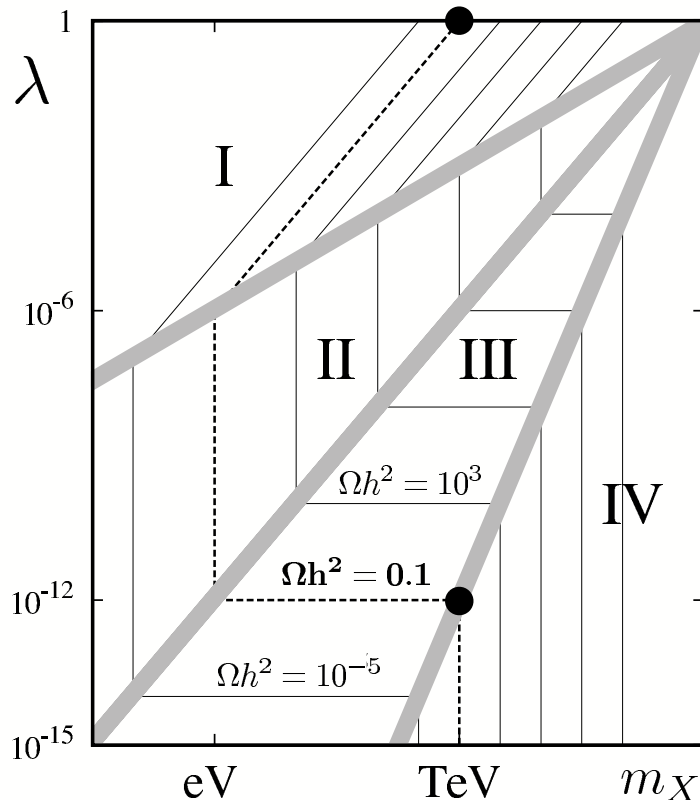


**Region I:** Coupling  $\lambda$  of  $X$  to thermal bath strong enough such that equilibrium  $\sim T^3$  density will be attained and at  $T < m_X$   $n_X \ll T^3$  will be frozen out  $\rightarrow$  **non-relativistic freeze-out**

**Region II:** Coupling  $\lambda$  of  $X$  to thermal bath strong enough such that equilibrium  $\sim T^3$  density will be attained – however when  $T < m_X$  no further reduction  $\rightarrow$  **relativistic freeze-out**

# A Unified View of Freeze-In and Freeze-Out

$$\mathcal{L} \sim \lambda X B_1 B_2 \text{ and } M_x \lesssim M_{B_1}$$



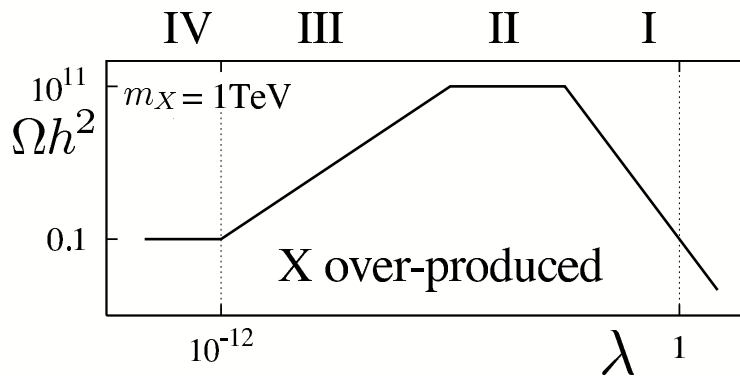
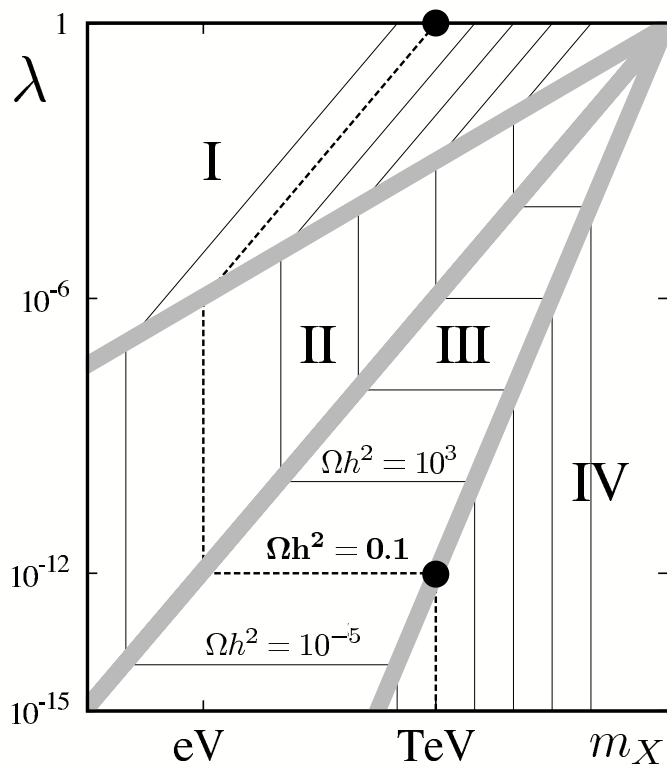
**Region III:** Coupling to thermal bath **NOT** strong enough to attain equilibrium density  $\sim T^3$  – freeze-in – abundance of  $X$  dominated by freeze-in

**Region IV:** Coupling to thermal bath **NOT** strong enough to attain equilibrium density  $\sim T^3$  – freeze-in – abundance of  $X$  dominated by freeze-out of bath particles  $B$  and subsequent decay to  $X$

freeze-in completes the lower half of the diagram

# A Unified View of Freeze-In and Freeze-Out

$$\mathcal{L} \sim \lambda X B_1 B_2 \text{ and } M_x \lesssim M_{B_1}$$



freeze-in completes the lower half of the diagram

## Detectability of FIMPs ?

Production via  $B_1 \rightarrow B_2 + X$

$$\Omega_X h^2 \approx \frac{1.09 \times 10^{27} g_{B_1} m_X \Gamma_{B_1}}{g_*^S \sqrt{g_*^p} m_{B_1}^2}$$

$$\tau_{B_1} = 7.7 \times 10^{-3} \text{sec}$$

$$g_{B_1} \left( \frac{m_X}{100 \text{ GeV}} \right) \left( \frac{300 \text{ GeV}}{m_{B_1}} \right)^2 \left( \frac{10^2}{g_*(m_{B_1})} \right)^{3/2} \left( \frac{\Omega_X h^2}{0.011} \right)^{-1}$$

*direct test of production mechanism in lab*  
**!!!!**

# Production of Dark Matter via Freeze-In of FIMPs

so far, have assumed FIMP is the dark matter particle

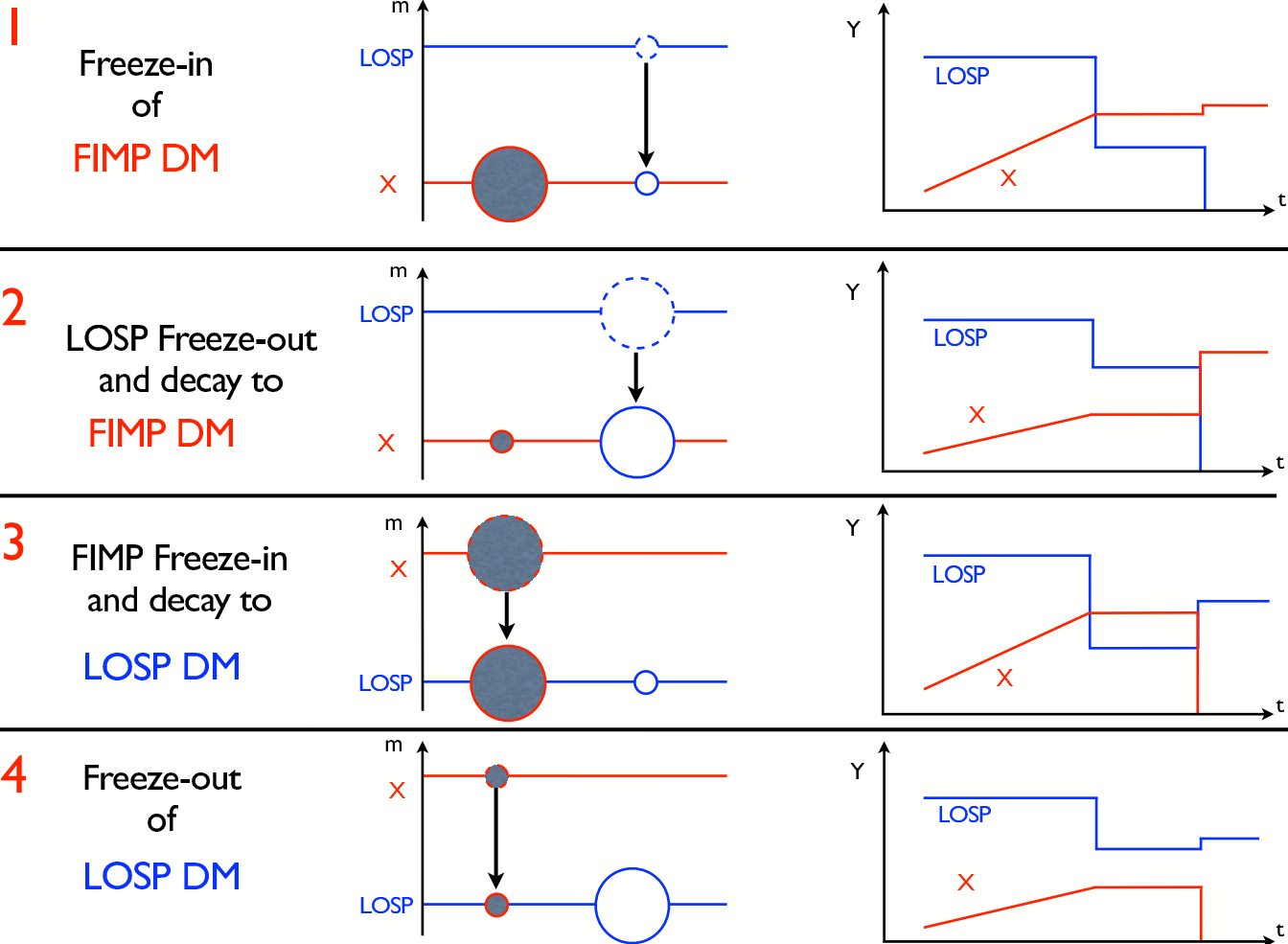
- need some (at least approximate) symmetry which stabilizes the dark matter particle, call it parity
- the standard model particles have positive parity
- the dark matter particle and other yet undiscovered particles have negative parity, stabilizing them towards decay into standard model particles

**LOSP**  $\equiv$  "Lightest Observable Sector Particle" which carries negative parity

$m_{\text{LOSP}} < m_{\text{FIMP}}$  is possible  $\rightarrow$  the **LOSP** may be the dark matter particle

- FIMPs are produced by inverse decays, e.g.  $B + \text{LOSP} \rightarrow \text{FIMP}$ , which decay into LOSPs after LOSP freeze-out
- the LOSP self-annihilation cross section can be large

# Four possibilities



# Candidate Particles

- Moduli determining soft SUSY breaking parameters

$$\begin{array}{lll} m^2 \left(1 + \frac{T}{M}\right) (\phi^\dagger \phi + h^\dagger h) & \mu B \left(1 + \frac{T}{M}\right) h^2 & Ay \left(1 + \frac{T}{M}\right) \phi^2 h \\ m_{\tilde{g}} \left(1 + \frac{T}{M}\right) \tilde{g}\tilde{g} & \mu y \left(1 + \frac{T}{M}\right) \phi^2 h^* & \mu \left(1 + \frac{T}{M}\right) \tilde{h}\tilde{h}, \end{array}$$

- Dirac Neutrinos within weak scale supersymmetry

$$\lambda LN H_u,$$

- $\lambda \sim 10^{-13}$  for observed neutrino masses !! Right-handed sneutrino close to perfect candidate for FIMP (cf. Asaka *et al.* 06,07)



# LOSP/FIMP Decays during BBN ?

- two-body decay:  
 $\tau \sim 10^{-2} \text{ sec } (\Omega_X h^2 / 0.1)^{-1} g_{B_1}$
- for  $\Omega_X h^2 \sim 0.1$  and  $g_{B_1} \sim 1$   
 $\rightarrow$  no effect
- three-body decay:  
 $\tau \sim 3 \text{ sec } g^{-2} (\Omega_X h^2 / 0.1)^{-1} g_{B_1}$
- possible effect, especially when  
 $\Omega_X h^2 < 0.1$  and/or  $g_{B_1} \gg 1$
- three-body decay, for example, when  
 LOSP *not* directly coupled to FIMP

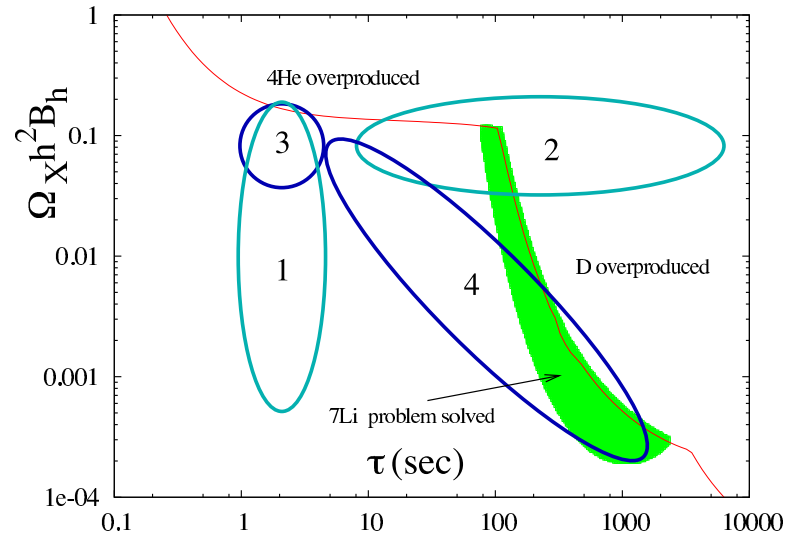


figure assumes LOSP/FIMP three-body decays but production of FIMPs via two-body (inverse) decays

# LHC Experiments to find metastable particles

Both, Atlas and CMS, are searching for **metastable particle decays** .... as of January 31, CMS has a *hint* of a  $\sim 700\text{GeV}$  "heavy stable charged particle" ...

# How to convince oneself that FIMPs constitute the dark matter ?

- the LOSP is charged and/or strongly interacting, *NOT* a neutralino
- it is metastable
- its life time falls in the right ballpark to fulfill the  $\tau_{\text{LOSP}} \gtrsim 10^{-2} \text{sec } m_X / m_{\text{LOSP}}$  relationship

FIMPs as dark matter is a very plausible scenario

how to really convince oneself

- one may determine  $m_{\text{LOSP}}$  and  $m_X \sim m_{\text{LOSP}}$  from kinematics
- the  $\tau_{\text{LOSP}} - \Omega_X$  relationship is consistent with/close to the WMAP value

# Summary

- dark matter production via freeze-out may occur in (plausible) thermodynamic equilibrium conditions, is UV insensitive, and accelerator testable !
- when looking at other dark matter production mechanism with such attributes one is led to the process of freeze-in
- in fact, freeze-in and freeze-out may be unified in a dark matter *interaction strength - mass* diagram
- candidate particles for *Feebly Interacting Massive Particles* as required in freeze-in do exist, in fact, the required interaction strength  $\lambda \lesssim 10^{-12}$  is suggestive
- freeze-in production may lead to a simple testable correlation between the life time of a new fundamental metastable particle and the abundance of the dark matter

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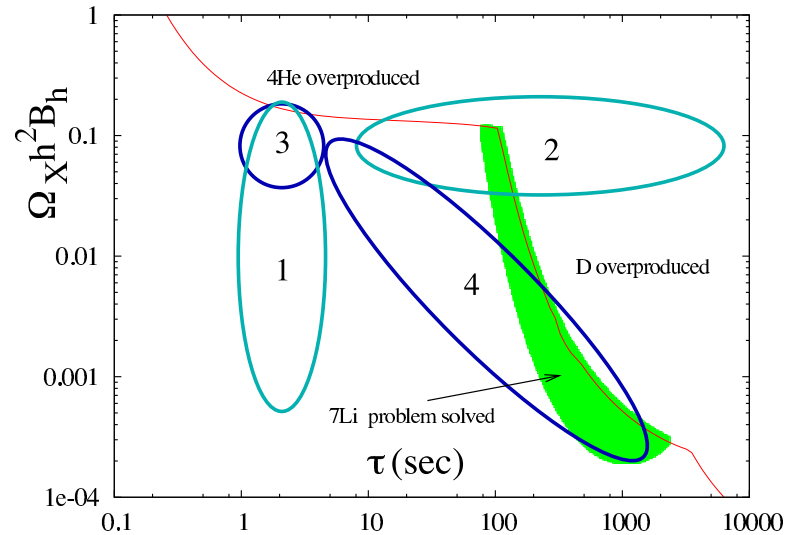


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