Polarization and nonlinear effects in the beam-beam interaction

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- Review beam-beam incoherent processes and the treatment of polarization
- develop expressions for, virtual photon polarization, Breit-Wheeler cross-section with full polarization, polarization of final states
- simulation runs to find pair numbers, E, Pt and luminosity
- Coherent processes resonant cross-sections in the second order processes

E⁺E⁻ beam-beam source of background pairs



Incoherent pair processes • $\gamma_{virtual} + \gamma_{real} \rightarrow E^{+} + E^{-} (\sim 71\%)$ • $\gamma_{virtual} + \gamma_{virtual} \rightarrow E^{+} + E^{-} (\sim 23\%)$ • $\gamma_{real} + \gamma_{real} \rightarrow E^{+} + E^{-} (\sim 6\%)$

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UNKNOWNS

- Polarised beams
- Virtual photon polarisation
- Virtual photon flux (inaccurate at low energies)
- Second order non-linear processes (contains resonant cross-sections)
- **IS THERE VARIATION IN:** pair #'s, P_t, luminosity?

Beam-beam simulators: CAIN mod polarizations

e+ 2 real γ 's - Breit-Wheeler 1 real and 1 virtual γ – Bethe Heitler 2 virtual γ 's – Landau Lifshitz same cross-section: $k_{b1} + k_{b2} \rightarrow e^- + e^+$

1 virtual \mathcal{Y} and 1 electron - Bremstrahlung

	Initial Polarization			Final Polarization		
	$\xi(\mathbf{X})$	ξ(y)	ξ(z)	$\boldsymbol{\zeta}(x)$	$\zeta(y)$	$\zeta(z)$
Breit_Wheeler	 Image: A second s		\checkmark	\checkmark	\checkmark	\checkmark
Bethe-Heitler	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark
Landau-Lifshitz	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark
Bremstrahlung	?	?	?	?	?	?

● Already existing ✓ added to CAIN ? yet to implement

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Existing CAIN structure: Beamstrahlung photon polarization

Energy spectra given by Sokolov-Ternov equation

$$dW = -i \frac{\alpha m}{\sqrt{3}\pi \gamma} \left[\int_{z}^{\infty} K_{5/3}(z) dz + \frac{x^{2}}{1-x} K_{2/3}(z) \right] dx \quad where \quad z = \frac{2B_{sch}}{3B\gamma} \frac{x}{1-x}$$

polarization first calculated with individual basis vectors then rotated to the same basis to be used as input to pair processes e_{7}



A version of the Sokolov-Ternov equation with the polarizations written explicitly gives beamstrahlung polarizations





Virtual photon polarization I

Spectral component of bunch electric field as a function of transverse position

$$E_{\omega}^{x,y} = -\frac{ie}{\pi v} \iint \frac{q_{x,y}}{q_{x}^{2} + q_{y}^{2}} F(q) e^{i x q_{x}} e^{i y q_{y}} dq_{x} dq_{y}$$

where the form factor is $F(q) = N \exp \left[-\frac{1}{2}(q_x \sigma_x)^2 - \frac{1}{2}(q_y \sigma_y)^2\right]$

and the polarization vector of virtual photons

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see Engel, Schiller & Serbo Z Phys C 71, 665 (1996)

and similar for

 E^{y}

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integration is performed by expanding in a taylor series and using the limit $\sigma_X >> \sigma_V$ for flat beams

$$E_{\omega}^{x} = -i \frac{x}{\sigma_{x}^{3}} \exp\left(-\frac{x^{2}}{2}\sigma_{x}^{2}\right) \left[\sigma_{y} \exp\left(-\frac{y^{2}}{2}\sigma_{y}^{2}\right) + \sqrt{\frac{\pi}{2}} y \operatorname{Erf}\left(\frac{y}{\sqrt{2}}\sigma_{y}\right)\right]$$

Virtual photon polarization II

X and Y spectral components of bunch electric field



Virtual photon polarization



Magnitude of y component of spectral electric field is much greater than x component. Has consequences for stokes parameters since

$$\xi_{1} = \hat{E}_{\omega}^{x} \hat{E}_{\omega}^{y} * + \hat{E}_{\omega}^{y} \hat{E}_{\omega}^{x} * ; \quad \xi_{3} = \hat{E}_{\omega}^{x} \hat{E}_{\omega}^{x} * - \hat{E}_{\omega}^{y} \hat{E}_{\omega}^{y} * \\ \xi_{2} = \Im [\hat{E}_{\omega}^{y} \hat{E}_{\omega}^{x} * - \hat{E}_{\omega}^{x} \hat{E}_{\omega}^{y} *] = 0$$

No circular polarization **BUT** processes occur in bunches undergoing pinch effect and other disruption, could use a more realistic form factor for the bunch field

Modified cross-section

• Breit-Wheeler cross-section, CAIN original:

$$\sigma_{orig} \propto 2 \left(1 - h + \frac{2\epsilon^2 - 1}{2\epsilon^4} \right) \sinh^{-1} p + \frac{p}{\epsilon} \left(3h - 1 - \frac{1}{\epsilon^2} \right) \qquad \text{where} \qquad p = \text{electron momentum} \\ \epsilon = \text{electron energy} \\ h = \xi_2 \xi'_2$$

full treatment due to Baier & Grozin hep-ph/0209361

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$$\frac{d\sigma}{d\cos(\theta)d\phi} = \frac{\alpha^2}{4s^2 x^2 y^2} \sum_{ii'jj'} F_{jj'}^{ii'} \xi_j \xi_{j'} \zeta_i \zeta_i'$$

F are functions of scalar products of 4-momenta

$$\sigma_{new} \propto 2\left(1-h+\frac{2}{\epsilon^2}(ha+\xi_1\xi_1')-\frac{ha}{\epsilon^4}\right)\sinh^{-1}p + \frac{p}{\epsilon}\left(3h-1-\xi_1\xi_1'-\xi_3\xi_3'-\frac{ha}{\epsilon^2}\right)$$

where $ha = 1 + \xi_3 + \xi'_3 + \xi_3 \xi'_3$ Full expression has similar structure to original CAIN form, so can utilise existing monte-carlo methods <u>CLIC WS CERN</u> 17.10.2007

Cross-section plots



Both Breit-Wheeler cross-sections peak near the electron rest mass, but Breit-Wheeler with full polarizations shows substantially lower values at low pair energies – suggests we are currently overestimating pair numbers

Bremstrahlung calculation has been changed to allow full polarizations, but
CAIN cross-section only terms containing products of one or two polarizations (there should be products of 3 and 4 as well)

• terms of up to products of two polarizations depend only on the second stokes parameter which is zero for virtual photons ... no change

 full cross-section is necessary (Grozin, Proc VIII workshop on HEP & QFT, 93) Tony Hartin
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Final pair polarizations $\boldsymbol{\zeta}^{(f)}$

ICBW initial state photon Stokes vector components (per bunch crossing)

$$\zeta_{i}^{(f)} = \frac{1}{F} \sum_{ii'jj'} F_{jj'}^{i0} \xi_{j} \xi_{j'}^{'} \text{ where } F = \sum_{jj'} F_{jj'}^{00} \xi_{j} \xi_{j'}^{'}$$

- Beamstrahlung photons have almost no circular polarization component – due to beam field having constant crossed field vectors
- Breit-Wheeler pair polarization depends heavily on the photon circular polarization component, therefore pairs almost completely depolarized



ICBW final state electron Spin vector components (per bunch crossing)



CAIN original & full pol – ILC data

Pair #s for seven parameter sets





- Full polarizations show 10-20% less pair particles aross all parameter schemes
- Would be good to compare with GUINEA-PIG inclusive of polarization effects
- Usually lower pair numbers indicates lower luminosity, so important to check the effect on luminosty of full polarizations
- However luminosity remains unchanged across all schemes

CLIC 3TeV pair background results

N=4x10⁹,
$$\sigma_x/\sigma_y$$
=53/1 nm, $\gamma \varepsilon_x/\gamma \varepsilon_y$ =660/20 nm rad

Background generator	CAIN 2.35	Guinea-Pig++	CAIN Full Pol
Bethe-Heitler pairs	216214	238105	186830
Landau-Lifshitz pairs	101838	128445	48392
Bremsstrahlung pairs	149545	152649	149580
Breit-Wheeler pairs	1164	1252	1358
NCOHERENT PAIRS TOTAL	468761	520451	386160
COHERENT PAIRS	5.126 x10 ⁸	5 ×10 ⁸	5.126 ×10 ⁸

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Results

- Landau-Lifshitz pairs reduced by 50%
- Bremsstrahlung yet to be adjusted
- Overall, incoherent pairs reduce by 20%
- Breit-Wheeler pairs increase but from low starting point
- Coherent pairs unchanged

...but there are changes to coherent processes

Pair data – Energy and Pt



- As expected from the cross-section data there are fewer low energy and low P_t pairs.
- There are no changes at higher energies or higher P_t so only 'good' surprises, no nasty ones... SO FAR...
- Other effects not yet included could potentially have a large effect on background pair parameters

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Coherent Breit-Wheeler (CBW) process





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 2nd order process contains twice as many Volkov E_n

• Bunch field can act in intermediate virtual state

fermions recieve a mass shift due to bunch field and the propagator can reach mass shell whenever $n\omega \sim \omega_{b}$

CBW cross-section with simplifications

$$\frac{d\sigma_{CBW}}{d\Omega} \approx \frac{d\sigma_{BW}}{d\Omega} \int_{\neg\omega_1/\omega}^{\infty} \frac{dn}{\left[\left(n\,\omega\pm\omega_1\right)^2 + \Gamma^2\right]^2} H$$

- Can write CBW diff x-section as the ordinary BW diff x-section times a function F and a resonance
- lower bound of integration is determined physically c of m energy must be at least 2x0.511 MeV
- F is an integration of products of Airy functions for crossed beam field numerically difficult
- Γ is a resonance width determined from a self energy calculation

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• No resonances for collinear beamstrahlung and bunch momentum, but appear quickly for slightly non-collin**ear**

More work required

Summary and things to do

- CAIN modified to include virtual photon polarization, Breit-Wheeler x-sect with full polarization and allowance for polarization of initial and final states in the Bremstrahlung process
- 10-20% reduction in low energy and low Pt pair particles across all parameter schemes – due mainly to virtual photon polarization which is related to beam electric field. No change in luminosity
- Circular polarization of initial states very low (beam field is a constant crossed field) consequently final states almost completely depolarised
- Second order non-linear processes may contribute significantly to pair numbers – more work needed

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