

Fifth International Workshop on Top Quark Physics

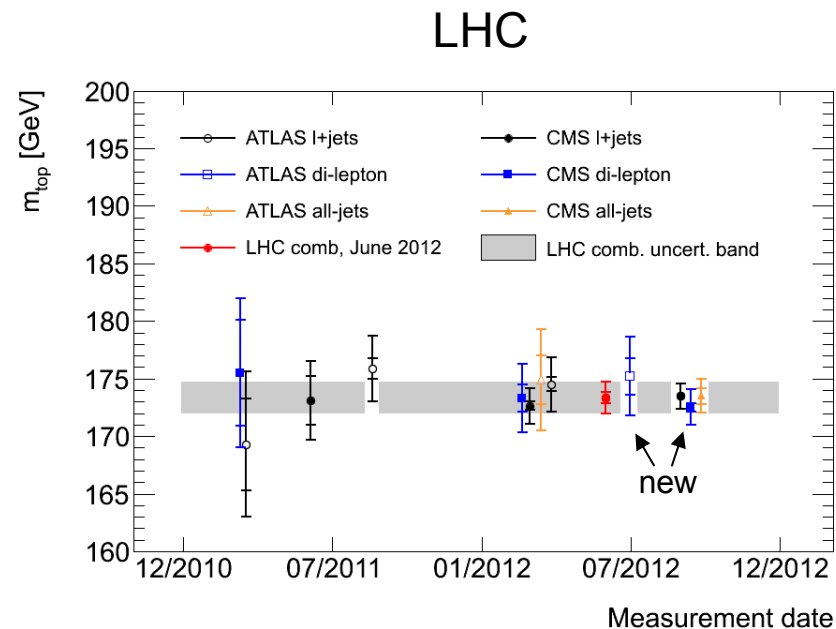
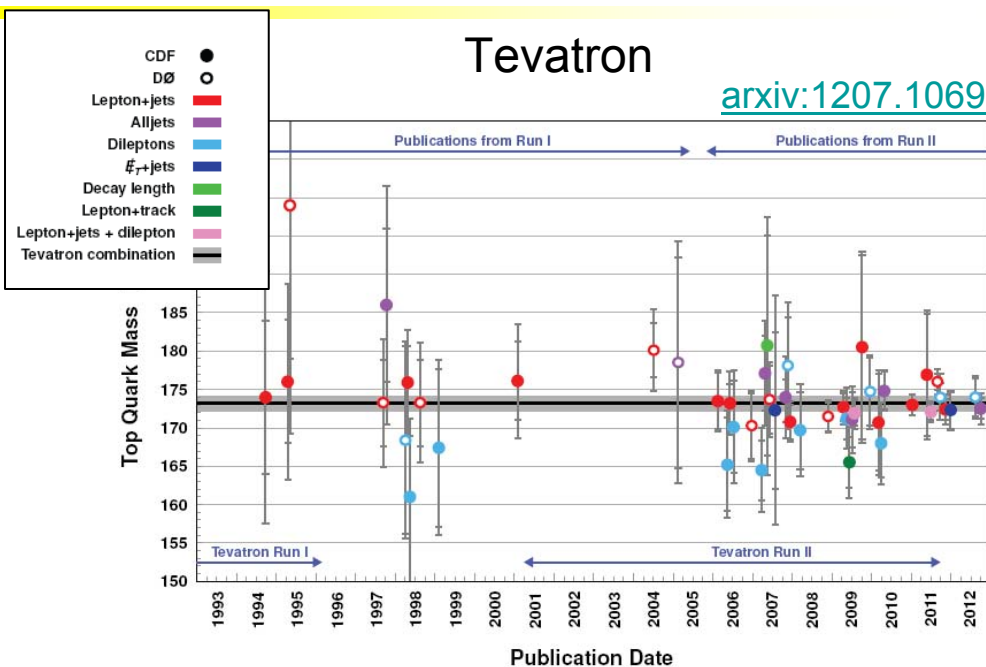
Winchester U.K. September 16-21, 2012

Experimental and theoretical uncertainties in top-quark mass measurements at Tevatron and LHC

Giorgio Cortiana MPP Munich



Measurements overview



Time evolution of the Tevatron and LHC m_{top} measurements

- Several measurements from 1995 to 2012
- Almost since the beginning (except for the very early Tevatron measurements) systematic uncertainties are dominating

	Tevatron		LHC	
[GeV]	CDF	D0	ATLAS	CMS
Stat.	0.5 – 10.3	0.8 – 12.3	0.6 – 4.0	0.3 – 4.6
Syst.	1.0 – 5.7	1.2 – 3.9	2.3 – 4.6	1.0 – 4.6
Tot	1.1 – 11.5	1.5 – 12.8	2.4 – 6.3	1.1 – 6.5



I will mainly discuss the uncertainties on the current [Tevatron](#) and [LHC](#), m_{top} combinations input measurements

▶ but also some updates (in red) as from talks by:
 Pavol Bartos (Tevatron)
 Markus Seidel (LHC)

Systematic uncertainties on top-quark mass measurements

- Theory uncertainties:
 - Signal simulation
 - Event modelling and environment

- Experimental uncertainties
 - Physics objects and detector modelling
 - identification, reconstruction, and calibration
 - Energy scales
 - (in particular for jets)

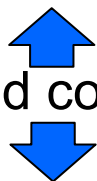
- How the systematics uncertainties are evaluated by different experiments

- What are the techniques adopted and/or prospects to reduce the uncertainties

I will discuss the uncertainties on direct m_{top} determination. m_{top} extraction from x-section will have in addition theoretical uncertainties related to the x-sec calculation (PDF, scales)

Systematic uncertainties on top-quark mass measurements

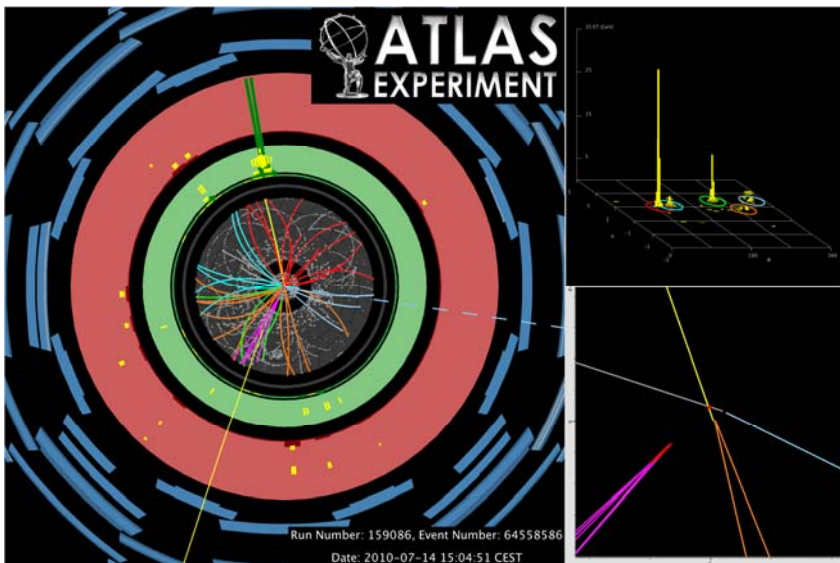
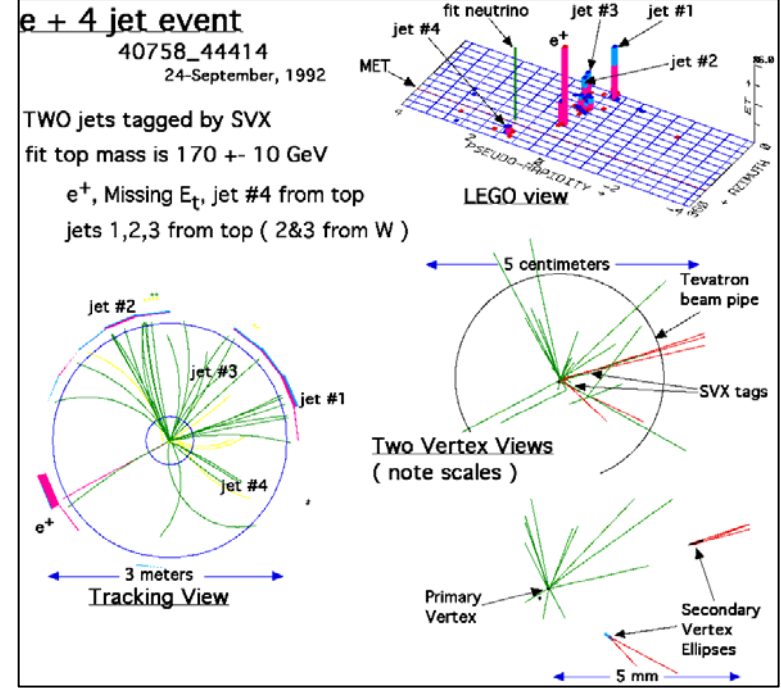
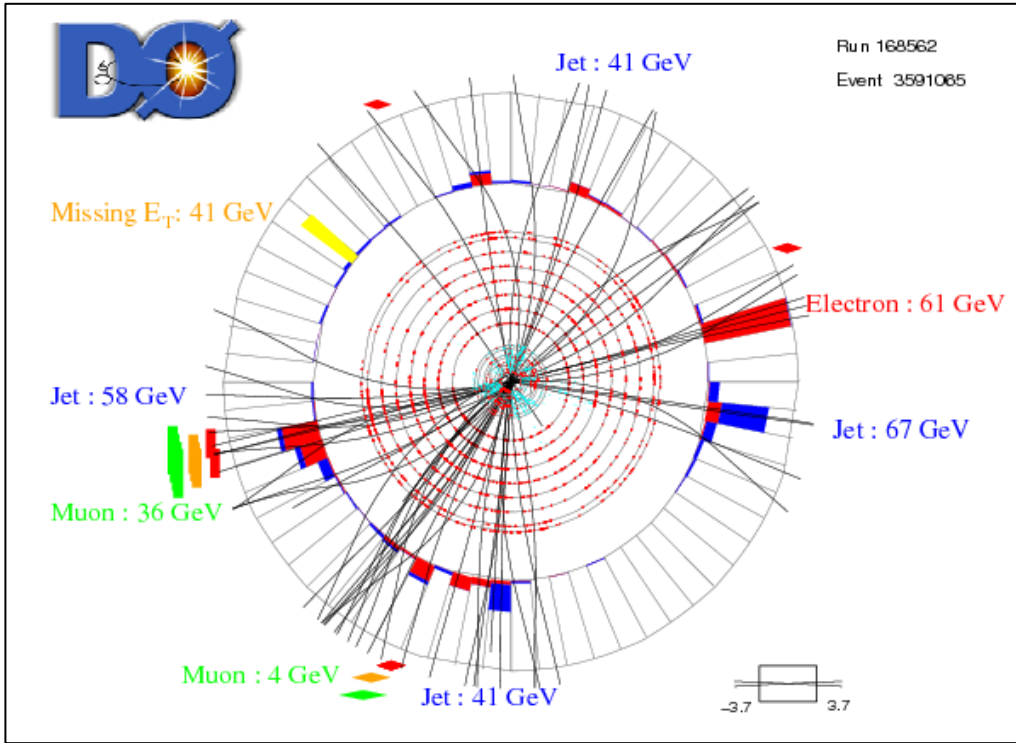
- Theory uncertainties:
 - Signal simulation
 - Event modelling and environment
- Background contamination
- Experimental uncertainties
 - Physics objects and detector modelling
 - identification, reconstruction, and calibration
 - Energy scales
 - (in particular for jets)



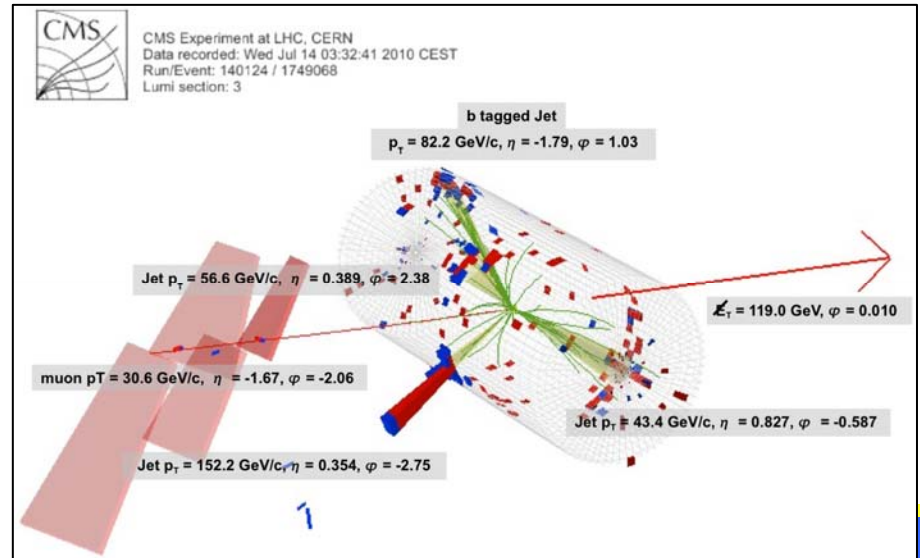
- How the systematics uncertainties are evaluated by different experiments
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I will discuss the uncertainties on direct m_{top} determination. m_{top} extraction from x-section will have in addition theoretical uncertainties related to the x-sec calculation (PDF, scales)

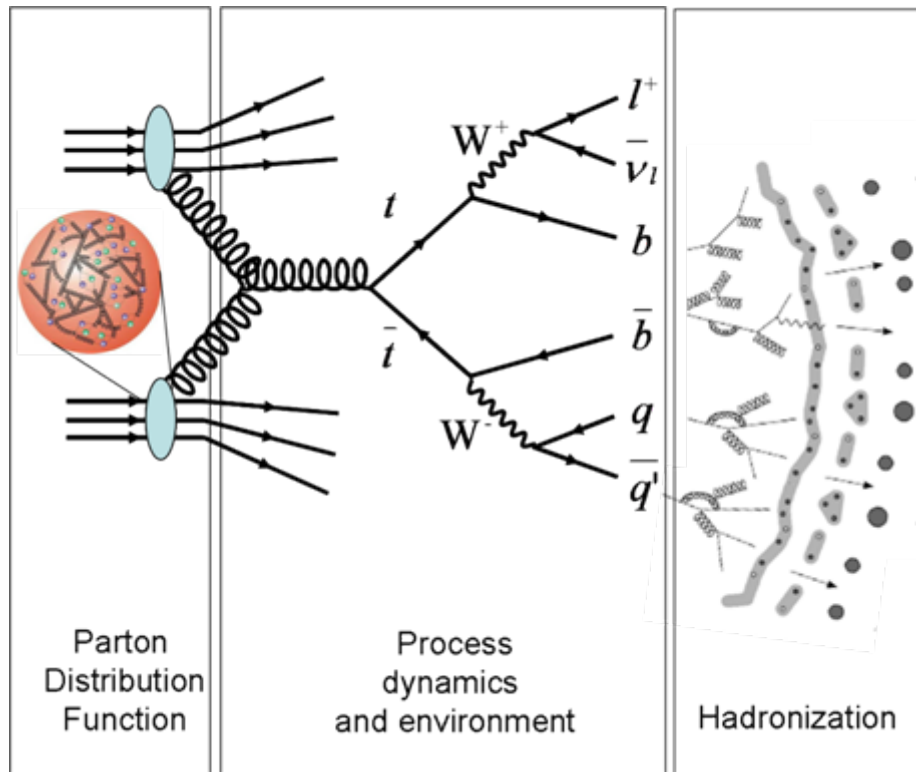
Event displays



ATLAS EventDisplay



Anatomy of a top-quark pair event



► see also talks by:
*Ivor Fleck (ATLAS),
Martijn Gosselin (CMS),
Viatcheslav Sharyy
(Tevatron)*

■ Signal simulation

- Proton PDFs
- MC generator
- Hadronization models

■ Event modelling and environment

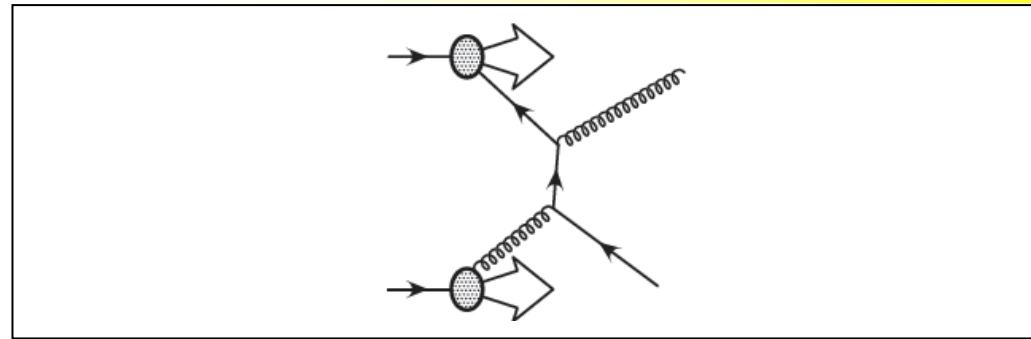
- Underlying event and MC tunes
- Color reconnection
- QCD radiation
- Pileup (multiple interactions)

■ Theory uncertainty evaluated by:

- using different generators / tools / inputs (ie PDF) for the simulation of the hard process and/or the parton shower (varied steering parameters in the simulation programs).
- May have stat components due to the limited MC sample sizes

Proton PDF

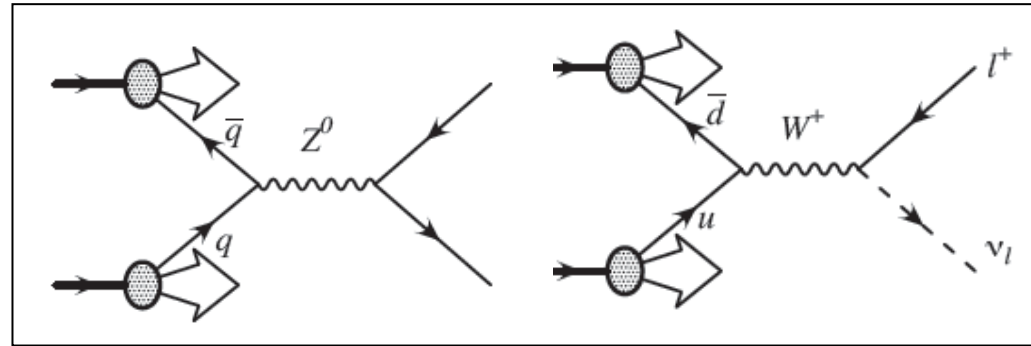
- How well do we understand the (anti) proton composition impacts the top-pair production x-sec and the event kinematics



- The PDF are obtained global fits from different groups.

Use mainly measurements from e-p DIS and in addition:

- Jets physics

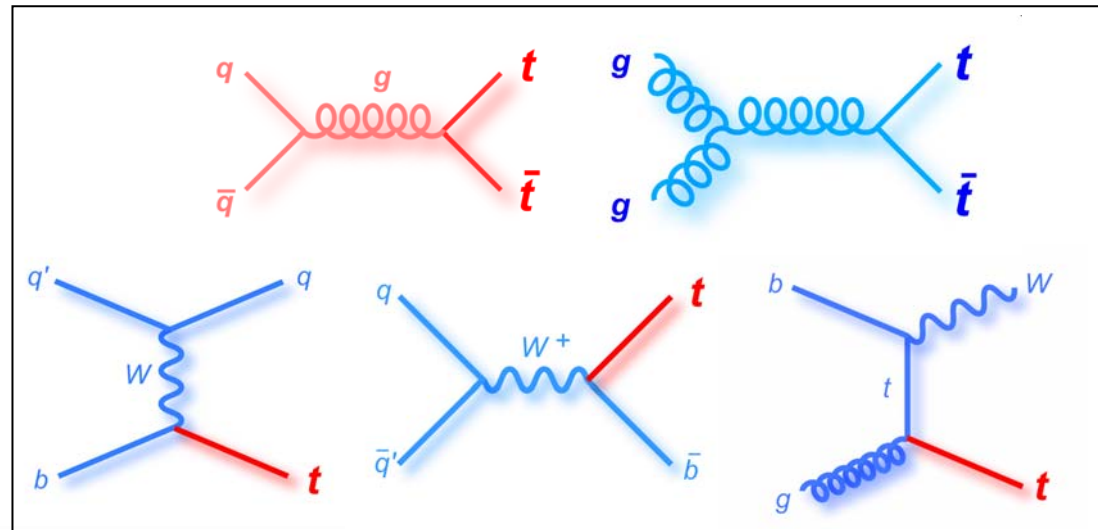


- Weak bosons production:

- Inclusive W/Z production
- W/Z asymmetries

- (Top-quark production)

- Top pair
- Single top



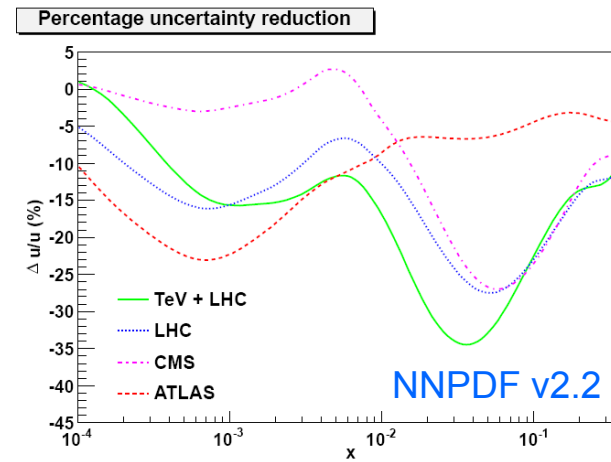
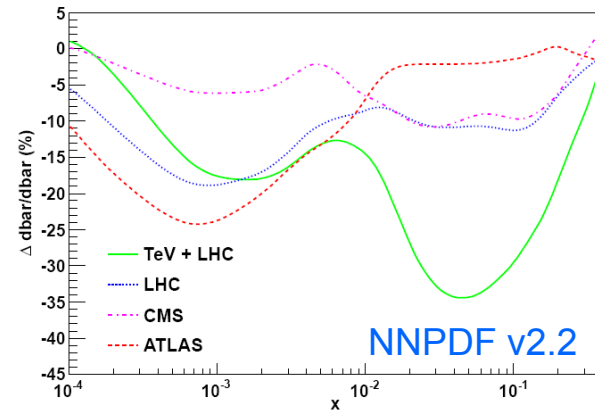
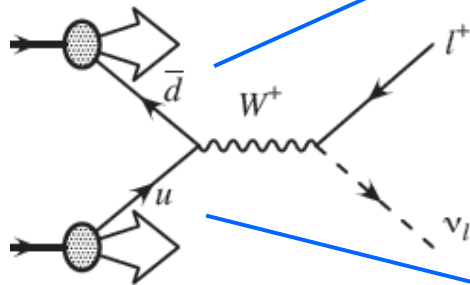
	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
PDF reweighting	<i>CTEQ6M / MRST98L</i>	<i>CTEQ6M</i>	<i>CTEQ6.6</i>	<i>MSTW08/ CTEQ6.6 / NNPDF2.0</i>
Unc. on m_{top} [GeV]	~ 0.1	~ 0.2	$0.1 - 0.6$	$0.1 - 0.5$

■ Uncertainties are in general small, but new data keeps helping reducing them
(LHC probes new x - Q^2 regions with respect to previous experiments)

■ Weak bosons production

■ W/Z asymmetries

- to reduce PDF uncertainty on valence and sea quarks

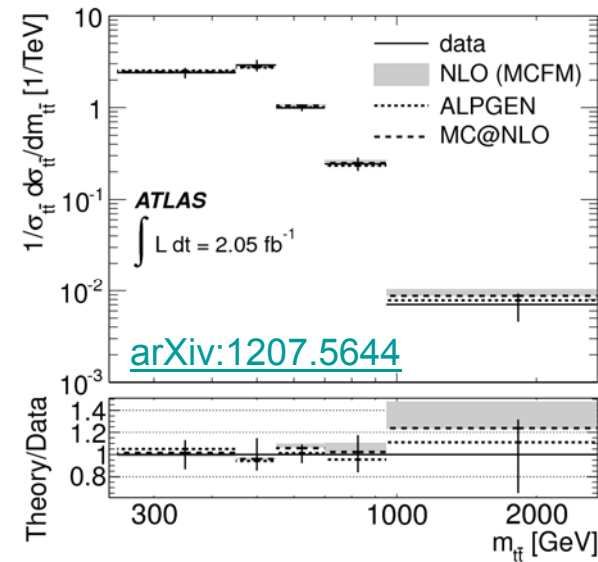
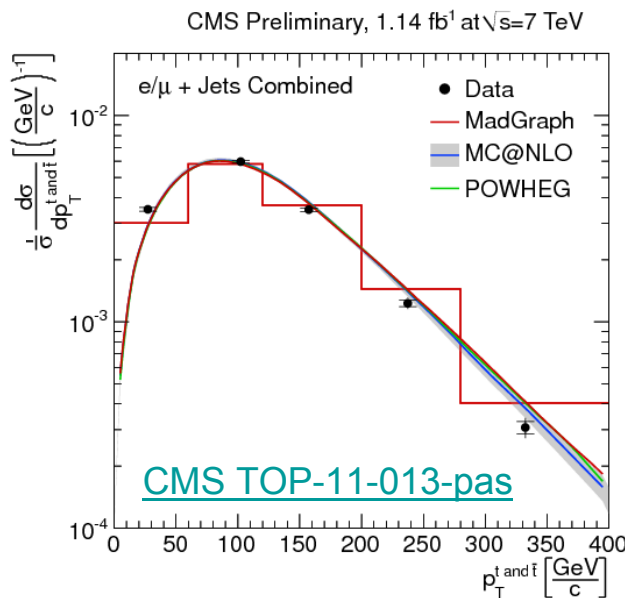
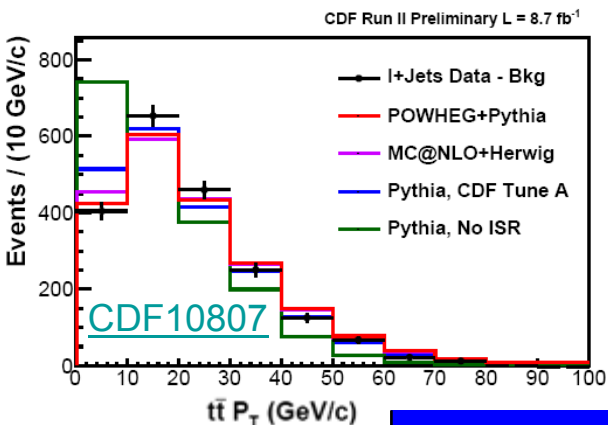
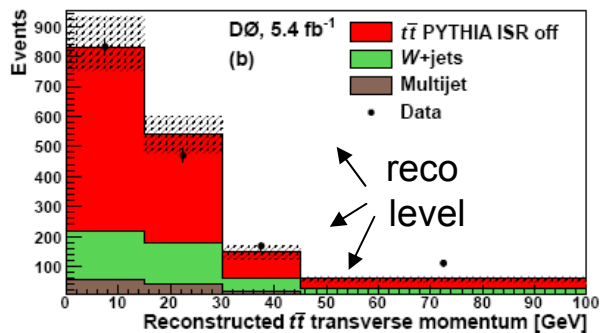
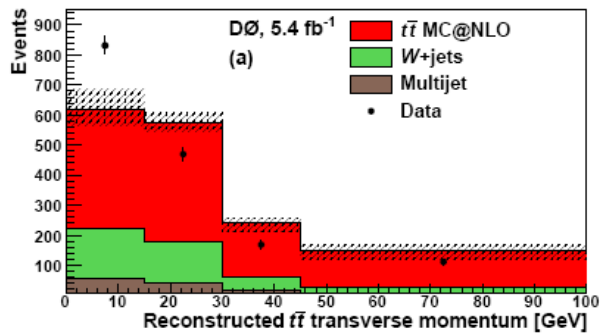


<http://arxiv.org/abs/1108.1758>

Similarly $W+1\text{jet}$ can be used to constrain the gluon PDF

Various simulation engines are available to simulate signal (and background) events.

- LO (Pythia, Herwig)
- NLO (MC@NLO, PowHeg)
- multi-leg generators (Alpgen, Madgraph)

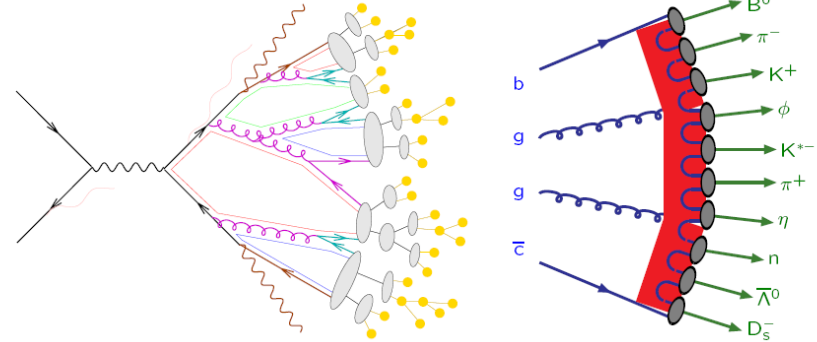


	Tevatron		LHC	
	CDF	D0	ATLAS	CMS (dilep only)
MC generator Systematics	<i>Pythia / MC@NLO</i>	<i>Alpgen / MC@NLO</i>	<i>MC@NLO / PowHeg</i>	<i>Madgraph/PowHeg/Alpgen</i>
Unc. on m_{top} [GeV]	0.1 – 0.7	0.25 – 0.6	0.3 – 1.3	<0.1 – 0.4

Hadronization models

- Describe the transition from final state partons to colorless hadrons

- Cluster and string hadronization models are implemented in Herwig and Pythia, respectively



- Hadronization systematics are considered in the determination of both the JES (Jet Energy Scale) and the MC modelling uncertainties:

- the component in the JES refers mainly to single isolated jets
- the hadronization unc. from top-pair MC accounts also for the multi-jet environment

there could be a sizeable double counting

	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
Hadronization	<i>Pythia/Herwig incl. UE syst</i>	<i>Alpgen Pythia/Herwig incl. UE syst</i>	<i>Powheg Pythia/Herwig</i>	<i>(considered in the JES syst determination)</i>
Unc. on m_{top} [GeV]	0.2 – 0.3	0.6	0.2 – 0.9	n/e

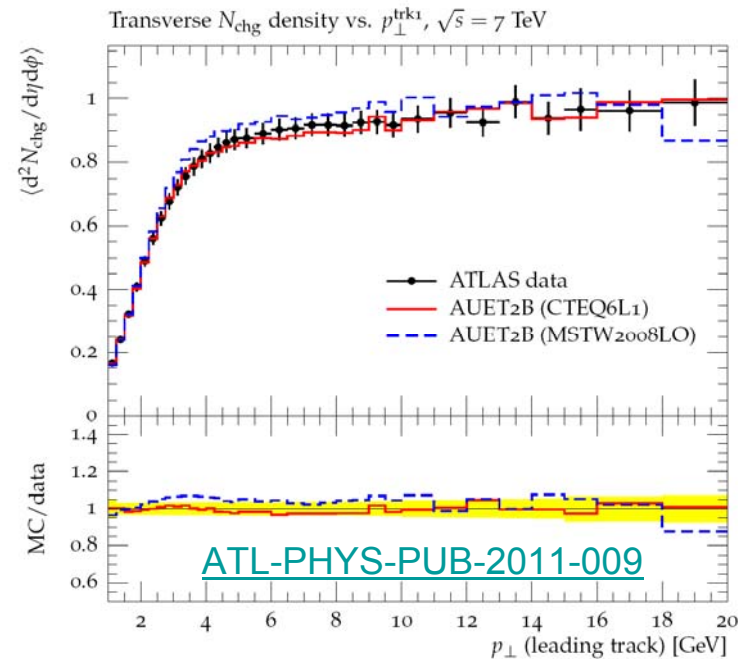
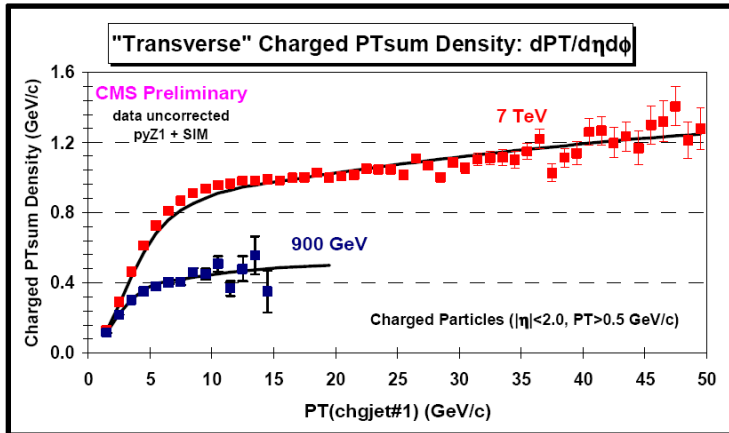
Event modelling and environment - 1

■ UE and MC tunes

- Additional physics (soft QCD) processes can contribute to the total observed activity in hard hadron collisions

- Interactions between proton remnants

<http://arxiv.org/pdf/1010.3558v1>

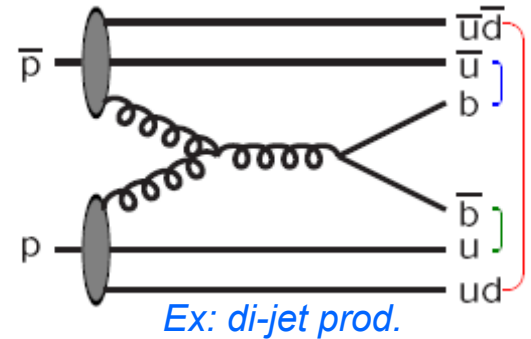


	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
UE tune	<i>Pythia/Herwig</i>		<i>Varying Pythia tunes</i>	
Unc. on m_{top} [GeV]	<i>Incl. in the hadroniz. syst</i>	<i>Incl. in the hadroniz. syst</i>	0.2 – 0.6	0.2 – 1.4

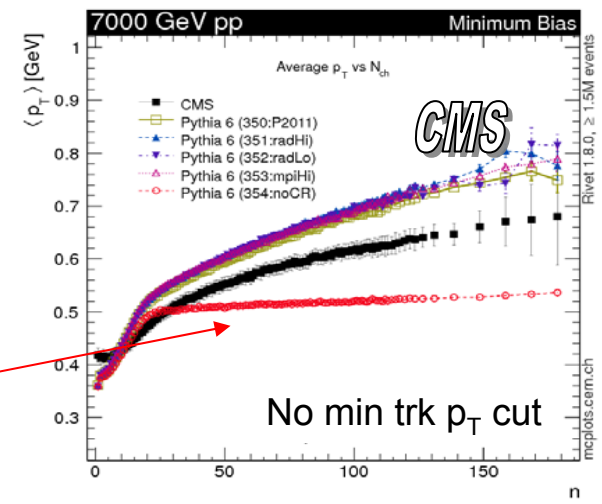
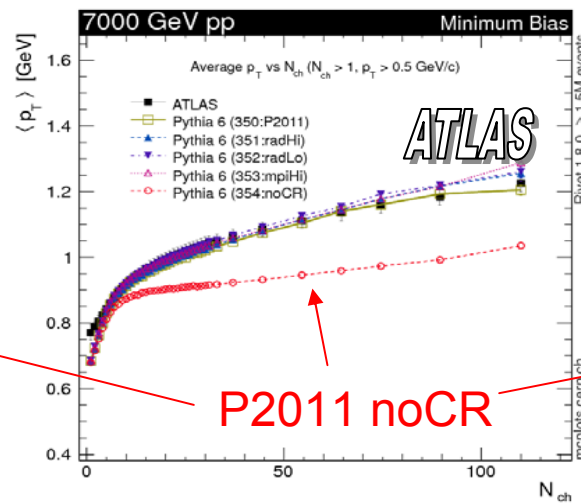
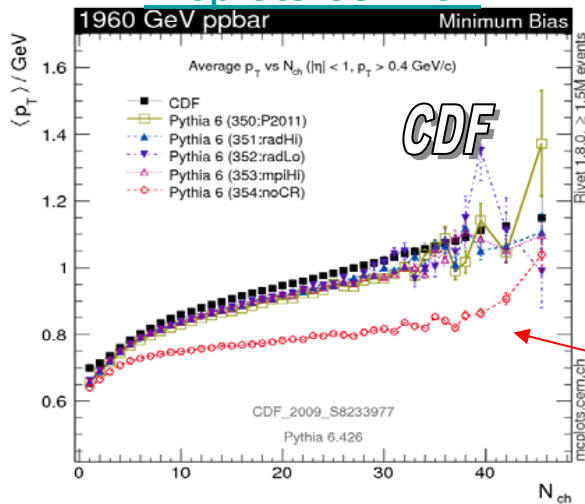
Event modelling and environment - 2

Color reconnection

- cross-talks between the hard interaction and the beam remnants in different color systems (see for example [hep-ph:0703081](https://arxiv.org/abs/hep-ph/0703081))



mcplots.cern.ch



P2011 noCR

No min trk p_T cut

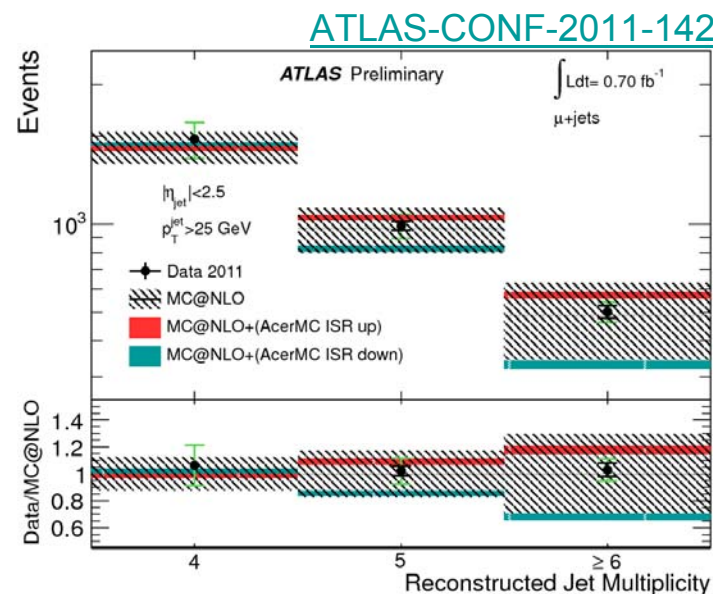
	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
Colour reconnection (CR or noCR)	<i>Pythia TuneA variation</i>		<i>Pythia Tune A and Perugia variations</i>	
Unc. on m_{top} [GeV]	0.2 – 0.5	0.3 – 0.5	0.6 – 1.2	0.1 – 0.5

Event modelling and environment - 3

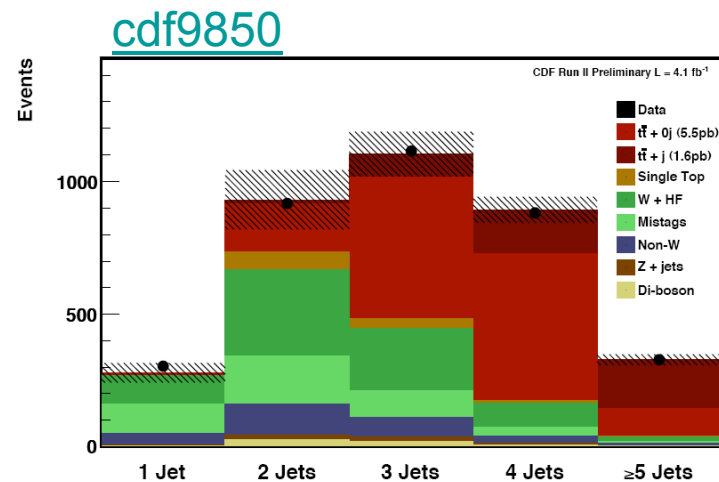
Radiation systematics (ISR/FSR)

- Increased/decreased levels of extra QCD radiation modify the event kinematics and topology (ie jet multiplicity) and alter the top-pair decay reconstruction

- Can be constrained using top-pair+jets: but in general large systematics due to detector effects (JES, b-tagging) and luminosity.



▶ see also poster by Sarah Allwood-Spiers



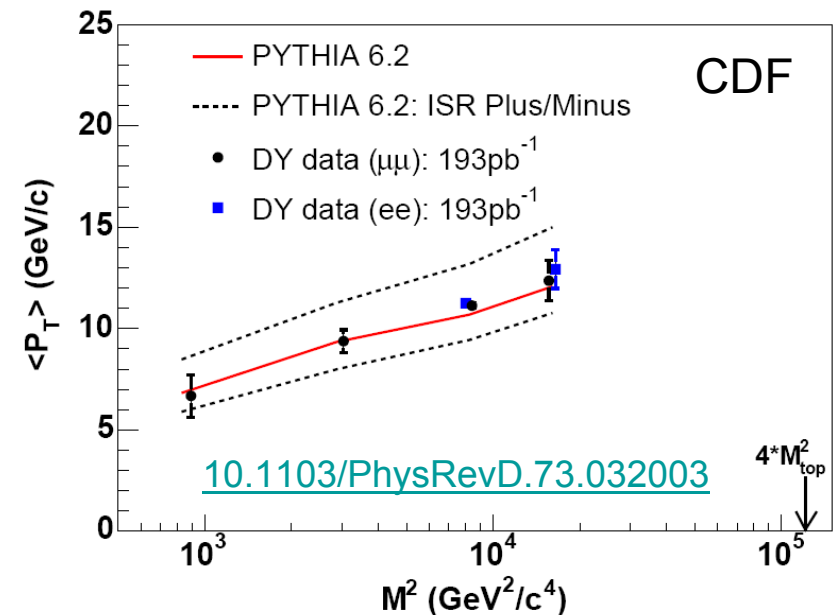
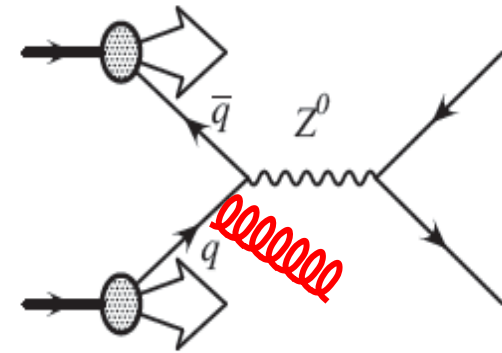
Constraining ISR with data

At Tevatron:

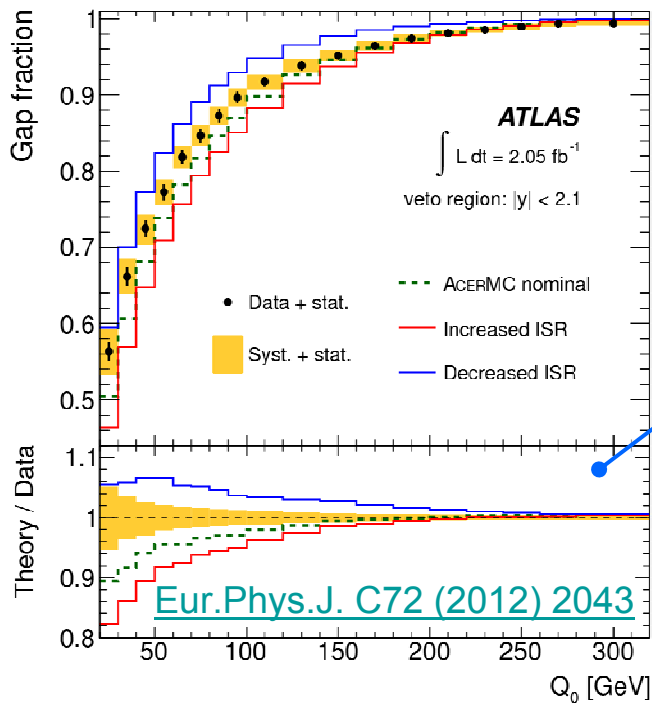
Drell-Yan events

- Same initial state as most top-pair events, but no final-state radiation
- The mean p_T of the produced di-lepton pairs is sensitive to the ISR settings
- The values of Λ_{QCD} and Q^2 in the MC that bracket the data are used to tune the ISR variations (extrapolated to the top-pair region)

- Since the initial and final-state shower algorithms are controlled by the same QCD evolution equation, the same parameter variations are used to estimate the effect of final-state radiation



Jet shapes can be used to complement FSR tunings



Unfolded data/MC comparison using AcerMC (LO) + Pythia ISR variations

At LHC (for example),

■ As an alternative to top-pair+jets, use the Gap Fraction in top-pair di-lepton events

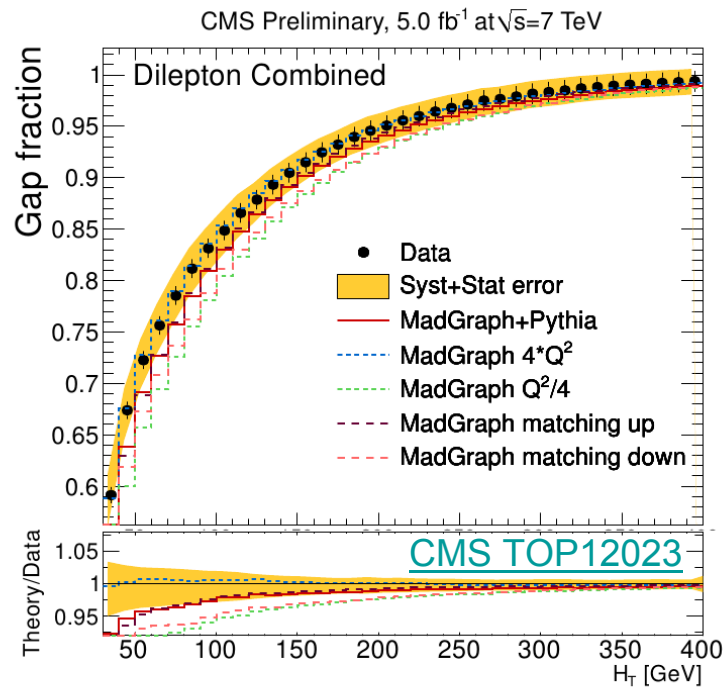
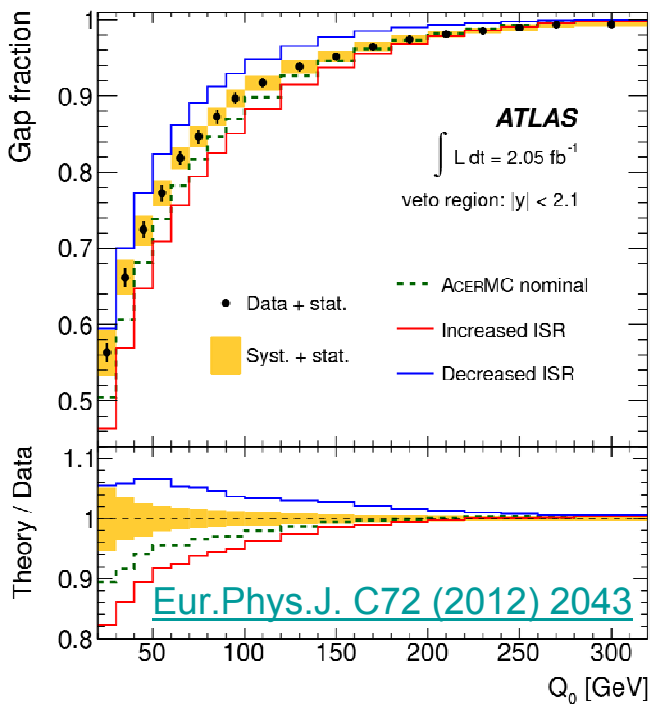
■ measure the fraction of top-pair dilepton events which do not contain an extra jet with $p_T > Q_0$

$$f(Q_0) = \frac{n(Q_0)}{N}$$

sensitive to the leading p_T emission

■ Many of the experimental systematic uncertainties cancel in the ratio

■ The data are used to constrain the modelling of quark and gluon radiation in top-quark pair events



Unfolded data/MC comparison using AcerMC (LO) + Pythia ISR variations

At LHC (for example),

As an alternative to top-pair+jets, use the Gap Fraction in top-pair di-lepton events

- measure the fraction of top-pair dilepton events which do not contain an extra jet with $p_T > Q_0$
- the veto criterion can be extended to probe jet activity beyond the leading additional jet.

$$f(Q_0) = \frac{n(Q_0)}{N}$$

sensitive to the leading p_T emission

$$f(Q_{\text{sum}}) = \frac{n(Q_{\text{sum}})}{N}$$

sensitive to all hard emissions accompanying the system ($Q_{\text{sum}} = H_T$)

Many of the experimental systematic uncertainties cancel in the ratio

- The data are used to constrain the modelling of quark and gluon radiation in top-quark pair events

Radiation systematics summary

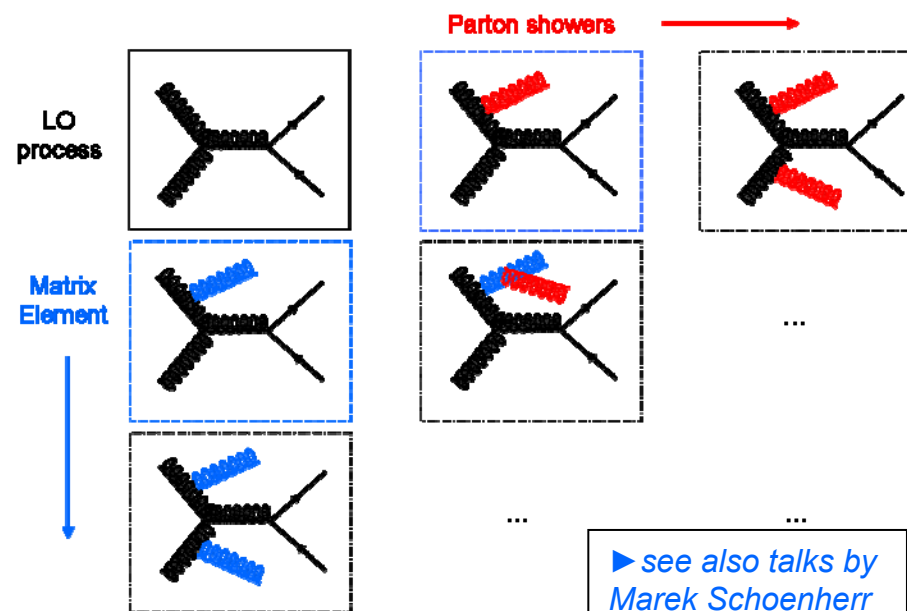
	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
QCD radiation (ISR/FSR)	<i>Pythia More/Less PS</i>		<i>ACEMC + Pythia More/Less PS</i>	<i>Q² variation in Madgraph+Pythia + variation in the ME/PS jet-to-parton matching scale (for initial measurements also Pythia More / Less PS)</i>
Unc. on m_{top} [GeV]	0.1 – 2.7	0.3 – 0.4	0.5 / 1.0 – 2.5 <i>(Before gap fraction analysis)</i>	0.3 – 1.2

■ In multi-leg MC, the same multi parton final state can be produced by the matrix element or by parton shower

- To avoid double counting need a matching procedure

■ matching scale + Q^2 scale variations describe to a large extent the same physics effect as ISR/FSR variations (preliminary investigation within the TopLHC WG)

- This approach is used with MadGraph + Pythia by the CMS collaboration

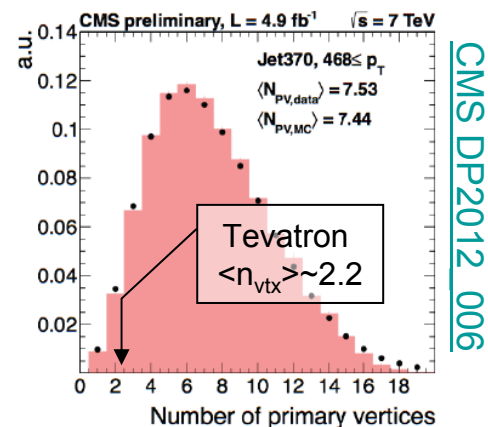
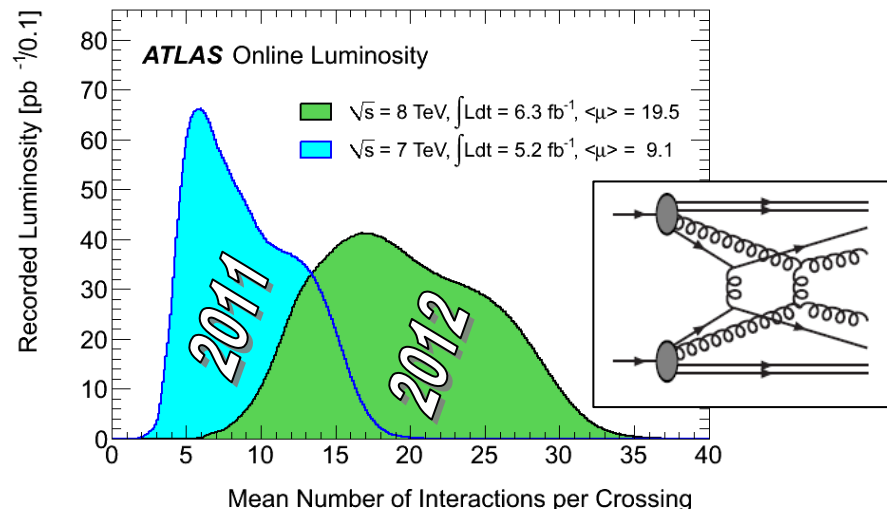


▶ see also talks by Marek Schoenherr

Event modelling and environment - 4

Pileup

- Overlay of additional low- p_T events (from MC or data for D0):
 - contribute to building up the total amount of energy and
 - cause color exchanges between the remnants, thereby increasing the number of particles produced in the hadronization stage.
- Affects both the JES determination (sep. component), the m_{top} resolution and the event kinematics
- Reweighting of MC event to match the data luminosity profile is required.

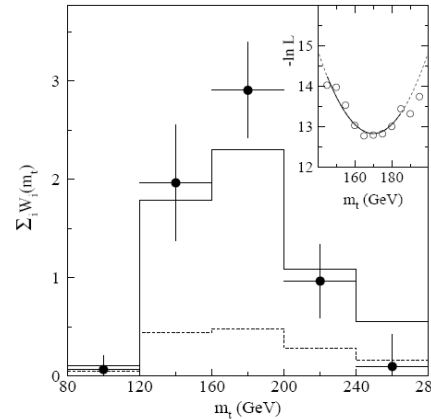


	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
Pileup	Data/MC diff in n_{vtx}	Unc in the reweighting	Data/MC after reweighting	
Unc. on m_{top} [GeV]	0.1 – 0.2	0.1 – 0.2 (1.2 for Run1 di-lep)	0.05 – 0.7	0.1 – 1.0

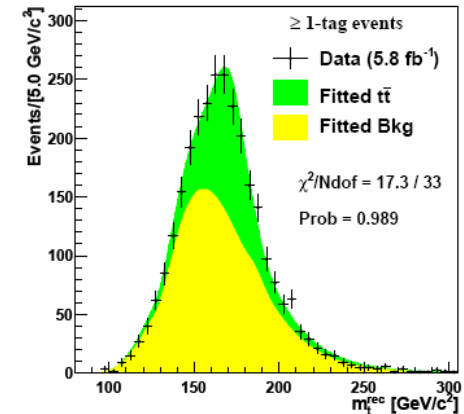
Background contamination

- Uncertainties on the background normalization and shape induce m_{top} unc (via the fitting procedure)
 - analysis, kin. sel. and decay channel dependent
- Uncertainties can be classified as
 - from MC: W/Z+jets shape, and sub-dominant contributions from single top and di-boson production
 - from data: mainly W+jets norm and QCD contributions and/or fake leptons (statistical unc. components can be reduced increasing the samples size)

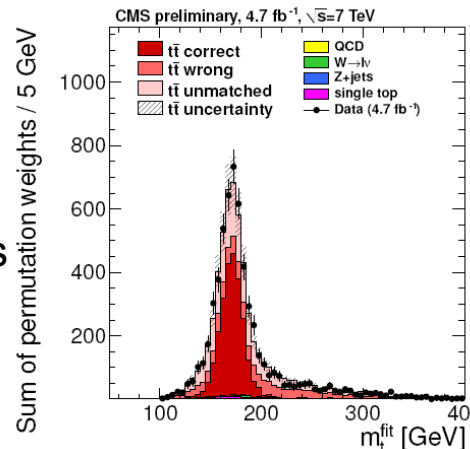
D0 di-lep Run I



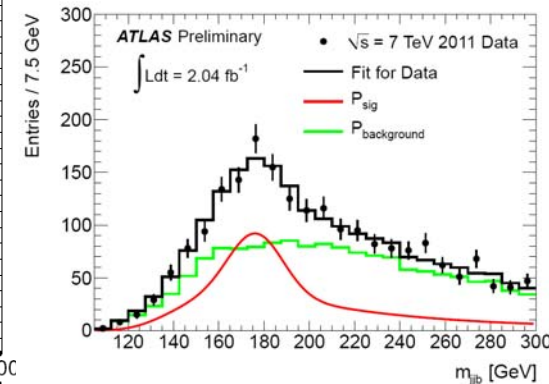
CDF all-jet Run II



CMS μ +jets



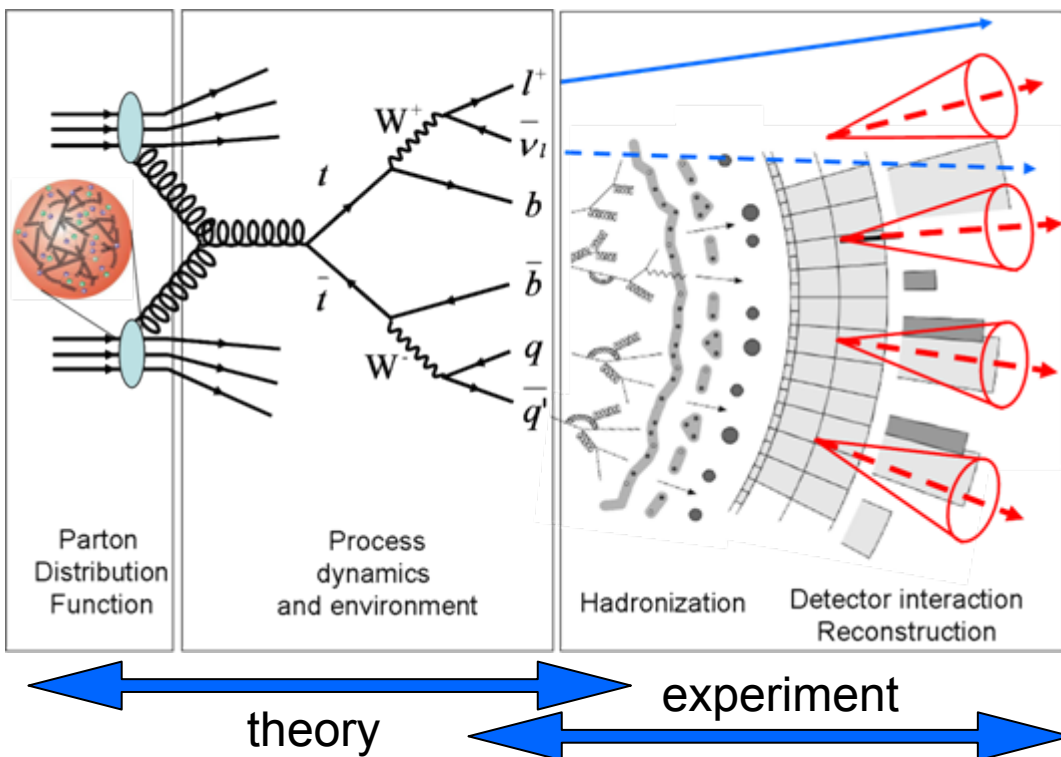
ATLAS all-jet



see also talks by: Frederic Deliot and Joseph Boudreau

Syst type / Values in GeV	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
Background from MC (theory)	<0.01 – 1.7	<0.01 – 1.0	0.1 – 1.8	0.1 – 0.2
Background from data	0.1 – 0.6	0.2	0.6 – 1.9	0.2 – 0.4

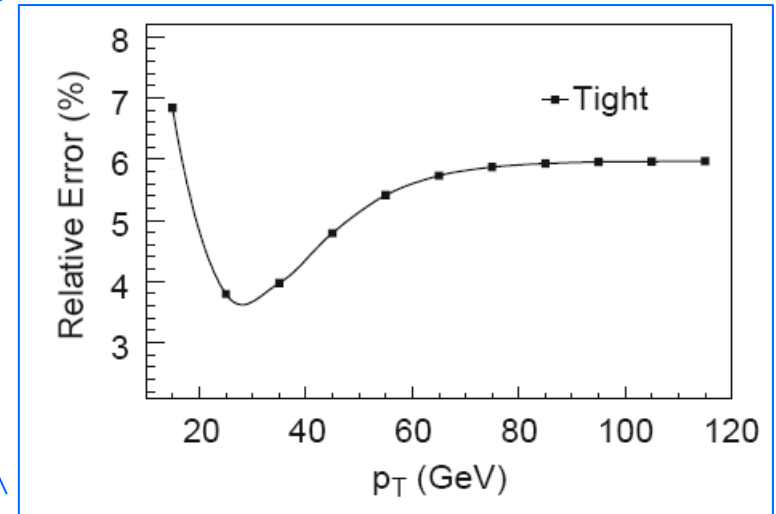
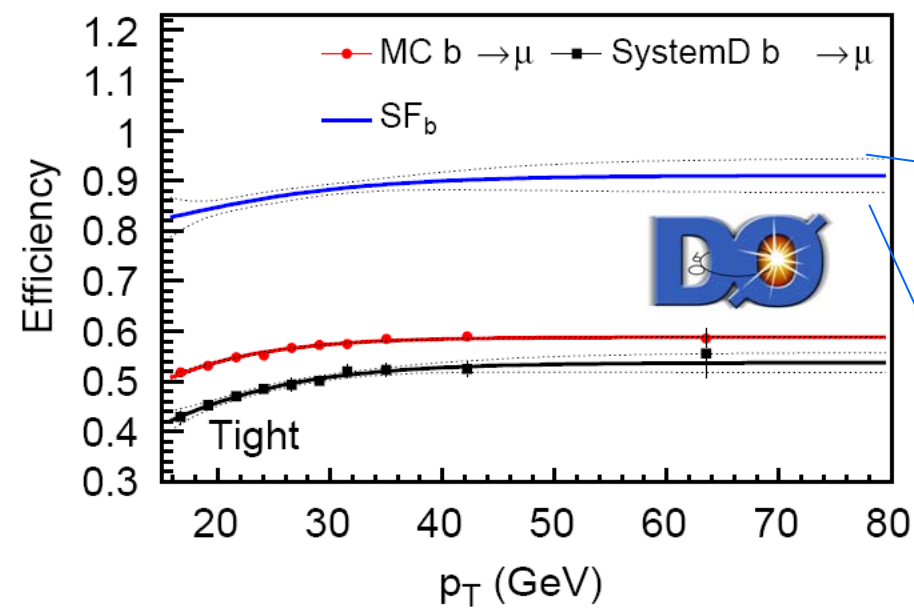
Physics objects and detector modelling



- b-tagging
- Jet energy resolution
- Jet reconstruction efficiency
- E_T^{miss} modelling
- Leptons momentum and energy scale
- **Jet Energy Scale**
[deserves separate treatment]

▶ see also talks by:
Carmen Diez Pardos (CMS)
Veronique Boisvert (ATLAS)

b-tagging calibration



***b*-tagging is an important handle for background reduction**

Algorithms need to be calibrated, to account for detector effects and possible mis-modeling of the jet/trk/event variables in the MC with respect to the data.

Data samples are used to determine:

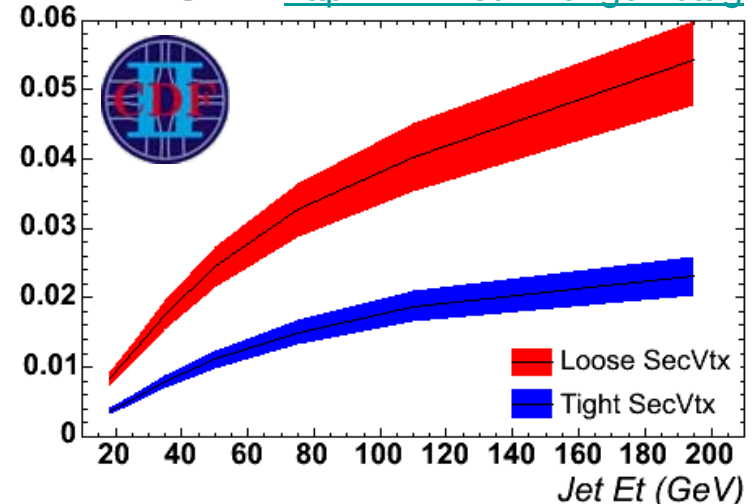
b-tagging efficiency

mis-tag rate

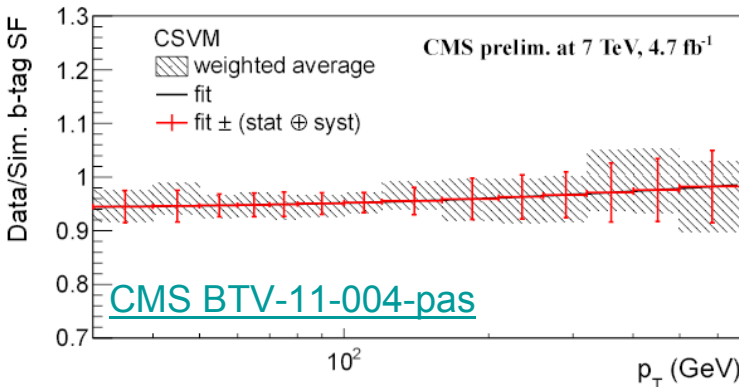
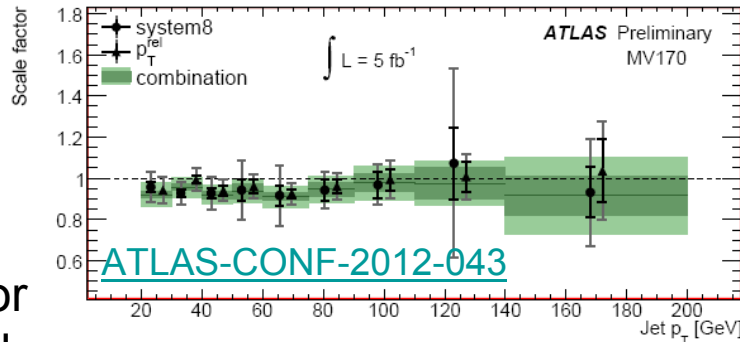
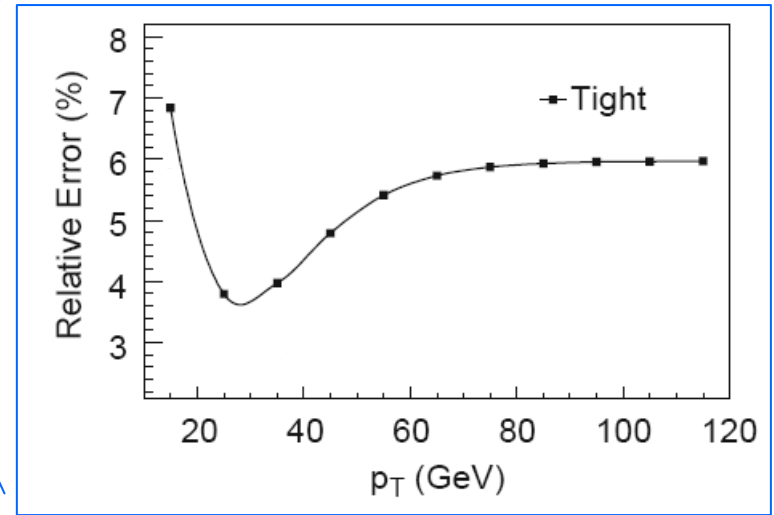
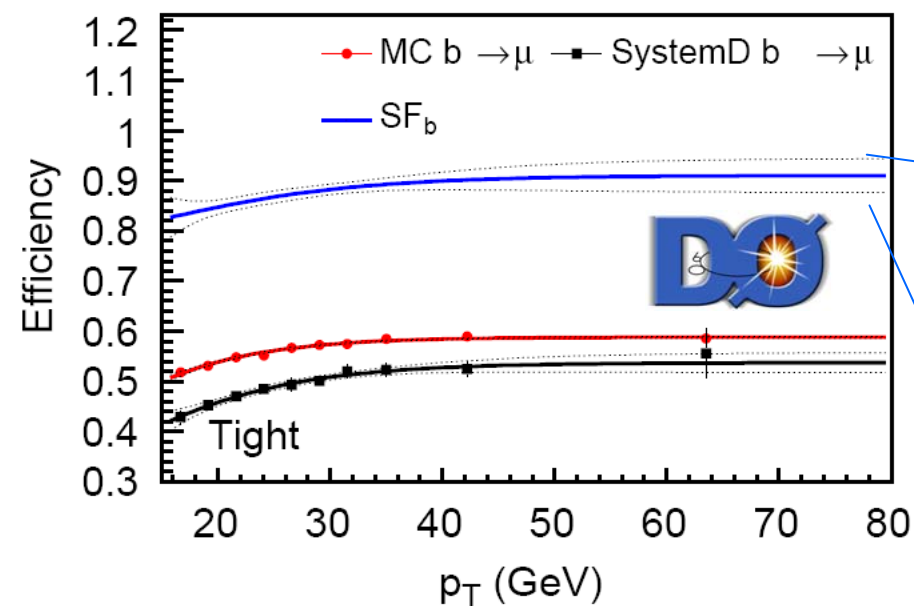
and the corresponding uncertainties.

Include *b*-jet modelling uncertainties such as semilep BR, *b*-jet fragmentation

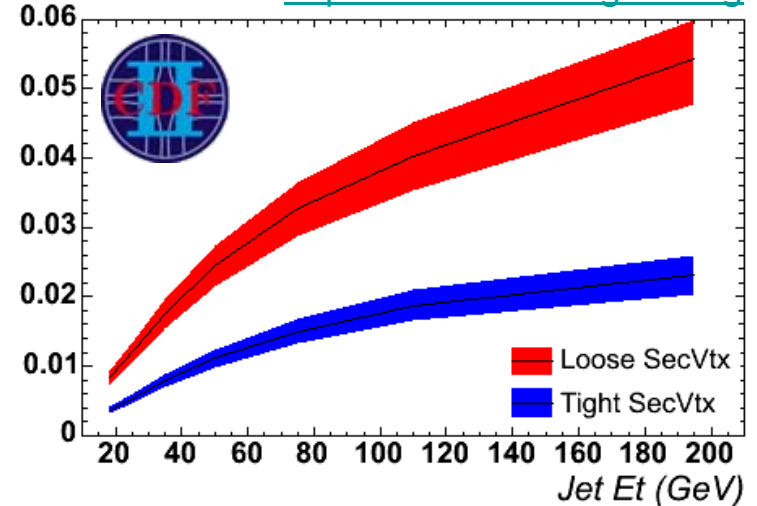
SecVtx Mistag Rate <http://www-cdf.fnal.gov/btag>



b-tagging calibration



SecVtx Mistag Rate <http://www-cdf.fnal.gov/btag>

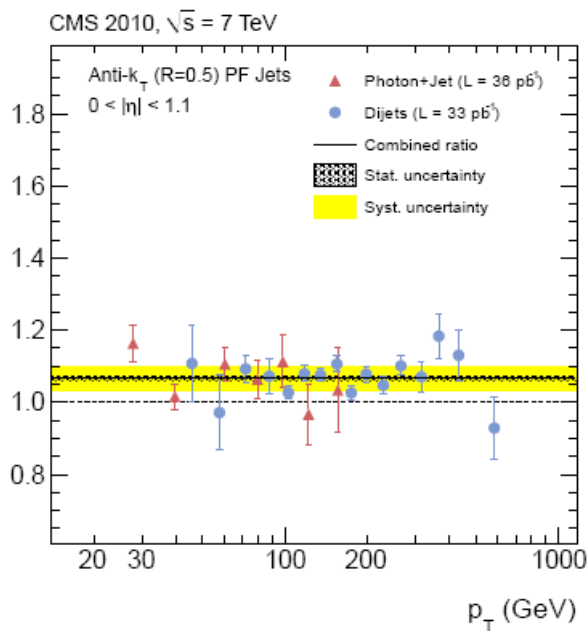
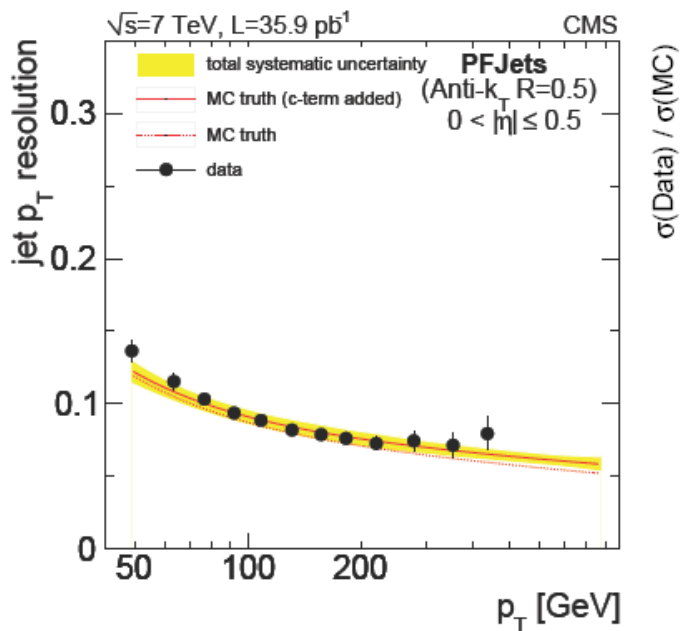


Similar concepts for ATLAS and CMS

Jet energy resolution and reco efficiency

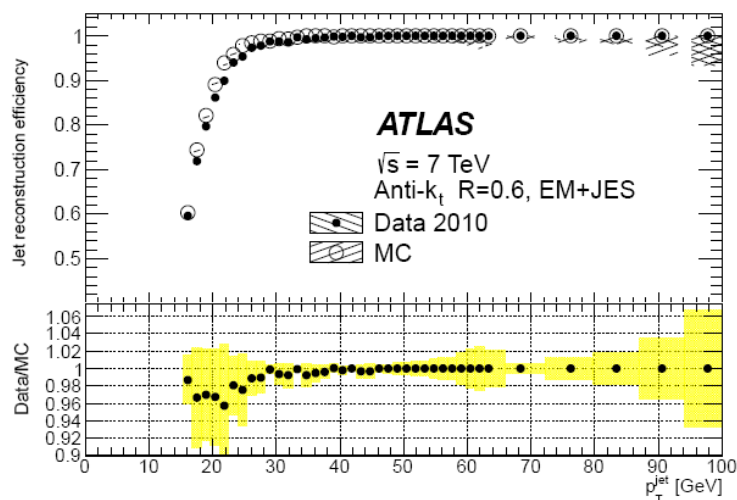
Jet Resolution

- Difference in the jet energy resolution between data and MC have to be accounted for
- ► can result in broader/narrower invariant mass distrib.



Jet Reconstruction efficiency

- The jet reco efficiency is measured relative to track jets in data and MC.
- The Data/MC difference, is propagated to the analyses



important only for low p_T jets

Summary of detector uncertainties

	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
B-tagging / b-jet modelling	<i>Data/MC scale factors varied within uncertainties</i>			
Unc. on m_{top} [GeV]	<i>0.2 – 2.5 (for L_{xy})</i>	<i>0.1 – 0.2</i>	<i>0.3 – 0.5</i>	<i>0.1 – 0.5</i>
Quite some statistical components in the b-tagging calibration, can be reduced by increased sample sizes				
Jet energy resolution / reco eff	<i>ne</i>	<i>MC Jet energy smearing and random jet drop</i>	<i>MC Jet energy smearing</i>	
Unc. on m_{top} [GeV]	<i>negligible</i>	<i>0.2 – 0.3</i>	<i>0.1 – 0.9</i>	<i>0.1 – 0.5</i>

Summary of detector uncertainties

Missing E_T modelling

■ the E_T^{miss} absorbs the scale and resolution uncertainty of all physics objects like e , γ , μ , τ and jets.

■ There are also “soft terms” including un-clustered or energy not associated to physics objects

	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
B-tagging / b-jet modelling	<i>Data/MC scale factors varied within uncertainties</i>			
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E_T^{miss} modelling (soft term)	<i>ne</i>	<i>ne</i>	<i>Variation of the soft term of E_T^{miss} within uncertainties</i>	
Unc. on m_{top} [GeV]	–	–	<i>0.1</i>	<i>0.1 – 0.4</i>

Summary of detector uncertainties

Missing E_T modelling

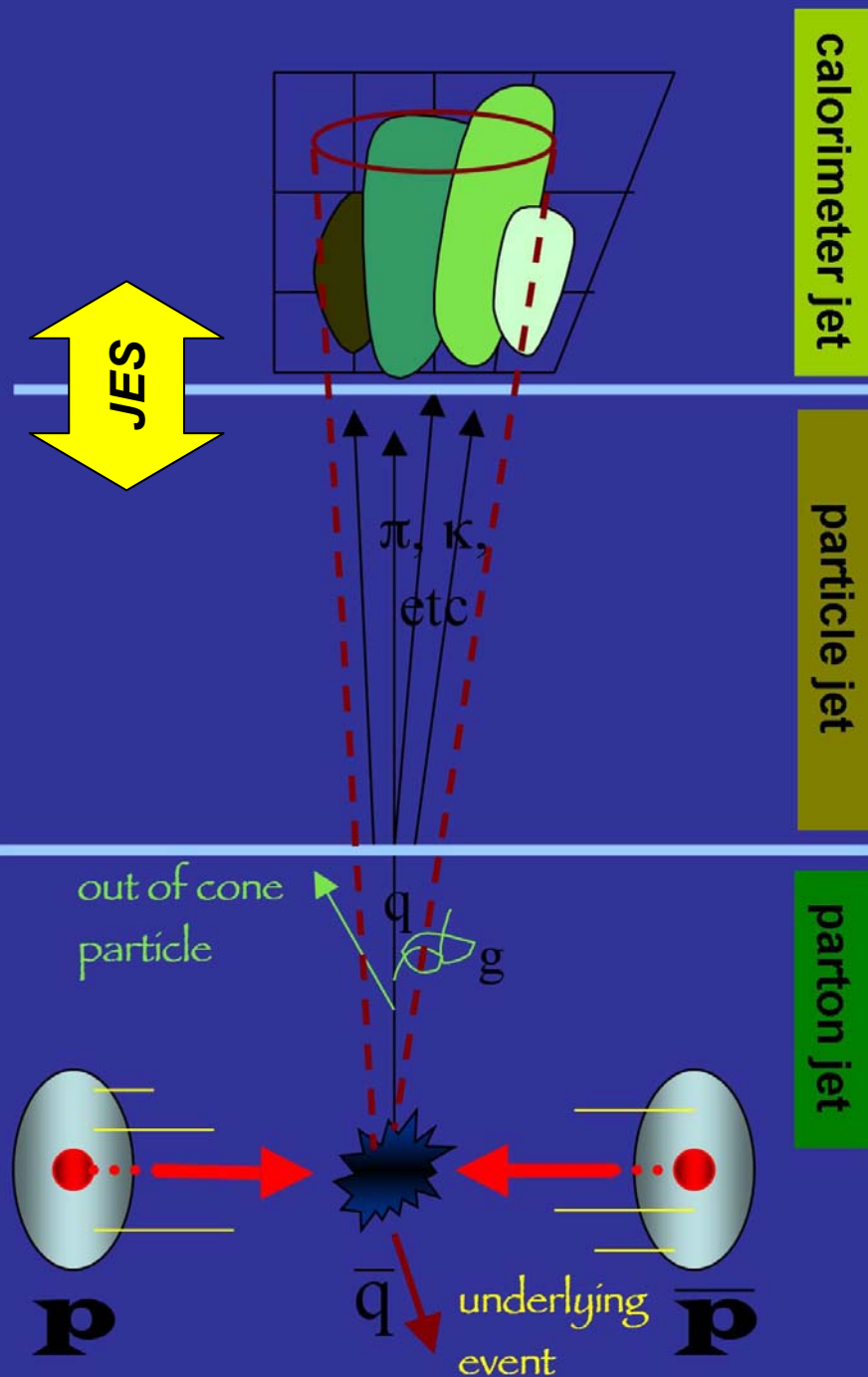
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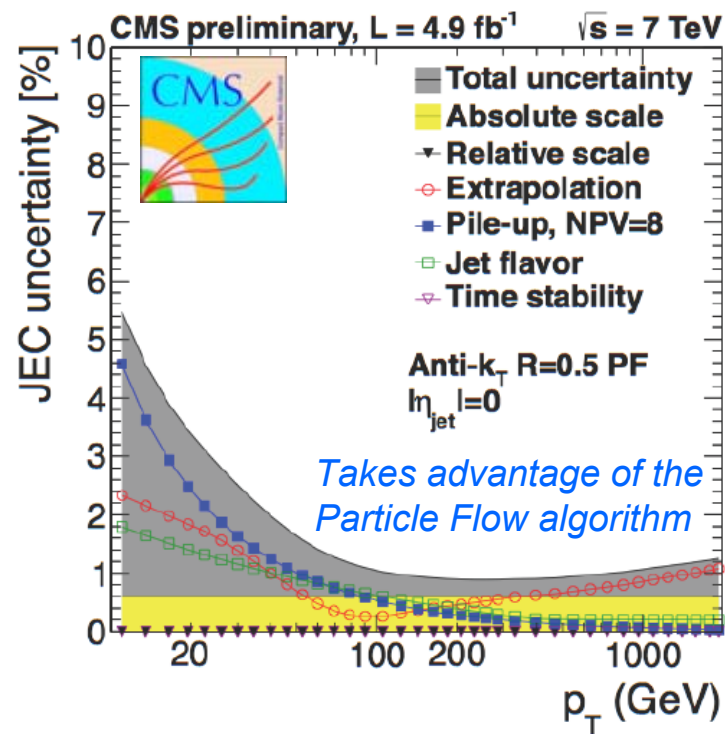
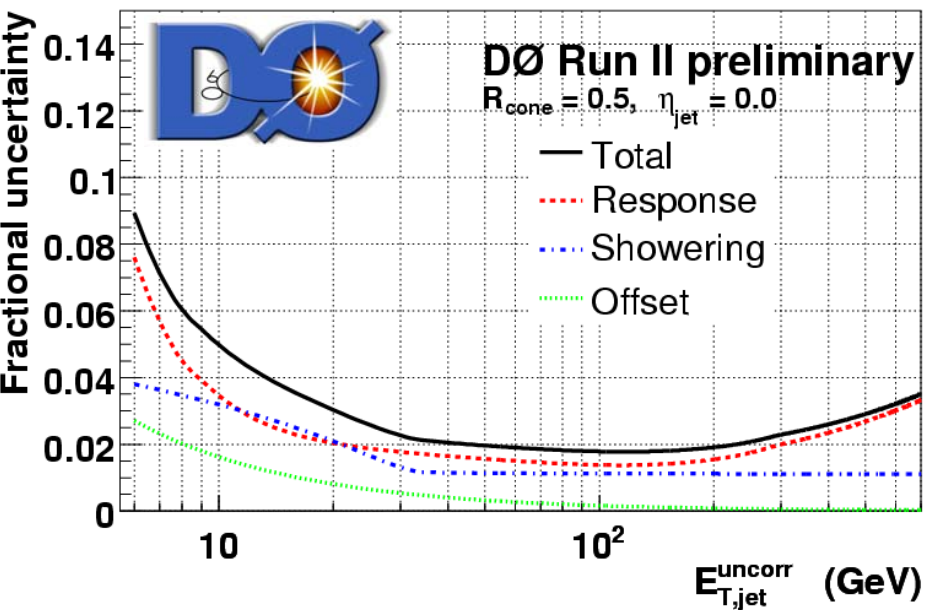
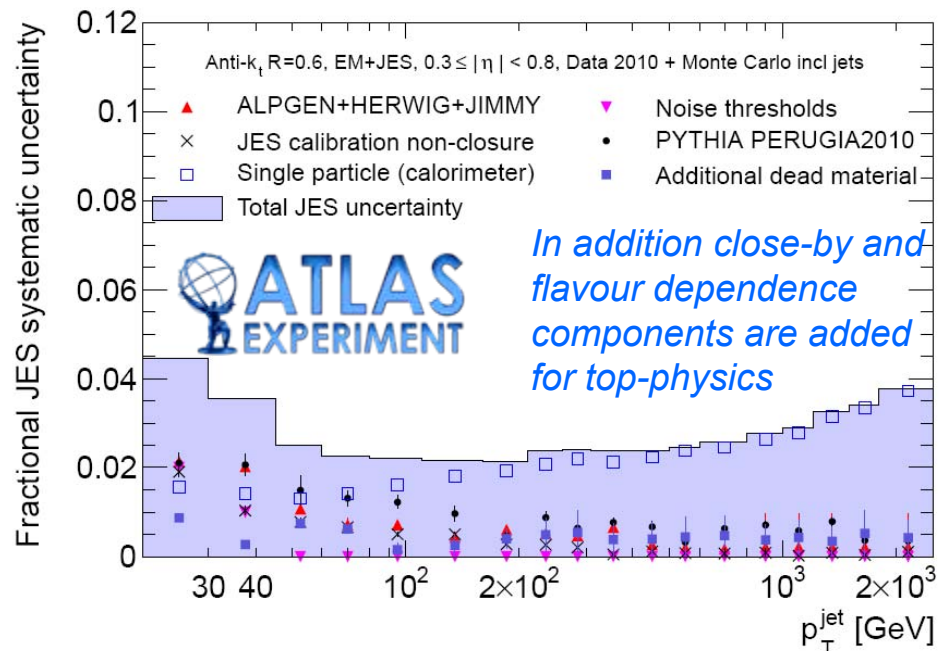
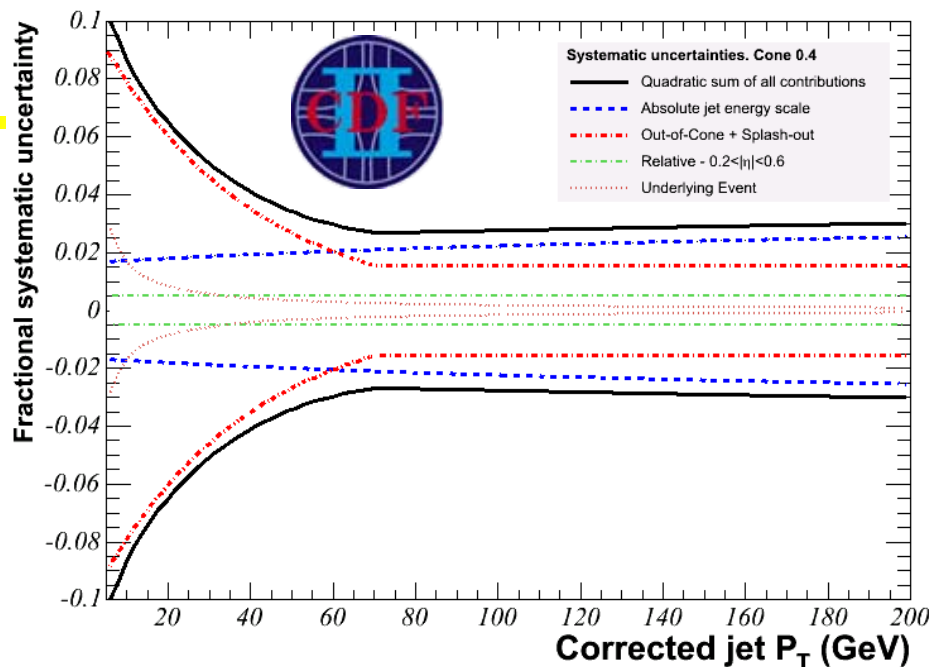
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Unc. on m_{top} [GeV]	–	–	<i>0.1</i>	<i>0.1 – 0.4</i>

Lepton momentum / energy scale

■ In general a sub-dominant uncertainty contribution that varies between <0.1 GeV (l+jets) to $0.2 – 0.3$ GeV (di-lep)



- Describes difference between reconstructed jet and particle jet (or partons in the CDF case)
- Non-trivial determination.
 - involve both MC and in-situ techniques (ie γ/Z +jet, di-jet balance)
- The jet energy scale uncertainty is derived from in-situ measurements along with systematic variations of the Monte Carlo simulation
- Typically uncertainties are $O(\text{few } \%)$ and vary with jet properties (p_T, η)
 - $\sim 100\%$ correlated with m_{top} (when using jet info for its reconstruction)
 - $\sigma(JES) = 3\% \rightarrow \sigma(m_{top}) \sim 5 \text{ GeV}$

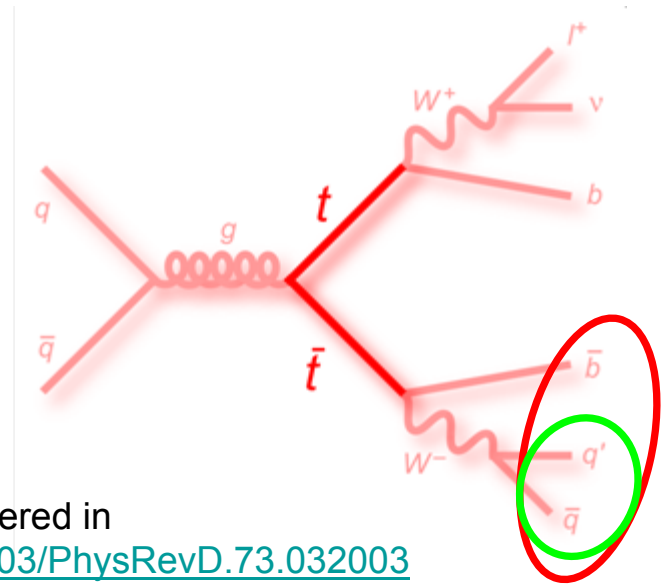


In-situ constraint of the light-quark JES

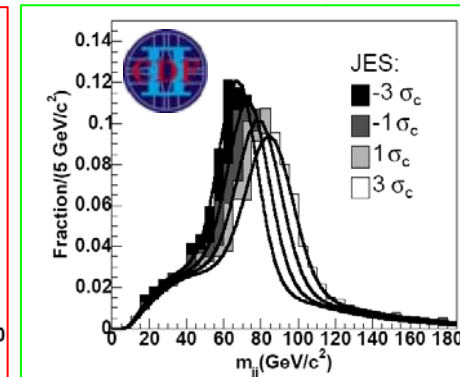
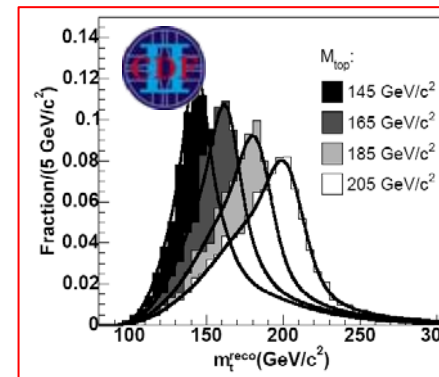
Most of the recent m_{top} measurements in the l+jets ($E_{\text{T}}^{\text{miss}}+\text{jets}$) and all-jets channels have developed techniques to constrain in-situ the light-quark jet energy scale using in top events information from the $W \rightarrow jj$ decay.

An example also exists for the transfer of the JES constraint in the l+jets channel to the di-lepton one (taking into account the environmental difference <http://arxiv.org/abs/1201.5172v1>)

This allow a large reduction of the impact of the JES uncertainty on the m_{top} determination, at the price of an additional contribution to the stat uncertainty (2D fit)



Pioneered in [10.1103/PhysRevD.73.032003](https://arxiv.org/abs/10.1103/PhysRevD.73.032003)



Statistical component

Syst type / Values in GeV	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
light JES	0.2 (Lxy) 0.5– 4.4	0.6 – 3.4	0.7 – 2.1	0.2 – 2.1
Jet Scale Factor (from in-situ $W \rightarrow jj$)	0.5 – 1.5	0.5 – 0.6	0.4	0.4

In-situ constraint of the light-quark JES - 2

- The method generally works well, however at the precision level of the latest measurements (and JES determinations) can become less powerful
 - In the recent CMS m_{top} result in the all-had channel ([CMS TOP-11-017-pas](#)), the gain of 2D method for the jet energy scale uncertainty reduction was not compensated by the increase of other systematics effects altering the m_W (and m_{top}) distribution, namely the CR and UE event syst.
 - ▶ The 1D result at the end was found to be more precise
 - This may be attributed to a sensitivity of the 2D method to additional relative differences in the light- and b-jet responses that can be introduced by the event modelling.

b- (or flavor-) specific JES

- The standard JES is set using light-quark/gluon jets (similar arguments for the 2D in-situ method)
- The jet response to heavy-flavor jets can be different:
 - Semileptonic decays
 - Differences in the shower and particle content of b-jet wrt the light-jets
- CDF measures the b-to-light-quark energy scale difference applying single particle responses measured in data or predicted by MC to samples of jets from top pair events.

$$\frac{\left(p_T^{data} / p_T^{MC}\right)_{b-jets}}{\left(p_T^{data} / p_T^{MC}\right)_{light-jets}} = 1.01$$

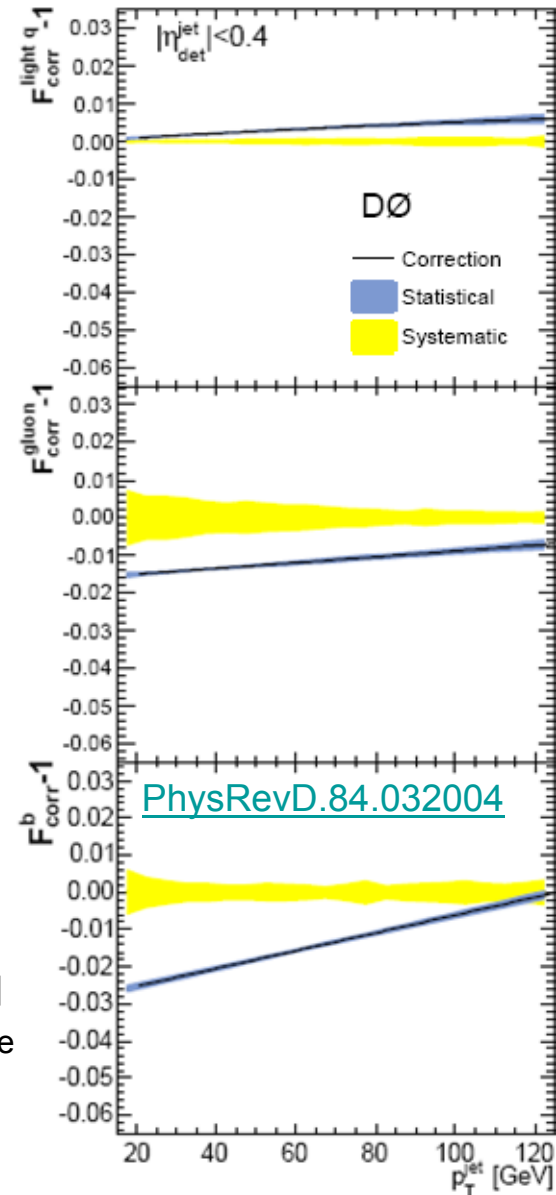
- D0 applies flavor dependent jet energy corrections to the MC to bring the response closer to that of the data and thus minimize the uncertainty.

$$D = \frac{\sum E_i \cdot R_i^{data}}{\sum E_i \cdot R_i^{MC}}$$

single particle response

$$F_{corr}^\beta = D_\beta / \langle D_{\gamma+jet} \rangle$$

JES for light jet mixture



Syst type / Values in GeV	Tevatron	
	CDF	D0
b-jet JES (response and model)	0.1 – 0.8	0.3 – 0.7

In addition, fragmentation and semilep BR uncertainties are considered (b-jet modelling)

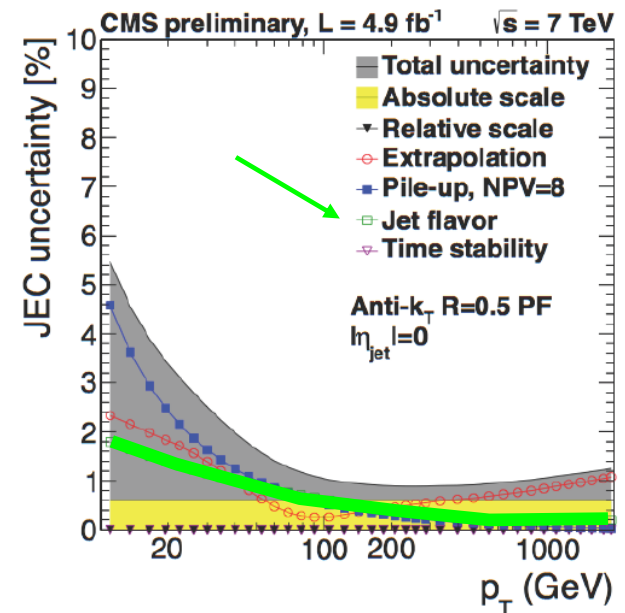
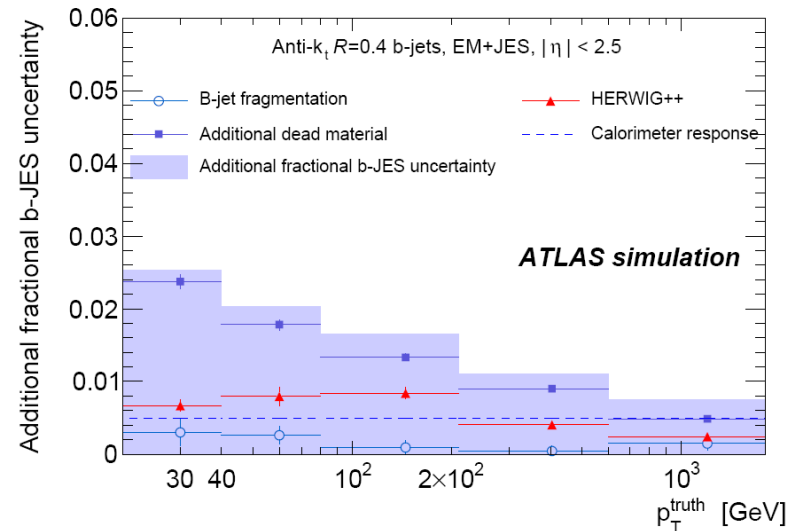
b- (or flavor-) specific JES at LHC

ATLAS:

- Additional b-JES uncertainty evaluated from:
 - calorimeter response (single particle)
 - b-jet fragmentation
 - Pythia/Herwig++
 - limited knowledge of the detector material

CMS:

- Use the max Pythia/Herwig response differences, for pure quark flavor at low p_T and for pure gluon flavor at high p_T , as a flavor uncertainty applicable to any jet flavor mixture.
- Detector effects are considered to be negligible due to the extensive use of tracking to reconstruct individual particles in the jet (Particle Flow)



Syst type / Values in GeV	Tevatron		LHC	
	CDF	D0	ATLAS	CMS
b-jet JES (response and model)	0.1 – 0.8	0.3 – 0.7	1.4 – 2.5	0.5 – 1.1

bJES checks with data?

■ $Z \rightarrow bb$ @ CDF used to check the b-jet energy scale. Very large background from QCD di-jet production.

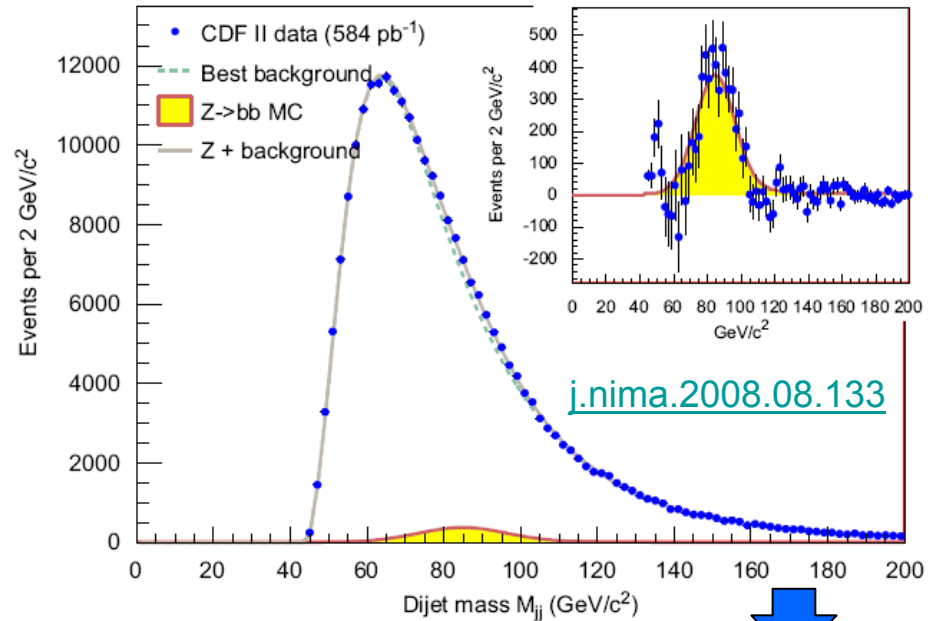
■ impractical @ LHC due to larger background cross section.

■ ATLAS used tracks in jets to validate the residual b-to-light-quark JES:

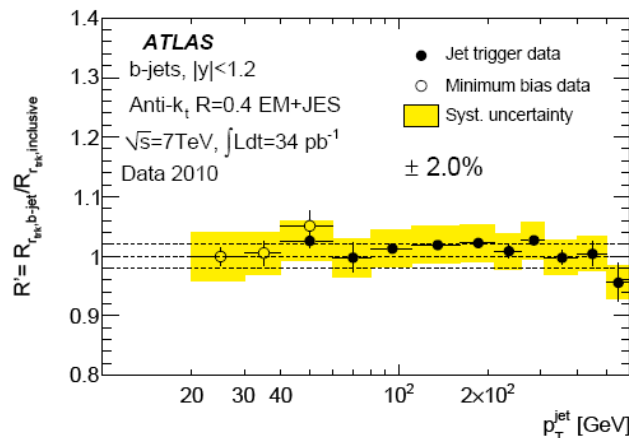
$$r_{trk} = \frac{\left| \sum \vec{p}_T^{track} \right|}{\left| \vec{p}_T^{jet} \right|}$$

$$R = \frac{\langle r_{trk} \rangle_{Data}}{\langle r_{trk} \rangle_{MC}}$$

$$R' = R^{b-jets} / R^{light-jets}$$



$$bJSF = 0.974 \pm 0.011(\text{stat})^{+0.017}_{-0.014}(\text{syst}) = 0.974^{+0.020}_{-0.018}(\text{total}).$$



It opens the possibility for an in-situ b-jet specific calibration!

Conclusions

- Experiments at Tevatron and LHC have been and are doing a great job in reducing the m_{top} uncertainty related to the physics objects and the detector description
 - This will continue in the future taking advantage of the large statistics available at the LHC.
 - In-situ constraint of the light jets energy scale helped in the mitigation of the JES (but also of the radiation) systematics on m_{top} .
 - Next steps will include beating down the flavor dependent JES uncertainty (b-JES), either by using data to constrain it, or (as done by D0) by employing flavour specific jet energy corrections.
- To reach and consolidate the sub-GeV m_{top} precision this must go hand in hand with the reduction of the theory uncertainties
 - Close collaboration between experiment and theory will be needed.
 - Reduce possible sources of double counting (in each systematic category try to separate out experimental and theoretical components)
 - Use available datasets to further tune MC generators
- Last but not least: make sure we understand what we are measuring
 - the m_{top} definition and interpretations ► [see talk by Sven-Olaf Moch](#)

Thanks for your attention!

- backup -

More info on the CMS ParticleFlow JES

CMS DP2012_006

