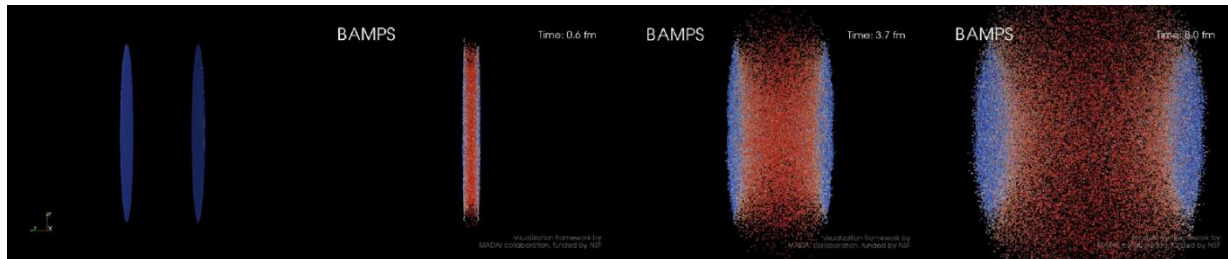


# D mesons and non-prompt J/Psi at LHC within a transport model

Jan Uphoff

with O. Fochler, Z. Xu and C. Greiner

Based on arXiv:1104.2295 and 1112.1559



# Motivation

Large heavy quark mass

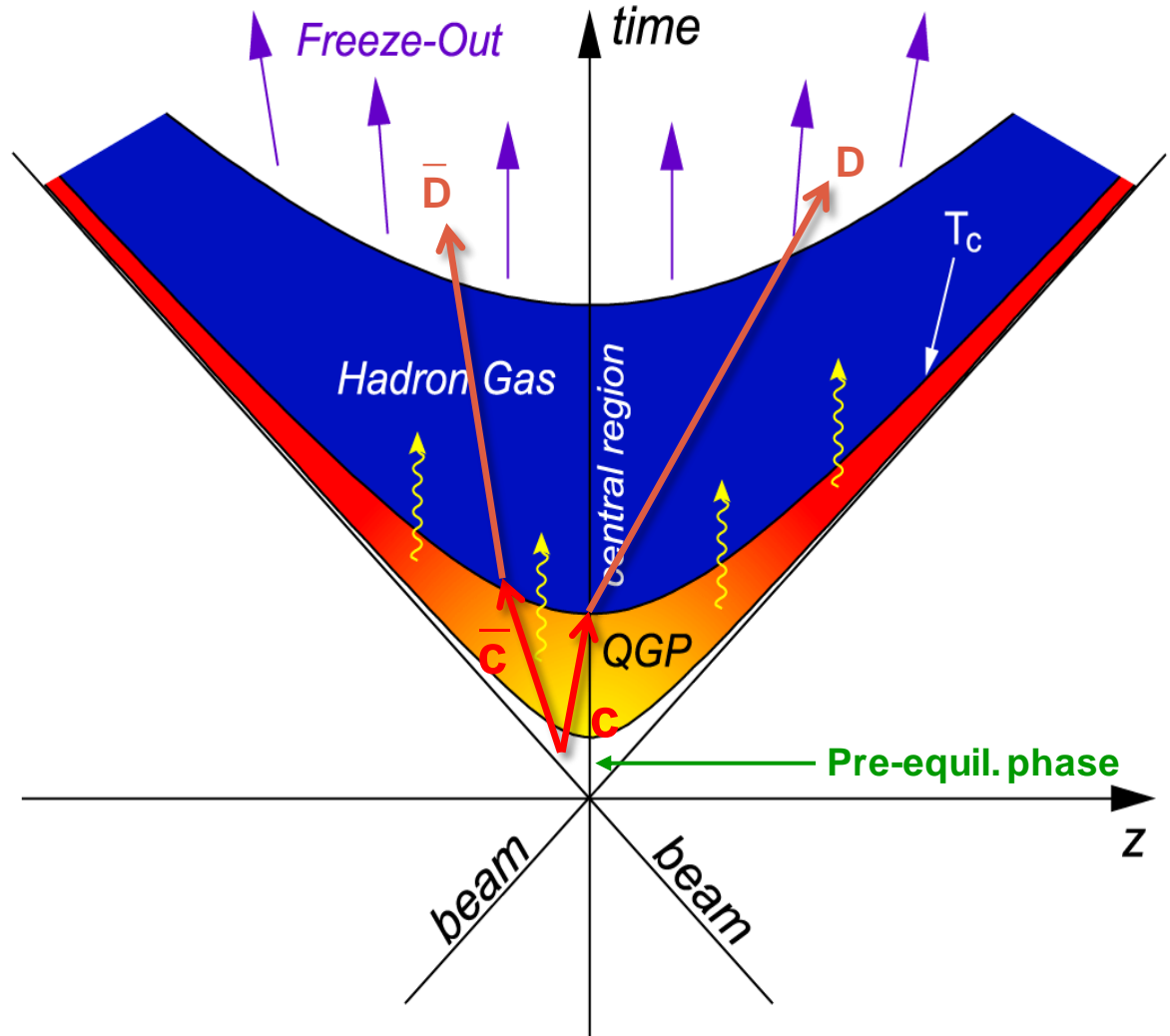
$$\gg \Lambda_{\text{QCD}}$$

Charm:  $M_c \approx 1.5 \text{ GeV}$

Bottom:  $M_b \approx 4.75 \text{ GeV}$

➔ Heavy quark production at early stage of collision

➔ ideal probe for this stage

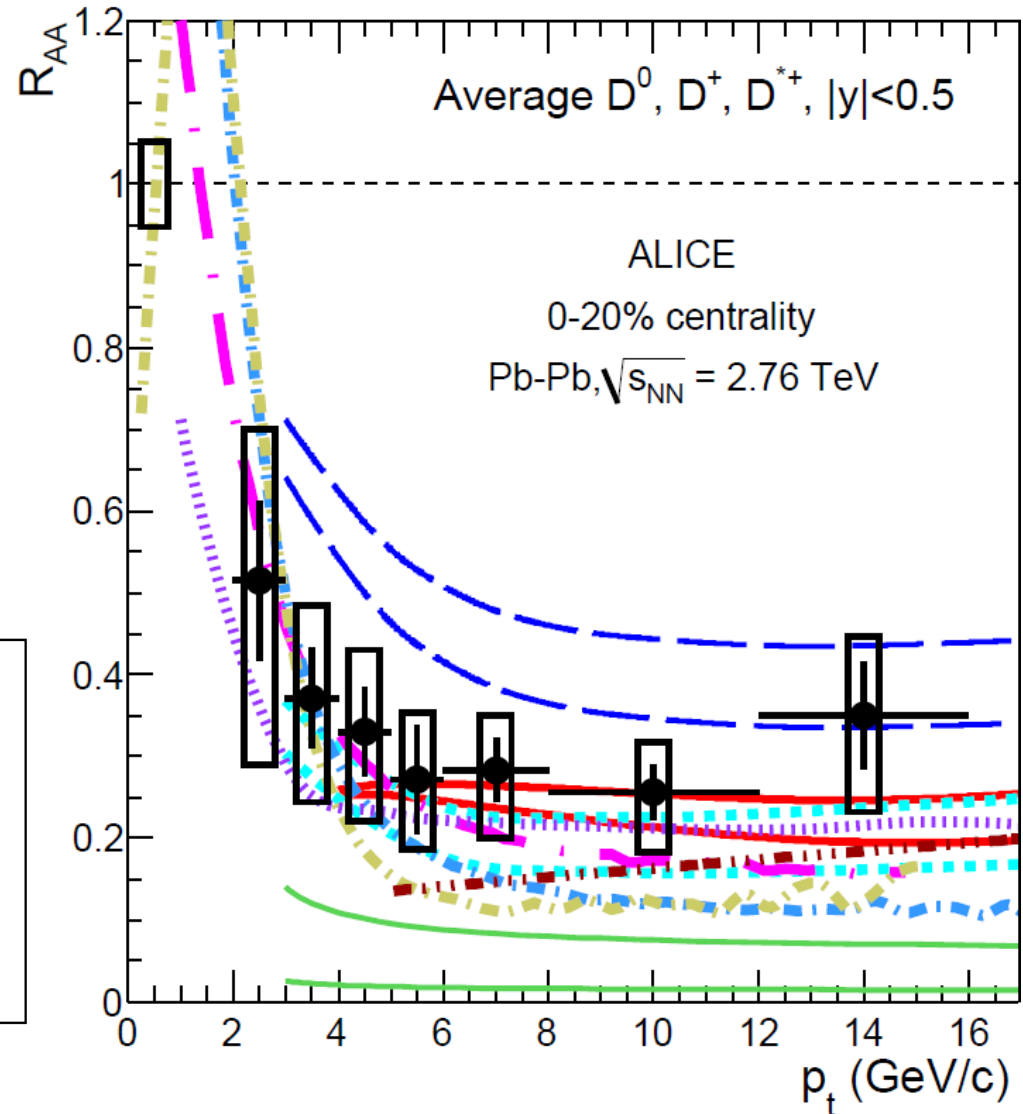


## Nuclear modification factor

$$R_{AA} = \frac{dN/dp_T dy|_{A+A}}{N_{\text{bin}} dN/dp_T dy|_{p+p}}$$

- Vitev rad (I)
- Vitev rad + dissoc (II)
- WHDG rad + coll (III)
- AdS/CFT Drag (IV)
- Langevin HTL2 (V)
- Coll + LPM rad (VI)
- BAMPS (VII)
- CUJET1.0 (VIII)
- BDMPS-ASW rad (IX)

ALICE, arXiv:1203.2160

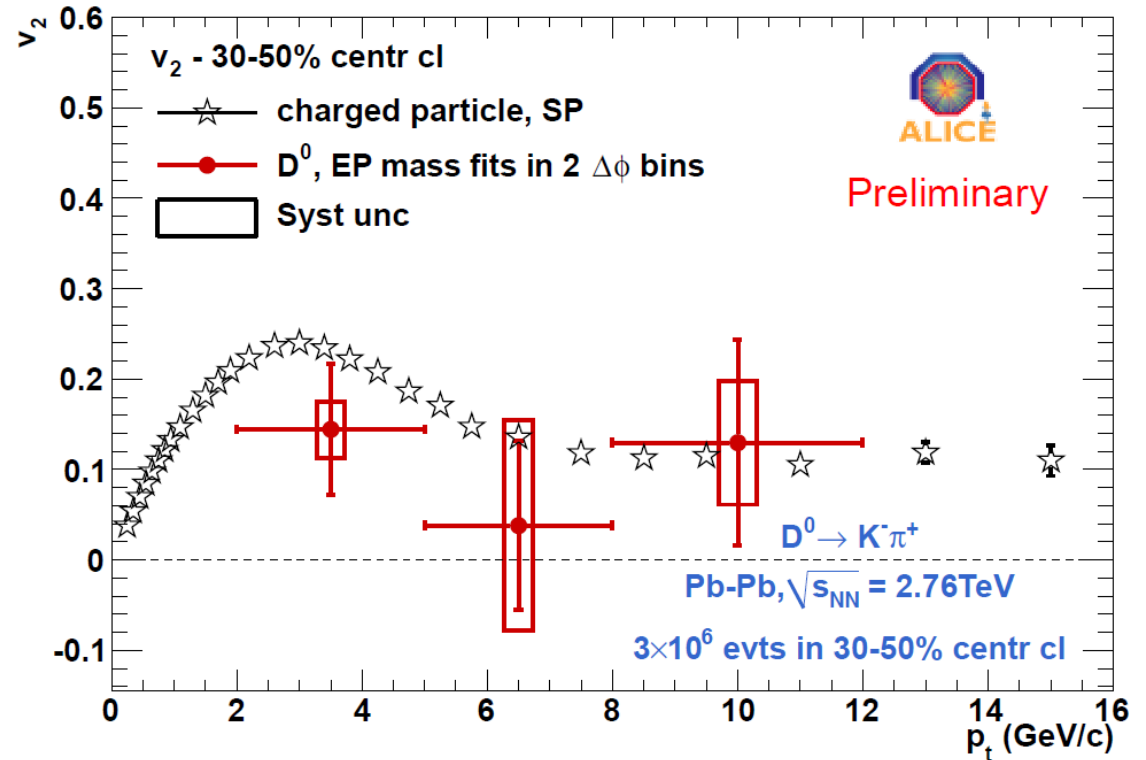
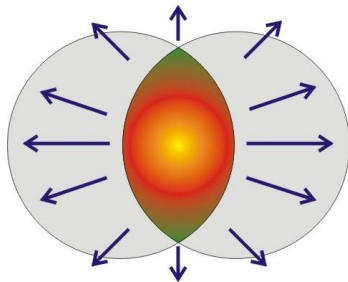
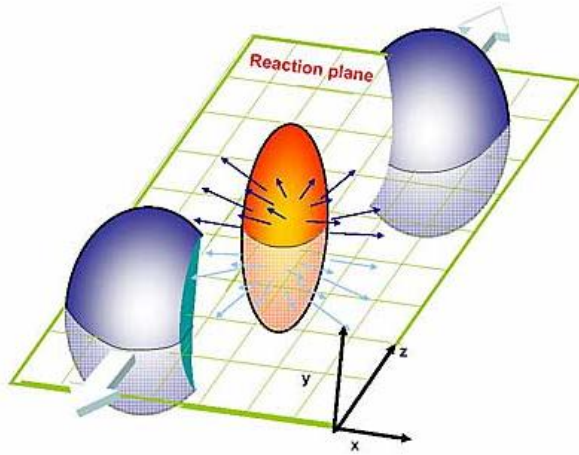


# Motivation

## Elliptic flow

$$v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$

$$\frac{d^3N}{p_T dp_T dy d\phi}(p_T, y, \phi) = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T dy} [1 + 2v_2(p_T, y) \cos(2\phi) + \dots]$$



## BAMPS: Boltzmann Approach of MultiParton Scatterings

- 3+1 dimensional, fully dynamic parton transport model
- solves the Boltzmann equations for on-shell partons with pQCD interactions

$$\left( \frac{\partial}{\partial t} + \frac{\mathbf{p}_i}{E_i} \frac{\partial}{\partial \mathbf{r}} \right) f_i(\mathbf{r}, \mathbf{p}_i, t) = \mathcal{C}_i^{2 \rightarrow 2} + \mathcal{C}_i^{2 \leftrightarrow 3} + \dots$$

Z. Xu & C. Greiner,  
Phys. Rev. C71 (2005)  
Phys. Rev. C76 (2007)

### Implemented processes:

$$g + g \rightarrow g + g$$

$$g + g \rightarrow g + g + g$$

$$g + g + g \rightarrow g + g$$

Light quarks have been implemented but are not included in all presented calculations

$$g + g \rightarrow Q + \bar{Q}$$

$$Q + \bar{Q} \rightarrow g + g$$

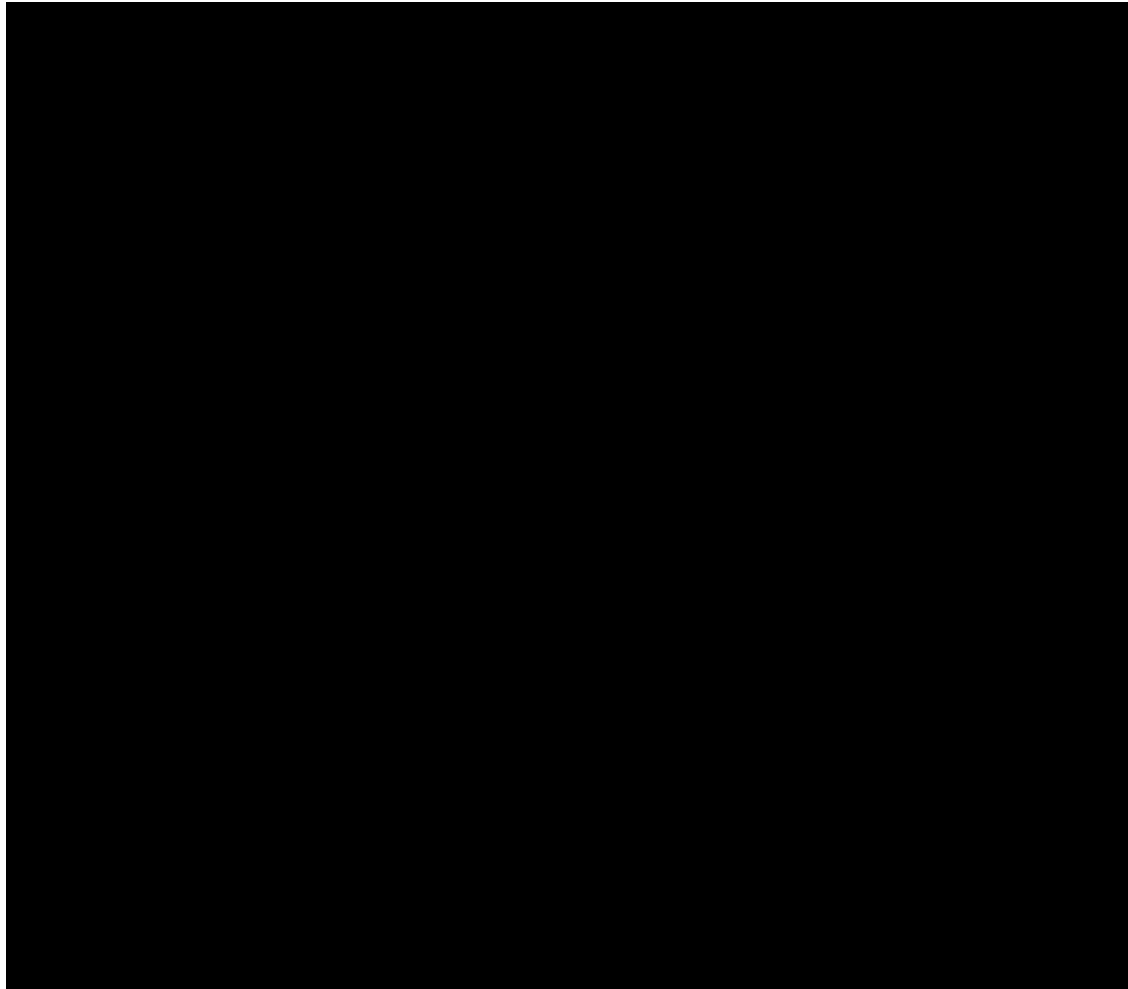
$$g + Q \rightarrow g + Q$$

$$g + \bar{Q} \rightarrow g + \bar{Q}$$

$$g + J/\psi \rightarrow c + \bar{c}$$

$$c + \bar{c} \rightarrow g + J/\psi$$

## BAMPS simulation of QGP phase at LHC at $\sqrt{s_{NN}} = 2.76$ TeV



Visualization framework  
courtesy MADAI  
collaboration, funded by  
the NSF under grant# NSF-  
PHY-09-41373

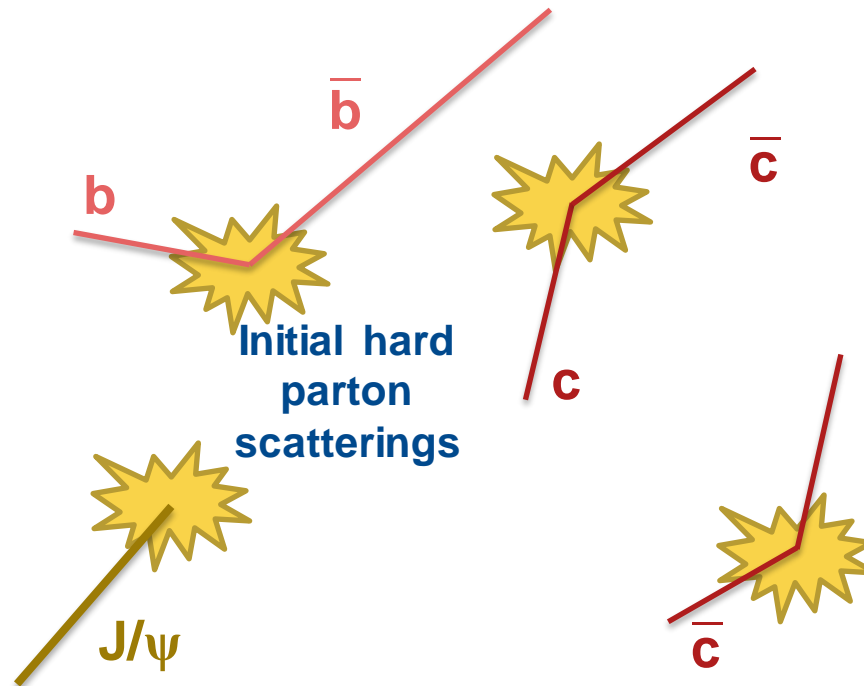
# Sketch of heavy-ion collision in BAMPS

## Heavy flavor in BAMPS



# Sketch of heavy-ion collision in BAMPS

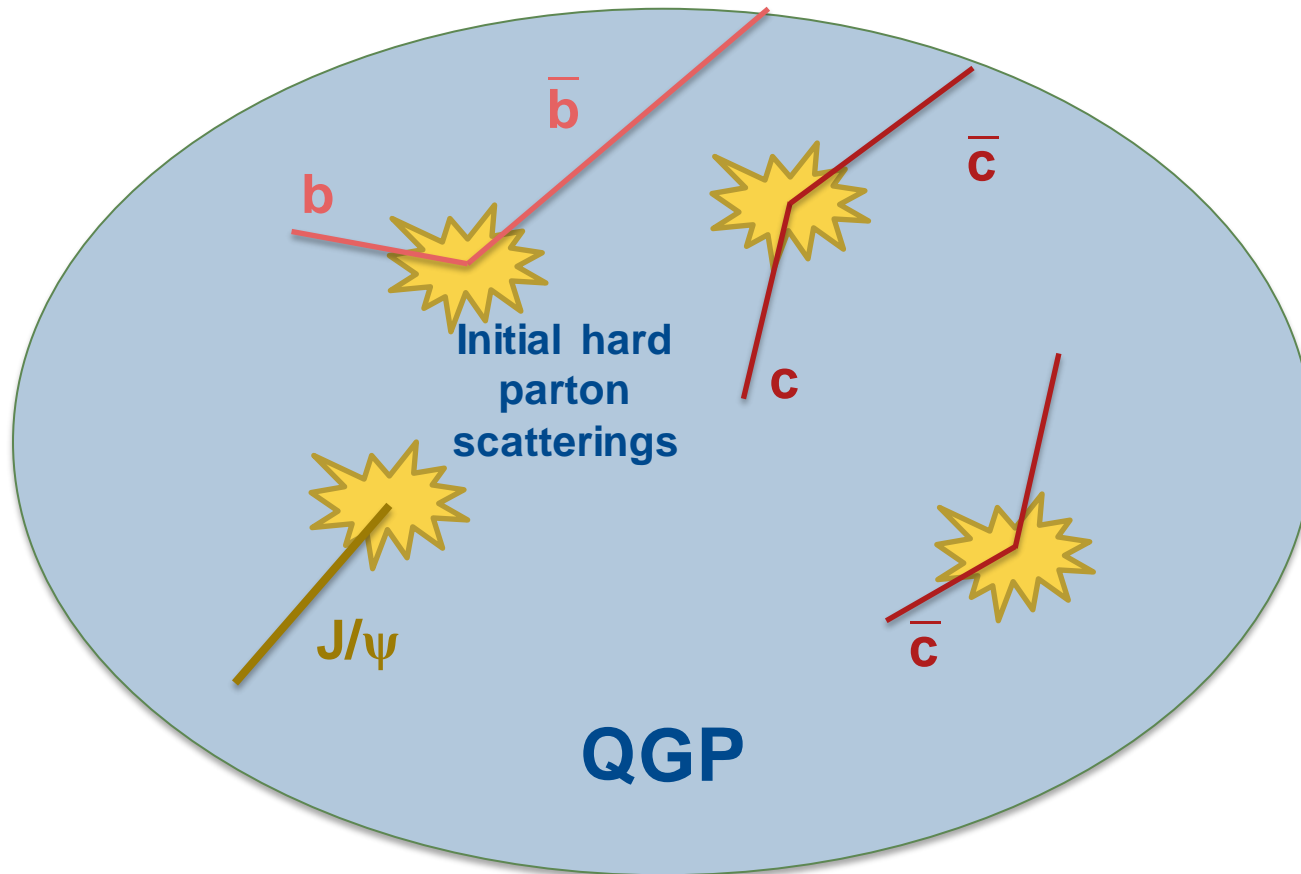
## Heavy flavor in BAMPS





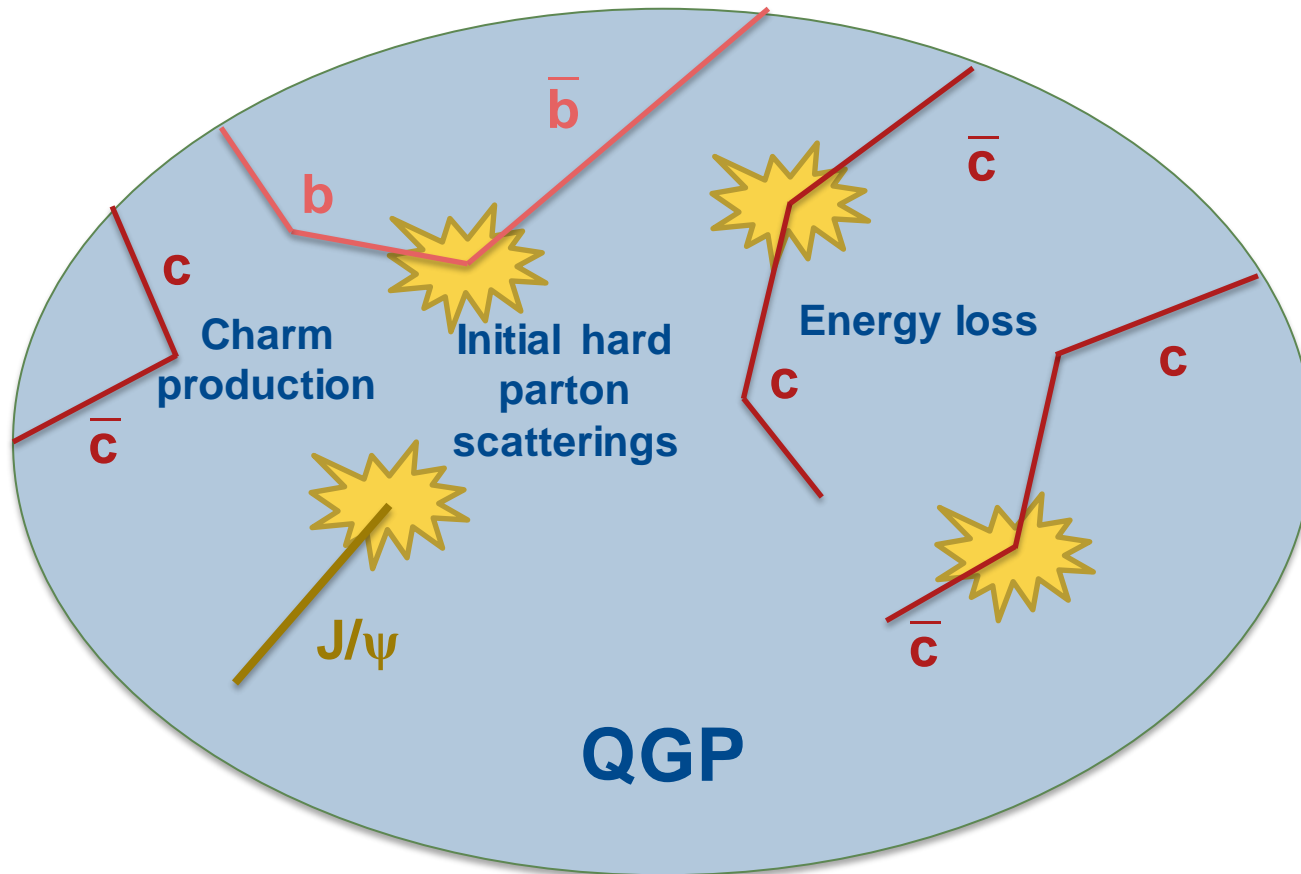
# Sketch of heavy-ion collision in BAMPS

## Heavy flavor in BAMPS



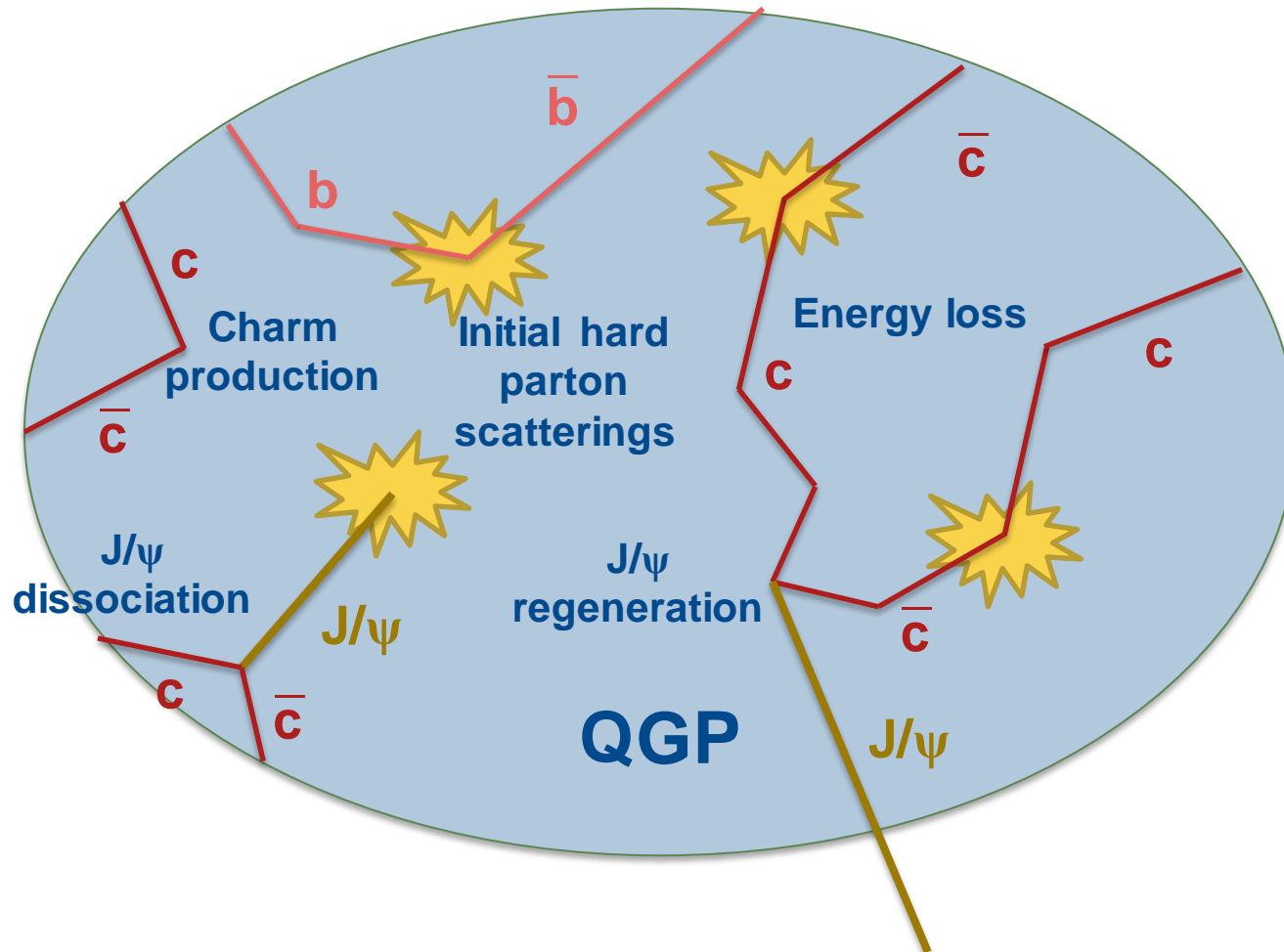
# Sketch of heavy-ion collision in BAMPS

## Heavy flavor in BAMPS



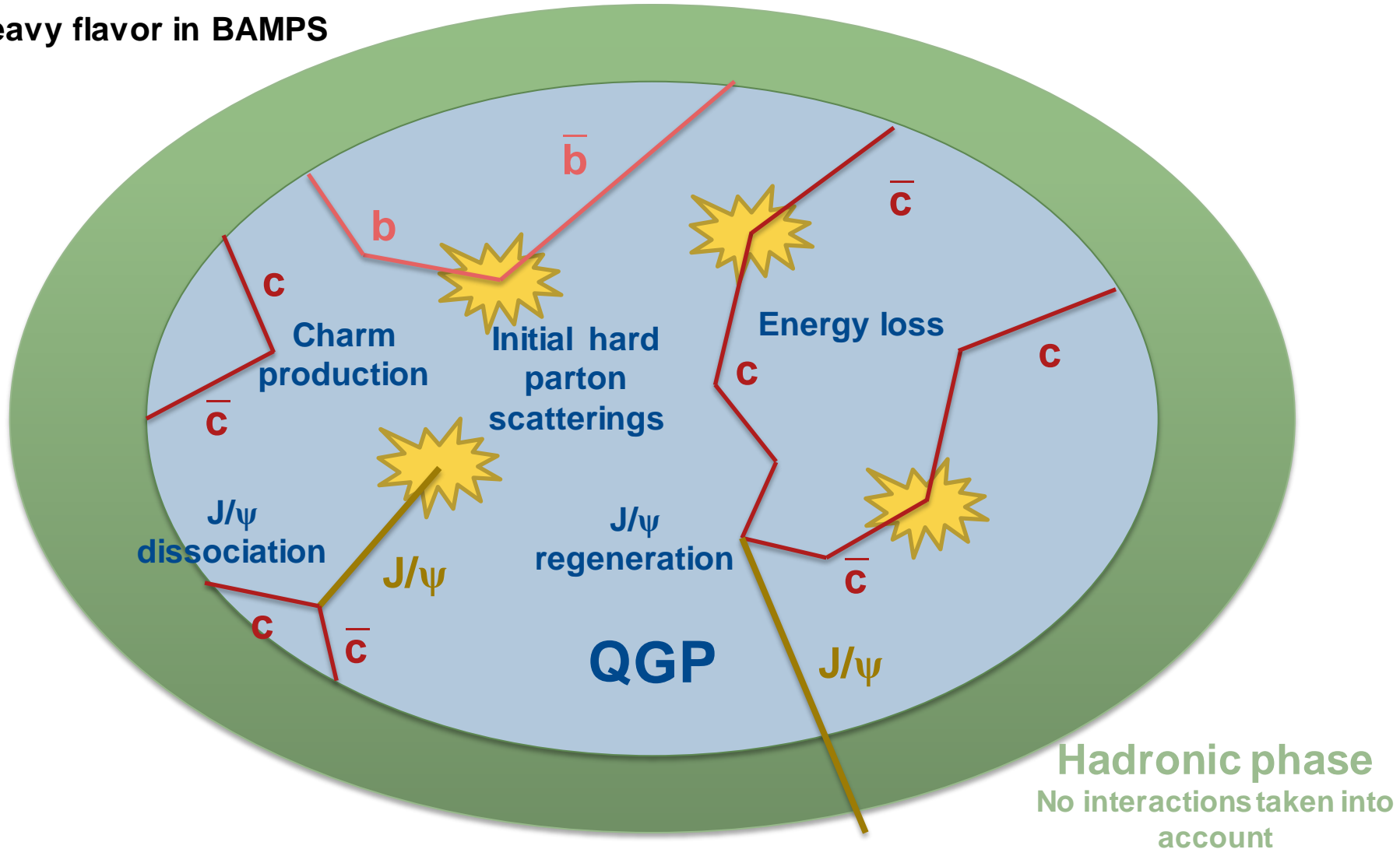
# Sketch of heavy-ion collision in BAMPS

## Heavy flavor in BAMPS



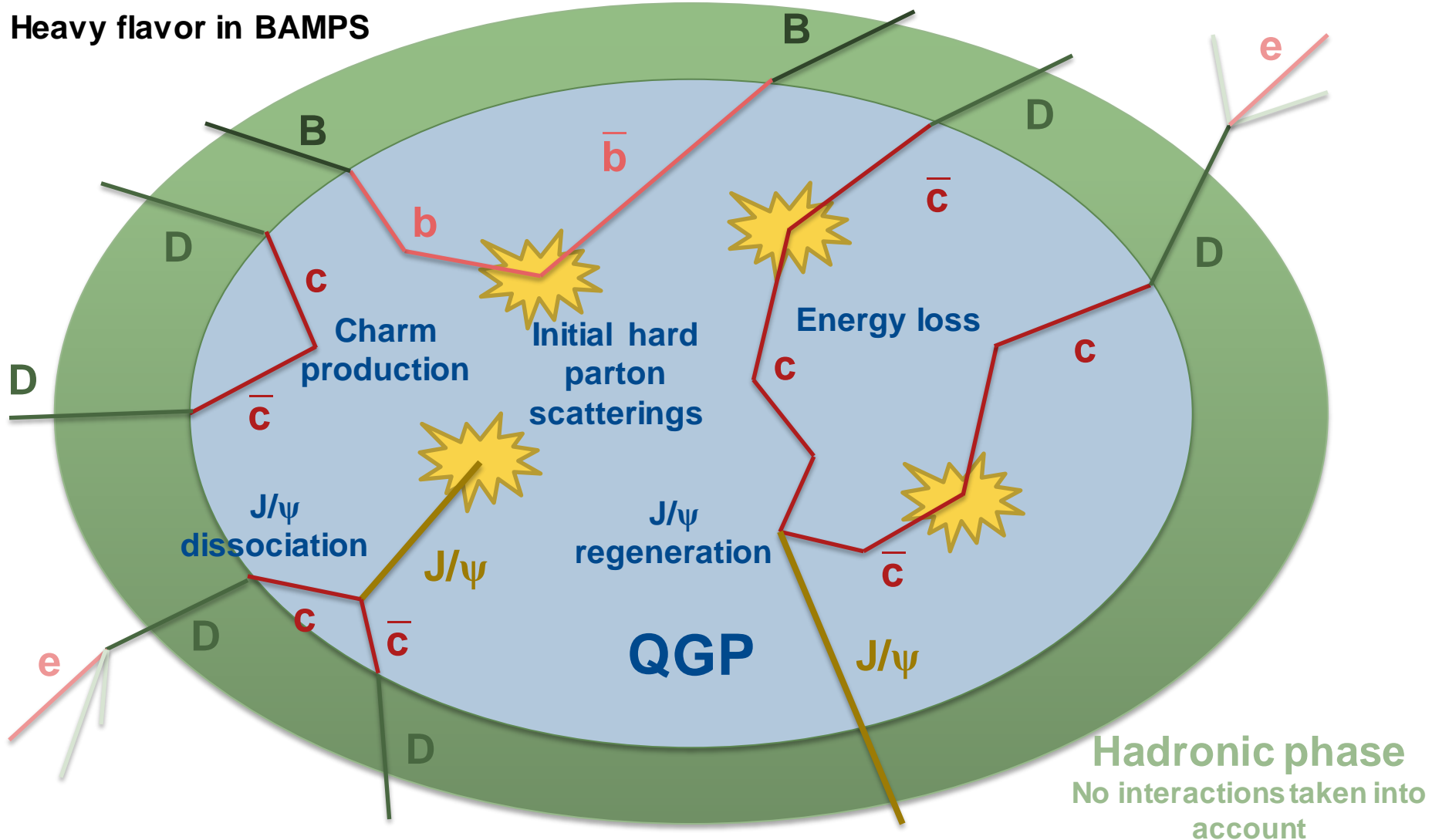
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# Sketch of heavy-ion collision in BAMPS

## Heavy flavor in BAMPS

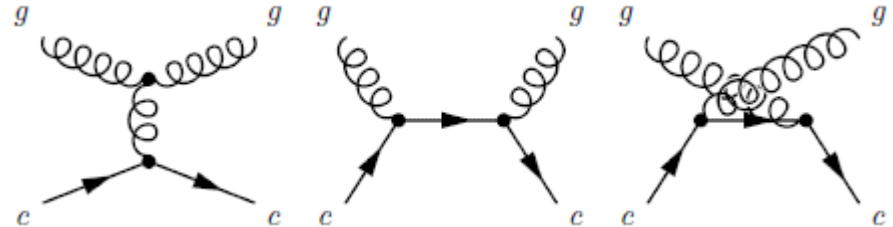


# Heavy quark scattering

Leading order perturbative QCD:

$$g + Q \rightarrow g + Q$$

$$g + \bar{Q} \rightarrow g + \bar{Q}$$



**t channel is divergent for small t**

$$\frac{1}{t} \rightarrow \frac{1}{t - \kappa m_D^2}$$

$\kappa$  can be fixed to

$$\kappa = \frac{1}{2e} \approx 0.184$$

by comparing  $dE/dx$  to  
HTL result beyond  
logarithmic accuracy

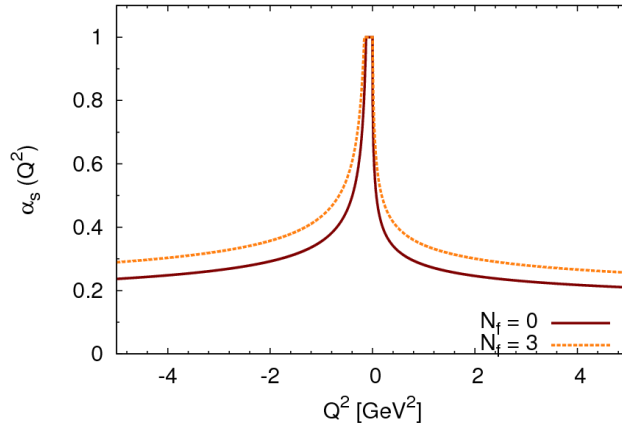
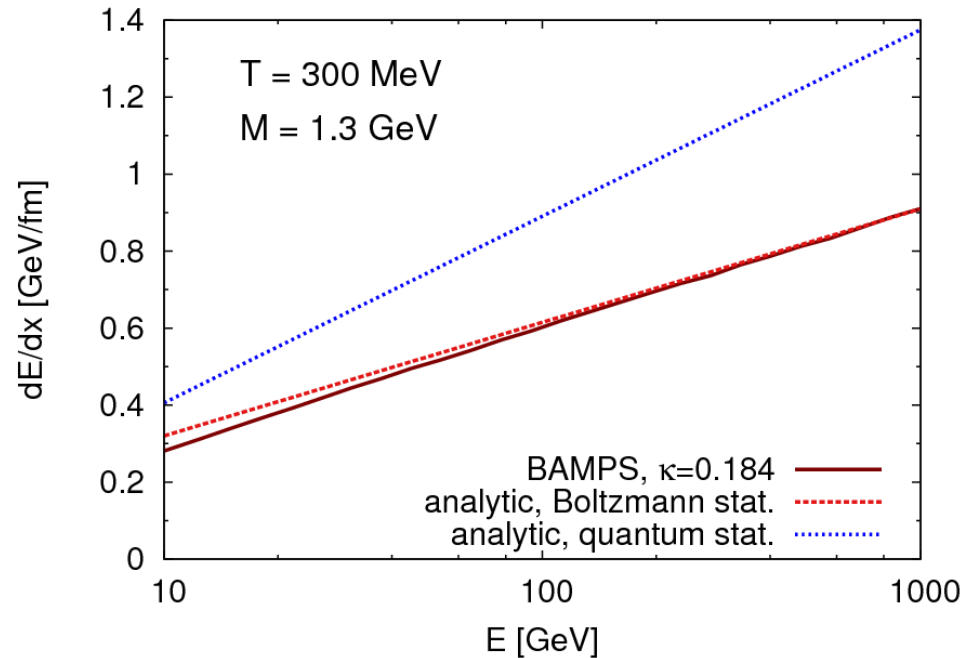
A. Peshier,  
arXiv:0801.0595  
[hep-ph]

P.B. Gossiaux,  
J. Aichelin,  
Phys.Rev.C78 (2008)

# Heavy quark scattering

## Compare to analytic formula

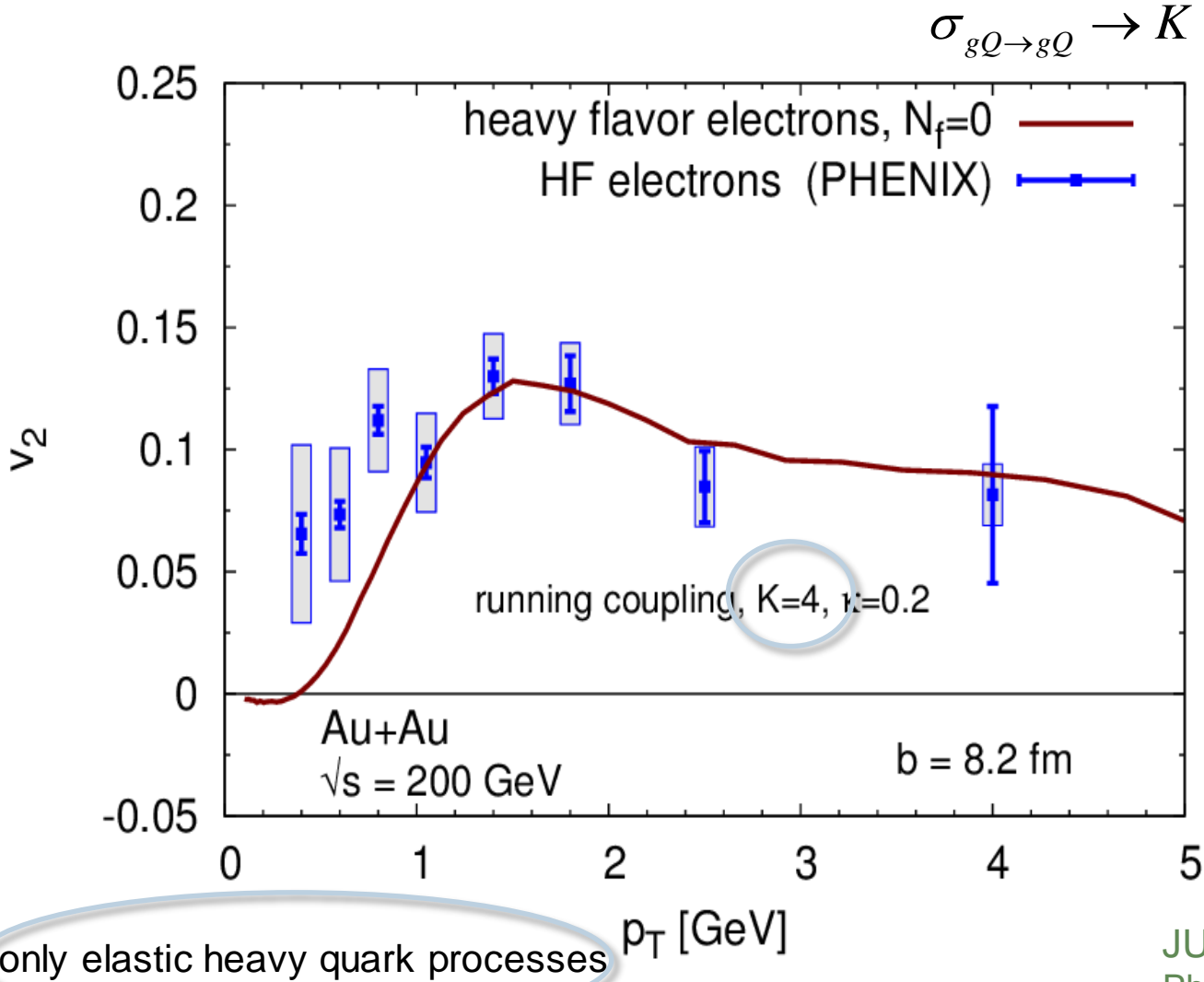
$$\frac{dE}{dx} = \frac{8\alpha_s^2 T^2}{\pi} \left[ \left(1 + \frac{n_f}{3}\right) \ln \frac{ET}{m_D^2} + \frac{2}{9} \ln \frac{ET}{M^2} + \left(\ln 2 - \frac{1}{4} - \frac{\gamma}{3}\right) n_f + \frac{31}{9} \ln 2 - \frac{101}{108} - \frac{11\gamma}{9} \right]$$



**Introduce a running coupling constant for all channels**

P.B. Gossiaux, J. Aichelin, Phys.Rev.C78 (2008)

# Heavy quark elliptic flow $v_2$ at RHIC



RHIC

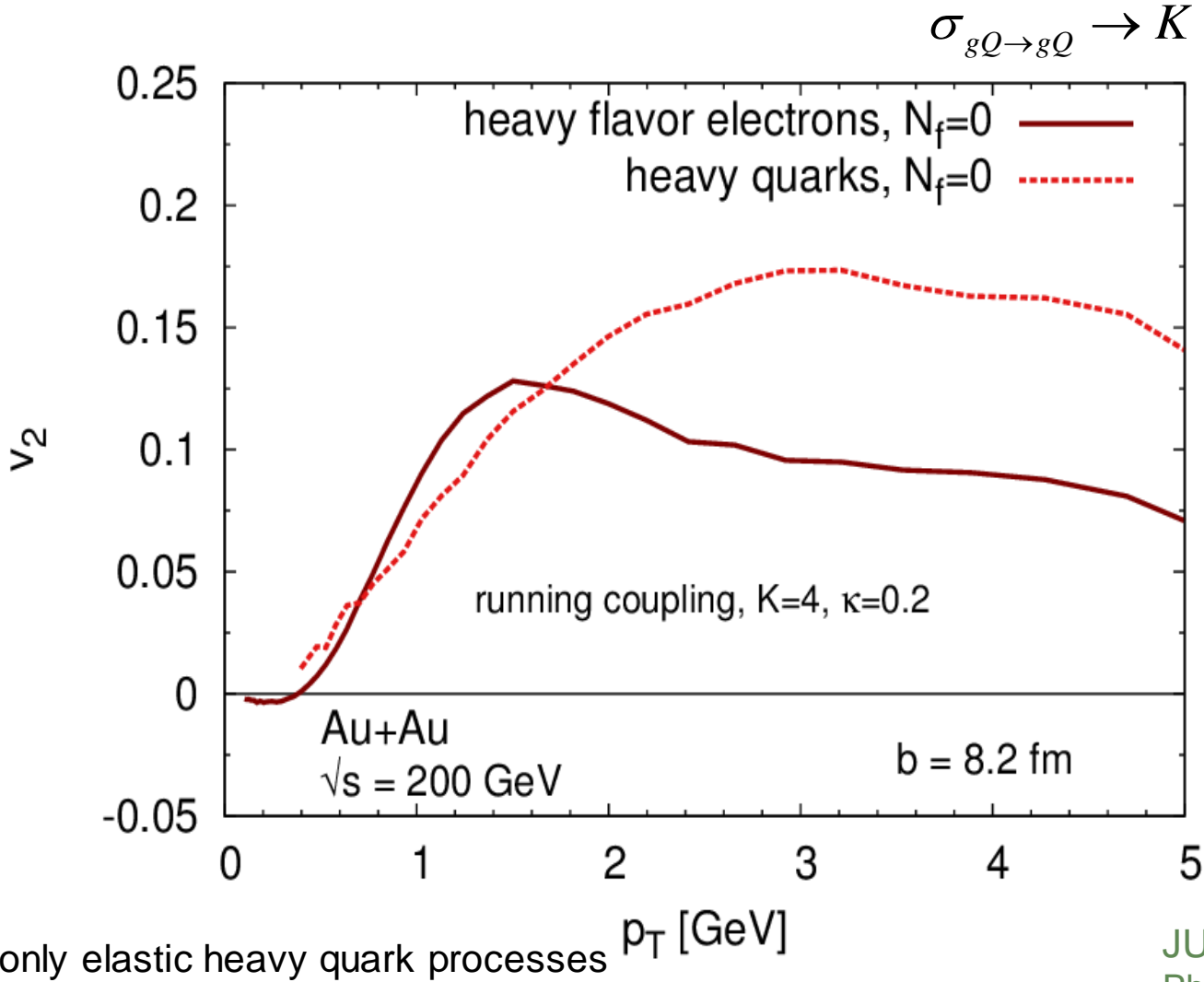
- What is missing:
- Radiative contributions
  - Quantum statistics

PHENIX,  
arXiv:1005.1627

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)



# Heavy quark elliptic flow $v_2$ at RHIC



RHIC

- What is missing:
- Radiative contributions
  - Quantum statistics

PHENIX,  
arXiv:1005.1627

only elastic heavy quark processes

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

# BAMPS with $N_{\text{flavor}} = 3+2$

## Implemented processes

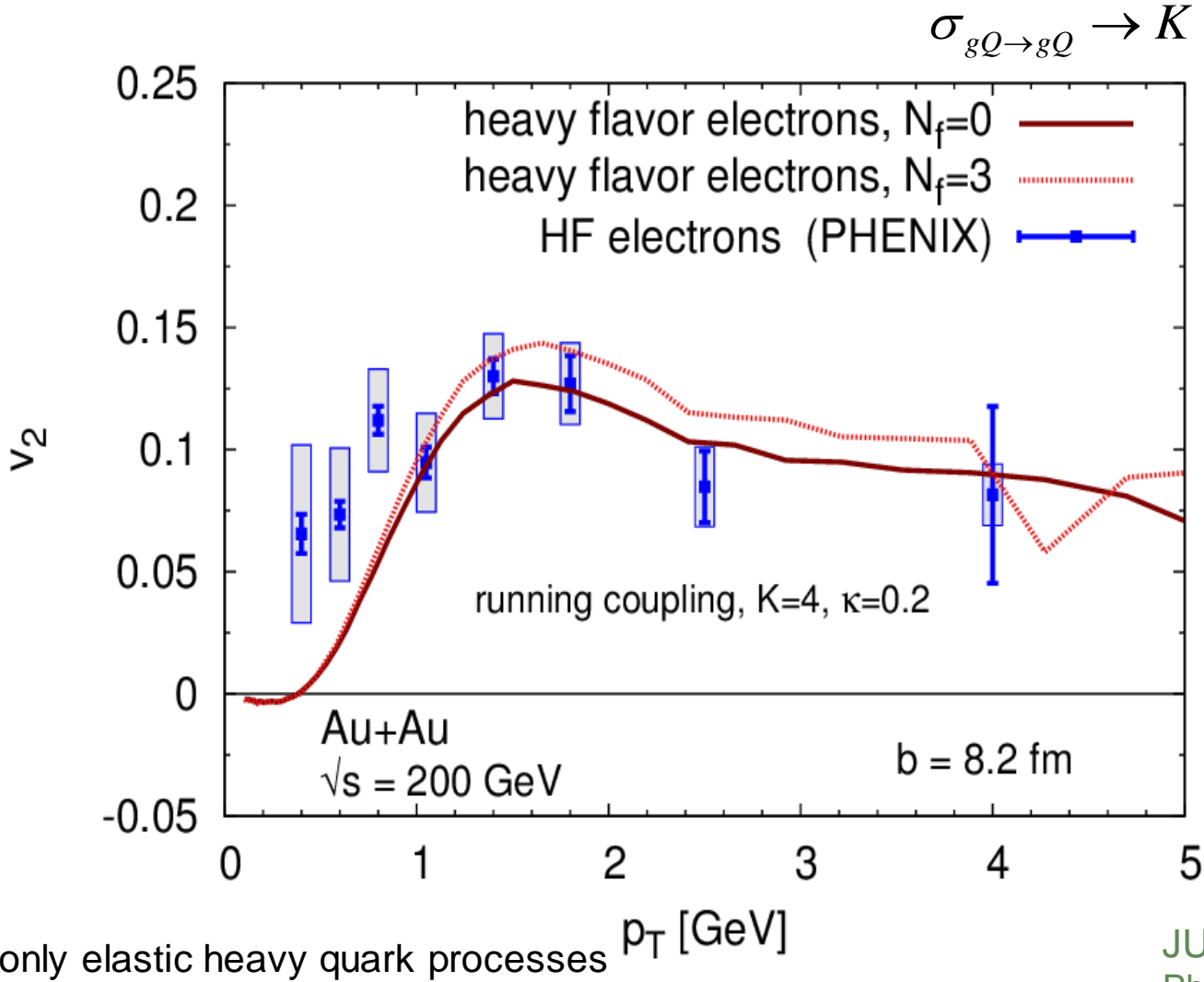
### Heavy Flavor

$$\begin{aligned}
 g + g &\rightarrow Q + \bar{Q} \\
 Q + \bar{Q} &\rightarrow g + g \\
 q + \bar{q} &\rightarrow Q + \bar{Q} \\
 Q + \bar{Q} &\rightarrow q + \bar{q} \\
 g + Q &\rightarrow g + Q \\
 g + \bar{Q} &\rightarrow g + \bar{Q} \\
 q + Q &\rightarrow q + Q \\
 q + \bar{Q} &\rightarrow q + \bar{Q} \\
 g + J/\psi &\rightarrow c + \bar{c} \\
 c + \bar{c} &\rightarrow g + J/\psi
 \end{aligned}$$

$$\begin{aligned}
 g g &\rightarrow g g && \mathbf{2 \rightarrow 2} \\
 g g &\rightarrow q \bar{q} \\
 q \bar{q} &\rightarrow g g && \text{and } q \bar{q} \rightarrow q' \bar{q}' \\
 q g &\rightarrow q g && \text{and } \bar{q} g \rightarrow \bar{q} g \\
 q \bar{q} &\rightarrow q \bar{q} \\
 q q &\rightarrow q q && \text{and } \bar{q} \bar{q} \rightarrow \bar{q} \bar{q} \\
 q q' &\rightarrow q q' && \text{and } q \bar{q}' \rightarrow q \bar{q}'
 \end{aligned}$$

$$\begin{aligned}
 g g &\leftrightarrow g g g && \mathbf{2 \leftrightarrow 3} \\
 q g &\leftrightarrow q g g && \text{and } \bar{q} g \leftrightarrow \bar{q} g g \\
 q \bar{q} &\leftrightarrow q \bar{q} g \\
 q q &\leftrightarrow q q g && \text{and } \bar{q} \bar{q} \leftrightarrow \bar{q} \bar{q} g \\
 q q' &\leftrightarrow q q' g && \text{and } q \bar{q}' \leftrightarrow q \bar{q}' g
 \end{aligned}$$

# Heavy quark elliptic flow $v_2$ at RHIC



RHIC

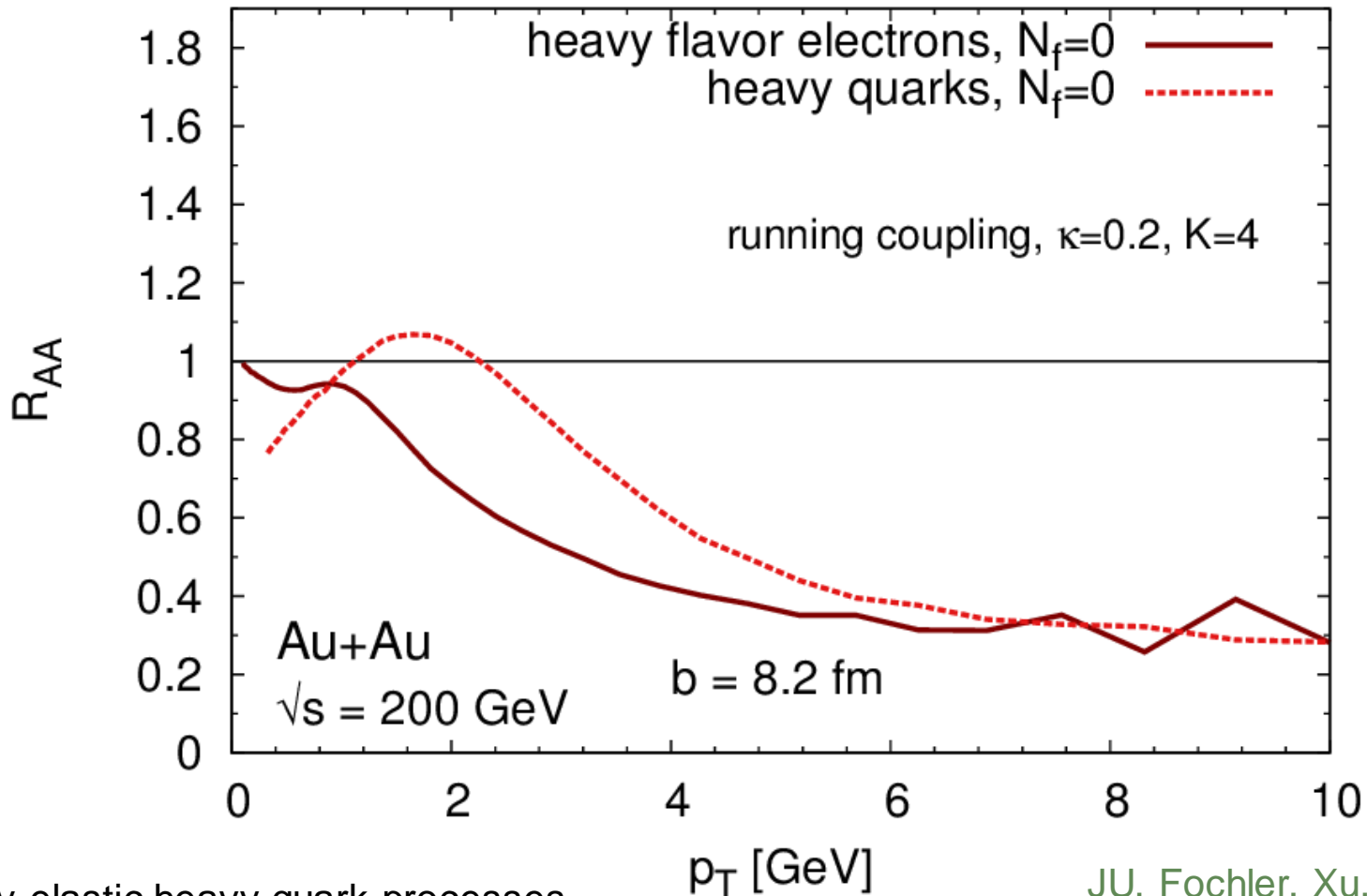
- What is missing:
- Radiative contributions
  - Quantum statistics

PHENIX,  
arXiv:1005.1627

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

# Heavy quark $R_{AA}$ at RHIC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



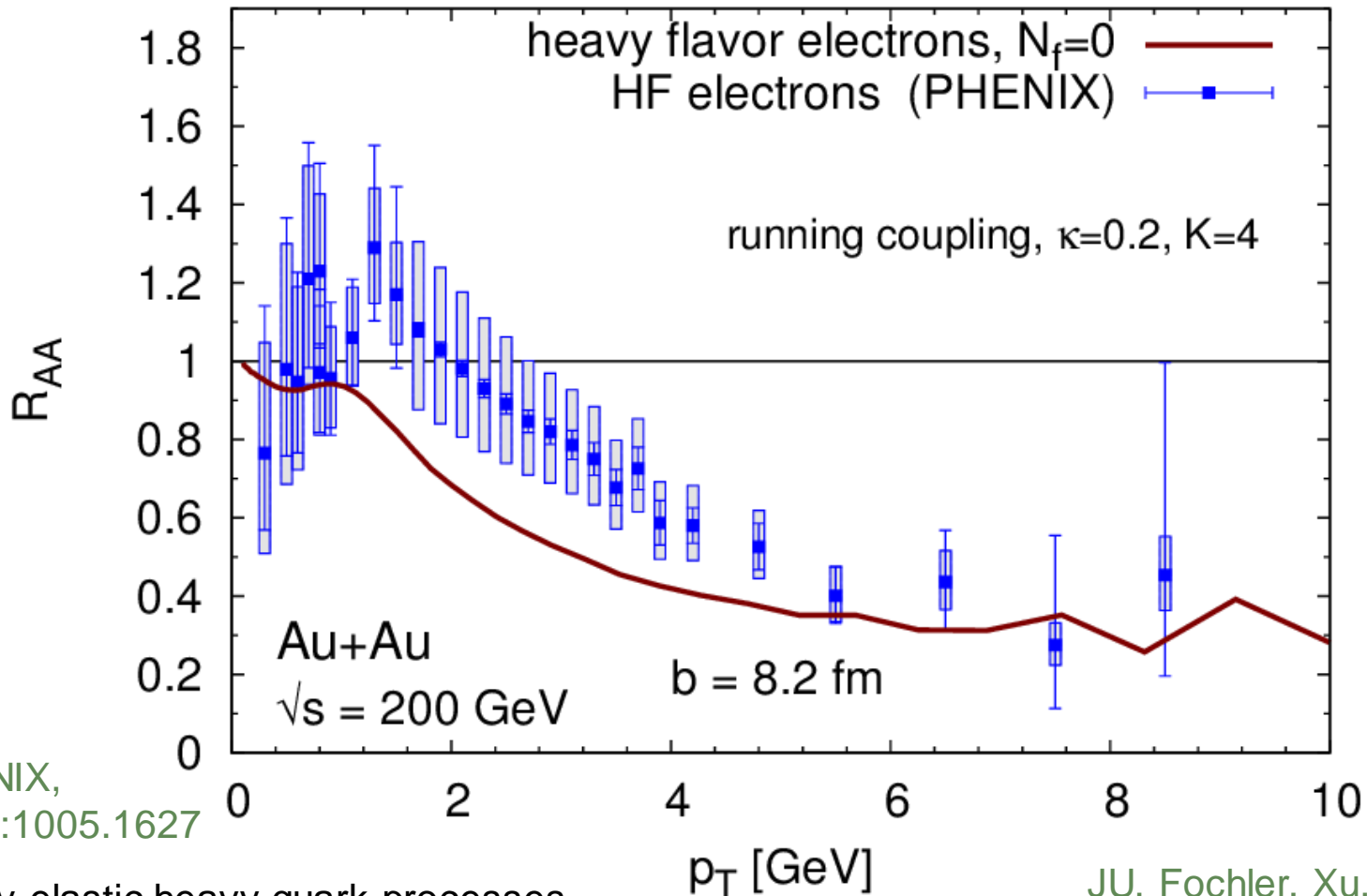
RHIC

only elastic heavy quark processes

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

# Heavy quark $R_{AA}$ at RHIC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



RHIC

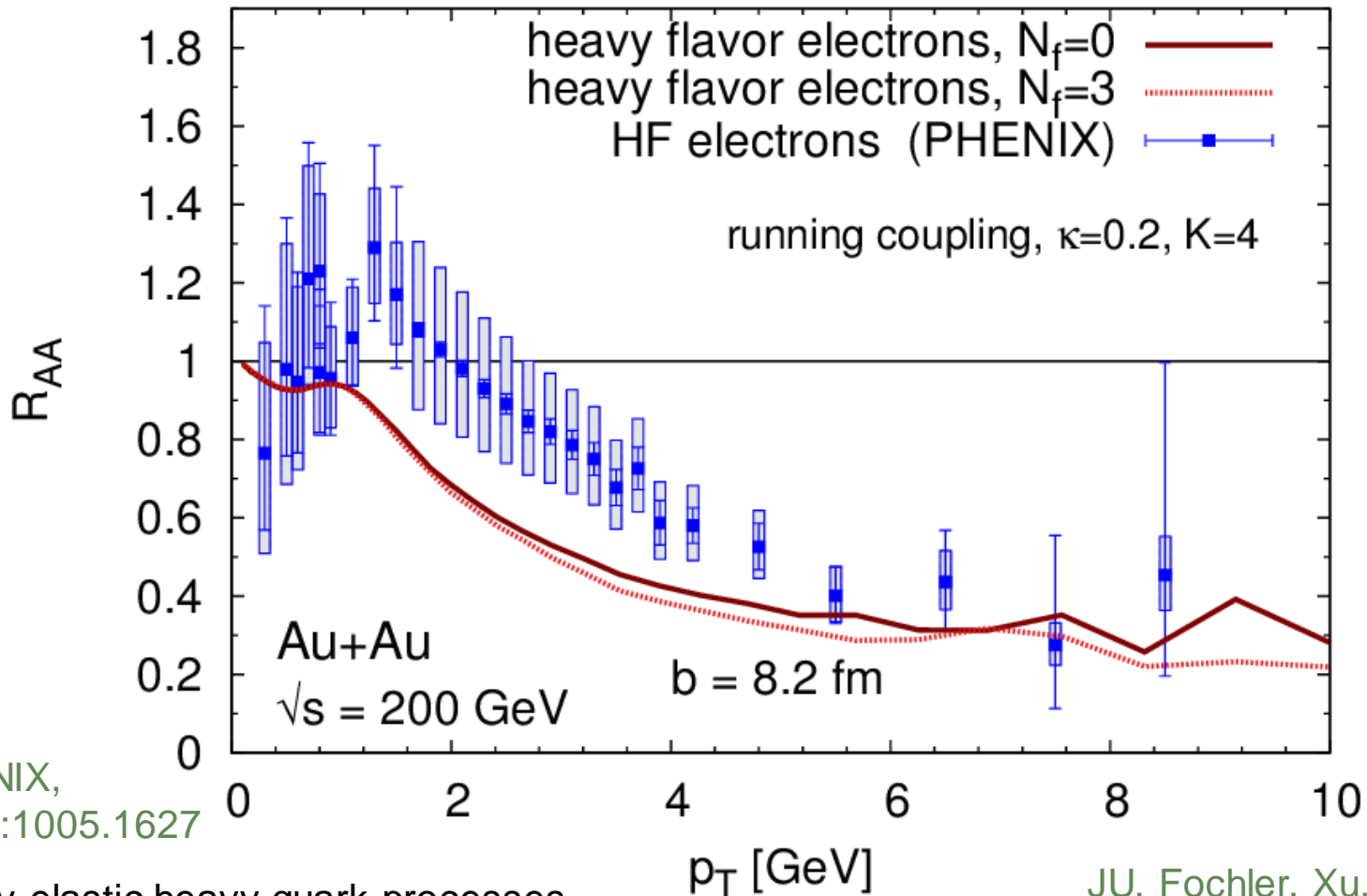
PHENIX,  
arXiv:1005.1627

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JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

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$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



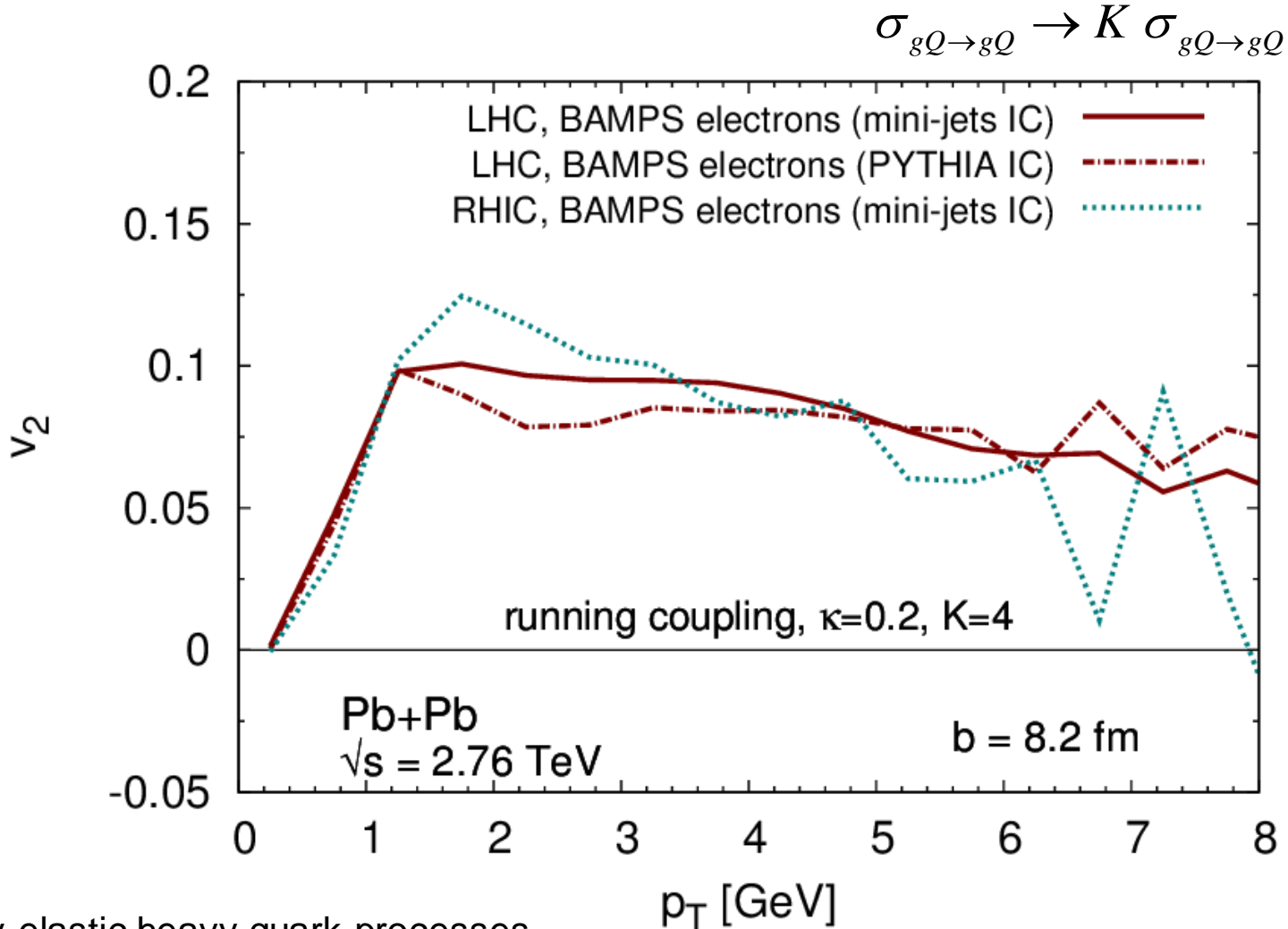
RHIC

PHENIX,  
arXiv:1005.1627

only elastic heavy quark processes

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

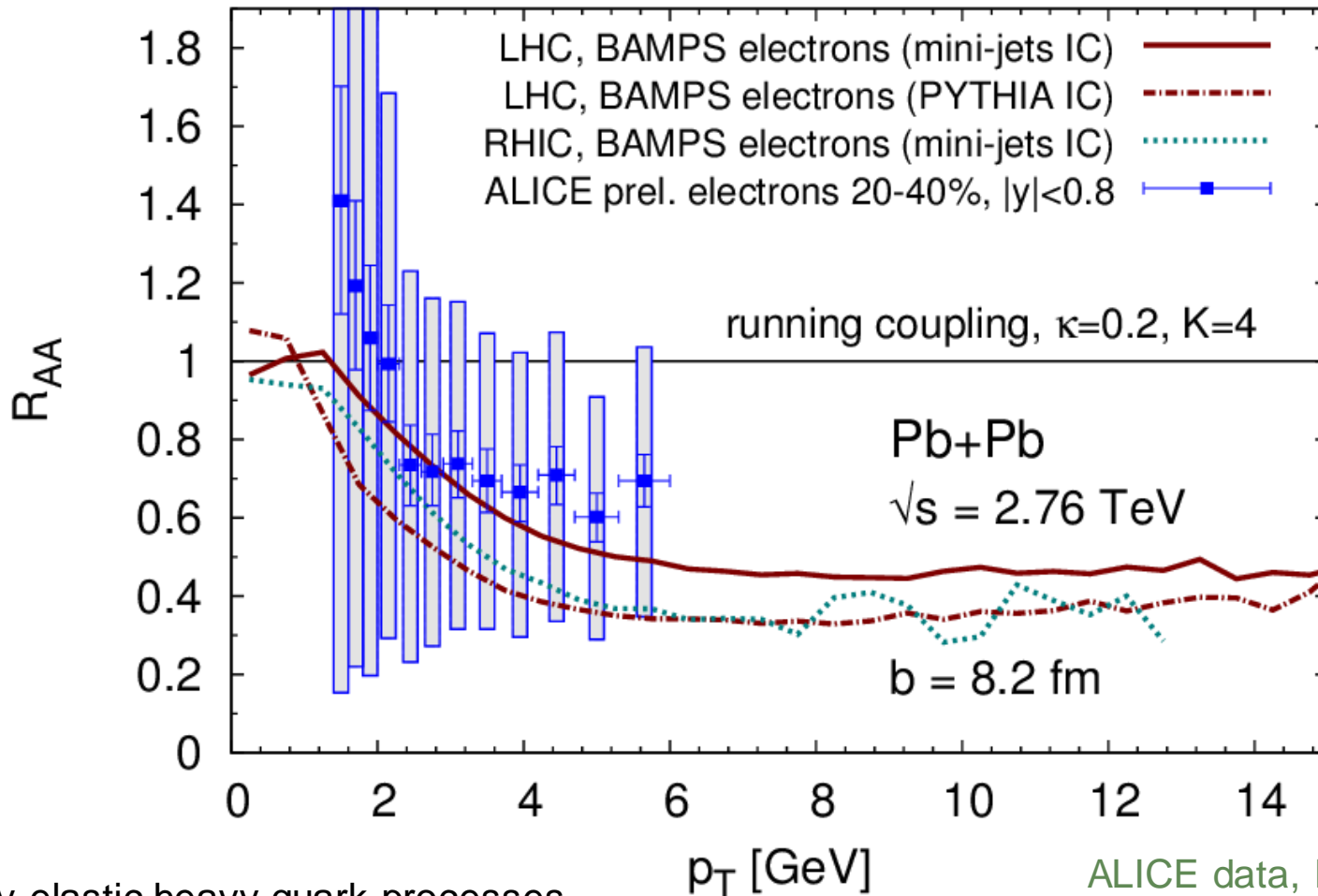
# Heavy quark $v_2$ at LHC



only elastic heavy quark processes

# Heavy quark $R_{AA}$ at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



LHC

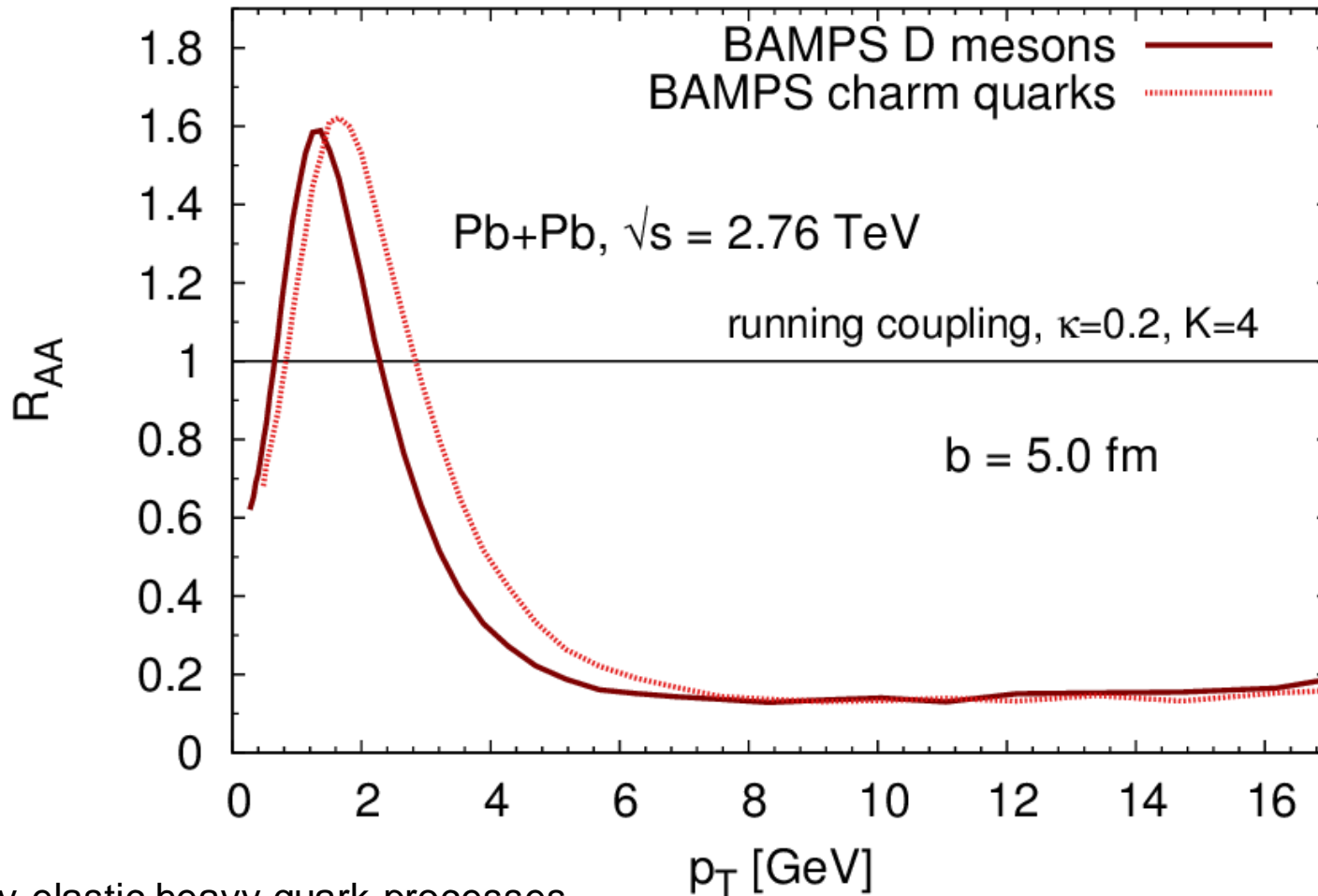
only elastic heavy quark processes

ALICE data, Pachmayer,  
QM 2011



# D meson $R_{AA}$ at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$

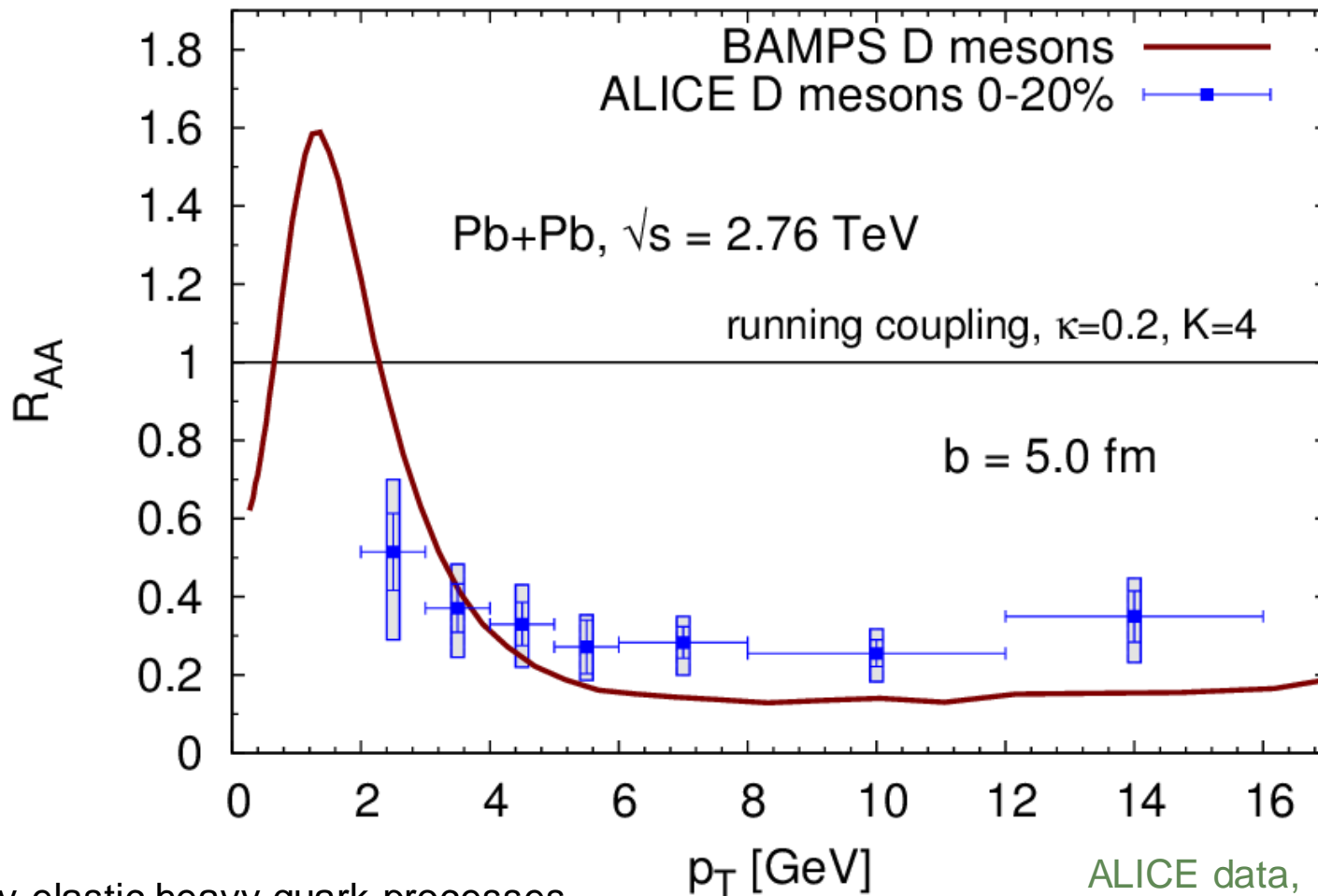


LHC

only elastic heavy quark processes

# D meson $R_{AA}$ at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$

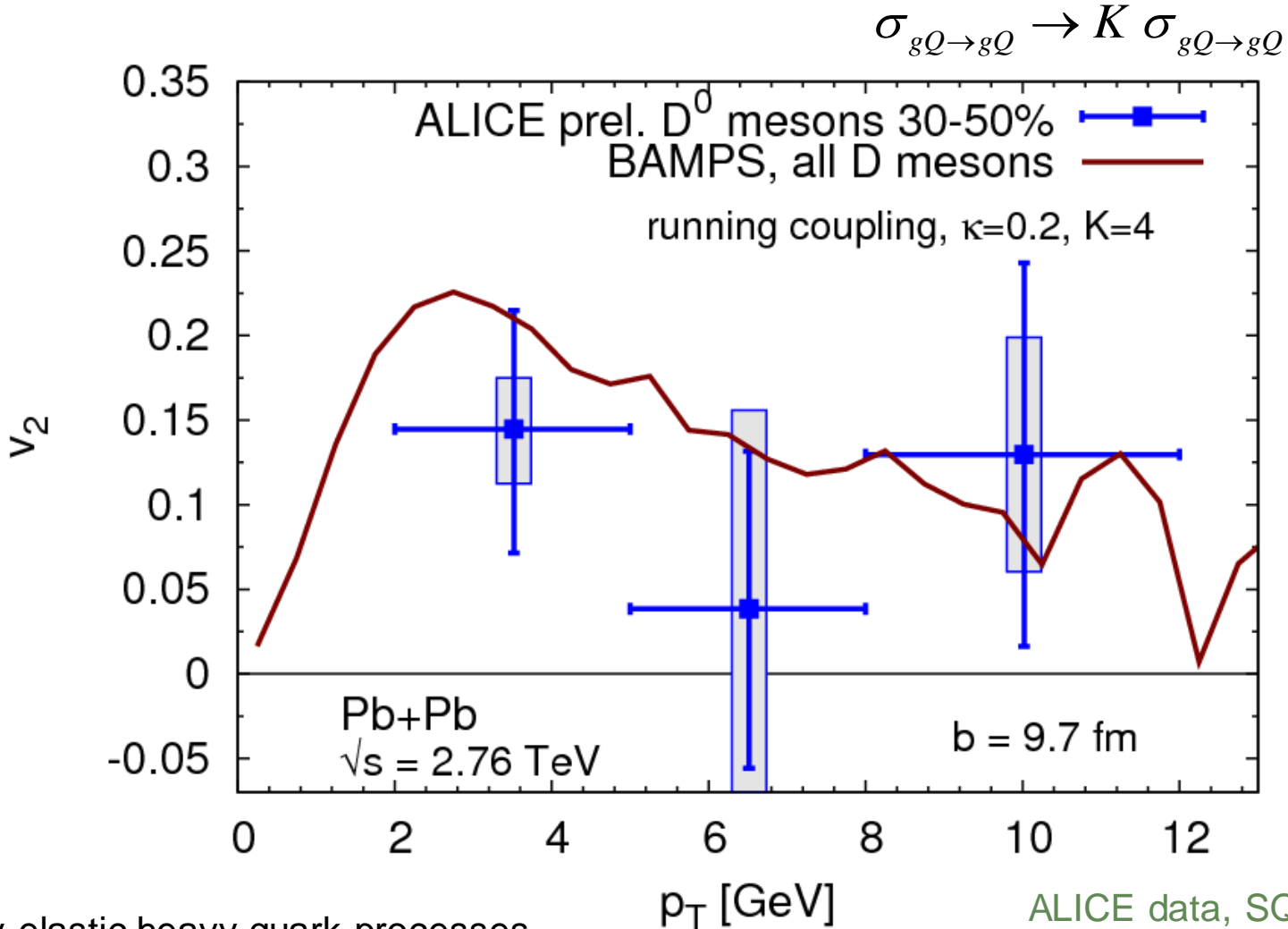


LHC

only elastic heavy quark processes

ALICE data,  
arXiv:1203.2160

# D meson $v_2$ at LHC

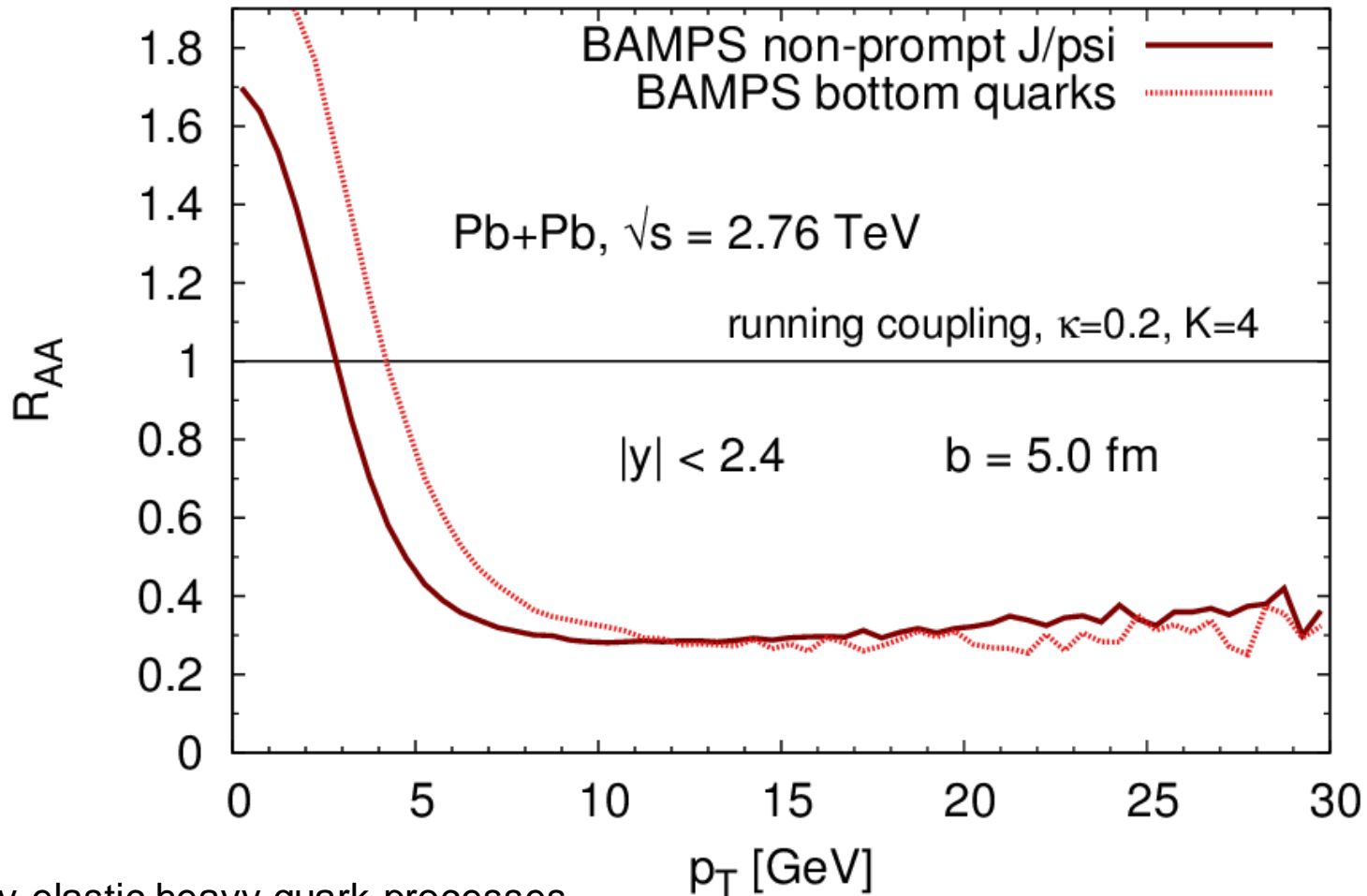


only elastic heavy quark processes

ALICE data, SQM 2011

# Non-prompt J/psi $R_{AA}$ at LHC

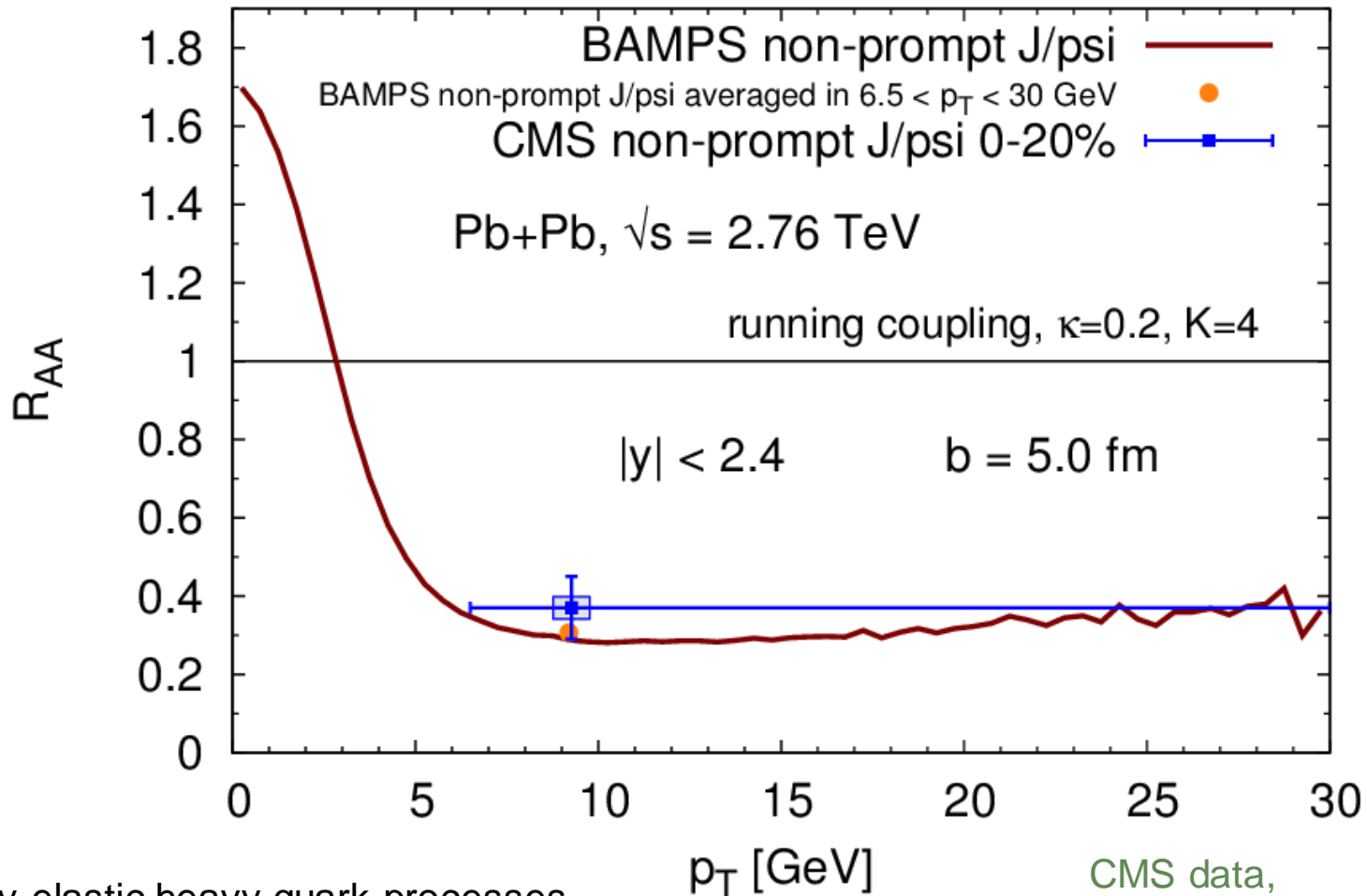
$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



only elastic heavy quark processes

# Non-prompt J/psi $R_{AA}$ at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$

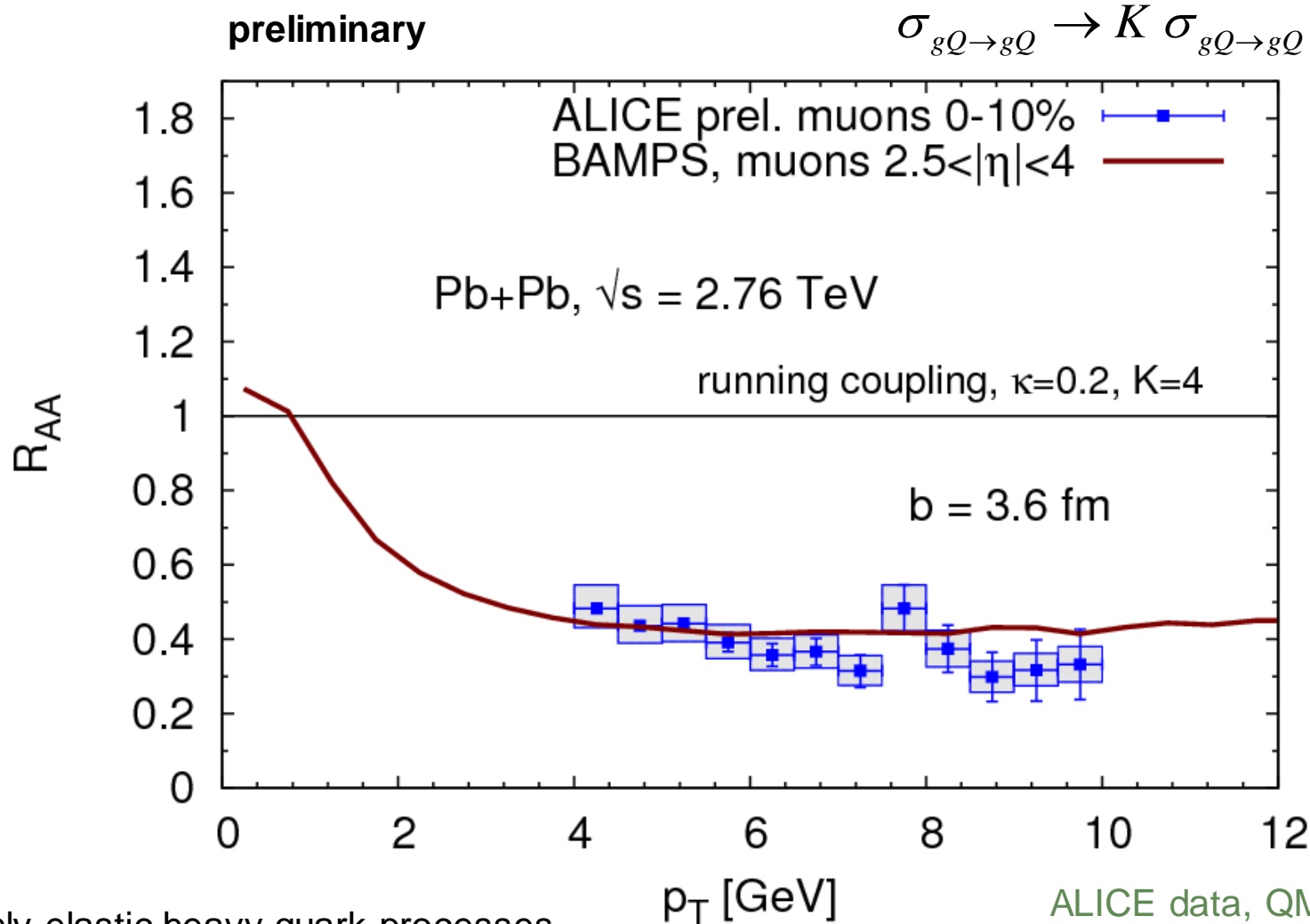


LHC

only elastic heavy quark processes

CMS data,  
arXiv:1201.5069

# Muon $R_{AA}$ at forward rapidity at LHC

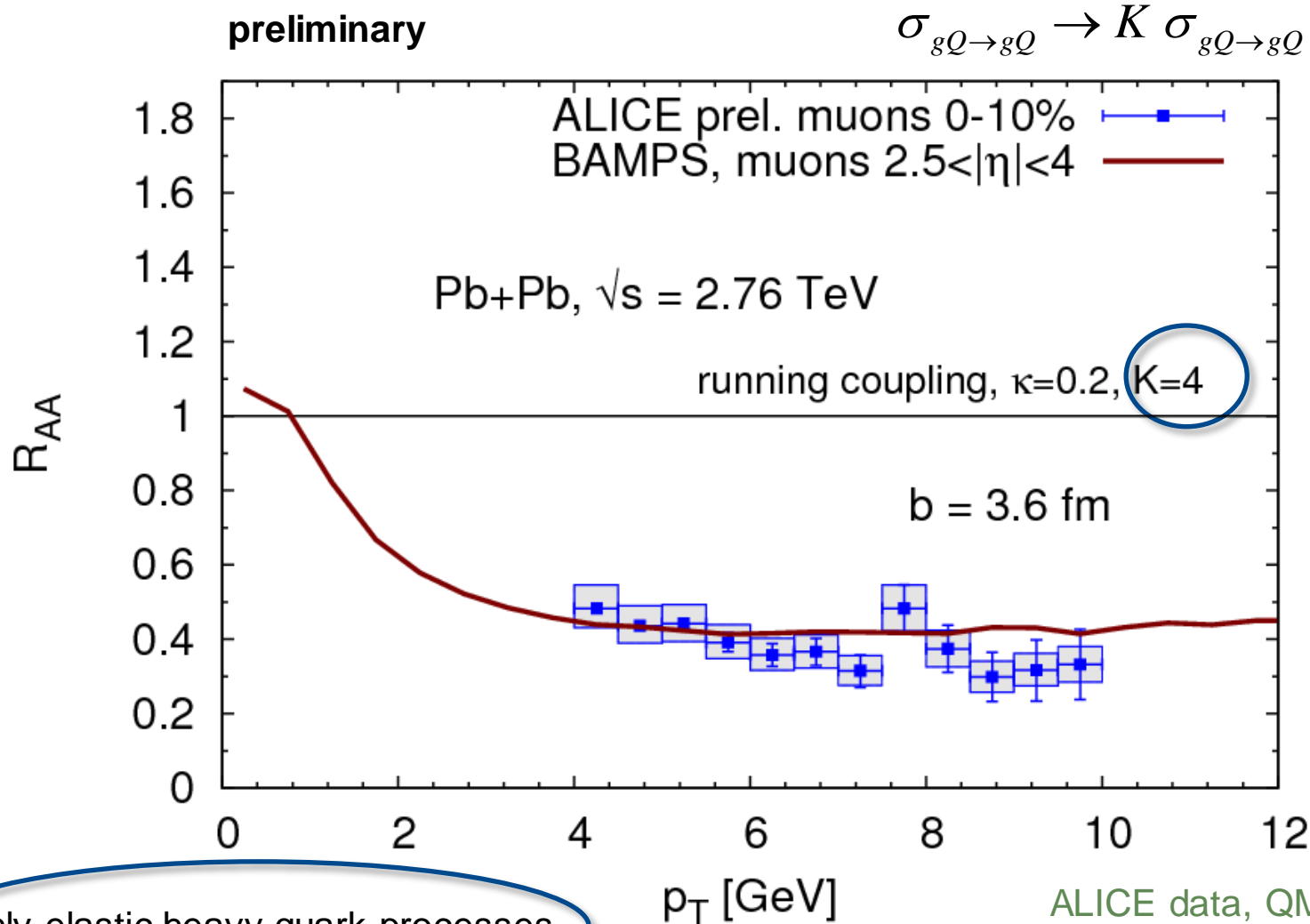


only elastic heavy quark processes

ALICE data, QM 2011,  
arXiv:1106.4042

# Muon $R_{AA}$ at forward rapidity at LHC

LHC

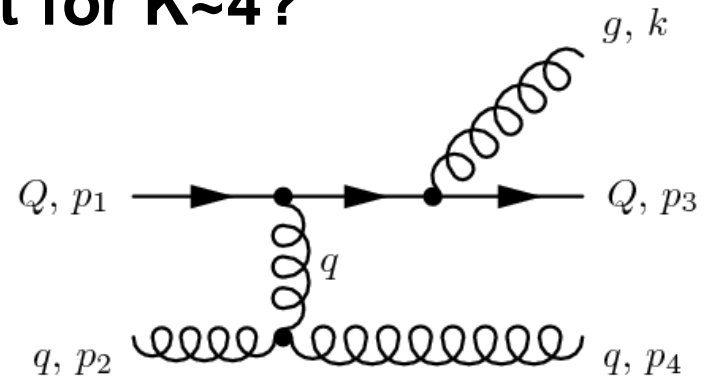


ALICE data, QM 2011,  
arXiv:1106.4042

# Radiative processes

Can radiative processes account for  $K \sim 4$ ?

$$g + Q \rightarrow g + Q + g$$



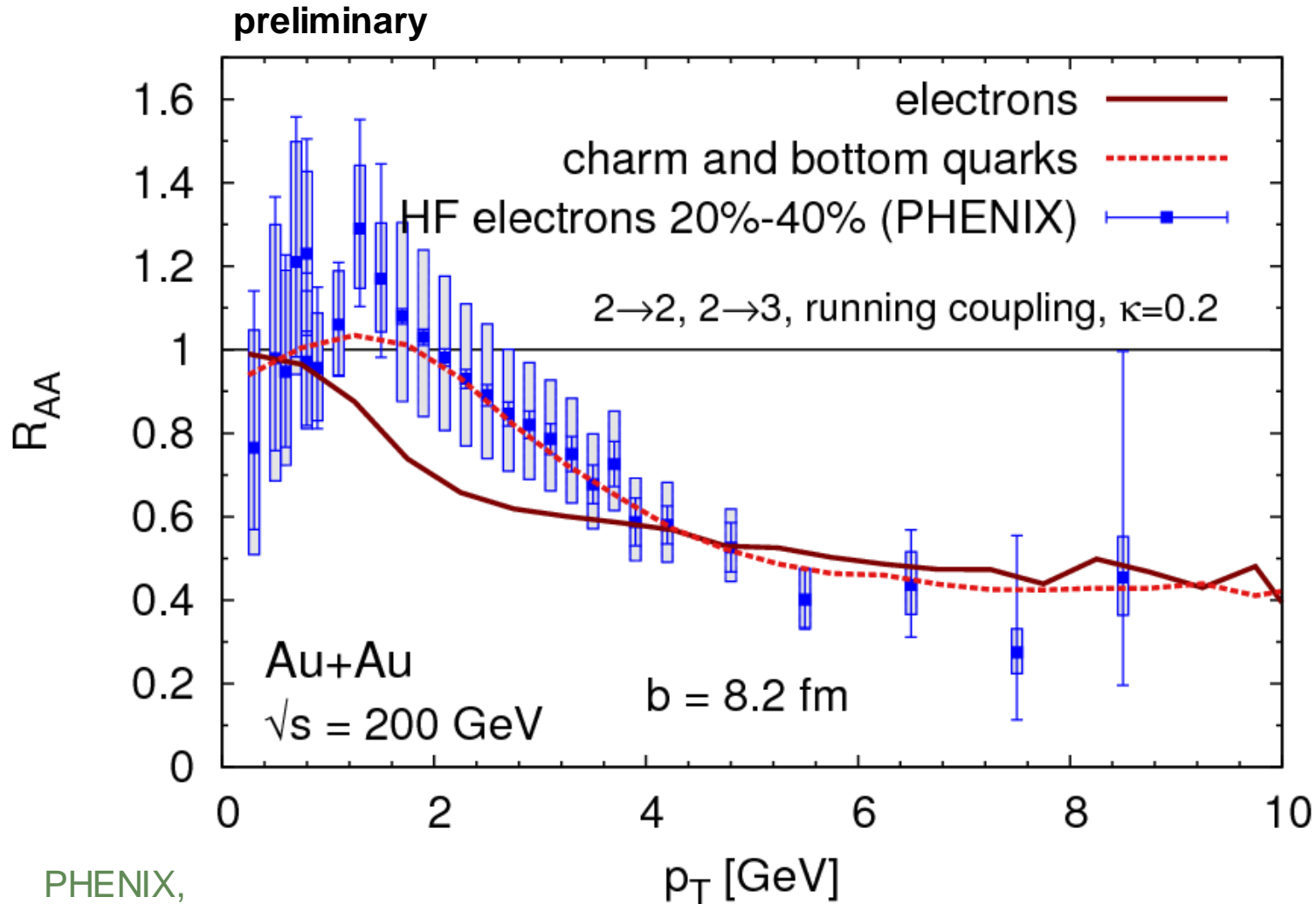
Gunion-Bertsch matrix element generalized to heavy quarks:

$$|\overline{\mathcal{M}}_{gQ \rightarrow gQg}|^2 = 12g^2 \left| \overline{\mathcal{M}}_0^{gQ} \right|^2 \left[ \frac{\mathbf{k}_\perp}{k_\perp^2 + x^2 M^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + x^2 M^2} \right]^2$$

In accordance to scalar QCD result from  
Gossiaux, Aichelin, Gousset, Guiho, J.Phys.G37 (2010)



# Heavy quark $R_{AA}$ at RHIC with 2- $\rightarrow$ 3



RHIC

PHENIX,  
arXiv:1005.1627

# Conclusions & outlook

## Full space-time evolution of QGP with charm and bottom quarks

- Running coupling and improved Debye screening yield results that can explain experimental  $v_2$  and  $R_{AA}$  at RHIC if  $K=4$  is introduced
- Good agreement with D meson  $v_2$  and  $R_{AA}$  at LHC
- Good agreement with non-prompt J/psi and muon  $R_{AA}$  at LHC
- Preliminary results with  $2 \rightarrow 3$  in full cascade are promising

Further details on [arXiv:1104.2295](https://arxiv.org/abs/1104.2295) and [1112.1559](https://arxiv.org/abs/1112.1559)

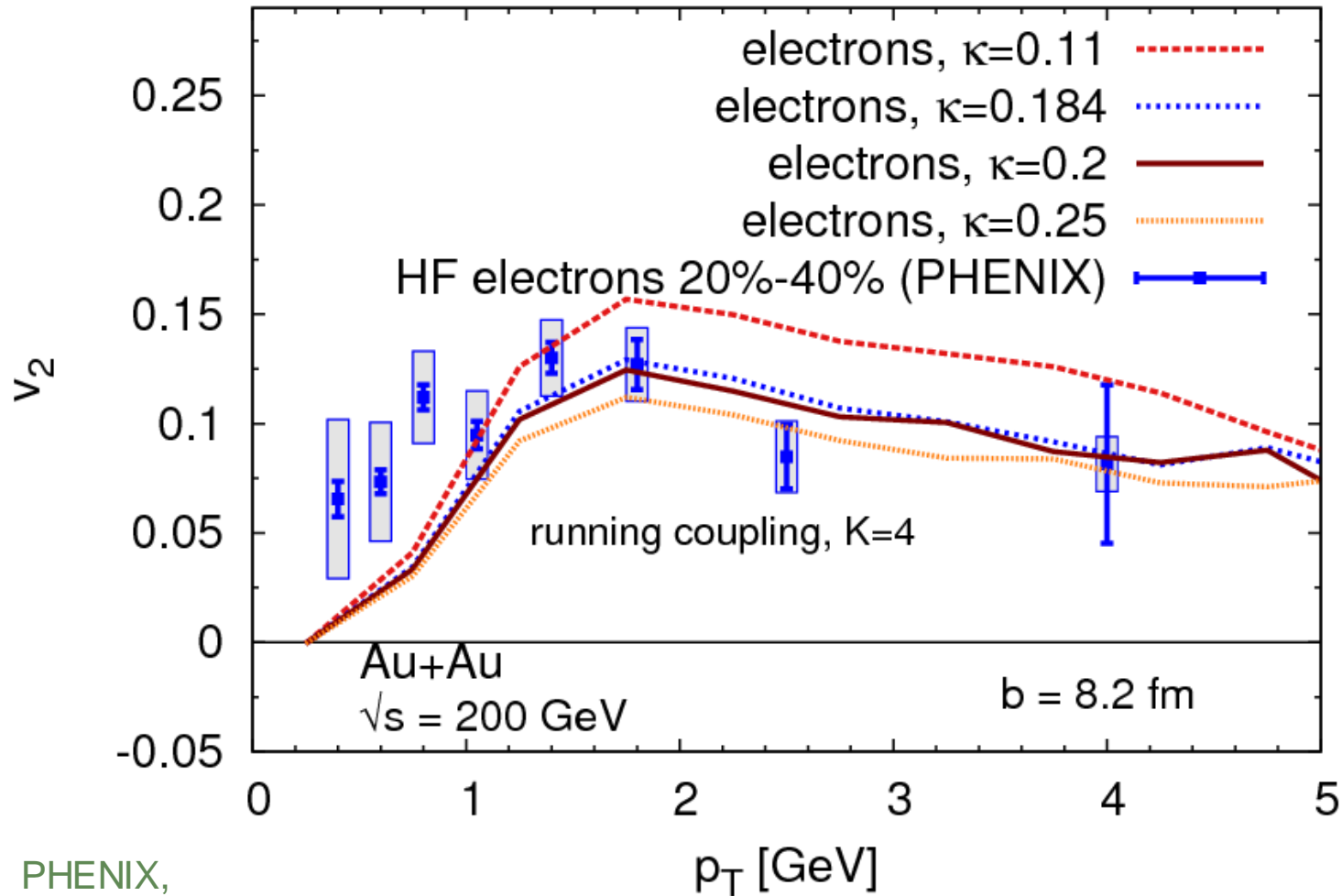
### Future tasks:

- Further study of radiative heavy quark scattering in full cascade
- Light quark interactions with heavy quarks

**Thank you for your attention.**

# Heavy quark elliptic flow $v_2$ at RHIC

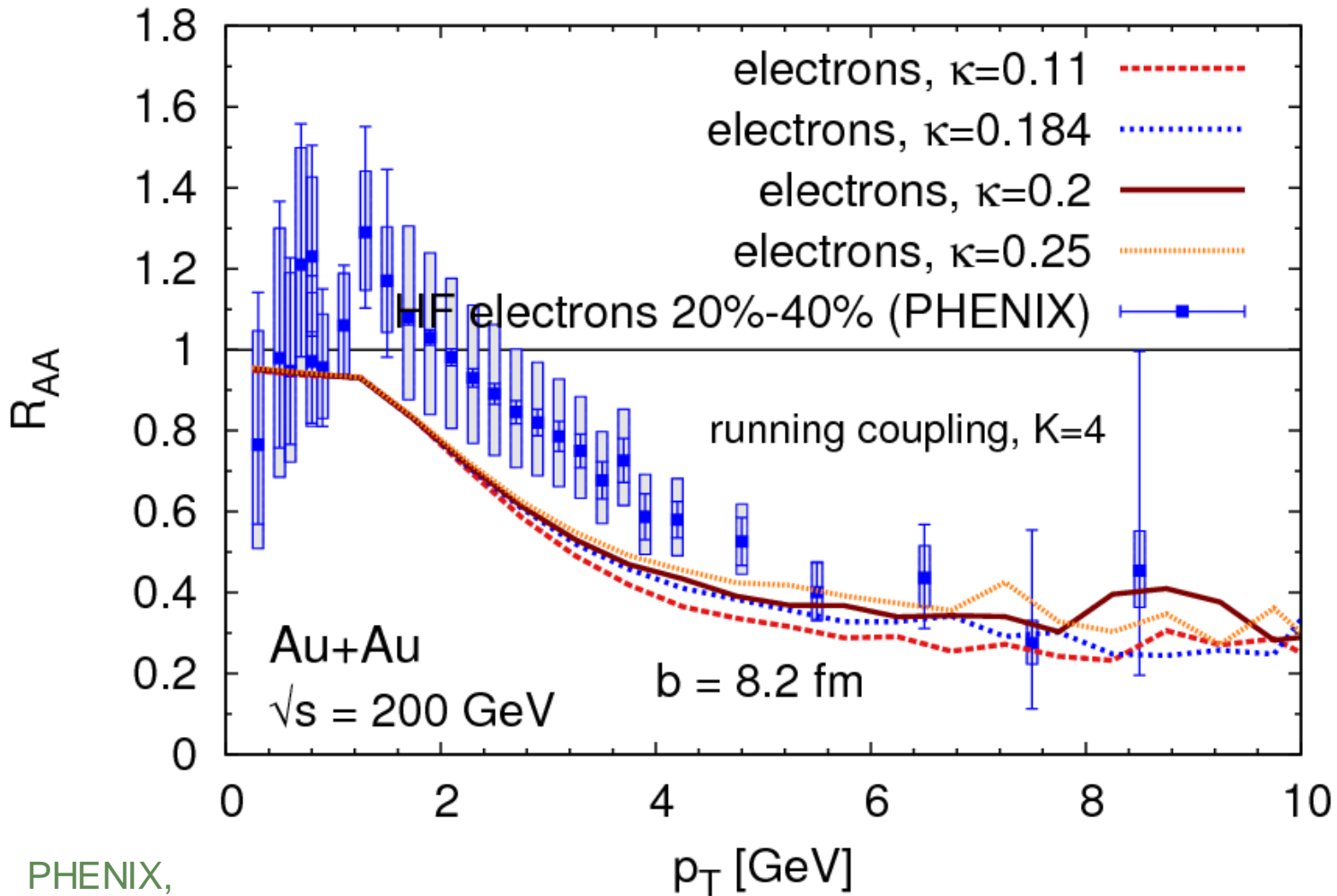
RHIC



PHENIX,  
arXiv:1005.1627

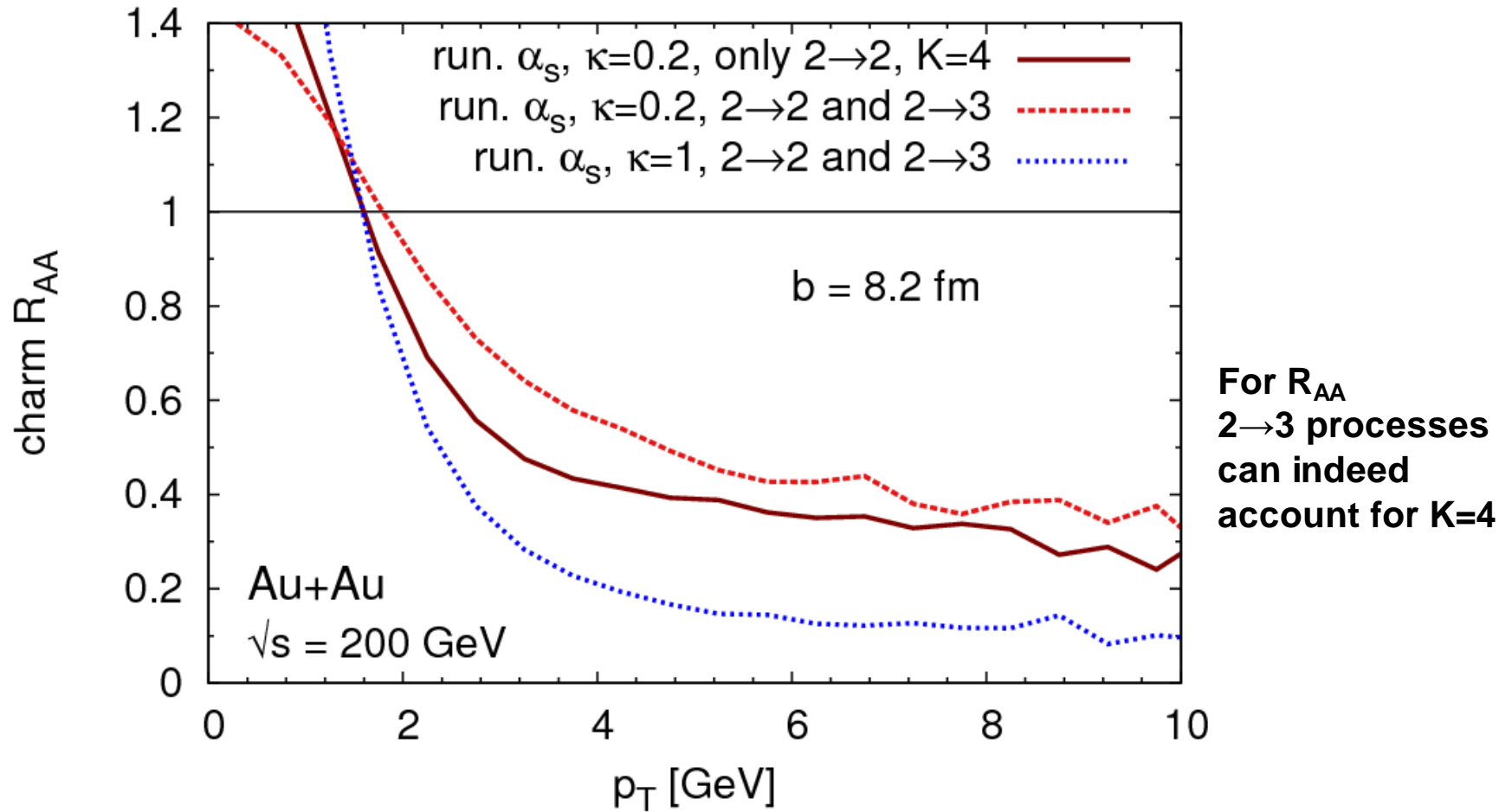
# Heavy quark elliptic flow $v_2$ at RHIC

RHIC



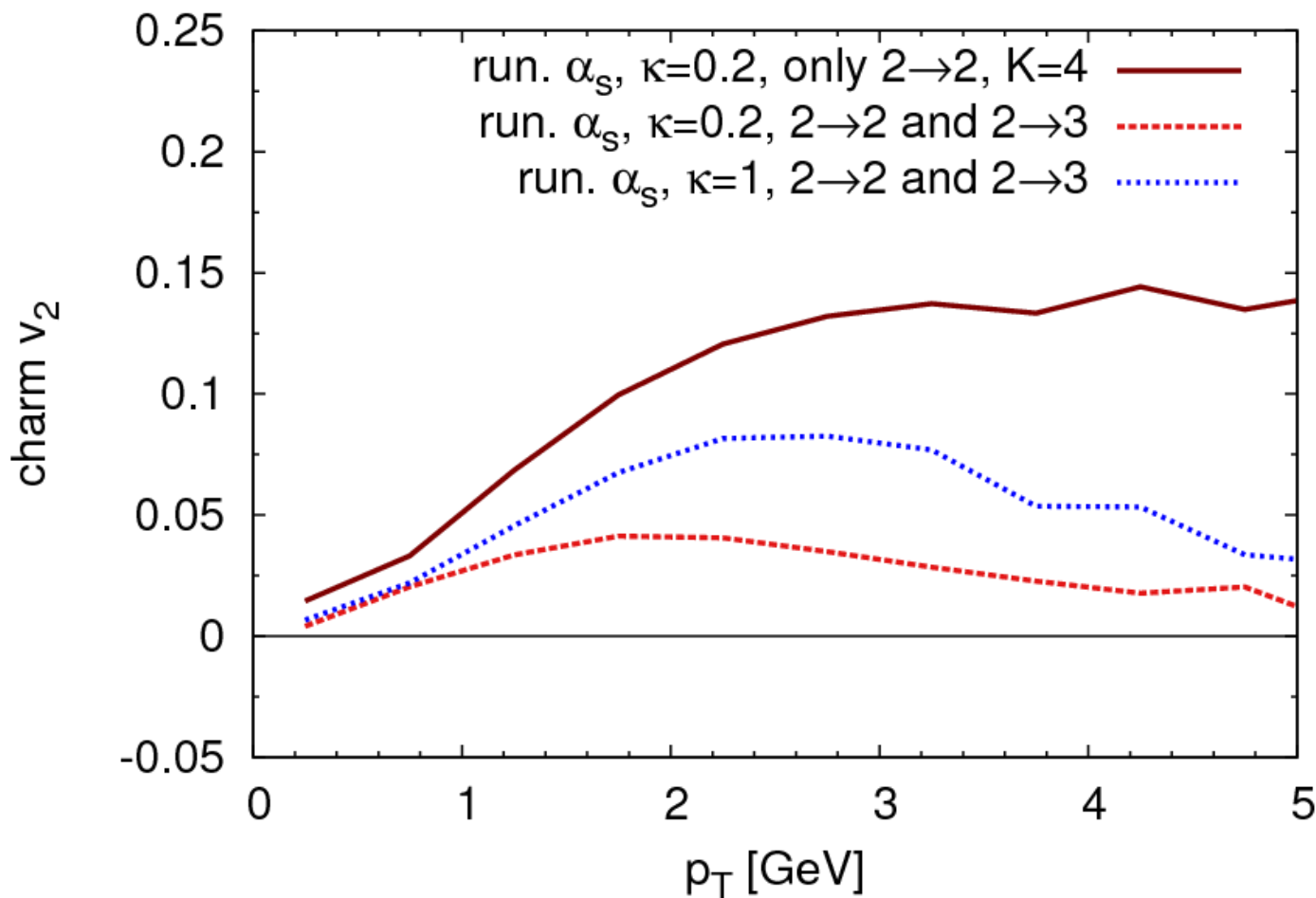
PHENIX,  
arXiv:1005.1627

# Charm $R_{AA}$ at RHIC



Only charm quarks (no heavy flavor electrons!) for better comparison

# Charm elliptic flow $v_2$ at RHIC



For  $v_2$   
 $2 \rightarrow 3$  processes  
cannot explain  
missing factor  
 $K=4$

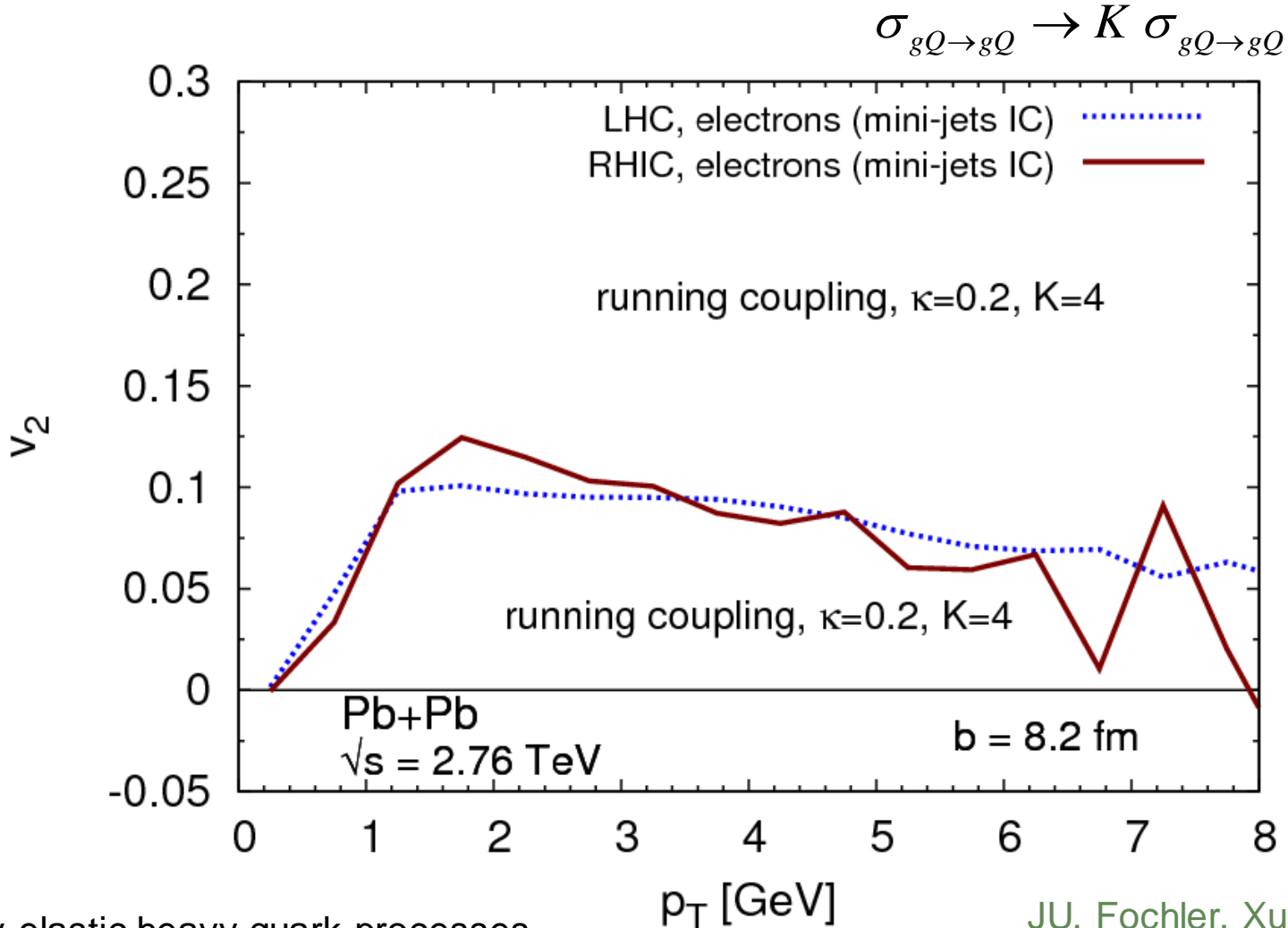
Different impact  
of  $2 \rightarrow 3$   
processes on  $v_2$   
and  $R_{AA}$

Reason:  
LPM effect

$\kappa=1$  is even  
better since  $2 \rightarrow 3$   
processes more  
important due to  
LPM effect

Only charm quarks (no heavy flavor electrons!) for better comparison

# Heavy quark elliptic flow $v_2$ at LHC



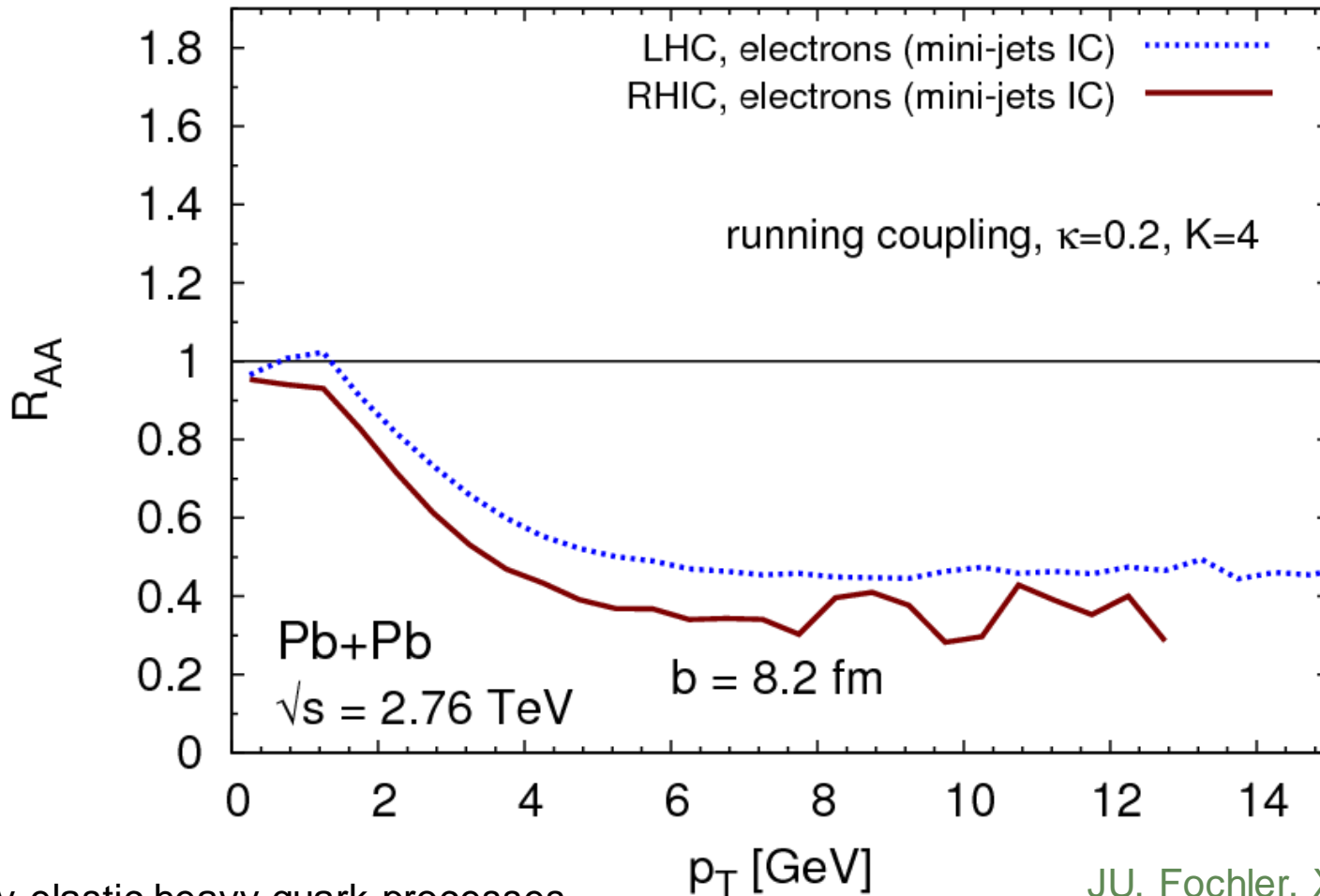
only elastic heavy quark processes

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)



# Heavy quark $R_{AA}$ at LHC

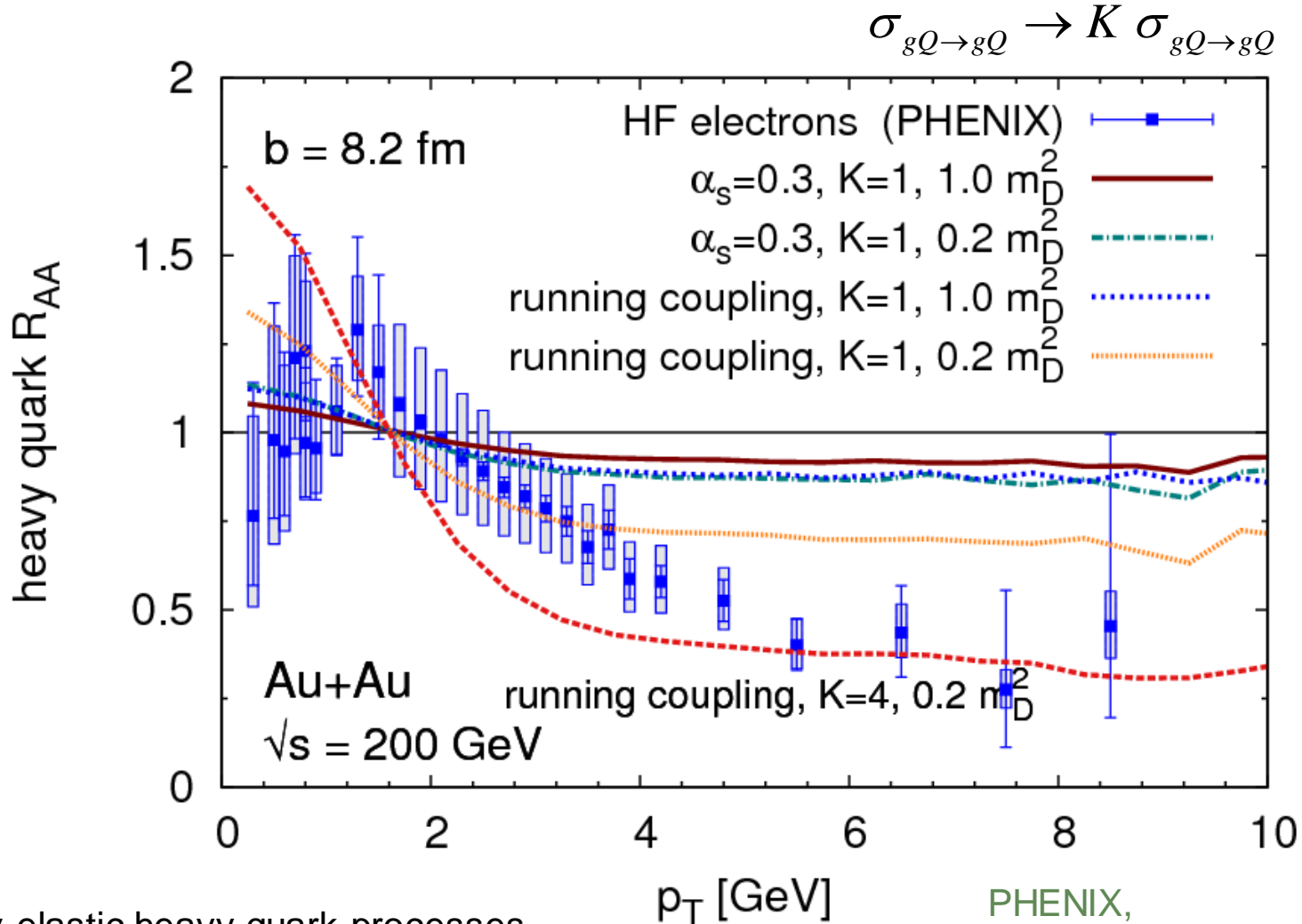
$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



only elastic heavy quark processes

JU, Fochler, Xu, Greiner  
Phys. Rev. C 84 (2011)

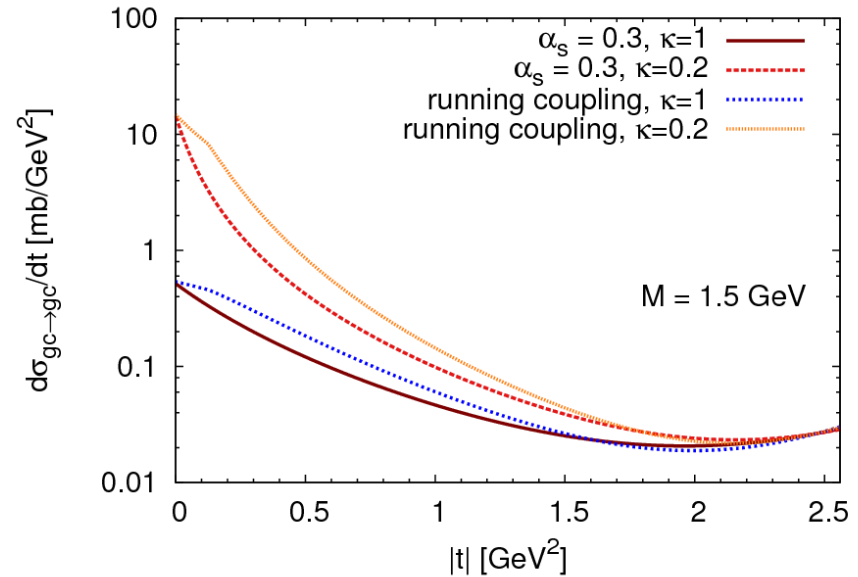
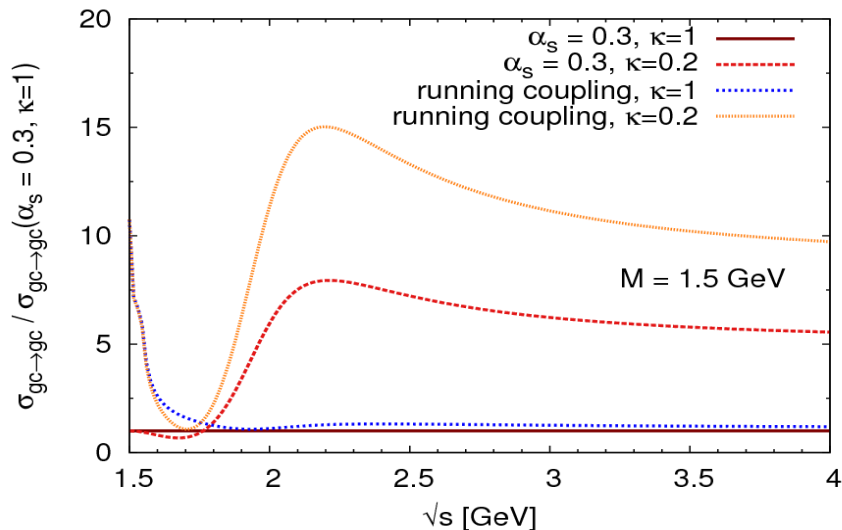
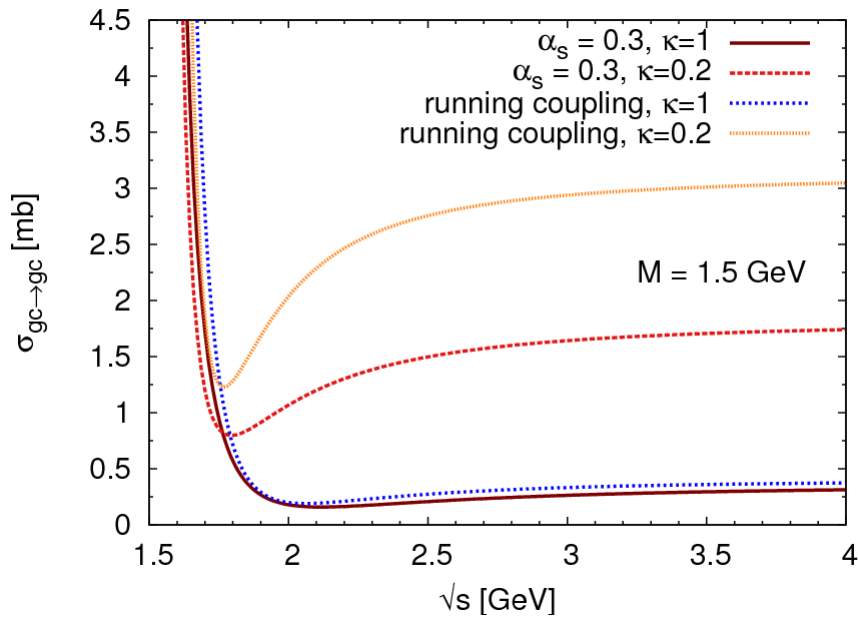
# Heavy quark $R_{AA}$ at RHIC



only elastic heavy quark processes

PHENIX,  
arXiv:1005.1627

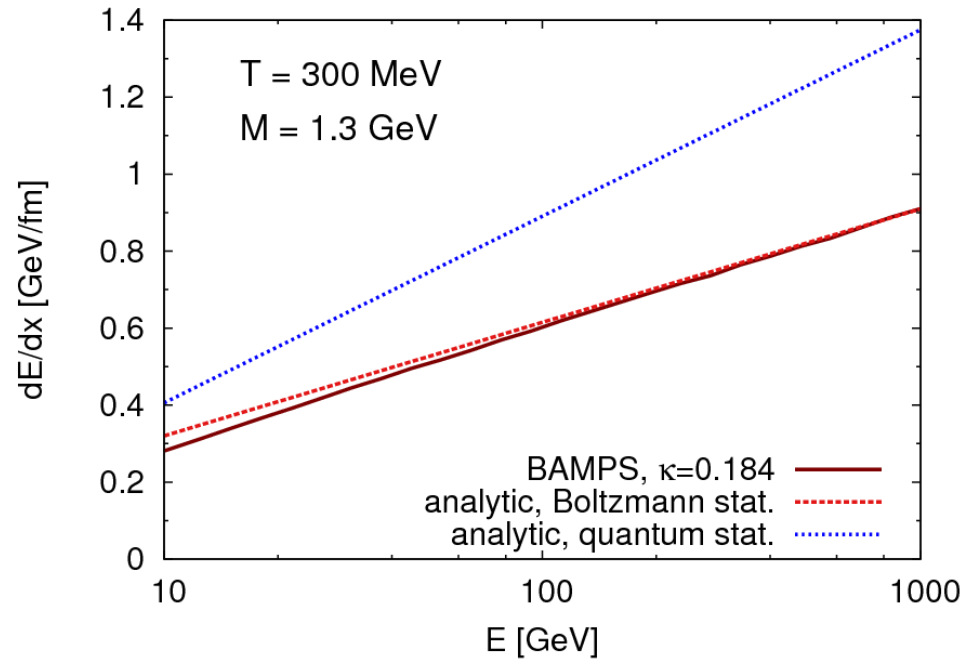
# Heavy quark scattering cross section



# Heavy quark scattering

## Compare to analytic formula

$$\frac{dE}{dx} = \frac{8\alpha_s^2 T^2}{\pi} \left[ \left(1 + \frac{n_f}{3}\right) \ln \frac{ET}{m_D^2} + \frac{2}{9} \ln \frac{ET}{M^2} \right. \\ \left. + \left(\ln 2 - \frac{1}{4} - \frac{\gamma}{3}\right) n_f \right. \\ \left. + \frac{31}{9} \ln 2 - \frac{101}{108} - \frac{11\gamma}{9} \right]$$



# Fragmentation and Decay

- Peterson fragmentation

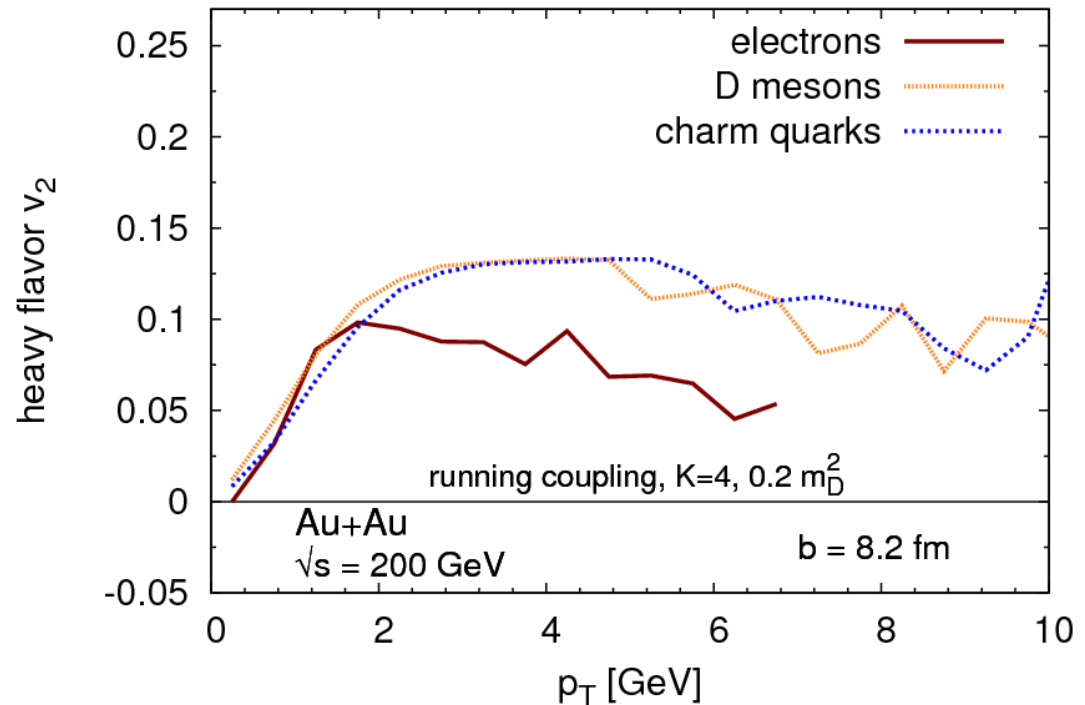
Peterson et al., Phys. Rev. D27 (1983)

$$D_{H/Q}(z) = \frac{N}{z \left( 1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z} \right)^2}$$

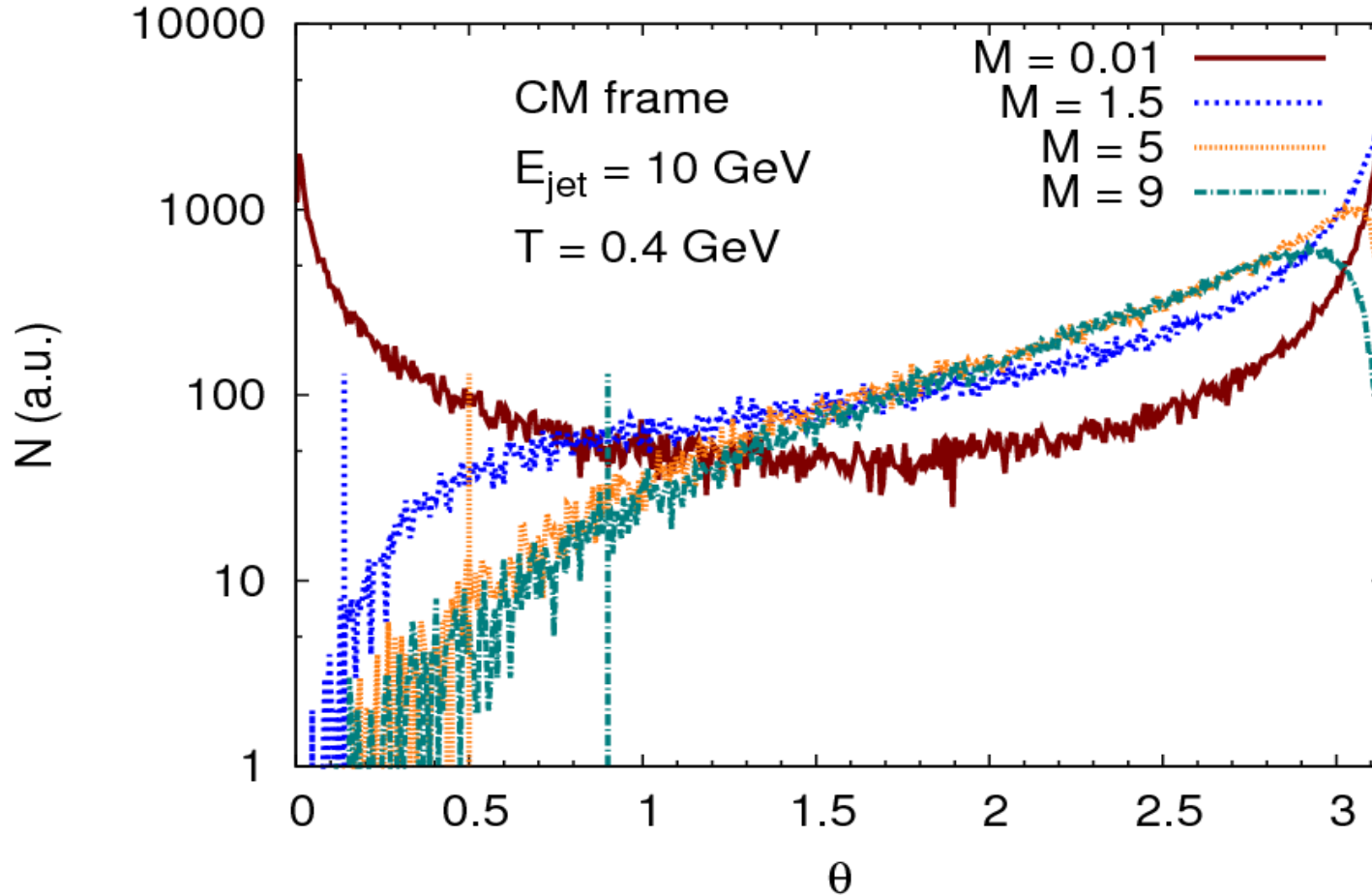
$$z = \frac{|\vec{p}_H|}{|\vec{p}_Q|} \quad \begin{array}{l} \epsilon_c = 0.05 \\ \epsilon_b = 0.005 \end{array}$$

- Decay to electrons with PYTHIA

Impact of hadronization and decay small

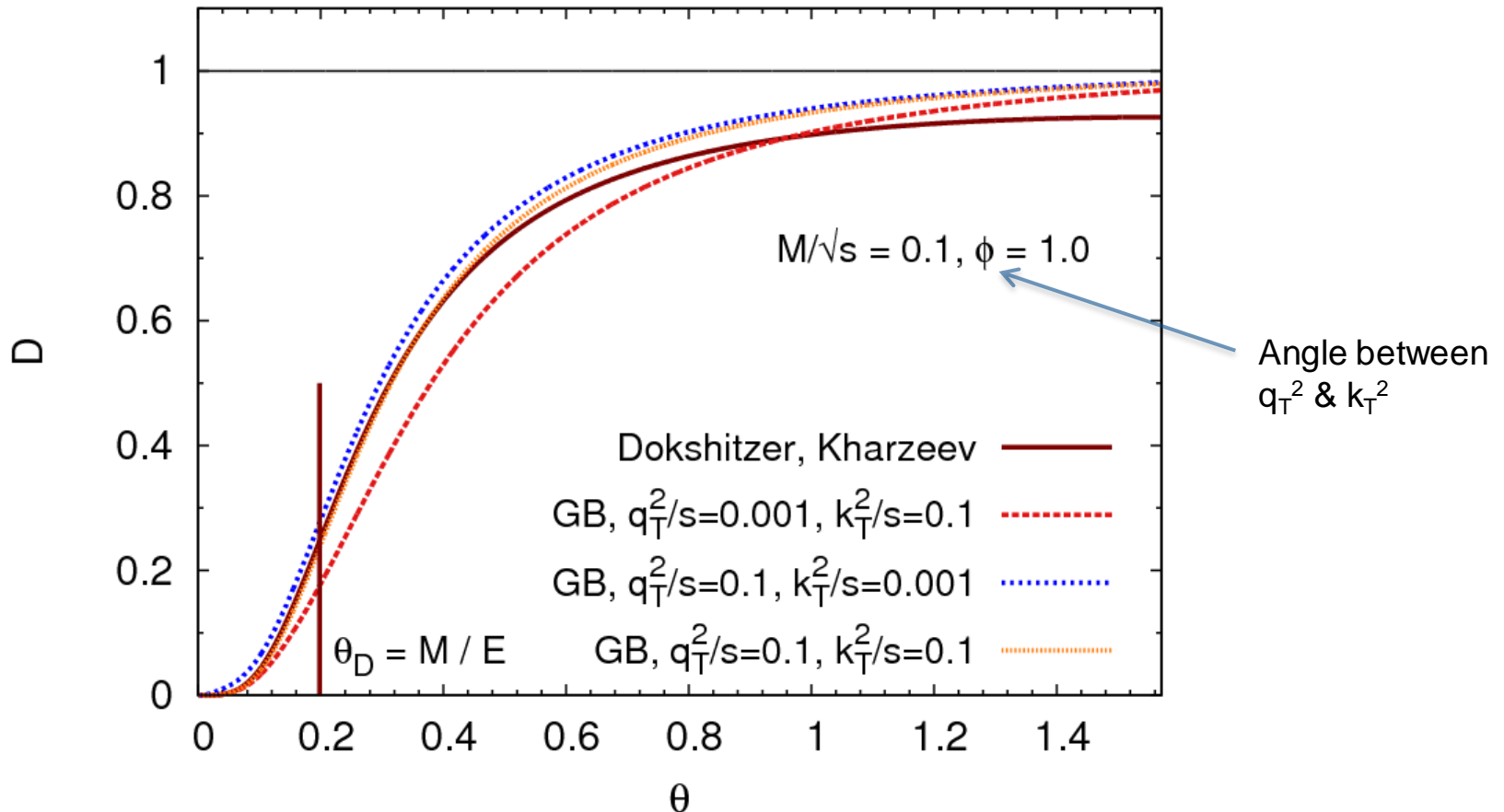


# $\Theta$ dependence in static medium

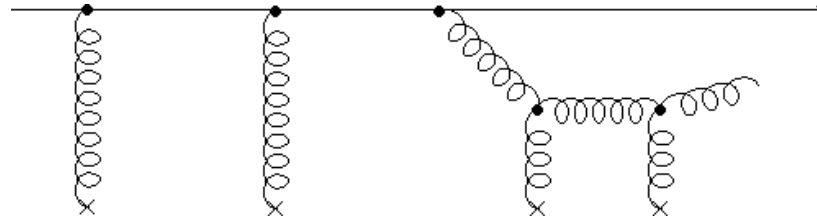


# Dead cone effect

$$|\overline{\mathcal{M}}_{gQ \rightarrow gQg}|^2 = 12g^2 |\overline{\mathcal{M}}_0^{gQ}|^2 \left[ \frac{\mathbf{k}_\perp}{k_\perp^2 + x^2 M^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + x^2 M^2} \right]^2$$



# LPM effect



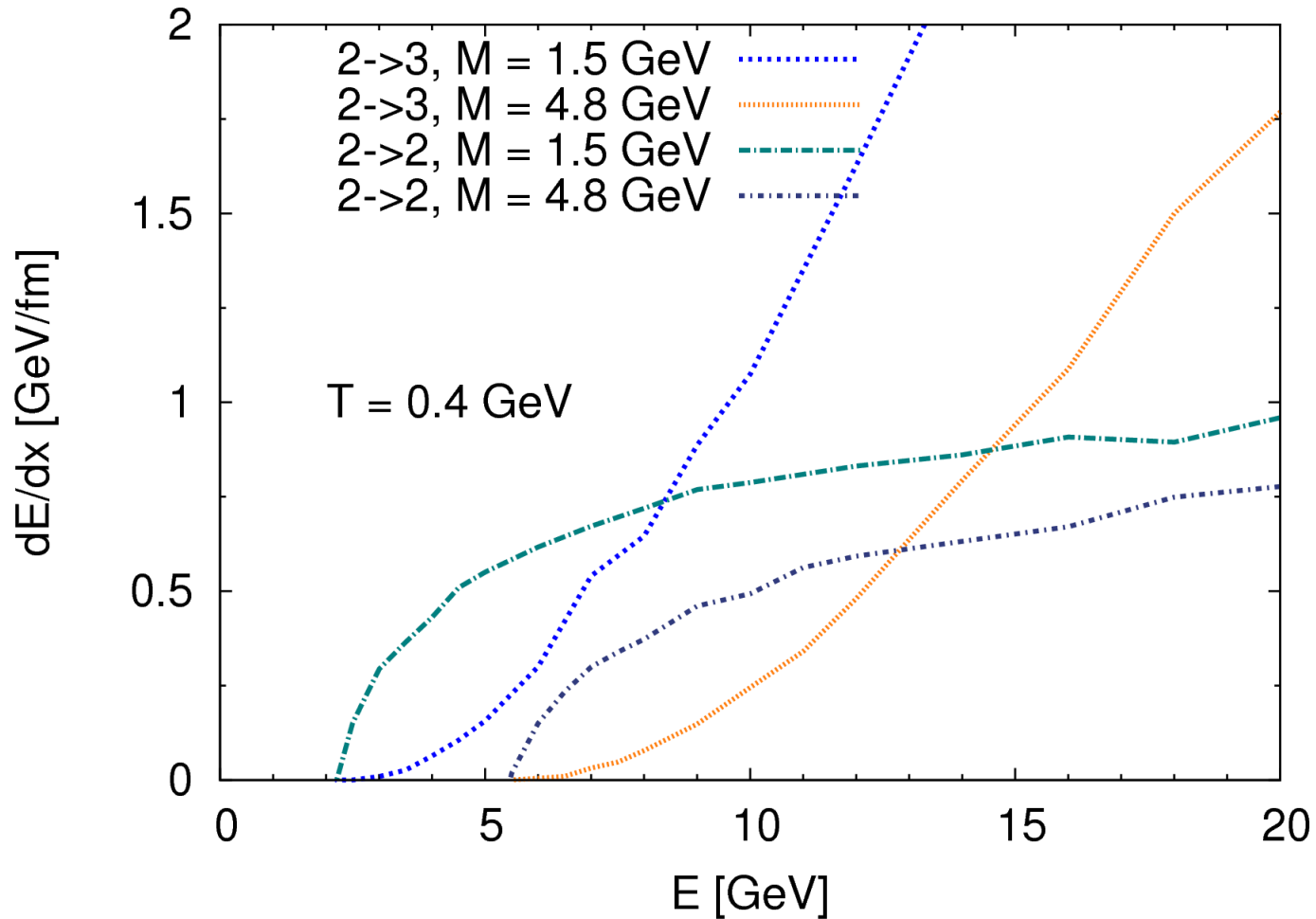
$$\lambda > \tau$$

2 → 3 only allowed if mean free path of jet larger than formation time of radiated gluon

➔ Bethe-Heitler regime, independent scatterings

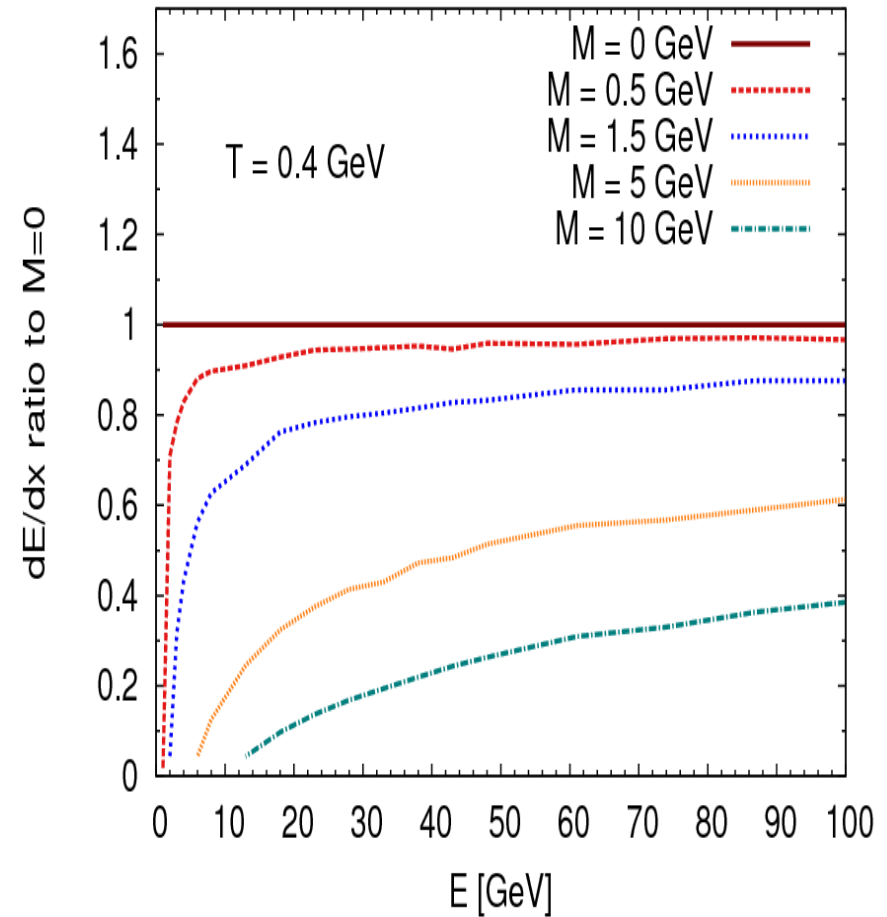
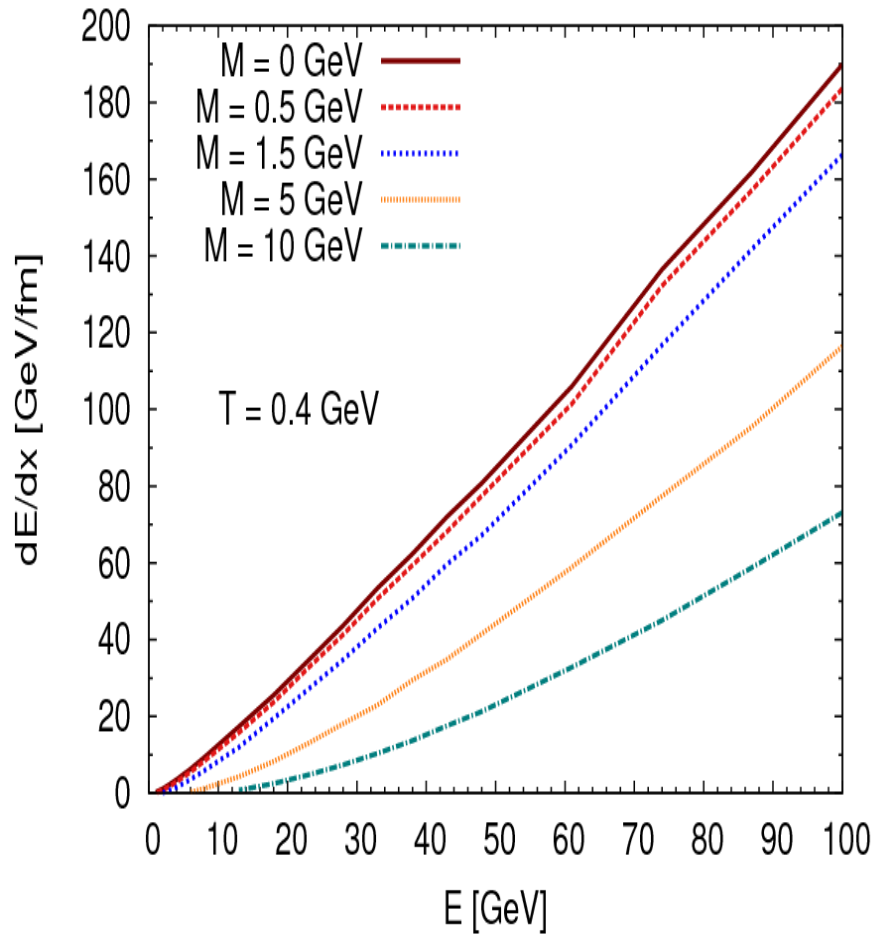


# Energy loss in static medium



Running coupling, with LPM effect

# Energy loss in static medium



**Fixed coupling, without LPM effect**

# Initial conditions

## Gluons:

- **PYTHIA**  
scaling to heavy-ion collisions with Glauber model (considering shadowing) and energy conservation
- **Minijets**  
(low  $p_T$  cut-off at 1.4 GeV)
- **Color glass condensate**  
H.J. Drescher & Y. Nara, Phys. Rev. C75 (2007)

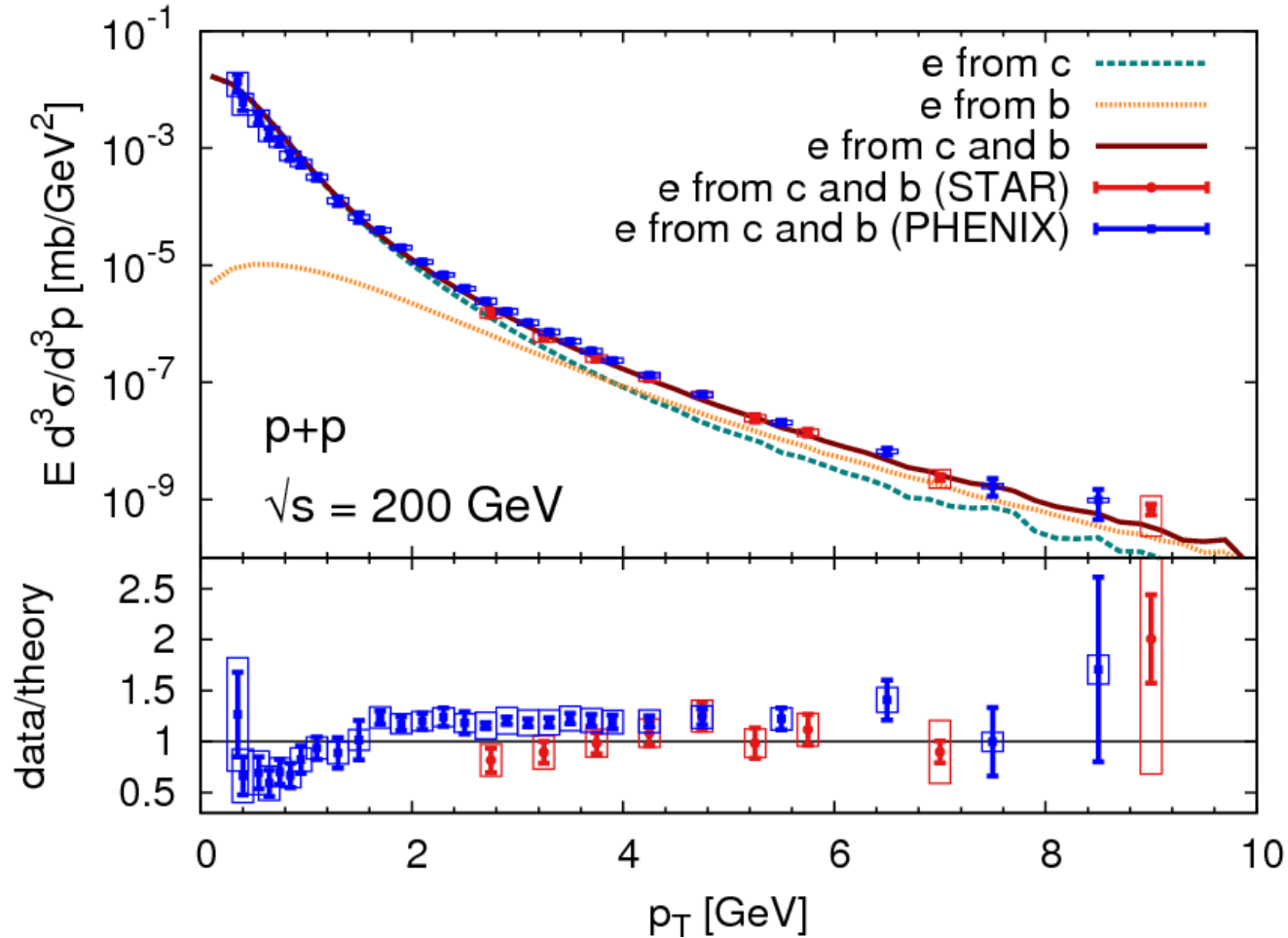
## Heavy quarks:

- **PYTHIA**  
Monte Carlo Event Generator for nucleon-nucleon collisions



- **NLO pQCD**  
Distributions from R. Vogt
- **MC@NLO**  
Next-to-leading order matrix elements

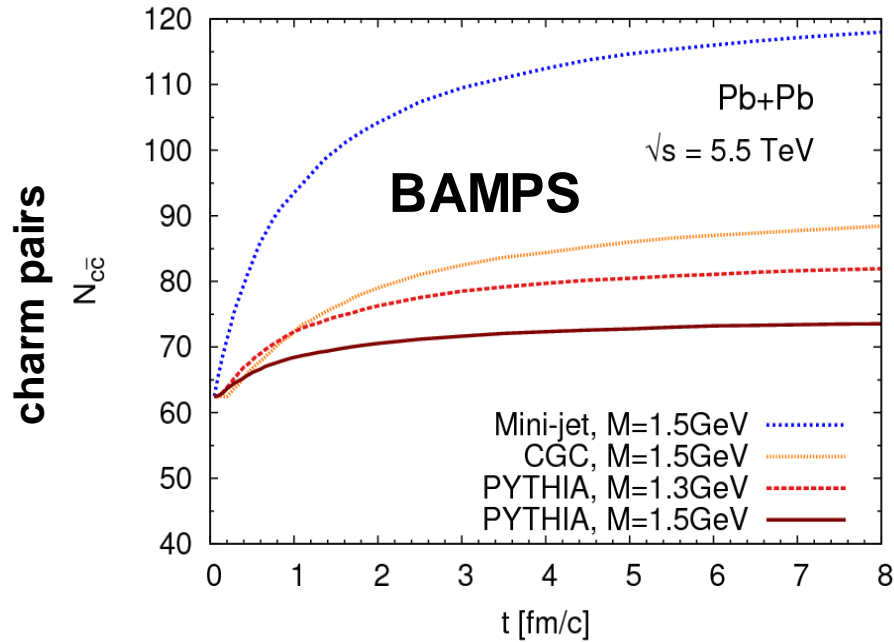
# Initial heavy quark distribution



$$\mu_F = \mu_R = 0.65 \sqrt{p_T^2 + M_c^2} \text{ for charm } (M_c = 1.3 \text{ GeV})$$

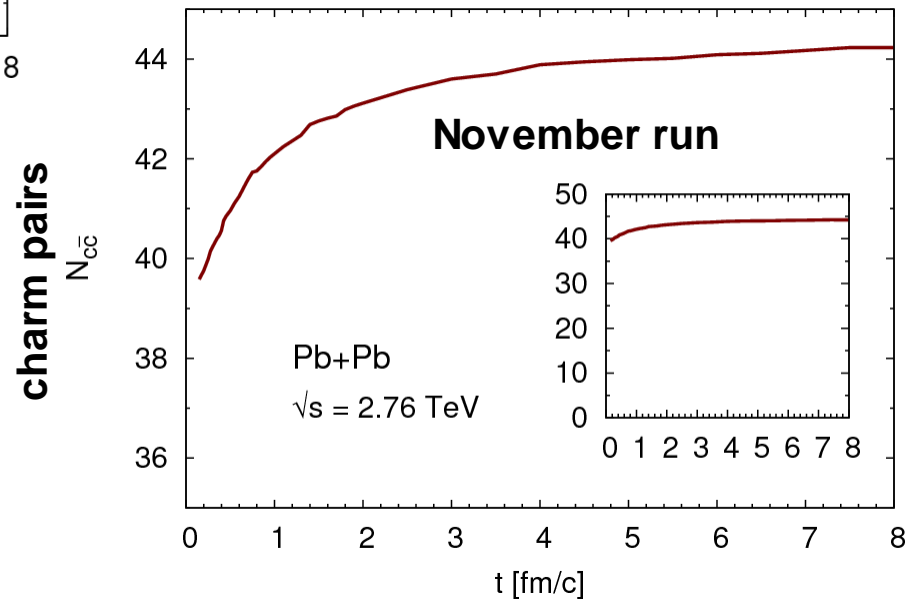
$$\mu_F = \mu_R = 0.4 \sqrt{p_T^2 + M_b^2} \text{ for bottom quarks } (M_b = 4.6 \text{ GeV})$$

# Charm production in the QGP at LHC



LHC

Large secondary production  
→ Can even be comparable to  
initial production



JU, Fochler, Xu, Greiner  
Phys. Rev. C 82 (2010)