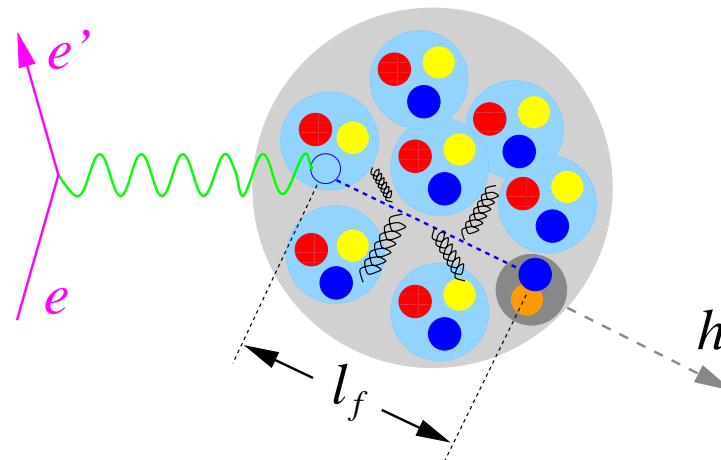


7th international workshop on High- p_T Physics at LHC, Frankfurt 2012

Color Transparency Effect in Deep Inelastic Scattering and at LHC era

- HERMES experiment and its kinematic
- Coherence length effect in exclusive ρ^0 electroproduction
- Color transparency effect
- Data on hadron formation and p_T -broadening at HERMES
- R_{AA} from ALICE and CMS data
- Summary and Outlook

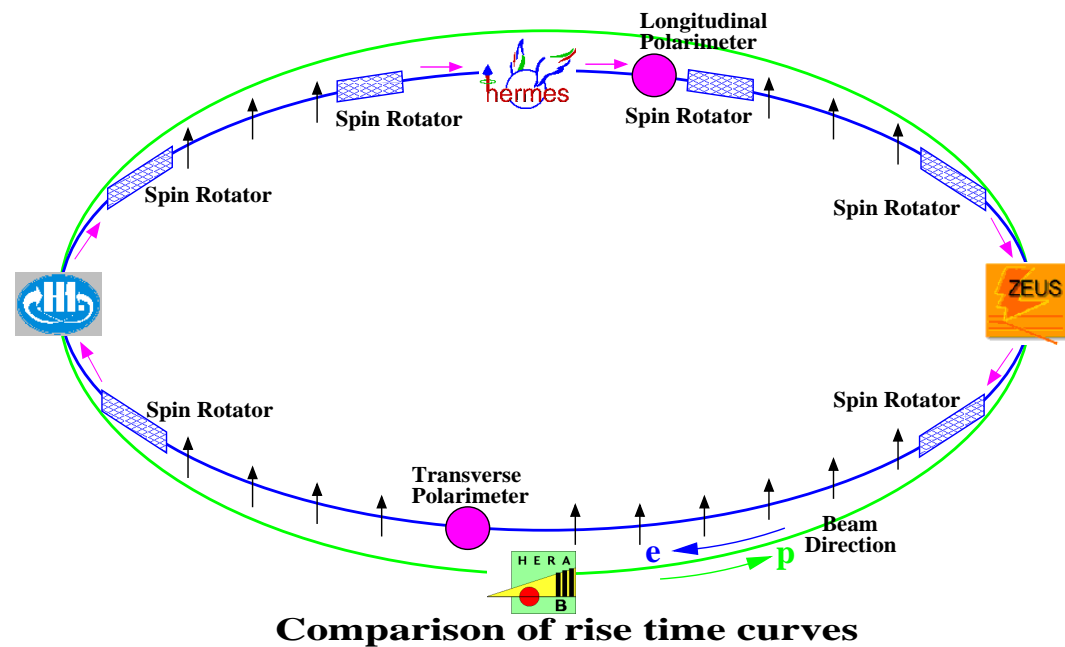


Preamble

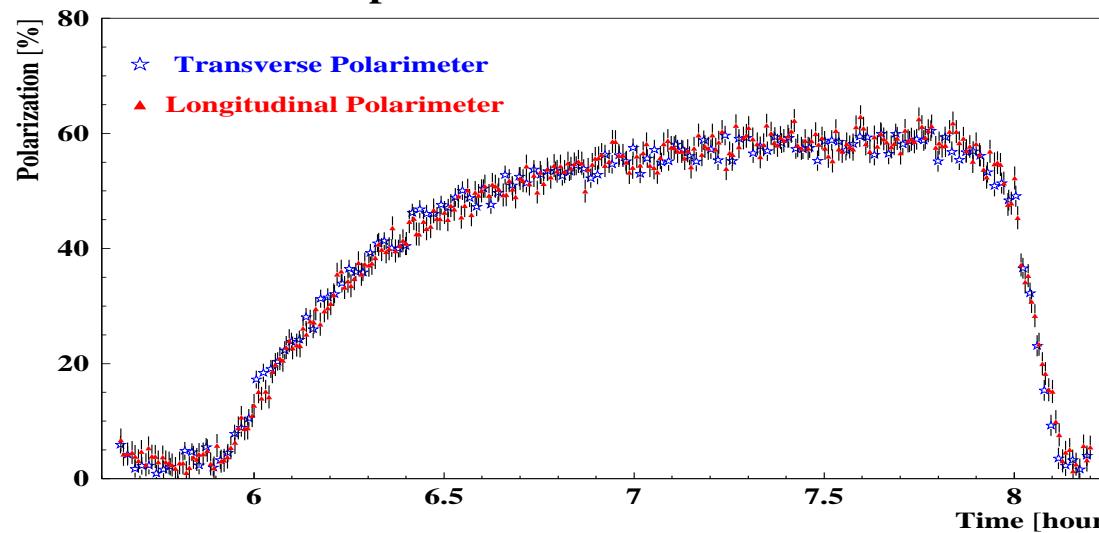
- Recent ALICE and CMS data on R_{AA} of charge hadrons are compared with the calculations from several models.
- Paper “Color transparency and suppression of high- p_T hadrons in nuclear collisions” (B.Z.Kopeliovich, I.K.Potashnikova and I.Schmidt, Phys.Rev. C. 83, 021901 (R), 2011) provides one of the model-dependent explanations.
- Described below HERMES data on exclusive ρ^0 (coherent length and color transparency effects) and semi-inclusive hadron electroproduction (hadron formation and p_T broadening) have been compared with the calculations and, in a such way, support its applicability.

Longitudinally Polarized $e^{+(-)}$ Beam at HERA

$P = 27.56$ GeV/c, current 50...100 mA, polarization of about $55 \pm 2\%$

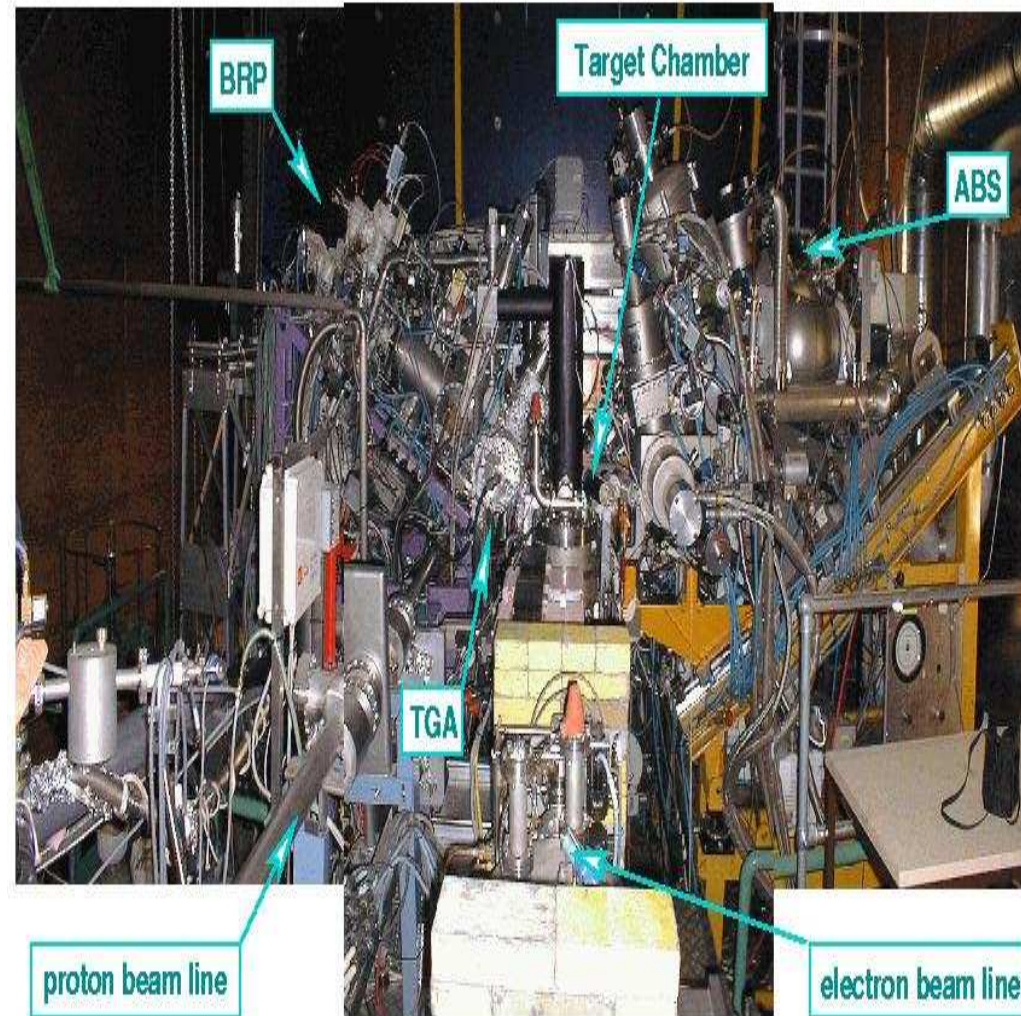
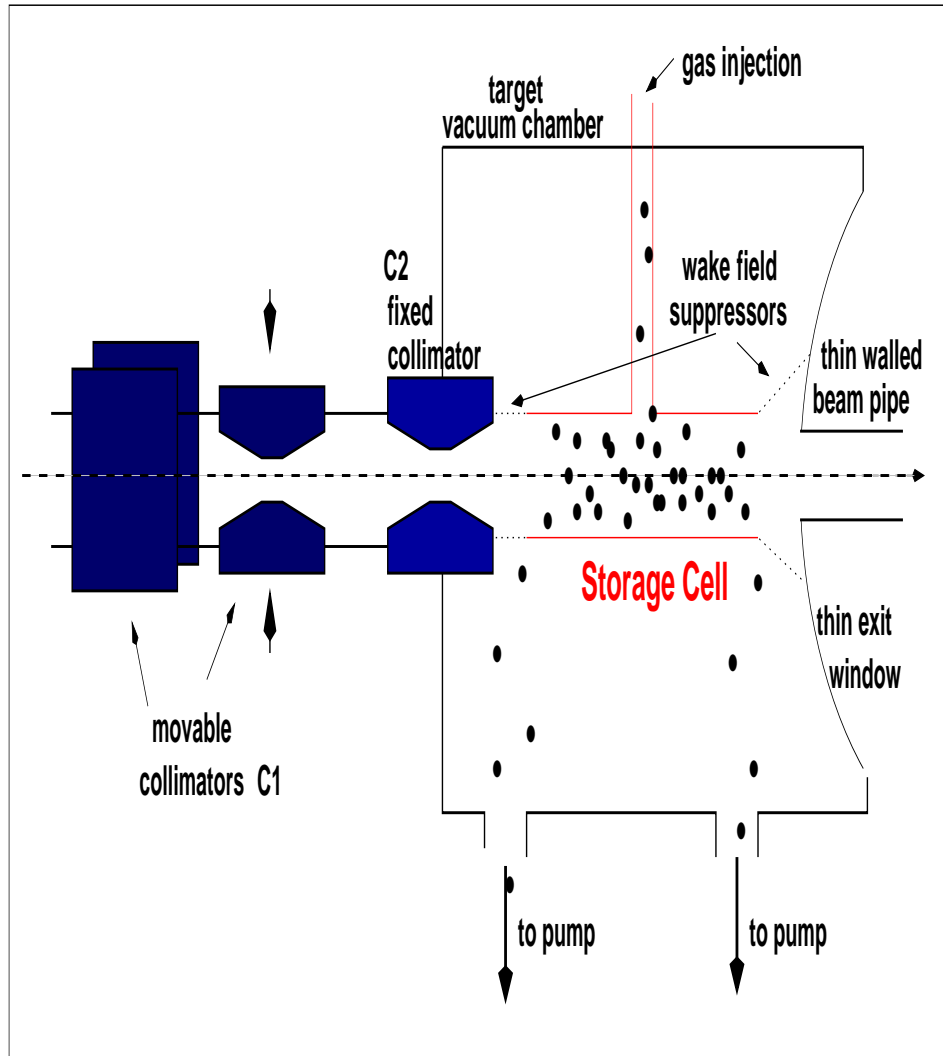


Comparison of rise time curves

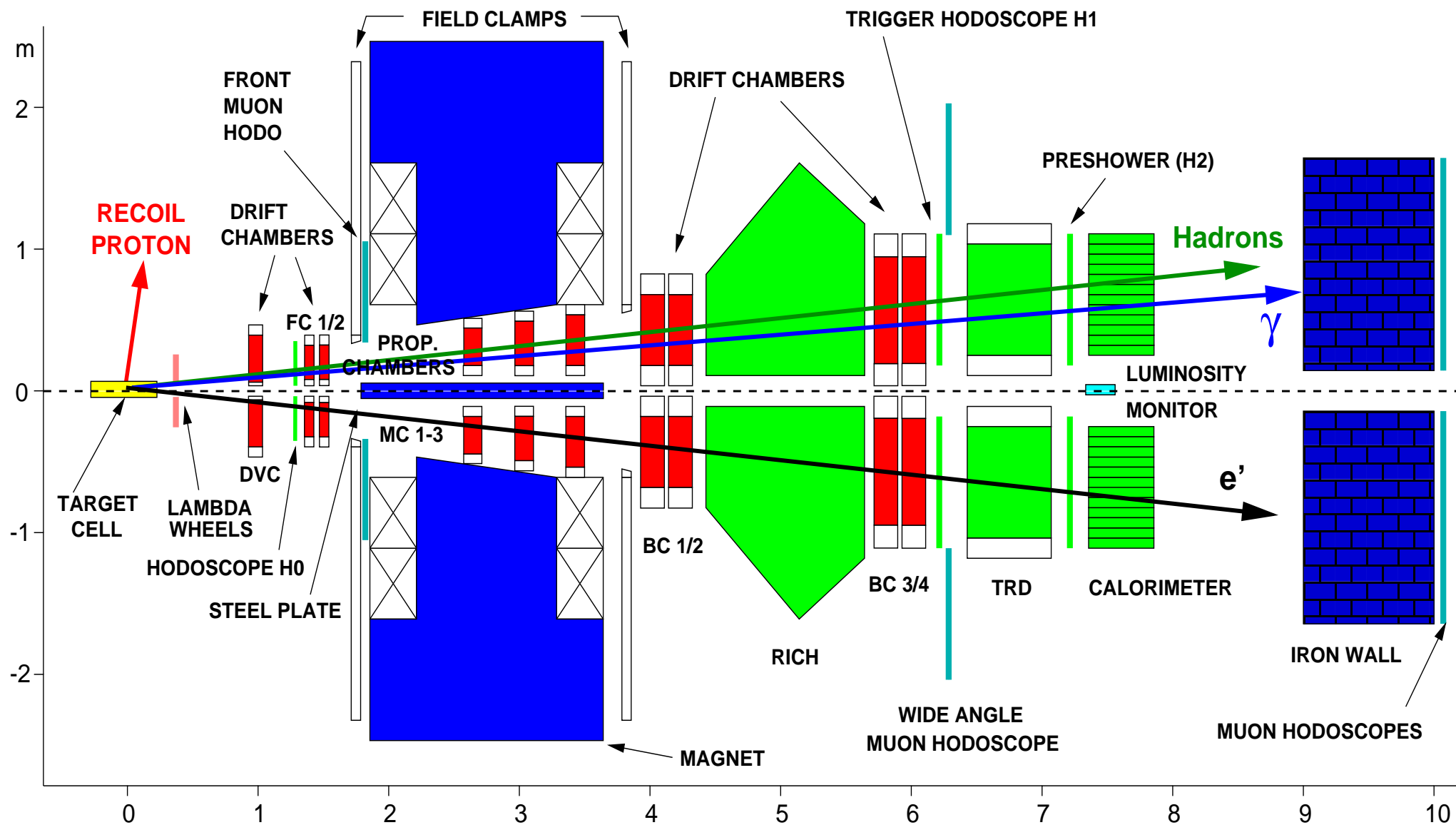


Internal Storage Cell Gas Target

polarized: $\sim 10^{14}$ nucl/cm², longitudinal polarization $\sim 88\%$: ¹H, ²H; transverse $\sim 78\%$: ¹H
unpolarized: $\sim 5 \cdot 10^{15}$ nucl/cm²: ¹H, ²H, ⁴He, ¹⁴N, ²⁰Ne, ⁸⁴Kr, ¹³¹Xe



HERMES Detector *was* (1995-2007) Two Identical Halves of Forward Spectrometer



- Acceptance: $40 < \Theta < 220$ mrad, $|\Theta_x| < 170$ mrad, $40 < |\Theta_y| < 140$ mrad
- Resolution: $\delta p/p \leq 1\%$, $\delta\Theta \leq 0.6$ mrad

Tracking and Particle Identification

Forward and Backward Drift Chambers:

- Ar(90 %), CO₂(5 %), CF₄(5 %) gases,
- drift time \leq 300 ns,
- 1877 FASTBUS TDCs readout,
- order of 250 μ m resolution

- Electron identification done by RICH, TRD, Preshower, Calorimeter: efficiency \geq 98%, hadron contamination \leq 1%

● Calorimeter

- lead glass (F101) blocks $9 \times 9 \times 50$ cm in array of 420 for top and bottom parts

- resolution:

$$\Delta E/E[\%] = 1.5_{\pm 0.5} + 5.1_{\pm 1.1}/\sqrt{E[\text{GeV}]}$$

- $M_{\pi^0} = 134.9 \pm 0.2$ MeV, $\sigma = 12.5 \pm 0.2$ MeV

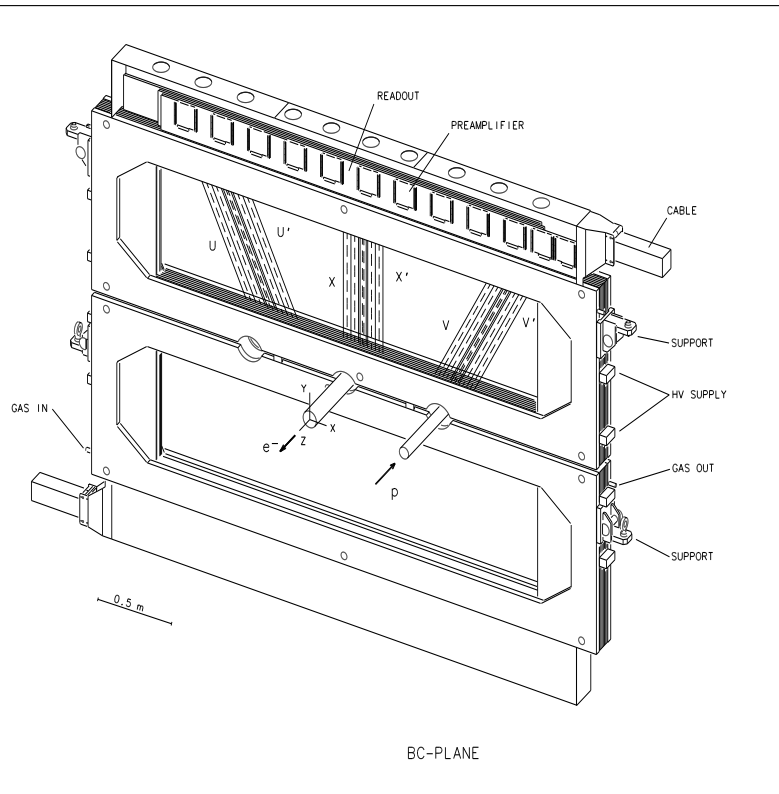
- RICH from 1998 for π^\pm vs K^\pm and $p(\bar{p})$ separation: \sim 4000 PMTs, aerogel and C₄F₁₀

- Dual rings of Cherenkov photons, $n=1.03$ for aerogel and $n=1.0014$ for C₄F₁₀

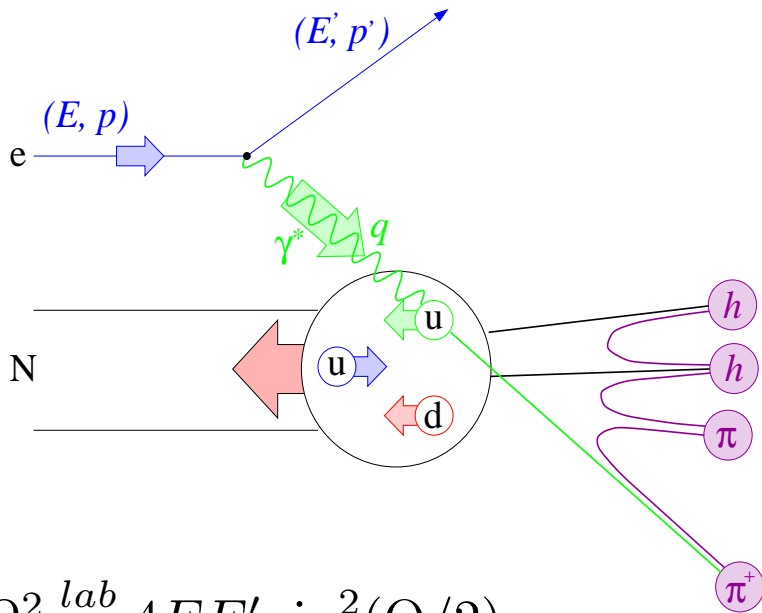
- Probability-Matrix for one track, no impact on p_t^2 of tracks

- RICH efficiencies over all kinematic region:

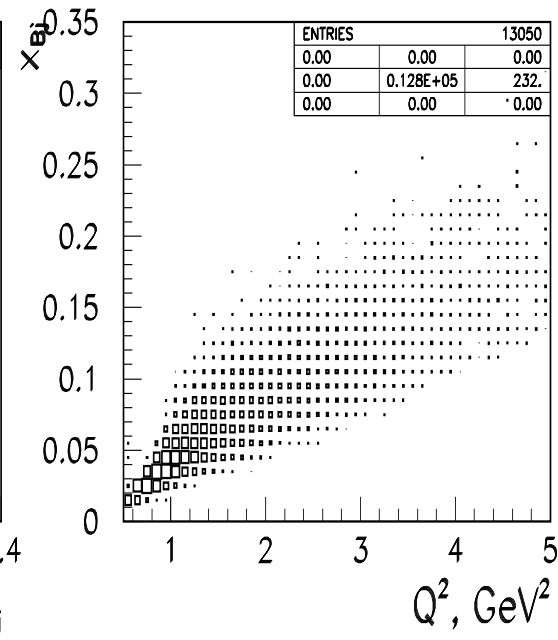
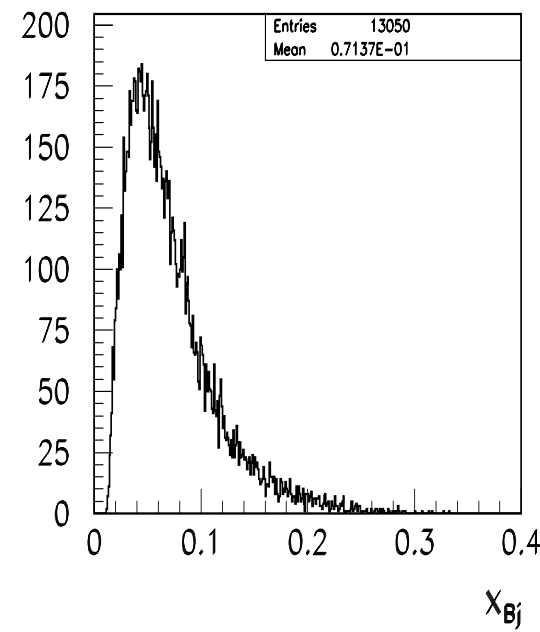
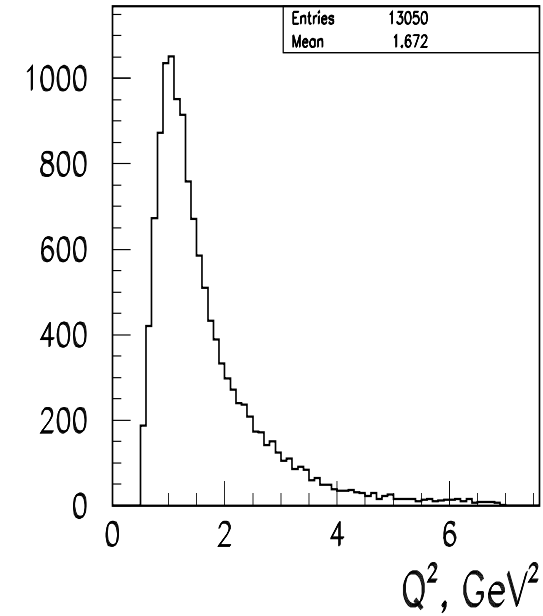
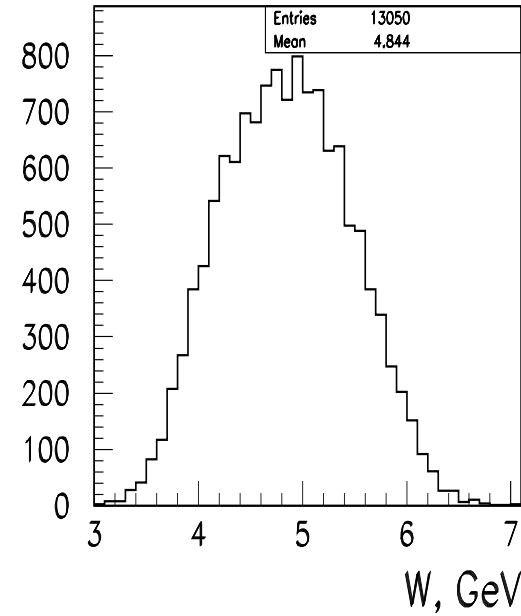
$$\pi: 98\%, K: 88\%$$



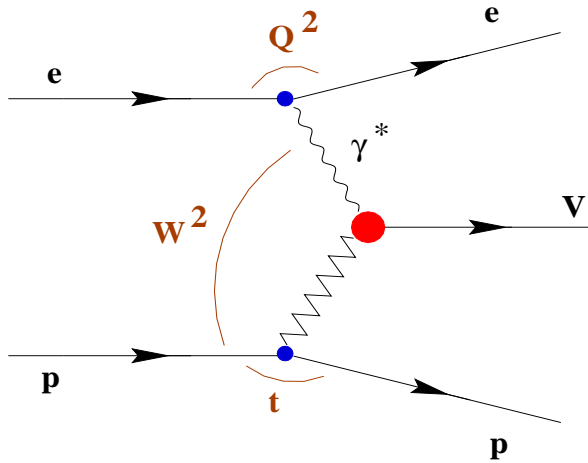
Deep Inelastic Scattering: Variables and Kinematic Distributions



- $Q^2 \stackrel{lab}{=} 4EE' \sin^2(\Theta/2)$
- $\nu \stackrel{lab}{=} E - E'$
- $x_{Bj} \stackrel{lab}{=} Q^2/2M\nu$
- $W^2 \stackrel{lab}{=} M^2 + 2M\nu - Q^2$
- $z \stackrel{lab}{=} E_h/\nu$

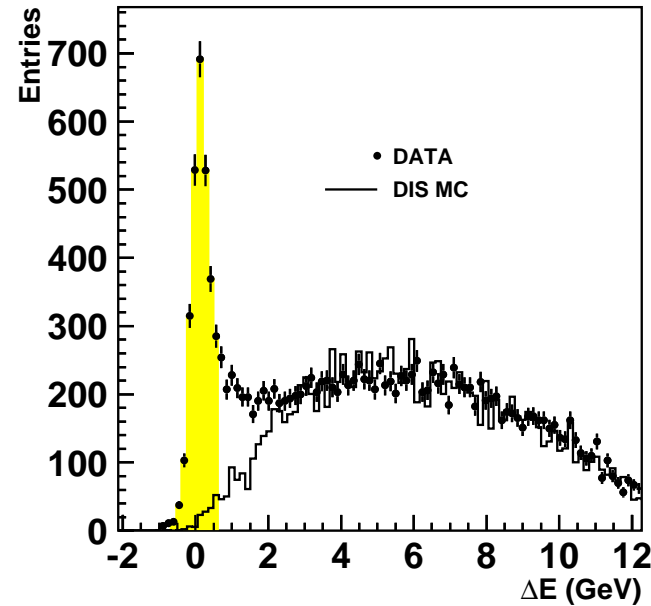
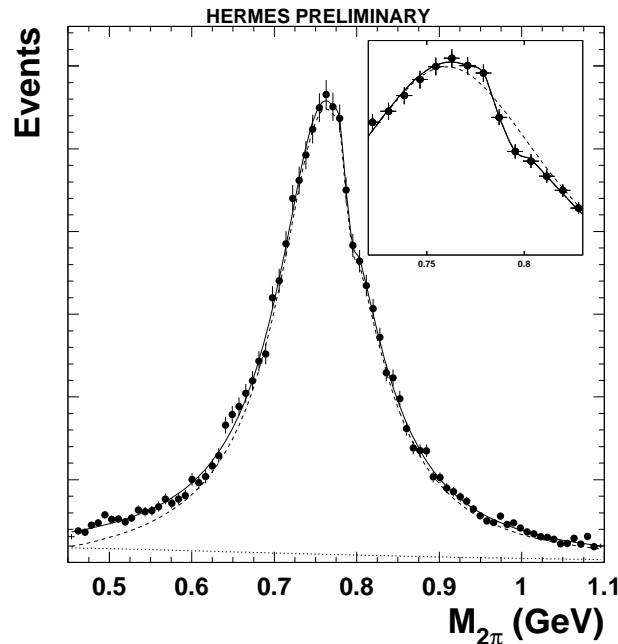


Kinematics of Exclusive ρ^0 Electroproduction



- $\nu = 5 \div 24$ GeV, $\langle \nu \rangle = 13.3$ GeV,
- $Q^2 = 1.0 \div 5.0$ GeV², $\langle Q^2 \rangle = 2.3$ GeV²
- $W = 3.0 \div 6.5$ GeV, $\langle W \rangle = 4.9$ GeV
- $x_{Bj} = 0.01 \div 0.35$, $\langle x_{Bj} \rangle = 0.07$

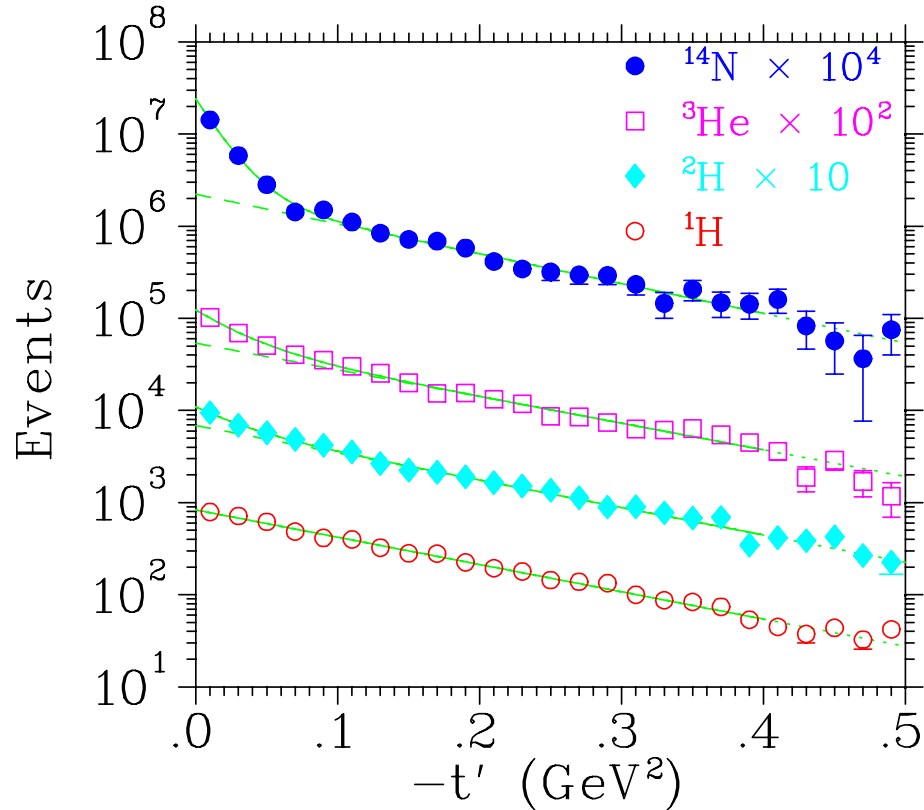
Clean mass and exclusive peaks for ρ^0 :



$$\Delta E = \frac{M_X^2 - M_p^2}{2M_p}, \quad M_X^2 = (p + q - V)^2, \quad \text{background is subtracted using PYTHIA MC.}$$

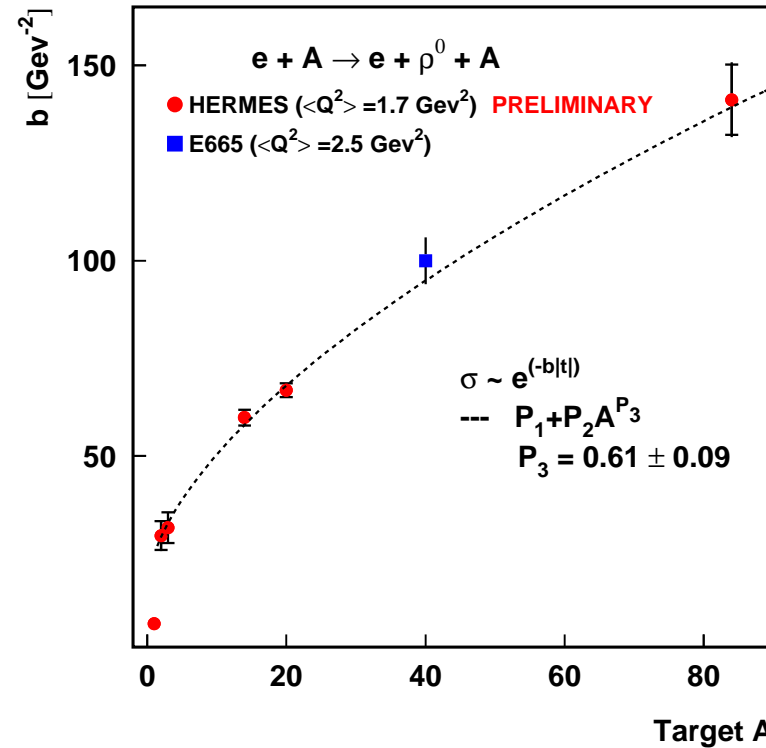
Nuclear Effects of Exclusive ρ^0 Production: Coherent and Incoherent ρ^0

HERMES collab., Phys.Lett.B 513 (2001) 301-310; Eur.Phys.J. C 29, 171 - 179 (2003)



$$d\sigma_{tot}/dt \propto a_1 \exp(b_{coh}t) + a_2 \exp(b_{inct}t)$$

at $-t \lesssim 0.045 \text{ GeV}^2$ coherent ρ^0 dominates,
 at $-t \gtrsim 0.1 \text{ GeV}^2$ incoherent.

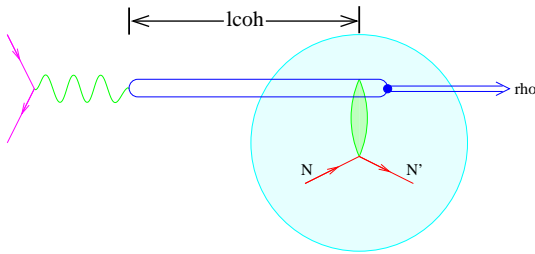


$b_{coh} \approx r_A^2/3$ is in agreement with world data of nuclear size measurements

(H.Alvensleben et al, Phys.Rev.Let. **24**,792 (1970)).

Kinematics of exclusive ρ^0 matches dimension of Nuclei

- radius of the nucleus: $r_{14N} \simeq 2.5$ fm
- coherence length: distance traversed by $q\bar{q}$ fluctuation of the virtual photon:



$$l_c = \frac{2 \cdot \nu}{Q^2 + m_V^2} = 0.6 \div 8 \text{ fm},$$

$$\langle l_c \rangle = 2.7 \text{ fm}$$

- ρ^0 formation length: after $q\bar{q}$ is put on-shell it will evolve into ρ^0 over distance:

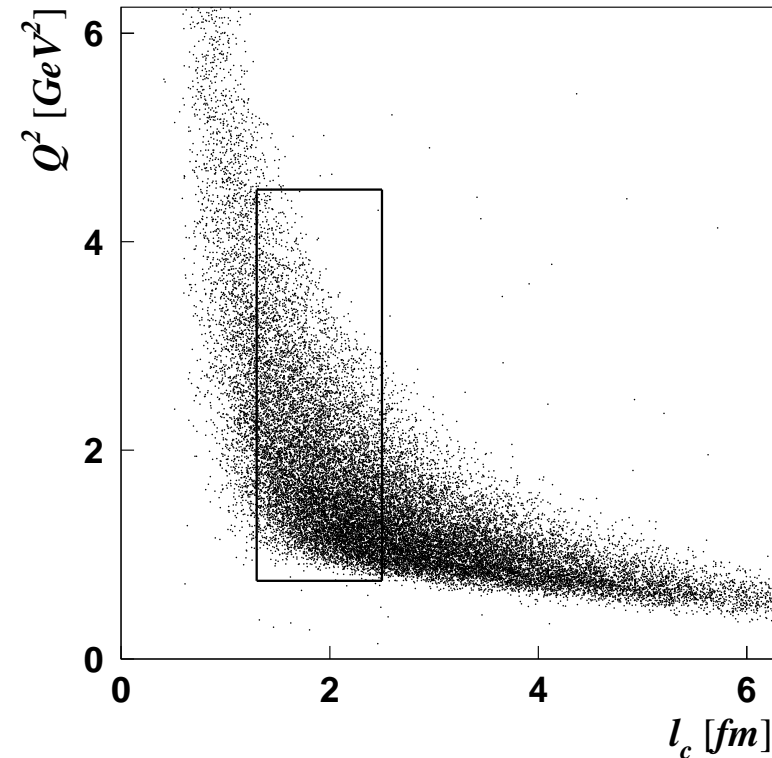
$$l_{form} = \frac{2 \cdot \nu}{m_V^2 - m_V^2} = 1.3 \div 6.3 \text{ fm}$$

$$\langle l_{form} \rangle = 3.47 \text{ fm}$$

- color transparency: transverse size of the $q\bar{q}$ of the wave packet

$$r_{q\bar{q}} \sim 1/\langle Q \rangle \simeq 0.4 \text{ fm}, r_p = 1 \text{ fm}$$

Q^2 vs l_c of exclusive ρ^0 at HERMES:



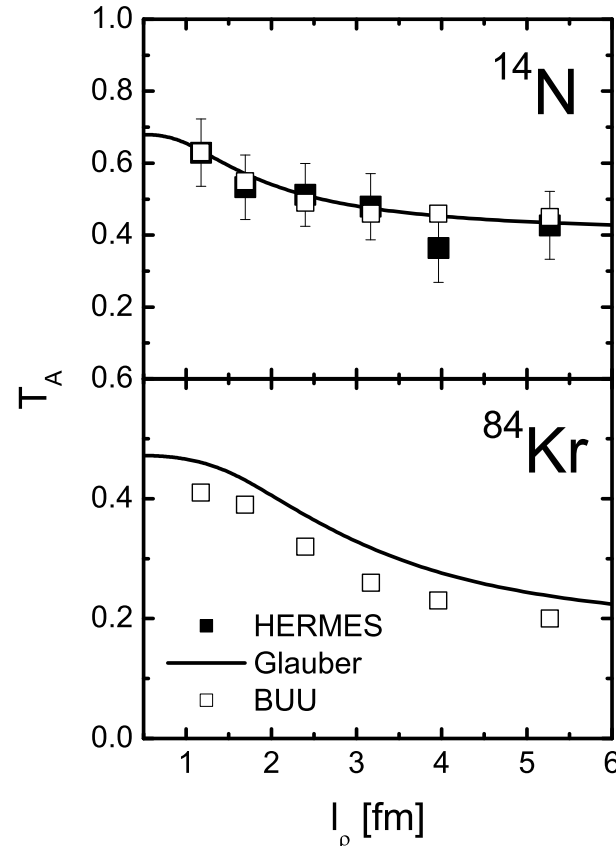
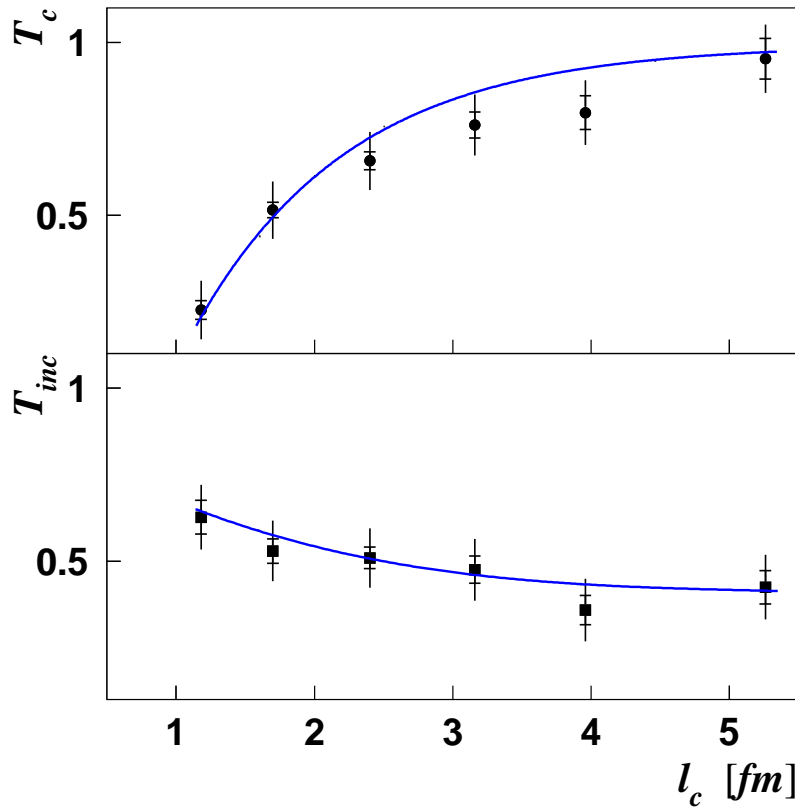
$\rightarrow \rho^0$ absorption at $l_c \leq r_{14N}$

\Rightarrow 2-dimensional analysis of Q^2 , l_c dependencies

Coherent Length Effect

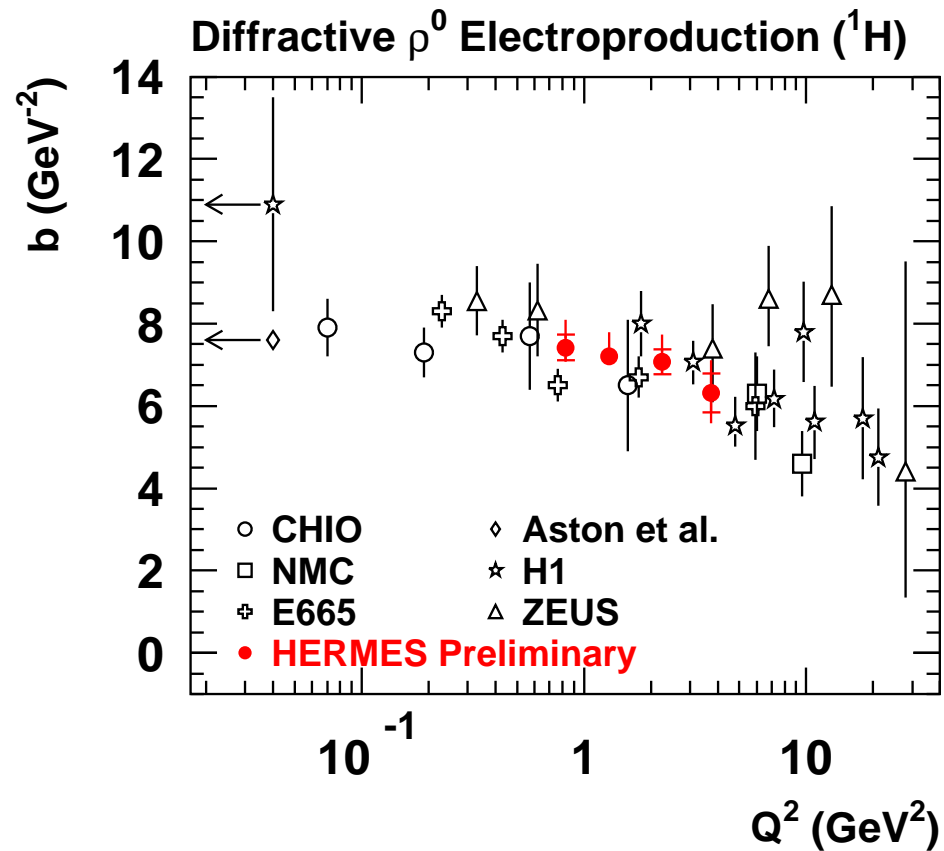
(HERMES collab., Phys. Rev. Let. **82**, 1999; Phys.Rev.Let., **90**, 5, 2003)

$$T_{c/inc}(l_c) = \frac{\sigma_{Ac/inc}}{A\sigma_H} = \frac{N_{Ac/inc} \cdot L_H}{A \cdot N_H \cdot L_A}, \quad A = {}^{14}\text{N}$$



- Left panel. Combined effect of initial and final state interactions for incoherent ρ^0 and additional effect of nuclear formfactor for coherent ρ^0 . Agreement with calculations (blue curves) based on CT approach (B.Z. Kopeliovich et al, Phys.Rev. C, **65**, 035201, 2002).
- Right panel. Calculations for incoherent production of semi-classical transport model without CT presented on right panel. (T.Falter, W.Cassing, K.Gallmeister and U.Mosel, nucl-th/0309057).

'Photon Shrinkage' ($b_{inc}(Q^2)$) a Prerequisite for Color Transparency



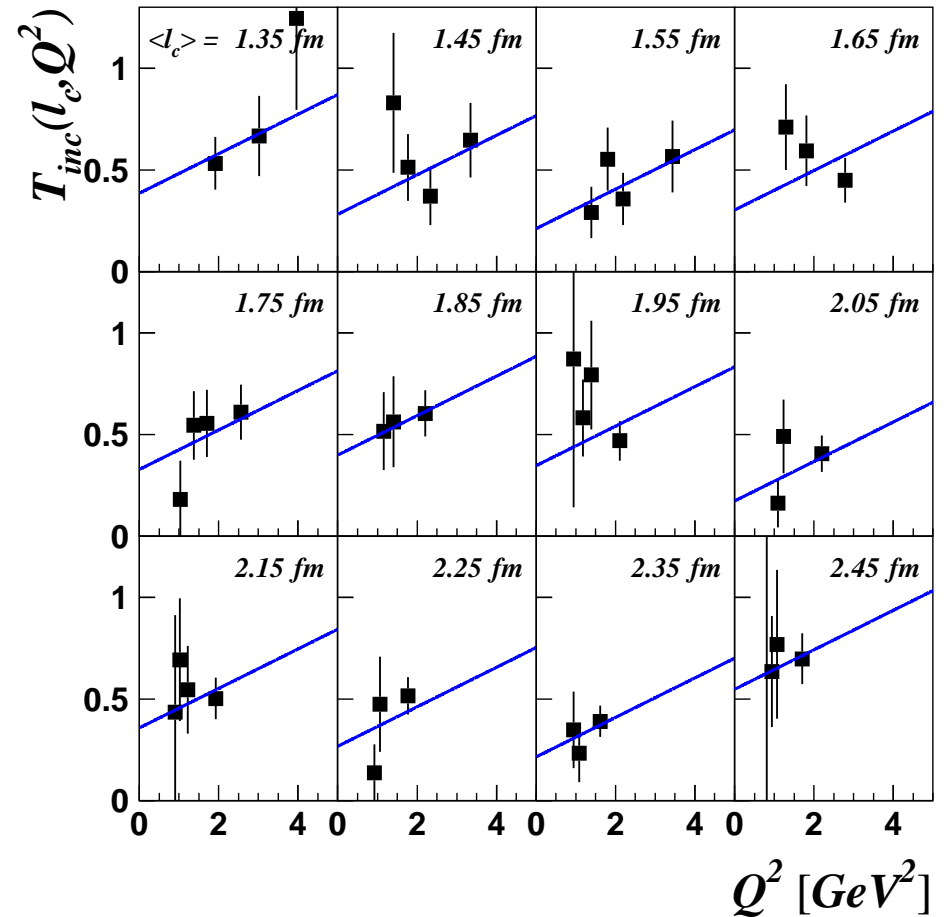
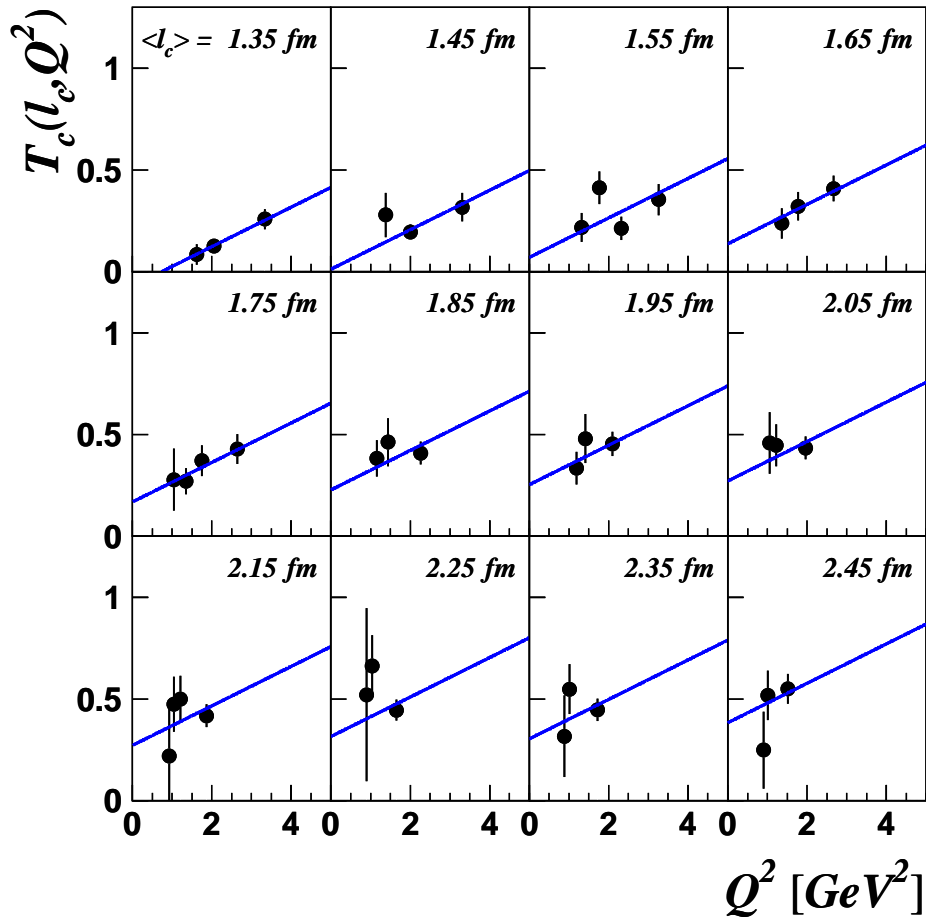
- Size of virtual photon controlled via Q^2
- No strong W -dependence

Color Transparency Effect (HERMES collab., Phys.Rev.Let.,90,5,052501,2003)

The QCD factorization theorem rigorously not possible without the onset of the color transparency: (J.Collins,L.L.Frankfurt,M.Strikman Phys.Rev.D56,2982 (1997); M.Strikman, Nucl.Phys.A663&664,64,2000)

- $r(qq)$ decreases with the increase of Q^2
- $Tr^A(Q^2, l_{coh}) = \sigma_{(in)coh}^A / \sigma^H$ grows with Q^2

At fixed l_{coh} :



data	Slope of Q^2 -dependence, GeV^{-2}	Prediction, GeV^{-2}
N incoh.	$0.089 \pm 0.046_{st} \pm 0.020_{syst}$	0.060
N coh.	$0.070 \pm 0.027_{st} \pm 0.017_{syst}$	0.048
N combined	0.074 ± 0.023	0.058

Agreement with theoretical calculations where positive slope of Q^2 -dependence was derived from the onset of the color transparency effect (B.Z. Kopeliovich et al, Phys.Rev. C, **65**, 035201, 2002)

\Rightarrow Different CT effects are predicted for ep , πp , pp and pA collisions

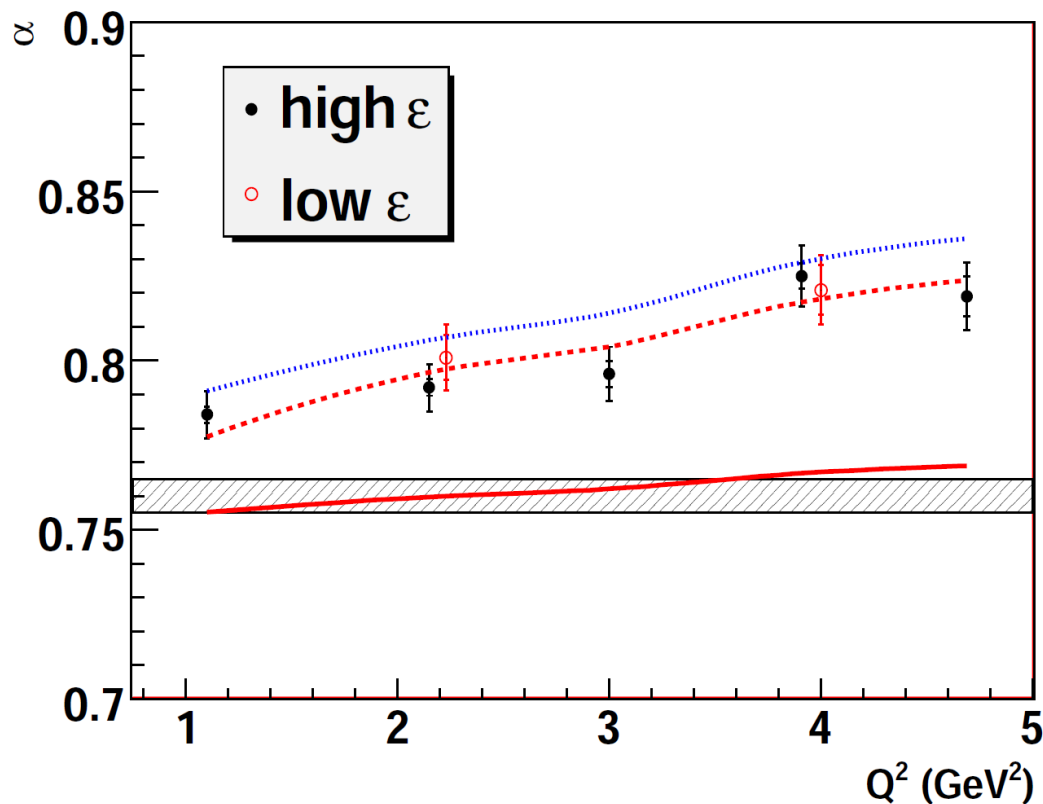
see, e.g.: G.Miller and M.Strikman talks on the workshop "Hard Exclusive Processes", Munich, 9-11 November 2009.

Color Transparency Effect from Jefferson Lab data: $e + p \rightarrow e' + \pi^+ + (n)$

(X.Qian et al, Phys.Rev.C, **81**, 055209, 2010) Kinematics:

$$4 \leq E^{beam} \leq 5.8 \text{ GeV} \quad 0.9 \leq E^{e'} \leq 1.7 \text{ GeV} \quad W \approx 2 \text{ GeV}$$

$$\text{Transparency } T = \frac{\sigma^A}{A\sigma^N} = \frac{A^\alpha \sigma^N}{A\sigma^N} = A^{\alpha-1}$$



Parameter α was measured from ^1H , ^2H , ^{12}C , ^{27}Al , ^{63}Cu and ^{197}Au data and compared with α extracted from pion-nucleus scattering data (hatched band).

\Rightarrow Noticeable dependence of α on Q^2 measured and compared with calculations: **Glauber (solid)**, **Glauber+CT (dashed)**, **Glauber+CT with short-range correlations (dotted)**.

ϵ is the longitudinal polarization of the virtual photon measured from e' kinematics

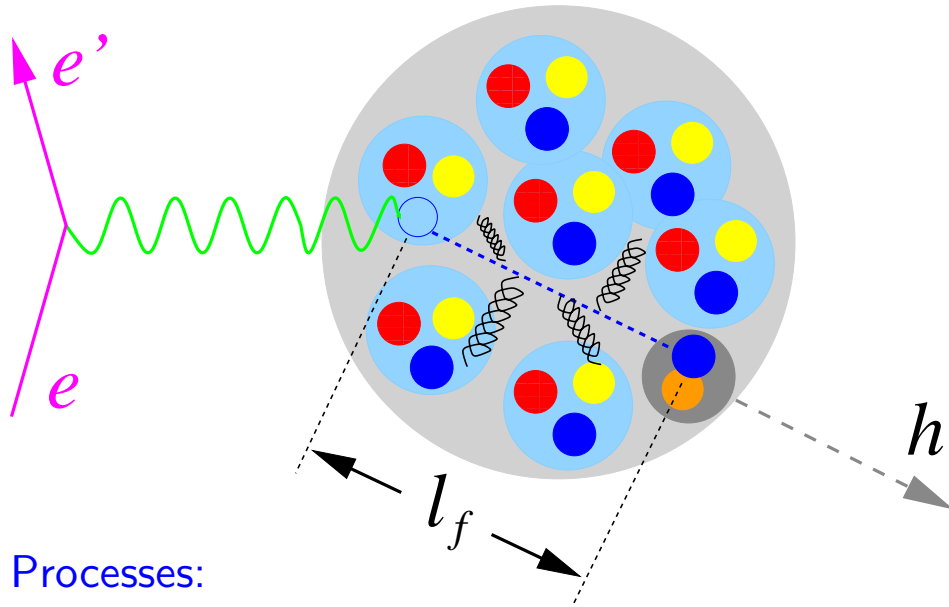
HERMES SIDIS Data on Heavy Targets

- Kinematic cuts:
 - $1 < Q^2 < 10 \text{ GeV}^2$
 - $W^2 > 10 \text{ GeV}^2$, no resonances
 - $\nu < 23.5 \text{ GeV}$, radiative effects
 - $0.2 < z < 1.0$

Target	$N^{\pi^+} \times 10^3$	$N^{\pi^-} \times 10^3$	$N^{K^+} \times 10^3$
^2D	1781	1445	356
^4He	134	107	27
^{20}Ne	380	303	82
^{84}Kr	321	260	72
^{131}Xe	193	157	44

- Corrections:
 - Acceptance and smearing
 - Radiative effects
 - Nuclear transparency for exclusive ρ^0 production at one π^\pm accepted
- To study contributions to the transverse momentum distributions from:
 - primordial transverse momentum
 - gluon radiation of the struck quark
 - multiple interactions of the “pre-hadron”
 - interactions of the formed hadrons

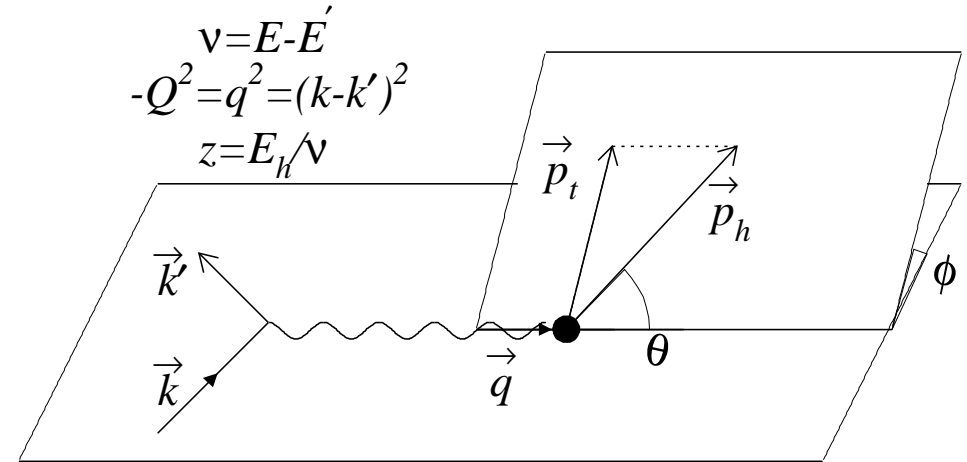
Hadron Electroduction in Cold Nuclear Matter



Processes:

- Scattering on a quark
- Quark energy loss in nuclear medium
- Fragmentation of quark to hadron (pre-hadron formation)
- Absorption of (pre-)hadron

Formation length from DIS to final hadron in various models for (pre-)hadron formation and absorption: $l_f \propto \nu(1 - z)/k$



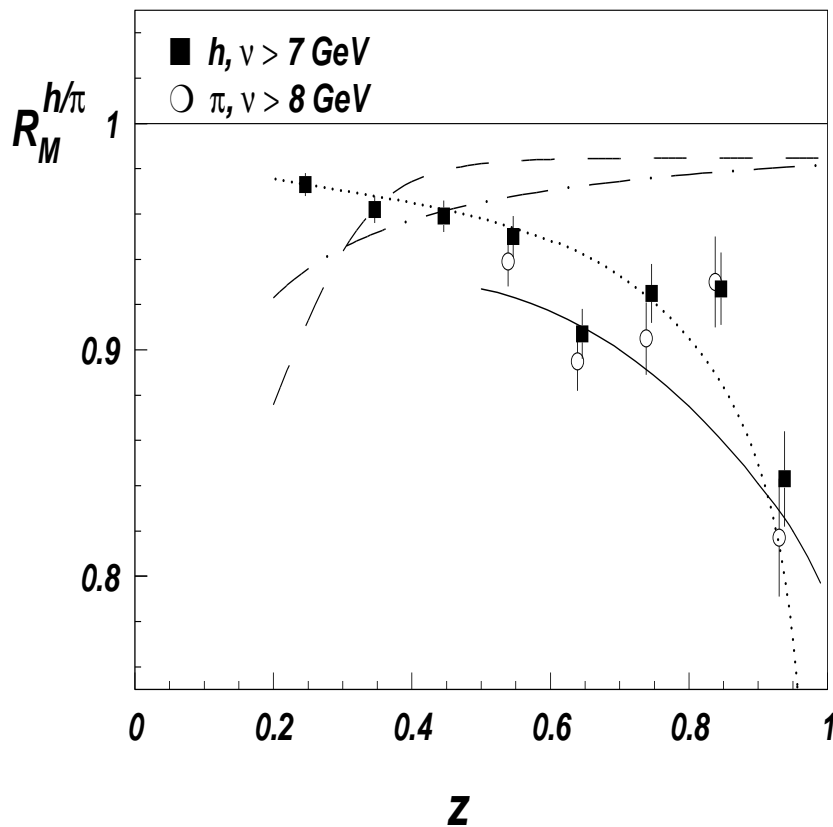
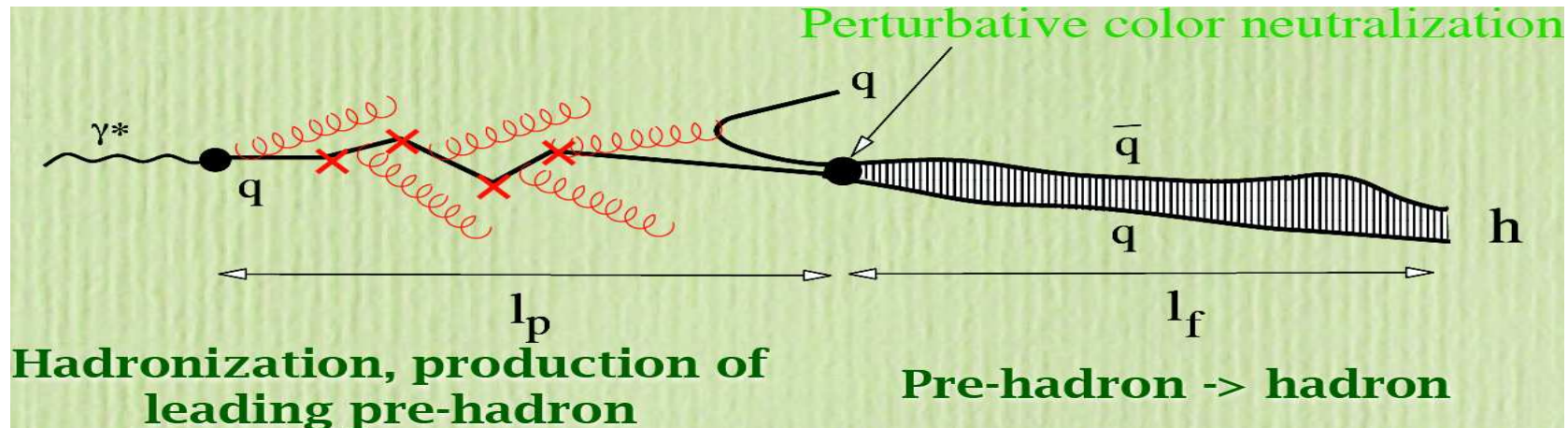
Experimental access to the hadronization through the measured multiplicity ratio:

$$R_A^h(\nu, Q^2, z, p_t^2) = \frac{\left(\frac{N^h(\nu, Q^2, z, p_t^2)}{N e'(\nu, Q^2)} \right)_A}{\left(\frac{N^h(\nu, Q^2, z, p_t^2)}{N e'(\nu, Q^2)} \right)_D}$$

data integrated over ϕ and Θ , see e.g.:

HERMES collab., Nucl. Phys. B 780 (2007) 1-27; arXiv:0704.3270

Hadron Formation in DIS in a Nuclear Environment



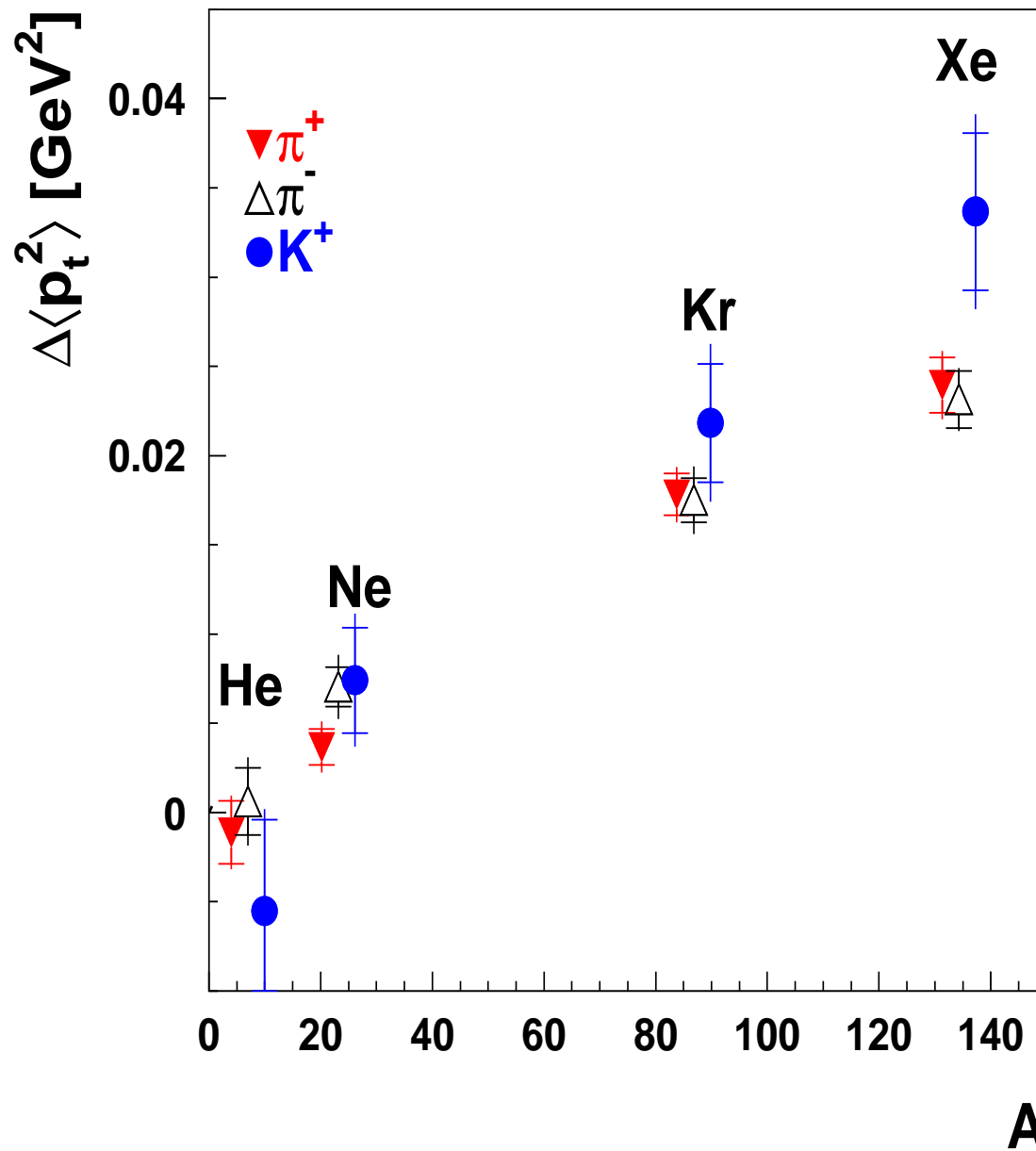
\Rightarrow Absorption is large at $z \rightarrow 1$, where $l_f \propto \nu(1-z)/k$ is short.

The fractional energy z_h is directly measured and compared with the prediction of B.Z.Kopeliovich, J.Nemchik, E.Predazzi in 1995, see S.Domdey, D.Gruenewald, B.Z.Kopeliovich and H.J.Pirner, Nucl.Phys. A 825, 200, 2009.

\Rightarrow Fast hadrons have a relatively short formation time, leading to the reduction of the ratio.

A-dependence of Nuclear p_t -broadening: $\Delta\langle p_t^2 \rangle^h = \langle p_t^2 \rangle_A^h - \langle p_t^2 \rangle_D^h$

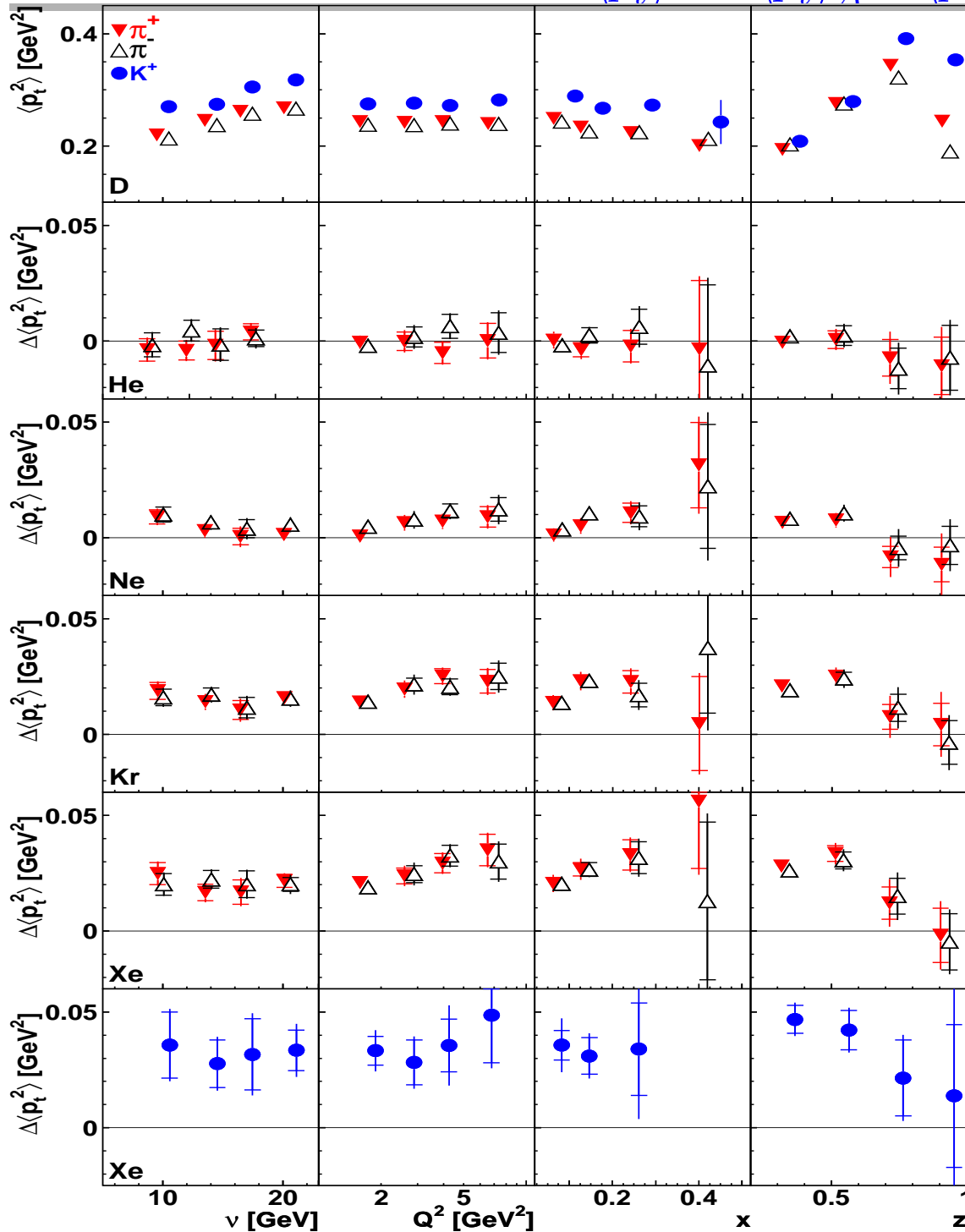
(HERMES collab., Phys. Lett. B 684 (2010) 114-118;; arXiv:0906.2478, hep-ex)



- Effect of 1-10 % relative $\langle p_t^2 \rangle_D^h \approx 0.25 \text{ GeV}^2$:
 \Rightarrow Direct evidence of p_t -broadening in leptonproduction

Kinematic Dependencies of Nuclear p_t -broadening:

$$\Delta \langle p_t^2 \rangle^h = \langle p_t^2 \rangle_A^h - \langle p_t^2 \rangle_D^h$$



- Broadening goes slightly down with increasing ν

- Broadening goes up with Q^2

- Clear $z = E_h/\nu$ dependence. At $z \rightarrow 1$

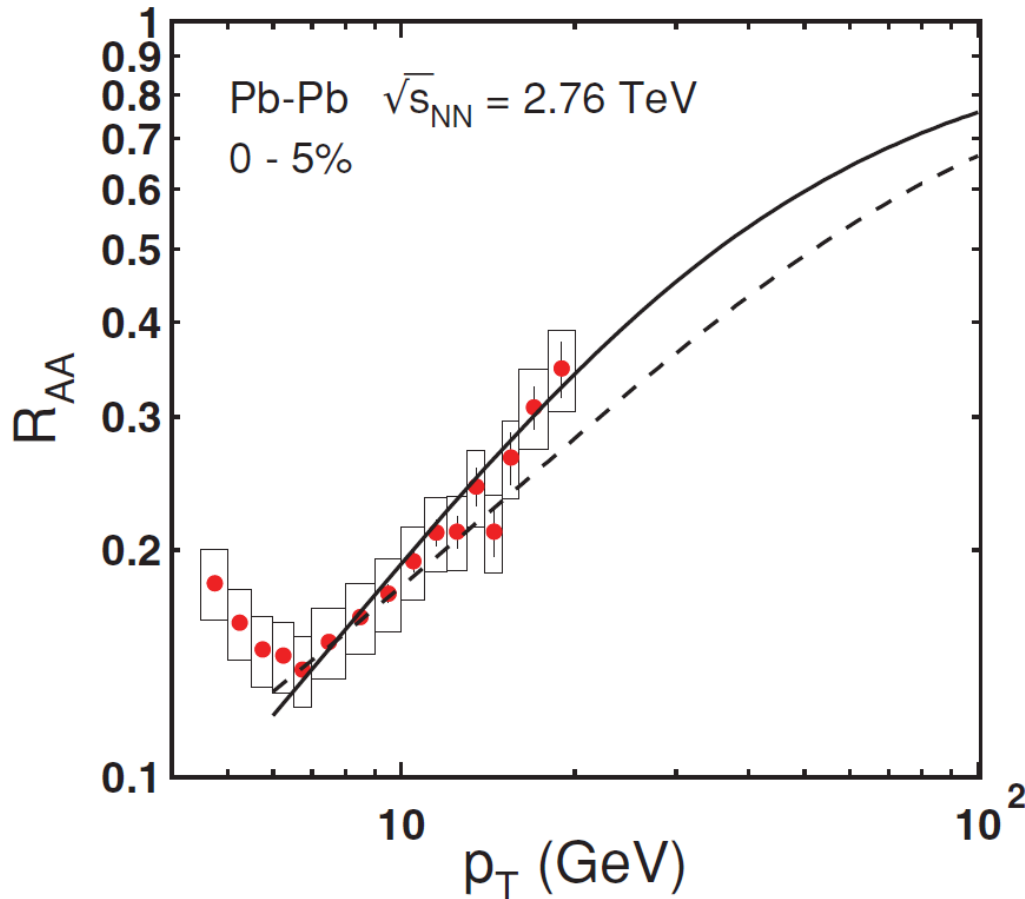
- p_t -broadening consistent with zero due to energy conservation: the struck quark cannot have lost energy
- it is not due to elastic scattering of (pre-)hadrons

(B.Z.Kopeliovich et al, Nucl.Phys. A740 (2004) 211; S.Domdey,D.Gruenewald, B.Z.Kopeliovich and H.J.Pirner, Nucl.Phys. A 825, 200, 2009.)

\Rightarrow Constraint on pre-hadron production mechanism

The R_{AA} data from ALICE as “Color transparency and suppression of high- p_T hadrons in nuclear collisions”

see B.Z.Kopeliovich, I.K.Potashnikova and I.Schmidt, Phys.Rev. C. 83, 021901 (R), 2011.



data from ALICE collab., Phys. Lett. B 696 (2011) 30-39

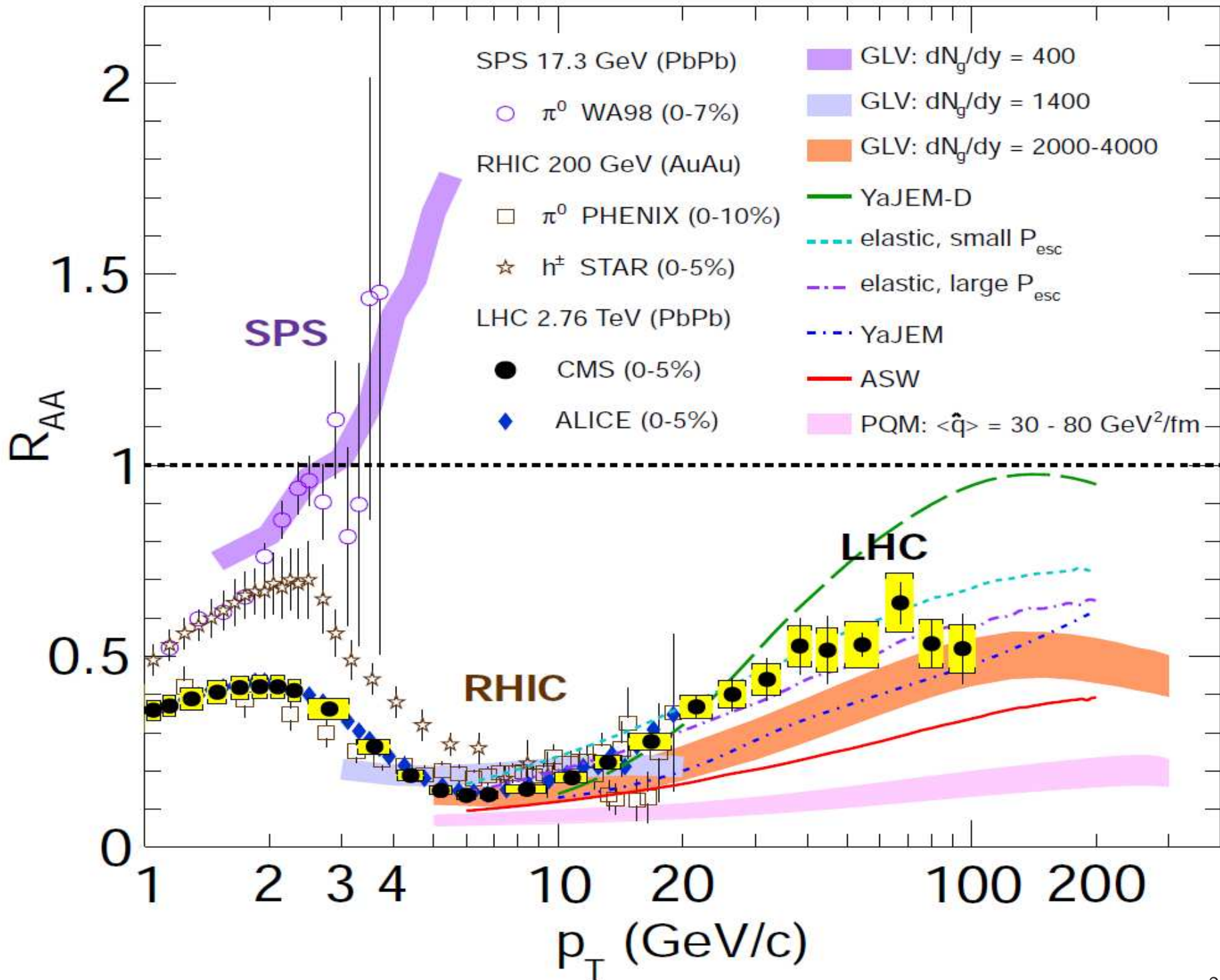
Calculations include:

- propagation of colorless dipole (pre-hadron) through a dense medium,
- its short production length,
- intensive energy loss at the early stage.

Data on high- p_T hadron production in heavy ion collisions include:

- the uncertainties from the variation in space and time of the medium properties,
- unmeasured fractional energy z_h convoluted with initial parton distribution, hard cross section and the fragmentation function,
- the contribution of the initial state effects in cold nuclear models.

⇒ Global analysis of experimental data to test the models is desirable...



Summary

- Data on nuclear transparency of exclusive ρ^0 production on ^{14}N are in agreement with Color Transparency based calculations.
 - \implies Phenomenon of color transparency, based on factorization theorem and small size configuration, could be used to study GPDs and their possible modifications in the cold or hot nuclear matter.
- HERMES data provide a strong support for shortness of the production length
- The first direct determination of p_T -broadening in SIDIS was performed on He, Ne, Kr and Xe targets. \implies It increases with A and remains constant with ν suggesting that effect is due to the “partonic” stage and that color neutralization happens near the surface or outside the nucleus.
- A colorless dipole created on a short production length evolves its size during propagation in medium.
- The energy dependence of this evolution and color transparency explain well the recent data from ALICE and CMS on nuclear suppression of light hadrons with large enough p_T .

Outlook: much more data on heavy targets are published by HERMES:

- *Multidimensional Study of **Hadronization** in Nuclei*, Eur. Phys. J. A 47 (2011) 113; arXiv:1107.3496.
- *Transverse momentum **broadening** of hadrons produced in semi-inclusive deep-inelastic scattering on nuclei*, Phys. Lett. B 684 (2010) 114-118; arXiv:0906.2478 (hep-ex).
- ***Nuclear-mass dependence** of beam-helicity and beam-charge azimuthal asymmetries in **DVCS***, Phys. Rev. C 81 (2010) 035202; arXiv:0911.0091 (hep-ex) and DESY-09-190
- ***Hadronization** in Semi-inclusive deep inelastic scattering on nuclei*, Nucl. Phys. B 780 (2007) 1-27; Eprint numbers: arXiv:0704.3270 and DESY-07-050
- ***Double hadron leptonproduction** in the nuclear medium*, Phys. Rev. Lett, 96 (2006) 162301; Eprint numbers: hep-ex/0510030 and DESY-05-205
- ***Quark Fragmentation** to $\pi^{+/-}$, π^0 , $K^{+/-}$, p and \bar{p} in the Nuclear Environment*, Phys. Lett. B 577 (2003) 37-46; Eprint numbers: hep-ex/0307023 and DESY-03-088
- *Q^2 -dependence of **Nuclear Transparency** for coherent and incoherent ρ^0 production*, Phys. Rev. Lett. 90 (2003) 052501; Eprint numbers: hep-ex/0209072 and DESY-02-152
- *Observation of a **Coherence Length Effect** in Exclusive ρ^0 Electroproduction*, Phys. Rev. Lett. 82 (1999) 3025-3029; Eprint numbers: hep-ex/9811011 and DESY-98-178