



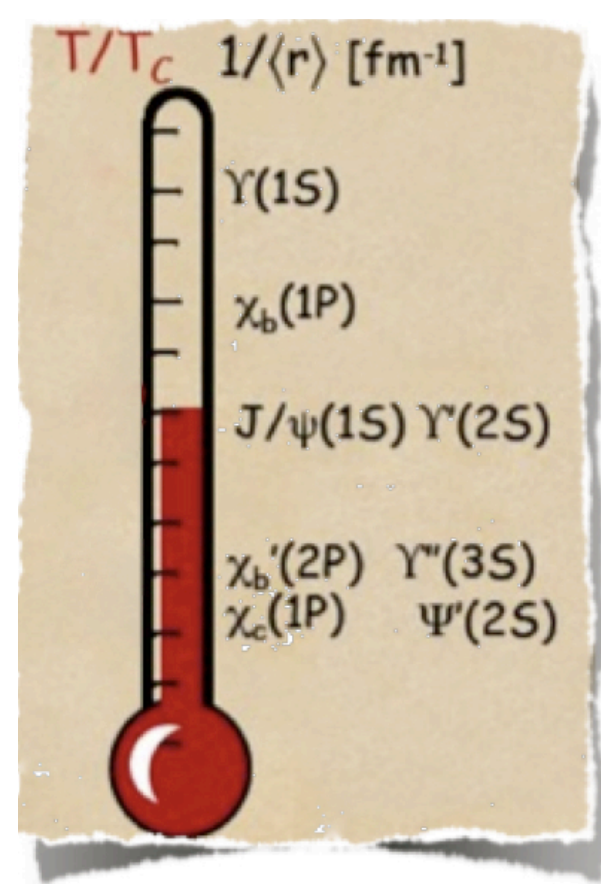
# Studies of $\Upsilon$ states in PbPb collisions at $\sqrt{s_{NN}} = 2.76$ TeV at LHC

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## Introduction

- In our universe today, quarks are always bound together by gluons to form "composite" particles. The Quark-Gluon Plasma (QGP) is a hot, dense state in which these quarks and gluons exist freely, unbound. This is thought to be the situation a few millionths of a second after the Big Bang.
- One of the predicted characteristics of the QGP is that its high temperature causes the "melting" of quarkonia. This melting manifests itself as the suppression of quarkonium production in heavy-ion collisions, compared to the number of quarkonia produced in collisions between protons.
- Detailed measurements of bottomonium production will help characterize the dense matter produced in heavy-ion collisions.



State	$J/\psi$ (1S)	$\chi_c$ (1P)	$\psi'$ (2S)
$m$ (GeV/c <sup>2</sup> )	3.10	3.53	3.68
$r_0$ (fm)	0.50	0.72	0.90

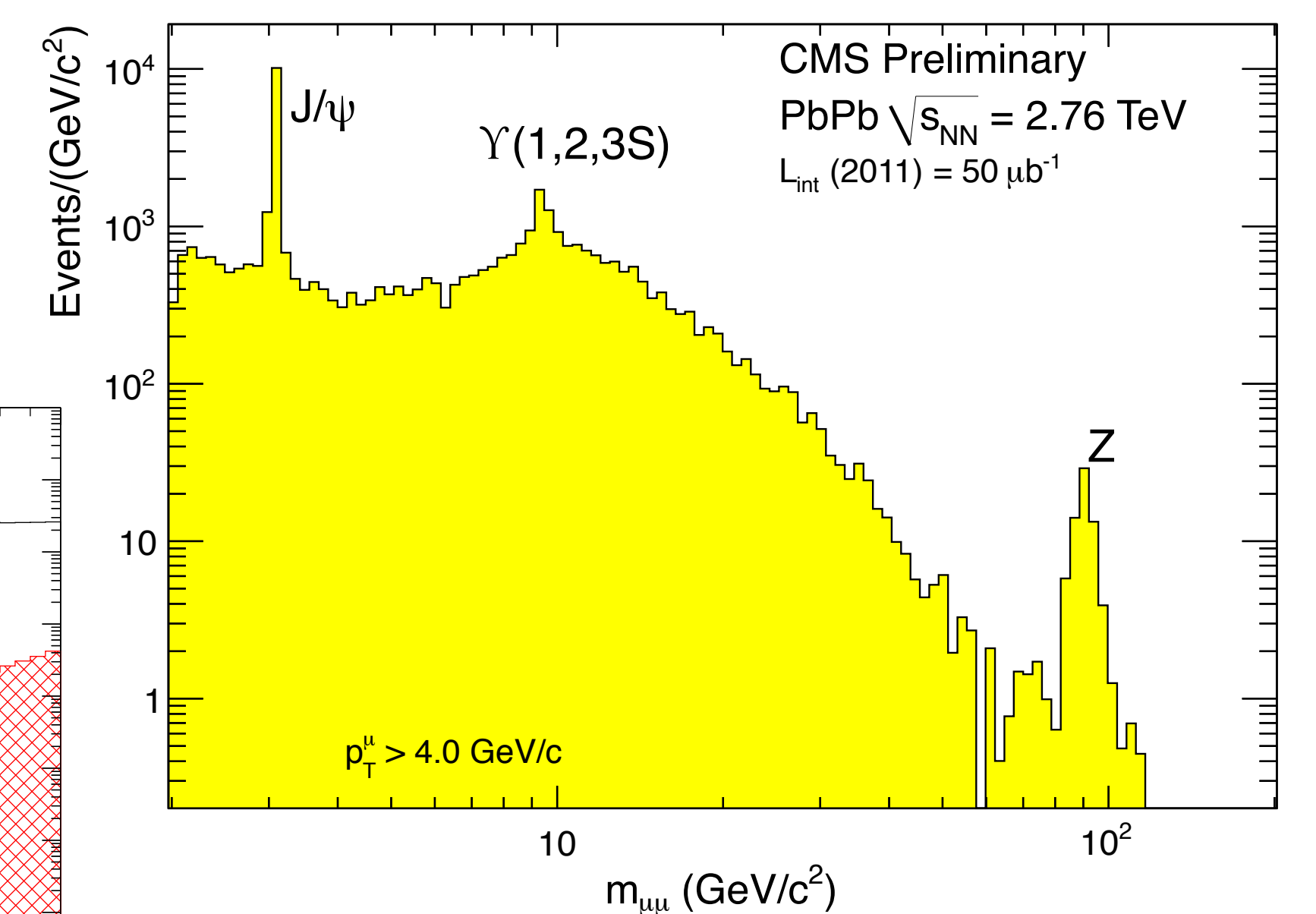
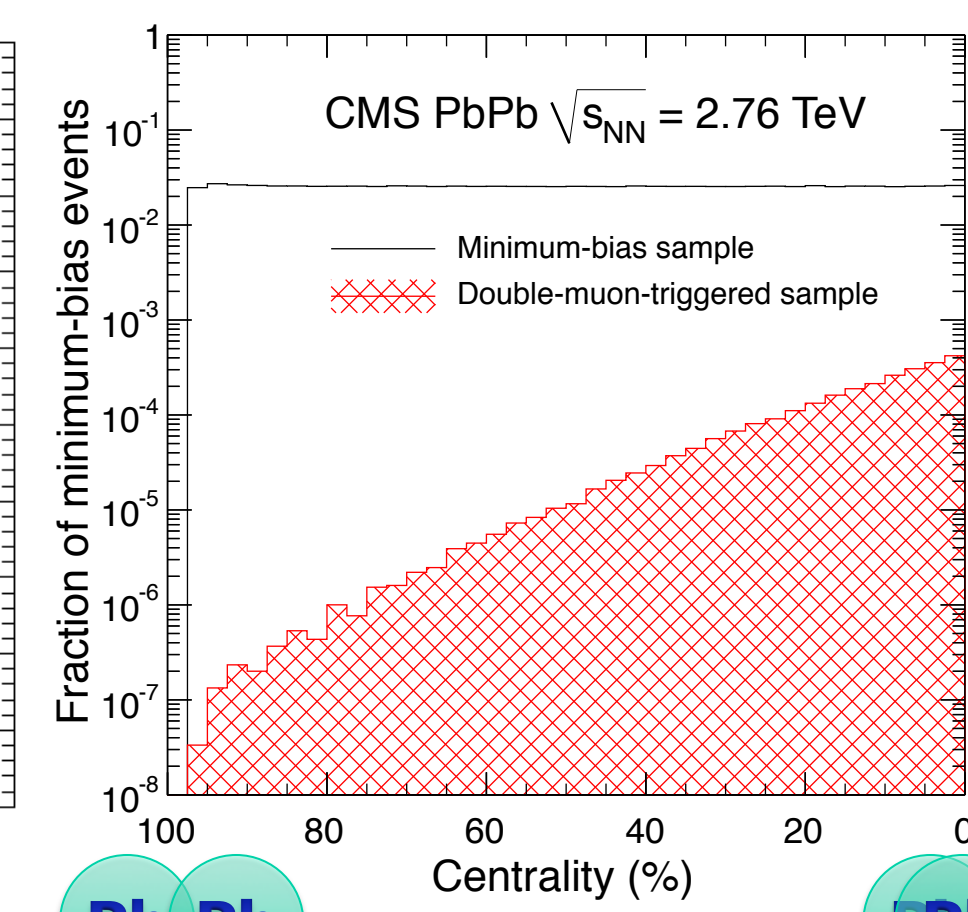
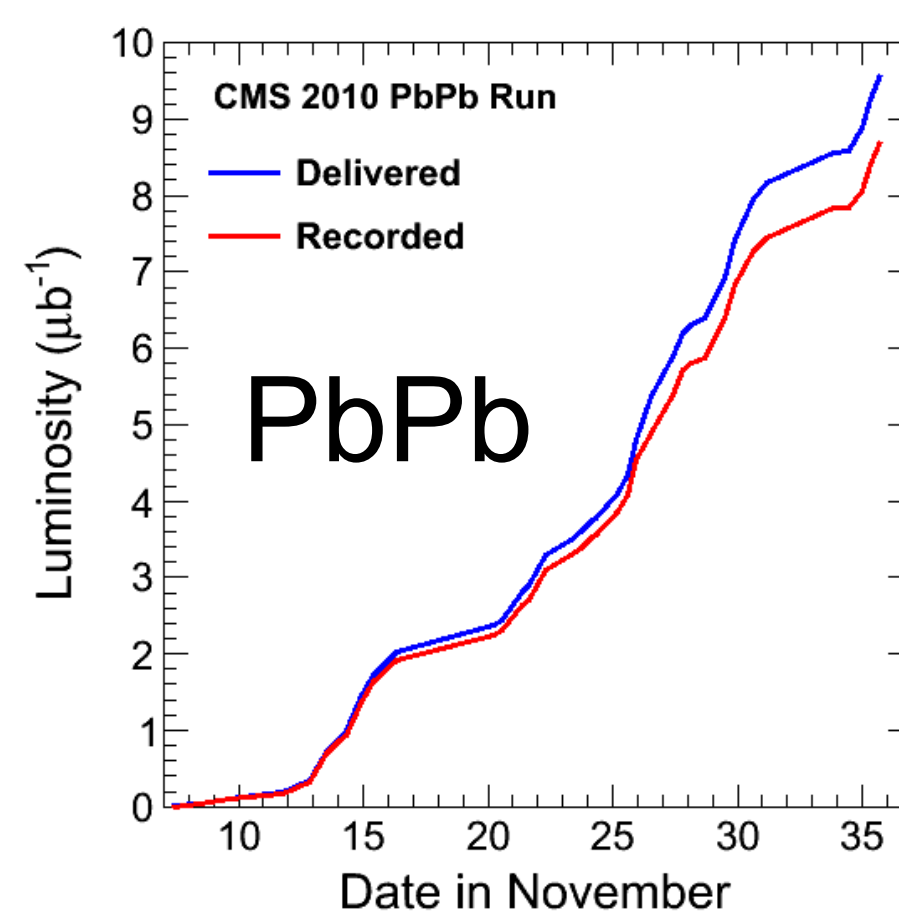
$\Upsilon$ (1S)	$\chi_b$ (1P)	$\Upsilon'$ (2S)	$\chi_b'$ (2P)	$\Upsilon''$ (3S)
9.46	9.99	10.02	10.26	10.36
0.28	0.44	0.56	0.68	0.78

decreasing binding energy

J.Phys.G32:R25,2006

## Quarkonia in PbPb Collisions in CMS

- PbPb data sample: PbPb run 2010 @  $\sqrt{s_{NN}} = 2.76$  TeV ( $L_{int} = 7.28 \mu\text{b}^{-1}$ )
- pp reference sample: pp run 2011 @  $\sqrt{s} = 2.76$  TeV ( $L_{int} = 231 \text{nb}^{-1}$ )
- Event selection:
  - Online dimuon trigger
  - Offline minBias selection
- Muon selection:
  - Muon quality cuts
  - Kinematic cut:  $p_T^\mu > 4$  GeV/c

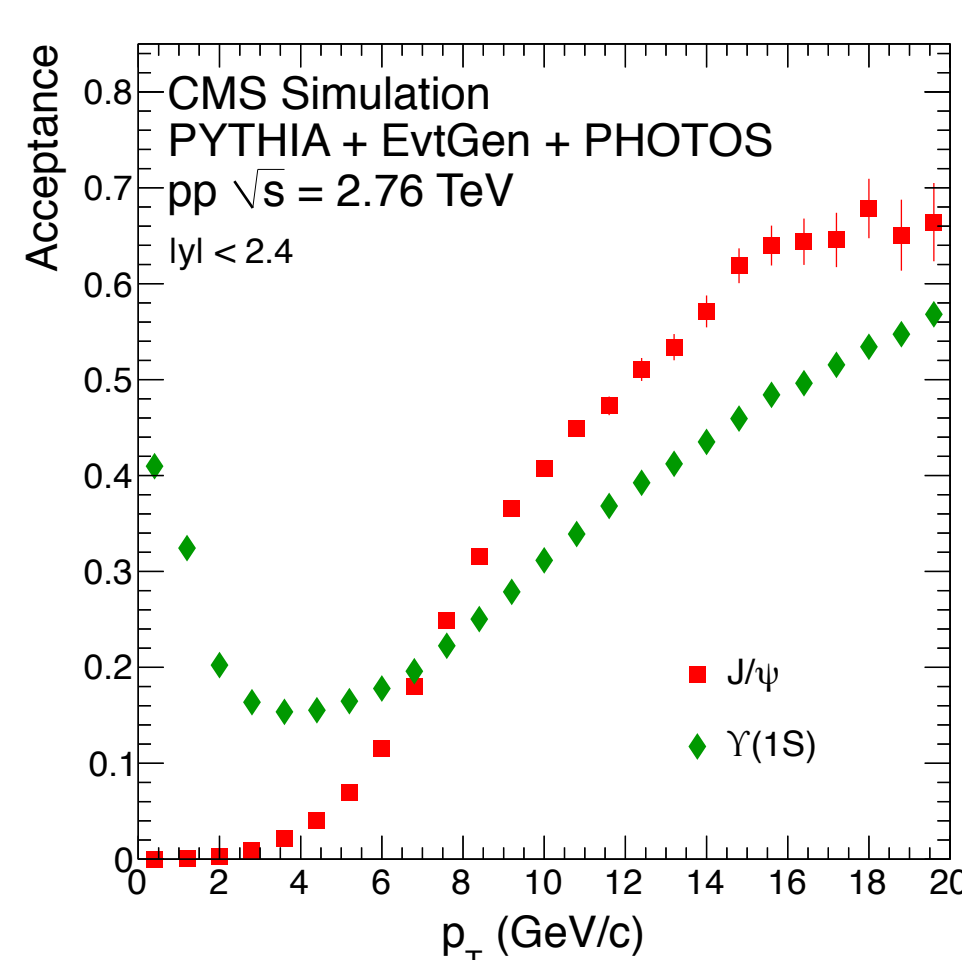


CMS has excellent mass resolution

## Nuclear Modification Factor $R_{AA}$ (arXiv:1201.5069)

### Acceptance

$$\alpha(p_T, y; \lambda_\theta) = \frac{N_{reconstructible, M}^{dimuon}(p_T, y; \lambda_\theta)}{N_{|y| < 2.4}^{dimuon}(p_T, y; \lambda_\theta)}$$



**Numerator:** the number of generated events in the MC simulation, declared detectable for given acceptance cuts  
**Denominator:** the number of dimuons generated within the muon stations coverage of the CMS detector ( $|\eta| < 2.4$ )

### Signal Extraction

Extended unbinned maximum likelihood fit

**Signal**

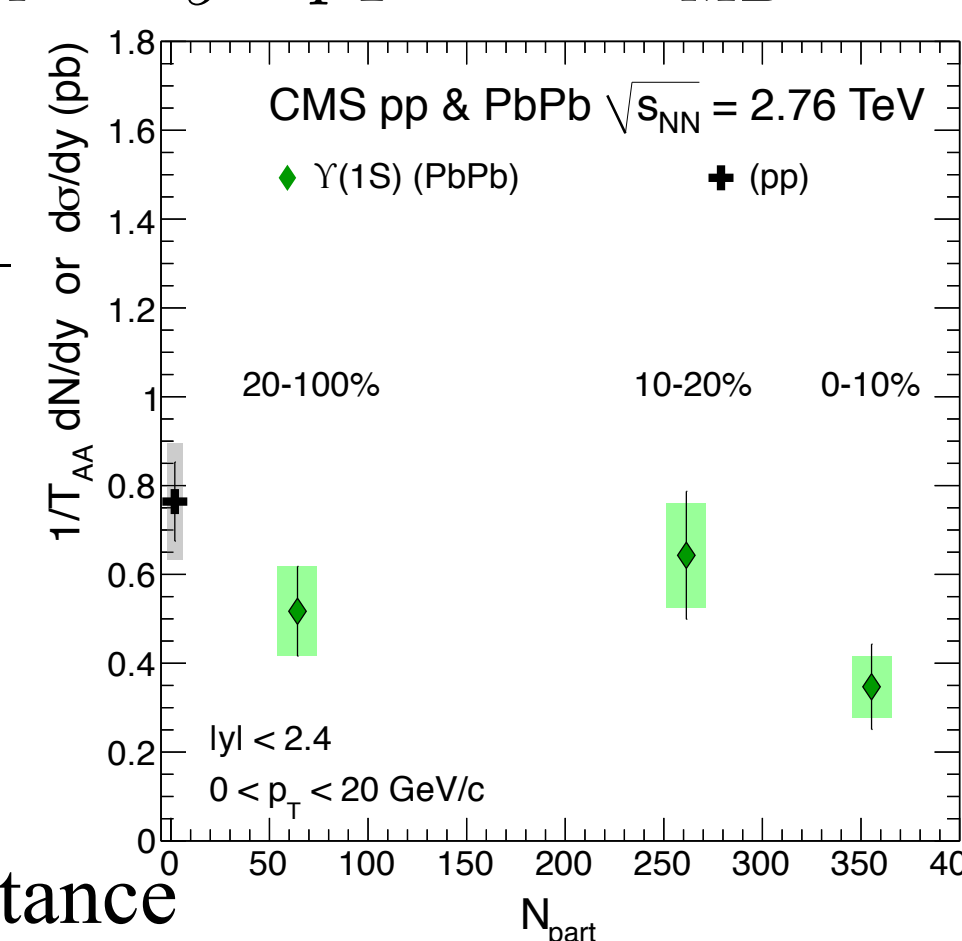
- Core Gaussian with power law tail for final state radiation
- Resolution fixed from MC simulation
- Peak separation fixed to PDG

**Background**

- Second order polynomial

### $\Upsilon(1S)$ invariant Yields in PbPb

$$\frac{1}{T_{AA}} \cdot \frac{d^2 N}{dp_T dy} = \frac{1}{T_{AA}} \cdot \frac{1}{\Delta y \Delta p_T} \cdot \frac{N_{Q\bar{Q}}}{\alpha \epsilon N_{MB}}$$



- $N_{Q\bar{Q}}$ : the number of measured  $\Upsilon$  in the  $\mu^+\mu^-$  decay channel
- $N_{MB}$ : the number of minimum bias events sampled by the event selection
- $\alpha$ : the geometric acceptance
- $\epsilon$ : the combined trigger and reconstruction efficiency
- $\Delta y$  and  $\Delta p_T$ : the bin width in rapidity and  $p_T$
- $T_{AA}$ : the nuclear overlap function (varies with the centrality of the collision and has units of  $\text{mb}^{-1}$ )

### $\Upsilon(1S) R_{AA}$

Comparing with the pp @ 2.76 TeV data, we can measure  $R_{AA}$

$$R_{AA} = \frac{\mathcal{L}_{int}^{pp} N_{MB}^{pp}}{T_{AA} N_{MB}^{PbPb}} \cdot \frac{N_{Q\bar{Q}}^{PbPb}}{N_{Q\bar{Q}}^{pp}} \cdot \frac{\epsilon_{pp}}{\epsilon_{PbPb}(\text{cent})}$$

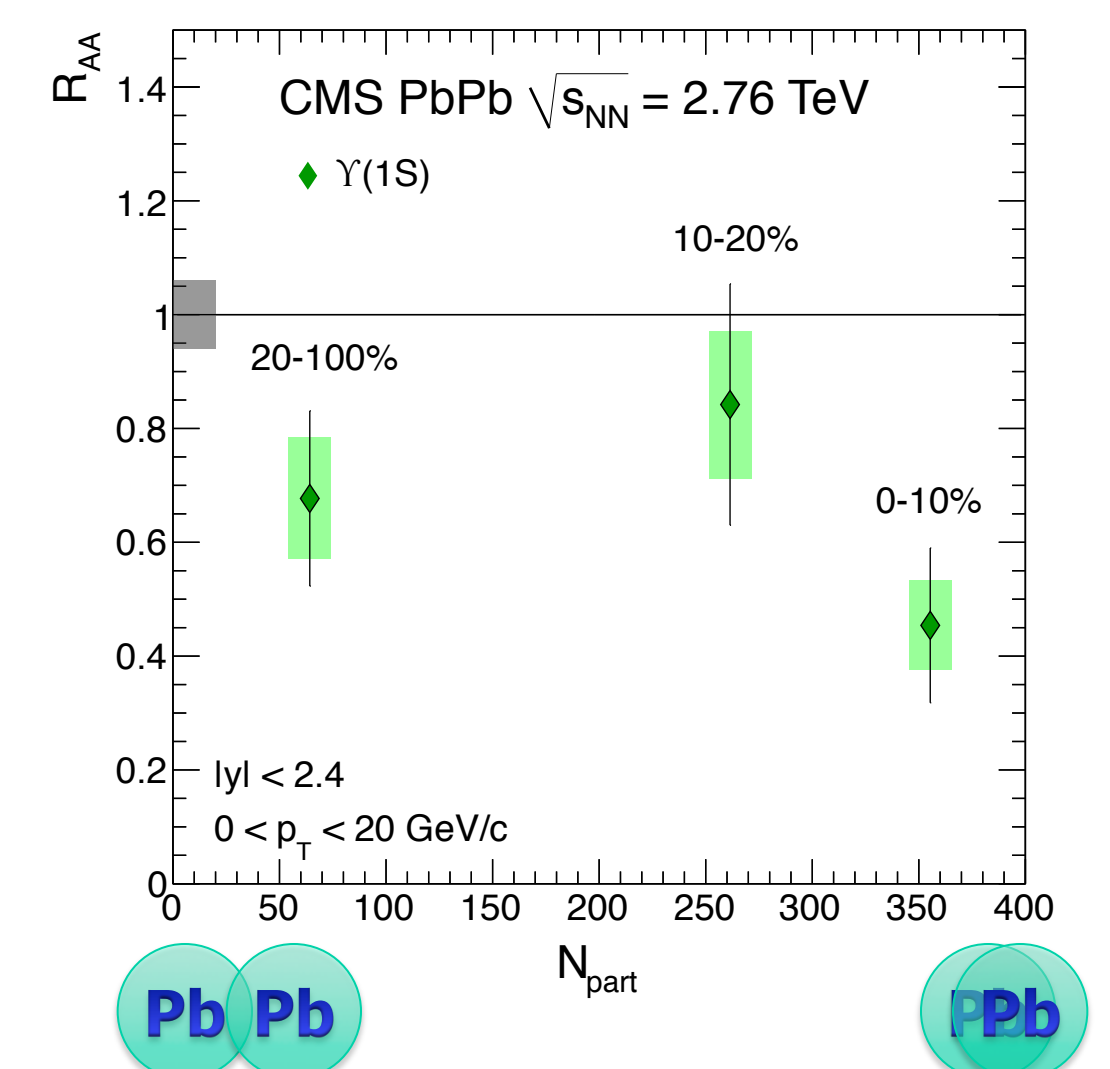
- $N_{PbPb} = 86 \pm 12[\text{stat}] \pm 3[\text{syst}]$
- $N_{pp} = 101 \pm 12[\text{stat}] \pm 3[\text{syst}]$
- $T_{AA} = 5.66 \text{mb}^{-1}$
- $N_{MB} = 55.7 \text{M MB PbPb collisions}$
- $L_{pp} = 231 \text{nb}^{-1}$

$$R_{AA} = \frac{1}{T_{AA}} \frac{dN_{AA}}{d\sigma_{pp}} \begin{cases} > 1: \text{enhancement} \\ = 1: \text{no medium effect} \\ < 1: \text{suppression} \end{cases}$$

$\Upsilon(1S) R_{AA}$  in the most central bin (0-20%):  $0.60 \pm 0.12(\text{stat.}) \pm 0.10(\text{syst.})$

$R_{AA} < 1$ : suppression

Consistent with the suppression of excited states (50% feed-down contribution measured by CDF)



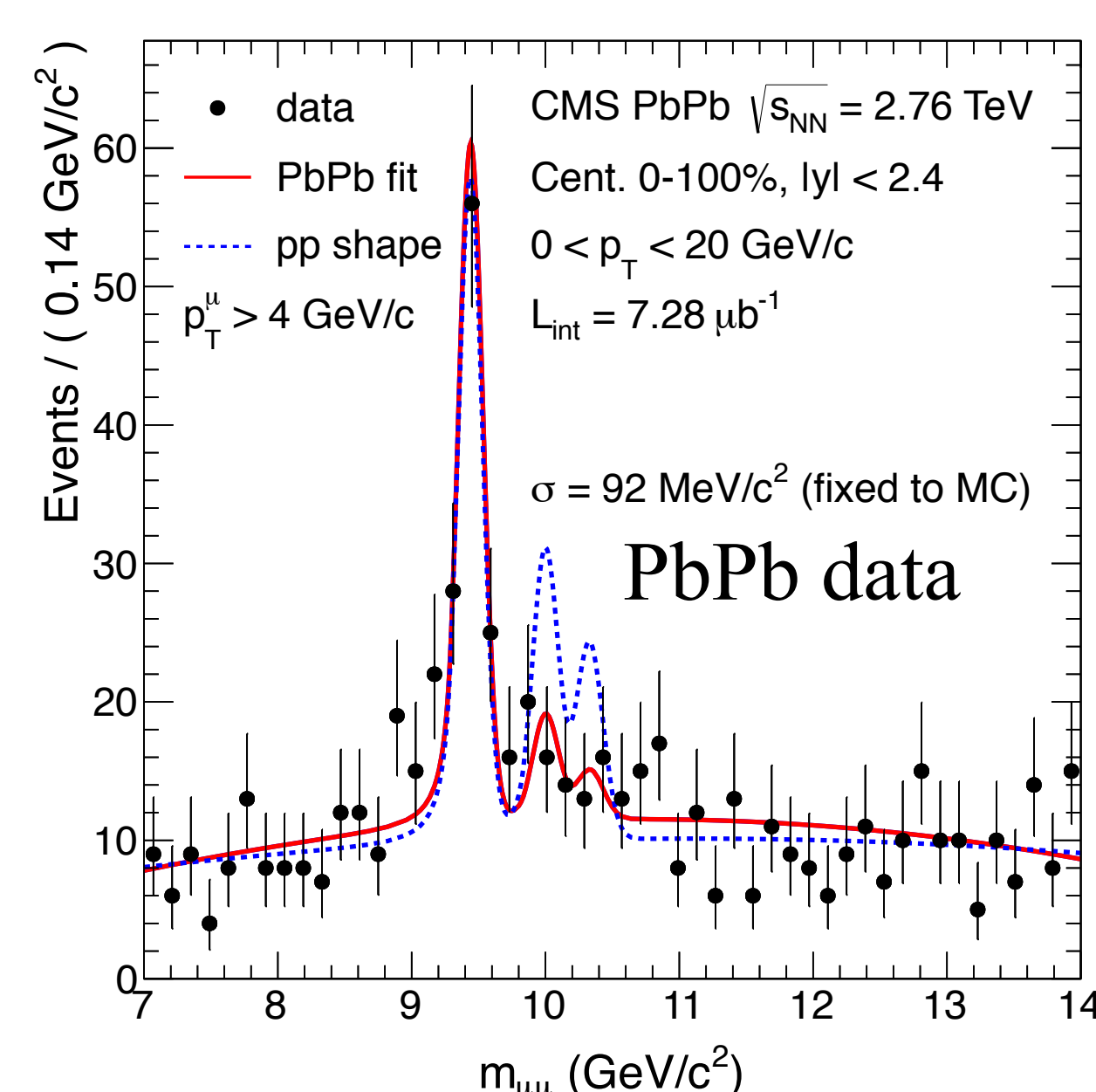
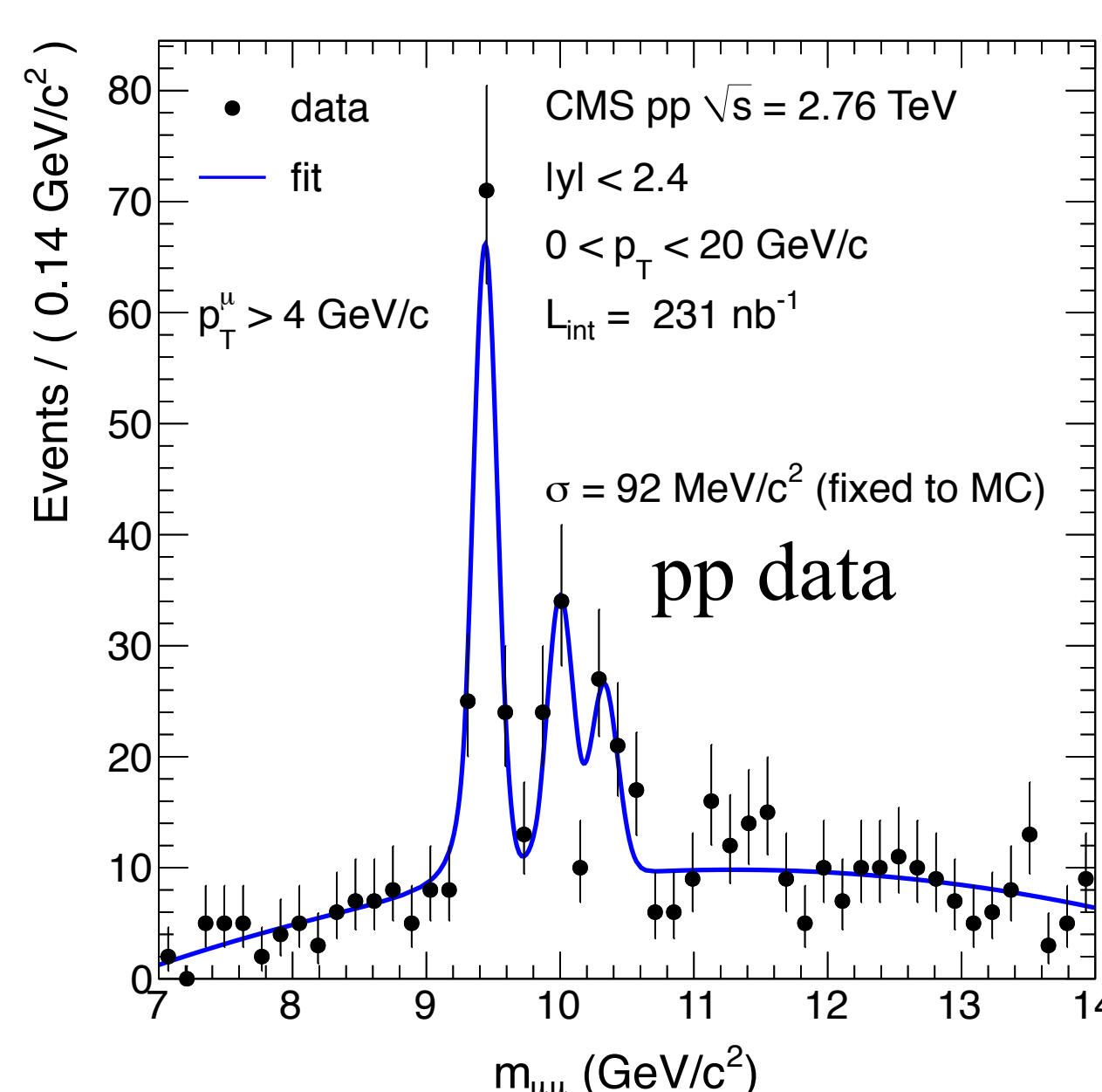
## Suppression of Excited $\Upsilon$ States in PbPb Collisions (Phys. Rev. Lett. 107, 052302 (2011))

### $\Upsilon(2S+3S)$ vs $\Upsilon(1S)$

- Measure the fraction of excited states  $\Upsilon(2S+3S)$  relative to  $\Upsilon(1S)$
- Fraction extracted directly from the fit to the PbPb and pp data sample (both at 2.76 TeV)

$$\Upsilon(2S+3S)/\Upsilon(1S)|_{pp} = 0.78^{+0.16}_{-0.14} \pm 0.02$$

$$\Upsilon(2S+3S)/\Upsilon(1S)|_{PbPb} = 0.24^{+0.13}_{-0.12} \pm 0.02$$



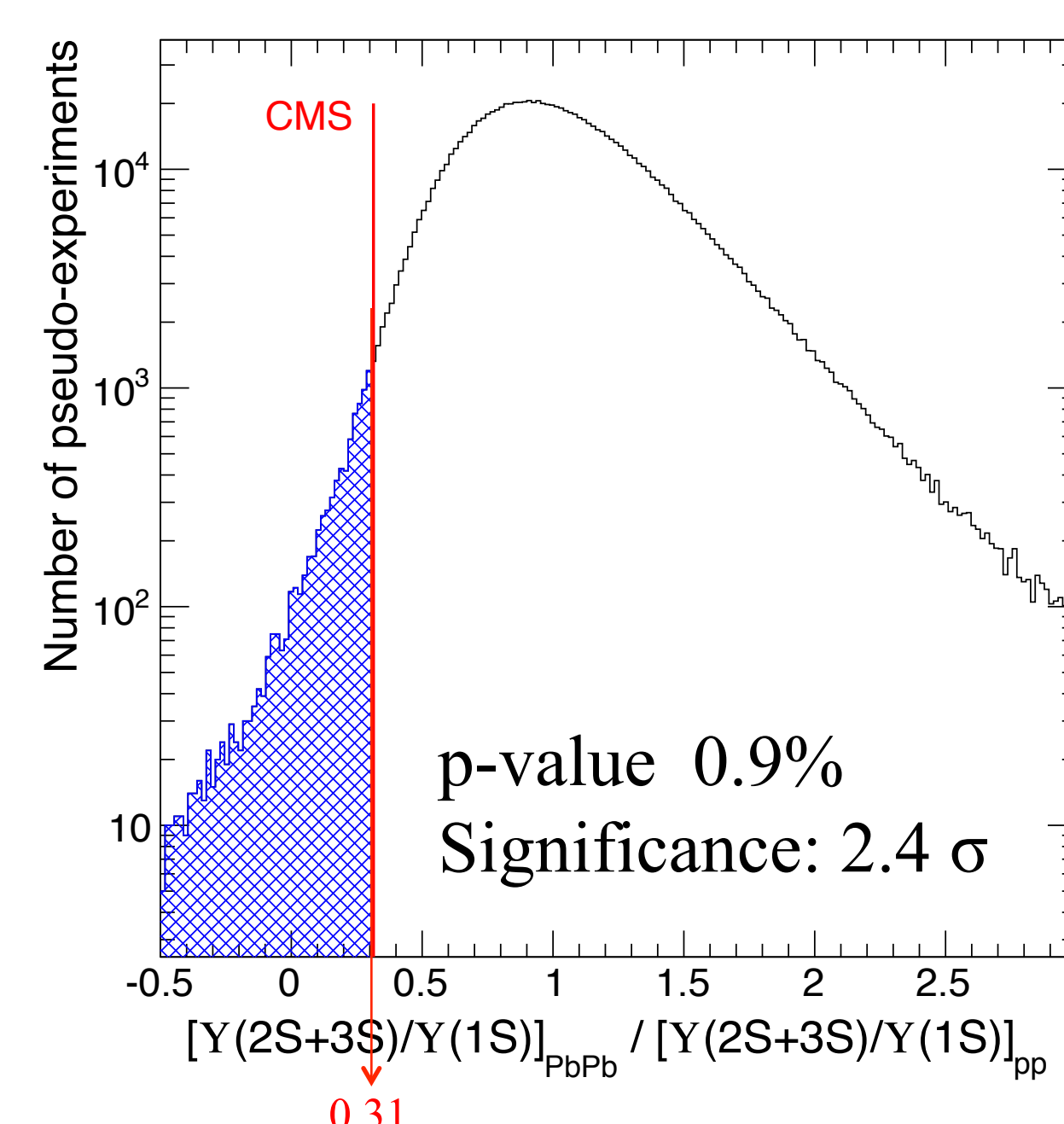
### Double Ratio

Compare ratios of  $\Upsilon(2S+3S)$  relative to  $\Upsilon(1S)$  (ground state) in PbPb & pp:

- Extract double ratio directly from simultaneous fit to both samples

$$\frac{\Upsilon(2S+3S)/\Upsilon(1S)|_{PbPb}}{\Upsilon(2S+3S)/\Upsilon(1S)|_{pp}} = 0.31^{+0.19}_{-0.15} \pm 0.03$$

- Generate pseudo-experiments following the *null-hypothesis* (i.e. no suppression) to get **p-value**



### Systematic for Double Ratio

Pros of a double ratio:

- Acceptance cancels
- Efficiency cancels

Possible differences dominated by systematic uncertainty from the fit model:

- Signal shape
- Mass resolution
- Background PDF and fit range

Total systematic uncertainty: 9.1%

### Cold Nuclear Effect

Cold nuclear matter may affect  $\Upsilon$  suppression:

- Smaller nuclear absorption cross section than at lower energy and for  $J/\psi$  (smaller size)
- Shadowing cancels in the  $\Upsilon(2S+3S)/\Upsilon(1S)$  ratio at least to the first order

### References

- "Suppression of non-prompt  $J/\psi$ , prompt  $J/\psi$ , and  $\Upsilon(1S)$  in PbPb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV", arXiv:1201.5069
- "Indications of Suppression of Excited  $\Upsilon$  States in Pb-Pb Collisions at  $\sqrt{s_{NN}} = 2.76$  TeV", PRL 107, 052302(2011)