

Hydrodynamic Fluctuations and Two Point Correlations

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Motivation

We are interested in the impact of *local* fluctuations on two particle correlation data from RHIC and LHC. In contrast to typical initial state fluctuations, here the fluctuations are localized in both transverse *and* longitudinal extent. Related recent work includes [1], which demonstrated that initial state hotspots placed at a particular position in the transverse plane could reproduce many of the features of the two-particle azimuthal correlations.

Point-like fluctuations occur naturally due to thermal noise in hydrodynamics. In [2] it was shown that local hydrodynamic fluctuations occurring throughout the evolution of the system could induce long range rapidity correlations.

Here, we extend these investigations by examining the energy-energy correlations (induced by a local perturbation propagating according to hydrodynamics) as a function of both azimuthal angular difference $\Delta\phi$, and rapidity difference, $\Delta\xi$.

Background Flow

We add perturbations on top of the Gubser flow, a generalization of Bjorken flow which includes expansion in the transverse plane [3,4]. The fluid is at rest in the following coordinate system

$$ds^2 = \tau(\rho, \theta)^2 [-d\rho^2 + \cosh^2 \rho (d\theta^2 + \sin^2 \theta d\phi^2) + d\xi^2]$$

Conformal coordinates ρ and θ are related to proper time τ and the radial distance r by the transformations. We denote the space-time rapidity as ξ .

$$\sinh \rho = -\frac{1 - (q\tau)^2 + (qr)^2}{2q\tau} \quad \tan \theta = \frac{2qr}{1 + (q\tau)^2 - (qr)^2}$$

Important properties of the Gubser flow:

- Exact, boost invariant solution of the hydro equations which includes radial flow.
- Free parameter: q , which is related to the finite size of the system. Bjorken flow is recovered by taking $q \rightarrow 0$.
- Radial velocity of the fluid is

$$v_r = \frac{2q^2 r \tau}{1 + (q\tau)^2 + (qr)^2}$$

- Requires azimuthal symmetry; applicable to central collisions only.
- Requires *conformal* symmetry; equation of state is fixed at $\varepsilon = 3P$.

Perturbing The Flow

- Introduce small perturbations of energy density and fluid velocity

$$\varepsilon(\rho, \theta, \phi, \xi) = \frac{\hat{\varepsilon}_0(\rho)}{\tau^4} [1 + \lambda \delta(\rho, \theta, \phi, \xi)] + \mathcal{O}(\lambda^2),$$

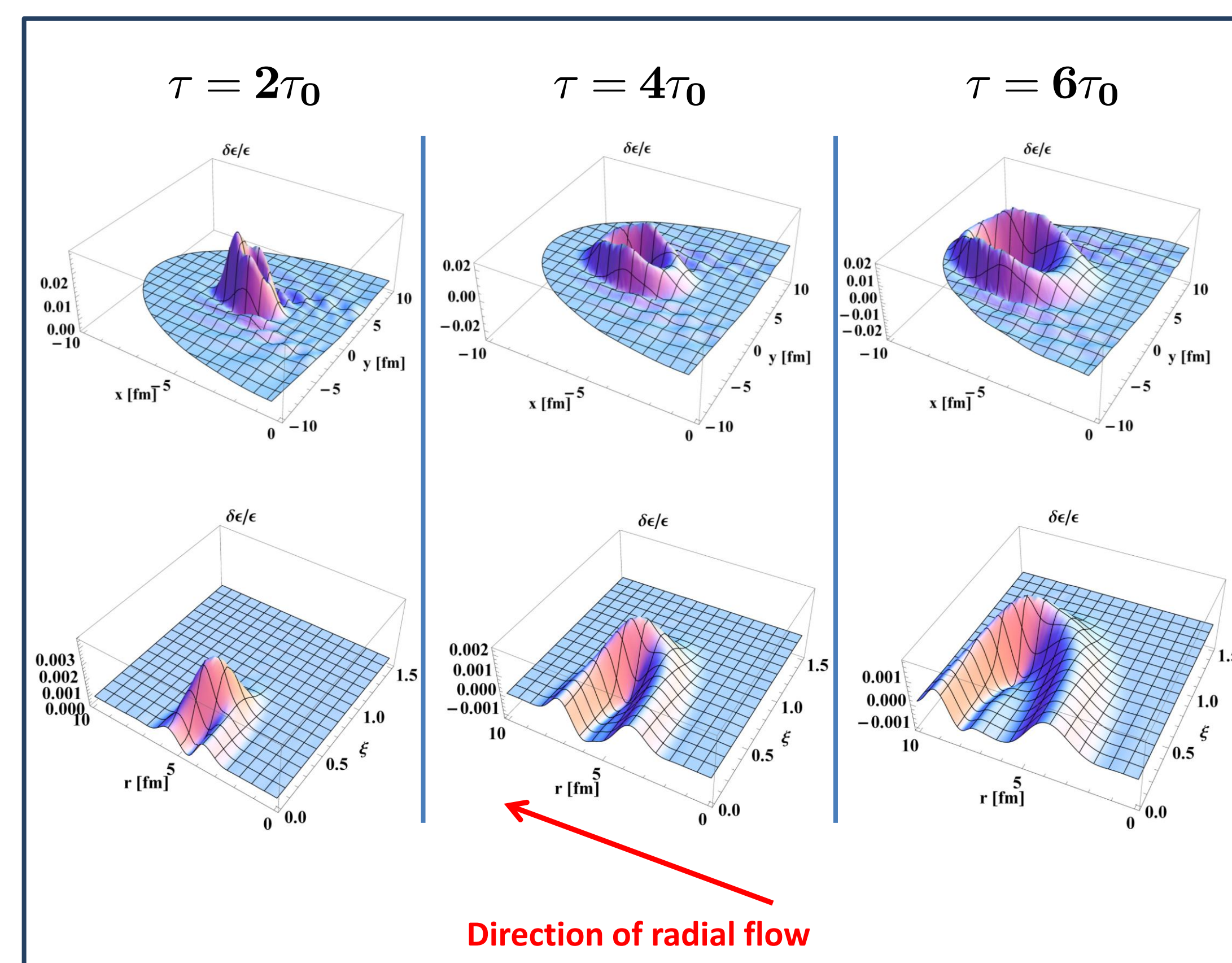
$$w^\mu(\rho, \theta, \phi, \xi) = \frac{1}{\tau} (1, \lambda u^i(\rho, \theta, \phi, \xi)) + \mathcal{O}(\lambda^2).$$

- Linearize the resulting hydrodynamic equations to first order in the small parameter λ .
- Fourier transform in ξ , and employ a spherical harmonic expansion in the angular (θ, ϕ) directions.
- Initial condition is a Gaussian hotspot in energy density with vanishing fluid velocity:

$$\delta(\rho_0, \theta, \phi, \xi) = A \exp \left\{ -\frac{\theta^2 + \theta_0^2 - 2\theta\theta_0 \cos(\phi - \phi_0) + \xi^2}{2\Sigma^2} \right\}.$$

- Solve the linearized equations numerically to evolve the perturbation in ρ .
- Transform the coordinates from ρ, θ to τ, r .

Time Evolution



Correlations

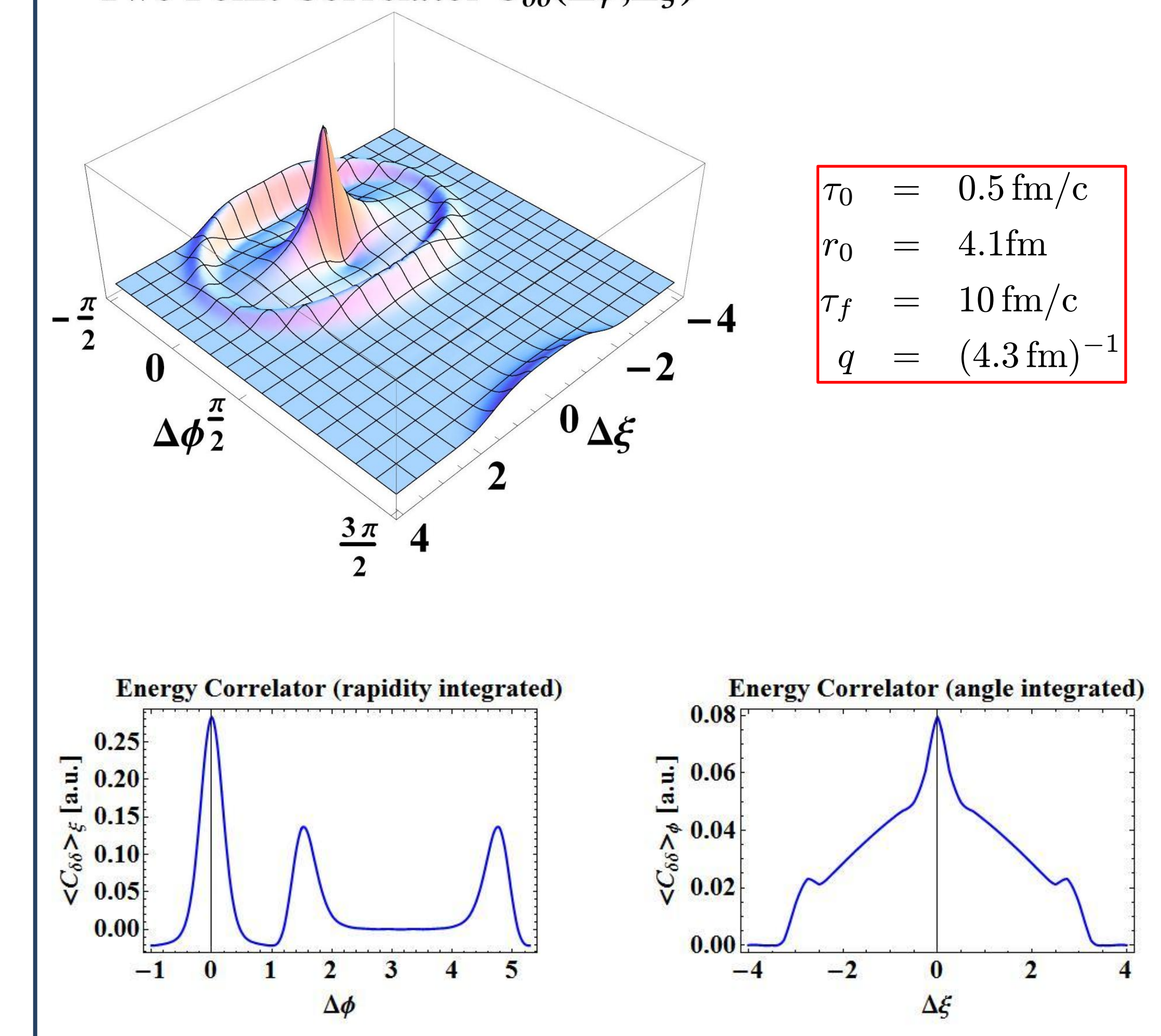
- Integrate over the transverse plane,

$$\bar{\delta}(\tau, \phi, \xi) \equiv \frac{1}{r_{\max}} \int_0^{r_{\max}} r dr \delta(\tau, r, \phi, \xi).$$

- Compute correlation function by averaging over the initial azimuthal angle and rapidity,

$$C_{\delta\delta}(\tau, \Delta\phi, \Delta\xi) \equiv \int d\chi d\eta \bar{\delta}(\tau, \phi_1 - \chi, \xi_1 - \eta) \bar{\delta}(\tau, \phi_2 - \chi, \xi_2 - \eta).$$

Two Point Correlator $C_{\delta\delta}(\Delta\phi, \Delta\xi)$



Discussion

- Perturbations lead to propagating sound waves. Sound front peaks are visible in both azimuth and rapidity.
- Central peak is broadened in rapidity direction due to expansion.
- In progress: replacing the hotspots presented here with stochastic hydrodynamic fluctuations. What impact do these fluctuations have on the two-particle correlation data, or equivalently the flow coefficients v_n ?

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