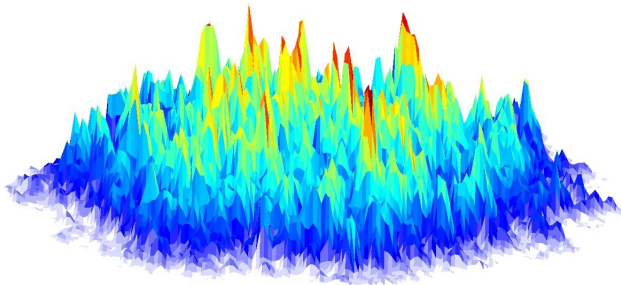




Imprinting Quantum Fluctuations on Hydrodynamic Initial Conditions

J. Scott Moreland, Z. Qiu; Advisors S. Bass & U. Heinz

- *Adding quark and gluon degrees of freedom to Monte Carlo
Color-Glass Condensate initial conditions*



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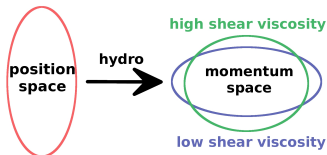
*Special thanks to U. Heinz and S. Bass for their guidance and to the Krell Institute for funding my future graduate studies.

Brief background: extracting the QGP shear viscosity

- Hybrid models using **viscous hydrodynamics** + **Boltzmann cascade** provide useful tool for extracting QGP shear viscosity η/s .

$$\eta/s \leftrightarrow \frac{V_n}{\epsilon_n} \leftarrow \begin{array}{l} \text{anisotropic flow} \\ \text{spatial anisotropy} \end{array}$$

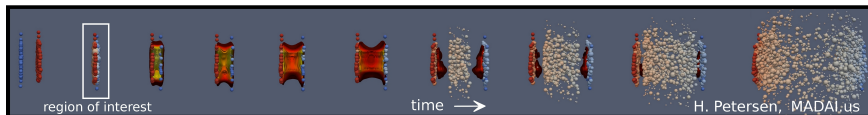
viscosity η/s dampens conversion of spatial anisotropy into flow anisotropy



- Initial conditions (IC) simulate transverse energy (entropy) density
- Hydrodynamic transport equations evolve IC down to T_c (critical temperature)
- Boltzmann cascade evolves hadron resonance gas from T_c to freeze-out
- Particle interactions cease and particles free-stream

(1) Initial Condition

(3) Hadron Resonance Gas (Boltzmann)



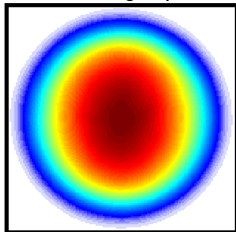
(2) QGP (Hydrodynamics)

(4) Kinetic Freezeout

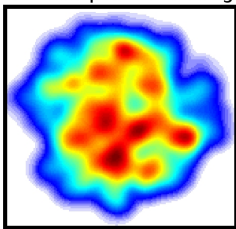
Concentrating on initial conditions

- ▶ At present, largest source of error for extracting η/s is uncertainty in initial conditions: MC-Glauber? DIPSY? AMPT? MC-KLN? MC-rcBK? NeXuS? UrQMD? IP-Sat? IP-Glasma?... It's a zoo.
- ▶ Current models not yet converging, only discovering new sources of uncertainty
- ▶ Let's concentrate on developments in fluctuations
Classic Example: Optical nucleus \rightarrow Monte Carlo nucleus w/finite nucleons

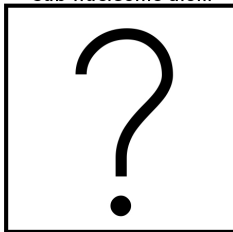
event averaged profile



nucleon pos. fluctuating



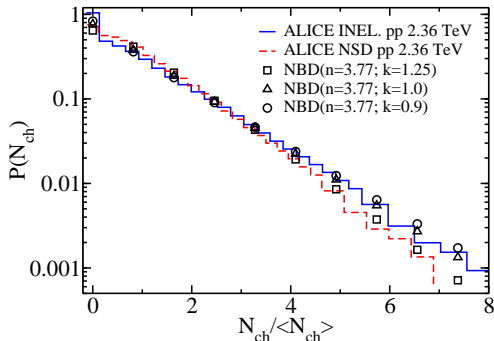
sub-nucleonic d.o.f.



What's next? Event-by-event gluon-field fluctuations

- ▶ Large event-by-event fluctuations in pp charged particle multiplicity N_{ch} evidence for sub-nucleonic initial state fluctuations

K. Aamodt et al. (ALICE Collaboration) Eur.Phys.J. C68 (2010) 89-108

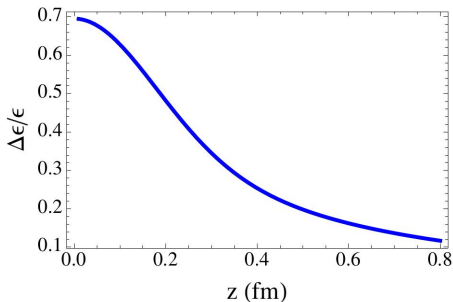


- ▶ pp multiplicity fluctuations **cannot be explained** using Glauber model for pp interaction where, $P(b) = 1 - \exp[-\sigma_{gg} T_{pp}(b)]$
- ▶ Sub-nucleonic fluctuations clearly exist in the initial state, but where to begin?

Müller & Schäfer calculate transverse energy density fluctuations in Color-Glass Condensate model

- ▶ Müller-Schäfer paper predicts structure of quantum fluctuations!
- ▶ Paper calculates mean normalized covar.
$$\text{Cov}[\epsilon(r)/\epsilon_0] = (\Delta\epsilon(r)/\epsilon_0)^2$$
 for collision of two infinite sheets of nuclear matter
- ▶ Gaussian color distribution (color content contribution from each nucleon is independent).
- ▶ Correlation length *proportional* to $1/Q_s \approx 0.14$ fm (relevant Q_s at RHIC)
- ▶ Use Müller-Schäfer covariance of gluon-field energy density fluctuations to texture MC-KLN initial conditions

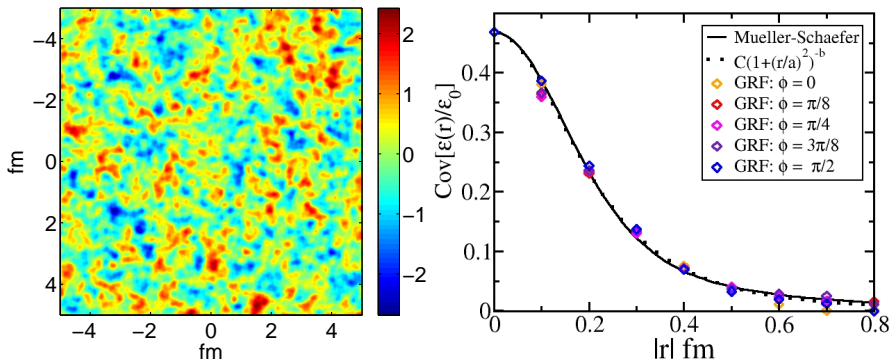
B. Müller, A. Schäfer, Phys.Rev.D 85, 114030 (2012)



Fitting the Müller-Schäfer Gaussian random field (GRF)

We use GRF **T**urning **B**and **S**IMulator TBSIM developed by X. Emery and C. Lantuéjoul to generate a large 4000x4000 cell GRF w/ the Müller-Schäfer covariance

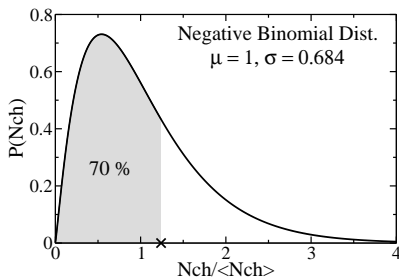
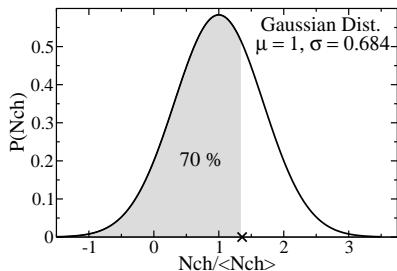
X. Emery and C. Lantuéjoul, Computers and Geosciences 32, 1615 (2006)



Fit to Müller-Schäfer covariance is nearly exact. One small problem: GRF allows for fluctuations $\Delta\epsilon/\epsilon_0 < -1$, negative energy densities which are, of course, non-physical.

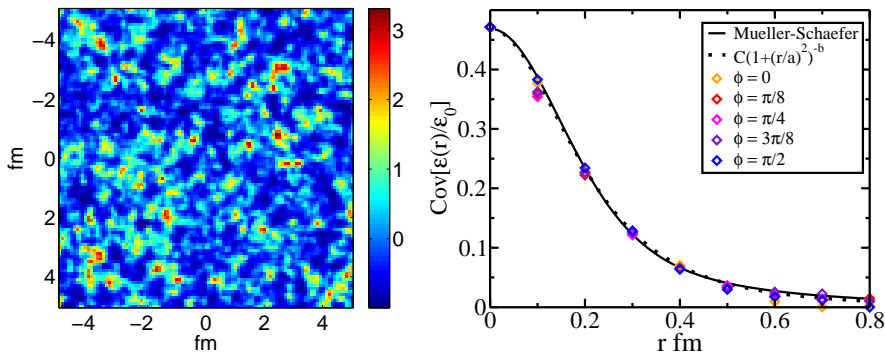
Mapping to positive definite Neg. Binom. random field

- ▶ 7.2% chance Gaussian fluctuations are negative (non-physical)
- ▶ Send Gaussian \rightarrow Negative Binomial Distribution (NBD) motivated by p-p multiplicity fluctuations
- ▶ $\text{NBD}(\bar{n}, k; n) = \frac{\Gamma(k+n)}{\Gamma(k)\Gamma(n+1)} \frac{\bar{n}^n k^k}{(\bar{n}+k)^{n+k}}$ 2-parameter function on \mathbb{Z}
- ▶ Generate single NBD with mean and standard deviation equal to Gaussian
- ▶ Equate random variables via their cumulative distribution functions (CDF's)
i.e. iff $\int_{-\infty}^x \text{Gaussian}(x) dx = \int_{-\infty}^y \text{NBD}(y) dy$ then $x \leftrightarrow y$



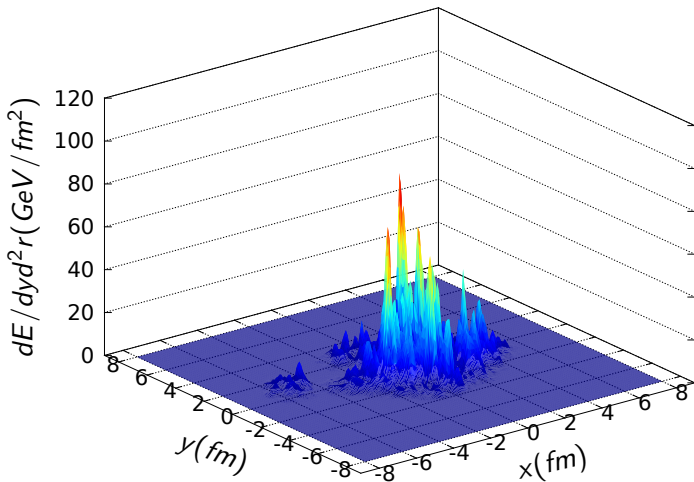
Did this transformation destroy our covariance?

- ▶ Nope. Almost indistinguishable from GRF fit to Müller-Schäfer covariance

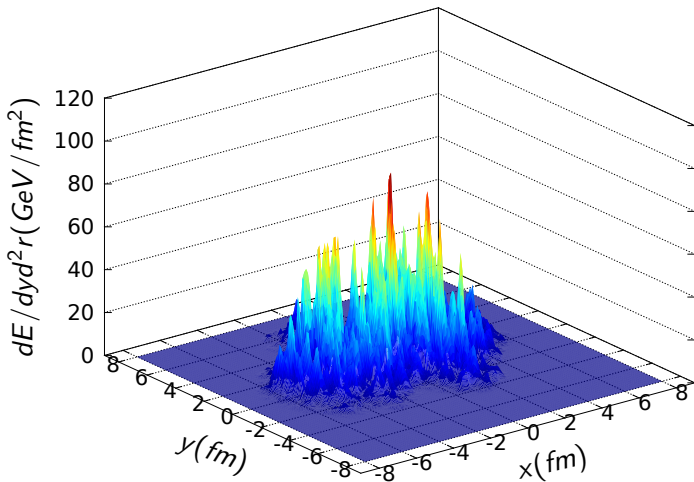


- ▶ Bulk of profile is pulled down and hot spots are more sharply peaked
- ▶ Texture can be thought of as percent fluctuation about mean gluon-field energy density

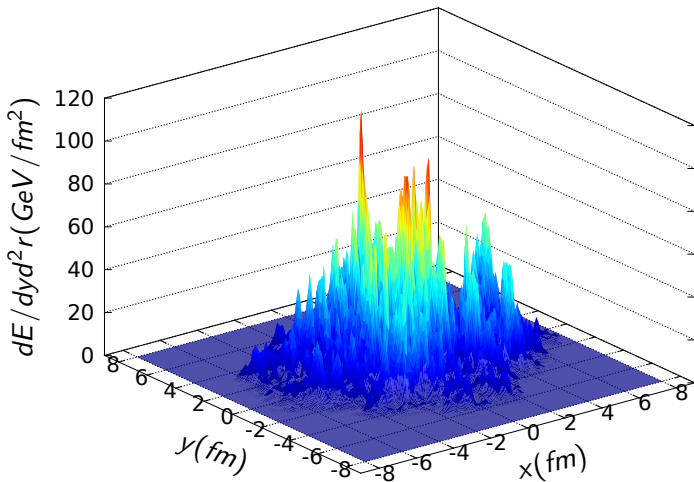
Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=100



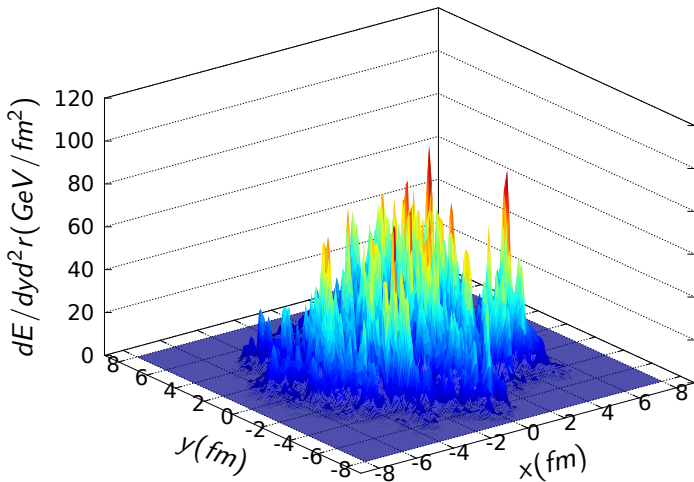
Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=200



Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=300



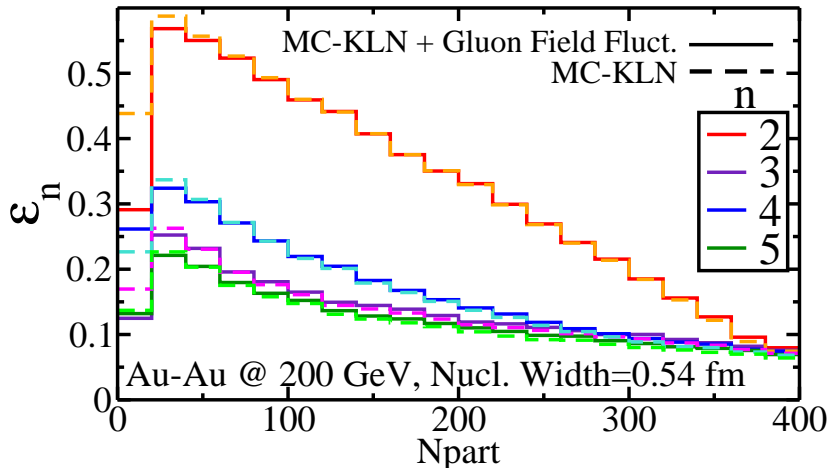
Fluctuated MC-KLN, Au-Au, 200 GeV; Npart=380



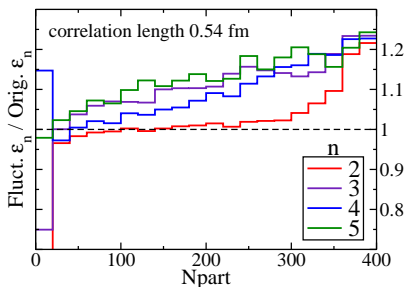
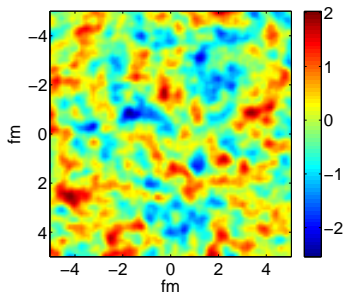
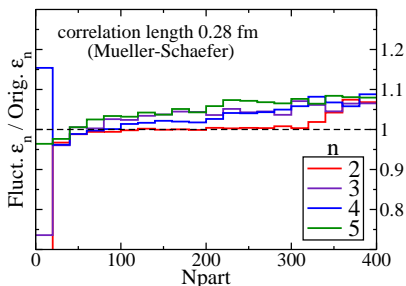
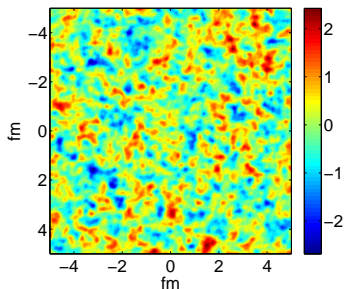
Analysis: Calculating the eccentricity harmonics

$$\text{Using definition: } \epsilon_n e^{in\Psi_n^{PP}} = -\frac{\int dx dy r^2 e^{in\phi} \rho(x,y)}{\int dx dy r^2 \rho(x,y)}$$

Quantum Fluct. w/Mueller-Schaefer Covar.



Analysis: Dependence on correlation length



Comparison with similar studies

Related Publications:

- ▶ A. Dumitru & Y. Nara
“KNO scaling of fluctuations in pp and pA, and eccentricities in heavy-ion collisions” Phys. Rev. C 85, 034907 (2012)
[arXiv:1201.6382 [nucl-th]].
- ▶ B. Schenke, P. Tribedy & R. Venugopalan
“Fluctuating Glasma initial conditions and flow in heavy ion collisions” arXiv:1202.6646 [nucl-th].

Strengths of the present analysis:

- ▶ Realistic covariance applicable in central (hot) regions of the fireball
- ▶ Fluctuations can be turned on and off for comparison (factorized) and applied to any suitable Color-Glass Condensate model

Concluding remarks

- ▶ We have implemented gluon-field fluctuations in MC-KLN initial conditions with a uniform Müller-Schäfer covariance.
- ▶ We see a small increase in $\epsilon_2 - \epsilon_5$ in central collisions $\approx 5 - 10\%$.
- ▶ Increasing the correlation length to the width of a nucleon (as may happen for smaller Q_s values) results in $\approx 20 - 25\%$ increase of ϵ_n in central collisions.
- ▶ Increases in ϵ_2 vanish for $N_{part} < 300$.

Looking forward

- ▶ Refine model to better account for fluctuations in dilute regions, i.e. fitting pp multiplicity fluctuations
- ▶ Converge with parallel studies (B. Schenke, P. Tribedy, R. Venugopalan & A. Dumitru, Y. Nara)
- ▶ Run fluctuated profiles through hydro

Special thanks to Berndt Müller, Adrian Dumitru, Hannah Petersen, Chun Shen and Jonah Bernhard for constructive advice and illuminating discussions.