

# Initial conditions for fluid dynamics from a parton cascade model

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## 1. Motivation

We investigate the pre-equilibrium evolution of the QGP in ultrarelativistic heavy-ion collisions. We first parametrize the initial particle distribution function at  $t = 0$  fm/c. Then the subsequent time-evolution of the distribution function, up to  $t = 1$  fm/c, is modeled by solving the relativistic Boltzmann equation.

## 2. Parametrization

The initial single-particle distribution function reads as follows:

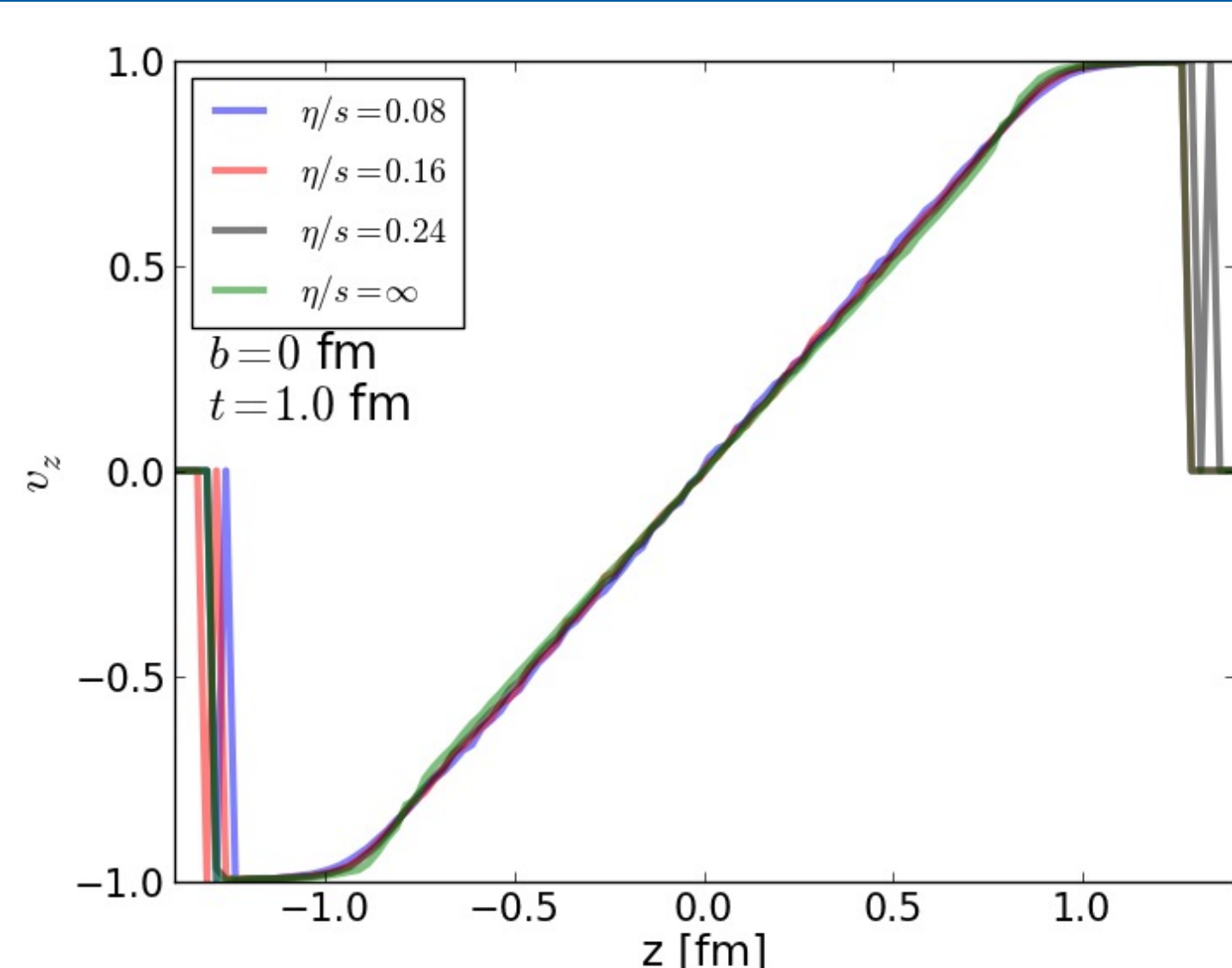
$$f(\mathbf{x}, \mathbf{p}) = K \frac{1}{E} \left( \frac{Q^n}{Q^n + p_T^n} \right)^n e^{-\frac{y^2}{\sigma_y^2}} \times e^{-\frac{z^2}{\sigma_z^2}} T_A(\mathbf{x}_T - \frac{\mathbf{b}}{2}) T_B(\mathbf{x}_T + \frac{\mathbf{b}}{2})$$

- $T_A$  and  $T_B$  are the nuclear thickness functions
- Parameters are:  $Q = 0.6$  GeV,  $n = 1.4$ ,  $m = 5.5$ ,  $\sigma_y = 2.25$  and  $\sigma_z = 0.065$  fm
- The normalization: we use  $K = 0.6$ ; for  $\eta/s = 0.08$  and  $\mathbf{b} = 0$ ,  $e_{\max} = 30$  GeV/fm<sup>3</sup> at  $0.6$  fm/c and  $z = 0$  fm

## 3. The Transport model BAMPS

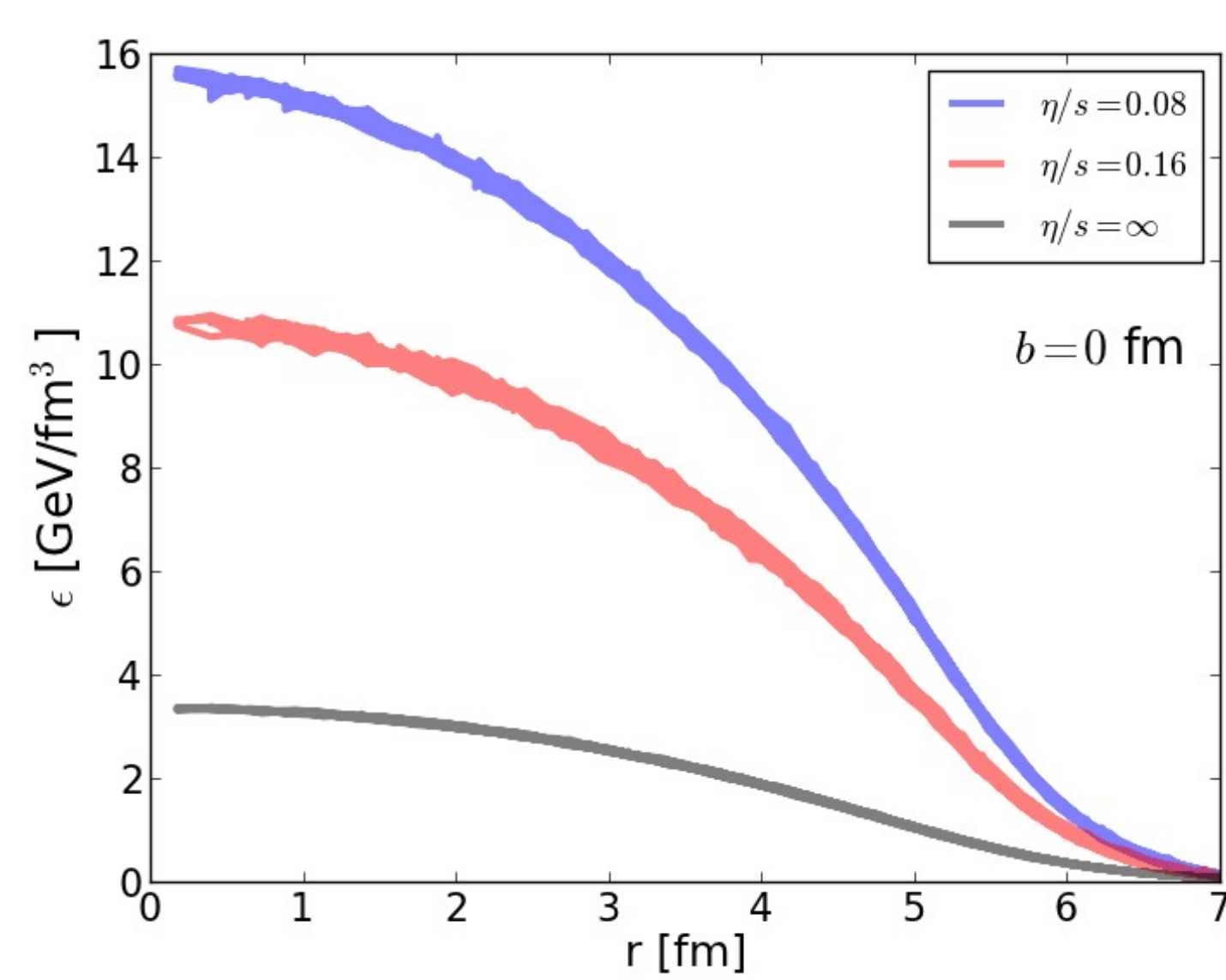
- The time evolution is solved by the Boltzmann approach for multi parton scatterings (BAMPS) [1] up to times of the order of 1 fm/c
- Cross-section adjusted to give a constant  $\eta/s$ ; we use two different impact parameters  $\mathbf{b}$
- we consider only massless particles, binary collisions and an isotropic cross section

## 4. Scaling Flow

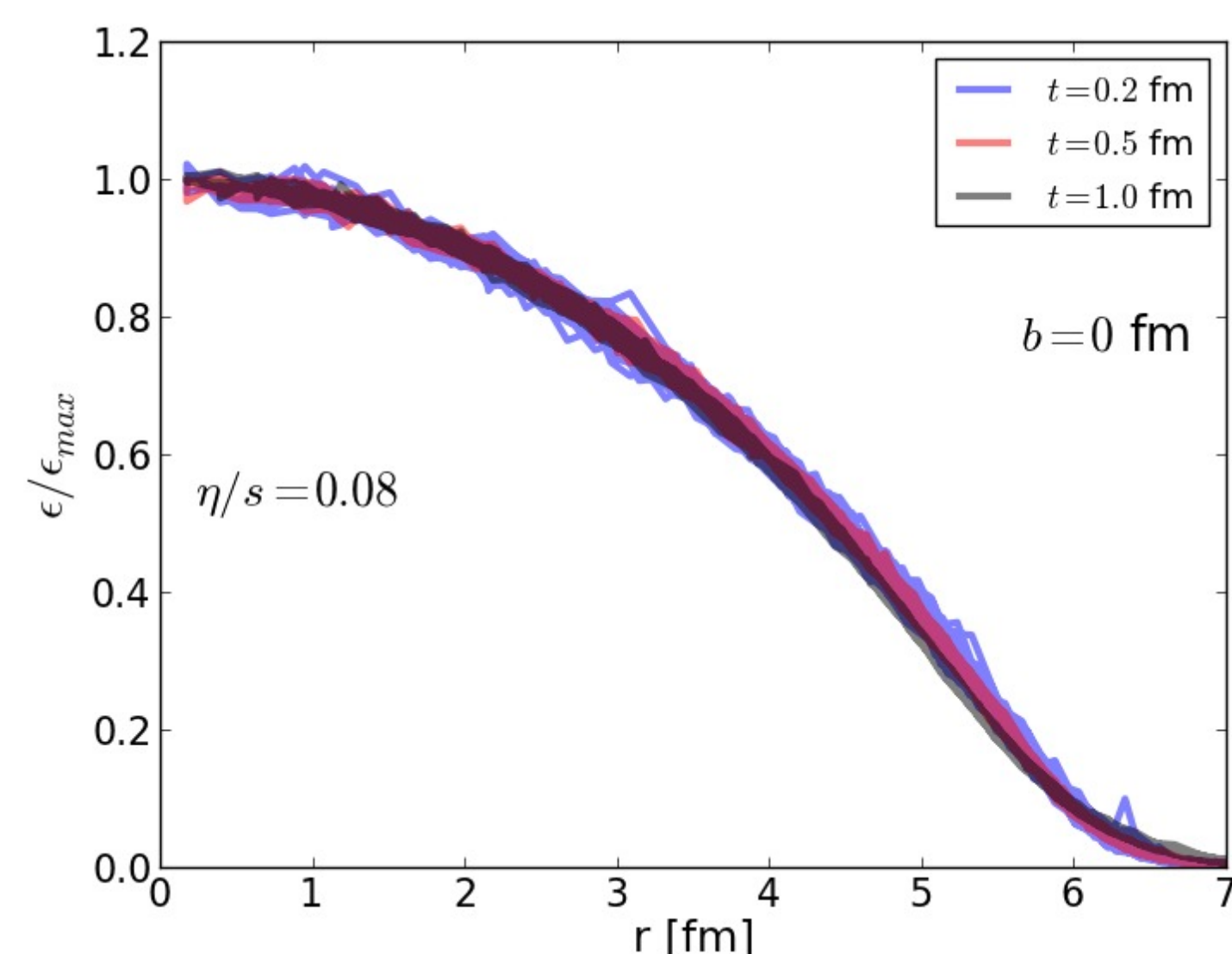
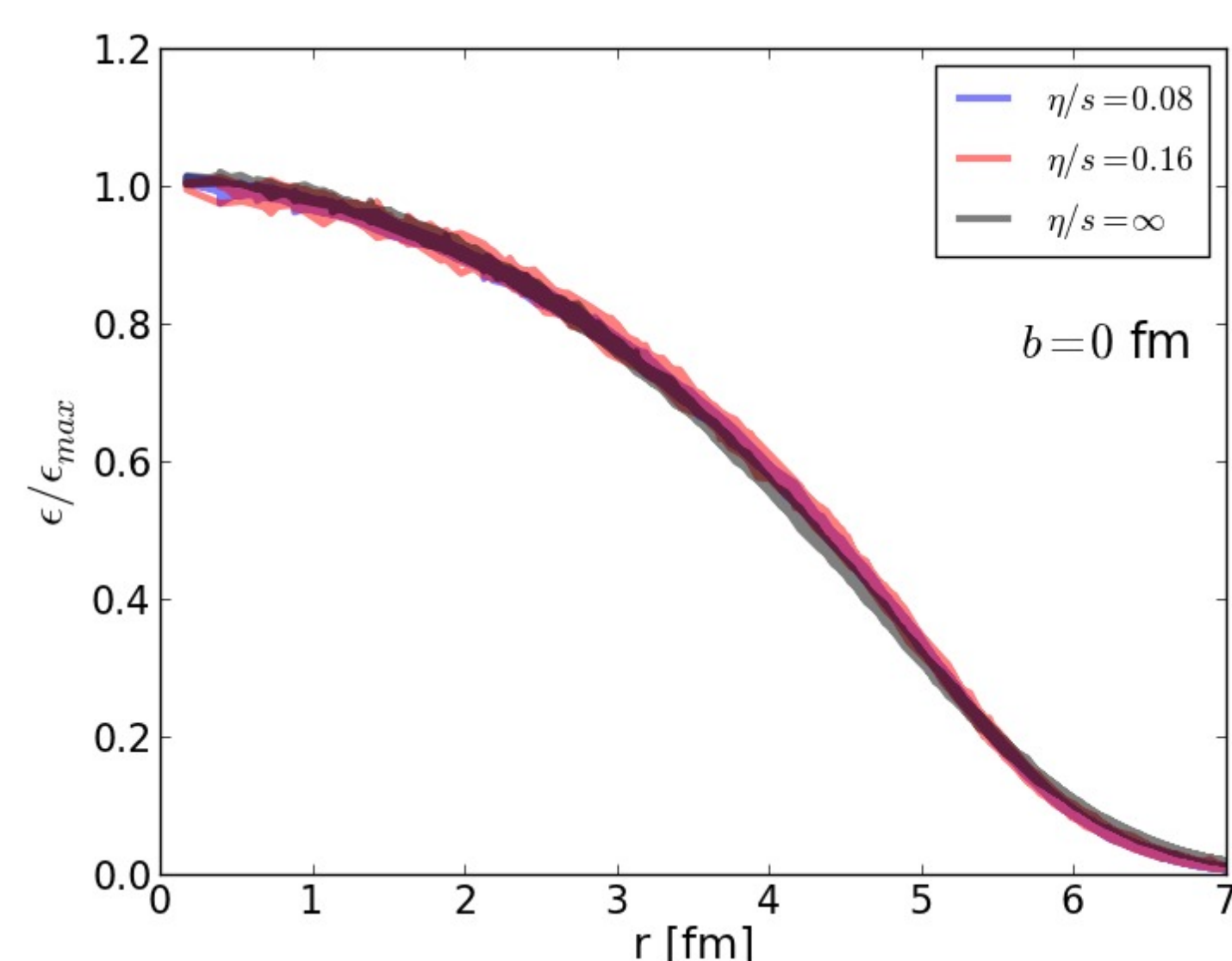


- Longitudinal velocity  $v_z$  at  $r = 0$  fm and  $t = 1.0$  fm/c and  $\mathbf{b} = 0$
- Scaling flow ( $v_z(z) = z/t$ ) reached for all choices of  $\eta/s$

## 5. Energy Density

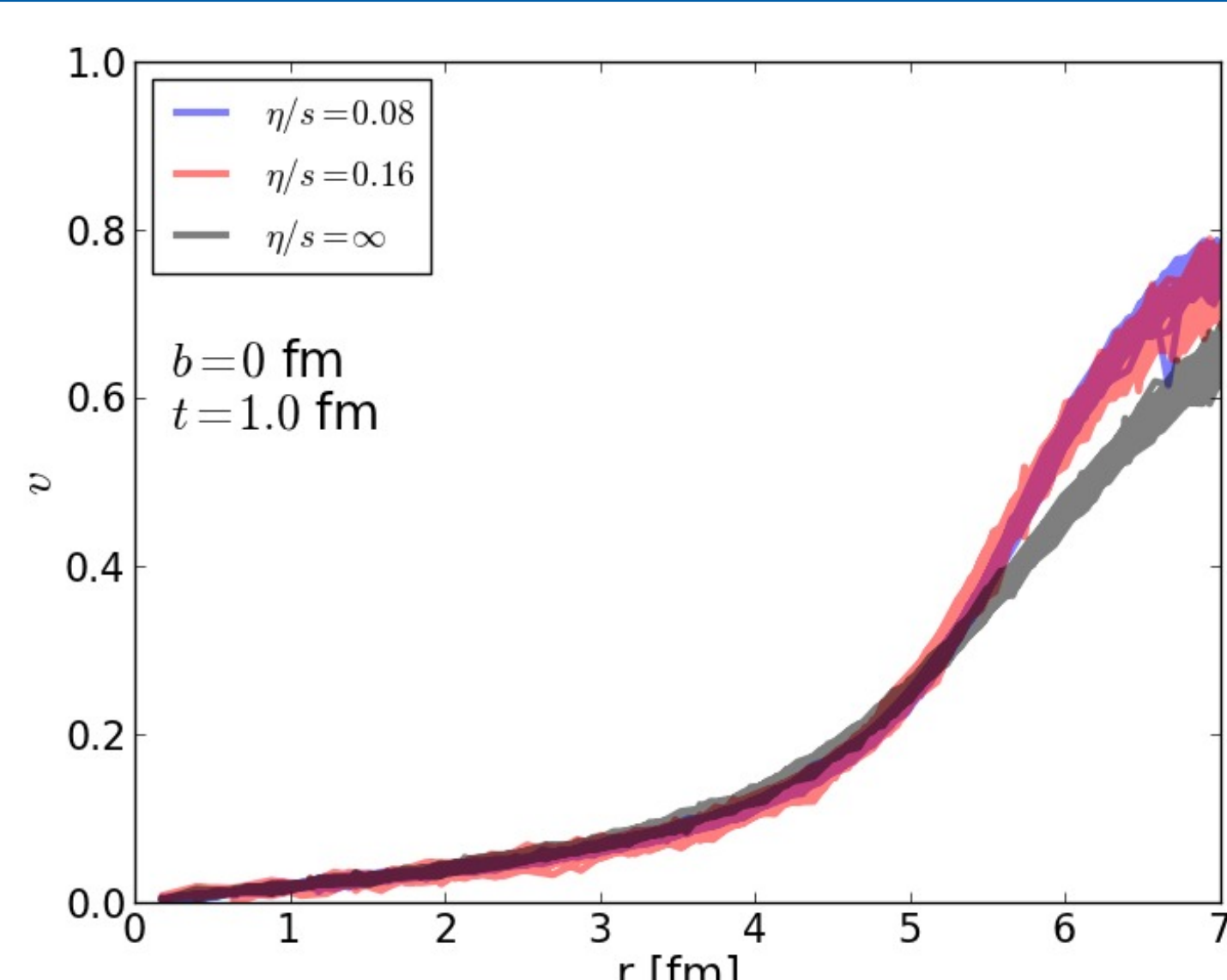


- energy density  $e$  at  $\mathbf{b} = 0$  and  $z = 0$  fm
- Strong dependence of  $e_{\max}$  on the value of  $\eta/s$
- The larger  $\eta/s$ , the faster energy density drops



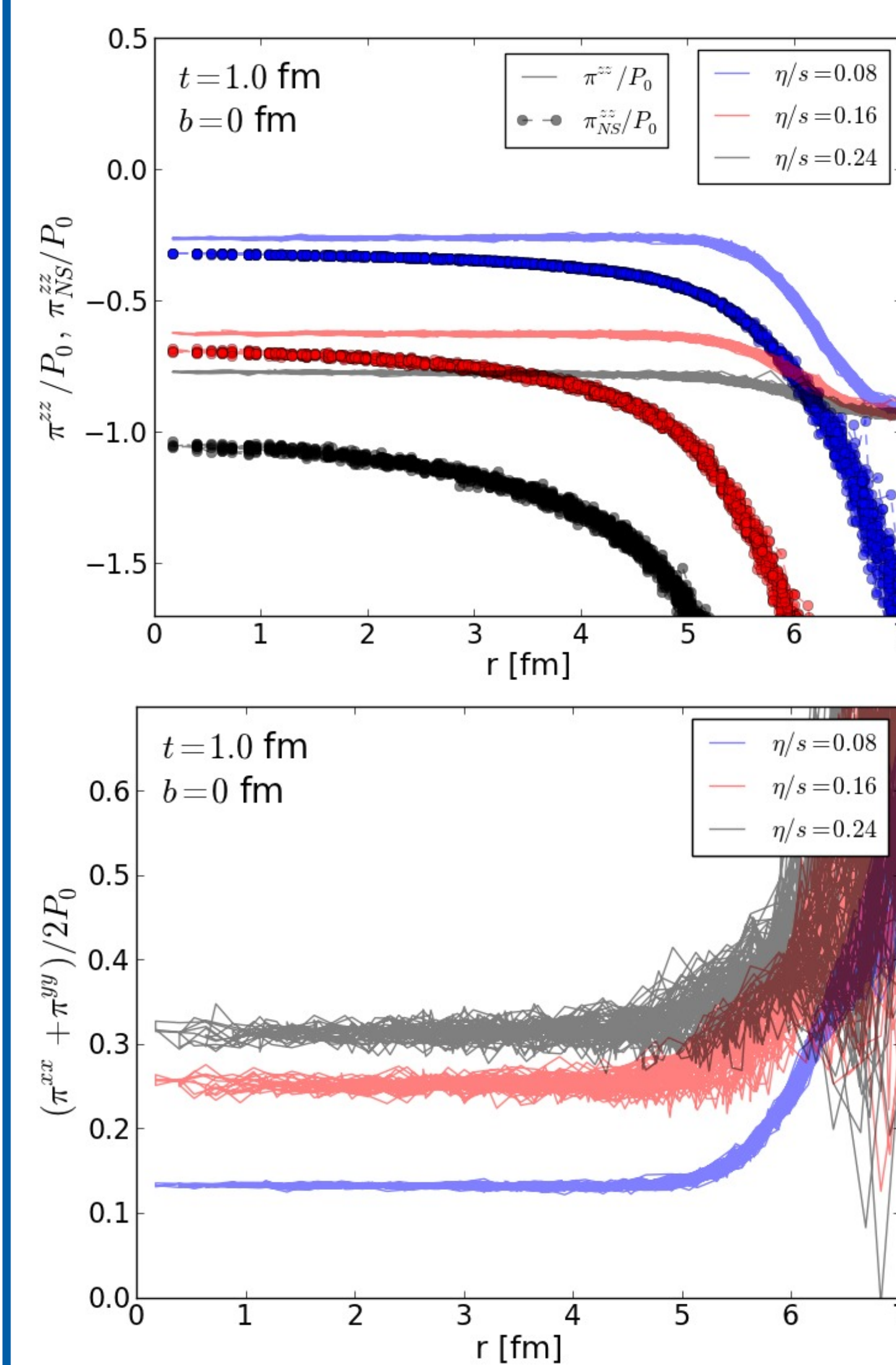
- Shape ( $e/e_{\max}$ ) of the energy density profile at  $t = 1.0$  fm/c practically independent of  $\eta/s$
- Shape of the profile practically unchanged between  $t = 0$  fm/c and  $t = 1$  fm/c

## 6. Transverse velocity

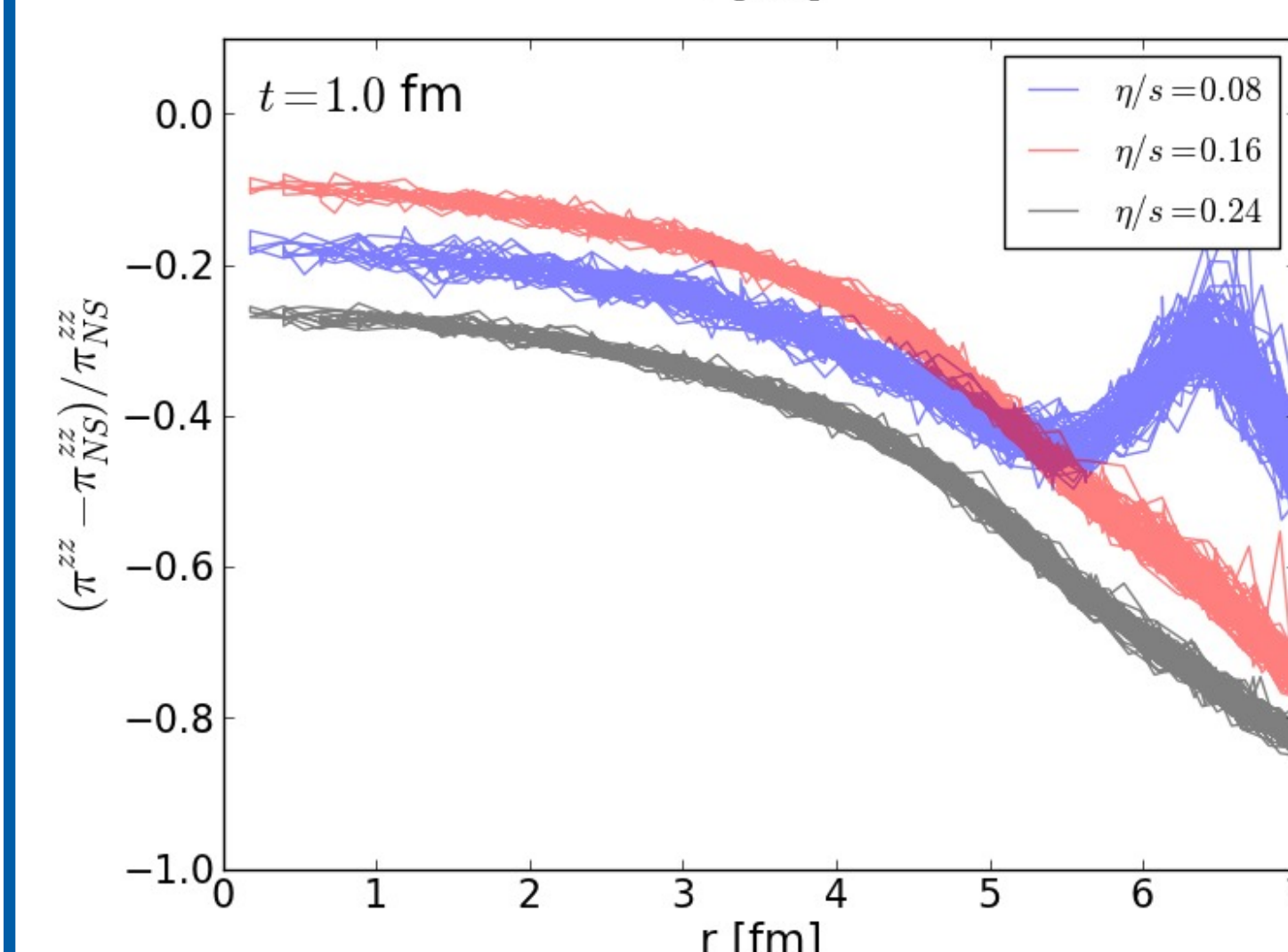
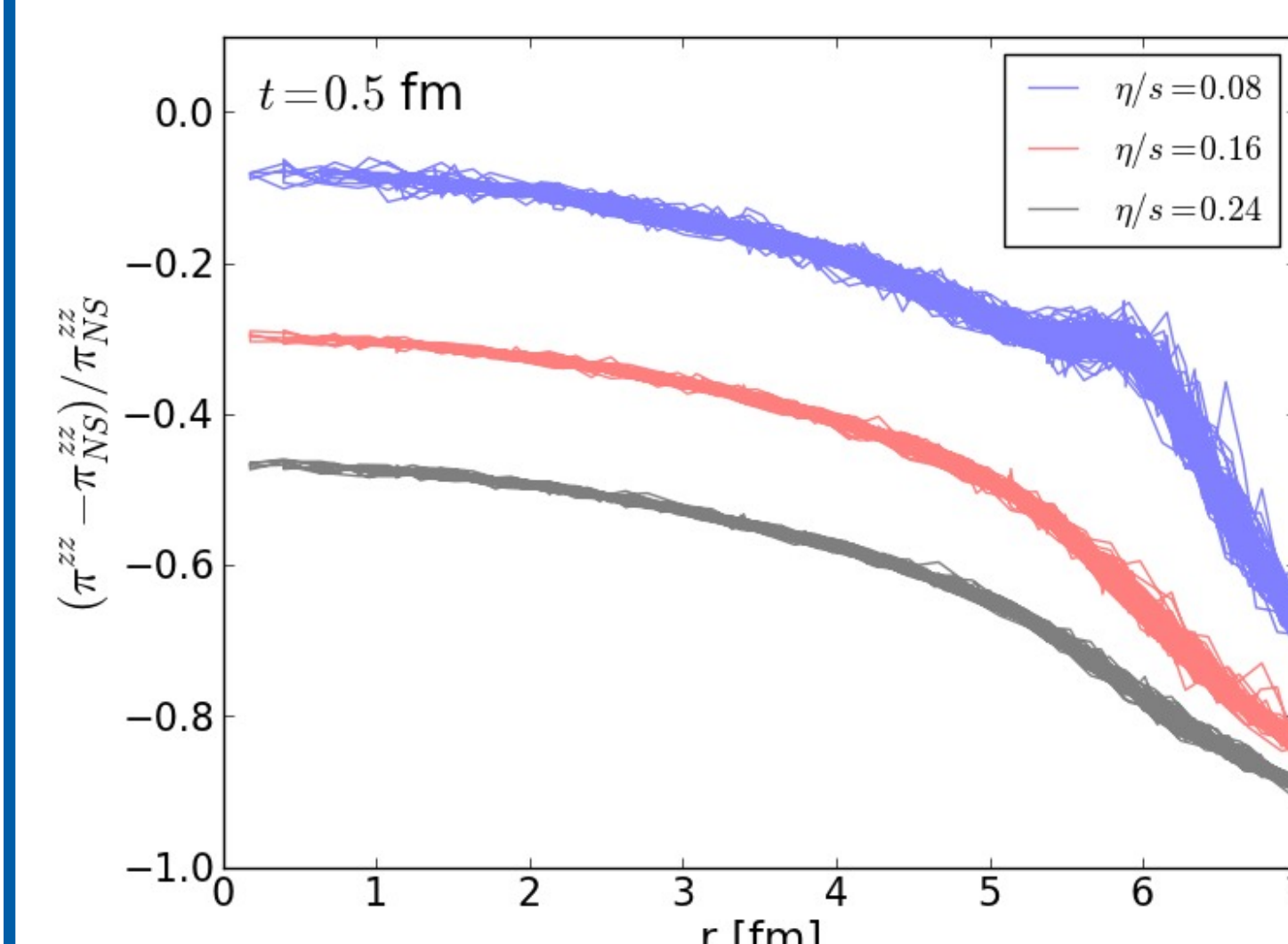


- Transverse velocity at  $z = 0$  fm,  $t = 1.0$  fm/c and  $\mathbf{b} = 0$
- Significant transverse velocity established (especially near the edges)
- Transverse velocity profiles almost independent of  $\eta/s$

## 7. Shear-Stress Tensor

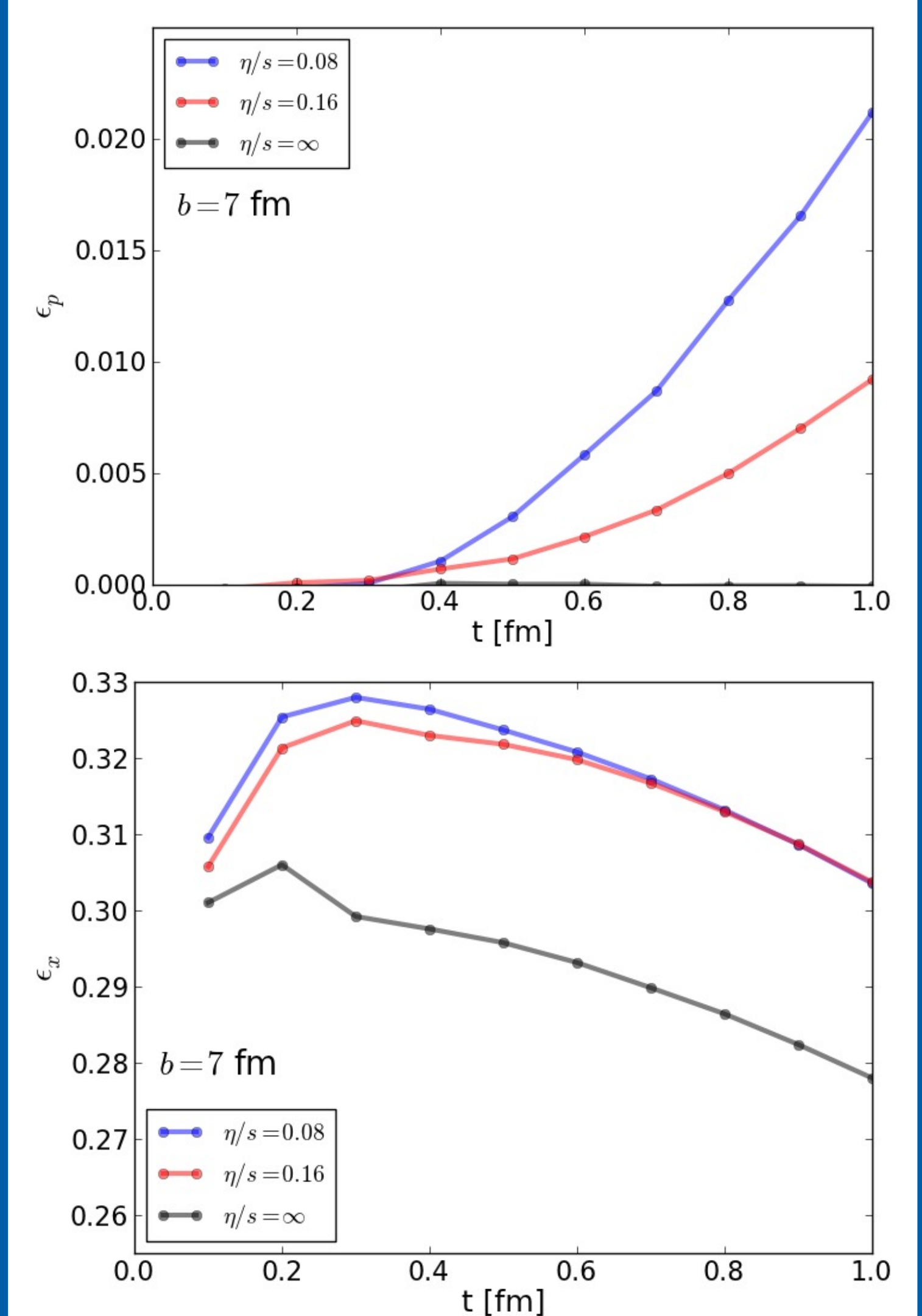


- Longitudinal and transverse components of the shear-stress tensor (over thermodynamic pressure) at  $\mathbf{b} = 0$  and  $z = 0$  fm at  $t = 1$  fm/c
- Strong deviations from the Navier-Stokes values (especially near the edges)
- NS-values overestimate the shear-stress tensor.



- Relative deviation from the NS-values at two different times

## 8. Eccentricities at early times



- time-evolution of momentum space  $\epsilon_p$  and spatial  $\epsilon_x$  eccentricities for non-central collisions  $\mathbf{b} = 7$  at  $t = 1$  fm/c
- $\epsilon_p = \frac{\int dx dy (T^{xx} - T^{yy})}{\int dx dy (T^{xx} + T^{yy})}$
- $\epsilon_x = \frac{\int dx dy e(\tau, x, y) (y^2 - x^2)}{\int dx dy e(\tau, x, y) (y^2 + x^2)}$
- Spatial eccentricity changes only slightly
- momentum eccentricity depends strongly on value of  $\eta/s$  (although transverse velocity almost independent of  $\eta/s$ )

## Conclusions

- We studied the pre-thermal evolution of the QGP using BAMPS
- Transverse velocity and the shape of the energy density profile after the pre-equilibrium phase are practically independent of  $\eta/s$
- The shear stress tensor developed in the first  $fm$  of evolution is non-zero and can deviate from the Navier-Stokes prediction
- For non-central collisions we found that a considerable momentum anisotropy can already be generated in the pre-thermal evolution of the QGP

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## References

- [1] Z. Xu and C. Greiner, Phys. Rev. C **71** (2005) 064901; Phys. Rev. C **76** (2007) 024911