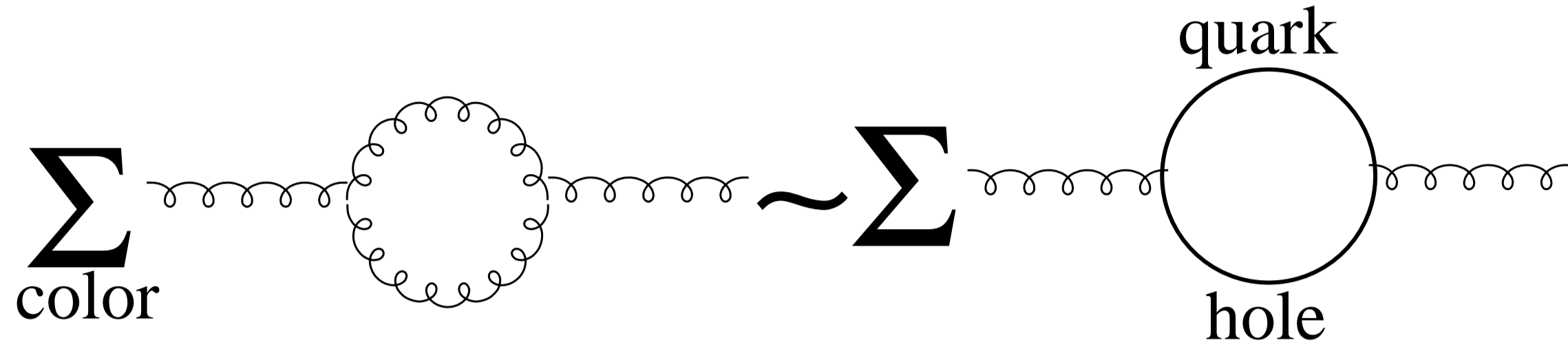


## Introduction: quarkyonic matter

Let us think about QCD at moderate temperature and chemical potential ( $T \leq T_c, \mu_Q \sim \Lambda_{QCD}$ ) at the large  $N_c \gg N_f \gg 1$  limit

**In momentum space**, infrared slavery should still be maintained between  $\mu_Q \sim \Lambda_{QCD}$  and  $\mu_Q \sim \sqrt{N_c/N_f} \Lambda_{QCD}$ . This is because quark-hole screening  $\sim \mu_Q^2 N_f N_c$  and gluon antiscreening  $\sim N_c^2$ .



higher order diagrams will not change this power

**In configuration space**, baryons should “touch”/“overlap” with each other, since  $N_c$  quarks in a baryon, quarks of neighboring baryons at  $\sim N_c^{-1/3} \Lambda_{QCD} \rightarrow 0$  apart

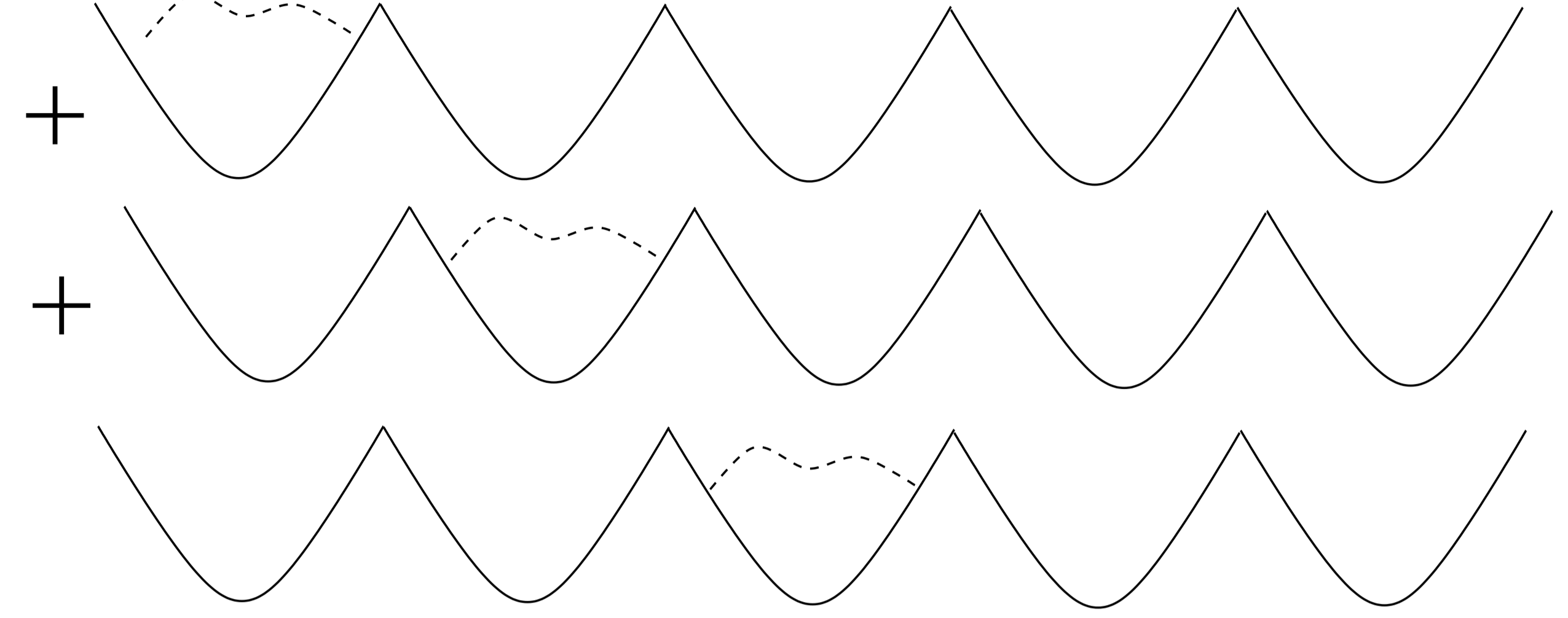
So matter seems to be both confining and asymptotically free. potentially very interesting.

[1] suggested the “quarkyonic” matter  $\Lambda_{QCD} \geq \mu_Q \leq \sqrt{N_c/N_f} \Lambda_{QCD}$  has quark degrees of freedom below the fermi surface and baryonic degrees of freedom at the Fermi surface. This would imply pressure and entropy  $\sim N_c$  even in confining matter.

**Interesting, but...** there is no controlled approximation in this regime: lattice has sign problem, EFTs diverge if  $T, \mu_Q \sim \Lambda_{QCD}$ , and  $N_c \rightarrow \infty$  for semiclassical Gauge/string.

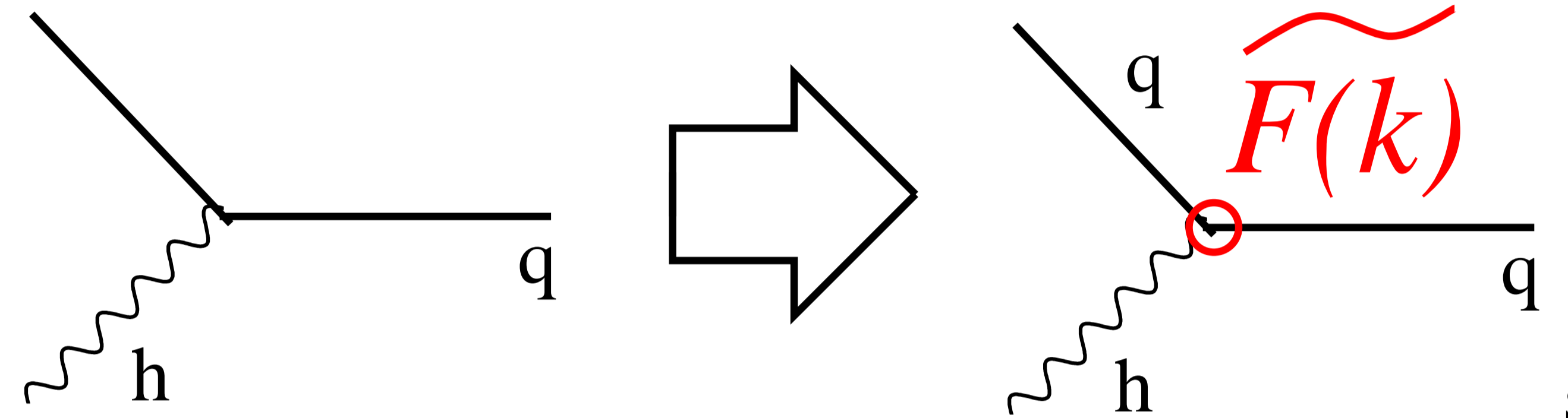
## Theory of quarkyonic percolation

We will try to outline what an EFT of quarkyonic matter looks like from the point of view of large  $N_c$  and percolation. A ready analogy is the electron gas model of the metal, where atoms are semi-classical potential wells (analogous to large  $N_c$  baryons) and electrons are quantum particles moving in them (analogous to quarks). The conductor-insulator phase transition (roughly similar to percolation) is defined by delocalization.



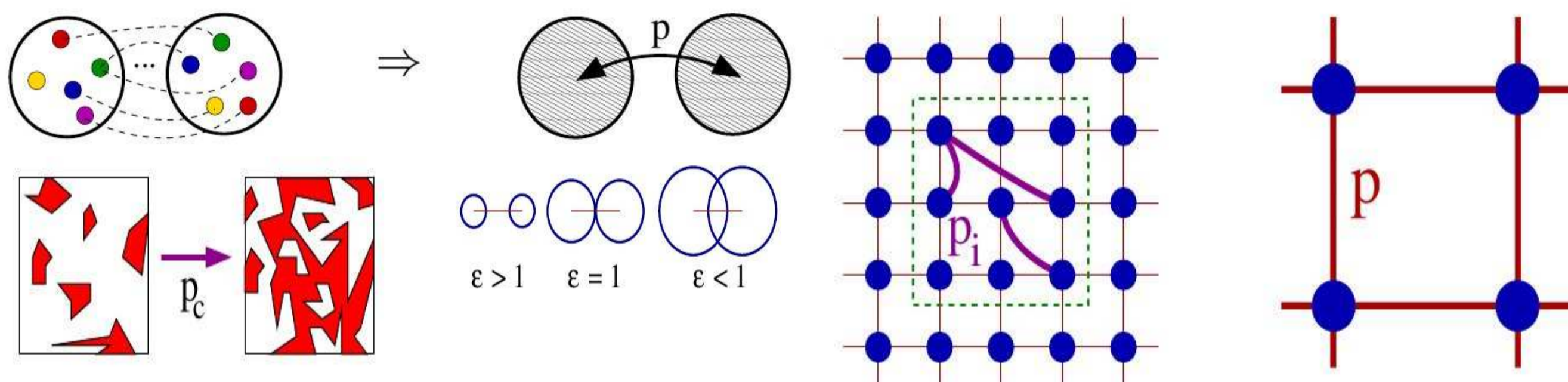
.... In the percolating phase, quark wavefunctions are delocalized w.r.t. baryons. Color is local on scales  $\geq \Lambda_{QCD}^{-1}$  but, because of delocalization, this does not mean quark wavefunctions at large distances are necessarily correlated (models like WZW can be modified to explicitly eliminate such long-range correlations).

Dynamics could look like pQCD, but baryon mean fields remain as a form factor



So, confinement is encoded in the baryonic form-factors, but baryons the partition function is still pQCD like, a la [1]. In particular,  $p, s \sim N_c$ . In baryon matter,  $e \sim N_c$  but  $p, s \sim N_c^0$ , so these two phases should be distinguishable.

## Some insight from percolation [3,4]



Look for a percolation type transition: integrate out quark degrees of freedom into baryon-baryon interactions

$$p = 1 - (q_{(i,j)})^{(N_c)^a} \quad q_{(i,j)} = \int f_A(x_i) dx_i \int f_B(x_j) dx_j (1 - F(|x_i - x_j|))$$

Then decimate a lattice into a super-lattice, and look for a fixed point

We assume [3] a “hard-sphere” distribution for  $f_{A,B}(x)$ , and a propagator respecting the scaling expected from QCD, a coupling constant  $\lambda/N_c$  and a range  $r_T \sim \Lambda_{QCD} \sim N_c^0$ . Propagator choices  $g(y) = (\lambda N_c^{-1}) F_{K,S,T}(y)$  include popular functions found in the literature [2]

$$F_T(y) = \Theta\left(1 - \frac{y}{r_T/\Lambda_{QCD}}\right), F_K(y) = \frac{2r_T^2}{\pi y^2} \sin^2\left(\frac{y}{r_T/\Lambda_{QCD}}\right), F_S(y) = F_K(y) e^{-M|y|}$$

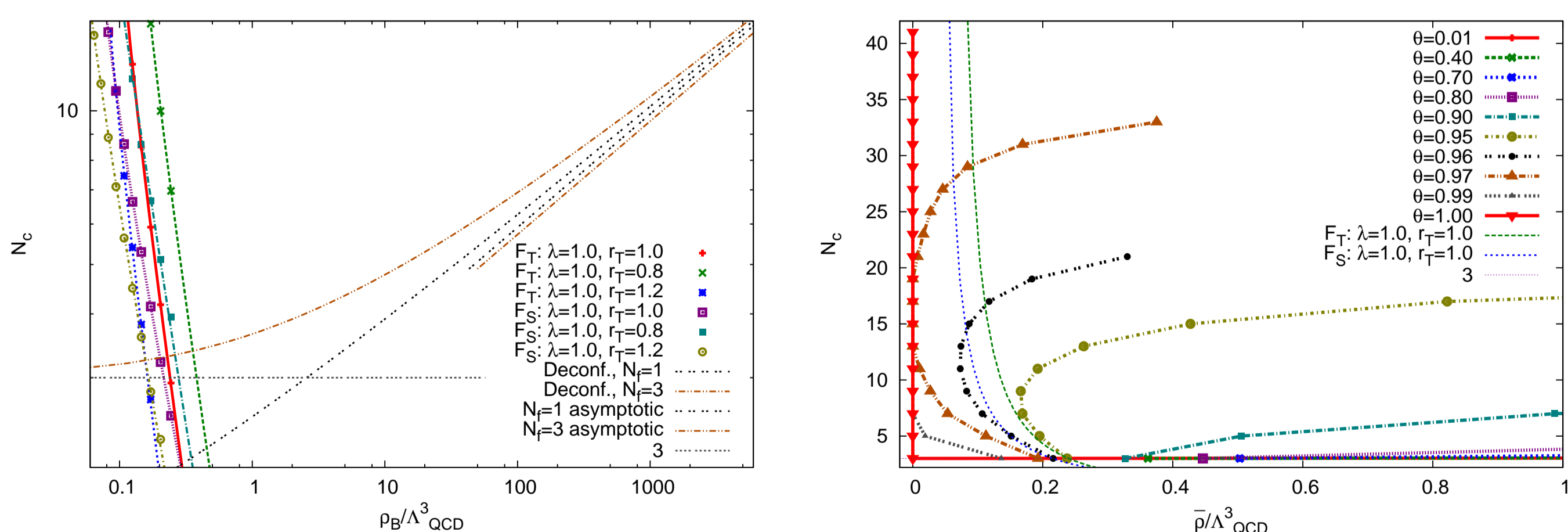
The first is a  $\Theta$ -function in configuration space, the second a  $\Theta$ -function in momentum space and the third a Debye-screened  $\Theta$ -function in momentum space.

**Not deconfinement** ( $\mu_q^{deconf} \sim N_c \Lambda_{QCD}$ ) . **but** percolation, deconfinement cross @  $N_c^{crit}$ . What is  $N_c^{crit}$ ? Relationship to ( $N_c^{crit}$ ) confinement?

## Percolation and deconfinement

We model  $T \ll \Lambda_{QCD}, \mu_Q \sim \Lambda_{QCD}$  matter by an ideal gas EoS (interactions, while not weak, impact pressure rather than density). For  $\mu_Q \sim N_c \Lambda_{QCD}$ , the baryonic density is

$$\rho_B^{conf}(T=0) = \Lambda_{QCD}^3 \frac{1}{6\pi^2} g_f g_s \frac{N_c^3}{N_f^{3/2}} (N_c - N_f)^{3/2} \sim \frac{g_f g_s}{6\pi^2} \Lambda_{QCD}^3 \frac{N_c^{9/2}}{N_f^{3/2}}$$



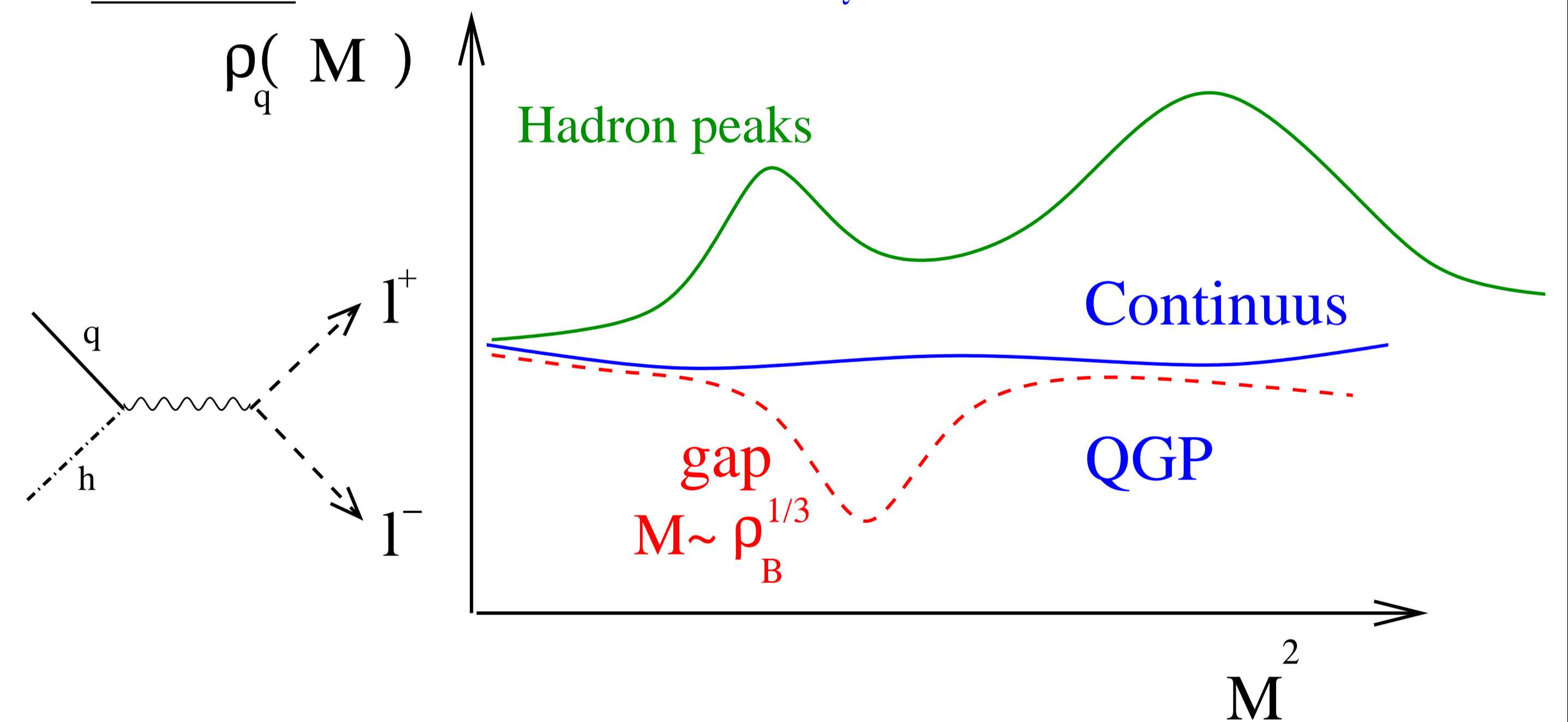
$T \sim \mu_Q \sim \Lambda_{QCD}$  confinement  $\sim N_c^0$ . Account for curvature in  $T - \mu_Q$  space by parametrizing phase diagram  $1 - \theta^2 = \left(\frac{\mu_B^{conf}}{\mu_0}\right)^2$ ,  $\theta = \frac{T}{T_c} \simeq \frac{3}{2} \frac{T}{\Lambda_{QCD}}$  finite density  $\rho_B$  is

$$\rho_B^{conf} = \sum_{\eta=0,1,\dots}^Q (2\eta + 2) \left\{ \int \frac{\alpha^2 d\alpha}{1 + \exp\left[\frac{3}{2} \frac{N_c - 1}{\sqrt{N_f} \theta} \left(\sqrt{\alpha^2 + N_f} + \eta \frac{\sqrt{N_f}}{N_c^2} - \sqrt{N_c} \sqrt{1 - \theta^2}\right)\right]} - \int \frac{\alpha^2 d\alpha}{1 + \exp\left[\frac{3}{2} \frac{N_c - 1}{\sqrt{N_f} \theta} \left(\sqrt{\alpha^2 + N_f} + \eta \frac{\sqrt{N_f}}{N_c^2} + \sqrt{N_c} \sqrt{1 - \theta^2}\right)\right]} \right\} \frac{4\pi g_f}{(2\pi)^3} \frac{N_c^3}{N_f^{3/2}} \Lambda_{QCD}^3$$

It looks like space for a sliver of confining but percolating quark matter remains @  $N_c = 3, \Lambda_{QCD} \mu_Q < \mathcal{O}(2) \Lambda_{QCD}$

## Phenomenology of quarkyonic percolation

**FAIR/RHIC/NICA** the logical place to see the form factor is in dilepton spectral functions,  $q\bar{q} \rightarrow ll$ . Note that periodicity of form factor will lead to suppression of signal at  $M^2 \sim \Lambda_{QCD}^2$  (usually signal is enhancement)



**neutron stars**  $P \sim N_c$  means a “quasi-QGP” is formed at  $\mu_B \sim m_B$ , but without mixed phase (percolation is second order.  $P \sim N_c \sim \mathcal{O}(3)$  in proto-neutron stars means proto-neutron star much stiffer than “normal” nuclear proto-neutron star. **Could quarkyonic matter make stars explode?**

## Conclusions, discussion and outlook

- A percolation-type “quarkyonic” transition, distinct from deconfinement, appears natural in  $T - \rho - N_c$  space
- At  $\mu_Q \sim \Lambda_{QCD}$  the critical  $N_c \sim \mathcal{O}(10)$ , so nuclear density not percolation. however, percolation could happen at  $\Lambda_{QCD} < \mu_Q < \sqrt{N_c/N_f} \Lambda_{QCD}$
- A natural analogy for constructing an EFT of percolating matter is the conductor-insulator transition, baryons fulfil the role of “atomic potential wells”, quarks of the “free electron gas”
- This picture has phenomenological consequences: A dip in the spectral function accessible at FAIR, and the effect of “electron gas pressure” in neutron and proto-neutron stars

## References

- [1] L. McLerran and R. D. Pisarski, Nucl. Phys. A **796**, 83 (2007)
- [2] T. Kojo, Y. Hidaka, L. McLerran and R. D. Pisarski, Nucl. Phys. A **843**, 37 (2010)
- [3] S. Lottini and G. Torrieri, Phys. Rev. Lett. **107**, 152301 (2011)
- [4] S. Lottini and G. Torrieri, arXiv:1204.3272 [nucl-th].