Dynamic event-by-event fluctuations and domain formation at the phase transition of QCD

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with:

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Quark Matter 2012, Washington DC

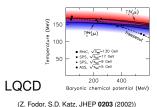


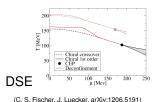




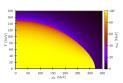
Approaches to the QCD phase diagram

QCD calculations in the nonperturbative regime:





Effective models of QCD:



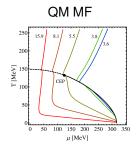
Heavy-ion collisions:



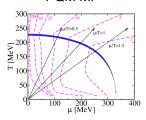




Phase diagrams and fluctuations in effective models

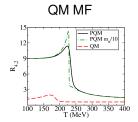


POM MF

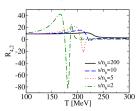


- Crossover (CO), critical point (CP), first order phase transition (FO) in (P)QM models.
- Location (T_c, μ_c) of the CP depends on model parameters, it is non-universal.
- At $\mu_B=0$ the phase transition can be tuned from CO to FO by changing the quark-meson coupling constant.

Phase diagrams and fluctuations in effective models



PQM MF



- Crossover (CO), critical point (CP), first order phase transition (FO) in (P)QM models.
- Location (T_c, μ_c) of the CP depends on model parameters, it is non-universal.
- At μ_B = 0 the phase transition can be tuned from CO to FO by changing the quark-meson coupling constant.
- Enhanced fluctuations in the quark number susceptibilities/kurtosis, increasing from CO to CP.

(V. Skokov, B. Stokic, B. Friman and K. Redlich, PRC 82 (2010), V. Skokov, B. Friman and K. Redlich, PRC 83 (2011), V. Skokov, B. Friman and K. Redlich, PRD 83 (2012))

Phase diagram and fluctuations in heavy-ion collisions

Coupling of the order parameter to pions $g\sigma\pi\pi$ and protons $G\sigma\bar{p}p\Rightarrow$ fluctuations in multiplicity distributions

$$\langle (\delta N)^2 \rangle \propto \langle (\Delta \sigma)^2 \rangle \propto \xi^2$$

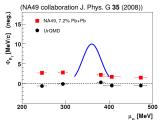
 ξ : correlation length of fluctuations of the order parameter, diverges at the CP

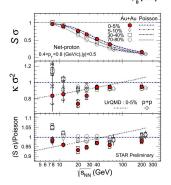
(M. Stephanov, K. Rajagopal, E. Shuryak, PRL 81 (1998), PRD 60 (1999))

Higher cumulants are more sensitive to the CP

$$\begin{split} &\langle (\delta \textit{N})^3 \rangle \propto \xi^{4.5} \\ &\langle (\delta \textit{N})^4 \rangle - \langle (\delta \textit{N})^2 \rangle^2 \propto \xi^7 \end{split}$$

(M. Stephanov, PLB 102 (2009), PRL 107 (2011))





(STAR collaboration, QM2012)

Phase transitions in dynamic systems

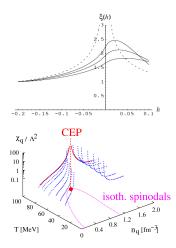
 Critical point: large relaxation times limit the growth of ξ (critical slowing down)

(B. Berdnikov, K. Rajagopal, Phys. Rev. D 61 (2000))

 First order phase transition: instability of slow modes at the spinodal lines (spinodal decomposition)

(I. Mishustin, PRL 82 (1999); C. Sasaki, B. Friman,

K. Redlich, PRD 77 (2008))



Fluctuations are different, but all are interesting!

- Crossover: remnants of the $\mathcal{O}(4)$ criticality.
- Critical point: divergent event-by-event fluctuations in thermodynamic equilibrium.
- First order phase transition: large nonstatistical fluctuations in η/p_T spectra in individual events.

Motivation: Heavy-ion collisions are dynamic, inhomogeneous and finite in space and time.

- ? Can nonequilibrium effects become strong enough to develop signals of the first order phase transition?
- ? Do enhanced equilibrium fluctuations at the critical point survive the dynamics?

Nonequilibrium chiral fluid dynamics - $N\chi$ FD

 Langevin equation for the sigma field: damping and noise from the interaction with the quarks

$$\partial_{\mu}\partial^{\mu}\sigma + rac{\delta U}{\delta\sigma} + g
ho_{s} + \eta\partial_{t}\sigma = \xi$$

- For PQM: phenomenological dynamics for the Polyakov-loop (poster by C. Herold).
- Fluid dynamic expansion of the quark fluid = heat bath, including energy-momentum exchange $\partial_{\mu}T_{\mathbf{q}}^{\mu\nu}=S^{\nu}=-\partial_{\mu}T_{\sigma}^{\mu\nu}$
- Nonequilibrium equation of state $p = p(e, \sigma)$.
- Fluid dynamic initial conditions from UrQMD (www.urqmd.org).

Selfconsistent approach within the 2PI effective action!

(MN, S. Leupold, C. Herold, M. Bleicher, PRC 84 (2011); MN, S. Leupold, M. Bleicher, PLB 711 (2012); MN, M. Bleicher, S. Leupold,

I. Mishustin, arXiv:1105.1962; and work in progress)

Not included (yet)

- Simulations with pion fields
- Quantum dynamics (e. g. J. Berges et al., PRL 107 (2011) 061301)
- Fluid dynamic fluctuations and viscosities
 - (J. Kapusta, B. Mueller, M. Stephanov, PRC 85 (2012); Acta Phys. Polon. B 43 (2012))
- Final state interaction

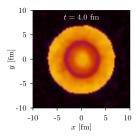
Results - overview

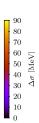
- $\mu_B = 0$ (no fluctuations in net quark number)
 - Domain formation at the first order phase transition.
 - Event-by-event fluctuations at the critical point.
- μ_B ≠ 0
 - Trajectories in the phase diagram.
 - Formation of domains of high quark number density at the first order phase transition.

Dynamic domain formation

First order phase transition

Sigma field fluctuations:
$$\Delta \sigma = \sqrt{(\sigma - \sigma_{eq})^2}$$

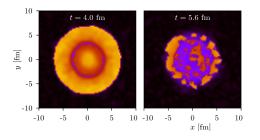


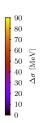


• Highly supercooled state at t = 4.0 fm/c.

Dynamic domain formation

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$$\Delta\sigma=\sqrt{(\sigma-\sigma_{eq})^2}$$



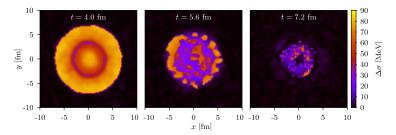


- Highly supercooled state at t = 4.0 fm/c.
- Dynamic formation of domains at t = 5.6 fm/c.

Dynamic domain formation

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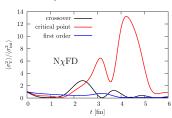


- Highly supercooled state at t = 4.0 fm/c.
- Dynamic formation of domains at t = 5.6 fm/c.
- Dynamic decay of domains at t = 7.2 fm/c.

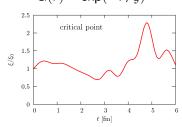
This could lead to non-statistical fluctuations in hadron multiplicities.

Dynamic enhancement of event-by-event fluctuations

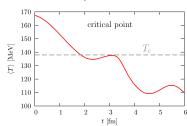
Event-by-event fluctuations



correlation length from $G(r) \propto \exp(-r/\xi)$

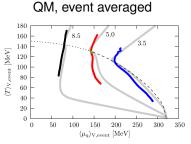


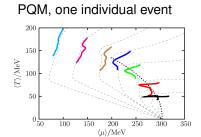
temperature



- Enhanced event-by-event fluctuations of the order parameter, $\langle \sigma_{\rm V}^2 \rangle \propto \xi^2$.
- Initial fluctuations are washed out during the first 1 fm.
- Delay between the averaged T_c and the peak in ξ and ⟨σ_V²⟩.

Trajectories and isentropes at finite μ_B

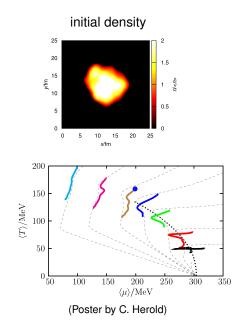




- Grey: equilibrium isentropes, color: T-μ_B trajectories from simulations.
- Fluid trajectories differ from the (equilibrium) isentropes due to interaction with the fields.
- Including the effective Polyakov-loop field leads to more structure of the trajectories, in particular at the first order phase transition (Poster by C. Herold).

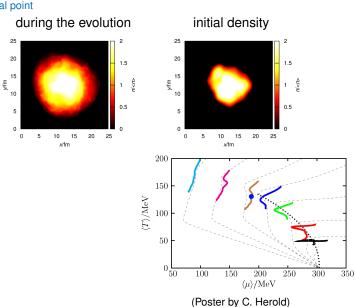
Evolution of quark number density

Crossover - critical point



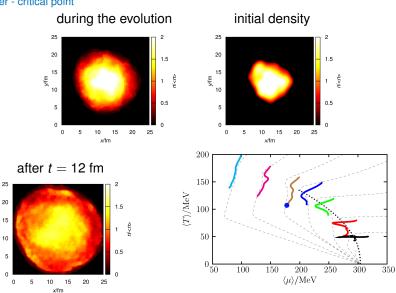
Evolution of quark number density

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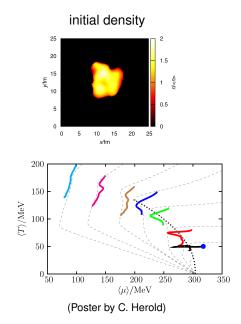
Evolution of quark number density

Crossover - critical point

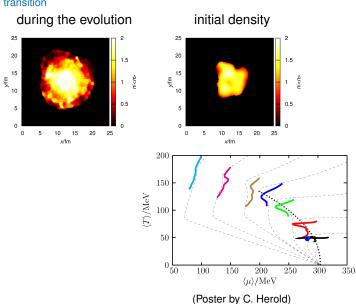


(Poster by C. Herold)

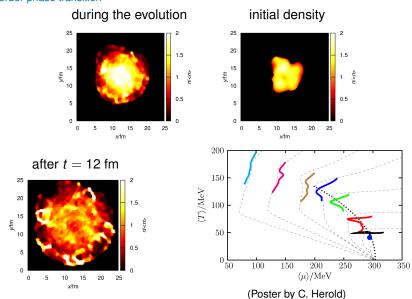
Formation of high quark number density domains



Formation of high quark number density domains



Formation of high quark number density domains



Summary



- Dynamic domain formation (σ and n) at the first order phase transition.
- ? Can nonequilibrium effects become strong enough to develop signals of the first order phase transition?
- Dynamic correlation length ξ grows up to \simeq 2.5 fm.
- Dynamic enhancement of event-by-event-fluctuations of the order parameter (σ) at the critical point.
- ? Do enhanced equilibrium fluctuations at the critical point survive the dynamics?



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- Dynamic domain formation (σ and n) at the first order phase transition.
- √ Yes, nonequilibrium dynamics can lead to signals at the first order phase transition!
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- √ Yes, nonequilibrium dynamics can lead to signals at the first order phase transition!
 - Dynamic correlation length ξ grows up to \simeq 2.5 fm.
 - Dynamic enhancement of event-by-event-fluctuations of the order parameter (σ) at the critical point.
- √ Yes, the critical fluctuations develop even in the situation of a nonequilibrium, dynamic simulation of heavy-ion collisions!



BACKUP

Semiclassical equation of motion for the sigma field

$$\partial_{\mu}\partial^{\mu}\sigma+rac{\delta extbf{ extit{U}}}{\delta\sigma}+ extbf{ extit{g}}
ho_{ extbf{ extit{s}}}+\eta\partial_{ extit{t}}\sigma=\xi$$

damping term η and noise ξ for $\mathbf{k}=0$

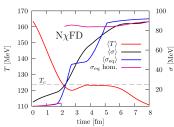
$$\eta = g^2 rac{d_q}{\pi} \left(1 - 2n_{
m F} \left(rac{m_\sigma}{2}
ight)
ight) rac{\left(rac{m_\sigma^2}{4} - m_q^2
ight)^{rac{3}{2}}}{m_\sigma^2} \int\limits_{rac{40}{500}}^{rac{40}{500}} \int\limits_{rac{100}{500}}^{rac{40}{500}} \int\limits_{rac{100}{500}}^{rac{40}{500}} \int\limits_{rac{700}{500}}^{racc} \int\limits_{rac{700}{500}}^{rac{400}{500}} \int\limits_{rac{700}{500}}^{racc} \int\limits_{rac{700}{500}}^{racc} \int\limits_{
ho}^{200}} \int\limits_{
ho}^{200}} \int$$

below T_c damping by the interaction with the hard pion modes, apply $\eta=2.2/\mathrm{fm}$

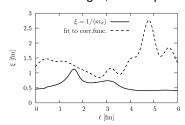
(T. S. Biro and C. Greiner, PRL 79 (1997))

Influence of the inhomogeneous system

first order phase transition

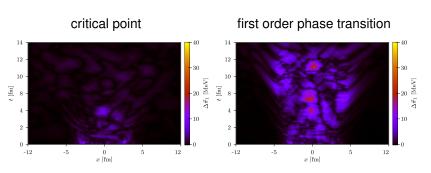


correlation length, critical point



Pion fluctuations

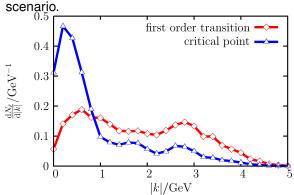
So far: pion fluctuations were not considered and $\vec{\pi}=\langle\vec{\pi}\rangle=0$. Propagate pion fluctuations, too:



Larger pion fluctuations in a scenario with a first order phase transition!

Box: after equilibration

Low mode fluctuations of the Polyakov-loop are enhanced in a CP scenario



(Poster by C. Herold)