

Instantons and Sphalerons in a Magnetic Field

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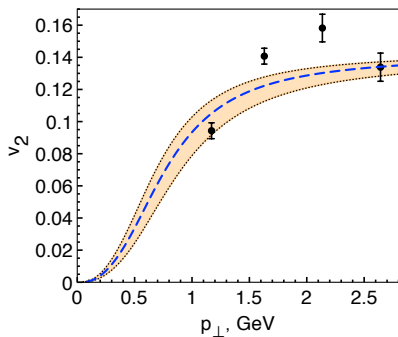
Quark Matter 2012, Washington D.C.

GB, G.Dunne & D. Kharzeev , arXiv:1112.0532, PRD **85** 045026

GB, D. Kharzeev, arXiv:1202.2161, PRD **85** 086012

Magnetic field generated in heavy ion collisions $\sim m_\pi^2$
combined with:

- ▶ Axial Anomaly \Rightarrow C.M.E. (charge separation)
C.M.W. (charge dependent v_2)
- ▶ Conformal Anomaly \Rightarrow photon v_2
(Başar, Kharzeev, Skokov, arXiv:1206.1334)



(talk by V. Skokov at xQCD, 08/22)

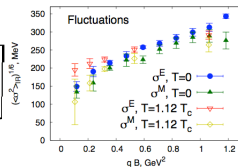
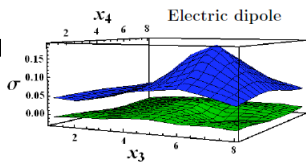
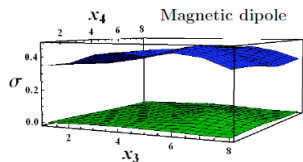
Part I

Instanton in a Magnetic Field

Motivation & some lattice results

Interplay between topology & magnetic field

- ▶ Chiral magnetic effect $\vec{J} \propto \mu_5 \vec{B}$
- ▶ What sources μ_5 ? sphalerons, η domains, etc..
- ▶ Instanton + magnetic field
- ▶ Lattice results
 - ▶ ITEP group (electric & dipole moments)
 - ▶ T. Blum et al. (zero modes $\propto B$)
 - ▶ A. Yamamoto (C.M. conductivity)



(Polikarpov et al. '09)

Notation & conventions

work in: \mathbb{R}^4

chiral basis: $\gamma_\mu = \begin{pmatrix} 0 & \alpha_\mu \\ \bar{\alpha}_\mu & 0 \end{pmatrix}$, $\gamma_5 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$$\alpha_\mu = (1, -i\vec{\sigma}) \quad , \quad \bar{\alpha}_\mu = (1, i\vec{\sigma}) = \alpha_\mu^\dagger$$

Dirac operator: $\mathcal{D} = \begin{pmatrix} 0 & \alpha_\mu \mathcal{D}_\mu \\ \bar{\alpha}_\mu \mathcal{D}_\mu & 0 \end{pmatrix} \equiv \begin{pmatrix} 0 & D \\ -D^\dagger & 0 \end{pmatrix}$

gauge field: $\mathcal{A}_\mu = A_\mu + a_\mu$

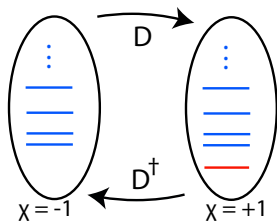
Notation & conventions

diagonal form: $(i\mathcal{D})^2 \psi_\lambda = \begin{pmatrix} DD^\dagger & 0 \\ 0 & D^\dagger D \end{pmatrix} \psi_\lambda = \lambda^2 \psi_\lambda$

$$\chi = +1 : \quad DD^\dagger = -\mathcal{D}_\mu^2 - \mathcal{F}_{\mu\nu} \bar{\sigma}_{\mu\nu}$$

$$\chi = -1 : \quad D^\dagger D = -\mathcal{D}_\mu^2 - \mathcal{F}_{\mu\nu} \sigma_{\mu\nu}$$

”supersymmetry:” for $\lambda \neq 0$, DD^\dagger and $D^\dagger D$ has identical spectra



Instanton & magnetic field

$$DD^\dagger = -\mathcal{D}_\mu^2 - B\sigma_3, \quad , \quad D^\dagger D = -\mathcal{D}_\mu^2 - F_{\mu\nu}\sigma_{\mu\nu} - B\sigma_3$$

Zero modes: Both spins, both chiralities

Index thm: $\text{tr} \left(\mathcal{F}_{\mu\nu} \tilde{\mathcal{F}}_{\mu\nu} \right) = \text{tr} \left(F_{\mu\nu} \tilde{F}_{\mu\nu} \right)$

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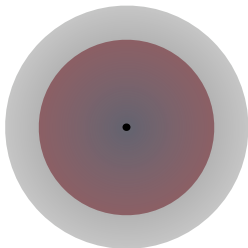
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Competition between instanton and magnetic field

Instanton zero mode: $|\psi_0|^2 = \frac{64\rho^2}{(x^2+\rho^2)^3}$

Topological charge: $q_5(x) = \frac{192\rho^4}{(x^2+\rho^2)^4}$



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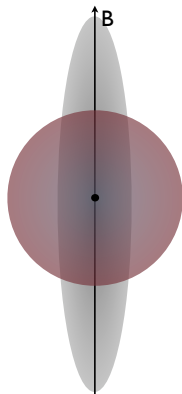
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B field zero mode: $|\psi_0|^2 \propto f(x_1 + ix_2)e^{-B|x_1+ix_2|^2}$

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Large instanton limit

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after appropriate gauge rotation & Lorentz transformation:

$$\mathcal{A}_\mu = -\frac{F}{2}(-x_2, x_1, -x_4, x_3)\tau^3 + \frac{B}{2}(-x_2, x_1, 0, 0)\mathbb{1}_{2 \times 2}$$

quasi-abelian, covariantly constant \rightarrow soluble!

Large instanton limit

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$$\mathcal{F}_{12} = \begin{pmatrix} B - F & 0 \\ 0 & B + F \end{pmatrix}$$

$$\mathcal{F}_{34} = \begin{pmatrix} -F & 0 \\ 0 & F \end{pmatrix}$$

Landau problem with field strengths \mathcal{F}_{12} & \mathcal{F}_{34}

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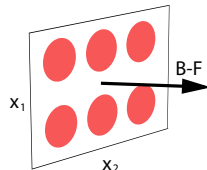
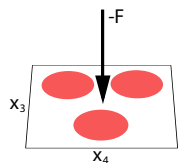
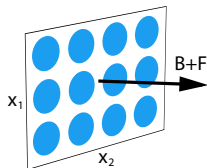
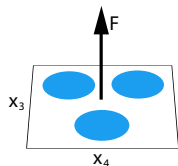
Landau problem with field strengths \mathcal{F}_{12} & \mathcal{F}_{34}

Zero modes

$$\tau = -1, \chi = -1, \text{spin} \uparrow, n_- = \frac{(B+F)}{2\pi} \frac{F}{2\pi}$$

$$\tau = +1, \chi = +1, \text{spin} \uparrow, n_+ = \frac{(B-F)}{2\pi} \frac{F}{2\pi}$$

$$n_+ + n_- = B \frac{F}{2\pi^2}, \quad n_+ - n_- = -\frac{F^2}{2\pi^2}$$



Dipole moments

$$\sigma_3^M = \frac{1}{2} \langle \bar{\psi} \Sigma_{12} \psi \rangle \quad , \quad \sigma_3^E = \langle \bar{\psi} \Sigma_{34} \psi \rangle$$

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$$m \langle \bar{\psi} \Sigma_{12} \psi \rangle = \text{tr}_{2 \times 2} \left(\sigma_3 \frac{m^2}{m^2 + DD^\dagger} \right) + \text{tr}_{2 \times 2} \left(\sigma_3 \frac{m^2}{m^2 + D^\dagger D} \right)$$

$$m \langle \bar{\psi} \Sigma_{34} \psi \rangle = -\text{tr}_{2 \times 2} \left(\sigma_3 \frac{m^2}{m^2 + DD^\dagger} \right) + \text{tr}_{2 \times 2} \left(\sigma_3 \frac{m^2}{m^2 + D^\dagger D} \right)$$

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$$m \langle \bar{\psi} \Sigma_{12} \psi \rangle \approx \frac{BF}{2\pi^2}$$

$$m \langle \bar{\psi} \Sigma_{34} \psi \rangle \approx \frac{F^2}{2\pi^2}$$

$$\langle \bar{\psi} \Sigma_{34} \psi \bar{\psi} \Sigma_{34} \psi \rangle \approx \left(\frac{F}{2\pi^2 m^2 L^4} \right) B$$

► $\sigma_3^M > \sigma_3^E$

Part II

Sphaleron Rate in a Magnetic Field

Sphaleron rate (basics)

$$\Gamma_{CS} = \frac{(\Delta Q_5)^2}{Vt} = \int d^4x \left\langle \frac{g^2}{32\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}(x) \frac{g^2}{32\pi^2} F_{\alpha\beta}^a \tilde{F}_a^{\alpha\beta}(0) \right\rangle$$

Diffusion of topological charge: $\frac{dN_5}{dt} = -c N_5 \frac{\Gamma_{CS}}{T^3}$

- ▶ CP odd effects in QCD (CME)
- ▶ Baryon number (B+L) violation in E.W.

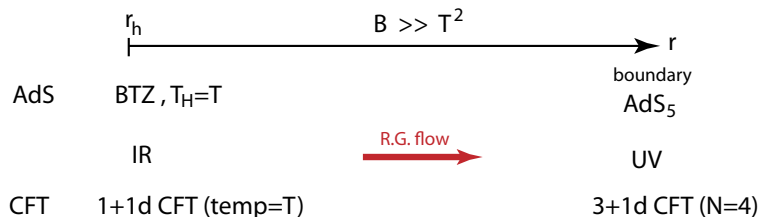
Weak coupling: $\Gamma_{CS} = \kappa g^4 T \log(1/g) (g^2 T)^3$ (Bödeker '98)

Strong coupling: $\Gamma_{CS} = \frac{(g^2 N)^2}{256\pi^3} T^4$ (Son, Starinets '02)

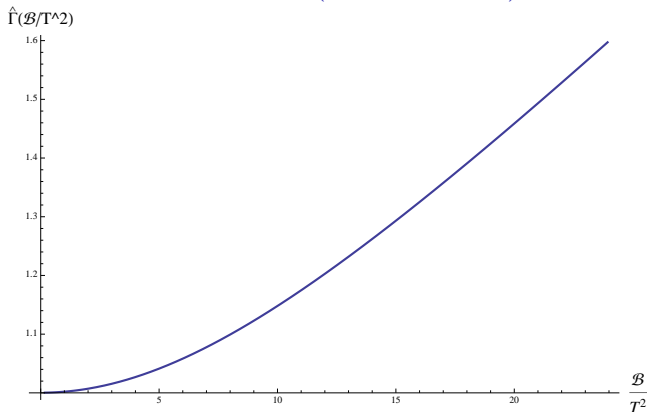
Sphaleron rate with B field (holography)

Gauge theory in magnetic field \Leftrightarrow Einstein-Maxwell theory

Dynamics with magnetic field \Leftrightarrow Self-consistent solutions



Sphaleron rate with B field(holography)



$$\Gamma_{CS} = \begin{cases} \frac{(g^2 N)^2}{256\pi^3} \left(T^4 + \frac{1}{6\pi^4} B^2 + \mathcal{O}\left(\frac{B^4}{T^2}\right) \right) & , \quad B \ll T^2 \\ \frac{(g^2 N)^2}{384\sqrt{3}\pi^5} \left(B T^2 + 15.9 T^4 + \mathcal{O}\left(\frac{T^6}{B}\right) \right) & , \quad B \gg T^2 \end{cases}$$

Sphaleron rate with B field (holography)

$$B \sim T^2$$

- ▶ for $B = T^2$ the effect is $\sim 0.2\%$
- ▶ it is safe to ignore the effects of B field on Γ_{CS} for CME estimates

Sphaleron rate with B field (holography)

$$B \gg T^2$$

$$\Gamma_{CS} = \frac{(g^2 N)^2}{384\sqrt{3}\pi^4} \frac{B}{\pi} \times T^2$$

Sphaleron rate with B field (holography)

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Landau level density

$$\Gamma_{CS} = \frac{(g^2 N)^2}{384\sqrt{3}\pi^4} \overset{\uparrow}{\frac{B}{\pi}} \times T^2$$

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↑
↓
diffusion scale in 1+1d

Conclusions

- ▶ Anomalies + magnetic field have a rich structure
- ▶ Electric and magnetic dipole moments
- ▶ Zero modes play a crucial role
- ▶ 1st order derivative expansion captures some lattice results
- ▶ Confinement ? (instantons with nonzero holonomy), CSB ?
- ▶ At strong coupling:
 - ▶ Magnetic field always increases the sphaleron rate
 - ▶ Back-reaction of magnetic field into non-abelian sector
 - ▶ Strong magnetic field leads to dimensional reduction
- ▶ Weak coupling?