

THERMAL MESON PROPERTIES AND CHIRAL SYMMETRY



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Recent and ongoing work on thermal meson properties relevant for the hadron gas regime include transport coefficients, chemical nonequilibrium, susceptibilities, isospin breaking and different aspects of chiral symmetry restoration. Chiral Perturbation Theory ensures the model independency of the low- T regime. This is combined with unitarity when an accurate analytical description of particle scattering is needed, as in the case of resonance studies and transport coefficients. Results agree with several phenomenological, theoretical and lattice analysis.

TRANSPORT COEFFICIENTS FOR A PION GAS IN CHPT

Kubo formula for current-current spectral function:

$$\text{Im} \int \omega \rightarrow 0^+, \vec{0} \rightarrow G^a(\omega, |\vec{p}|) G^a(\omega, |\vec{p}|) \approx \frac{\pi}{2E_p \Gamma_p} \delta(\omega^2 - E_p^2) E_p^2 = |\vec{p}|^2 + m_\pi^2$$

Chiral Power Counting for TC modified:

Lines with same 4-momentum (doubled-lined) prop to inverse pion width:

$$\text{Diagram} \rightarrow Y \sim \frac{n_B(E)}{E_p \Gamma_p} \quad (\text{instead of ChPT } O(p^2))$$

Diagrams potentially dominant:

Bubbles $\sim O(Y^{k+1})$

Ladders $\sim O(Y^{k+1} p^{2k})$

$T \ll m_\pi$: relevant three momentum $p \approx \sqrt{m_\pi T}$

$$\bar{n}_\pi \sim (\sqrt{m_\pi T})^3 e^{-m_\pi/T} \ll 1 \Rightarrow Y \sim \sqrt{\frac{m_\pi}{T}} \text{ potentially large}$$

$$\Gamma_p \sim \sigma_{\pi\pi} v_\pi \bar{n}_\pi \sim \sqrt{T} \bar{n}_\pi$$

loop spectral function $\sim \frac{p_{CM}}{m_\pi} \sim \frac{\sqrt{m_\pi T}}{m_\pi} = \frac{1}{Y}$

$T \approx m_\pi$: relevant $p \approx T \Rightarrow$ unitarized scattering in Γ relevant

$\sigma_{\pi\pi}^{\pi\pi}$ softer in $\Gamma_U \Rightarrow$ larger T effects in TC

Ladders still perturbative $Y \sim O(1)$

Higher T behaviour signals ChPT break-up and ladder resummation

SUSCEPTIBILITIES, ISOSPIN AND CHIRAL SYMMETRY

Conn/Disc separation of Scalar susceptibility for n_f light flavors:

$$\chi_S \equiv -\frac{\partial}{\partial m} \langle \bar{q}q \rangle_T = \frac{1}{3V} \frac{\partial^2}{\partial m^2} \log Z_{QCD} = n_f^2 \chi_{dis} + n_f \chi_{con} \quad (IL)$$

$$\chi_{dis} = \langle (\text{Tr} D^{-1})^2 \rangle_A - \langle \text{Tr} D^{-1} \rangle_A^2 \quad \chi_{con} = -\langle \text{Tr} D^{-2} \rangle_A$$

$$\chi_{con} \approx \frac{1}{2} (\chi_{uu} + \chi_{dd}) - \chi_{ud} = -\frac{1}{2} \partial_{m_s} \langle \bar{u}u - \bar{d}d \rangle_T$$

Fluctuations of $\langle \bar{q}q \rangle$. More sensitive to chiral restoration, maximum near critical region

$m_u \neq m_d$ allows to separate χ_{con}/χ_{dis} in the effective theory.

ChPT provides their model-independent T, m_q behaviour, relevant e.g. for scaling lattice results near $(T, m_q) \sim (T_c, 0^+)$.

For $n_f = 2$ and $m_u \neq m_d$:

$$\chi_S = \chi_{uu} + \chi_{dd} + 2\chi_{ud} = -\partial_{m_s} \langle \bar{u}u + \bar{d}d \rangle_T \quad \text{Isosinglet } (\sigma)$$

$$\chi_{con} = \frac{1}{2} (\chi_{uu} + \chi_{dd}) - \chi_{ud} = -\frac{1}{2} \partial_{m_s} \langle \bar{u}u - \bar{d}d \rangle_T \quad \text{Isotriplet } (a_3)$$

$$\chi_{dis} = \chi_{ud}$$

Sensitive to $O(m_s)$ corrections $\Rightarrow \chi_{con} \neq 0$ in the isospin limit contrary to naive expectation $\chi_{uu} = \chi_{dd} = \chi_{ud}$.

Actually, taking the IL after diff $\rightarrow \left\{ \begin{array}{l} \chi_{uu} \approx \chi_{dd} \approx \chi_S/4 - \chi_{con}/2 \\ \chi_{ud} \approx \chi_S/4 + \chi_{con}/2 \end{array} \right.$

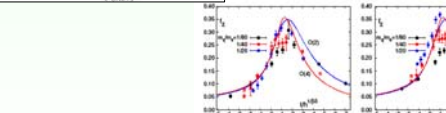
\Rightarrow conn.part cancels in χ_S but not in the one-flavour equivalent $4\chi_{uu}$.

Up to 30% corrections for $4\chi_{uu}/\chi_S$ at $T = 0$ coming from χ_{con}

"Connected contamination" in the definition may explain (partly) $O(N)$ lattice scaling violations

χ_{con} gets contributions from "false" GB-loops from extra tastes in staggered formalism

Condensates OK with scaling



T=0 subtracted Ejri et al '09 unsubtracted

Lattice points Aoki et al '09

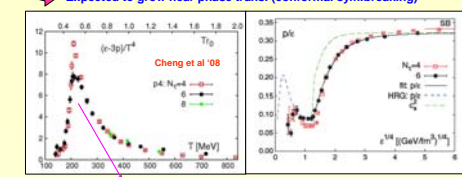
OK with lattice when normalizing to respective T_c

Bulk Viscosity

$$J \rightarrow -T_c^3/3 - c_s^2 T_{00} \quad c_s^2 = \partial P / \partial \epsilon$$

Related to QCD trace anomaly $\langle \theta \rangle \equiv (T_\mu^\mu) = \epsilon - 3P = T^S \frac{d}{dT} \left(\frac{P}{T^3} \right)$ suggested by finite- T sum rules (plus spectral function assumptions).

Expected to grow near phase trans. (conformal sym. breaking)



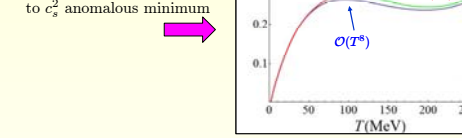
Glulon dominated: $(T_\mu^\mu)_{QCD} = \frac{\beta(g)}{2g} G_{\mu\nu}^a G_{\mu\nu}^a + (1 + \gamma_m(g)) \bar{q} M q$

In ChPT: $\zeta(T) = \int_0^\infty d^3p \frac{3p^2(p^2/3 - c_s^2 E_p^2)^2}{4\pi^2 T E_p^2} n_B(E_p) [1 + n_B(E_p)]$

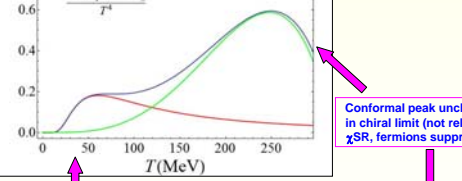
Pion gas in chemical nonequilibrium between CFO and TFO (elastic scattering only)

Chiral limit: $\zeta = 15(1/3 - c_s^2)^2 \eta T \gg m_\pi^2 \rightarrow 0$ (conformal limit w/o anomaly)

But ChPT to $O(T^8)$ sensitive to c_s^2 anomalous minimum

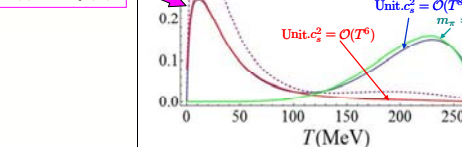


Conformal peak unchanged in chiral limit (not related to χ_S , fermions suppressed)



Low-T peak from fermion mass conformal breaking

$$\delta(\theta)_T \sim 2m_q \delta(\bar{q}q)_T$$



Strong correlation between bulk visco and trace anomaly established within ChPT (no further assumptions)

$\chi_{con}(T), \chi_{dis}(T)$ explicit expressions found in LO ChPT for $m \neq 0, N_f = 2 + 1$, including $m_u \neq m_d$ and EM corrections (virtual photons)

Relevant T behaviour can be approximated by the IR/Chiral + Isospin limit:

$$M_\pi \ll T \ll M_K, M_\eta \quad \text{and} \quad m_u \rightarrow m_d \text{ after differentiation}$$

$$\frac{4\bar{m}^2}{M_\pi^2} \chi_{dis}^R(T=0) = -\frac{3}{32\pi^2} \log \frac{M^2}{\mu^2} + 32L_6^*(\mu) + \frac{1}{288\pi^2} \left(-28 + 5 \log \frac{M_\eta^2}{\mu^2} \right)$$

$$\frac{4\bar{m}^2}{M_\pi^2} \chi_{con}^R(T=0) = 8 [H_2^*(\mu) + 2L_6^*(\mu)] - \frac{1}{16\pi^2} \left(1 + \log \frac{M_K^2}{\mu^2} + 2 \log \frac{M_\eta^2}{\mu^2} \right)$$

$$\frac{4\bar{m}^2}{M_\pi^2} [\chi_{dis}(T) - \chi_{dis}(0)]^{IR} = \frac{3T}{16\pi M_\pi} \quad \text{Critical}$$

$$\frac{4\bar{m}^2}{M_\pi^2} [\chi_{con}(T) - \chi_{con}(0)]^{IR} = \frac{3T}{18M_\pi^2} \quad \text{Regular}$$

Larger scale, comes from IB $\pi^0 \eta$ mixing

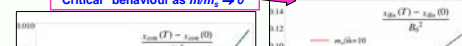
Deviations from chiral IR important!



"Critical" behaviour as $m/m_s \rightarrow 0$



Isospin corrections amplified $\sim O(T/\bar{m})$



A.Gómez Nicola, J.R.Peláez, J.Ruiz de Elvira, *Scalar susceptibilities and four-quark condensates in the meson gas*, in prep 2012.

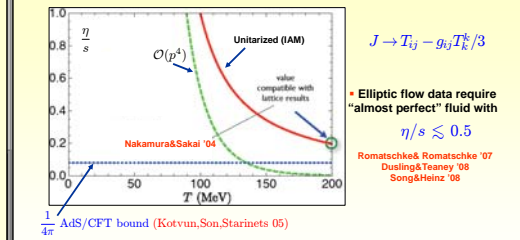
A.Gómez Nicola, R.Torres Andrés, *Isospin breaking and Chiral Symmetry Restoration*, Phys.Rev.D83:076005, 2011.

D.Fernández-Fraile, A.Gómez Nicola, *Transport coefficients and resonances for a meson gas in Chiral Perturbation Theory*, Eur.Phys.J. C62:37-54, 2009.

D.Fernández-Fraile, A.Gómez Nicola, *Bulk viscosity and the conformal anomaly in the pion gas*, Phys.Rev.Lett.102:121601, 2009.

D.Fernández-Fraile, A.Gómez Nicola, *The electrical conductivity of a pion gas*, Phys.Rev.D73:045025, 2006.

Shear Viscosity



Compatible with analysis based on KT (Dobado et al '04, Prakash et al '93)

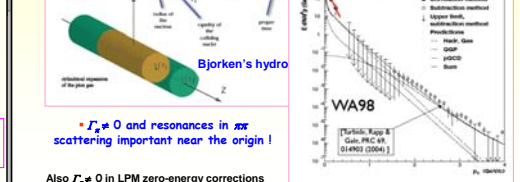
Minimum of η/s expected at phase trans: $O(N_c^2)$ and decreasing at high T $O(1)$ and increasing at high T Arnold et al

Electrical conductivity and photons

DC cond. related to soft-photon emitted by the meson (pion) gas: $J \rightarrow J^EM$

$$\frac{dR_1}{d^3q} = \frac{1}{8\pi} n_B(\omega) \rho_\mu^\mu(\omega, \vec{q}) \Rightarrow \frac{dR_1}{d^3q}(\omega \rightarrow 0^+, \vec{q} = 0) = \frac{1}{4\pi} 3T\sigma(T)$$

$$\frac{dN_1}{d^3q}(p \rightarrow 0) \approx 2\pi F_{\mu\nu}^a \frac{\partial}{\partial p} \langle T(r) \rangle(r) \quad \text{Ward identity: } q_\mu \rho_\mu^\mu = 0 \Rightarrow \rho_{00}(\omega \neq 0, \vec{q} = 0) = 0$$



Also $J_\mu \neq 0$ in LPM zero-energy corrections W.Liu, R.Rapp '07

$T \ll m_\pi$ limit for TC compatible with NR Kinetic Theory:

Electrical Conductivity:

$$\sigma(T) \approx 15e^2 \frac{f_\pi^2}{m_\pi^2 \sqrt{T}} \approx e^2 \frac{N_c \lambda}{m_\pi} \text{ with mean free path } \lambda \approx 1/\Gamma$$

$$\Gamma \approx n_\pi v_\pi \sigma_{\pi\pi} \approx n_\pi \left(\sqrt{\frac{m_\pi T}{m_\pi}} \right) \left(\frac{m_\pi}{f_\pi} \right)$$

Shear and bulk viscosities:

$$\zeta(T) \approx 13.3 \frac{f_\pi^4 \sqrt{T}}{m_\pi^2} \sim 0.36\eta(T) \approx m_\pi v_\pi n_\pi \lambda$$

Pseudoscalar/Scalar Degeneration @ Chiral Symmetry Restoration

Pseudoscalar Susceptibility $\rightarrow \int_0^\beta \int d^3x \langle T(\bar{q}\gamma_5 \tau^a q)(x) (\bar{q}\gamma_5 \tau^b q)(0) \rangle = \chi_P \delta^{ab}$

From Current Algebra (LO ChPT, $T = 0$):

$$\chi_P = 4B_0^2 \frac{f_\pi^2}{M_\pi^2} = \frac{\langle \bar{q}q \rangle}{m_q} \quad \text{whereas, e.g. from L}\sigma\text{M, } \chi_S = 4B_0^2 \frac{f_\pi^2}{M_\sigma^2}$$

Pion pole Signa pole

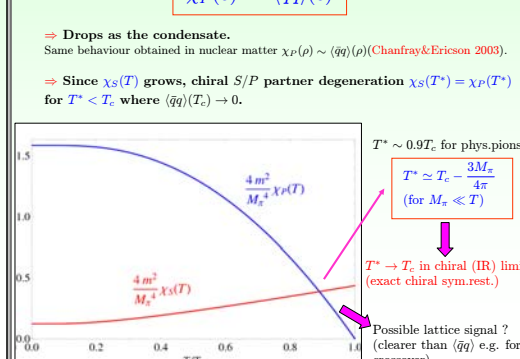
NLO (incl T corrections) ChPT Model-independent result: (coupling external p^0 sources)

$$\frac{\chi_P(T)}{\chi_P(0)} = \frac{\langle \bar{q}q \rangle(T)}{\langle \bar{q}q \rangle(0)}$$

\Rightarrow Drops as the condensate.

Same behaviour obtained in nuclear matter $\chi_P(\rho) \sim \langle \bar{q}q \rangle(\rho)$ (Chanfray&Ericson 2003).

\Rightarrow Since $\chi_S(T)$ grows, chiral S/P partner degeneration $\chi_S(T^*) = \chi_P(T^*)$ for $T^* < T_c$, where $\langle \bar{q}q \rangle(T_c) \rightarrow 0$.



Possible lattice signal? (clearer than $\langle \bar{q}q \rangle$ e.g. for crossover)