# A proposal to solve some puzzles in charmed semileptonic *B* decays

Phys.Rev. D85 (2012) 094033 or arXiv:1202.1834

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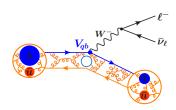
July 5, 2012



ICHEP 2012 Melbourne Australia



#### Outline

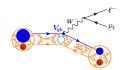


[Illustration by F. Tackmann]

- I. Introduction: Summary of the exp. and theo. situation
  - a Recap of incl. and excl. measurements
  - b Recap of the '1/2' vs '3/2' problem
- II. Discovery of potential 25 charmed state(s) by BABAR
- III. Our Proposal and its Viability
- IV. Prediction of  $\Gamma(B \to D'^{(*)} \ell \bar{\nu}_{\ell})$  using *light-cone sum rules*
- V. Summary

### I.a Experimental situation for $B o X_c \ell \bar{\nu}_\ell$

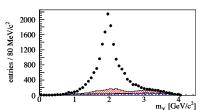
- BABAR and Belle:  $1.1 \text{ ab}^{-1}$  at  $\Upsilon(4S)$
- $\approx 25\%$  of all B decay semileptonic



Notation	$s_l^{\pi_l}$	J <sup>P</sup>	m (GeV)	Γ (GeV)
D	1 -	0-	1.87	
D*	$\frac{1}{2}$ -	$1^{-}$	2.01	
$D_0^*$	1 + 2	0+	2.40	0.28
$D_1^*$	1/2 +	1+	2.44	0.38
$D_1$	3 +	1+	2.42	0.03
$D_2^*$	$\frac{3}{2}$ +	$2^+$	2.46	0.04
D'	1 -	0_	2.54	0.13
D'*	$\frac{1}{2}$ -	1-	2.61	0.09

- Most abundant  $b \to c$ :  $\mathcal{B}(B^+ \to X_c \ell^+ \nu_\ell) = (10.92 \pm 0.16) \%$ 

 $X_c$ : charmed system; isospin averaged value from [HFAG]



- Major focus of experimental attention from B factories
- Inclusive  $X_c$  mass spectrum (not unfolded;  $p_l^* > 0.8$ ) [PRD81:032003]
- Presence of charm decays up to ≈ 3 GeV (resolution 0.36 GeV)
- $D^{(*)}$ ,  $D^{**}$ ,  $D'^{(*)} \leftrightarrow 1S$ , 1P, 2S

### I.a Experimental situation for $B o X_c \ell \bar{\nu}_\ell$

Charm state $X_c$	$\mathcal{B}(B^+ \to X_c \ell^+ \nu)$	-
D	$(2.31 \pm 0.09)\%$	
D*	$(5.63 \pm 0.18)\%$	
$\sum D^{(*)}$	$(7.94 \pm 0.20)\%$	
$D_0^*  o D \pi$	$(0.41 \pm 0.08)\%$	broad states
$D_1^* \rightarrow D^* \; \pi$	$(0.45 \pm 0.09)\%$	$(0.86 \pm 0.12)\%$
$\mathit{D}_1  \rightarrow  \mathit{D}^*   \pi$	$(0.43 \pm 0.03)\%$	narrow states
$\textit{D}_{2}^{*} \rightarrow \textit{D}^{(*)}\pi$	$(0.41 \pm 0.03)\%$	
$\sum D^{**} \rightarrow D^*\pi$	$(1.70\pm0.12)\%$	
$D \pi$	$(0.66 \pm 0.08)\%$	
$D^*$ $\pi$	$(0.87 \pm 0.10)\%$	
$\sum D^*\pi$	$(1.53\pm0.13)\%$	
$\sum D^{(*)} + \sum D^* \pi$	$(9.47 \pm 0.24)\%$	
$\sum D^{(*)} + \sum D^{**} \to D^{(*)} \pi$	$(9.64 \pm 0.23)\%$	
Inclusive $X_c$	$(10.92 \pm 0.16)\%$	

All values from [HFAG 2010]. For the values of  $B\to D~\pi~\ell~\bar{\nu}_\ell$  and  $B\to D^*~\pi~\ell~\bar{\nu}_\ell$  an uncertainty weighted average of both isospin modes was calculated assuming a 100% correlation between both values.

 $\Rightarrow$  'Gap' of  $(1.45 \pm 0.29)$  % emerges which is not accounted for

Uses semi-inclusive  $D^{(*)}\pi$  branching fractions; with measured 1P  $D^{**}\to D^{(*)}\pi\Rightarrow$  (1.28  $\pm$  0.29) %

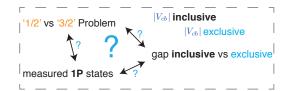
### I.b Theoretical situation of $B \to X_c \ell \bar{\nu}_\ell$

- Comparable rates for the narrow and broad  $D^{**}$  states problematic:

- Uraltsev's sum rule + covariant quark model estimate from [EPJ:C52975]  $\mathcal{B}(B^+ \to D_{1/2=\text{broad}}^{**} \, \ell^+ \, \nu)/\mathcal{B}(B^+ \to D_{3/2=\text{narrow}}^{**} \, \ell^+ \, \nu) \backsim 0.1 0.2$  i.e. clear dominance of narrow over broad.
- Experimental violation known as '1/2' vs '3/2' puzzle.
- Persistent  $\sim 2-3\sigma$  difference between  $|V_{cb}|$  from **inclusive** vs exclusive  $(41.9 \pm 0.4_{\rm exp.} \pm 0.6_{\rm theo.}) \times 10^{-3}$  vs  $(38.7 \pm 0.6_{\rm exp.} \pm 0.5_{\rm theo.}) \times 10^{-3}$

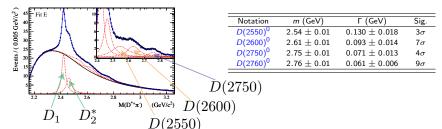
#### [ARNPS:201161119]

- Any connections?



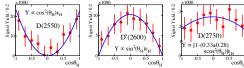
### II. Discovery of new charmed states at BABAR

- BABAR observed four new charmed states [PRD82:111101]:



The  $m_{D^*\pi}$  mass distribution for D(2550), D(2600), D(2750) is shown. D(2760) reconstructed in  $D\pi$  channel.

- Helicity angles of D(2550) and D(2600) helicity consistent with 25:



Helicity angle from  $D^* \to D\pi$  is defined as the angle between the primary pion and  $\pi_{slow}$  from  $D^* \to D\pi_{slow}$  (in the  $D^*$  rest frame)

- D(2750) candidate for 1D ( Likely no relevant for semileptonic decays ightarrow cf. Backup )

### II. Strong decays of 2S states D' & $D'^*$

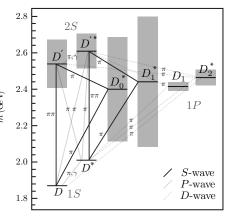
- Strong D' and  $D'^*$  decays:

$$2S \rightarrow 1S$$
or
$$2S \rightarrow 1P \rightarrow 1S$$

E.g. 
$$p ext{-wave} + \pi \to 1S$$
  
 $s ext{-wave} + \pi \to 1P_{ ext{broad}} \ (\to 1S)$   
Mom. of the emitted pion  $p_\pi \backsim 0.01 - 0.5$  GeV  
 $s ext{-wave} + \pi\pi \to 1S$   
 $d ext{-wave} + \pi \to 1P_{ ext{narrow}} \ (\to 1S)$   
Signature  $s ext{-wave}$ :  $D'^{(*)} \to D^{(*)}\pi\pi$   
Signature  $p ext{-wave}$ :  $D'^{(*)} \to D^{(*)}\pi$ 

More decays involving  $\rho$  and  $\eta$  in principle allowed

 Significant 2S → 1P<sub>broad</sub> cross feed plausible [PRD:83014009]



A selection of the allowed strong decays involving single or two pion emissions are illustrated.

### III. Our Proposal and its Viability

Proposal Explore possibility that the sum of  $D'^{(*)}$  rate is substantial,

$$\mathcal{B}(B^+ o \mathcal{D}'^{(*)} \ell^+ \nu_\ell) \backsim \mathcal{O}(1\%)$$

and show that this can help resolve the problems mentioned earlier without giving rise to new ones

1 This is a big enough contribution to the sum over exclusive states to close the gap between **inclusive** and exclusive without e.g. introducing **non-resonant**  $B^+ \to D^{(*)} \pi \ell^+ \nu_\ell$  contributions.

A large non-resonant rate at high  $D^*\pi$  invariant mass would disagree with the inclusive lepton spectrum and the measured semi-exclusive  $B^+ \to D^{(*)} \pi \, \ell^+ \, \nu_\ell$  rate

- 2 The  $D'^{(*)}$  states can decay with one pion in an *s*-wave to members of the  $s_l^{\pi} = \frac{1}{2}^+$  states, and could thus enhance the observed decay rate to the  $\frac{1}{2}^+$ , and thus give rise to the '1/2' vs '3/2' puzzle.
- 3 With the relatively low mass of the  $D'^{(*)}$  the lepton spectrum can stay quite hard, in agreement with the observations
- 4 The  $\mathcal{B}(B^+ \to D^*\pi \, \ell^+ \, \nu_\ell)$  semi inclusive measurement is not in conflict with our hypothesis, since the decay of the  $D'^{(*)}$  would yield two or more pions most of the time.

⇒ full details in Phys.Rev. D85 (2012) 094033 or arXiv:1202.1834

# IV. Prediction for $\Gamma(B^+ \to D'^{(*)} \ell^+ \nu_\ell)$

-  $D'^{(*)}$  and  $D^{(*)}$ : identical quantum numbers

i.e. same formulae for decay rate and definitions of form factors

$$\frac{\mathrm{d}\Gamma_{D'^*}}{\mathrm{d}w} = \frac{G_F^2 |V_{cb}|^2 m_B^5}{48\pi^3} r^3 (1-r)^2 \sqrt{w^2-1} (w+1)^2 \left[ \frac{\mathrm{d}\Gamma_{D'}}{\mathrm{d}w} = \frac{G_F^2 |V_{cb}|^2 m_B^5}{48\pi^3} r^3 (1+r)^2 (w^2-1)^{3/2} \left[ G(w) \right]^2, \right] \times \left[ 1 + \frac{4w}{w+1} \frac{1-2rw+r^2}{(1-r)^2} \right] \left[ F(w) \right]^2,$$

where  $r = m_{D'(*)}/m_B$  and  $w = v \cdot v'$  denotes the recoil parameter, where v denotes the velocity of the B meson, and v' of the D'(\*).

- In  $m_{b,c} \gg \Lambda_{\rm QCD}$  limit: 6 form factors  $\rightarrow$  single universal Isgur-Wise function  $\zeta_2(W)$  i.e.  $F(w) = G(w) = \zeta_2(w)$
- Heavy quark symmetry:  $\zeta_2(w=1)=0$ 
  - ightarrow Non-zero rate at zero recoil entirely due to  $\Lambda_{
    m QCD}/m_{b,c}$  corrections
- For w > 1 no power suppression, but low kinematic range of 1 < w < 1.3 role of  $\Lambda_{\rm QCD}/m_{b,c}$  corrections can be very large.
- Naive expectation:  $\frac{\mathrm{d} \zeta_2}{\mathrm{d} w}|_{w=1} > 0$  In quark model main effect of wave function of the brown muck is to increase the expectation value of the distance from the heavy quark of a spherically symmetric wave function. Overlap of initial and final state wave functions should increase as w increases.

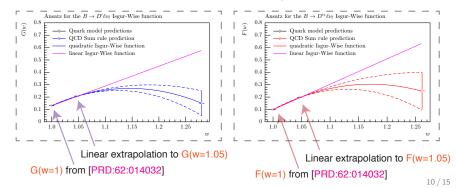
### IV. The $B^+ \to D'^{(*)} \ell^+ \nu_{\ell}$ form factors

Not easy to calculate the  $B^+ \to D'^{(*)} \ell^+ \nu_{\ell}$  form factors:

- a Quark model [PRD:62:014032] hoped to be trustable near w = 1.
- b Modify QCD light-cone sum rule calculation [PJC:60603] hoped to be reasonable near max. recoil

But Both models were developed, tuned, and tested for states that are the lightest within a given set of quantum numbers, thus take prediction with truck load of salt. But even rough estimates can be helpful!

Quark Model form factors at w = 1 and linear extrapolation to w = 1.05:



## IV. The $B^+ \to D'^{(*)} \ell^+ \nu_\ell$ form factors

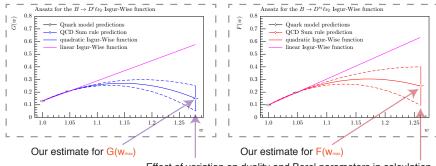
Modify QCD light-cone sum rules so that the 25 state can be projected out e.g. schematically for the decay constant

$$\left| \begin{array}{c} \frac{m_D^4 f_D^2}{m_C^2 (m_D^2 - q^2)} + \frac{m_D^4 f_{D'}^2}{m_C^2 (m_{D'}^2 - q^2)} + \int_{s_0^{D'}}^{\infty} \mathrm{d}s \frac{\rho(s)}{s - q^2} \end{array} \right|$$

where  $\rho$  is the spectral density function, and  $f_D$  and  $f_{D'}$  denote the 1S and 2S decay constant, respectively.

Modification of Borel transformation in [PJC:60603] non-trivial endeavor.

Form factors sensitive fo chosen decay constants, Borel, and duality parameters



# IV. Prediction for $\Gamma(B^+ \to D'^{(*)} \ell^+ \nu_\ell)$

Parametrize F(w) and G(w) which determine the  $D'^{(*)}$  as quad. polynom. i.e.

$$F(w) = \beta_0^* + (w - 1)\beta_1^* + (w - 1)^2 \beta_2^*,$$

$$G(w) = \beta_0 + (w - 1)\beta_1 + (w - 1)^2 \beta_2.$$

 $\Rightarrow$  Rough estimate for sum of two semileptonic  $B^+ \to D'^{(*)} \ell^+ \nu_{\ell}$  decays:

$$\mathcal{B}(B^+ \to D'^{(*)} \, \ell^+ \, \nu_\ell) \backsim (0.3 - 0.7) \%$$

Earlier quark models without accounting for  $\Lambda_{\rm QCD}/m_{b,c}$  effects obtained smaller rates, c.f.

[PRD:39799], [PTP:91757]. Including  $\Lambda_{\rm QCD}/m_{b,c}$  effects a value of 0.4% was obtained by [PRD:62:014032].

With a linear parametrization and the quark model result only:

$$\mathcal{B}(B^+ o D'^{(*)}\,\ell^+\,
u_\ell)\sim 1.4\,\%$$

We take this as an indication that a large radial contribution is plausible, and that  $B^+ \to D'^{(*)} \, \ell^+ \, \nu_\ell$  may account for a substantial part of the observed 'Gap' between inclusive and exclusive decays.

## V. Summary and Ideas

- Indication that hypothesis plausible and that B → D'(\*) ℓ v̄<sub>ℓ</sub> may account for a substantial part of the observed 'gap'.
- Interesting measurement for LHCb (or *B*-factories):  $B \to D'^{(*)} \pi = [D^{(*)} \pi^+ \pi^-] \pi^-$ Factorization [PRL:87201806] implies relation between these channels and semileptonic decay rate at  $w_{\text{max}}$ :

$$|\Gamma(B \to D'^{(*)} \pi) = \frac{3\pi^2 C^2 |V_{ud}|^2 f_{\pi}^2}{m_B m_{D'}(*)} \frac{d\Gamma(B \to D'^{(*)} \ell \bar{\nu}_{\ell})}{dw} \Big|_{w = w_{\text{max}}}$$

C combination of Wilson coefficients with C  $|V_{ud}| \approx$  1, and  $w_{\max}$  corresponds to  $q^2 = 0 \backsimeq m_\pi^2$ 

- If future measurement find a  $B \to D'^{(*)} \ell \bar{\nu}_\ell$  decay rate ... the precise determination of the branching fraction and form factors would impact other measurements and the theory of semileptonic decays, e.g. it may yield a better understanding . . .
  - i. ... of  $b \rightarrow c$  backgrounds and improve  $|V_{ub}|$  and  $|V_{cb}|$
  - ii. ... missing exclusive contributions to **inclusive**  $B \to X_c \ell \bar{\nu}_\ell$
  - iii. ... of the measured  $B \to D^{(*)} \, \tau \, \bar{\nu}_{\tau}$  and its tension with the SM

#### **Further**

- iv. Help improve the measurements of the semileptonic branching fractions of the  $s_l^{\pi} = \frac{1}{2}$  and  $\frac{3}{2}$  states, thus maybe help resolving the '1/2' vs '3/2' puzzle
- v. Help improve the sum rule bound on the  $B \to D^* \ell \bar{\nu}_{\ell}$  form factor.

#### Thank you for your attention!

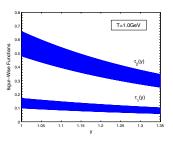
# Backup

### A. Prediction for 1D from QCD sum rules

 QCD sum rule result of [PRD:79034025] suggests that 1D contributions to the inclusive semileptonic decay rate are small

Decay	PRD:79034025	PLB:478408
$B  o D_1^* \ell \bar{\nu}_\ell$	$6 \times 10^{-6}$	
$B \to D_2' \ell \bar{\nu}_\ell$	$6 \times 10^{-6}$	
$B  o D_2  \ell  \bar{ u}_\ell$	$1.5 \times 10^{-4}$	$1 \times 10^{-5}$
$B  o D_3^* \ell \bar{\nu}_\ell$	$2.1 \times 10^{-4}$	$1 \times 10^{-5}$

The branching fractions for the four  ${\bf 1D}$  states are quoted. Note that the  $D_1^*$  is not identical with the  ${\bf 1P}$  state with the same name (which is sometimes denoted as  $D_1'$  to avoid this confusion)



The Isgur-Wise functions for the  $\frac{3}{2}$  and the  $\frac{5}{2}$  1D doublets as a function of the recoil param. y~(=w) are shown.