



Generalized Galileons for Particle Physics & Cosmology

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- We'll see - too early to know if these will be useful or not - but it is turning out to be great fun trying.



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$$\pi(x) \rightarrow \pi(x) + c + b_\mu x^\mu$$

The Galilean symmetry!



Galileons



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Can consider this symmetry as interesting in its own right

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Computing Feynman diagrams - terms of the galilean form cannot receive new contributions! More soon.

Luty, Porrati, Rattazzi (2003); Nicolis, Rattazzi (2004)



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So forces much smaller than gravitational strength within the Vainshtein radius - hence safe from 5th force tests.



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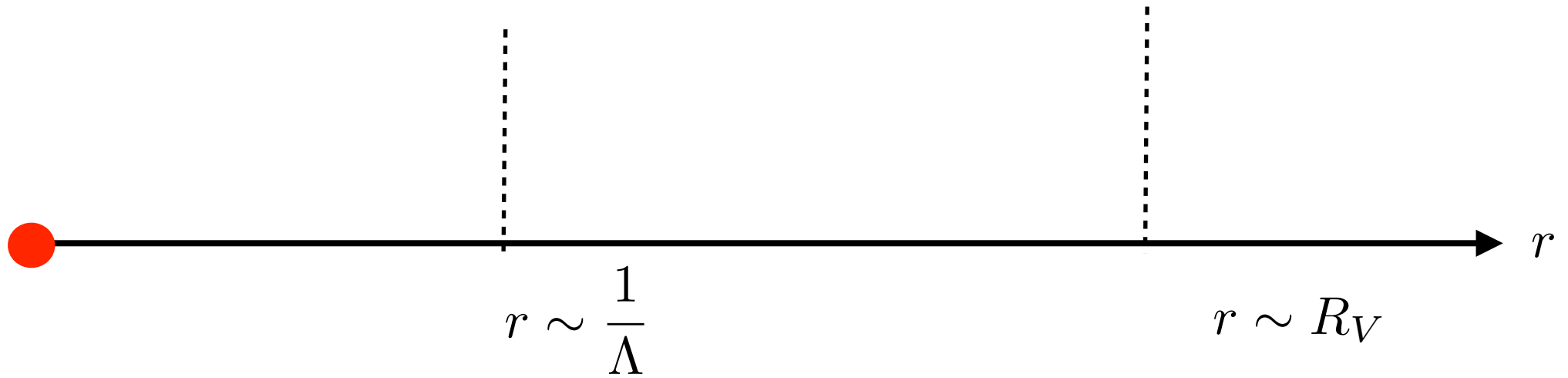
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Thus, if we canonically normalize the kinetic term of the perturbations, we raise the effective strong coupling scale, and, more importantly, heavily suppress the coupling to matter!



Regimes of Validity



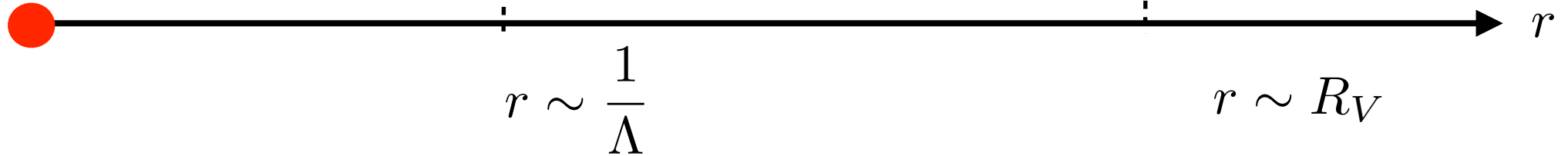


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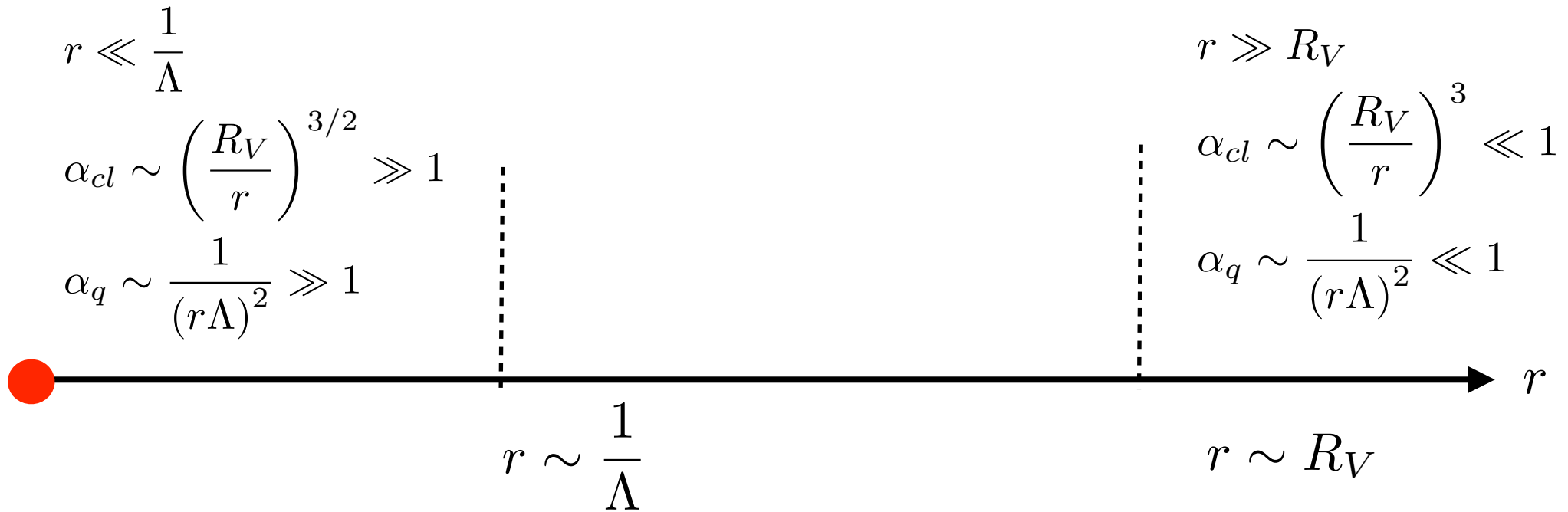
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Multi-field Galileons and Higher co-Dimension Branes



Higher co-Dimension Probe Branes

[K. Hinterbichler, M.T., D.Wesley, *Phys. Rev. D*82 (2010) 124018.]



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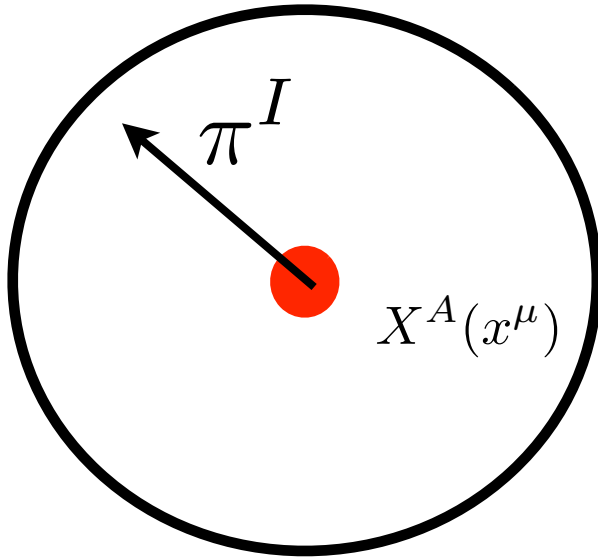


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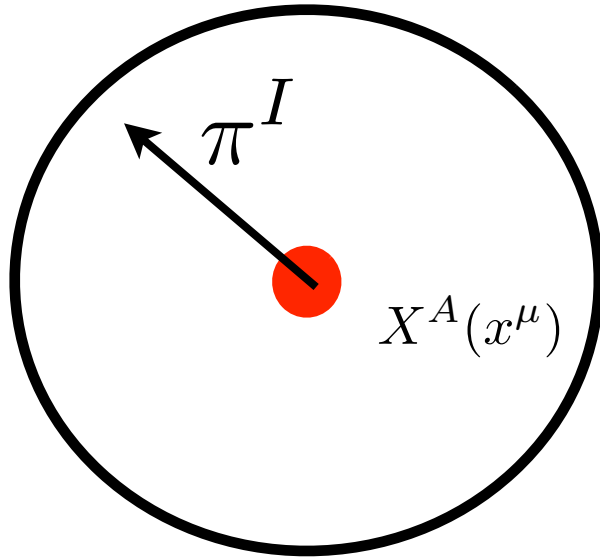
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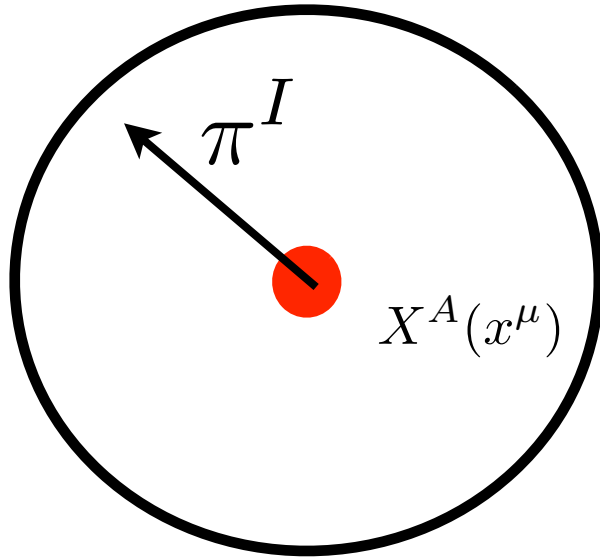
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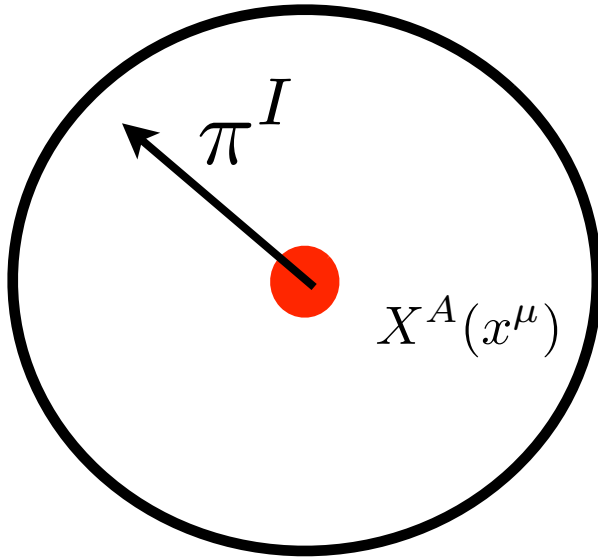
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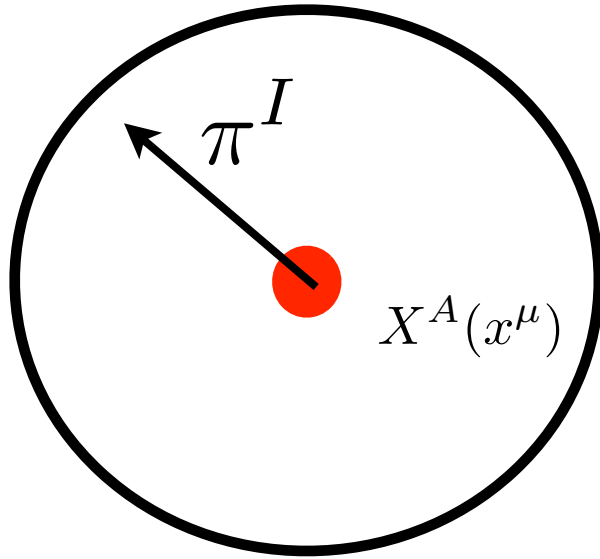
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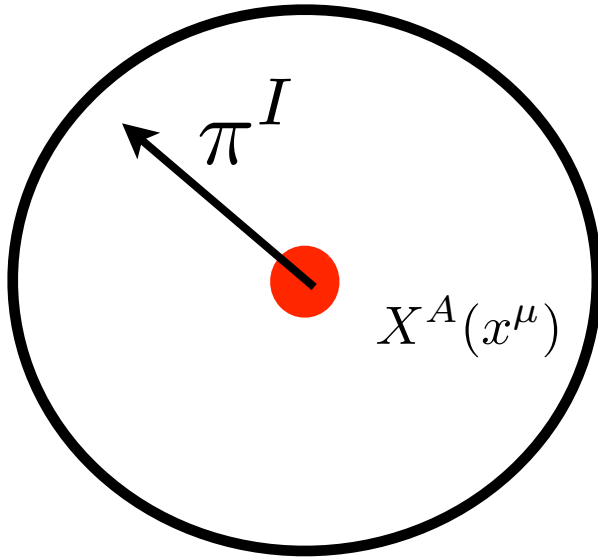
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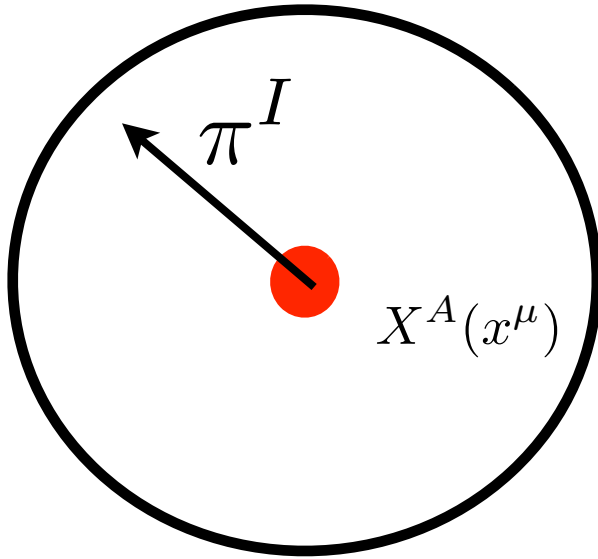
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Normal Bundle Curvature



Higher co-Dimension Probe Branes

[K. Hinterbichler, M.T., D.Wesley, *Phys. Rev. D*82 (2010) 124018.]



With some work, can extend probe brane construction to multiple co-dimensions

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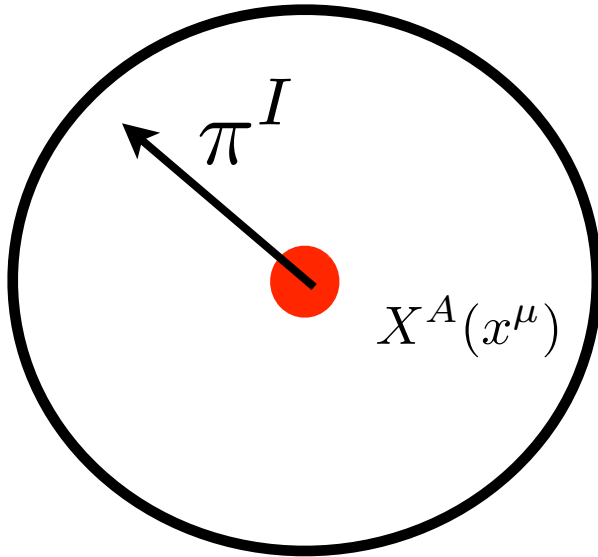
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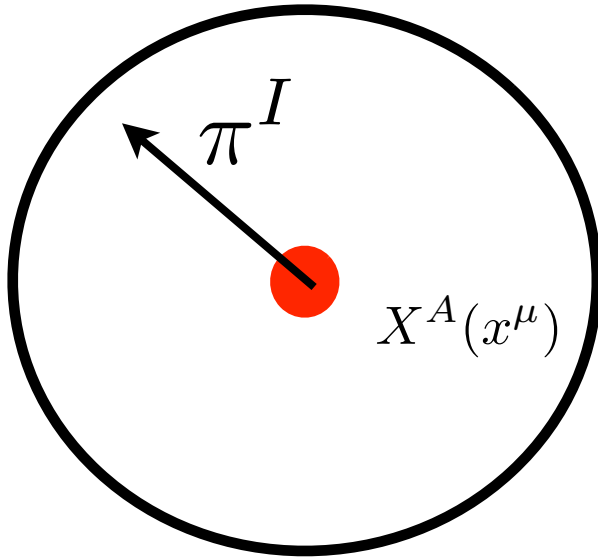
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Normal Bundle Curvature

In co-dimension I, for 2nd order equations, use Lovelock terms and associated boundary terms. Here, for 4d brane, prescription *depends on co-dimension*

If $N = 2$, boundary terms include only brane cosmological constant, and

$$\mathcal{L}_{N=2} = \sqrt{-g} (R[g] - (K^i)^2 + K^i_{\mu\nu} K_i^{\mu\nu})$$



The Multi-Galileon Limit

[K. Hinterbichler, M.T., D.Wesley, *Phys. Rev. D*82 (2010) 124018;

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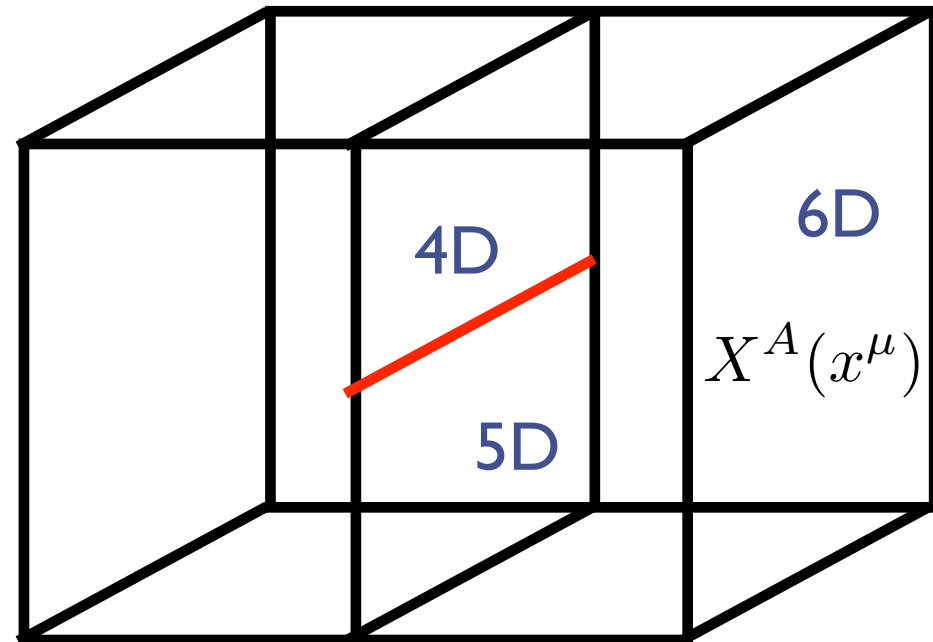
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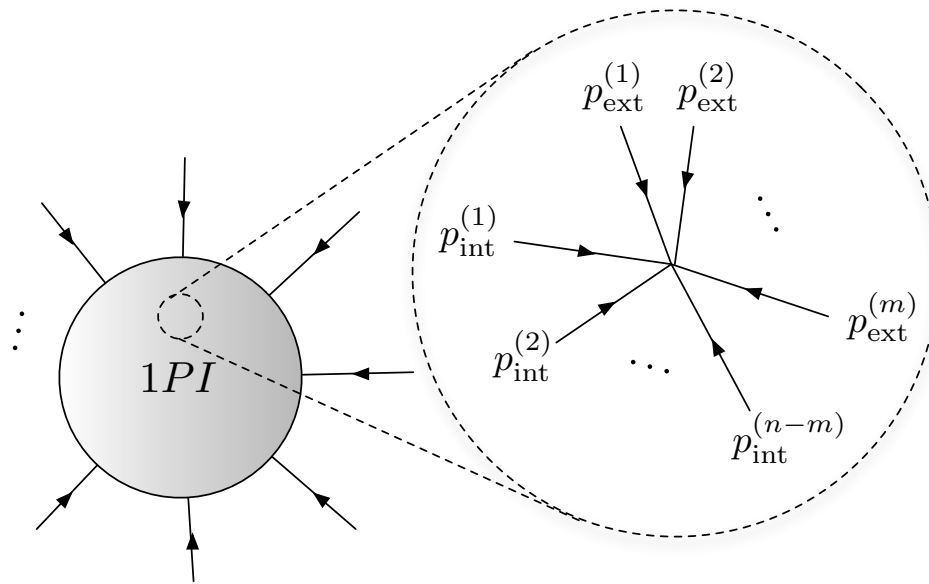
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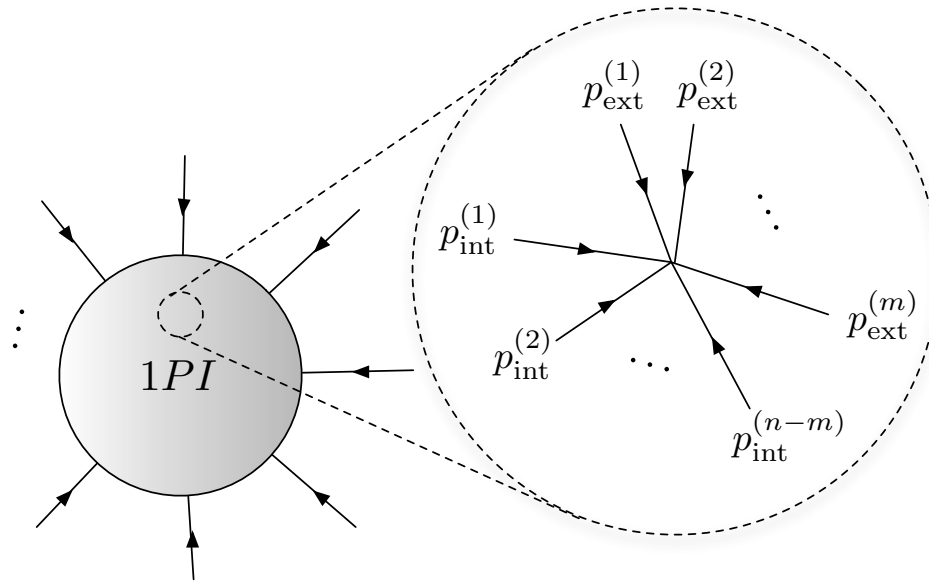




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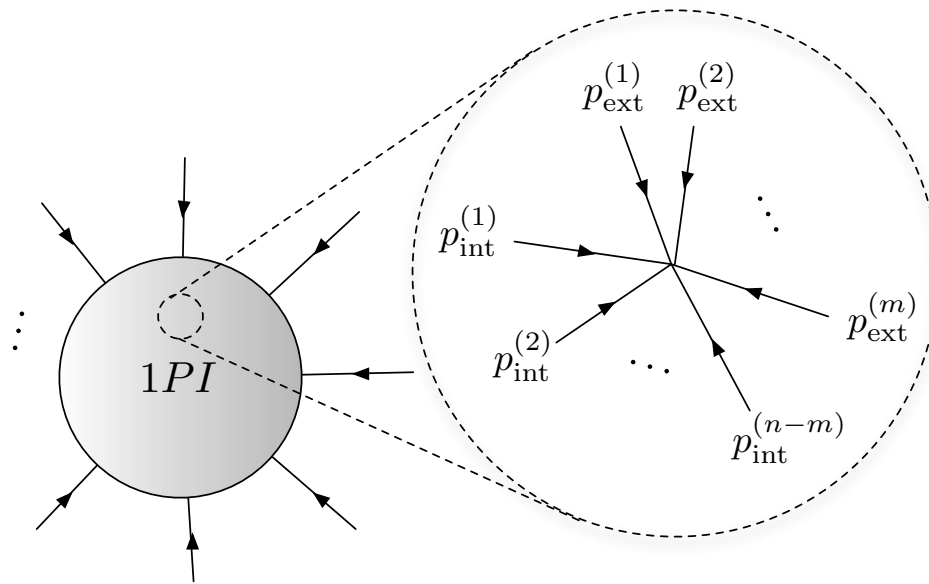
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Can even add a mass term and remains technically natural



Generalized Galileons on Curved Geometries



Galileons on General Backgrounds

[Goon, Hinterbichler, M.T., *Phys. Rev.Lett.* 106, 231102 (2011).
Goon, Hinterbichler, M.T., *JCAP* 1107, 017 (2011).]



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Main point:

- Have emphasized probe brane construction because it can be extended to more general geometries. e.g. other maximally-symmetric examples

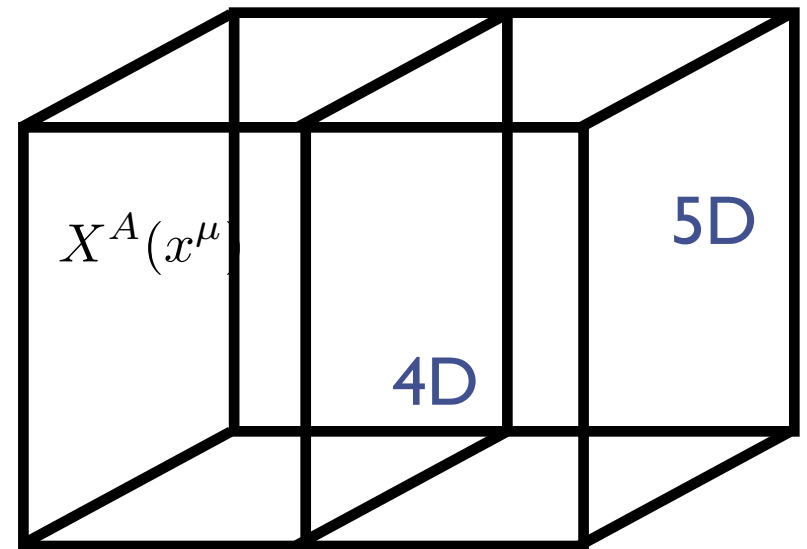


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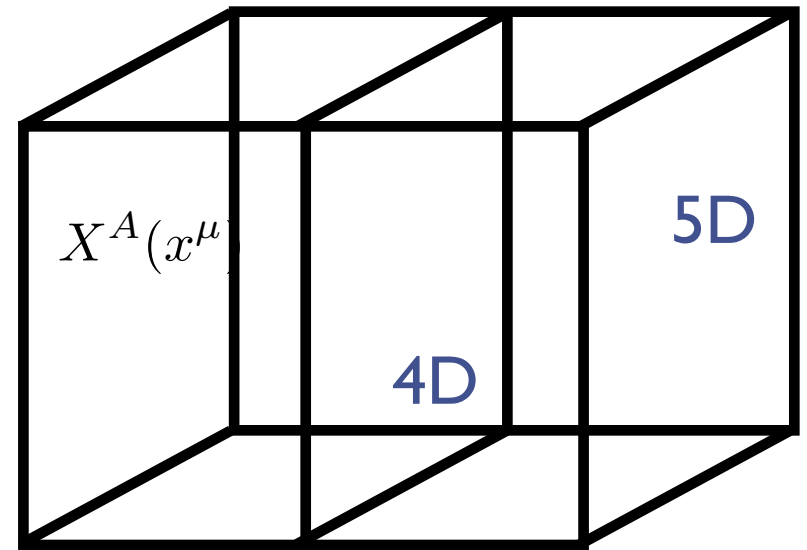
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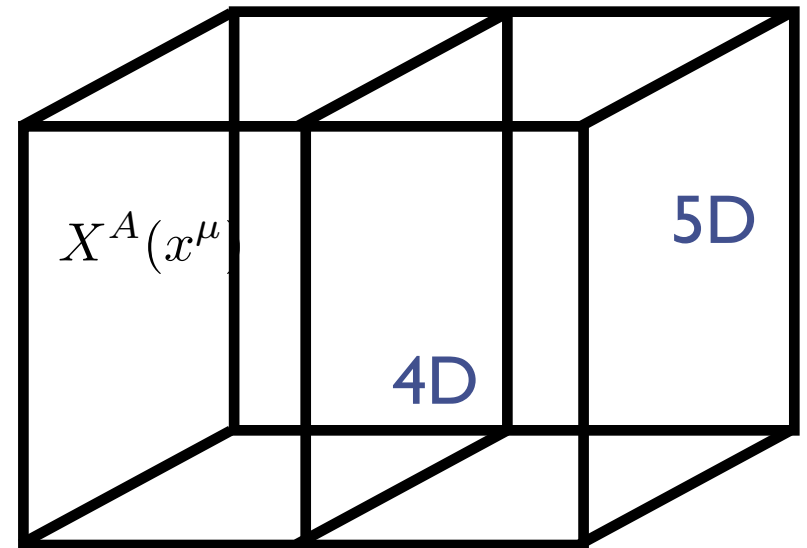
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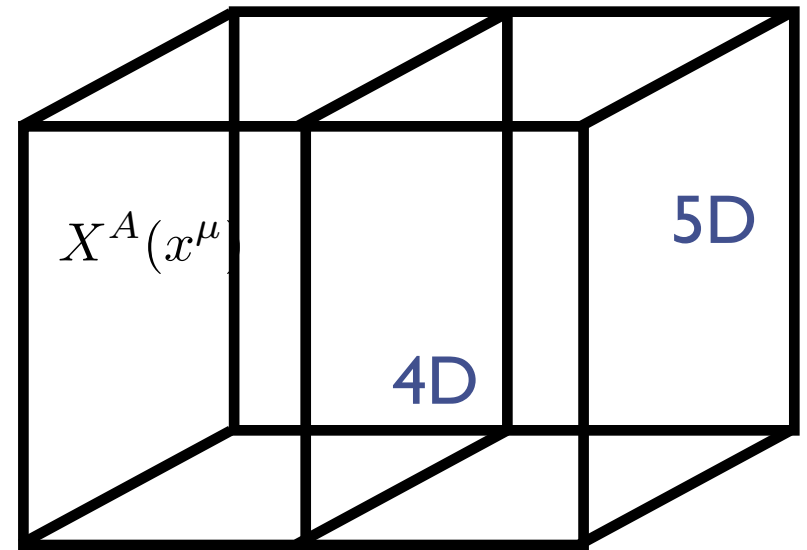
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Bulk Killing Vectors





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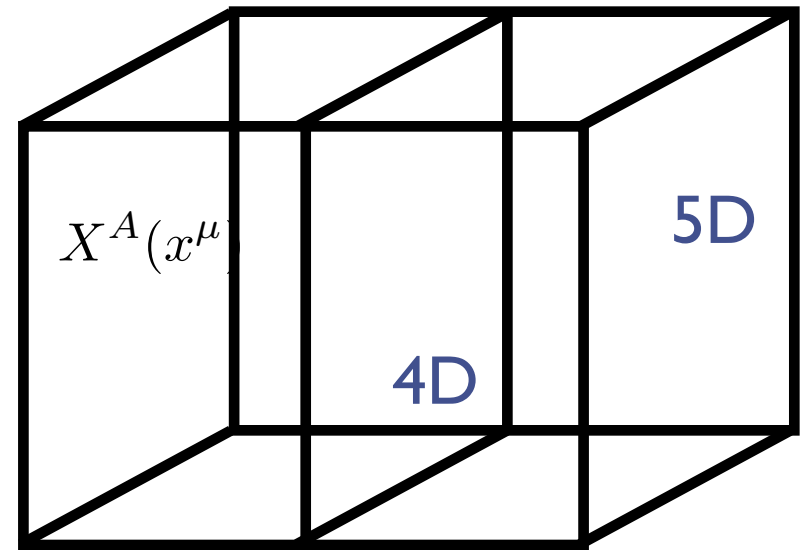
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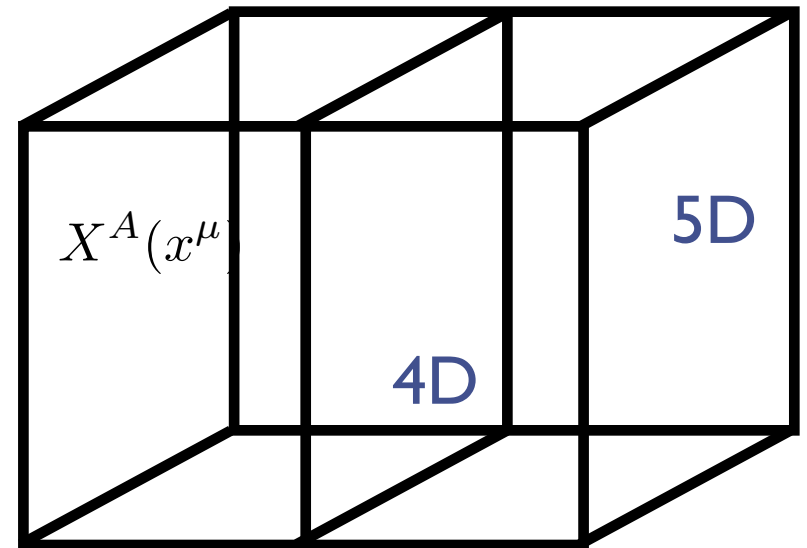
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Galileons with symmetry

$$(\delta_K + \delta_{g,\text{comp}})\pi = -a^i k_i^\mu(x) \partial_\mu \pi + a^I K_I^5(x, \pi) - a^I K_I^\mu(x, \pi) \partial_\mu \pi$$



The Maximally-Symmetric Taxonomy



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Potentially different Galileons corresponding to different ways to foliate a maximally symmetric 5-space by a maximally symmetric 4-d hypersurface



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		Brane metric		
		AdS_4	M_4	dS_4
Ambient metric	AdS_5	AdS DBI galileons $so(4, 2) \rightarrow so(3, 2)$ $f(\pi) = \mathcal{R} \cosh^2(\rho/\mathcal{R})$	Conformal DBI galileons $so(4, 2) \rightarrow p(3, 1)$ $f(\pi) = e^{-\pi/\mathcal{R}}$	type III dS DBI galileons $so(4, 2) \rightarrow so(4, 1)$ $f(\pi) = \mathcal{R} \sinh^2(\rho/\mathcal{R})$
	M_5	X	DBI galileons $p(4, 1) \rightarrow p(3, 1)$ $f(\pi) = 1$	type II dS DBI galileons $p(4, 1) \rightarrow so(4, 1)$ $f(\pi) = \pi$
	dS_5	X	X	type I dS DBI galileons $so(5, 1) \rightarrow so(4, 1)$ $f(\pi) = \mathcal{R} \sin^2(\rho/\mathcal{R})$

Small field limit

↓

AdS galileons

↓

normal galileons

↓

dS galileons



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 - Galilean genesis (alternative to inflation); and in general as a way to violate the null energy condition.



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 - An appearance in the decoupling limit of some massive gravity theories (I'll mention a little more in my plenary)



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- Couplings to matter and stability still need investigating in generality - some simple couplings display instabilities.



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Thank You!



Acknowledgements & References



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- G.Goon, K.Hinterbichler, A.Joyce and M.T., ``Galileons as Wess-Zumino Terms," arXiv:1203.3191 [hep-th].
- G.Goon, K.Hinterbichler, A.Joyce and M.T., ``Gauged Galileons From Branes," arXiv:1201.0015 [hep-th].
- G.Goon, K.Hinterbichler and M.T., ``Galileons on Cosmological Backgrounds," JCAP 1112, 004 (2011) [arXiv:1109.3450 [hep-th]].
- G.Goon, K.Hinterbichler and M.T., ``A New Class of Effective Field Theories from Embedded Branes," PRL 106, 231102 (2011) [arXiv:1103.6029 [hep-th]].
- G.Goon, K.Hinterbichler and M.T., ``Symmetries for Galileons and DBI scalars on curved space," JCAP 1107, 017 (2011) [arXiv:1103.5745 [hep-th]].
- G.Goon, K.Hinterbichler and M.T., ``Stability and superluminality of spherical DBI galileon solutions," PRD 83, 085015 (2011) [arXiv:1008.4580 [hep-th]].
- M.Andrews, K.Hinterbichler, J.Khoury and M.T., ``Instabilities of Spherical Solutions with Multiple Galileons and $SO(N)$ Symmetry," PRD 83, 044042 (2011) [arXiv:1008.4128 [hep-th]].
- K.Hinterbichler, M.T. and D.Wesley, ``Multi-field galileons and higher co-dimension branes," PRD 82, 124018 (2010)[arXiv:1008.1305 [hep-th]].