NLO corrections to squark-squark production and decay at the LHC

in collaboration with W. Hollik and J. M. Lindert based on arXiv:1207.xxxx





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SUSY PARTICLES AT THE LHC

Production



Decay

Squark-Squark: NLO QCD: Beenakker et al. '96 NLO EW:Germer et al. '10 Beyond NLO (resummed): Beneke, Falgari, Schwinn '10 Beenakker et al '09

Squark: NLO QCD: Djouadi, Hollik, Junger '96 NLO EW:Guasch, Hollik, Sola '02

In order to study systematically the experimental signature emerging from SUSY particles, both corrections to production and decays must be included.

We study the experimental signature

 $2j + \not\!\!E_T(+X)$

via squark-squark production and direct decay into the lightest neutralino

 $pp \to \tilde{q}\tilde{q}' \to qq'\tilde{\chi}_1^0\tilde{\chi}_1^0(+X)$



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Standard procedure:

Producing events with a parton shower generator and rescaling the weights with the global K factor from NLO QCD corrections to the total cross-section of the squark-squark production. We study the experimental signature

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Our procedure:

Including fully differential NLO corrections to both the decay and production, where in the calculation all flavour and chirality configurations of intermediate squarks are treated independently.





NLO production

For every chirality and flavour configuration:



$$d\sigma_{pp\to\tilde{q}\tilde{q}'(+X)}^{(1)} = d\sigma_{pp\to\tilde{q}\tilde{q}'(g)}^{\text{virtual+soft}} + d\sigma_{pp\to\tilde{q}\tilde{q}'(g)}^{\text{coll}} + d\sigma_{pp\to\tilde{q}\tilde{q}'g}^{\text{hard}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{q}'\bar{q}'g}^{\text{real-quark}} + d\sigma_{pp\to\tilde{q}\tilde{$$

Fully differential cross-section

NLO production

For every chirality and flavour configuration:



Fully differential cross-section

NLO decay

$$d\Gamma^{(1)}_{\tilde{q}\to q\tilde{\chi}^0_1} = d\Gamma^{\text{virtual+soft}}_{\tilde{q}\to q\tilde{\chi}^0_1(g)} + d\Gamma^{\text{coll}}_{\tilde{q}\to q\tilde{\chi}^0_1(g)} + d\Gamma^{\text{hard}}_{\tilde{q}\to q\tilde{\chi}^0_1g}$$



Fully differential decay

NLO production

For every chirality and flavour configuration:



Fully differential cross-section

NLO decay

$$d\Gamma^{(1)}_{\tilde{q}\to q\tilde{\chi}^0_1} = d\Gamma^{\text{virtual+soft}}_{\tilde{q}\to q\tilde{\chi}^0_1(g)} + d\Gamma^{\text{coll}}_{\tilde{q}\to q\tilde{\chi}^0_1(g)} + d\Gamma^{\text{hard}}_{\tilde{q}\to q\tilde{\chi}^0_1g}$$



Fully differential decay

NLO total decay width

$$\Gamma_{\tilde{q}\to q'\tilde{\chi}_j^{\pm}}^{(0+1)} = \Gamma^{(0)} \left[1 + \frac{4}{3} \frac{\alpha_s}{\pi} F^{QCD} \left(\frac{m_{\tilde{\chi}_j^{\pm}}}{m_{\tilde{q}}}, \frac{m_{\tilde{q}}}{m_{\tilde{g}}} \right) \right]$$



COMBINATION

For all different combinations of light flavours and chiralities, weighted events for squark-squark production are produced in the LAB frame.

Weighted decay events are generated in the respective squark rest-frame.

boost of decay events + "master formula"

Fully differential distributions of factorizable NLO contributions in NWA.







NUMERICAL RESULTS

We cluster partons into jets with anti- k_T algorithm, R=0.4 ATLAS and R=0.5 CMS and we always select jets according to:

 $p_{j_{1/2}}^{\mathrm{T}} > 20 \text{ GeV} \quad |\eta_j| < 2.8,$ $p_{j_i}^{\mathrm{T}} > 50 \text{ GeV} \quad |\eta_j| < 3.0 \text{ (for CMS observables)}$

SPS1a (14 TeV)

Scale variation plots: $\mu_f = \mu_r = (m/2, m, 2m)$, m: average squark mass



CMSSM 10.1.5 (14 TeV)

qq

Comparison between NLO and LO rescaled by global K-factor: corrections purely in the shapes

| 5_ | | | | | | | | | | | | |
|----|-----------------------|-------------------|--------------------|--------------|-------------------|------------------------------------|-------------------|-----|----------------------------|-----|------|------|
| 5 | $^{1}_{0.8}10.1.5$ | $ig \ 	ilde u_L$ | $	ilde{u}_R$ | $	ilde{d}_L$ | $	ilde{d}_{R}$ | $\tilde{\boldsymbol{\Theta}}_{.8}$ | $	ilde{\chi}^0_1$ | | | | | |
| | $m \Theta ss (G = V)$ | 400437.760 | 0 138 82.03 | 1999.71 | 20 p 376.9 | 156896 | 2 9 903 | 400 | 600 p ^T [GeV | 800 | 1000 | 1200 |
| | | | | | | | - | | J | | | |

D 10-1



b/GeV]

NLO

CMSSM 10.1.5 (14 TeV)

Comparison between NLO and LO rescaled by global K-factor: corrections purely in the shapes



 $H_T = \sum_{i=1,2(,3)} p_i^{\mathrm{T}}$

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Effect on cut-and-count searches performed by ATLAS.

Signal Region:

 $p_{j_1}^{\mathrm{T}} > 130 \text{ GeV}, \ p_{j_2}^{\mathrm{T}} > 40 \text{ GeV}, \ |\eta_{j_{1/2}}| < 2.8, \ \Delta \phi(j_{1/2}, \not{\!\!E_T}) > 0.4$ $m_{\mathrm{eff}} > 1 \text{ TeV}, \ \not{\!\!E_T}/m_{\mathrm{eff}} > 0.3$

| benchmarkpoint | Energy [TeV] | $N_{ m ATLAS}^{(0)}$ | $N_{ m ATLAS}^{(0+1)}$ | $K_{N_{ m ATLAS}}$ | $K_{pp ightarrow	ilde{q}	ilde{q}'}$ |
|----------------|--------------|----------------------|------------------------|--------------------|-------------------------------------|
| | 7 | 0.066 pb | 0.083 pb | 1.26 | 1.37 |
| SPS1a | 8 | $0.097\mathrm{pb}$ | $0.121\mathrm{pb}$ | 1.25 | 1.35 |
| | 14 | $0.347\mathrm{pb}$ | $0.424\mathrm{pb}$ | 1.22 | 1.28 |
| | 7 | $0.313\mathrm{fb}$ | $0.503{ m fb}$ | 1.61 | 1.57 |
| 10.1.5 | 8 | $0.861{ m fb}$ | $1.344\mathrm{fb}$ | 1.56 | 1.52 |
| | 14 | $13.82{ m fb}$ | $19.77\mathrm{fb}$ | 1.43 | 1.40 |
| | 7 | $0.140\mathrm{fb}$ | $20.76\mathrm{fb}$ | ~ 150 | 1.40 |
| p19MSSM1 | 8 | $0.339{ m fb}$ | $37.96\mathrm{fb}$ | ~ 110 | 1.39 |
| | 14 | $0.0044\mathrm{pb}$ | $0.264\mathrm{pb}$ | ~ 60 | 1.34 |

CONCLUSION

We provide a consistent fully differential calculation of factorizable NLO QCD corrections in NWA for squark-squark production and decay.

NLO corrections are, in general, important for precise description of physical observables.

In particular cases they can be essential for realistic description.

Fully differential NLO QCD prediction of production and decay for all squark/gluino channels are desirable.

Thank you for the attention.

Effect on cut-and-count searches performed by CMS.

| benchmarkpoint | Energy [TeV] | $N_{ m CMS}^{(0)}$ | $N_{ m CMS}^{(0+1)}$ | $K_{N_{ m CMS}}$ | $ig K_{pp ightarrow 	ilde{q}	ilde{q}'}$ |
|----------------|--------------|-------------------------------|----------------------|---------------------|--|
| | 7 | 0.112 pb | 0.141 pb | 1.26 | 1.37 |
| SPS1a | 8 | $0.157\mathrm{pb}$ | $0.197\mathrm{pb}$ | 1.25 | 1.35 |
| | 14 | $0.488\mathrm{pb}$ | $0.614\mathrm{pb}$ | 1.26 | 1.28 |
| | 7 | 0.201 pb | $0.261\mathrm{pb}$ | 1.30 | 1.57 |
| 10.1.5 | 8 | $0.542\mathrm{fb}$ | $0.674{ m fb}$ | 1.24 | 1.52 |
| | 14 | $8.129\mathrm{fb}$ | $8.884\mathrm{fb}$ | 1.09 | 1.40 |
| | 7 | $10^{-6}{\rm pb}$ | $0.095\mathrm{pb}$ | $\mathcal{O}(10^4)$ | 1.40 |
| p19MSSM1 | 8 | $10^{-6}{\rm pb}$ | $0.151\mathrm{pb}$ | $\mathcal{O}(10^4)$ | 1.39 |
| | 14 | $2 \cdot 10^{-5} \mathrm{pb}$ | $0.687\mathrm{pb}$ | $\mathcal{O}(10^4)$ | 1.34 |

Signal Region:

 $p_{j_{1/2}}^{\mathrm{T}} > 100 \text{ GeV}, \ |\eta_{j_1}| < 2.5, \ |\eta_{j_2}| < 3.0,$ $H_T > 350 \text{ GeV}, \ H_T/E_T < 1.25, \ \alpha_T > 0.55.$

CMSSM 10.1.5 (14 TeV) Contributions of the different flavour and chirality configurations

| channel | $\sigma^{(0)}_{pp ightarrow 	ilde{q}	ilde{q}'}$ | $\sigma^{(0+1)}_{pp ightarrow 	ilde{q}	ilde{q}'}$ | $K_{pp ightarrow	ilde{q}	ilde{q}'}$ | $\sigma_{2j+E_T(+X)}^{(0)}$ | $\sigma^{(0+1)}_{2j+{ ot\!\!/} E_T(+X)}$ | $K_{2j+\not\!\!\!E_T(+X)}$ |
|---------------------------|---|---|-------------------------------------|-----------------------------|--|----------------------------|
| | [fb] | [fb] | | [fb] | [fb] | |
| $	ilde{u}_L 	ilde{u}_L$ | 7.08 | 9.44 | 1.33 | $1.22 \cdot 10^{-3}$ | $1.68 \cdot 10^{-3}$ | 1.38 |
| $\tilde{u}_R \tilde{u}_R$ | 8.64 | 11.5 | 1.33 | 8.25 | 11.36 | 1.38 |
| $	ilde{d}_L 	ilde{d}_L$ | 1.07 | 1.44 | 1.36 | $2.82 \cdot 10^{-4}$ | $3.96 \cdot 10^{-4}$ | 1.40 |
| $	ilde{d}_R 	ilde{d}_R$ | 1.39 | 1.88 | 1.35 | 1.33 | 1.84 | 1.39 |
| $\tilde{u}_L \tilde{u}_R$ | 6.00 | 8.49 | 1.42 | $7.78 \cdot 10^{-2}$ | $11.33 \cdot 10^{-2}$ | 1.45 |
| $	ilde{d}_L 	ilde{d}_R$ | $8.20 \cdot 10^{-1}$ | 1.19 | 1.45 | $1.32 \cdot 10^{-2}$ | $1.96 \cdot 10^{-5}$ | 1.49 |
| $	ilde{u}_L 	ilde{d}_L$ | 8.25 | 11.9 | 1.44 | $1.76 \cdot 10^{-3}$ | $2.62 \cdot 10^{-3}$ | 1.49 |
| $	ilde{u}_R 	ilde{d}_R$ | 10.5 | 15.1 | 1.44 | 10.00 | 14.92 | 1.49 |
| $	ilde{u}_L 	ilde{c}_L$ | $3.28 \cdot 10^{-1}$ | $4.33 \cdot 10^{-1}$ | 1.32 | $5.65 \cdot 10^{-5}$ | $7.73 \cdot 10^{-5}$ | 1.37 |
| $	ilde{u}_R 	ilde{c}_R$ | $4.29 \cdot 10^{-1}$ | $5.74 \cdot 10^{-1}$ | 1.34 | $4.09 \cdot 10^{-1}$ | $5.68 \cdot 10^{-1}$ | 1.39 |
| $	ilde{d}_L 	ilde{s}_L$ | $1.95 \cdot 10^{-1}$ | $2.75 \cdot 10^{-1}$ | 1.41 | $5.16 \cdot 10^{-5}$ | $7.5097 \cdot 10^{-5}$ | 1.46 |
| $	ilde{d}_R 	ilde{s}_R$ | $2.71 \cdot 10^{-1}$ | $3.87 \cdot 10^{-1}$ | 1.42 | $2.59 \cdot 10^{-1}$ | 3.82 | 1.48 |
| $	ilde{u}_L 	ilde{d}_R$ | 2.44 | 3.50 | 1.44 | $3.16 \cdot 10^{-2}$ | $4.67 \cdot 10^{-2}$ | 1.48 |
| $	ilde{u}_R 	ilde{d}_L$ | 2.40 | 3.46 | 1.44 | $3.87 \cdot 10^{-2}$ | $5.70 \cdot 10^{-2}$ | 1.48 |
| $	ilde{u}_L	ilde{c}_R$ | $1.69 \cdot 10^{-1}$ | $2.39 \cdot 10^{-1}$ | 1.41 | $2.19 \cdot 10^{-3}$ | $3.18 \cdot 10^{-3}$ | 1.46 |
| $	ilde{d}_L 	ilde{s}_R$ | $9.51 \cdot 10^{-2}$ | $1.39 \cdot 10^{-1}$ | 1.46 | $1.52 \cdot 10^{-3}$ | $2.29 \cdot 10^{-3}$ | 1.50 |
| sum | 50.04 | 69.86 | 1.40 | 20.41 | 29.32 | 1.44 |

p19MSSM1 (14 TeV)

do/o

Comparison between NLO and LO rescaled by global K-factor: corrections purely in the shapes



 $H_T = \sum_{i=1,2(,3)} p_i^{\mathrm{T}}$

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