

1. Motivation

We consider in general processes at the LHC with $\sqrt{s} = 8$ TeV. Generic Features

- High center of mass energy
- $\alpha_{\text{ew}} \ln^2[M_{W/Z}^2/s]$ potentially large

⇒ Electroweak corrections become important

Interesting processes are e.g.

1. $W/Z + n$ jet
2. Gauge boson pair production

Especially the first one is important

- Testing environment for pQCD calculations
- New physics searches
- Huge and irreducible SM background

⇒ Need a thorough understanding

A lot of effort has been put in QCD calculations. Electroweak corrections are as important as some of the known QCD corrections.

- High complexity ⇒ Monte Carlo generators
- ⇒ Need a framework to automatically provide NLL corrections to various processes

2. Setup of the Theory

1. Match full SM $SU(3)_C \otimes SU(2)_W \otimes U(1)_Y$ onto soft collinear effective theory, still with the unbroken SM

2. Run down to scale of EW symmetry breaking

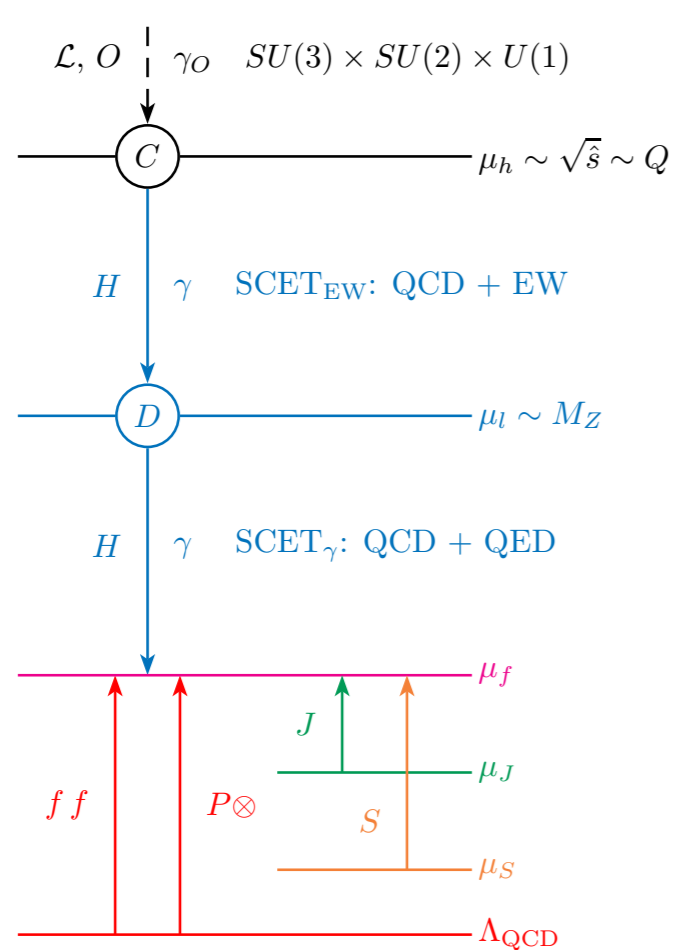


Figure 1: Schematic overview of the EFT

1. Fields are described by collinear and soft Wilson lines ⇒ valid for energetic particles

2. Regulator choice:

⇒ Factorizes soft and collinear calculations

Ingredients for the Calculation

1. Collinear corrections

- Depends on field, only
- Matching accounts for symmetry breaking effects, e.g. $\gamma - Z$ mixing, $M_W \neq M_Z$

2. Soft corrections

- Depends on gauge structure, only
- Matching accounts for blow up of operator basis ⇒ Electroweak symmetry breaking

Master Formula

$$\mathcal{C}_{\text{low}} = \mathcal{M}_{\text{low}} \mathcal{M}_{\text{col}} \mathcal{M}_{\text{soft}} \mathcal{C}_{\text{high}} \mathcal{M}_{SU(3)}^T$$

with the individual matrices below (Special case $SU(3)$: Separate gauge space, only running)

For the electroweak matrices we have, where R is a running and M a matching matrix

$$\mathcal{M}_{\text{low}} = \exp \left[\int_{\mu}^{M_Z} \frac{d\mu}{\mu} R_{U(1)_{\text{EM}}}^{\text{soft+col}} \right]$$

$$\mathcal{M}_{\text{col}} = \exp \left[M_{U(1)_Y \otimes SU(2)_W}^{\text{col}} \right]$$

$$\times \exp \left[\int_{M_Z}^{\sqrt{s}} \frac{d\mu}{\mu} R_{U(1)_Y \otimes SU(2)_W}^{\text{col}} \right]$$

$$\mathcal{M}_{\text{soft}} = \exp \left[\int_{M_Z}^{\sqrt{s}} \frac{d\mu}{\mu} R_{U(1)_Y}^{\text{soft}} \right] M_{SU(2)}^{\text{break}}$$

$$\times \exp \left[\int_{M_Z}^{\sqrt{s}} \frac{d\mu}{\mu} R_{SU(2)_W}^{\text{soft}} \right]$$

⇒ Wilson Coefficient

– Columns: Different $SU(3)_C$ Op.

– Rows: Different $SU(2)_W \otimes U(1)_Y$ Op.

Operators

- High Scale

$$\mathcal{O} = \bar{Q}_L \tau^a T^A Q_L W^a G^A$$

- Low scale

$$\mathcal{O}_1 = \bar{U}_L T^a D_L W^- G^a$$

$$\mathcal{O}_2 = \bar{D}_L T^a U_L W^+ G^a$$

$$\mathcal{O}_{3-6} = \text{Involving } Z \text{ and } \gamma$$

⇒ Non-square matrices

Parton Distribution Functions (PDF)

- Partons inside proton collide

⇒ Match onto non-perturbative set of PDFs

- We use CTEQ5 for the plots

Phase-Space

- Two additional variables x_A and x_B

⇒ Accounts for momentum distribution of partons within the protons

- Use special parameterisation: Factorization of single final state plus extra radiation
- Bar components: Defined in 1 particle PS

$$d\phi_2 = d\bar{\phi}_1 d\phi_{\text{rad}}$$

$$d\bar{\Phi}_1 = d\bar{x}_A d\bar{x}_B d\bar{k}_V (2\pi)^4 \delta^4(\bar{k}_A + \bar{k}_B - \bar{k}_V) = \frac{2\pi}{S} d\bar{Y}$$

$$d\Phi_{\text{rad}} = \frac{s}{(4\pi)^3} \frac{\xi}{1-\xi} \theta(\xi - \xi_{\text{max}}) d\xi d\gamma d\phi$$

1. $|\bar{Y}| \leq \log \frac{\sqrt{s}}{M}$: Rapidity of vector boson
2. $0 \leq \xi \leq \xi_{\text{max}}$: Energy fraction of the jet
3. $y = \cos \theta \in [-1, 1]$ of radiated parton
4. $\phi \in [0, 2\pi]$ is the azimuthal angle

3. Observables

$pp \rightarrow W^\pm j$

- In total 16 subprocesses contribute
- $W^+ \leftrightarrow W^-$: up and down quarks change role
- Taking into account the first and second generation as well as CKM factors
 1. $\bar{q}_u q_d \rightarrow W^+ G$ and $1 \leftrightarrow 2$
 2. $q_u G \rightarrow W^+ q_d$ and $1 \leftrightarrow 2$
 3. $\bar{q}_d G \rightarrow W^+ \bar{q}_u$ and $1 \leftrightarrow 2$

- Investigate p_T spectrum of gauge boson

⇒ Effectively we resum $\alpha_{\text{ew}} \ln^2[M_{W/Z}^2/p_{T,W/Z}^2]$

Plots (Preliminary!)

- Statistical error only (4M events)
- p_T binning: 25 GeV
- Corrections are colored as
 1. Black: Tree
 2. Red[Light]: LL[with QCD corrections]
 3. Blue[Light]: NLL[with QCD corrections]

- $pp \rightarrow W^- j$

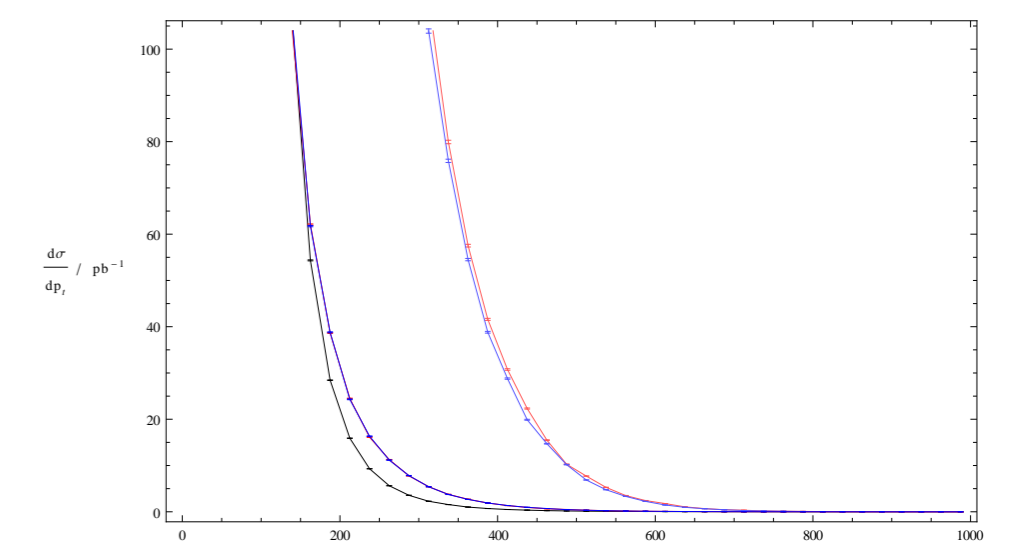


Figure 2: No real Gluon emission considered!

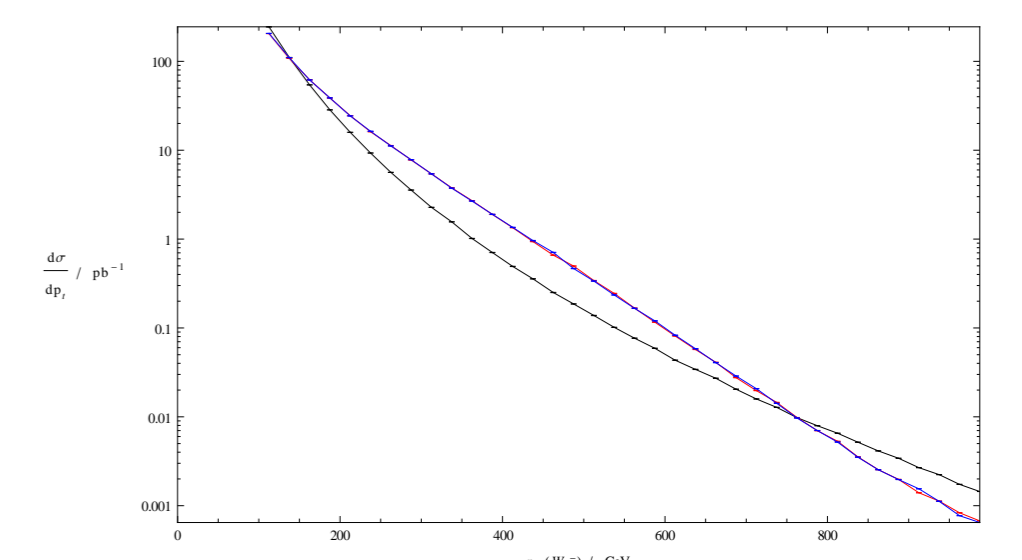


Figure 3: No real infrared divergence

- $pp \rightarrow W^+ j$

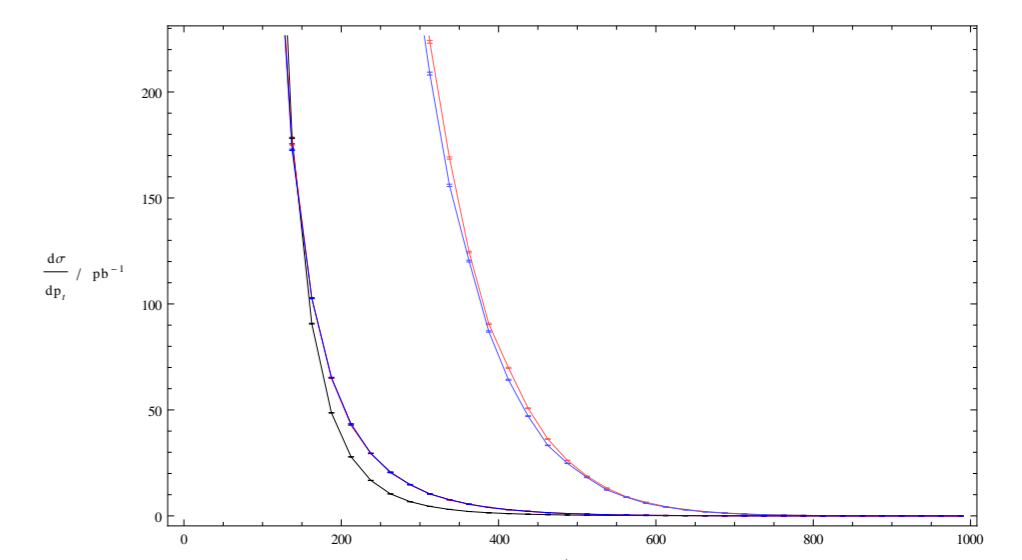


Figure 4: No real Gluon emission considered!

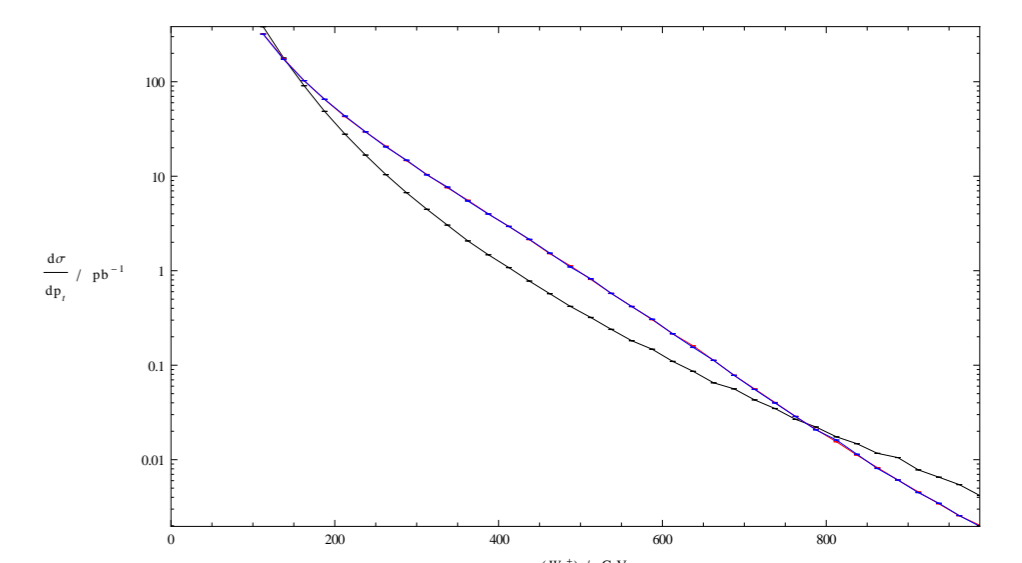


Figure 5: No real infrared divergence

- Combined

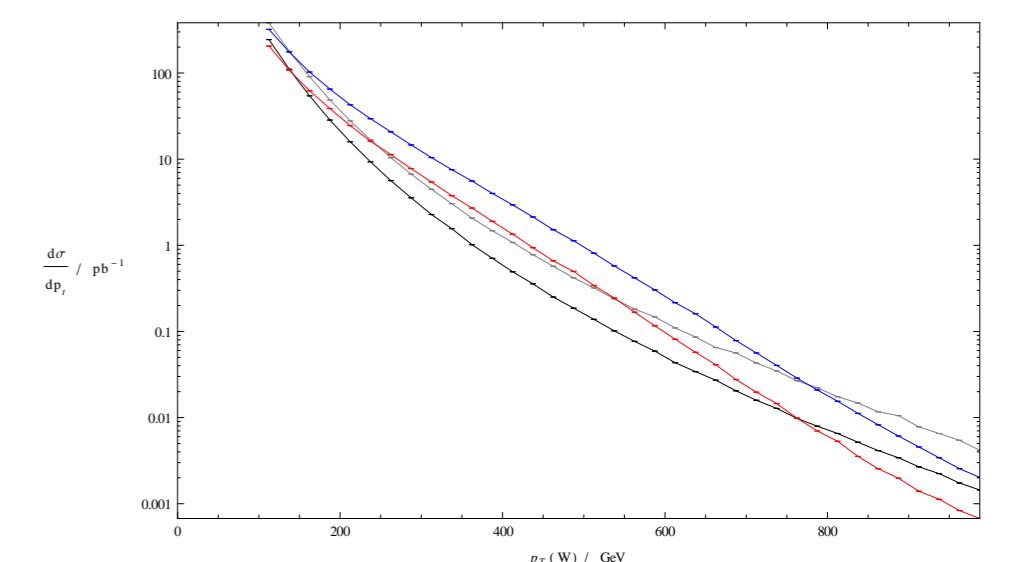


Figure 6: Black[Gray]: Tree $W^- [W^+]$, Red[Blue]: NLL w/o QCD $W^- [W^+]$

References

- [1] J. -y. Chiu, F. Golf, R. Kelley and A. V. Manohar, Phys. Rev. Lett. **100** (2008) 021802 [arXiv:0709.2377 [hep-ph]].
- [2] J. -y. Chiu, F. Golf, R. Kelley and A. V. Manohar, Phys. Rev. D **77** (2008) 053004 [arXiv:0712.0396 [hep-ph]].
- [3] J. -y. Chiu, A. Fuhrer, R. Kelley and A. V. Manohar, Phys. Rev. D **81** (2010) 014023 [arXiv:0909.0947 [hep-ph]].