

# $J/\psi$ and $\Upsilon$ production in proton-nucleus collisions: lessons from RHIC for the 2012 proton-lead LHC run

**Jean-Philippe Lansberg**

IPN Orsay, Université Paris-Sud

July 6, 2012 – **ICHEP 2012** – Melbourne, Australia



on behalf of E.G. Ferreiro, F. Fleuret, N. Matagne and A. Rakotozafindrabe

# Part I

## Background

# Quarkonium production cross section in pA collisions

We shall **assume** here that

- collinear factorisation applies: use of **nPDF**  
(for instance, saturation, if any, to be mimicked by gluon shadowing)
- pA collisions are described by a **probabilistic Glauber** approach
- other nuclear effects in an **effective absorption x-section**  
(should be compatible with the bound-state size, though)

# Quarkonium production cross section in $pA$ collisions

We shall **assume** here that

- collinear factorisation applies: use of **nPDF**  
(for instance, saturation, if any, to be mimicked by gluon shadowing)
- $pA$  collisions are described by a **probabilistic Glauber** approach
- other nuclear effects in an **effective absorption x-section**  
(should be compatible with the bound-state size, though)

What we want to **learn**:

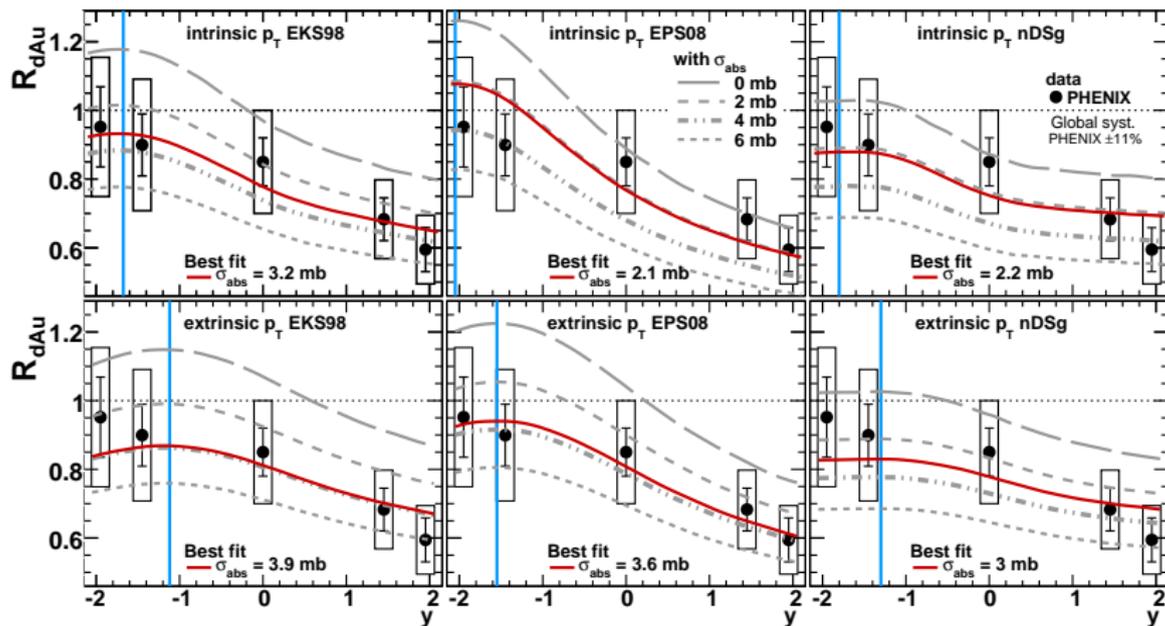
- whether all the nuclear effects can be encompassed in nPDFs  
(+ maybe an (effective) x-section for nuclear absorption)
- how large the nuclear effects –likely to affect  $AA$  production– are  
**baseline to be subtracted**
- genuine new information on QCD

## Part II

### $J/\psi$ production in $(p, d)A$ and $AA$ collisions

# Nuclear modification factor for $J/\psi$ in $dAu$ collisions at RHIC

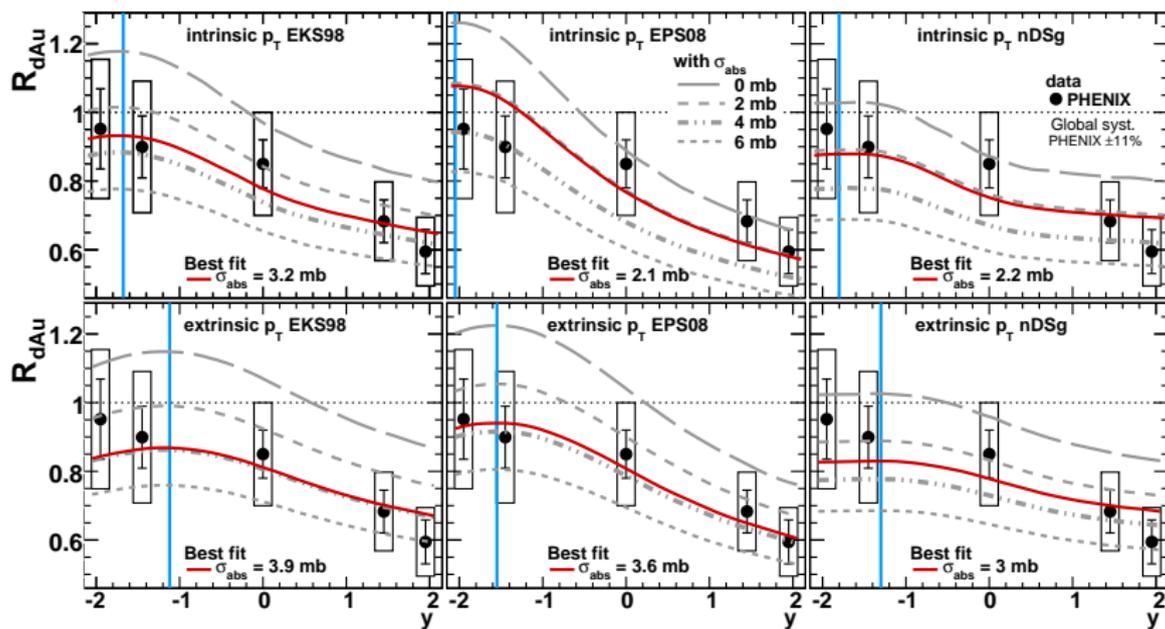
E.G. Ferreiro, F. Fleuret, J.P.L., A. Rakotozafindrabe, PLB 680:50,2009, PRC 81:064911, 2010; PHENIX PRC 77: 024912, 2008



- The shadowing impact does depend on the kinematics:

# Nuclear modification factor for $J/\psi$ in $dAu$ collisions at RHIC

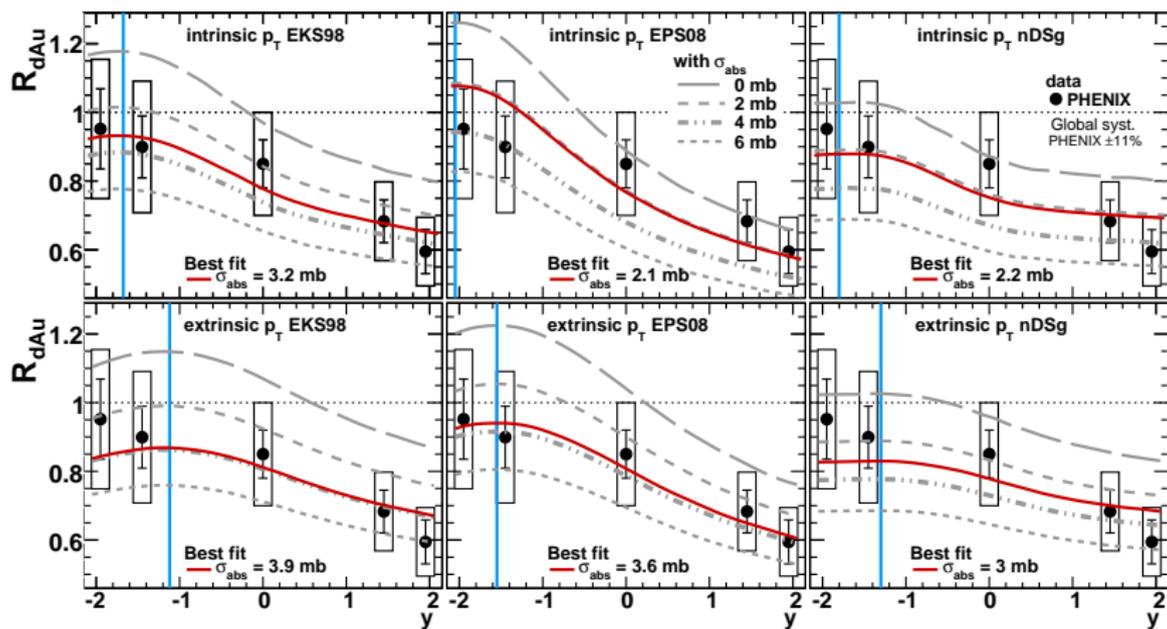
E.G. Ferreiro, F. Fleuret, J.P.L., A. Rakotozafindrabe, PLB 680:50,2009, PRC 81:064911, 2010; PHENIX PRC 77: 024912, 2008



- The shadowing impact does depend on the kinematics:  $2 \rightarrow 1$  vs  $2 \rightarrow 2$   
intrinsic extrinsic

# Nuclear modification factor for $J/\psi$ in $dAu$ collisions at RHIC

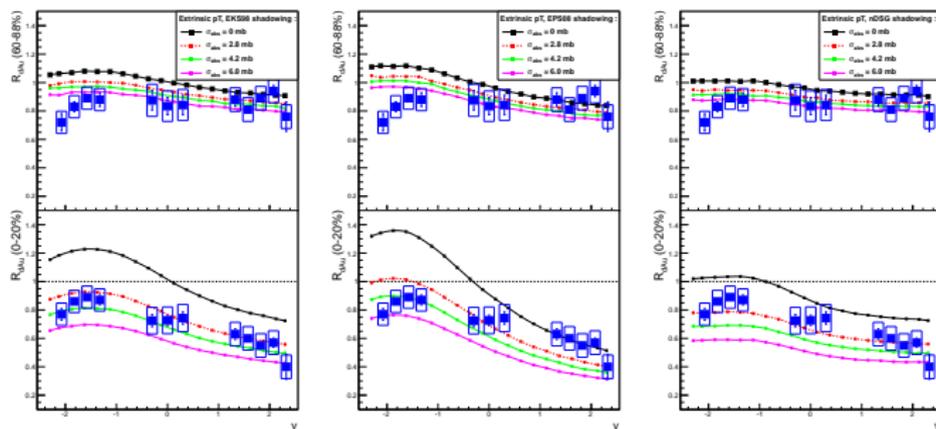
E.G. Ferreiro, F. Fleuret, J.P.L., A. Rakotozafindrabe, PLB 680:50,2009, PRC 81:064911, 2010; PHENIX PRC 77: 024912, 2008



- The shadowing impact does depend on the kinematics:  $2 \rightarrow 1$  vs  $2 \rightarrow 2$   
intrinsic extrinsic
- Shift of the rapidity distribution (see the vertical blue line)

# Comparison with the most recent PHENIX data

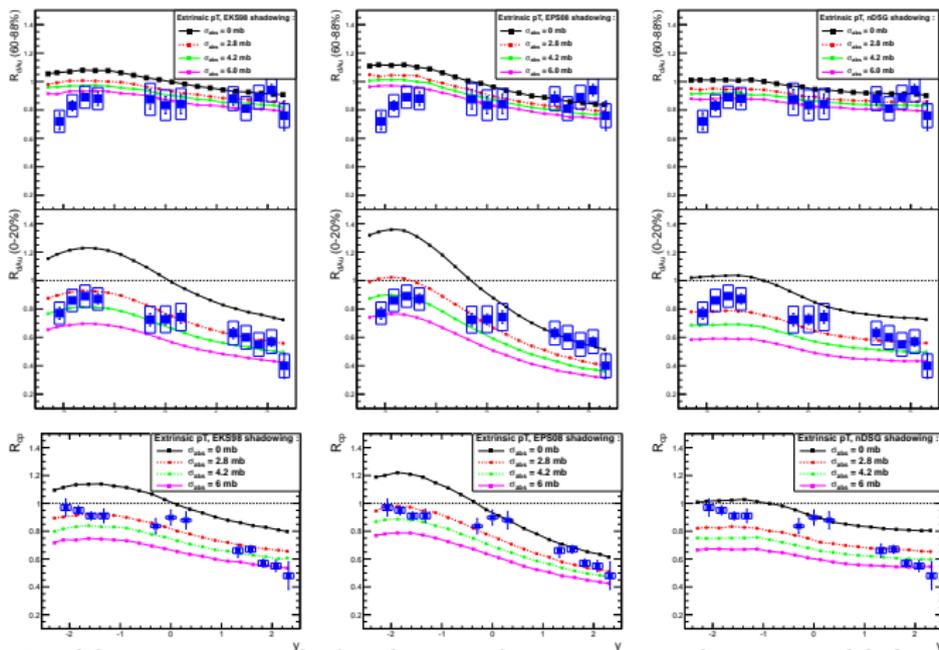
E.G. Ferreiro, F. Fleuret, J.P.L., N. Mataane A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX PRL 107: 142301, 2011



- EKS98 with  $\sigma_{abs} \simeq 3\text{mb}$  (red curve) seems to do a good job

# Comparison with the most recent PHENIX data

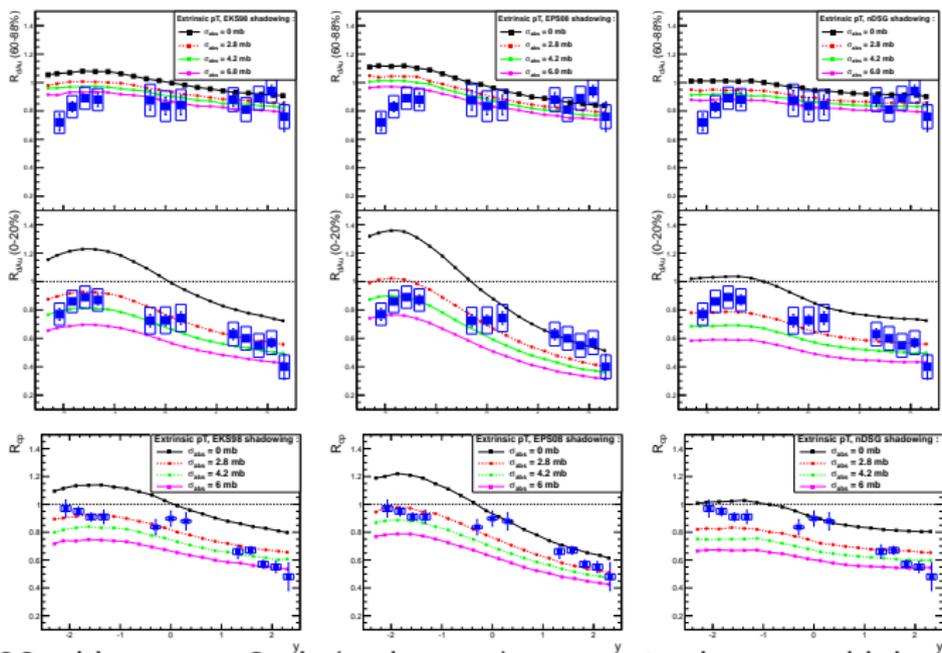
E.G. Ferreiro, F. Fleuret, J.P.L., N. Mataane A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX PRL 107: 142301, 2011



- EKS98 with  $\sigma_{abs} \simeq 3$ mb (red curve) seems to do a good job
- Less true when one looks at  $R_{CP}$  (EPS08 (i.e. strong shadowing) better)

# Comparison with the most recent PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Mataane A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX PRL 107: 142301, 2011

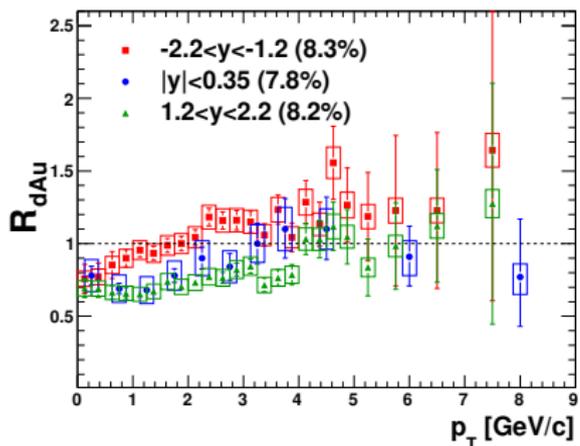


- EKS98 with  $\sigma_{abs} \simeq 3\text{mb}$  (red curve) seems to do a good job
- Less true when one looks at  $R_{CP}$  (EPS08 (i.e. strong shadowing) better)
- Lesson for the LHC:  $R_{CP}$  can be quite instructive

◀ even when one has  $R_{pA}$

# $R_{dAu}$ vs $P_T$ : new PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX: 1204.0777



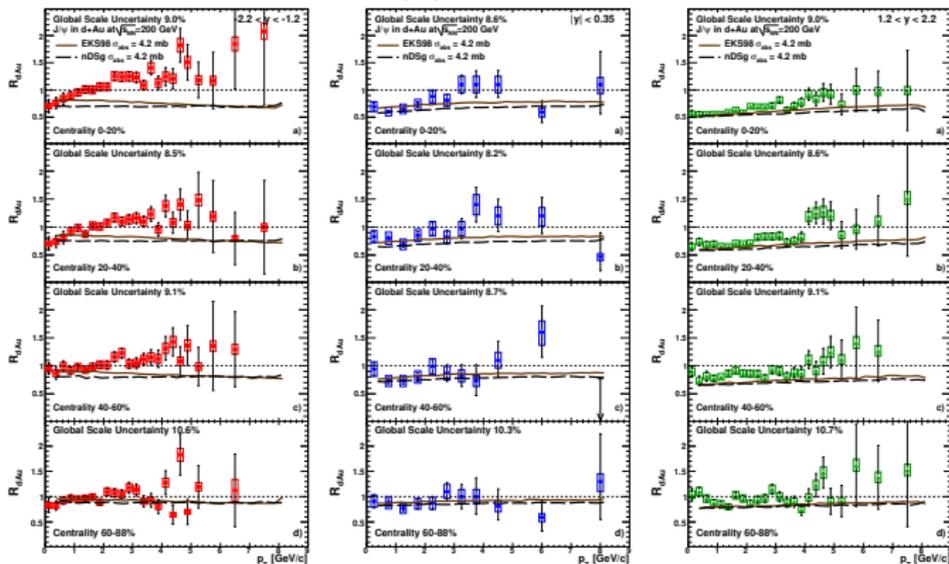
# $R_{dAu}$ vs $P_T$ : new PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX: 1204.0777

$-2.2 < y < -1.2$

$|y| < 0.35$

$1.2 < y < 2.2$



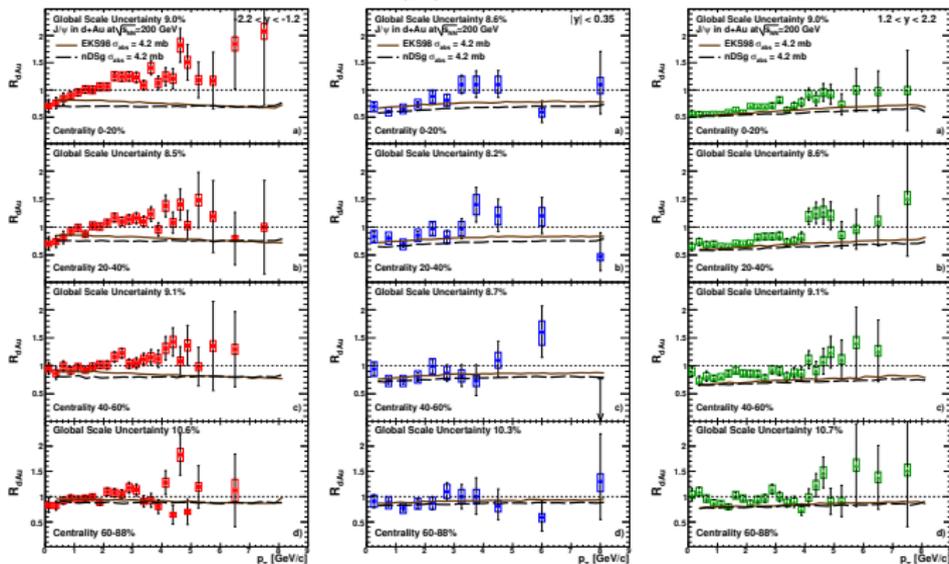
# $R_{dAu}$ vs $P_T$ : new PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX: 1204.0777

$-2.2 < y < -1.2$

$|y| < 0.35$

$1.2 < y < 2.2$



● Note that our curves do not include any specific Cronin effect

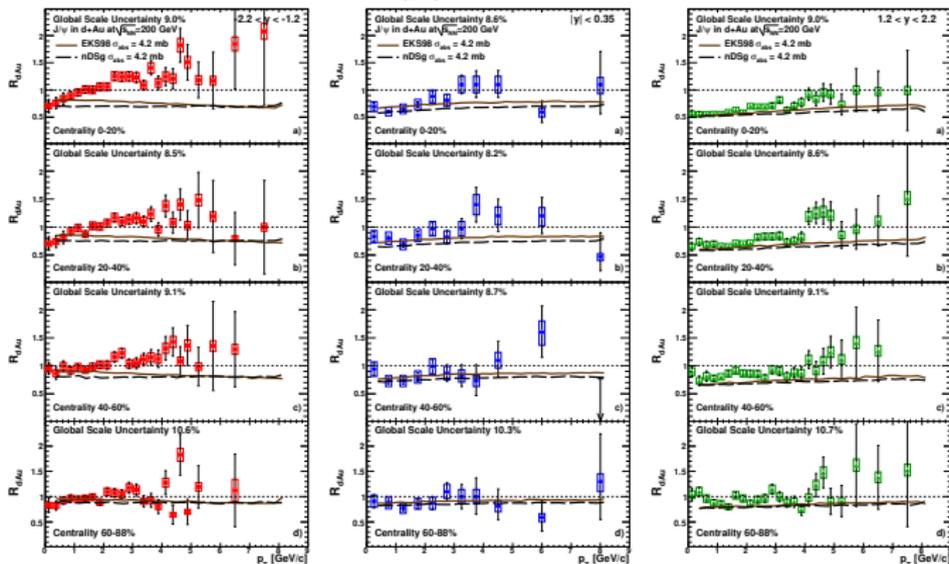
# $R_{dAu}$ vs $P_T$ : new PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX: 1204.0777

$-2.2 < y < -1.2$

$|y| < 0.35$

$1.2 < y < 2.2$



- Note that our curves do not include any specific Cronin effect
- Unexpected –and interesting– behaviour in the backward region

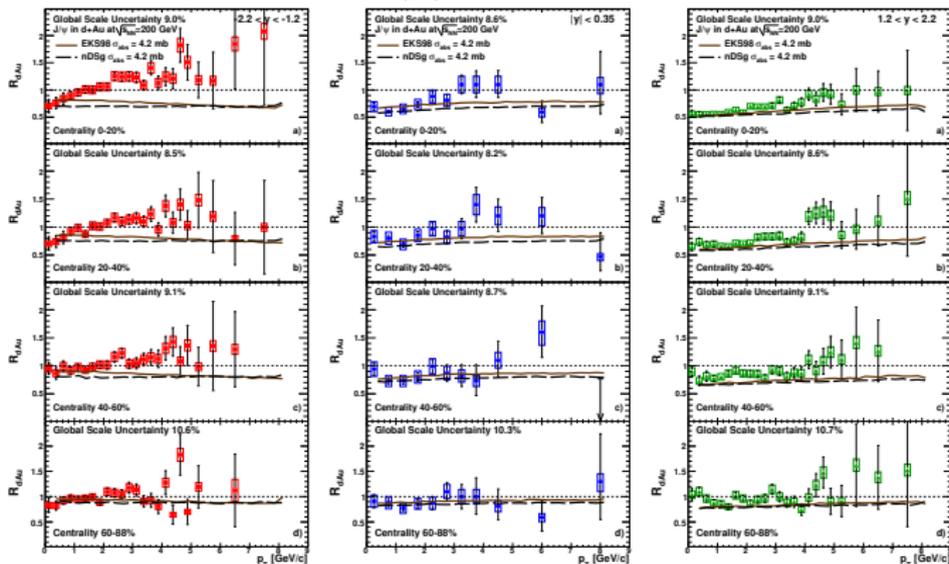
# $R_{dAu}$ vs $P_T$ : new PHENIX data

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne A. Rakotozafindrabe, FBS 53: 27, 2012; PHENIX: 1204.0777

$-2.2 < y < -1.2$

$|y| < 0.35$

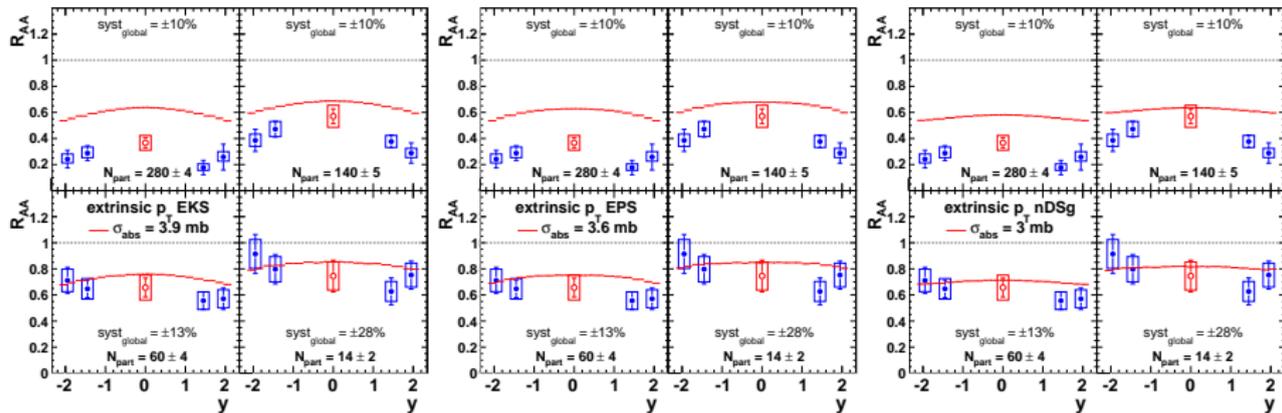
$1.2 < y < 2.2$



- Note that our curves do not include any specific Cronin effect
- Unexpected –and interesting– behaviour in the backward region
- Cronin effect stronger in the backward region ?
- Unexpected nPDF effect ? It would be nice to see  $[-2.2:-1.7]$  &  $[-1.7:-1.2]$

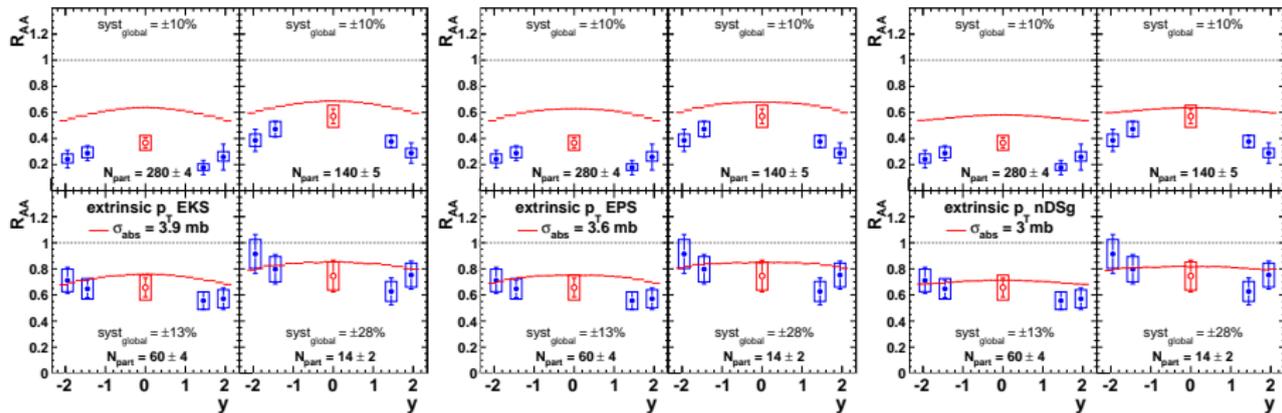
# Rapidity dependence in AuAu

E.G. Ferreiro, F. Fleuret, JPL, A. Rakotozafindrabe, PLB 680:50,2009, PRC 81, 064911 (2010); PHENIX: PRL 98: 232301,2007



# Rapidity dependence in AuAu

E.G. Ferreiro, F. Fleuret, JPL, A. Rakotozafindrabe, PLB 680:50,2009, PRC 81, 064911 (2010); PHENIX: PRL 98: 232301,2007

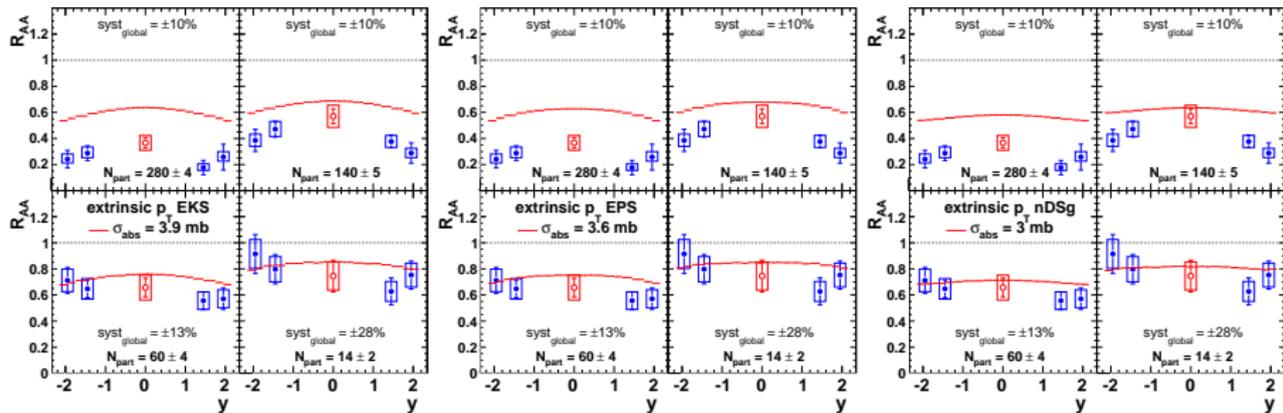


- The  $J/\psi$  suppression from shadowing is  $y$ -dependent

Reminder:  $\sigma_{\text{abs}}$  taken as constant

# Rapidity dependence in AuAu

E.G. Ferreiro, F. Fleuret, JPL, A. Rakotozafindrabe, PLB 680:50,2009, PRC 81, 064911 (2010); PHENIX: PRL 98: 232301,2007



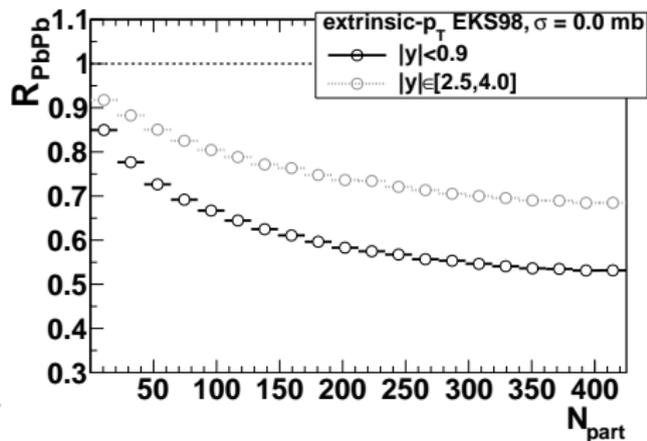
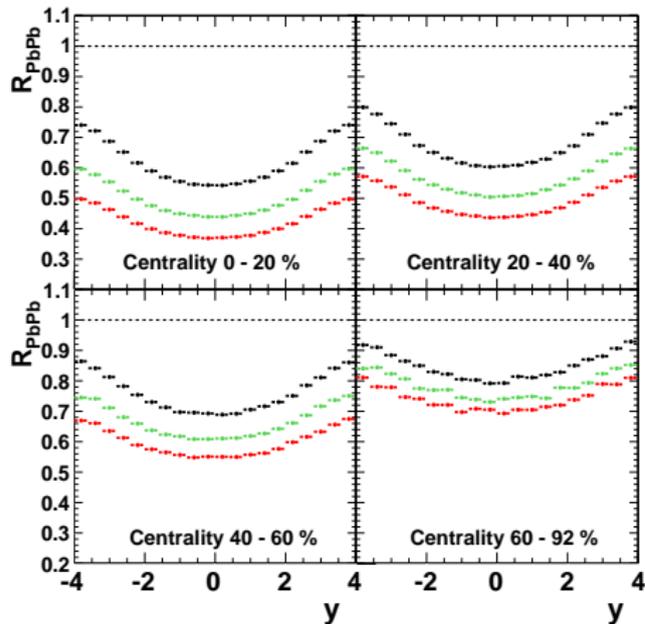
- The  $J/\psi$  suppression from shadowing is  $y$ -dependent  
Reminder:  $\sigma_{abs}$  taken as constant
- Can account for a part of the  $y$ -dependence of the most central data

# Rapidity dependence in PbPb collisions at $\sqrt{s_{NN}} = 2.76$ TeV

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330

# Rapidity dependence in PbPb collisions at $\sqrt{s_{NN}} = 2.76$ TeV

E.G. Ferreira, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330



In black,  $\sigma_{abs}^{eff} = 0.0$  mb, in green,  $\sigma_{abs}^{eff} = 1.5$  mb, in red,  $\sigma_{abs}^{eff} = 2.8$  mb

Strong rapidity dependence of  $R_{PbPb}$ !

Opposite to that at RHIC energy!

## CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV

Let us compare

- 2 nPDF sets: EKS98 (blue) and nDSg (red)

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV

Let us compare

- 2 nPDF sets: EKS98 (blue) and nDSg (red)
- two production models different at low  $P_T$  : one to be combined with  $k_T$  smearing (model 1,  $\sim$  CEM NLO), the other not (model 2)

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV

Let us compare

- 2 nPDF sets: EKS98 (blue) and nDSg (red)
- two production models different at low  $P_T$  : one to be combined with  $k_T$  smearing (model 1,  $\sim$  CEM NLO), the other not (model 2)
- taking into account the uncertainty on the scale  $Q \leftrightarrow \mu_F$ :  
width of the bands

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV

Let us compare

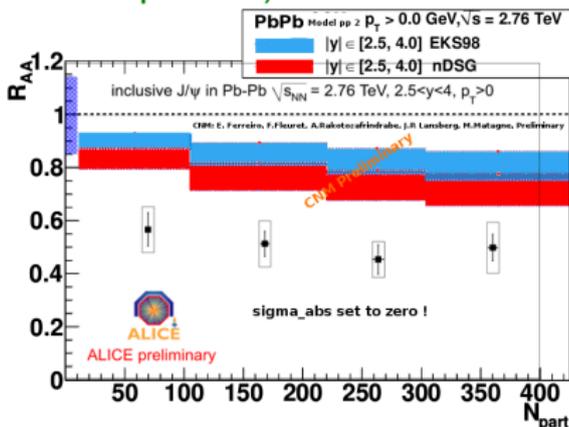
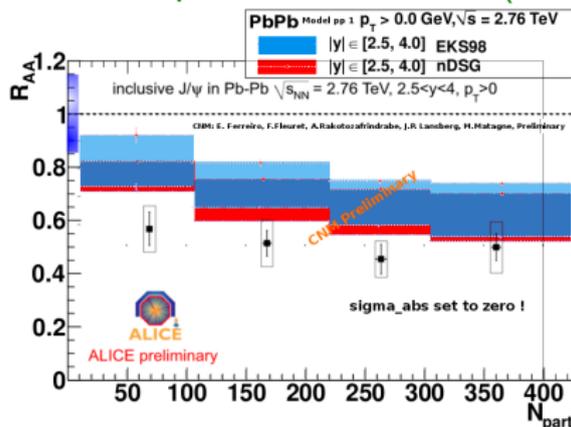
- 2 nPDF sets: EKS98 (blue) and nDSg (red)
- two production models different at low  $P_T$  : one to be combined with  $k_T$  smearing (model 1,  $\sim$  CEM NLO), the other not (model 2)
- taking into account the uncertainty on the scale  $Q \leftrightarrow \mu_F$ :  
width of the bands
- no  $\sigma_{abs}^{eff}$  for now (is it really negligible at the LHC ?)

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV

Let us compare

- 2 nPDF sets: EKS98 (blue) and nDSg (red)
- two production models different at low  $P_T$ : one to be combined with  $k_T$  smearing (model 1,  $\sim$  CEM NLO), the other not (model 2)
- taking into account the uncertainty on the scale  $Q \leftrightarrow \mu_F$ :  
width of the bands
- no  $\sigma_{abs}^{eff}$  for now (is it really negligible at the LHC ?)

Without  $P_T$  cut and forward (ALICE acceptance):



# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV with a $P_T$ cut

With  $P_T > 6.5$  GeV cut and mostly central (CMS/ATLAS acceptance)

CMS, JHEP 1205 (2012) 063

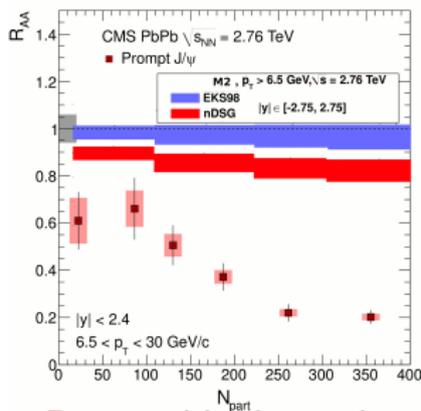
---

<sup>1</sup>Note: we consider the shadowing with M1 to be exaggerated:  $Q^2$  too small, too much weight on small  $P_T$ , i.e. small  $x$  

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV with a $P_T$ cut

With  $P_T > 6.5$  GeV cut and mostly central (CMS/ATLAS acceptance)

CMS, JHEP 1205 (2012) 063



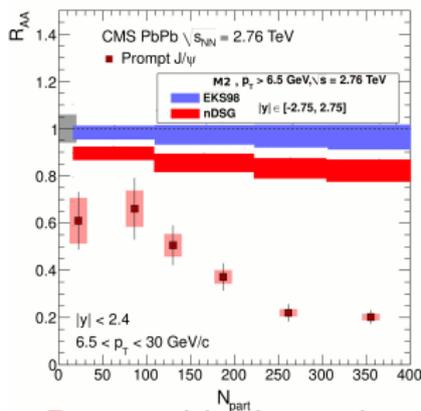
- Non trivial effect of the  $P_T$  cut: this depends on the nPDF set

<sup>1</sup>Note: we consider the shadowing with M1 to be exaggerated:  $Q^2$  too small, too much weight on small  $P_T$ , i.e. small  $x$

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV with a $P_T$ cut

With  $P_T > 6.5$  GeV cut and mostly central (CMS/ATLAS acceptance)

CMS, JHEP 1205 (2012) 063



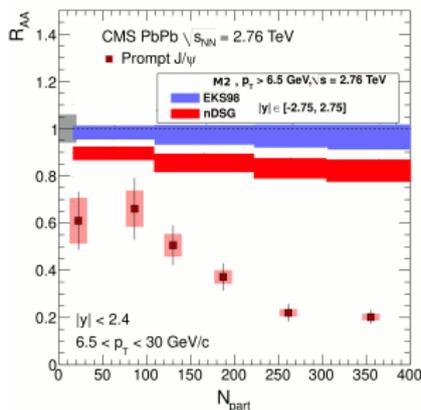
- **Non trivial effect of the  $P_T$  cut:** this depends on the nPDF set
- Without  $P_T$  cut with M1<sup>1</sup>, EKS98  $\sim$  nDSg and strong suppression

<sup>1</sup>Note: we consider the shadowing with M1 to be exaggerated:  $Q^2$  too small, too much weight on small  $P_T$ , i.e. small  $x$

# CNM vs data at $\sqrt{s_{NN}} = 2.76$ TeV with a $P_T$ cut

With  $P_T > 6.5$  GeV cut and mostly central (CMS/ATLAS acceptance)

CMS, JHEP 1205 (2012) 063



- **Non trivial effect of the  $P_T$  cut:** this depends on the nPDF set
- Without  $P_T$  cut with M1<sup>1</sup>, EKS98  $\sim$  nDSg and strong suppression
- With a  $P_T$  cut with M2, EKS98 (slight/no suppression)  $\approx$  nDSg (still significant suppression)

<sup>1</sup>Note: we consider the shadowing with M1 to be exaggerated:  $Q^2$  too small, too much weight on small  $P_T$ , i.e. small  $x$

## Part III

# $\Upsilon$ production in $p/dA$ and $AA$ collisions

# Absorption: $\Upsilon$ vs. $J/\psi$

$\sigma_{abs}^{\Upsilon}$  should be small

Absorption:  $\Upsilon$  vs.  $J/\psi$ 

- at  $y > 0$ ,  $t_f = \gamma \times 0.4\text{fm} \gg r_{\text{Au}}$ :  $\sigma_{abs}^Y$  should be small  
pre-resonant state exiting the nucleus

$$\sigma_{abs}^{b\bar{b}} \simeq 0.1 \sigma_{abs}^{c\bar{c}}$$

Absorption:  $\Upsilon$  vs.  $J/\psi$  $\sigma_{abs}^Y$  should be small

- at  $y > 0$ ,  $t_f = \gamma \times 0.4\text{fm} \gg r_{\text{Au}}$ : pre-resonant state exiting the nucleus

$$\sigma_{abs}^{b\bar{b}} \simeq 0.1 \sigma_{abs}^{c\bar{c}}$$

- at  $y < 0$ ,  $t_f \leq r_{\text{Au}}$ , fully formed in the nucleus: e.g.  $\sigma_{abs}^{Y(2S)} \geq 4\sigma_{abs}^{Y(1S)}$   
Yet, equal suppression found by E772 in the backward region:

$$\Rightarrow \sigma_{abs}^{Y(2S)} - \sigma_{abs}^{Y(1S)} \text{ small} \Rightarrow \sigma_{abs}^{Y(1S)} \text{ small}$$

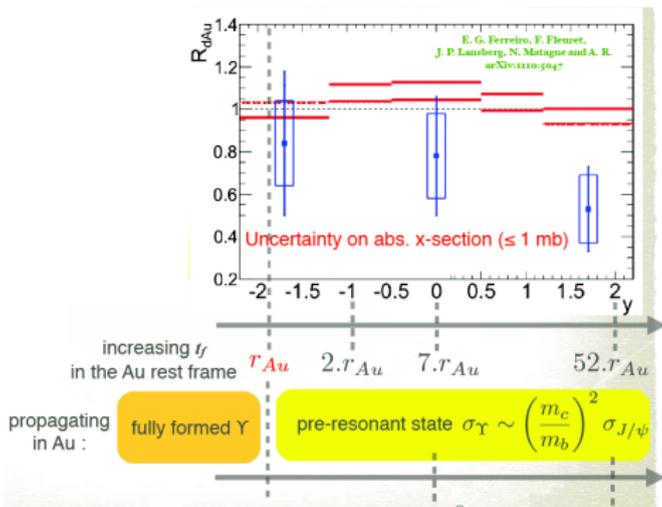
Absorption:  $\Upsilon$  vs.  $J/\psi$  $\sigma_{abs}^Y$  should be small

- at  $y > 0$ ,  $t_f = \gamma \times 0.4\text{fm} \gg r_{Au}$ : pre-resonant state exiting the nucleus

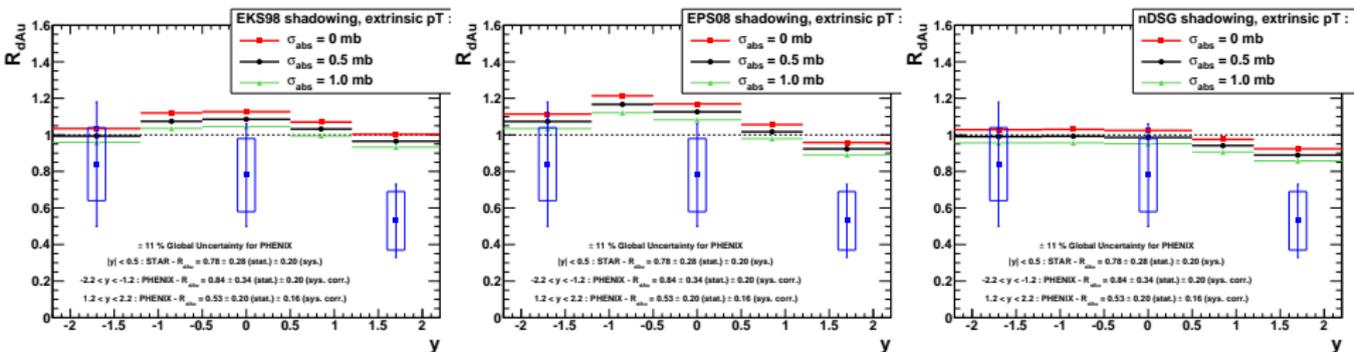
$$\sigma_{abs}^{b\bar{b}} \simeq 0.1 \sigma_{abs}^{c\bar{c}}$$

- at  $y < 0$ ,  $t_f \leq r_{Au}$ , fully formed in the nucleus: e.g.  $\sigma_{abs}^{Y(2S)} \geq 4\sigma_{abs}^{Y(1S)}$   
Yet, equal suppression found by E772 in the backward region:

$$\Rightarrow \sigma_{abs}^{Y(2S)} - \sigma_{abs}^{Y(1S)} \text{ small} \Rightarrow \sigma_{abs}^{Y(1S)} \text{ small}$$

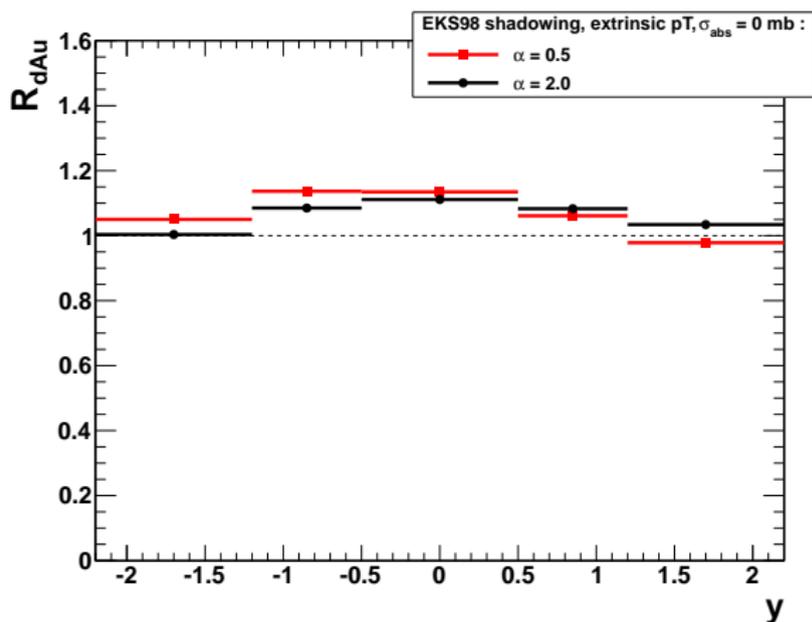


## nPDF uncertainty on Y production in dAu collisions



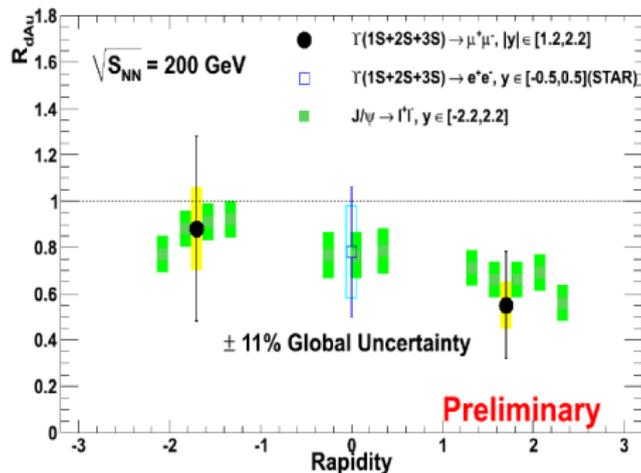
- Similar trend for the three nPDFs (EKS98, EPS08, nDSG) (nDSG shadowing as strong as EKS since its  $Q^2$  evolution is slow)
- Uncertainty from the nPDFs larger than our conservative estimate for that from the absorption x-section (0-1 mb)
- A priori, uncertainties smaller than for  $J/\psi$

## (Factorisation) scale dependence of Y production in dAu

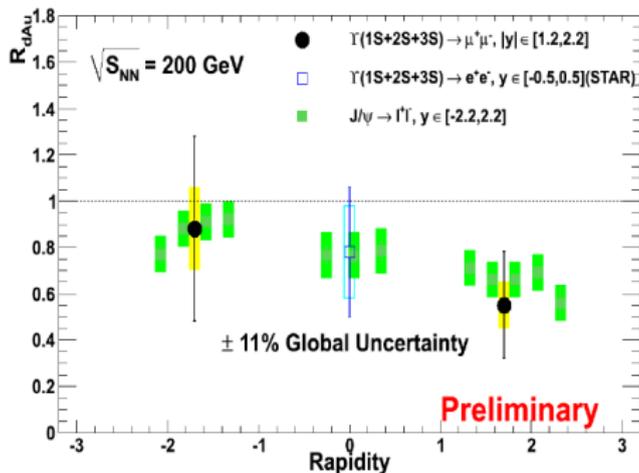


$$Q = \mu_F = \alpha \times m_T$$

→ Scale uncertainty for Y is smaller than for  $J/\psi$

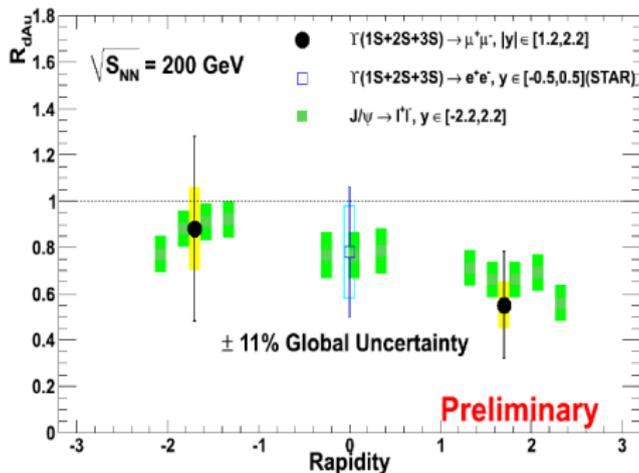
Y vs  $J/\psi$ : different effects, same suppression ?

T. Frawley, Hard Probes 2012

Y vs  $J/\psi$ : different effects, same suppression ?

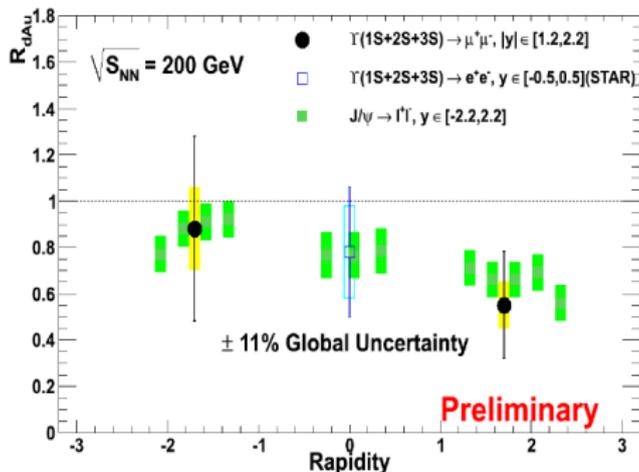
T. Frawley, Hard Probes 2012

- A priori same behaviour, but

Y vs  $J/\psi$ : different effects, same suppression ?

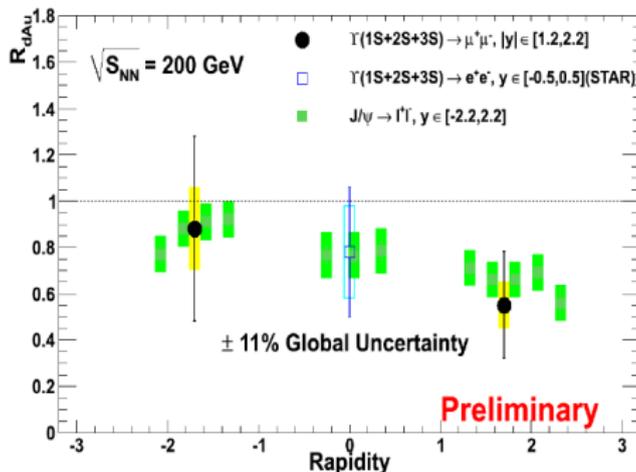
T. Frawley, Hard Probes 2012

- A priori same behaviour, but
- **Absorption** should be definitely **smaller**

Y vs  $J/\psi$ : different effects, same suppression ?

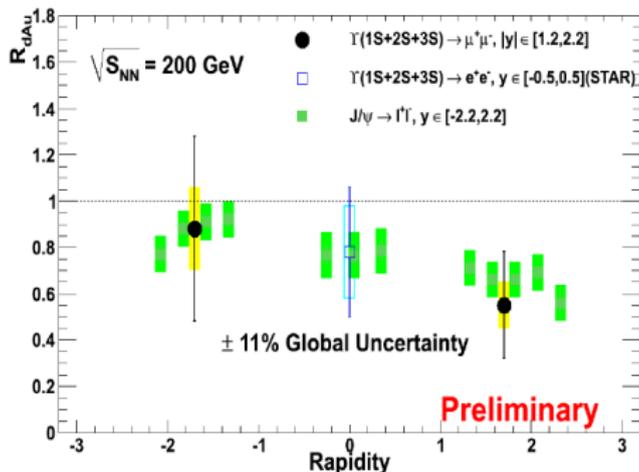
T. Frawley, Hard Probes 2012

- A priori same behaviour, but
- **Absorption** should be definitely **smaller**
- **Shadowing** should be **weaker** ( $x_2$  is 3 times and  $Q^2$  is 10 times larger)

Y vs  $J/\psi$ : different effects, same suppression ?

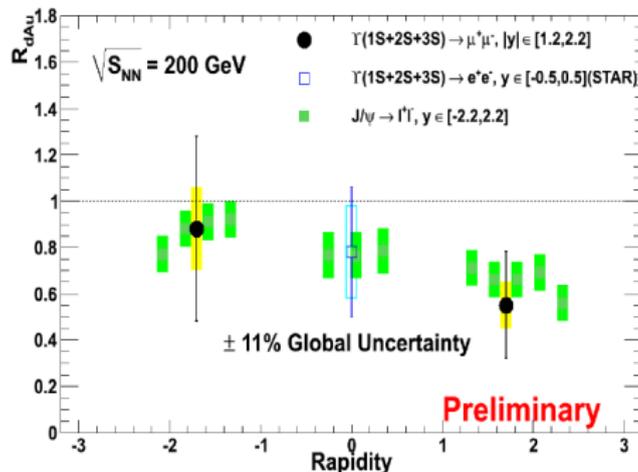
T. Frawley, Hard Probes 2012

- A priori same behaviour, but
- **Absorption** should be definitely **smaller**
- **Shadowing** should be **weaker** ( $x_2$  is 3 times and  $Q^2$  is 10 times larger)
- **Antishadowing** for  $y^Y \simeq -0.5$  instead for  $y^\psi \simeq -1.5$

Y vs  $J/\psi$ : different effects, same suppression ?

T. Frawley, Hard Probes 2012

- A priori same behaviour, but
- **Absorption** should be definitely **smaller**
- **Shadowing** should be **weaker** ( $x_2$  is 3 times and  $Q^2$  is 10 times larger)
- **Antishadowing** for  $y^Y \simeq -0.5$  instead for  $y^\psi \simeq -1.5$
- **Backward** Y's sit in the **EMC region**

Y vs  $J/\psi$ : different effects, same suppression?

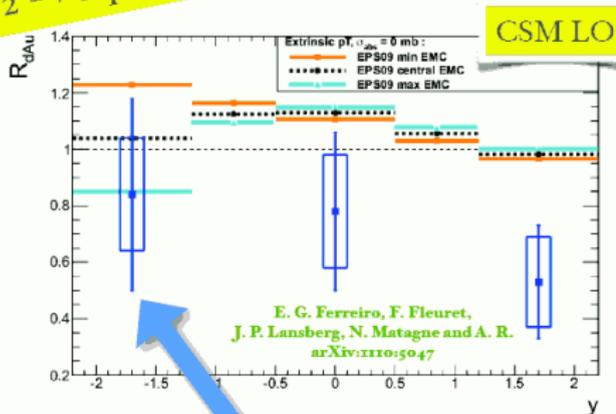
T. Frawley, Hard Probes 2012

- A priori same behaviour, but
- **Absorption** should be definitely **smaller**
- **Shadowing** should be **weaker** ( $x_2$  is 3 times and  $Q^2$  is 10 times larger)
- **Antishadowing** for  $y^Y \simeq -0.5$  instead for  $y^\psi \simeq -1.5$
- **Backward** Y's sit in the **EMC region**
- Note, however, than the error bars are large

Y production in  $dAu$  at RHIC and the gluon EMC effect

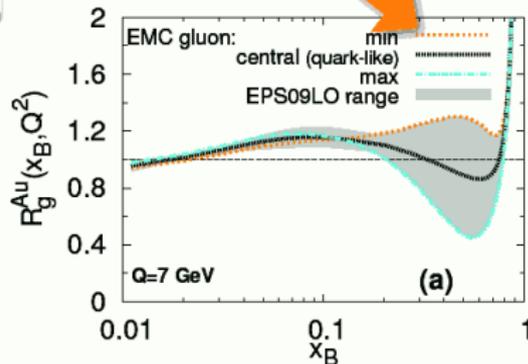
Let us focus in the EMC region and pick the EPS09 sets that are the limiting cases in this region :

$2 \rightarrow 2$  process



EMC effect stronger  
for  $g$  than for  $q$ ?

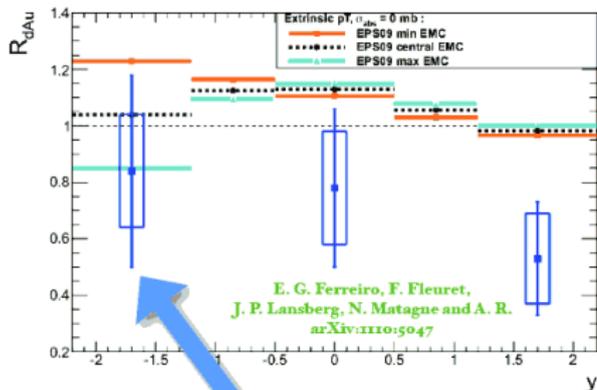
min. disfavoured



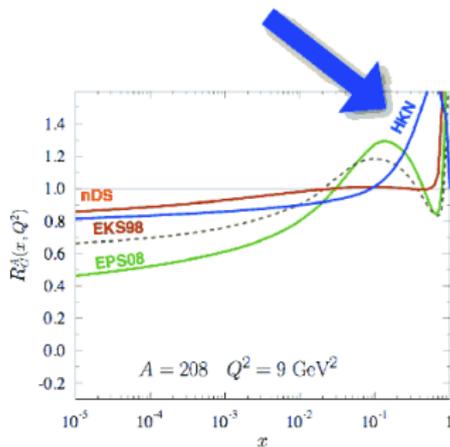
Y production in  $dAu$  at RHIC and the gluon EMC effect

Let us focus in the EMC region and pick the EPS09 sets that are the limiting cases in this region :

HKN disfavoured



EMC effect stronger  
for  $g$  than for  $q$ ?



Y production in  $dAu$  at RHIC in the mid and forward region

E. G. Ferreira, F. Fleuret,  
J. P. Lansberg, N. Matagne and A. R.  
arXiv:1110.5047

Typical gluon nPDF  
parametrisations induce a  
**flat rapidity dependence**  
w.r.t. data

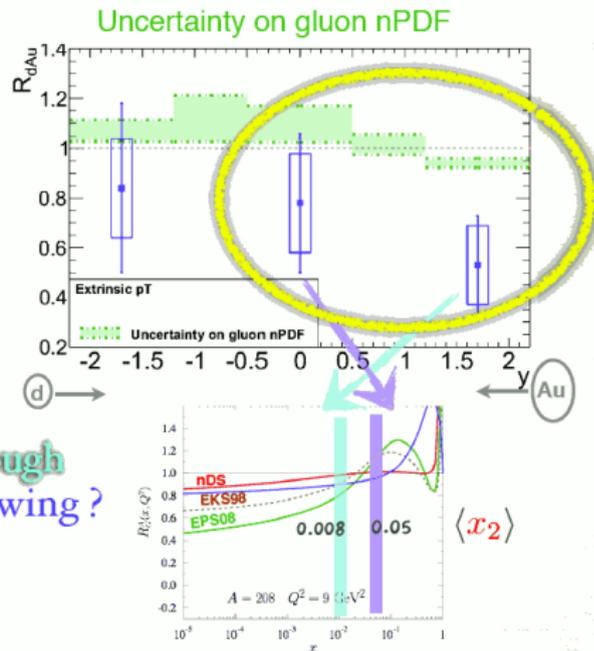
shadowing not strong enough  
absence of antishadowing?

Data:

STAR Preliminary, Nucl. Phys. A855 (2011) 440,

PRD 82 (2010) 012004.

PHENIX Preliminary, PoS DIS2010 (2010) 077.



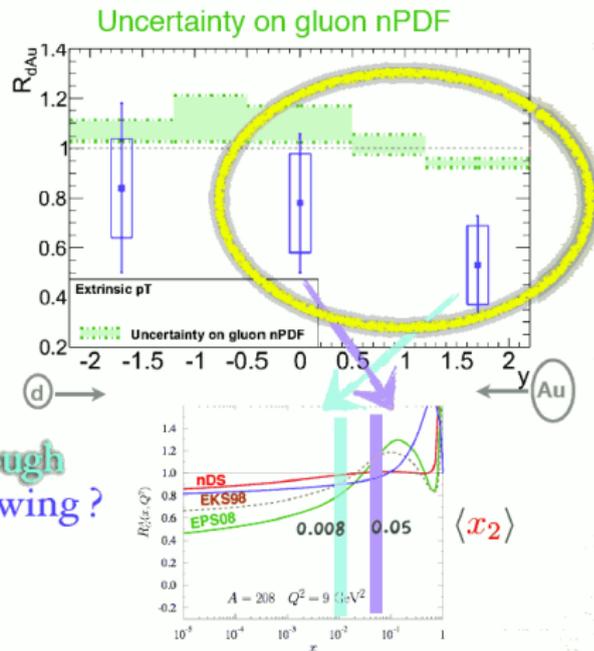
Y production in  $dAu$  at RHIC in the mid and forward region

E. G. Ferreira, F. Fleuret,  
J. P. Lansberg, N. Matagne and A. R.  
arXiv:1110.5047

Typical gluon nPDF  
parametrisations induce a  
**flat rapidity dependence**  
w.r.t. data

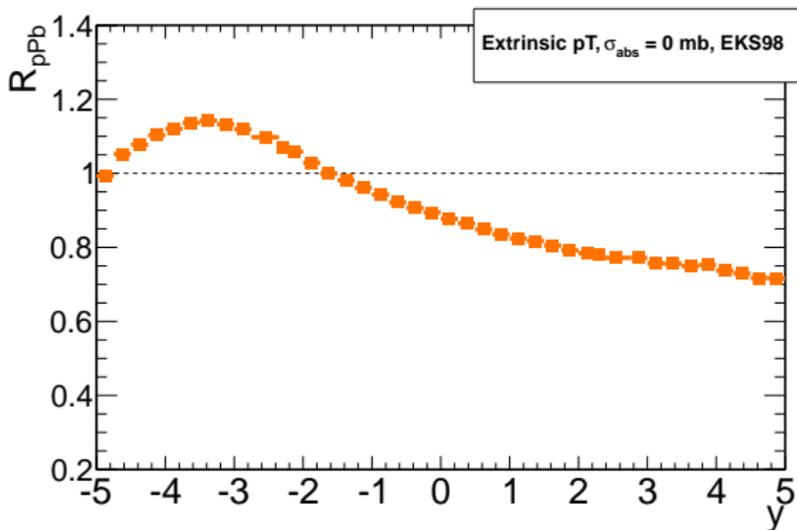
shadowing not strong enough  
absence of antishadowing?

Data:  
STAR Preliminary, Nucl. Phys. A855 (2011) 440,  
PRD 82 (2010) 012004.  
PHENIX Preliminary, PoS DIS2010 (2010) 077.



→ Fractional Energy loss à la Arleo-Peigné ?

# What's next: Y at 5 TeV



- Shadowing and antishadowing effects: **expected to be significant**
- Rules of thumb:  $Q_Y^2 \simeq 10 \times Q_{J/\psi}^2$  but also  $x_2^Y \simeq 3 \times x_2^{J/\psi}$
- Let's keep in mind that  $Q$  is nothing but  $\mu_F$  which **we do not know**

# Part IV

## Conclusions

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges
- One can take the **kinematics correctly** into account ( $2 \rightarrow 2$  at least)
- However, remaining **ambiguities due to  $k_T$  smearing**

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges
- One can take the **kinematics correctly** into account ( $2 \rightarrow 2$  at least)
- However, remaining **ambiguities due to  $k_T$  smearing**
- The  **$P_T$  cuts matter** (can be implemented with a  $2 \rightarrow 2$  kinematics)

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges
- One can take the **kinematics correctly** into account ( $2 \rightarrow 2$  at least)
- However, remaining **ambiguities due to  $k_T$  smearing**
- The  **$P_T$  cuts matter** (can be implemented with a  $2 \rightarrow 2$  kinematics)
- Conventional wisdom in  **$pp$**  is to **vary  $\mu_F$**  by 2 about a “natural” scale  
Very rarely done for **nuclear** PDF...

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges
- One can take the **kinematics correctly** into account ( $2 \rightarrow 2$  at least)
- However, remaining **ambiguities due to  $k_T$  smearing**
- The  **$P_T$  cuts matter** (can be implemented with a  $2 \rightarrow 2$  kinematics)
- Conventional wisdom in  $pp$  is to **vary  $\mu_F$**  by 2 about a “natural” scale  
Very rarely done for **nuclear PDF...**
- Not a perfect description of  $J/\psi$  data at RHIC and large uncertainties  
 **$pA$  data at the LHC** are awaited for !
- **$R_{CP}$**  should not be overlooked: no need for a  $pp$  baseline, avoid  $\Delta y$   
& info on impact-parameter d.o.f.

# Conclusions

- **Uncertainties within one set of nPDF** can be well taken into account (EPS09 for instance)
- However, **different nPDF sets** have **different uncertainty** ranges
- One can take the **kinematics correctly** into account ( $2 \rightarrow 2$  at least)
- However, remaining **ambiguities due to  $k_T$  smearing**
- The  **$P_T$  cuts matter** (can be implemented with a  $2 \rightarrow 2$  kinematics)
- Conventional wisdom in  **$pp$**  is to **vary  $\mu_F$**  by 2 about a “natural” scale  
Very rarely done for **nuclear PDF...**
- Not a perfect description of  **$J/\psi$**  data at RHIC and large uncertainties  
 **$pA$  data at the LHC** are awaited for !
- **$R_{CP}$**  should not be overlooked: no need for a  **$pp$**  baseline, avoid  $\Delta y$   
& info on impact-parameter d.o.f.
- **$Y$**  offer a new playground: **smaller** uncertainties and a priori **new** effects
- At RHIC,  **$Y$  suppression** shows the same trend as for  **$J/\psi$** ,

◀ ▶ ◀ ▶ this is **unnatural** ↻ 🔍

# Part V

## Backup

# On the kinematics of quarkonium production

If  $\mathcal{F}_g^A(x, \vec{r}, z, \mu_f)$  gives the **distribution of a gluon** of mom. fract.  $x$  at a **position  $\vec{r}, z$  in a nucleus  $A$** , the differential cross-section reads:

$$\frac{d\sigma_{AB}^Q}{dy dP_T d\vec{b}} =$$

**2**  $\rightarrow$  **1** kinematics with **intrinsic**  $p_T$

**2**  $\rightarrow$  **2** kinematics with **extrinsic**  $p_T$

# On the kinematics of quarkonium production

If  $\mathcal{F}_g^A(x, \vec{r}, z, \mu_f)$  gives the **distribution of a gluon** of mom. fract.  $x$  at a **position  $\vec{r}, z$  in a nucleus  $A$** , the differential cross-section reads:

$$\frac{d\sigma_{AB}^Q}{dy dP_T d\vec{b}} =$$

**2 → 1** kinematics with **intrinsic**  $p_T$

$$\begin{aligned} & \int d\vec{r}_A dz_A dz_B \\ & \times \mathcal{F}_g^A(x_1^0, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2^0, \vec{r}_B, z_B, \mu_f) \\ & \times \sigma_{gg}^{\text{Intr.}}(x_1^0, x_2^0) \\ & \times \mathcal{S}_A(\vec{r}_A, z_A) \mathcal{S}_B(\vec{r}_B, z_B) \end{aligned}$$

**2 → 2** kinematics with **extrinsic**  $p_T$

$$\begin{aligned} & \int dx_1 dx_2 \int d\vec{r}_A dz_A dz_B \\ & \times \mathcal{F}_g^A(x_1, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2, \vec{r}_B, z_B, \mu_f) \\ & \times 2\hat{s}P_T \frac{d\sigma_{gg \rightarrow Q+g}}{d\hat{t}} \delta(\hat{s} - \hat{t} - \hat{u} - M^2) \\ & \times \mathcal{S}_A(\vec{r}, z_A) \mathcal{S}_B(\vec{r}_B, z_B) \end{aligned}$$

# On the kinematics of quarkonium production

If  $\mathcal{F}_g^A(x, \vec{r}, z, \mu_f)$  gives the **distribution of a gluon** of mom. fract.  $x$  at a **position  $\vec{r}, z$  in a nucleus  $A$** , the differential cross-section reads:

$$\frac{d\sigma_{AB}^Q}{dy dP_T d\vec{b}} =$$

**2 → 1** kinematics with **intrinsic**  $p_T$

$$\begin{aligned} & \int d\vec{r}_A dz_A dz_B \\ & \times \mathcal{F}_g^A(x_1^0, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2^0, \vec{r}_B, z_B, \mu_f) \\ & \times \sigma_{gg}^{\text{Intr.}}(x_1^0, x_2^0) \\ & \times \mathcal{S}_A(\vec{r}_A, z_A) \mathcal{S}_B(\vec{r}_B, z_B) \end{aligned}$$

$$x_{1,2} = \frac{m_T}{\sqrt{s_{NN}}} \exp(\pm y) \equiv x_{1,2}^0(y, P_T)$$

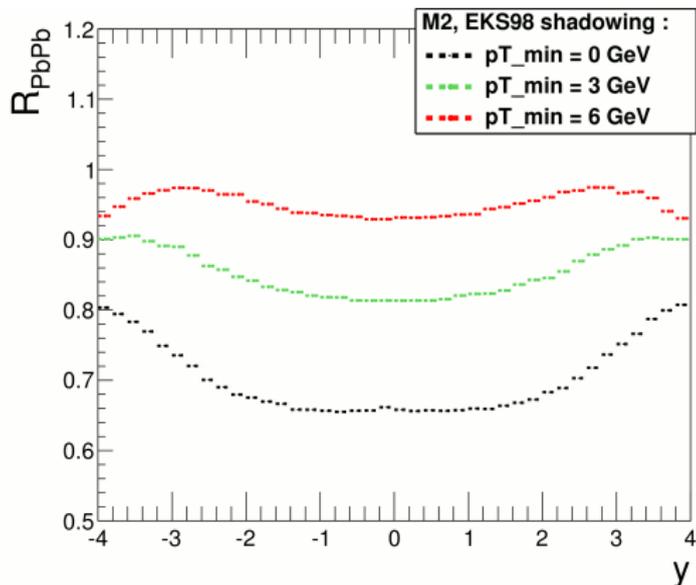
**2 → 2** kinematics with **extrinsic**  $p_T$

$$\begin{aligned} & \int dx_1 dx_2 \int d\vec{r}_A dz_A dz_B \\ & \times \mathcal{F}_g^A(x_1, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(x_2, \vec{r}_B, z_B, \mu_f) \\ & \times 2\hat{S}P_T \frac{d\sigma_{gg \rightarrow Q+g}}{d\hat{t}} \delta(\hat{s} - \hat{t} - \hat{u} - M^2) \\ & \times \mathcal{S}_A(\vec{r}, z_A) \mathcal{S}_B(\vec{r}_B, z_B) \end{aligned}$$

$$\delta(\dots) \rightarrow x_2 = \frac{x_1 m_T \sqrt{s_{NN}} e^{-y} - M^2}{\sqrt{s_{NN}} (\sqrt{s_{NN}} x_1 - m_T e^y)}$$

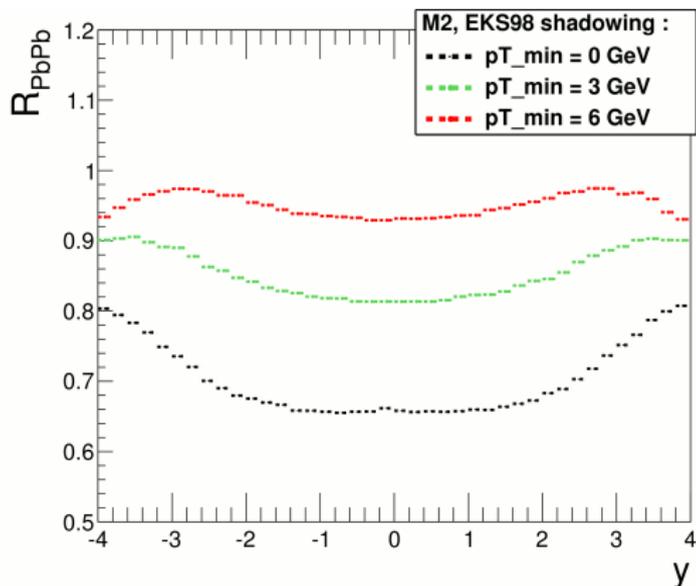
# The $P_T$ cut matters

E.G. Ferreira, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330 + in progress



# The $P_T$ cut matters, the nPDF as well ...

E.G. Ferreira, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330 + in progress

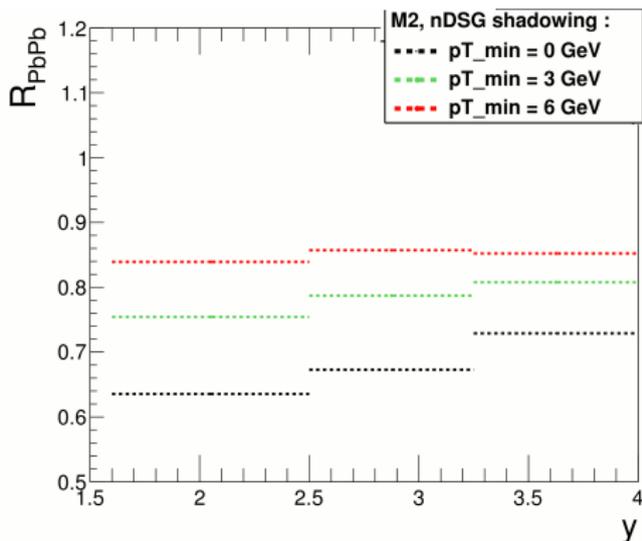
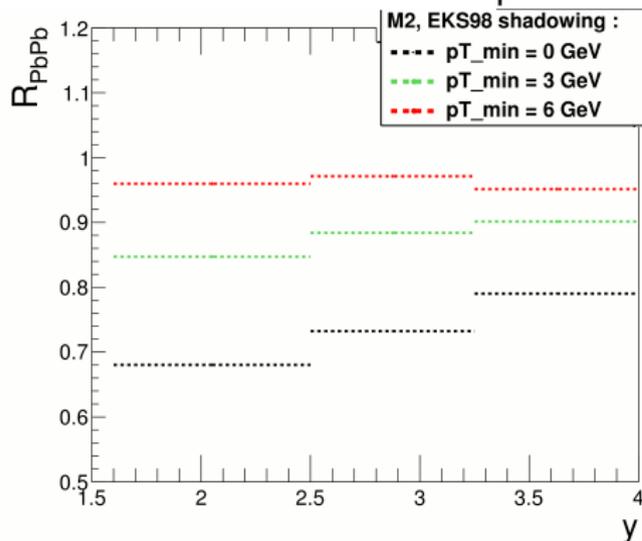


**WARNING:** The behaviour depends on the nPDF

# The $P_T$ cut matters, the nPDF as well ...

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330 + in progress

**WARNING:** The behaviour depends on the nPDF



- For larger  $P_T$  cuts, the shadowing decreases (less true with nDSg)

- $Q^2 \equiv \mu_F^2 \sim \sqrt{M^2 + P_T^2}$

nDSg mild  $Q^2$  evolution

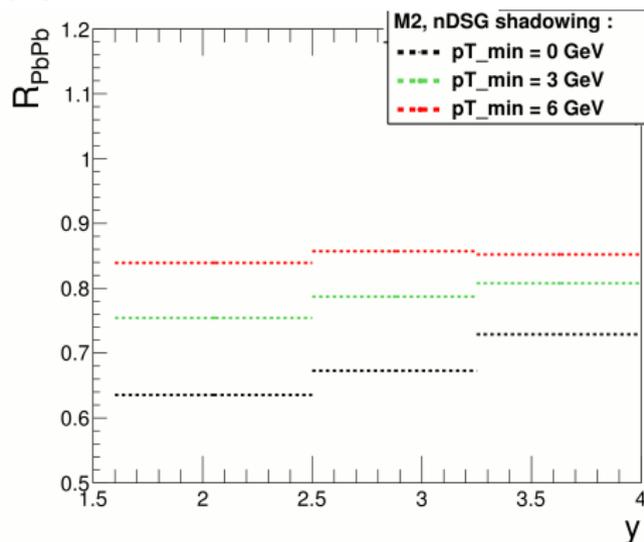
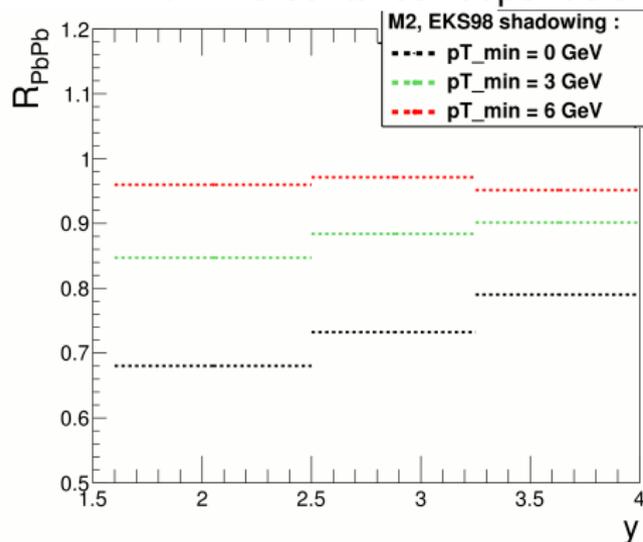
- But also, when  $P_T \nearrow$ ,  $x_{1,2} \nearrow$

nDSg flatter than EKS98 for  $x \leq 0.1$

# The $P_T$ cut matters, the nPDF as well ...

E.G. Ferreiro, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330 + in progress

**WARNING:** The behaviour depends on the nPDF



- For larger  $P_T$  cuts, the shadowing decreases (less true with nDSg)

- $Q^2 \equiv \mu_F^2 \sim \sqrt{M^2 + P_T^2}$

nDSg mild  $Q^2$  evolution

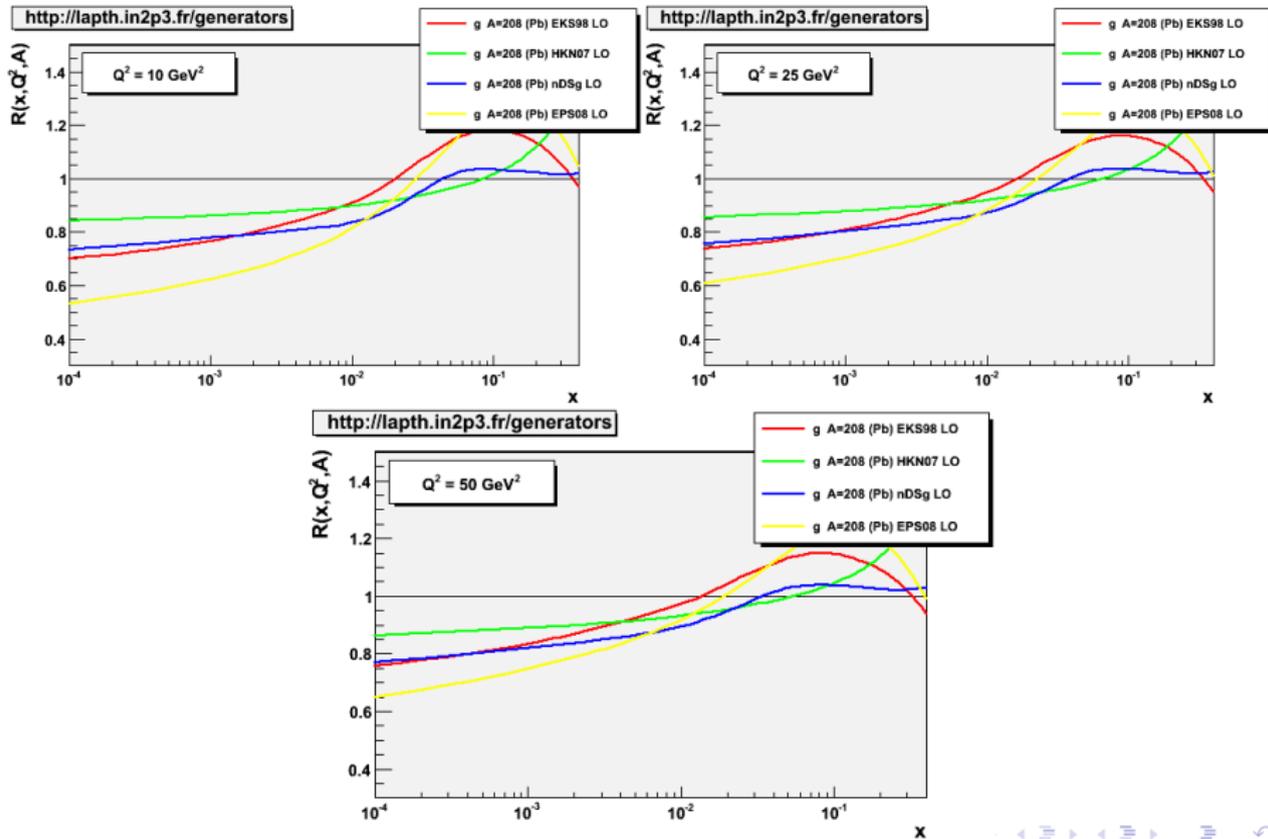
- But also, when  $P_T \nearrow$ ,  $x_{1,2} \nearrow$

nDSg flatter than EKS98 for  $x \leq 0.1$

- Non trivial modification of the  $y$ -dependence



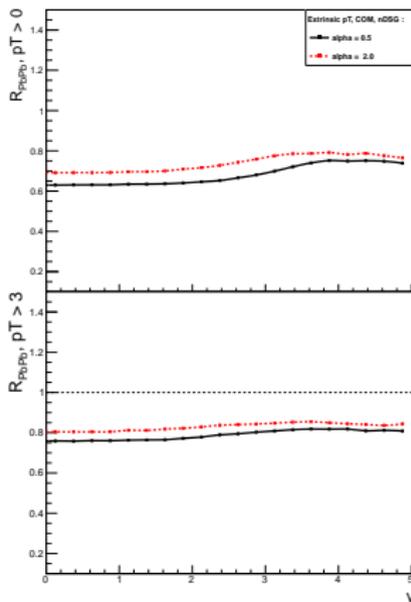
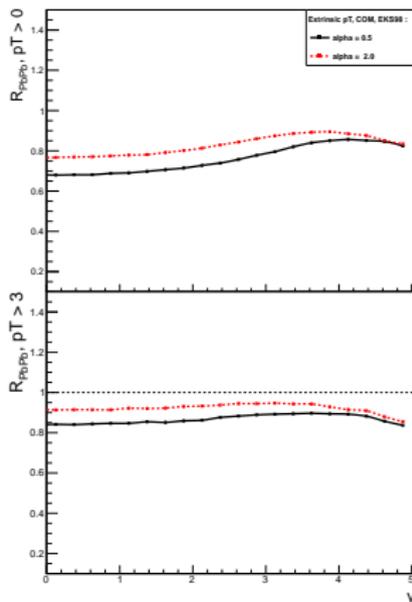
# Gluon shadowing at different scales for Pb ions



# The $P_T$ cut matters

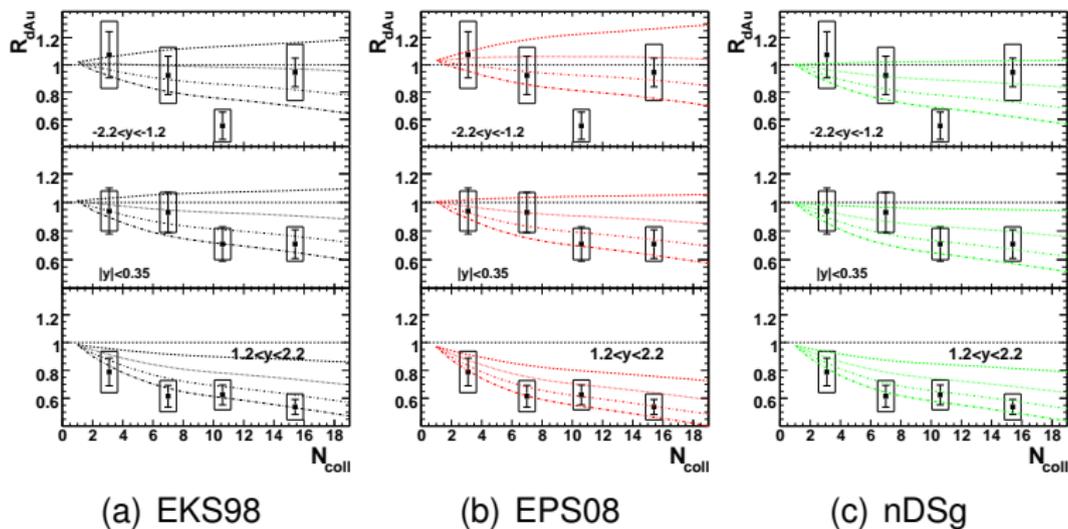
E.G. Ferreira, F. Fleuret, J.P.L., N. Matagne, A. Rakotozafindrabe, Nucl. Phys. A 855 (2011) 327-330 + in progress

- CMS could not go below  $P_T = 3$  GeV, and only in the slightly forward region
- Can one expect a big difference for the CNM when a cut  $P_T \geq 3$  GeV is applied ? not really ...



# Centrality dependence in $dAu$

E.G. Ferreiro, F. Fleuret, J.P.L., A. Rakotozafindrabe, PLB 680:50,2009, PRC 81, 064911 (2010)



- Backward region  $-2.2 < y < 1.2$ : the gluon in the Au is very energetic (large  $x$ )  
→ antishadowing
- Central region  $-0.35 < y < 0.35$ : the gluon in the Au is energetic (mid  $x$ )  
→ slight shadowing
- Forward region  $1.2 < y < 2.2$ : the gluon in the Au is not energetic (small  $x$ )  
→ stronger shadowing