

# Progress with applications of the stochastic LapH method to excited states

Keisuke Jimmy Juge  
University of the Pacific

John Bulava (CERN)

Brendan Fahy (CMU)

Justin Foley (U Utah)

You-Cyuan Jhang (CMU)

David Lenkner (CMU)

Colin Morningstar (CMU)

Chik Him Wong (UC San Diego)

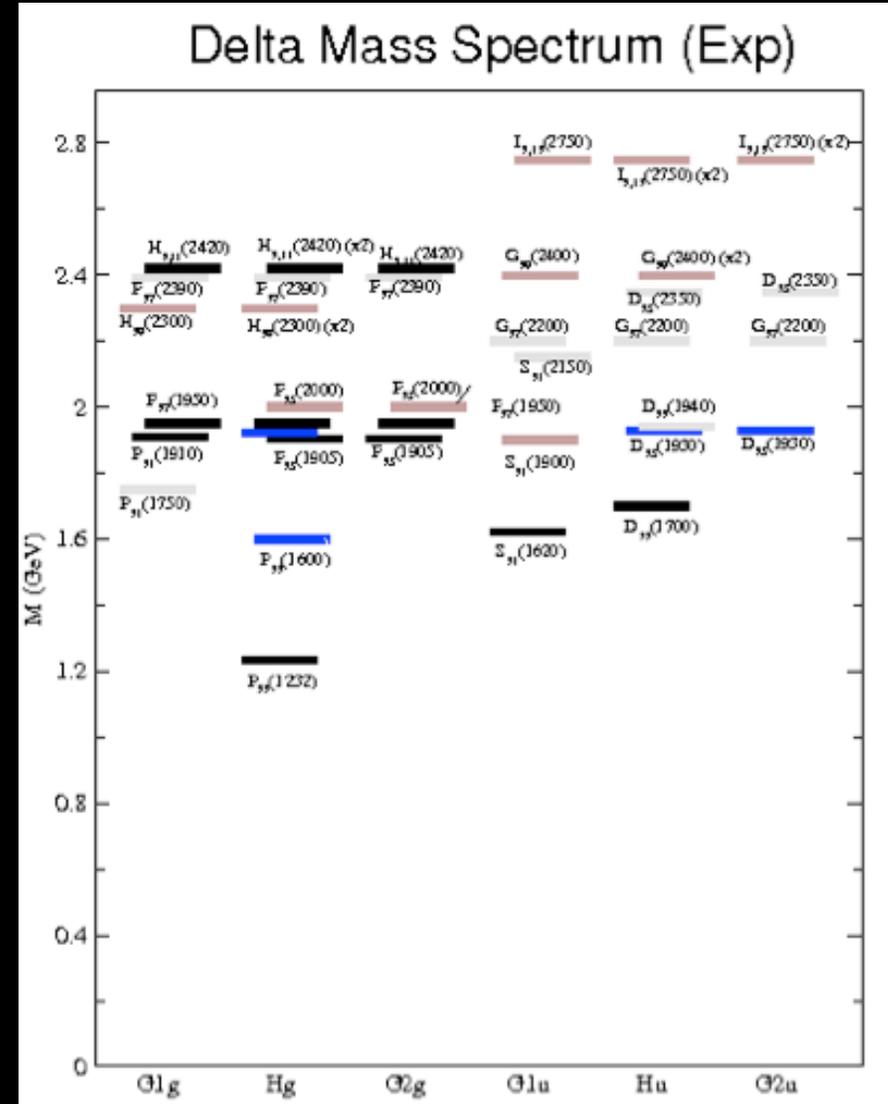
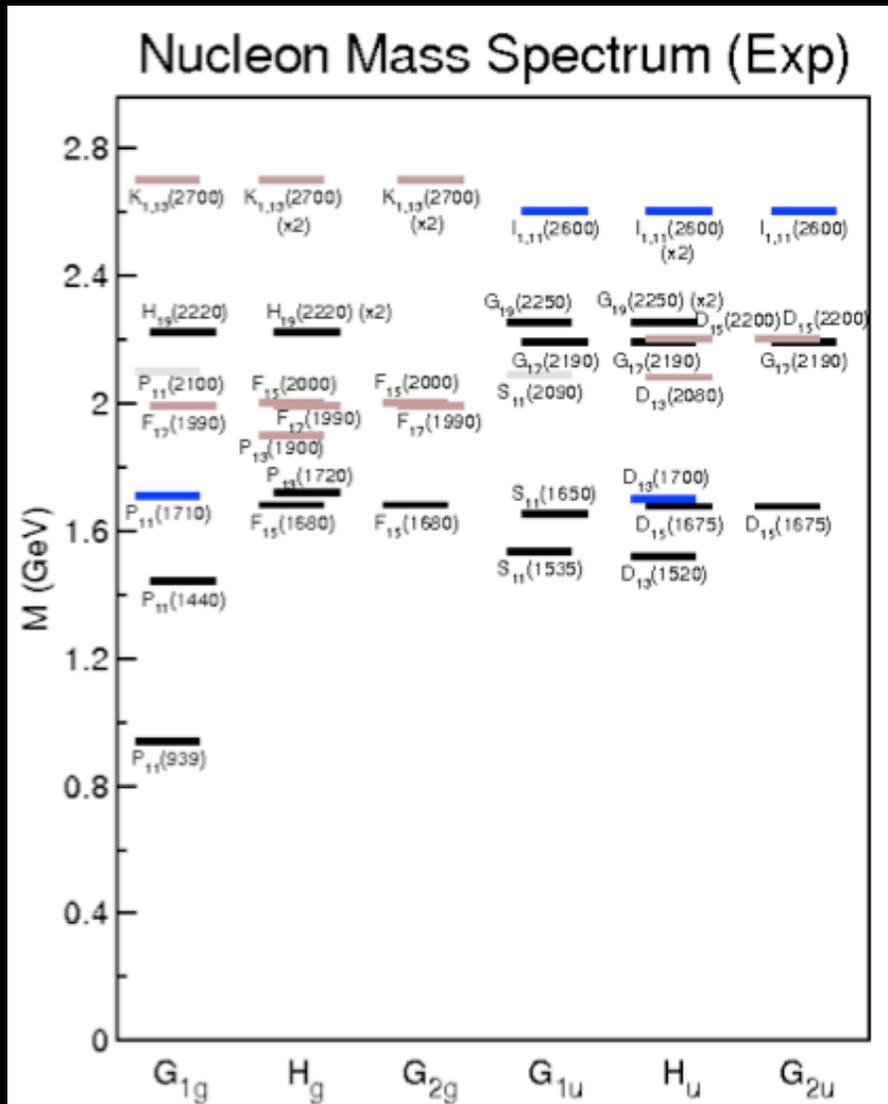
# Outline

- Introduction
  - goal, setup/approach
- Stochastic LqH
  - single meson correlator construction
- Two particle channel
  - Preliminary partial results for  $I=0,2$   $\pi$ - $\pi$
- $I=1$ , glueball mixing, baryons
- Summary

# Introduction

Goal: First principles calculation of the low-lying light hadron spectra

HadSpec Phys.Rev. D83 (2011)  
 PACS-CS Phys.Rev. D81 (2010)  
 BMW Phys.Lett. B701 (2011)  
 RBC,UKQCD Phys.Rev.D83 B701 (2011)  
 BGR Phys.Rev.D83 B701 (2011)



- We want to construct good interpolating operators that couple to the physical states in question which can be used for spectroscopy and other lattice calculations

## Require:

- anisotropic, dynamical lattices

Nf=2+1 anisotropic clover

$$a_s \simeq 0.12 \text{ fm} \quad a_t^{-1} \simeq 6.0 \text{ GeV} \quad \xi = 3.5$$

$$m_\pi \simeq 240 \text{ MeV}, 390 \text{ MeV}$$

HadSpec Phys.Rev. D79, 2009

- "large" volumes: L=1.8fm, L=2.3fm, L=2.8fm, L=3.7fm  
most recent being generated
- good interpolating operators

Fight noise!

cubical symmetry, group theory (Basak et al)  
(cubical lattice in the continuum limit)

- multi-particle channel analysis  
mixings and scattering

# Anisotropic Dynamical Lattices

Edwards et al. PRD78 (2008)

$$16^3 \times 128$$

$$m_\pi = 380 \text{ MeV}$$

$$m_\pi L \simeq 3.8$$

small and heavy  
but useful for  
pruning

$$24^3 \times 128$$

$$m_\pi = 240 \text{ MeV}$$

$$m_\pi L \simeq 5.7$$

$$m_\pi L \simeq 3.5$$

$\sim 350$  configs

$$32^3 \times 256$$

$$m_\pi = 240 \text{ MeV}$$

$$m_\pi L \simeq 4.7$$

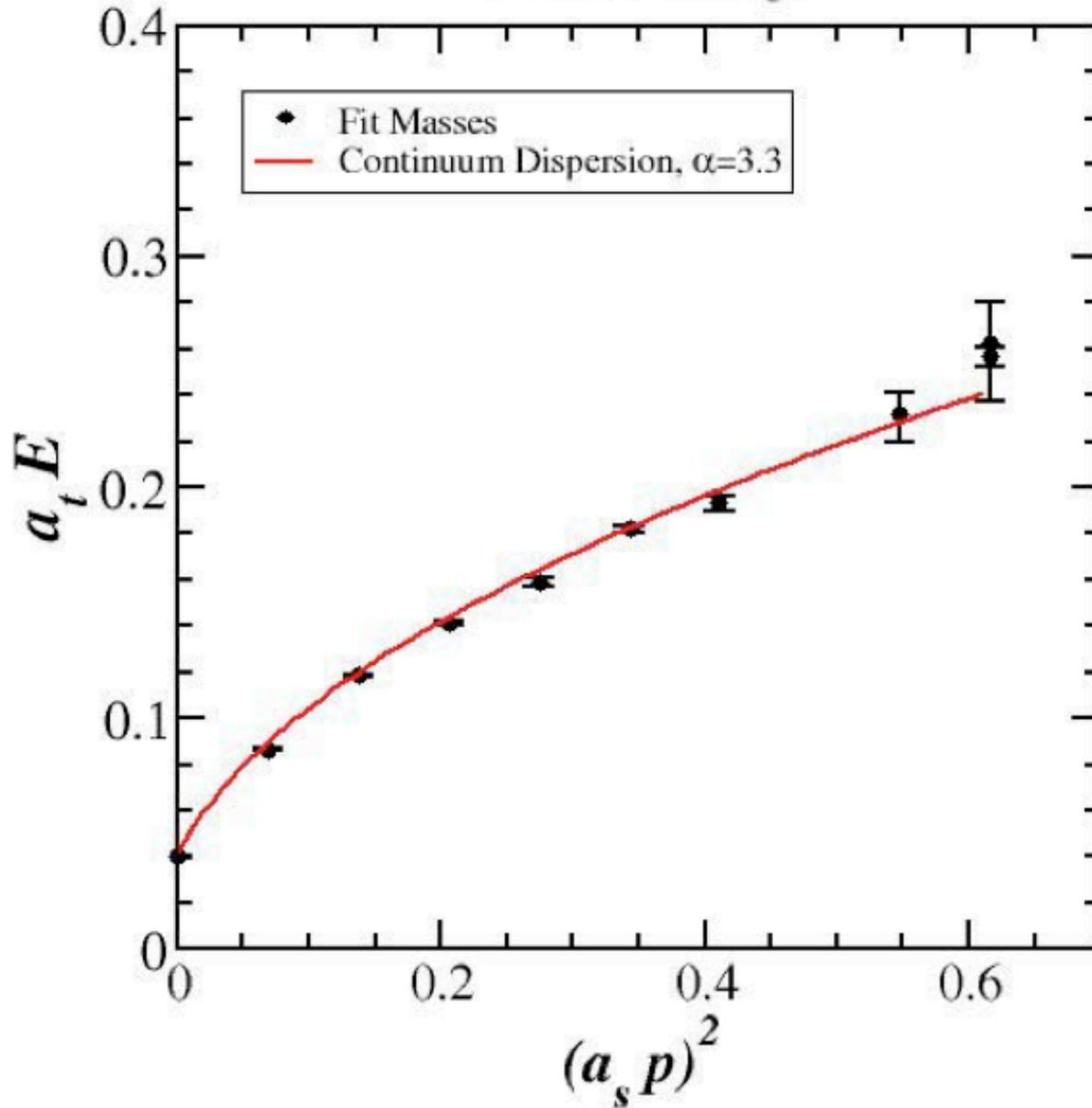
new

anisotropy = 3.459(4) Dudek et al (2011)

lightest quark mass = 3.3 (preliminary)

# Pion Dispersion Relation

$24^3 \times 128$ , 584 configs



$24^3 \times 128$   
 $m_\pi = 240 \text{ MeV}$

# Single hadron operators

- Smearing, covariantly-displaced “source” fields and “sink” fields
  - links: stout-smearing ... **noise reduction**
  - quarks: LapH smearing ... **better overlap**
- Separation of “source” and “sink” operators via **stochastic LapH** (all-to-all) quark propagators (PRD83,2011)
- Color-singlet, elemental operators which transform as one of the irreps of interest
- **Group theoretical projection** of the elemental operators (Basak et al PRD72, 2005)

$$H^{\Lambda\lambda F}(t) = \frac{d_{\Lambda}}{g_G} \sum_{R \in G} D_{\lambda\lambda}^{(\Lambda)}(R^*) U_R H_i^F(t) U_R^\dagger$$

# Stochastic LapH Propagators

- LapH smearing (Distillation PRD80,2009)

- smear quarks and propagate the low modes

$$\tilde{\psi}(x, t) = S(x, y; t)\psi(y, t)$$

$$= \mathbf{V}_s \mathbf{V}_s^\dagger \psi(y, t)$$

$\mathbf{V}_s$  matrix of the eigenvectors of  $\nabla^2$

control amount of smearing (overlap) by varying number of eigenvectors that are retained

3D Laplacian

- 3-D eigenvector calculations vs 4-D
- **Stochastic LapH**: combine LapH with stochastic method
  - introduce only in the complementary space
  - use dilution projectors to partition the noise 'wisely'

# Stochastic LapH

Quark Propagator on a given config

Diluted Z4 noise source  
(in LapH subspace)

→ time, spin  
LapH eigenvector #  
(Nev=112)

Foley et al. Com.Phy.Com.172 (2005)

dilution projectors  $\mathcal{P}^{(i)}$  {  
- full (**F**)  
- interlace (**I#**) "skip"  
- block (**B#**)

example:  $\mathcal{P}^{(t=3)}$  picks out t=3 timeslice

full noise source  $\varrho(x, t) = \sum_{i=1}^{N_d} \mathcal{P}^{(i)} \varrho(x, t)$   
 $= \sum_{i=1}^{N_d} \varrho^{(i)}(x, t)$

**Source:** Smear and Displace

$$\tilde{\varrho}_k^{(d)} = \tilde{D}_k \mathbf{V}_s \varrho^{(d)}$$

**Sink:** Solve for solution, project  
(solution) on eigenvector and smear

$$\tilde{\varphi}_j^{(d)} = \tilde{D}_j \mathbf{V}_s \mathbf{V}_s^\dagger M^{-1} \tilde{\varrho}_k^{(d)}$$

## Quark Propagator

$$Q(x, t; x_0, t_0) = \sum_{d=1}^{N_d} \tilde{\varphi}_j^{(d)}(x, t) \tilde{\varrho}_k^{(d)\dagger}(x_0, t_0)$$

$$\bar{Q}_{jk} = \sum_a \gamma_5 \gamma_4 \tilde{\varrho}_k^{(a)}(t_0) \tilde{\varphi}_j^{(a)\dagger}(t) \gamma_4 \gamma_5$$

## Elemental baryon (sink) operator

$$\mathcal{B}_l^F(t) = \sum_x e^{-ipx} \phi_{ABC}^F \epsilon_{abc} \left( \tilde{D}_i^{(d)} \tilde{\varphi}(x, t)_{Aa\alpha} \right) \left( \tilde{D}_j^{(d)} \tilde{\varphi}(x, t)_{Bb\beta} \right) \left( \tilde{D}_k^{(d)} \tilde{\varphi}(x, t)_{Cc\gamma} \right)$$

$\tilde{\varphi}$  solution to  $\mathbf{M}\tilde{\varphi}(x, t) = \tilde{\varrho}(x_0, t_0)$

## Elemental meson (sink) operator

$$\mathcal{M}_l^F(t) = \sum_x e^{-ip \cdot x} \phi_{AB}^F \delta_{ab} \left( \tilde{D}_i^{(d)} \tilde{\varrho}(x, t) \right)_{Aa\alpha}^\dagger \left( \tilde{D}_j^{(d)} \tilde{\varphi}(x, t) \right)_{Bb\beta}$$

'tilde' operators are stout-link smeared

Morningstar, Peardon PRD69 (2004)

Correlate the creation and annihilation operators

$$\langle \mathcal{B}_l(t) \mathcal{B}_{l'}^{src}(t_0) \rangle \quad \langle \mathcal{M}_l(t) \mathcal{M}_{l'}^{src}(t_0) \rangle$$

to form the correlation matrix

## Diagonalize (solve generalized eigenvalue problem)

$$\mathbf{C}(t_0)^{-1/2} \mathbf{C}(t) \mathbf{C}(t_0)^{-1/2} v = \lambda(t, t_0) v$$

The system determines the optimal linear combination of the operators for best overlap

We typically fix  $t_0$  and choose  $t^*$  (and  $t_0$ ) large where noise is still manageable

(Alpha Collaboration 2009)

Typical optimization range

$$t_0 = 3 \quad t^* = 8$$

but varies for each simulation

# Lattice Technology Summary

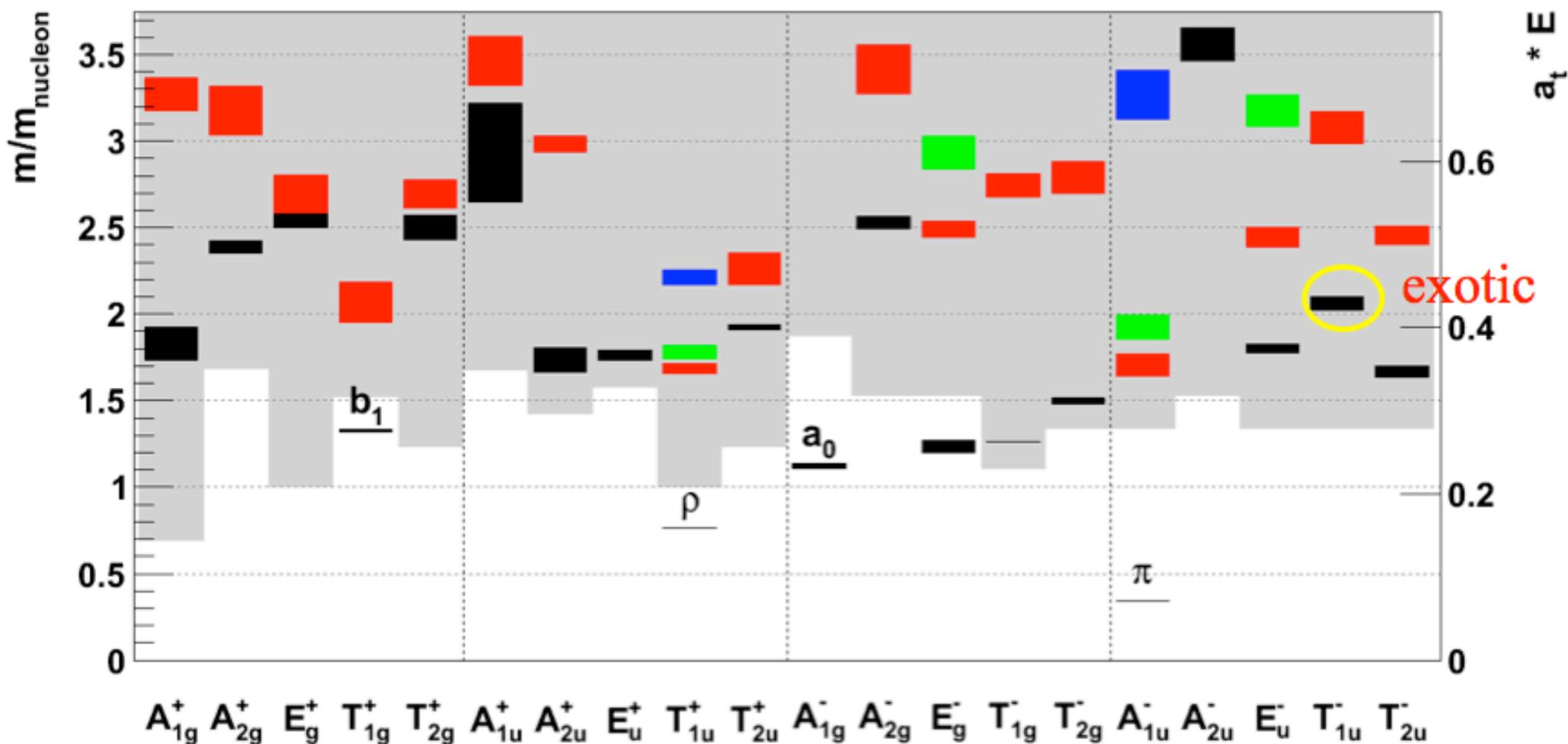
- anisotropic, 2+1 dynamical lattices (HadSpec Coll)
  - anisotropy 3.5
    - cutoff 6 GeV but still have 3 fm lattices
  - pion masses (240 MeV, 400 MeV)
- stout link smearing (Morningstar,Peardon)
  - link noise reduction
- stochastic LapH (all-to-all) quark propagators
  - disconnected diagrams, finite momentum operators
  - source and sink separation (in correlation functions)

# Lattice Technology Summary

- group theory projections (cubic group)
- cubical lattice even in the continuum limit
- optimization of operators via generalized eigenvalue method for use in other calculations

# Mesons

Isovector spectrum presented at Lattice '11  
 $24^3 \times 128$   $m_\pi = 380$  MeV 170 cfgs  
 shaded region needs multi-particle analysis



Multi-particle analysis requires meson operators with finite momentum

Injecting momentum is simple with all-to-all propagators

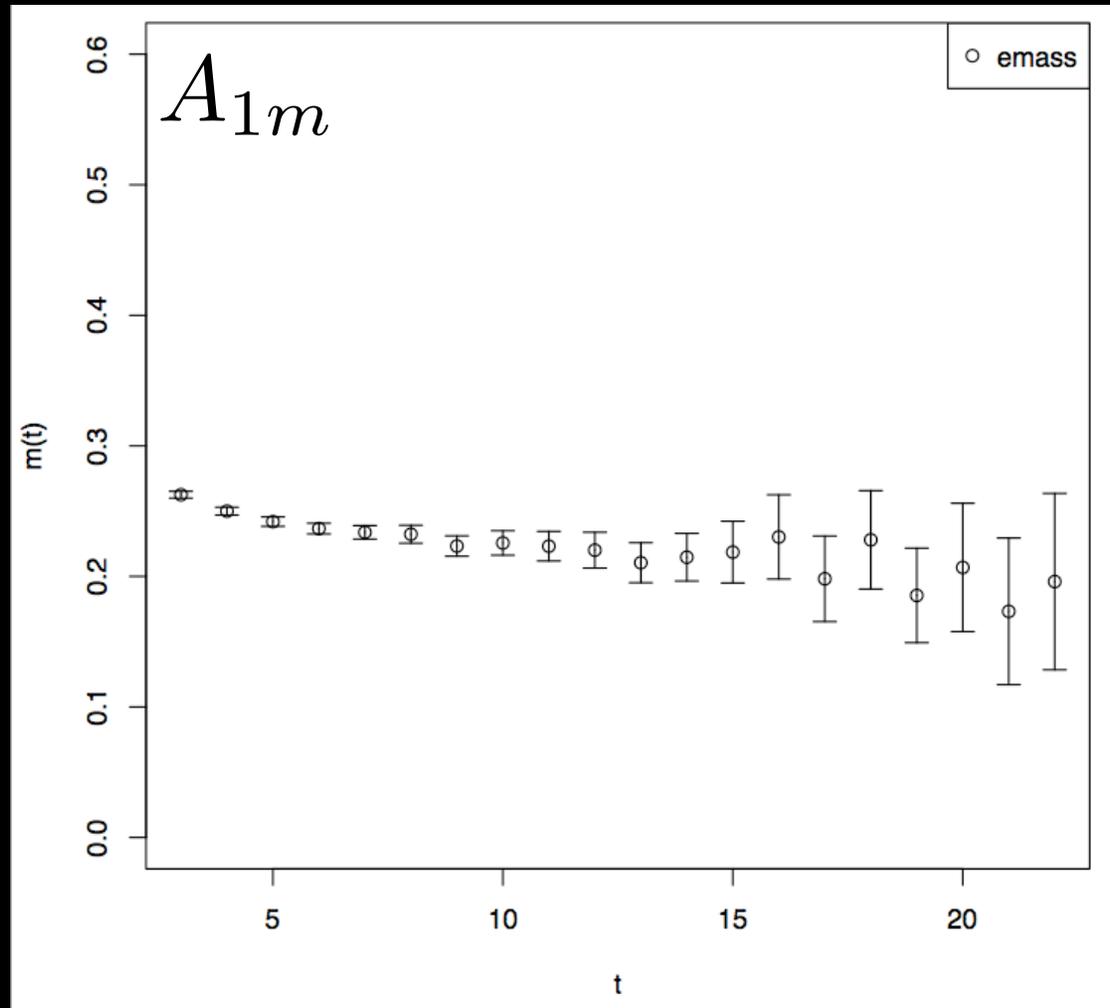
recall

$$\mathcal{B}_l^F(t) = \sum_x e^{-ipx} \phi_{ABC}^F \epsilon_{abc} \left( \tilde{D}_i^{(d)} \tilde{\varphi}(x, t)_{Aa\alpha} \right) \left( \tilde{D}_j^{(d)} \tilde{\varphi}(x, t)_{Bb\beta} \right) \left( \tilde{D}_k^{(d)} \tilde{\varphi}(x, t)_{Cc\gamma} \right)$$

$$\mathcal{M}_l^F(t) = \sum_x e^{-ip \cdot x} \phi_{AB}^F \delta_{ab} \left( \tilde{D}_i^{(d)} \tilde{\varrho}(x, t) \right)_{Aa\alpha}^\dagger \left( \tilde{D}_j^{(d)} \tilde{\varphi}(x, t) \right)_{Bb\beta}$$

Effective masses  $m_{\text{eff}} = -\ln \frac{C(t+1)}{C(t)} \rightarrow E$

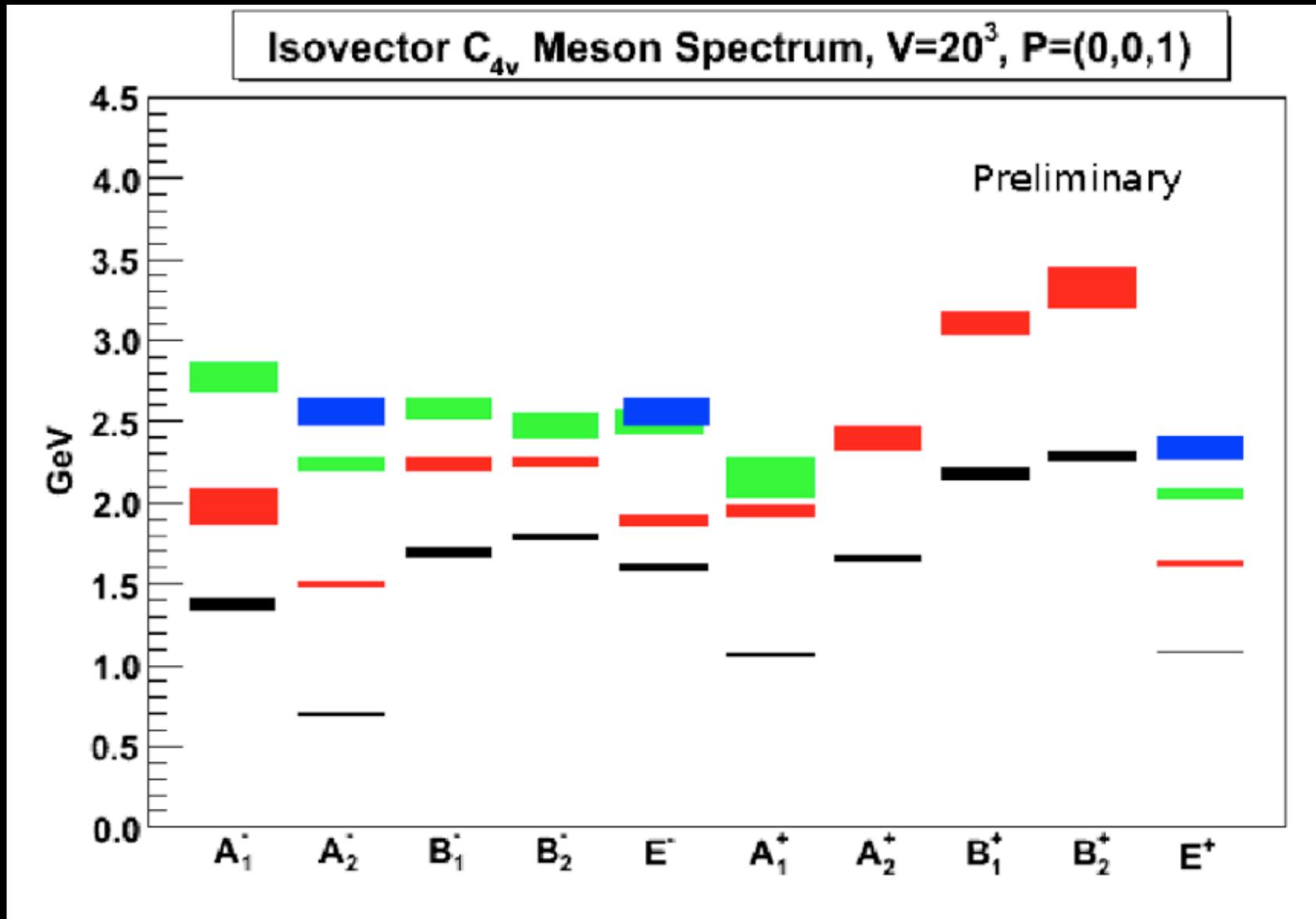
Example of an isovector with momentum (0,1,1)



$$24^3 \times 128$$
$$m_\pi = 240 \text{ MeV}$$

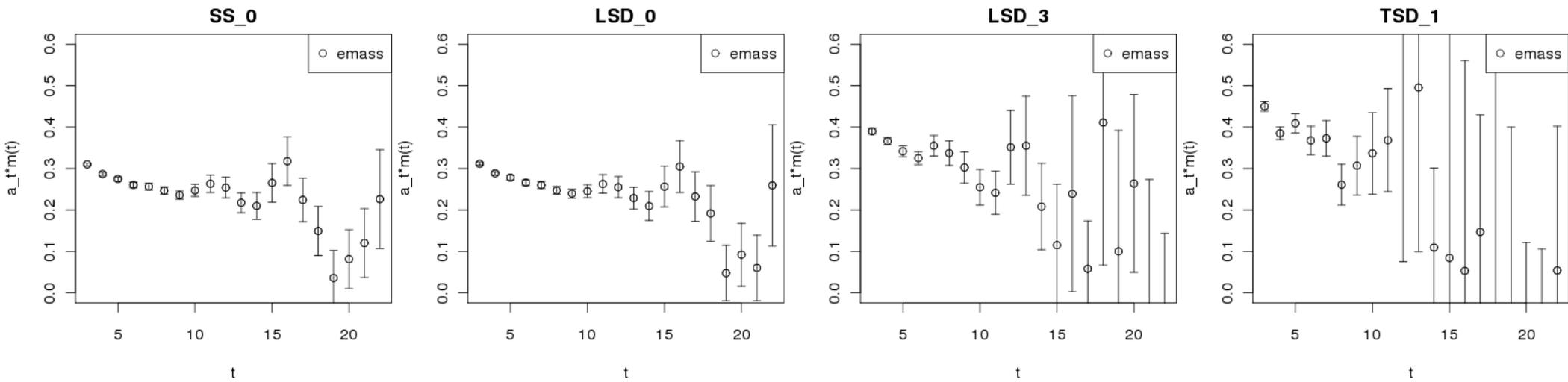
# Moving Mesons

$$20^3 \times 128 m_\pi = 380 \text{ MeV} \quad p = (0, 0, 1)$$

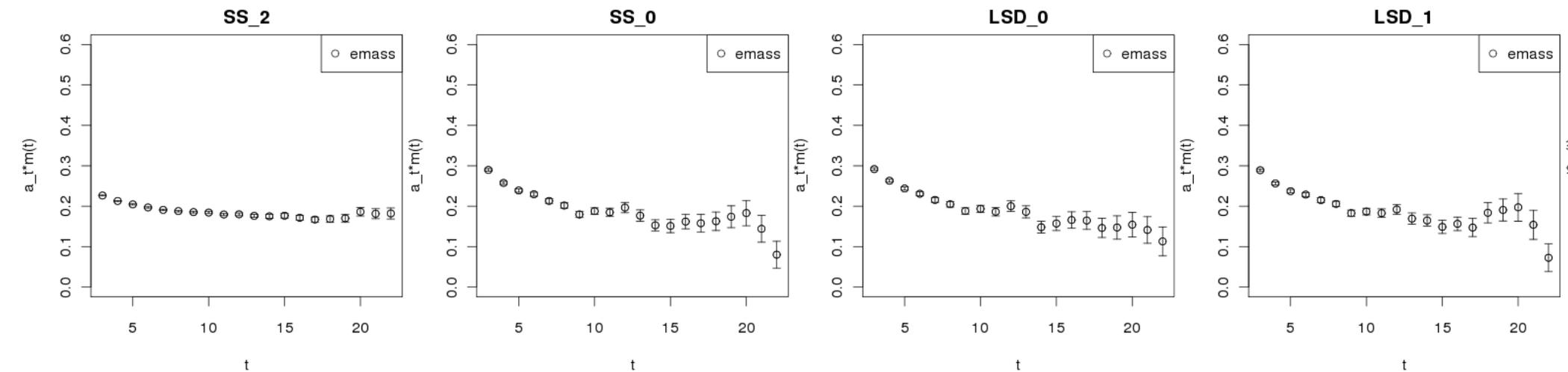


$24^3 \times 128$   
 $m_\pi = 240 \text{ MeV}$

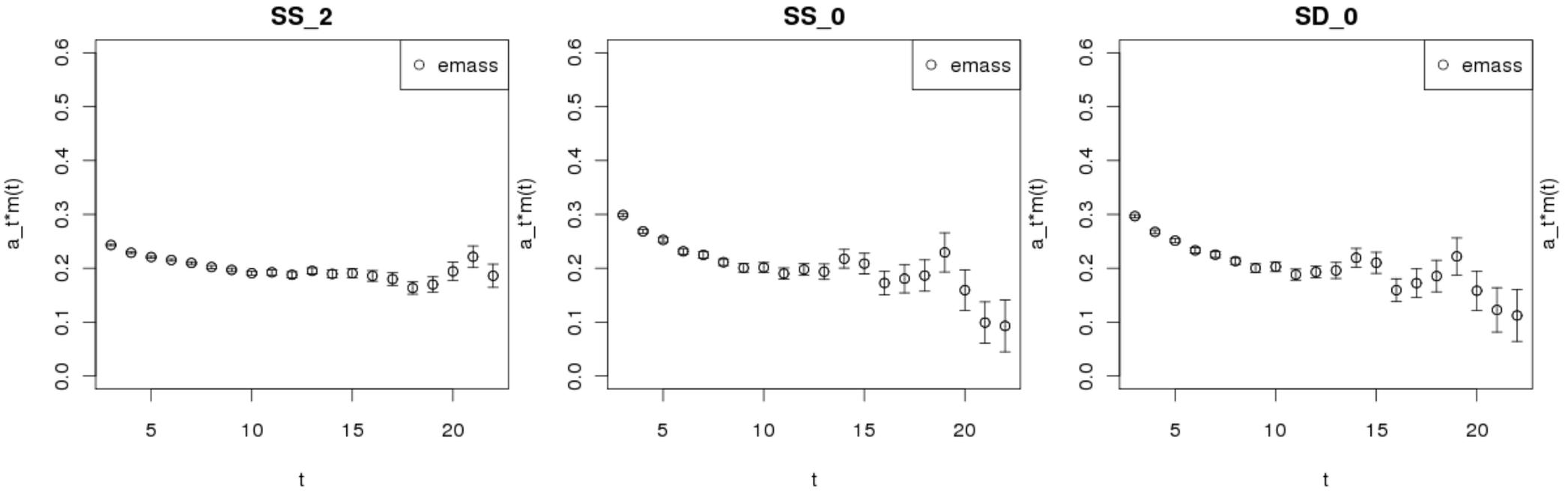
$A1m \langle 0,0,2 \rangle$  ← momentum



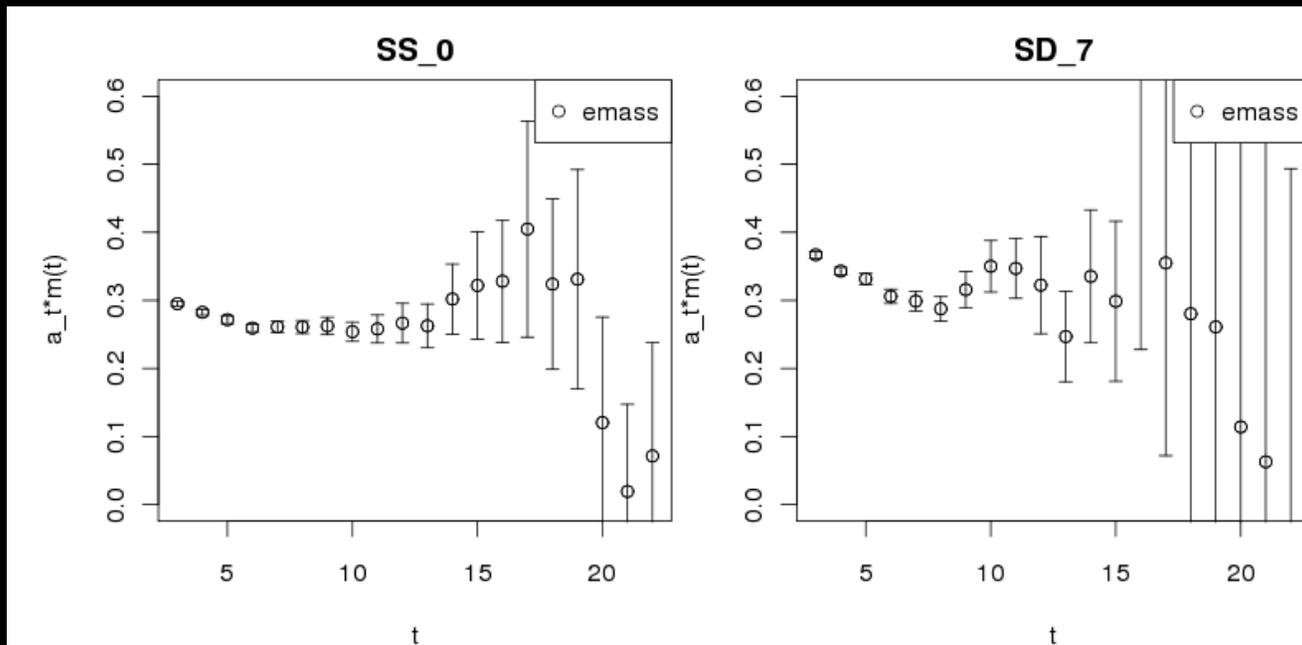
$A1p \langle 0,1,1 \rangle$



# A1p <1,1,1>

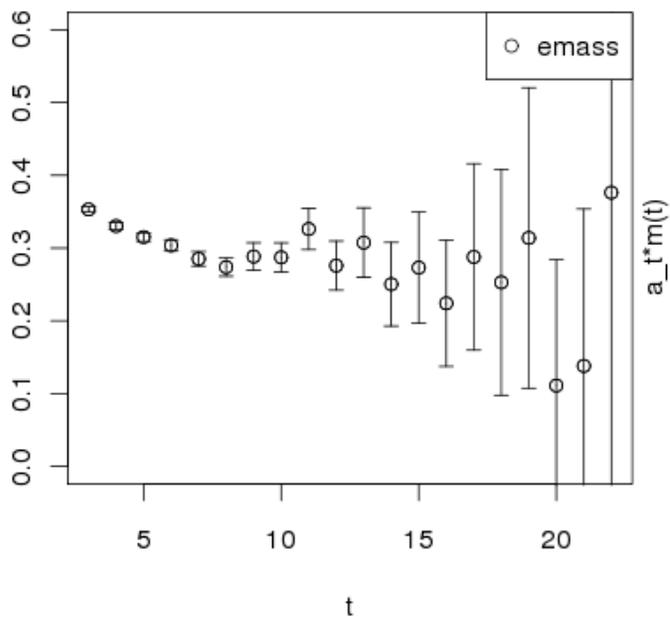


# A1m <1,1,1>

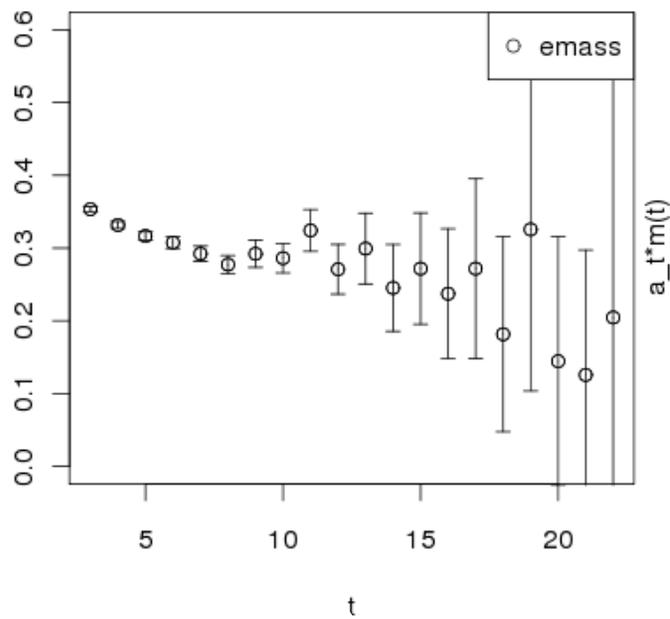


# B2p <1,1,1>

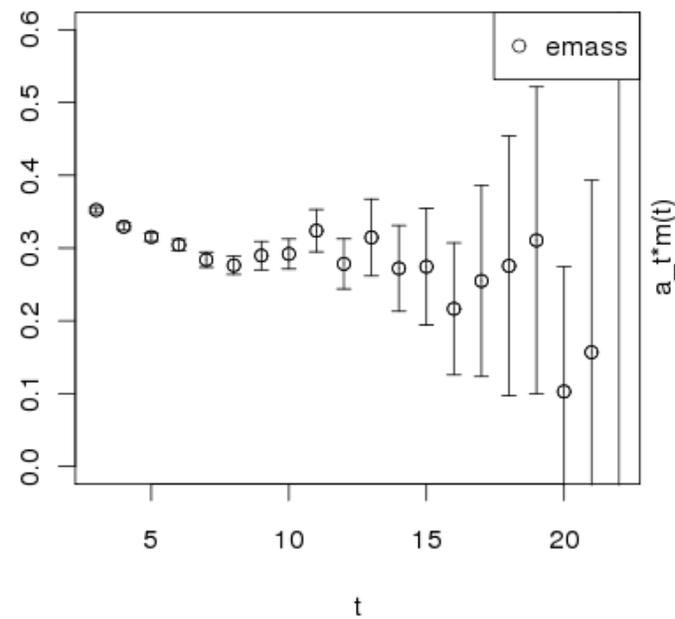
## SS\_0



## LSD\_2



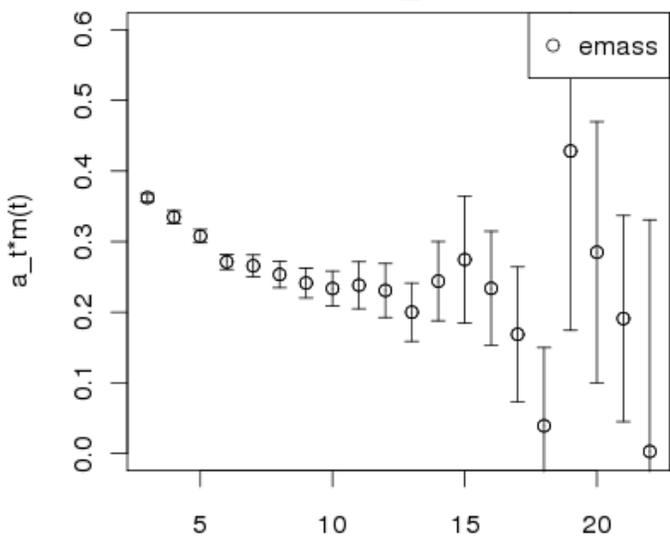
## LSD\_3



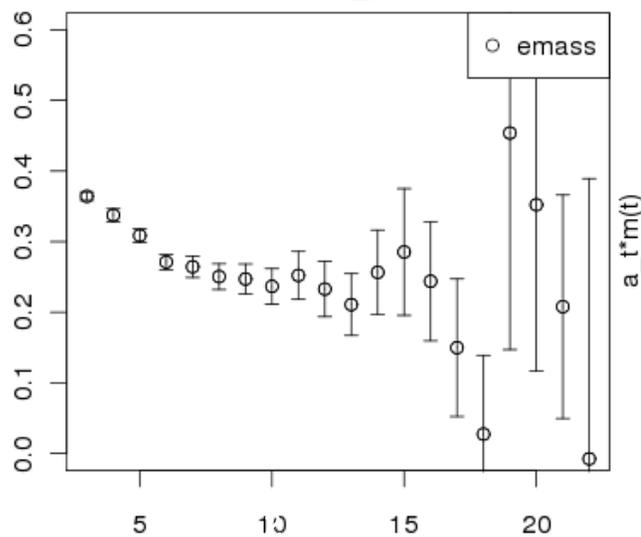
# A2m <0,2,2>

and more ...

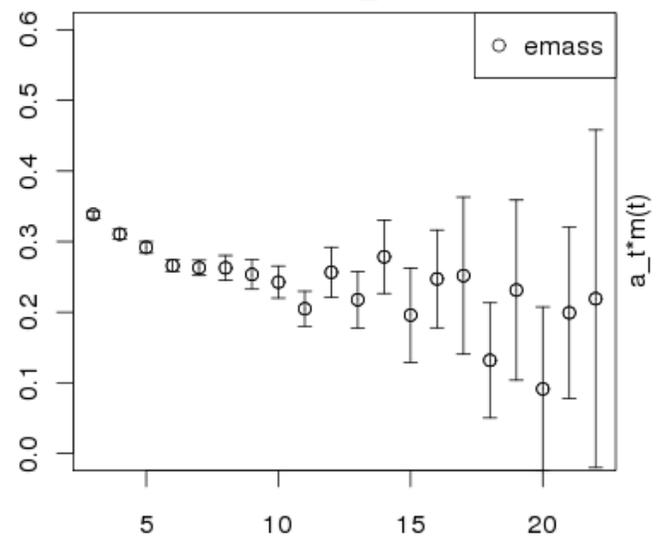
## SS\_0



## LSD\_0



## TSD\_1

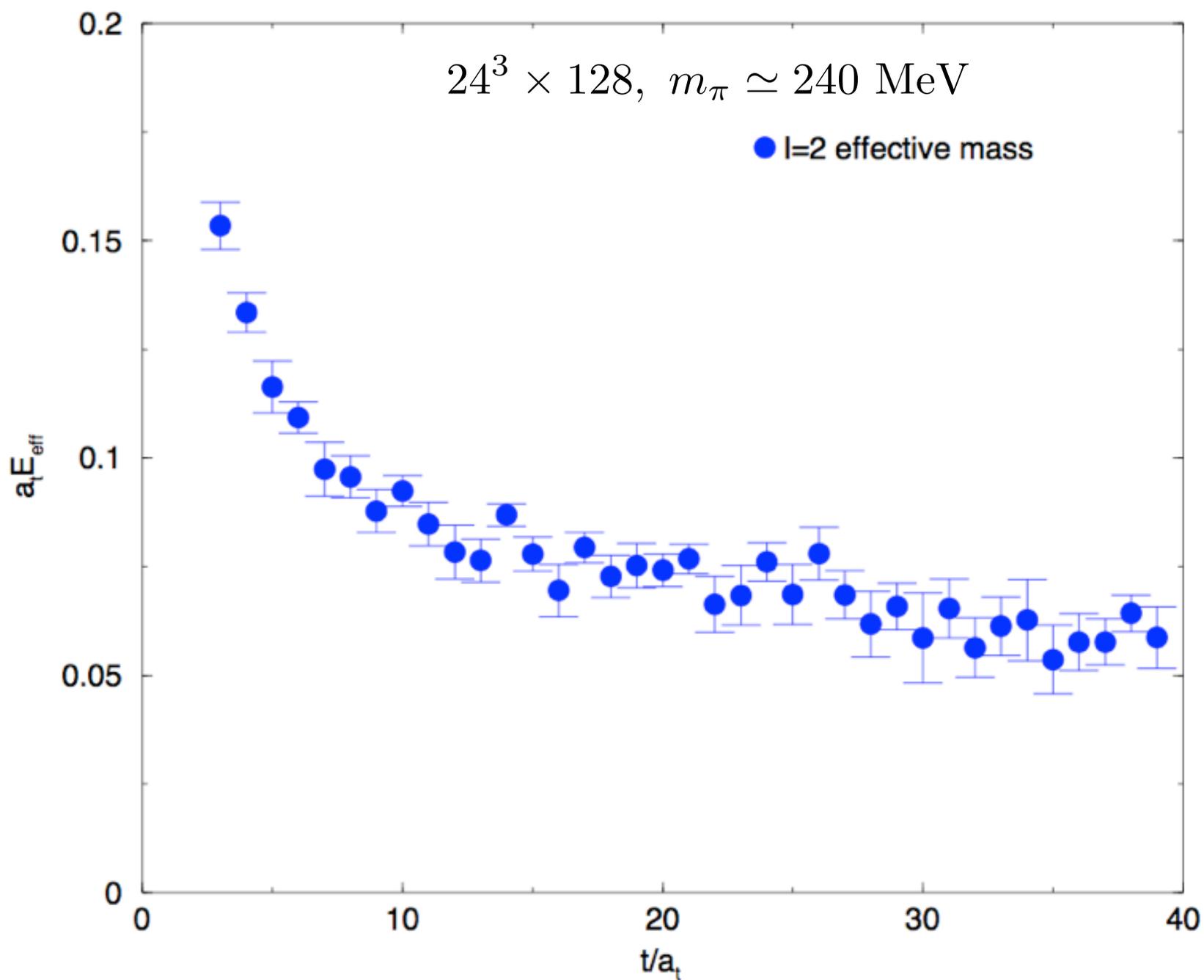


Multi-particle operators constructed from  
simple elemental operators with finite momentum

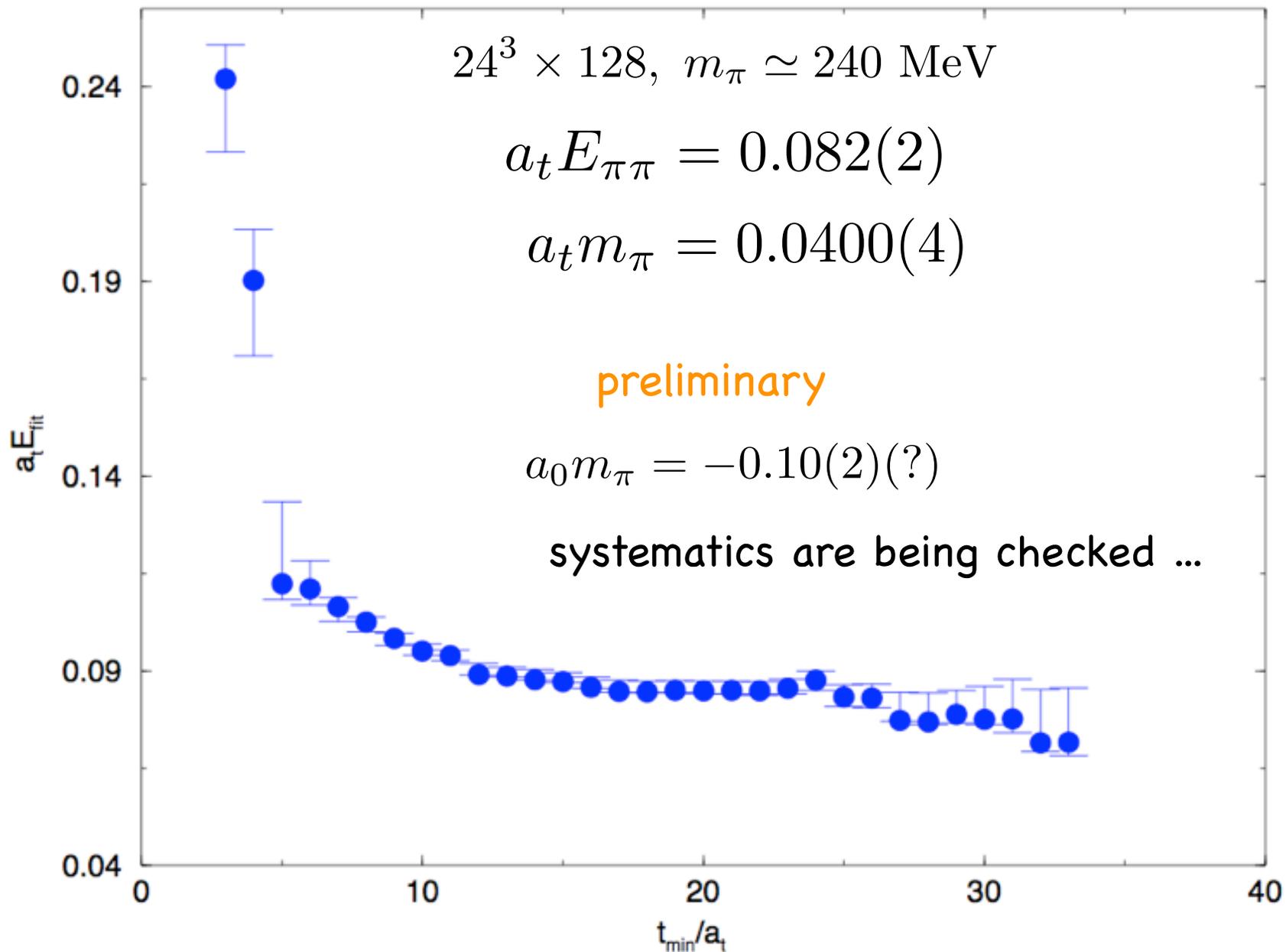
“Simplest” two-particle state

→ two-pion state

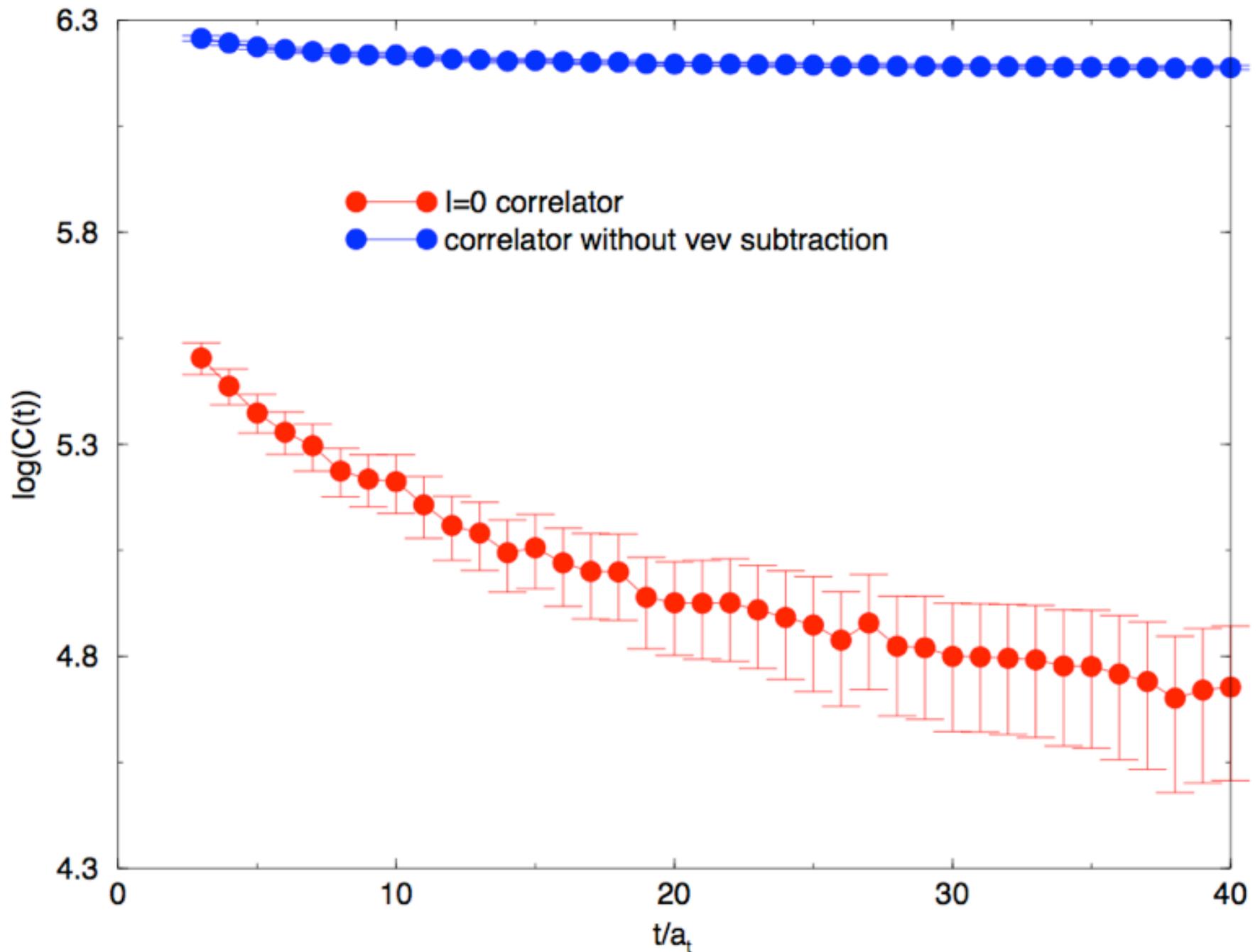
# I=2 Effective Mass



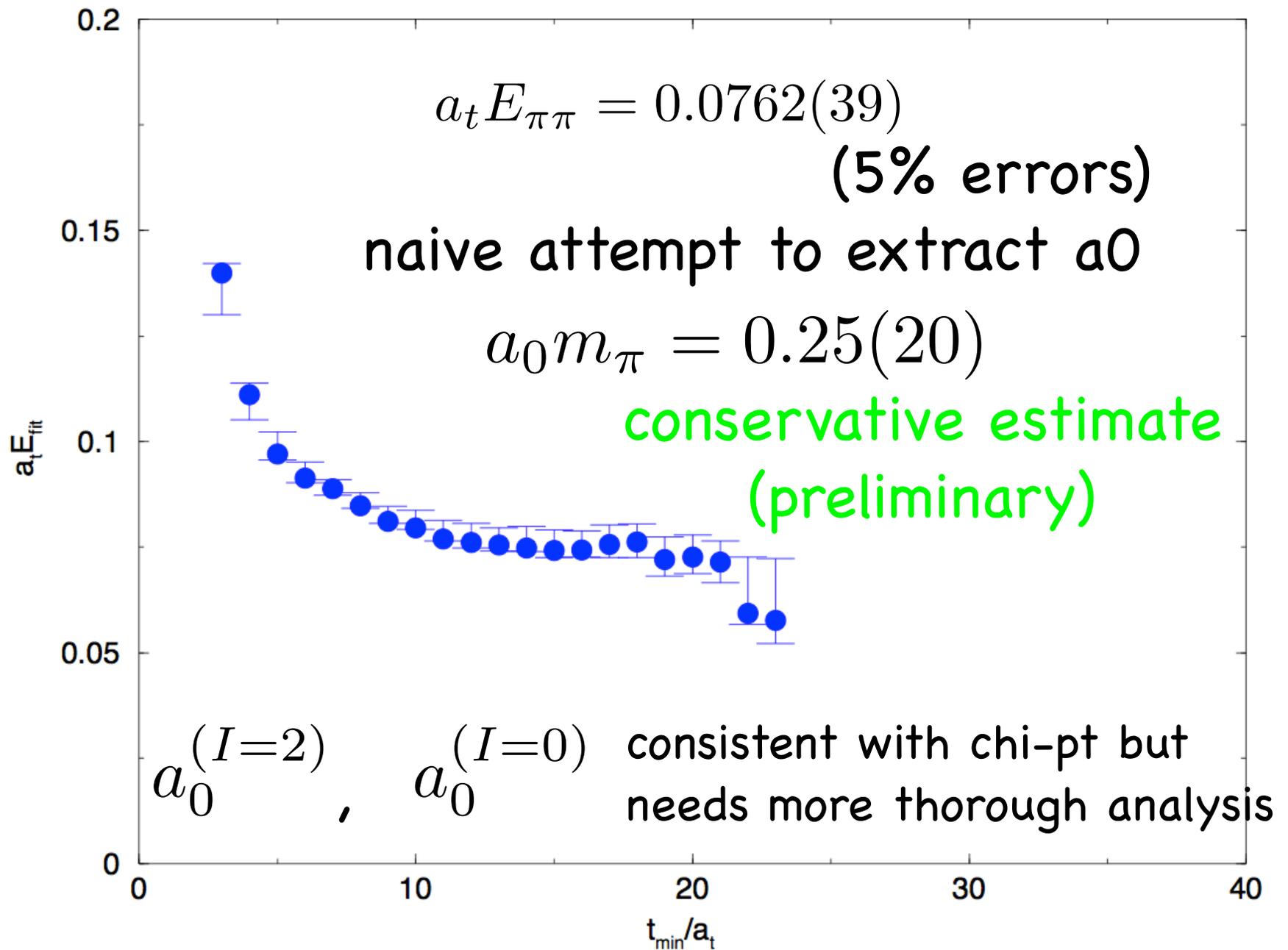
# I=2 tmin plot

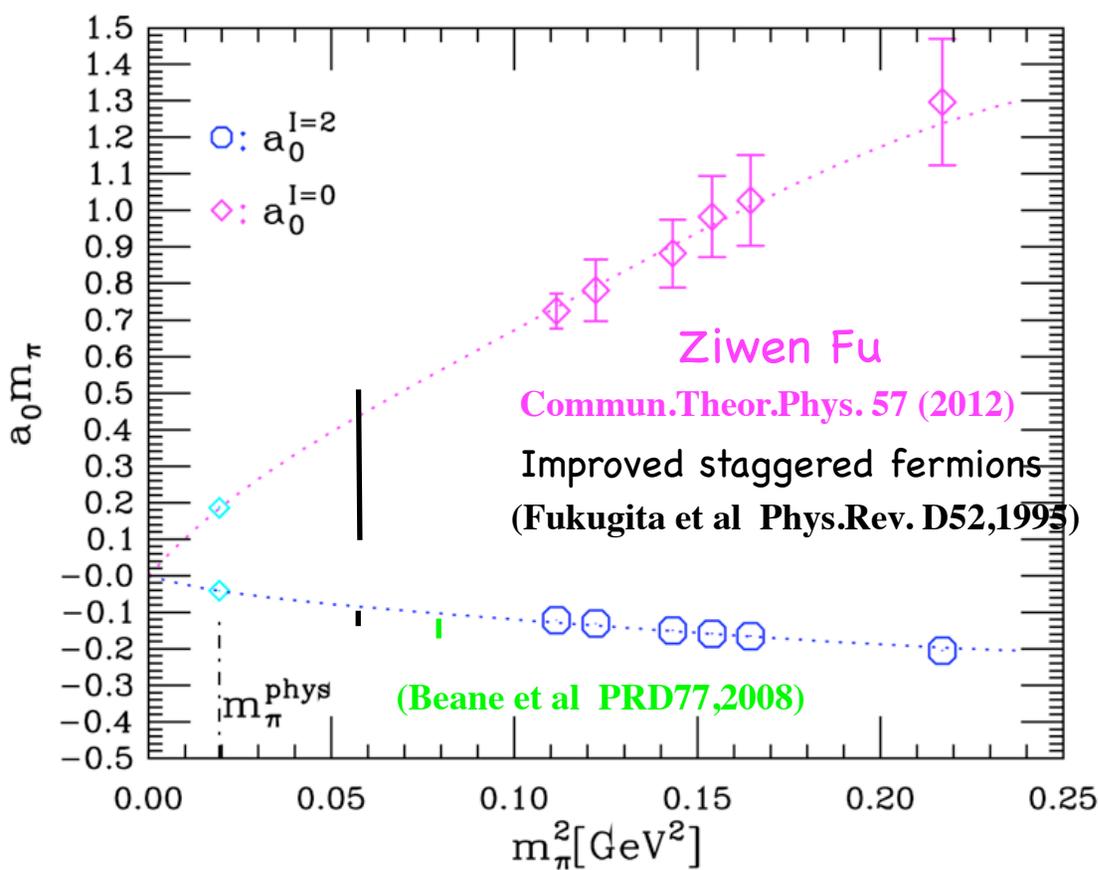


# I=0 pi-pi correlator



# I=0 tmin plot

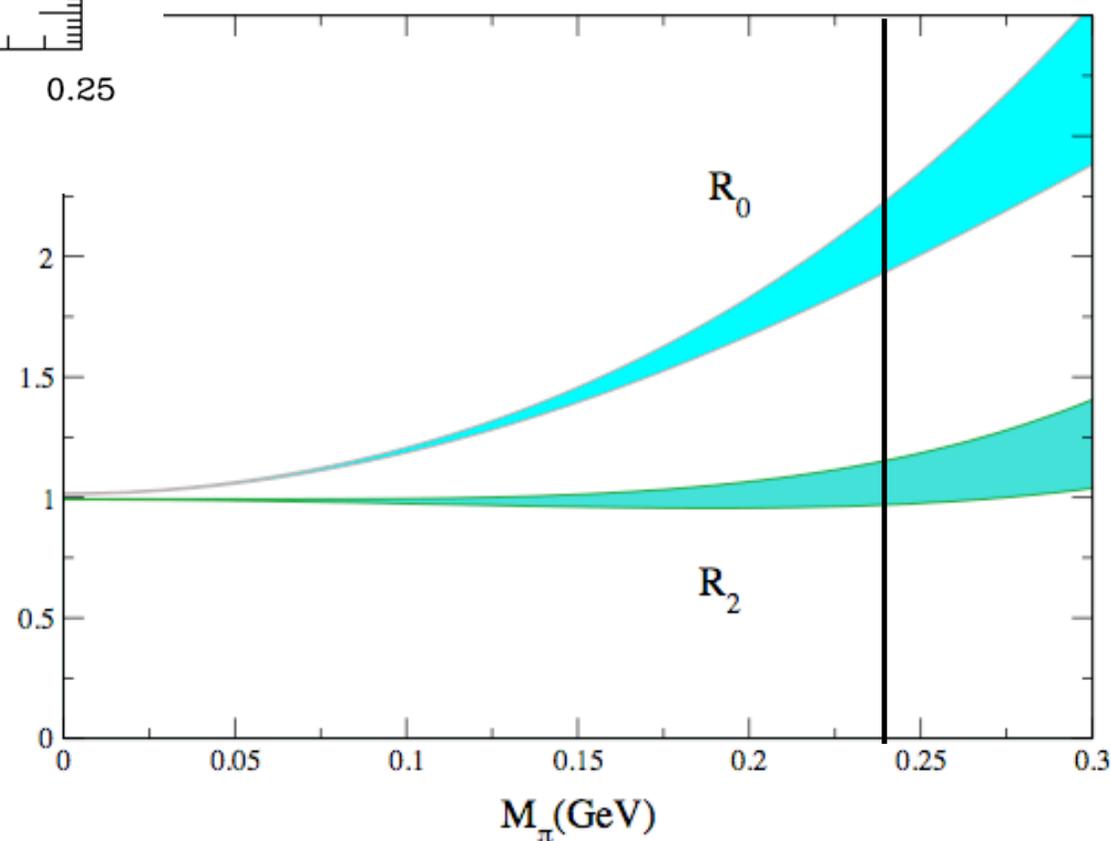




Fully correlated analysis underway with the new value of the renormalized anisotropy ...

$$a_0^{(I=0)} = \frac{7m_\pi^2}{32\pi f_\pi^2} R_0$$

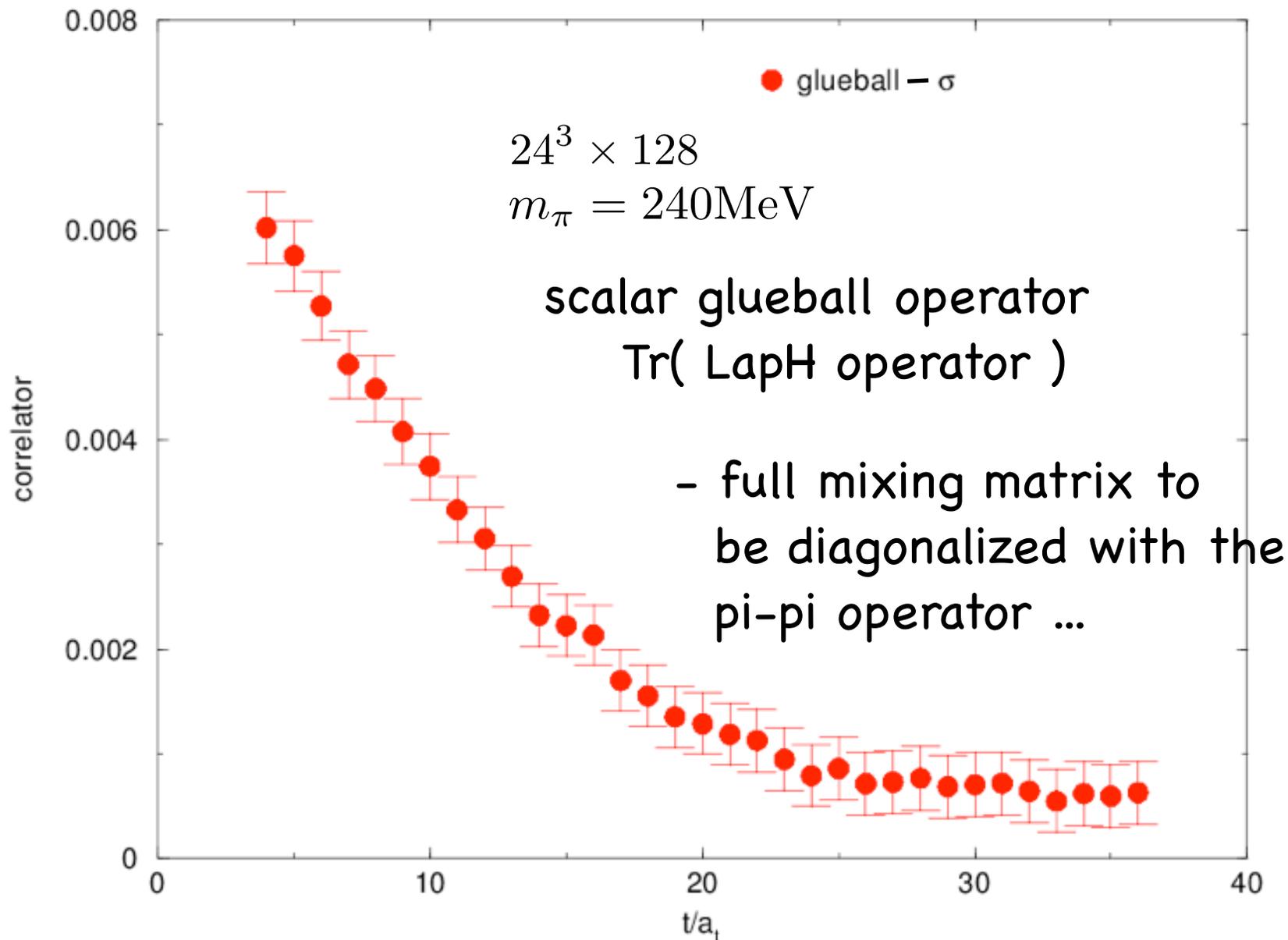
$$a_0^{(I=2)} = -\frac{m_\pi^2}{16\pi f_\pi^2} R_2$$



Other channels:

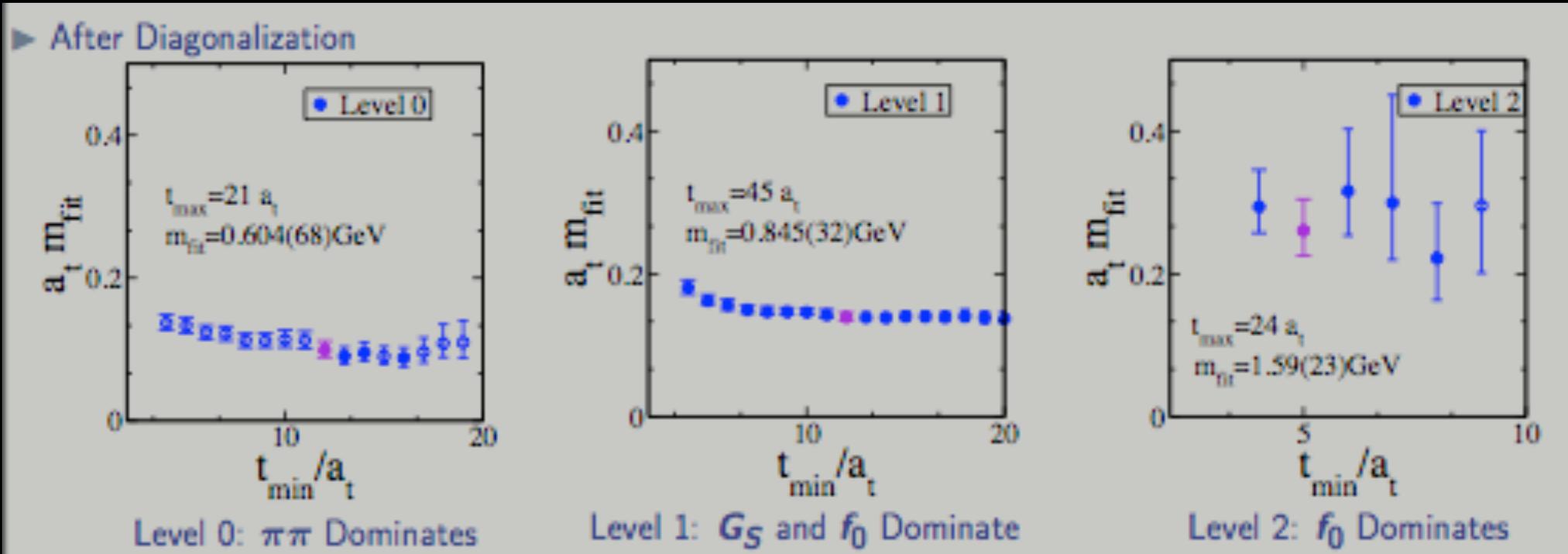
$0^{++}$

scalar glueball mixing



$0^{++}$

$16^3$  lattice with 380 MeV pions



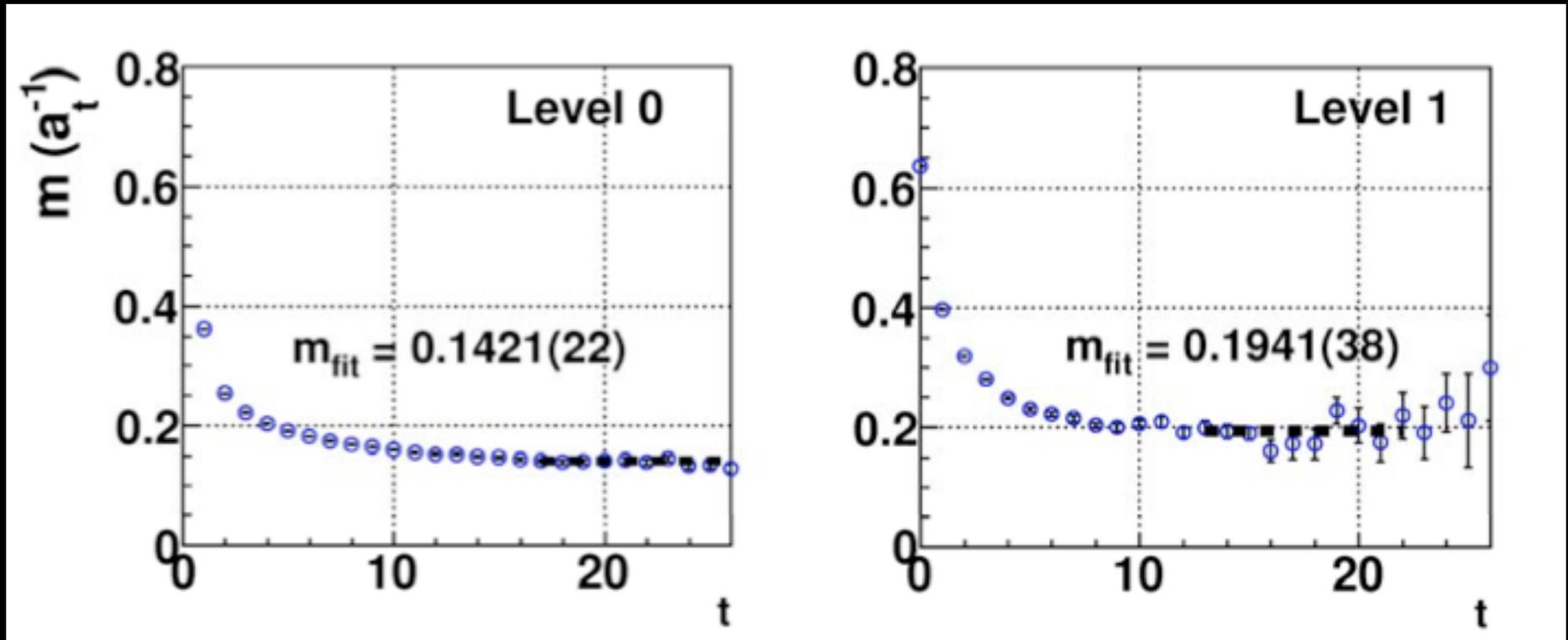
Ricky Wong  
(Lattice2012)

similar analysis on the  $24^3$  to follow ...

1--

$$\rho \rightarrow \pi\pi$$

diagonalize correlation matrix with single rho operator and pi-pi operator



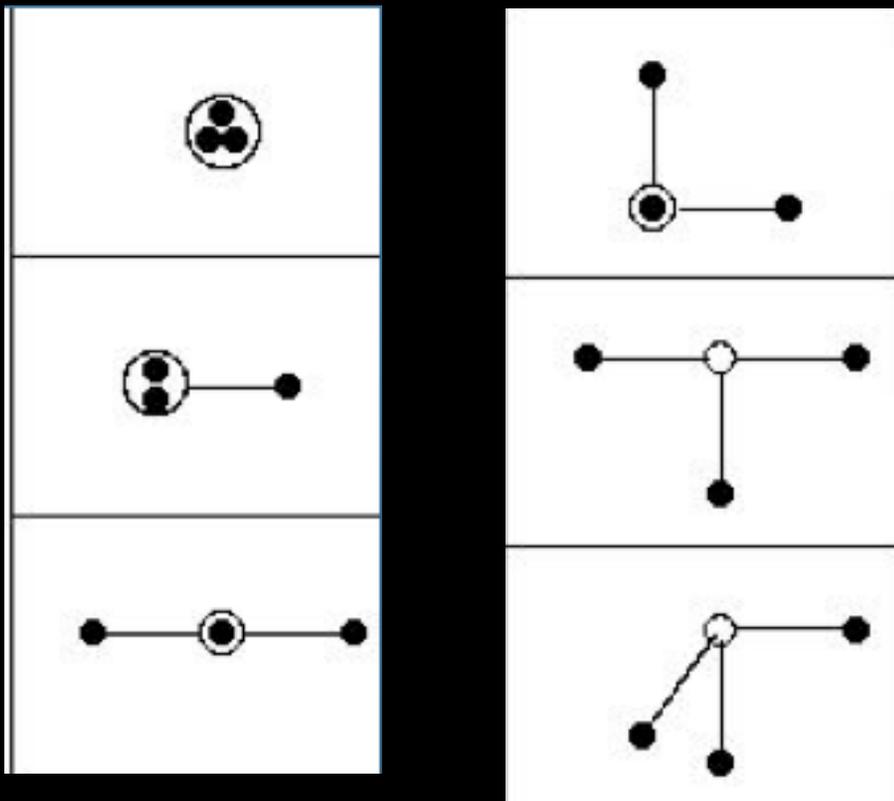
$16^3$  lattice with 380 MeV pions  
Similar analysis on the  $24^3$  lattice  
with 240 MeV pions being done ...

# Baryons

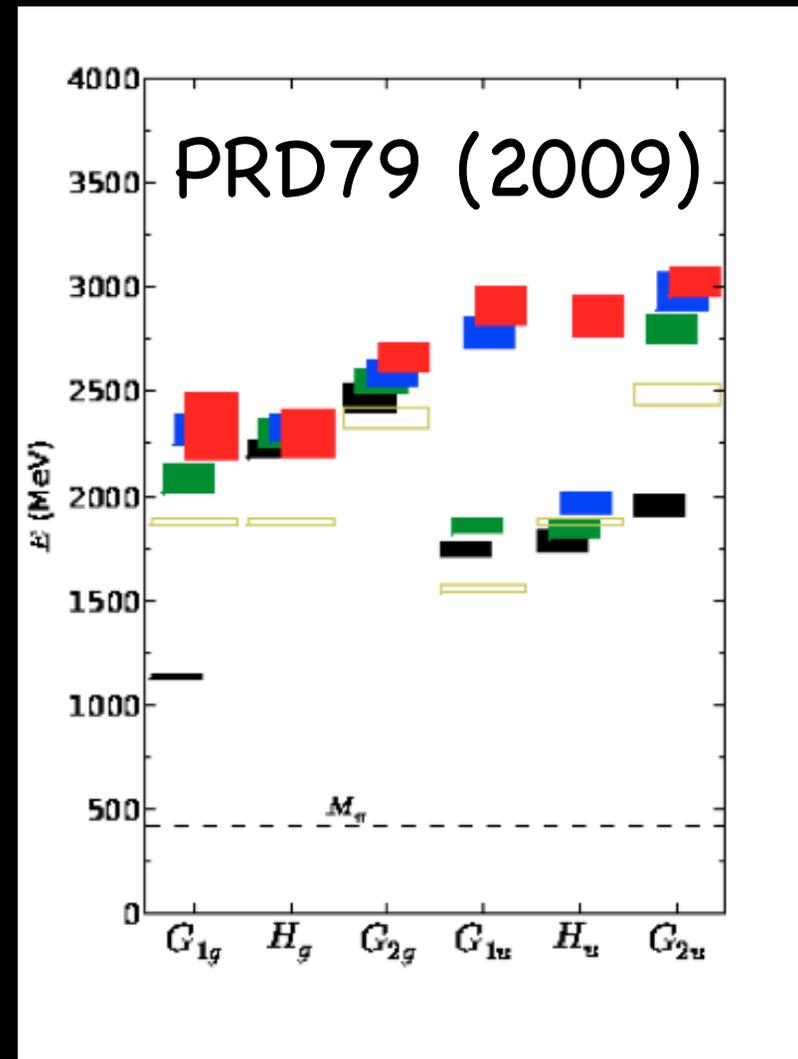
Pruning of the single baryon operator done on the  $16^3$  volume and heavy pions 380 MeV

Basic building blocks

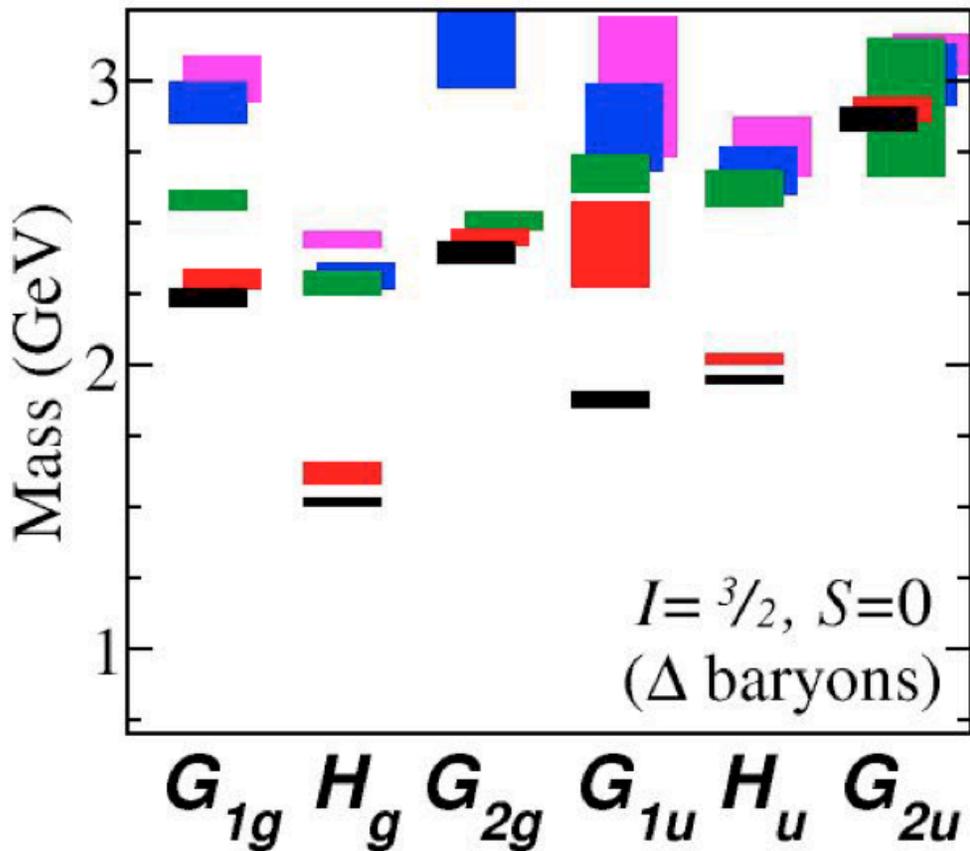
Nucleons done on  $24^3$



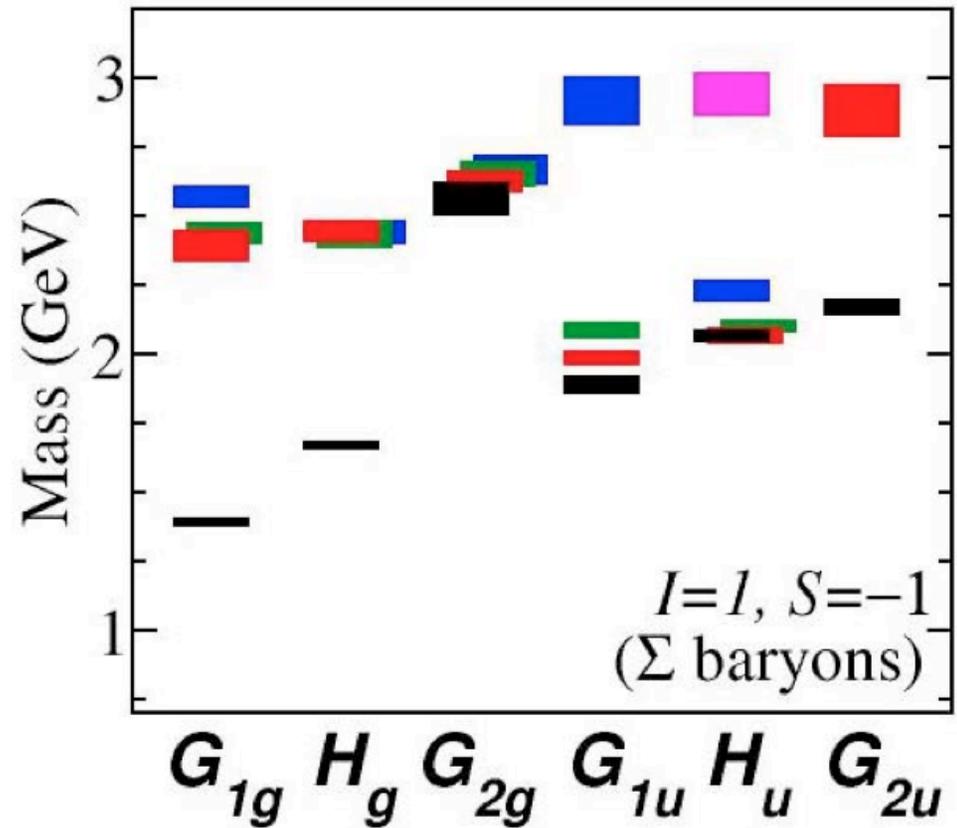
Group theoretical projections  
Basak et al. PRD72 (2005)



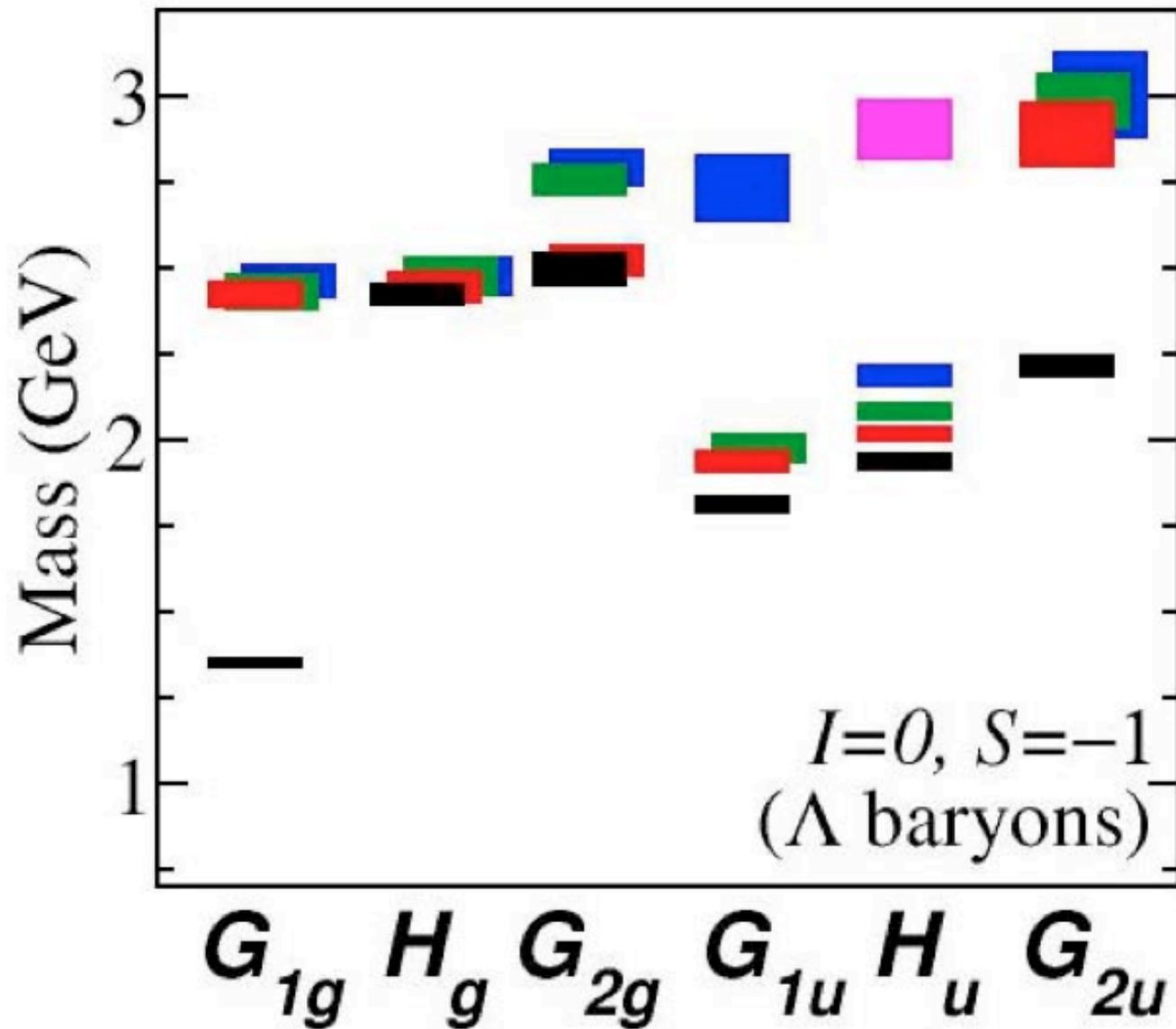
**Delta baryons**  
( $16^3$ , 380MeV)



**Sigma baryons**  
( $16^3$ , 380MeV)



# Lambda baryons ( $16^3$ , 380MeV)



# Conclusions

- Steady progress towards the dynamical simulation of the low-lying hadron spectrum
  - LapH eigenvectors and solutions computed on the  $24^3$ ,  $32^3$  lattices at the 240 MeV mass
  - renormalized anisotropy being finalized
  - single particle meson operators and correlation functions computed
    - isovector operators, finite momentum operators computed (spectrum computed on small lattices)
    - coefficients to construct two-particle states done

- two-particle correlation functions computed
  - $I=0$ , 2 final analysis is being finalized with more careful look at finite size effects and operator overlaps
  - $I=1$  coming soon
- glueball mixings on the large lattices to be completed with sigma, pi-pi.

- **Baryons** (Nucleon/Delta) to be done on the  $24^3$  lattices
  - operator pruning done for Nucleons, Delta, Sigma, Lambda on  $16^3$  lattices with 380 MeV pions
- inversions on the  $32^3$  lattices at the 240 MeV pion masses are done so that the spectrum can be computed on our (3.6fm) box

We are almost there to finally compute the optimized hadron lattice operators of the low-lying hadron spectrum (at  $m_{\pi}=240$  MeV)  
Phenomenology with these “good” operators will be the next step ...