

# Correlation and Flow:

Do we understand them as well as we claim?

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Pre-workshop:

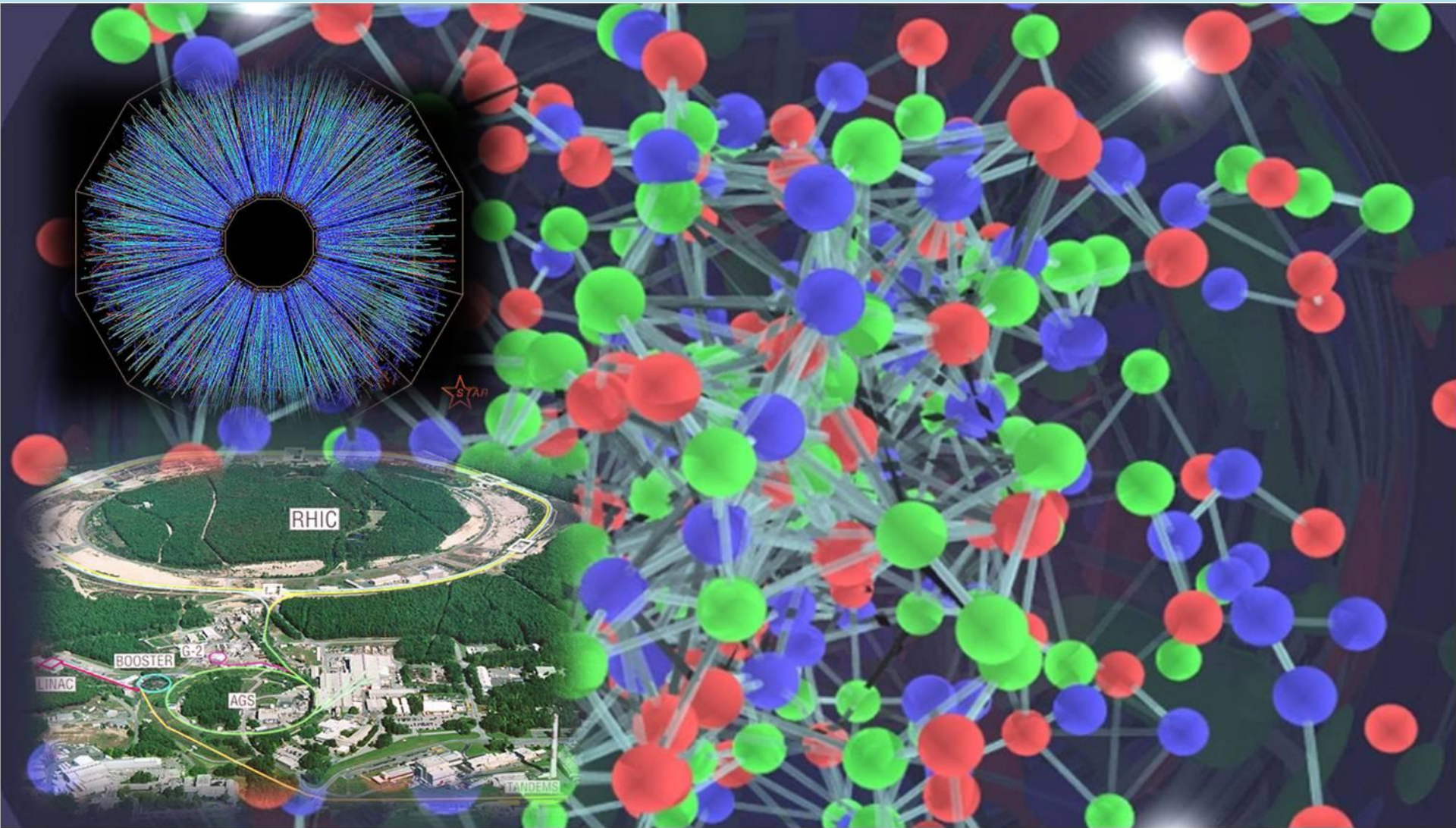


<http://conf.cmu.edu>





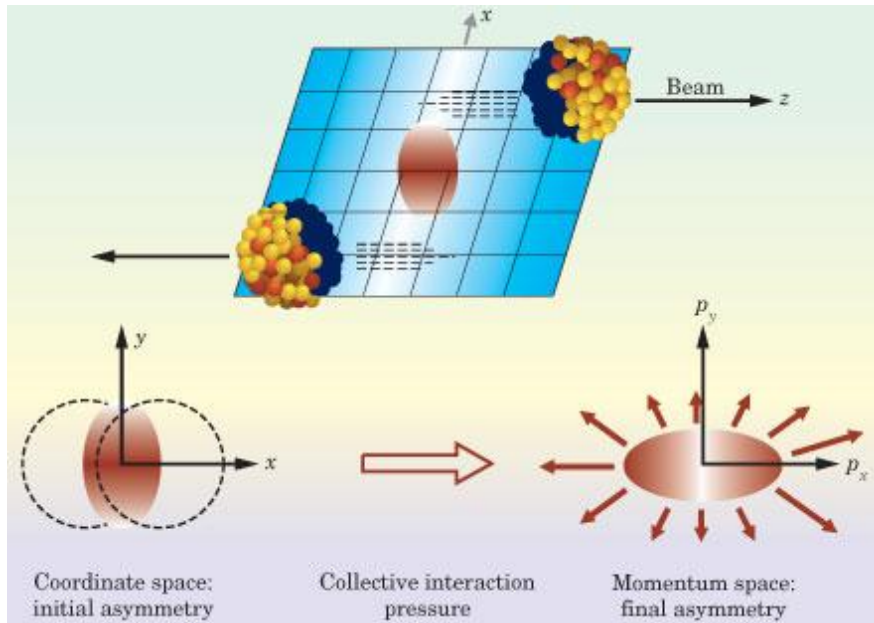
# Nearly Perfect Fluid @ RHIC





# Two main tools

## FLOW



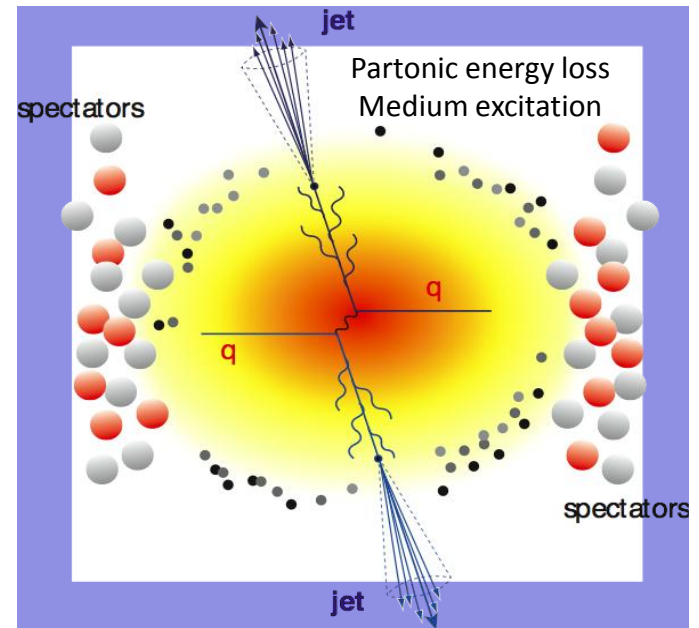
Hydrodynamics

Pressure, energy density, Equilibrium

Probe Equation of State

Low  $p_T$

## JET QUENCHING



Calibrated probe in pp

Jet-medium interactions, Energy loss

Probe QCD medium properties

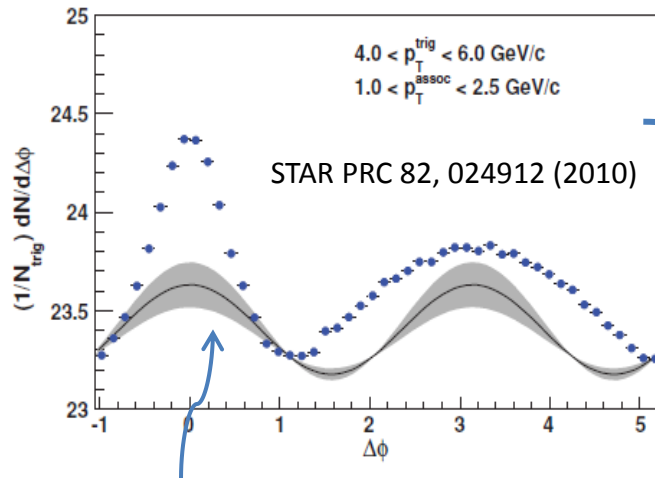
High  $p_T$



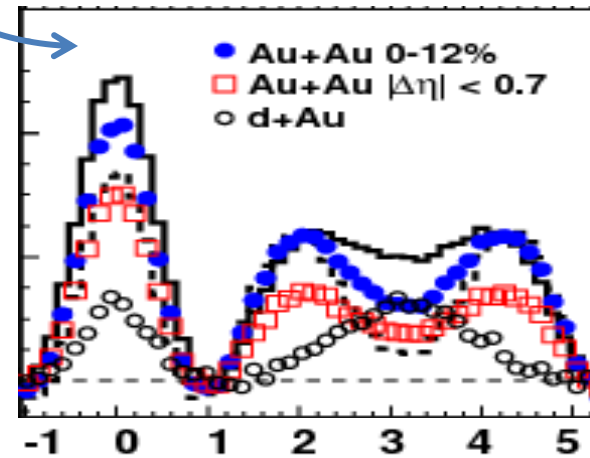
# Flows are background to jet-like correlations

Two-particle correlations contain both flow and nonflow

$$\text{CORREL SIGNAL} + \text{CORREL BKGD} = \text{NONFLOW} + \text{FLOW}$$

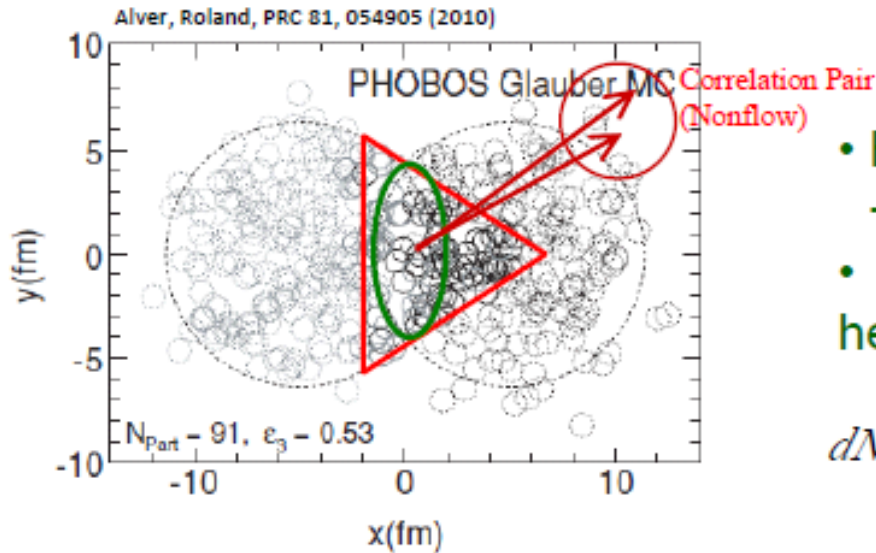


$v_2$  background subtraction



- Flow = background
- Nonflow = signal

# Triangular and higher order harmonics



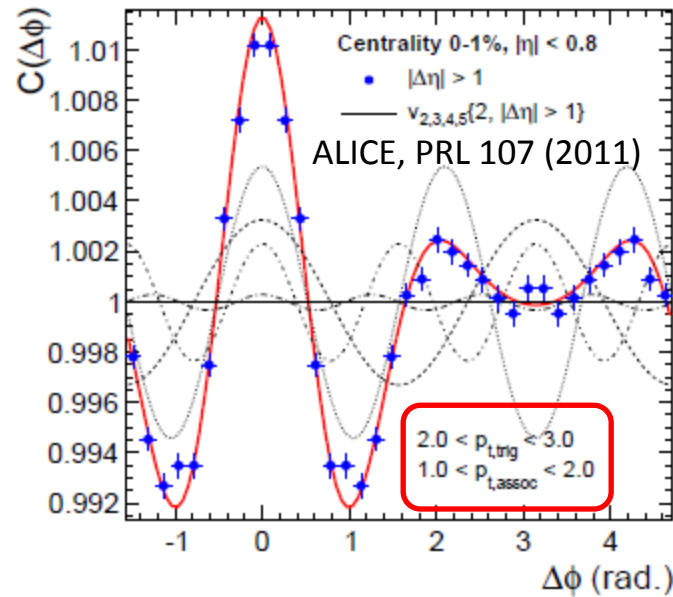
- Hydrodynamic expansion  
→ anisotropic flow;
- Flow is sensitive to early stage of heavy ions collisions

$$dN/d\varphi \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos(n(\varphi - \Psi_{nR}))$$

- Event-by-event initial geometry fluctuation  
→ odd harmonics
- The reaction plane azimuthal angle is unknown  
→ the measured anisotropies = flow( $v$ ) + flow fluctuation ( $\sigma$ ) + nonflow ( $\delta$ )

particle correlation unrelated to the reaction plane

# Some say: All are $v_n$ flow



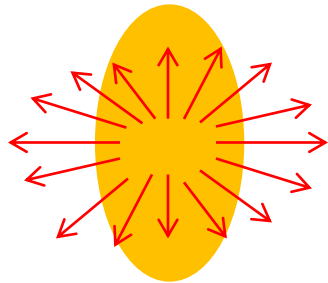
$p_T^{\text{ref}} = 0.2-5 \text{ GeV}/c$

If  $p_T^{\text{assoc}} = p_T^{\text{ref}}$ , then

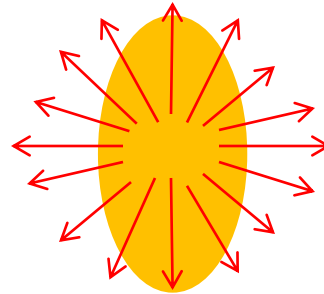
“flow”  $v_n(p_T^{\text{trig}})v_n(p_T^{\text{assoc}}) = \text{trig-assoc dihadron correlation}$   
 $\rightarrow \text{signal} = \text{data} - \text{data} = 0.$

**So, the real question is what's in  $v_n$ .**

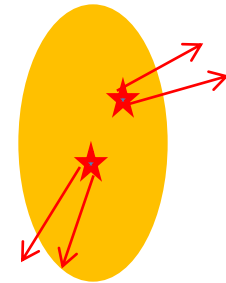
# What is in $v_n$ ?



Hydrodynamic  
pressure driven



Pathlength  
dependent  
energy loss



Nonflow



Flow: related to the geometry

# One support: Flow factorization

$$\langle \cos n(\phi_i - \phi_j) \rangle = \langle \cos n[(\phi_i - \psi_{EP}) - (\phi_j - \psi_{EP})] \rangle = \langle \cos n(\phi_i - \psi_{EP}) \rangle \langle \cos n(\phi_j - \psi_{EP}) \rangle$$

$$V_{n\Delta}(p_T^a, p_T^b) = v_n(p_T^a) v_n(p_T^b) \quad \leftarrow \text{Factorization}$$

$$v_n(p_T) = \sqrt{V_{n\Delta}(p_T, p_T)}$$

$$v_n(p_T) = \frac{V_{n\Delta}(p_T, p_T^{Ref})}{v_n(p_T^{Ref})}$$

Factorization expected due to global correlation with the reaction plane

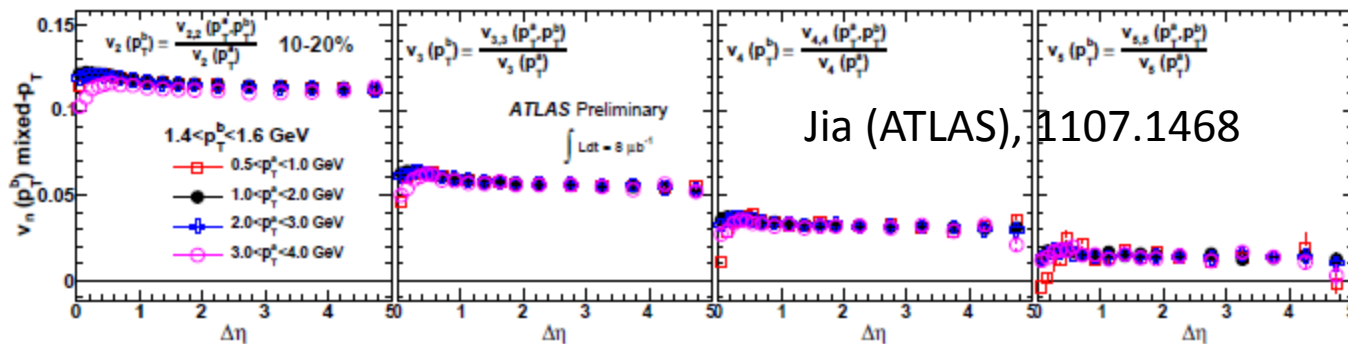
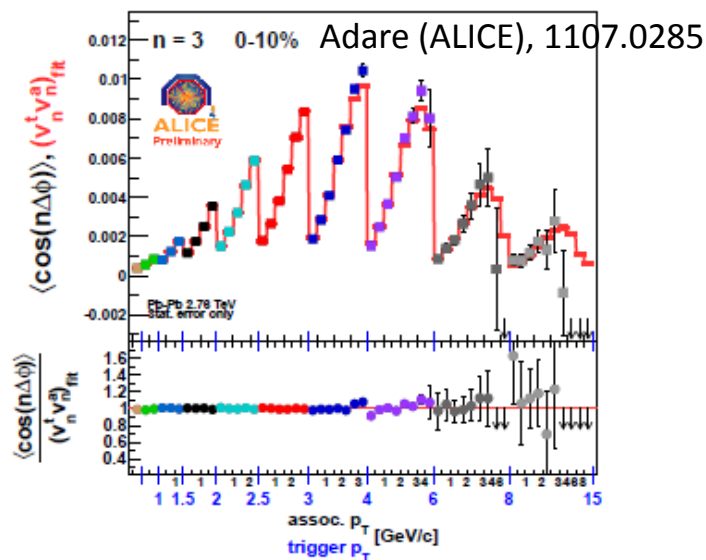
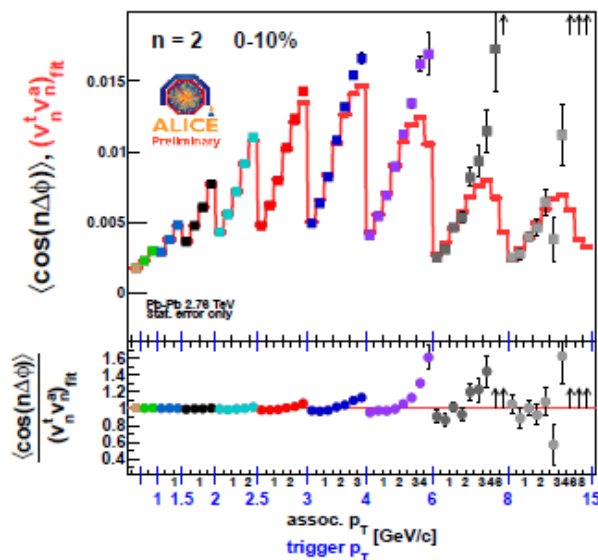
## Caveat:

- Due to fluctuations,  $V_{n\Delta}$  does not factorize in general
- Precise factorization if  $\sigma_n \propto \langle v_n \rangle$



# $v_n$ factorization

$$V_{n\Delta}(p_T^a, p_T^b) = v_n(p_T^a) v_n(p_T^b)$$



# Nonflow factorization?

Common perception:

Nonflow should not factorize

However:

**Independent jet fragmentation (except the common jet thrust axis)**

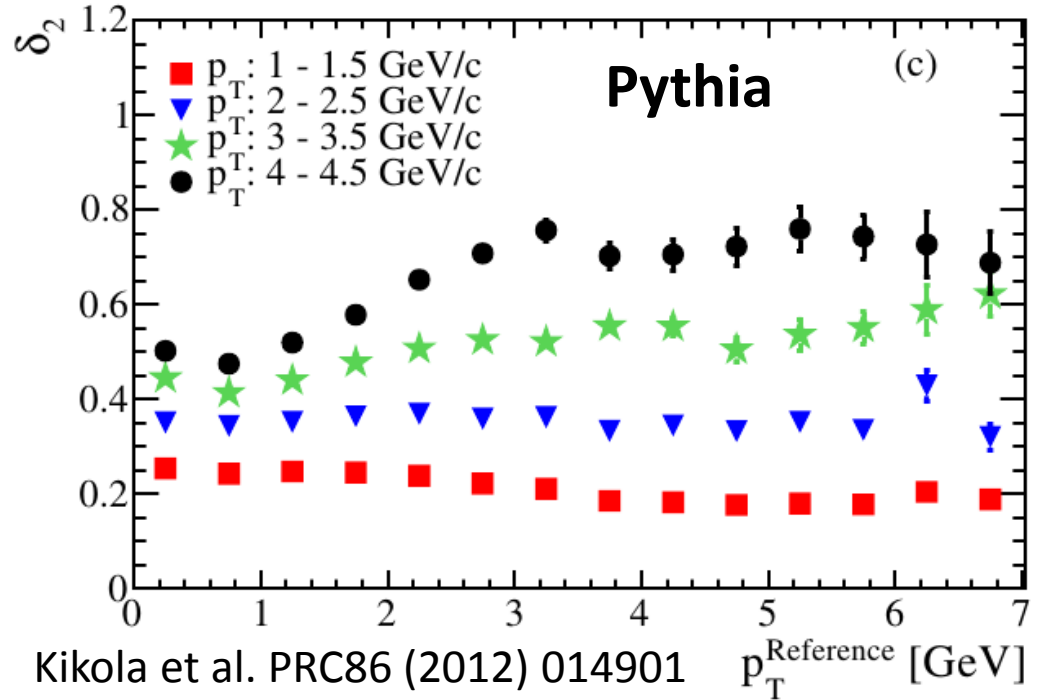
$$\langle \cos n(\phi_i - \phi_j) \rangle = \langle \cos n[(\phi_i - \psi_{jet}) - (\phi_j - \psi_{jet})] \rangle = \langle \cos n(\phi_i - \psi_{jet}) \rangle \langle \cos n(\phi_j - \psi_{jet}) \rangle$$

→ Jet correlation may approximately factorize !

# Nonflow approximate factorization

$$\delta_2(p_T^{\text{Ref}}) = \sqrt{V_{2\Delta}(p_T^{\text{Ref}}, p_T^{\text{Ref}})}$$

$$\delta_2(p_T) = \frac{V_{2\Delta}(p_T, p_T^{\text{Ref}})}{\delta_2(p_T^{\text{Ref}})}$$



$\delta_2(p_T^{\text{Ref}})$  approximately independent on  $p_T$

→ Nonflow approximately factorizes in a limited  $p_T$  range

→ factorization is not unique feature of flow



# Why is 'flow' factorization so good? Because it is bootstrapped!

Anisotropic flow + non-flow:  $V_{2\Delta}(p_T^a, p_T^b) = v_2(p_T^a)v_2(p_T^b) + \delta_2(p_T^a)\delta_2(p_T^b)$

$$\begin{aligned} & \frac{V_{n\Delta}(p_T^b, p_T^a)}{v_n'(p_T^a)v_n'(p_T^b)} - 1 \\ &= \frac{v_n(p_T^b)v_n(p_T^c) + \delta_n(p_T^b)\delta_n(p_T^c)}{\frac{v_n(p_T^a)v_n(p_T^c) + \delta_n(p_T^a)\delta_n(p_T^c)}{\sqrt{v_n^2(p_T^c) + \delta_n^2(p_T^c)}} \frac{v_n(p_T^b)v_n(p_T^c) + \delta_n(p_T^b)\delta_n(p_T^c)}{\sqrt{v_n^2(p_T^c) + \delta_n^2(p_T^c)}}} - 1 \\ &\approx \left( \frac{\delta_n(p_T^a)}{v_n(p_T^a)} - \frac{\delta_n(p_T^c)}{v_n(p_T^c)} \right) \left( \frac{\delta_n(p_T^b)}{v_n(p_T^b)} - \frac{\delta_n(p_T^c)}{v_n(p_T^c)} \right) \quad (14) \end{aligned}$$

- $\delta_n(p_T)/v_n(p_T) \sim 10\% \rightarrow \text{deviation} \sim 10^{-3}$
- $\delta_n(p_T) \propto v_n(p_T) \rightarrow \text{precise factorization even if nonflow is present}$



# Lessons learned

- Flow  $\rightarrow$  factorization; Factorization  ~~$\rightarrow$~~  flow
- Fourier components do not give further insights.
- We have to separate flow and nonflow in  $v_n$ .





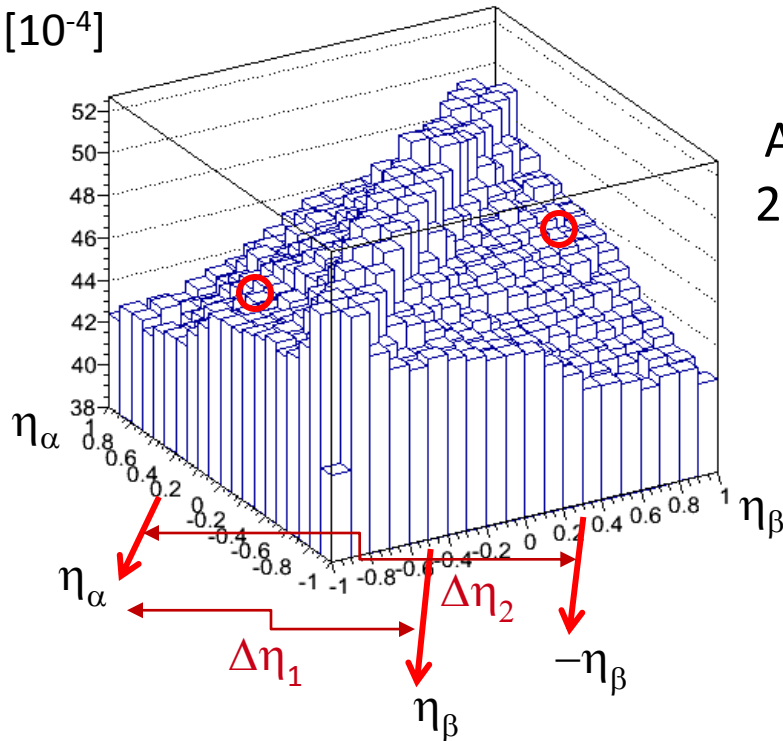
# Separate flow and nonflow

Lingshan Xu et al. PHYSICAL REVIEW C **86**, 024910 (2012)

$$v\{2\}^2 = \langle v_\alpha \rangle \langle v_\beta \rangle + \sigma_\alpha \sigma_\beta + \sigma'(\Delta\eta) + \delta(\Delta\eta)$$

$$v\{4\}^2 \approx \langle v_\alpha \rangle \langle v_\beta \rangle - \sigma_\alpha \sigma_\beta - \sigma'(\Delta\eta)$$

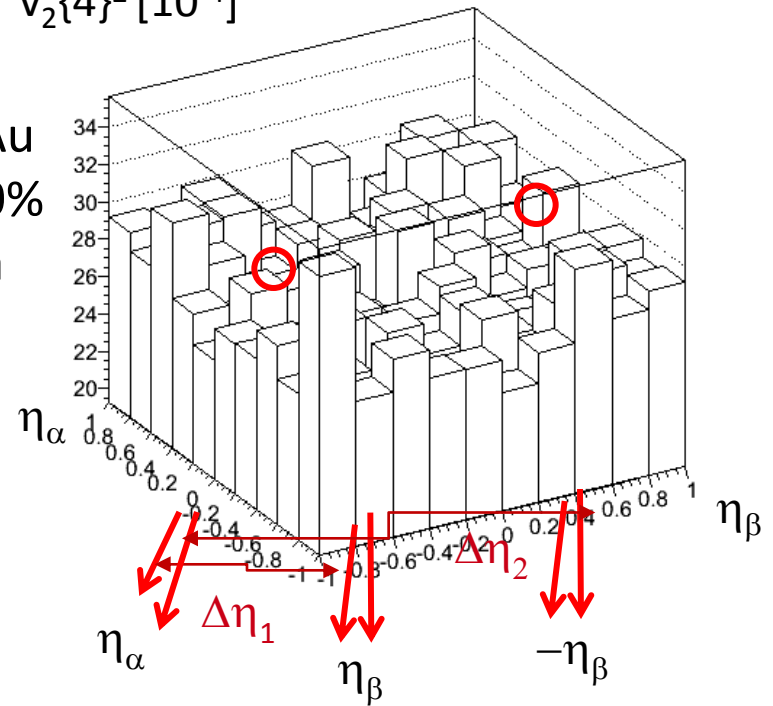
$v_2\{2\}^2 [10^{-4}]$



$\sigma'(\Delta\eta_1) + \delta(\Delta\eta_1)$  vs.  $\sigma'(\Delta\eta_2) + \delta(\Delta\eta_2)$

$v_2\{4\}^2 [10^{-4}]$

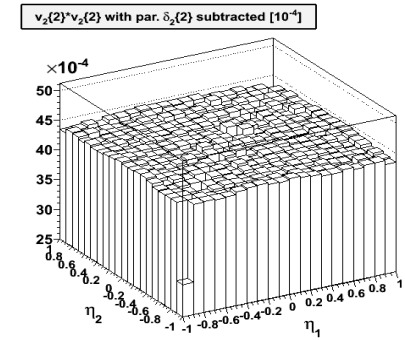
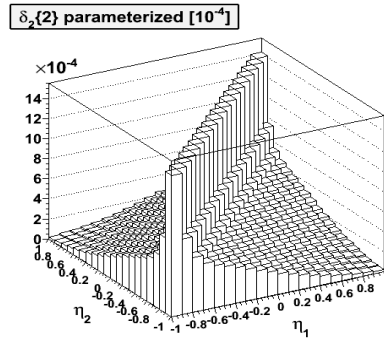
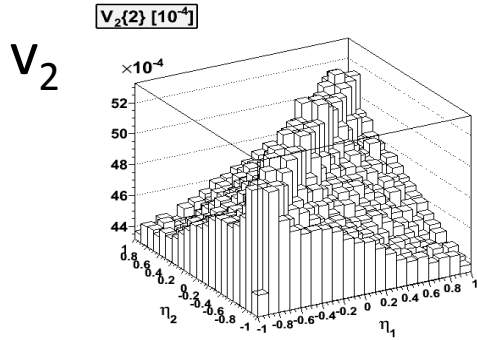
Au+Au  
20-30%  
data



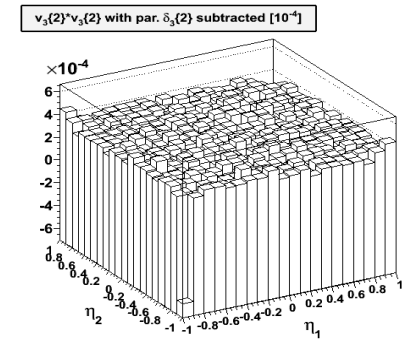
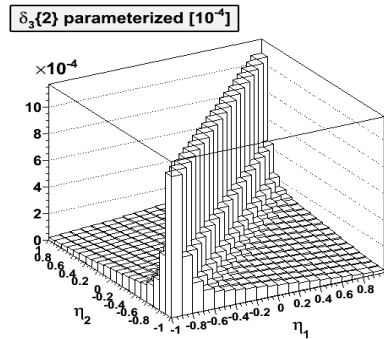
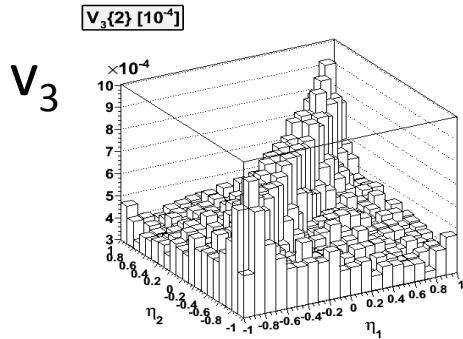
$\sigma'(\Delta\eta_1)$  vs.  $\sigma'(\Delta\eta_2)$



# 2-particle cumulant, nonflow $\delta$ , flow $v_n^2$



STAR preliminary



No Assumption about flow  $\eta$  dependence on our analysis!

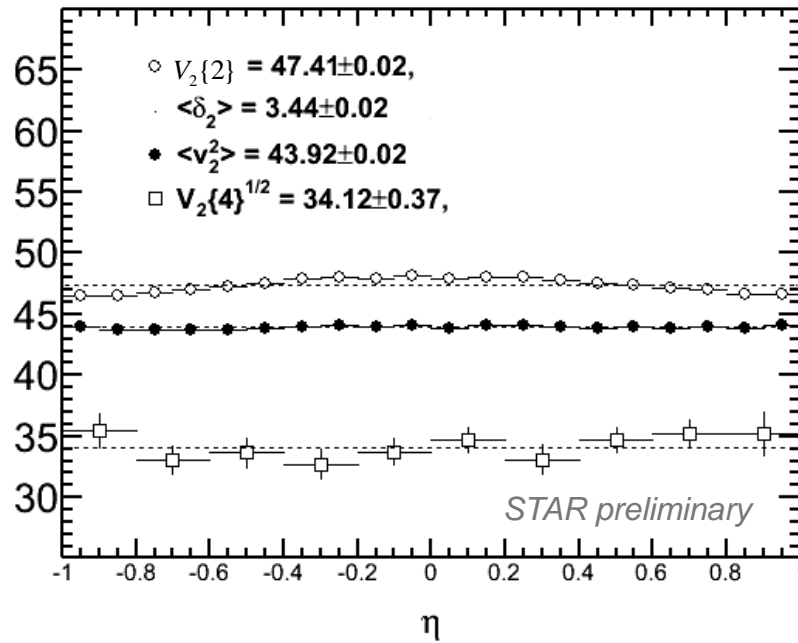
The decomposed 'flow' appears to be independent of  $\eta$ .

Run-4 Au+Au 20-30% data



# Flow vs $\eta$

$\times 10^{-4}$  Run-4 Au+Au 20-30% data

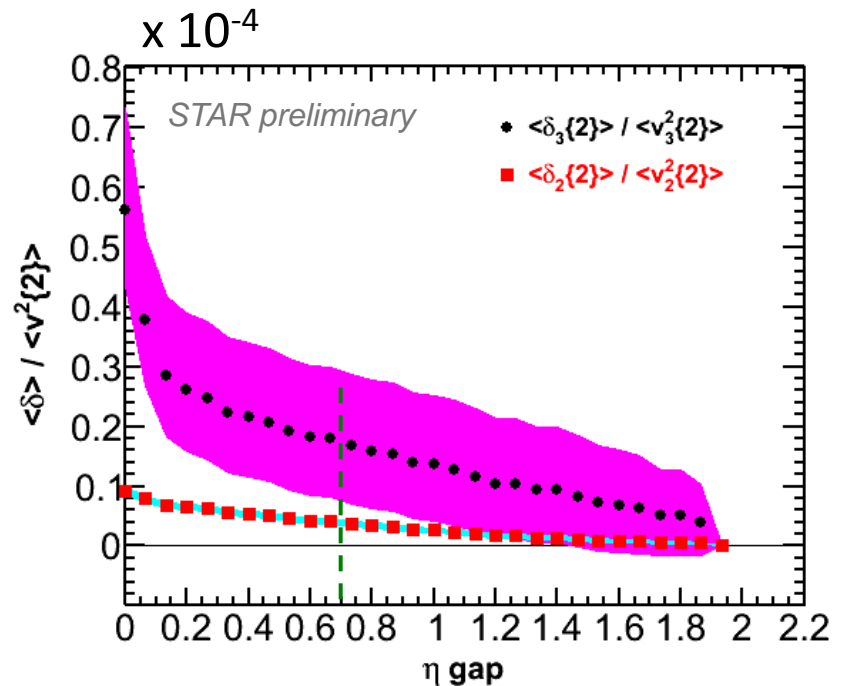
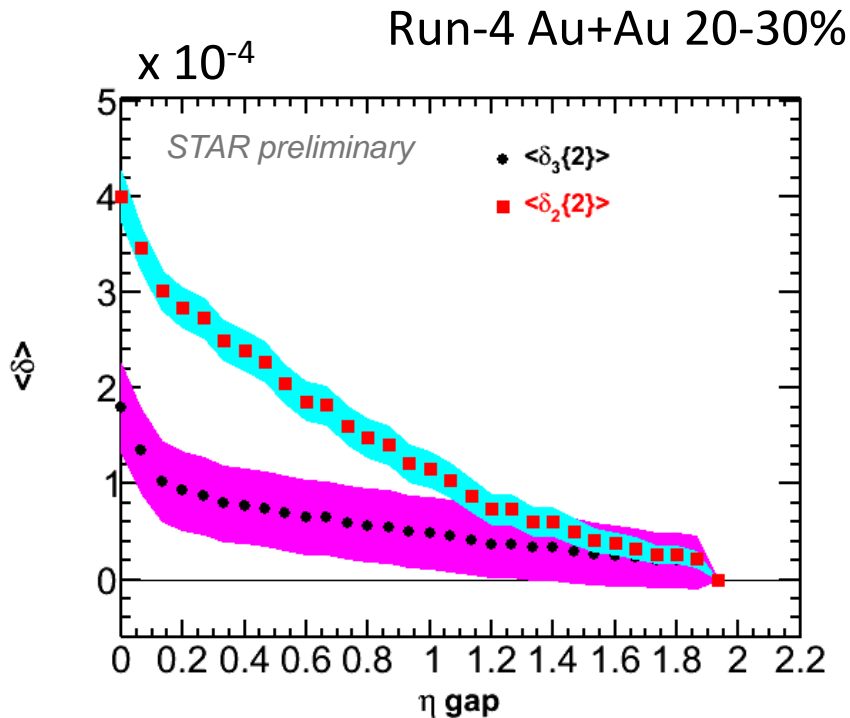


$$\frac{\sigma_2^2}{\langle v_2^2 \rangle^2} \sim \frac{\langle v_2^2 \rangle - V_2\{4\}^{1/2}}{\langle v_2^2 \rangle - V_2\{4\}^{1/2}} \sim 13\%$$

- Flow seems independent of  $\eta$ . Note no assumption of  $\eta$  dependence in our approach.
- Fluctuation / flow  $\sim 36\%$

# Near-side nonflow

- Calculate  $\langle \text{nonflow} \rangle$  of all  $(\eta_1, \eta_2)$  bins with  $x < \eta\text{-gap} < 2$ . ( $x$  = horizontal axis)
- $\Delta\eta > 0$ , nonflow / flow  $\sim 40\%$  for  $v_3$ ,  $10\%$  for  $v_2$
- $\Delta\eta > 0.7$ , nonflow / flow  $\sim 20\%$  for  $v_3$ ,  $5\%$  for  $v_2$





# $\Delta\eta$ -indep. away-side nonflow

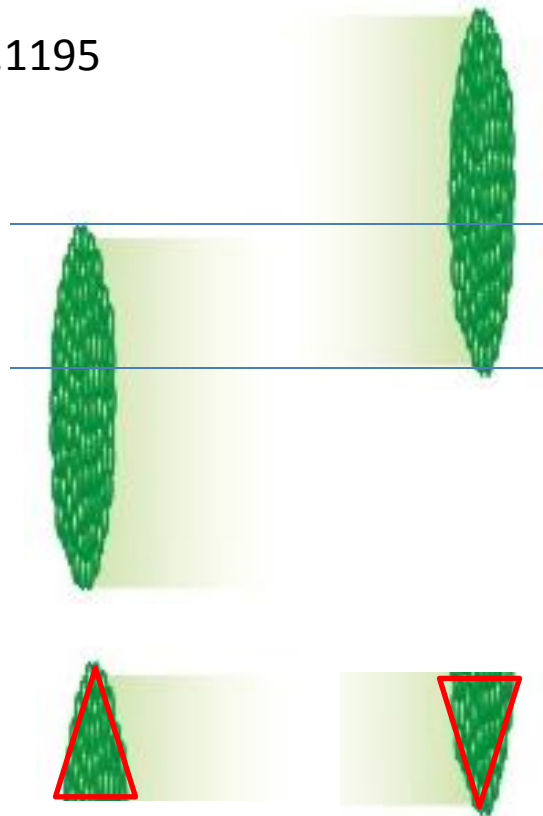
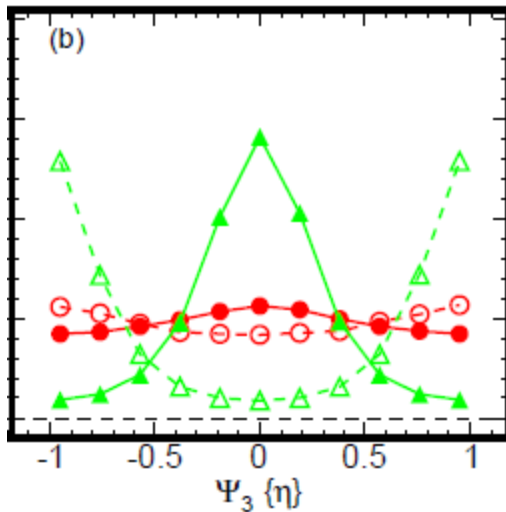
- Hijing: away-side/ near-side = 1 at  $\Delta\eta > 0$
- Assuming the same ratio in data:
  - $\Delta\eta > 0$ :  $\delta_2^{\text{Near Side}} / v_2^2 = 10\%$        $\delta_2^{\text{Away Side}} / v_2^2 = 10\%$
  - $\Delta\eta > 0.7$ :  $\delta_2^{\text{Near Side}} / v_2^2 = 5\%$        $\delta_2^{\text{Away Side}} / v_2^2 = 10\%$

Large enough  $\Delta\eta$  gap, near-side nonflow dies off and the away-side nonflow eventually also dies off. However...

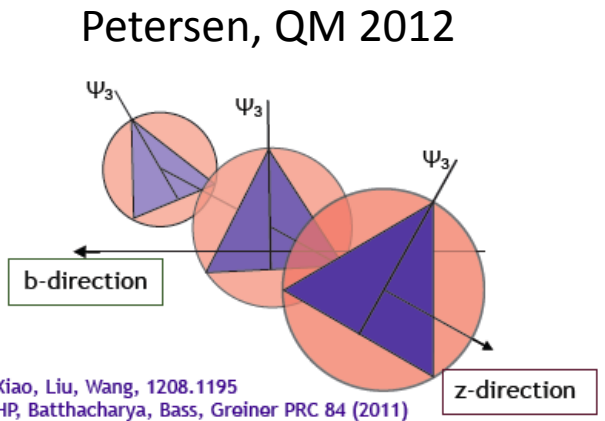


# Large $\Delta\eta$ to reduce nonflow?

Xiao et al. arXiv:1208.1195



Backward rapidity



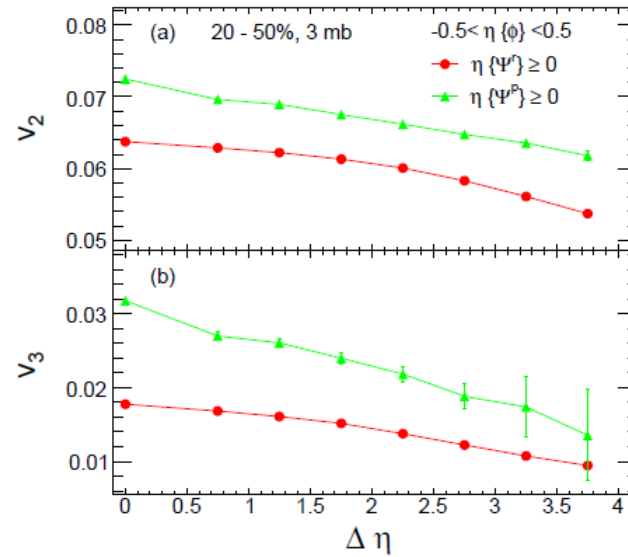
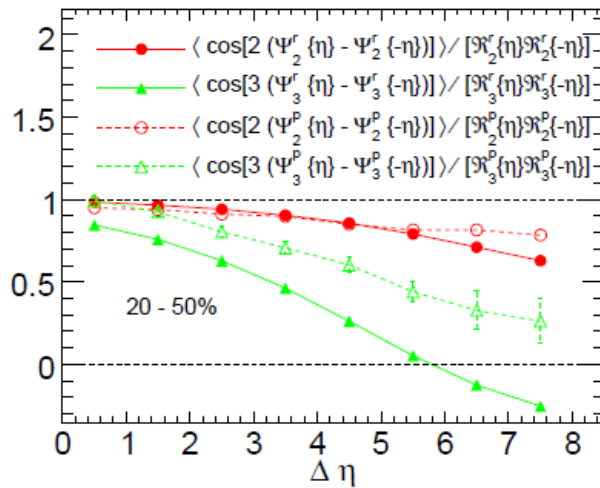
Forward rapidity

Harmonic planes may decorrelate over  $\Delta\eta$



# EP decorrelation over $\eta$

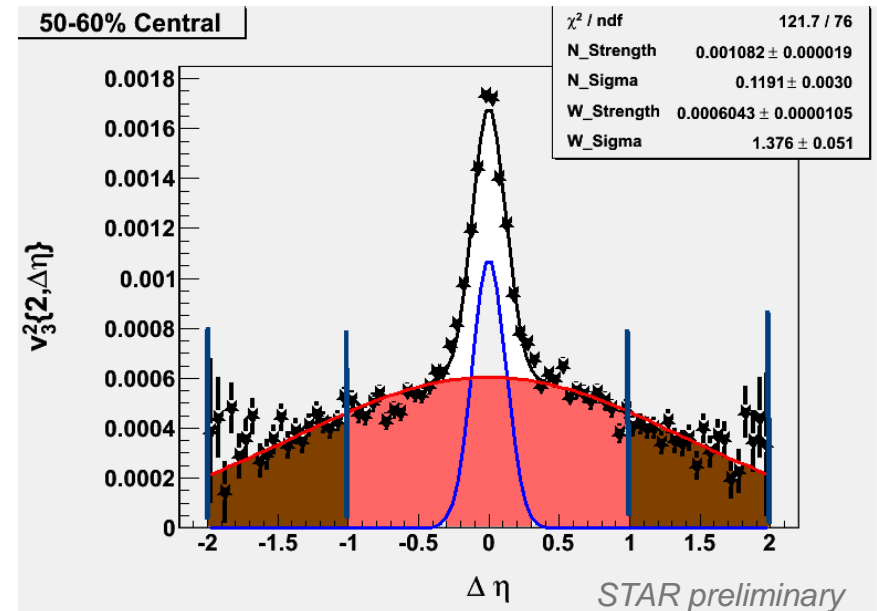
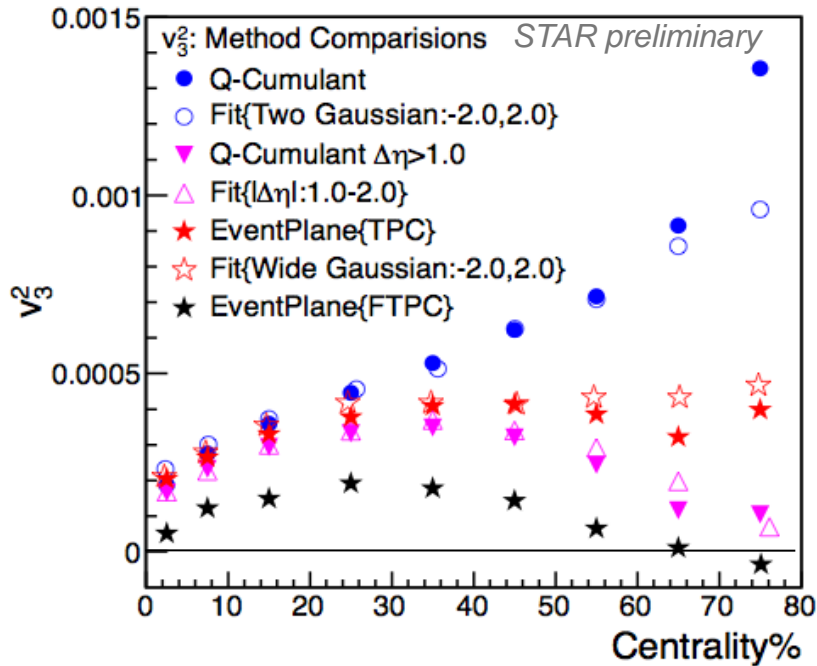
Xiao et al. arXiv:1208.1195



- EP's decorrelate with  $\Delta\eta$ .
- $\eta$ -gap reduces nonflow, but also under-measures flow.



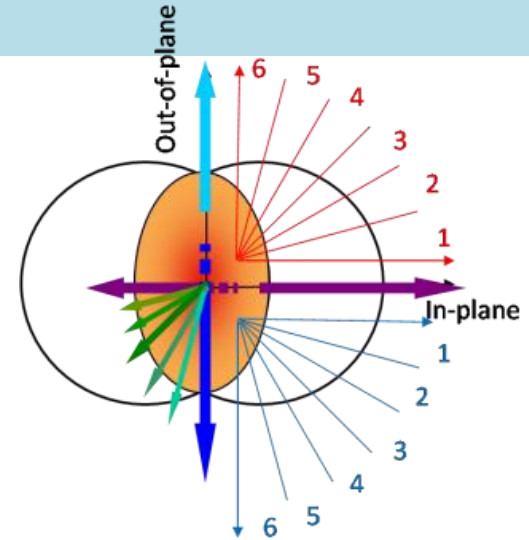
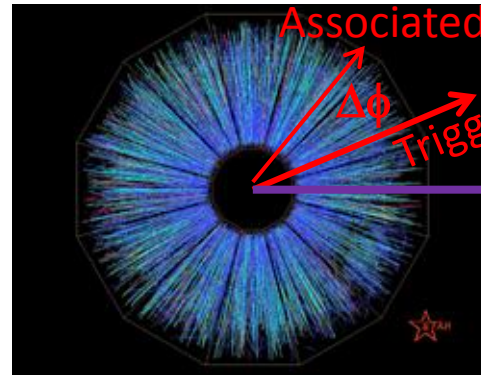
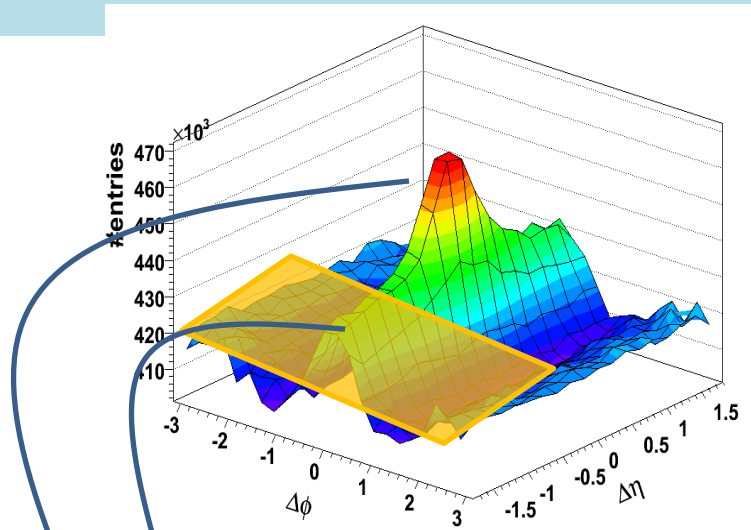
# $v_n$ depends on $\Delta\eta$



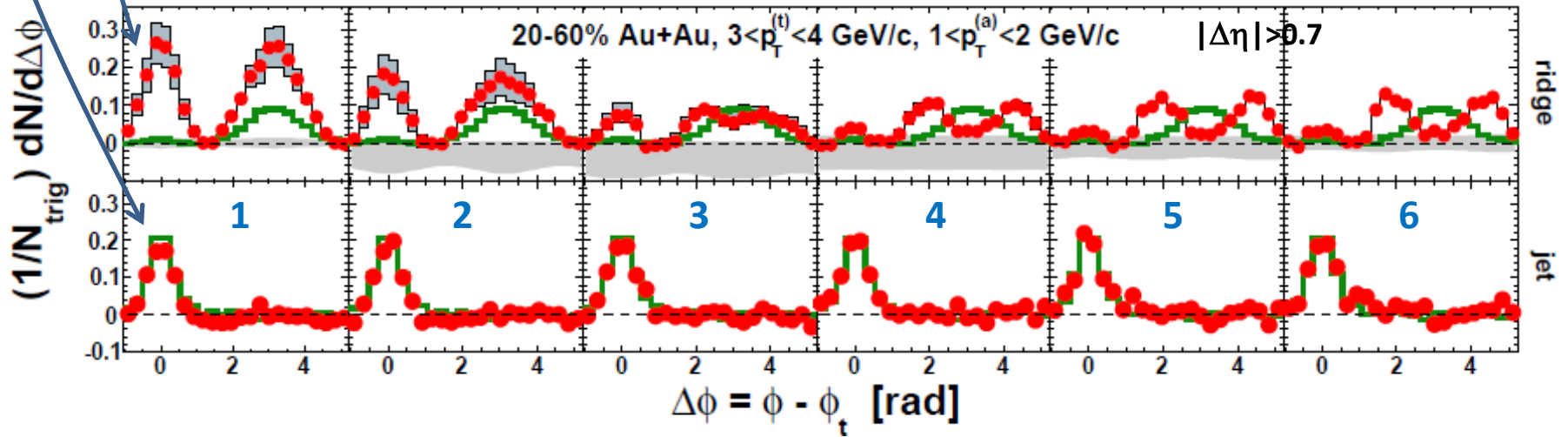
Li Yi (STAR) QM 2012.

The difference in different methods are due to different  $\Delta\eta$  window used to measure  $v_3$ .

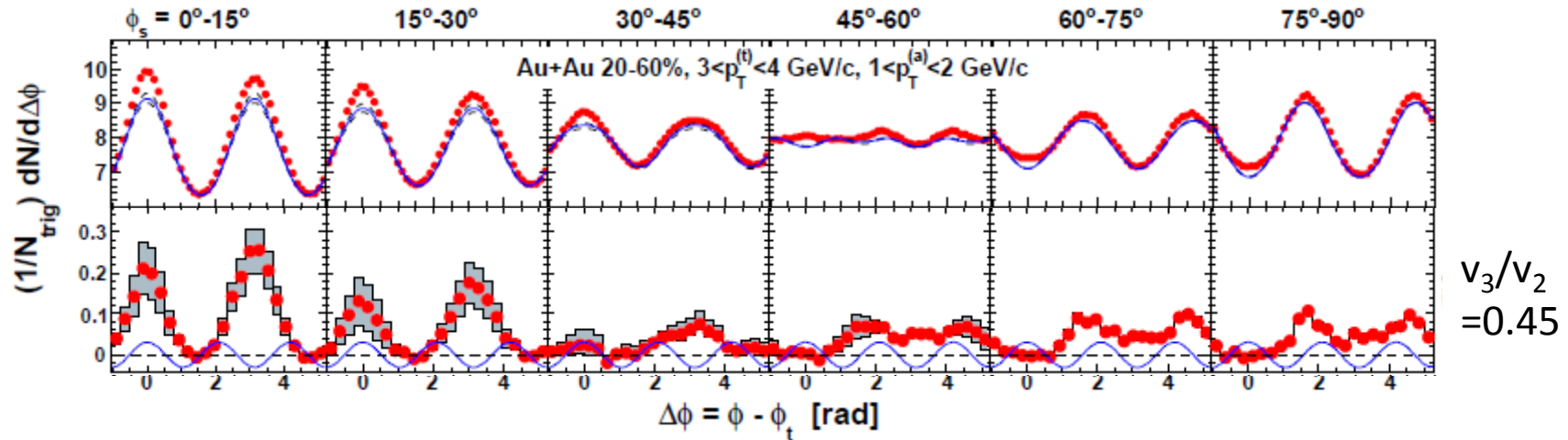
# RP-dependent dihadron correlations



STAR, arXiv:1010.0690



# $v_3$ does not change conclusion

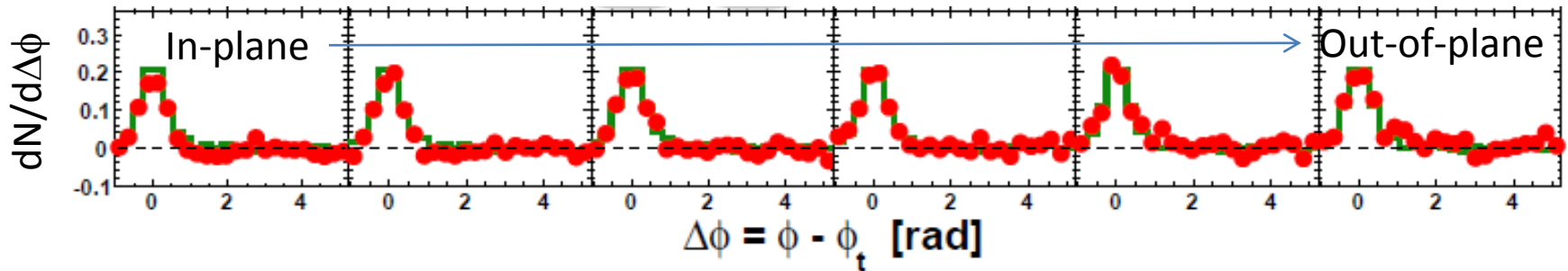


- Flow may depend on trigger particle direction  $\phi_s$  wrt RP, e.g.  $v_2$  may decrease with  $\phi_s$ .
- The in-plane ridge may not be as prominent. The in-plane away-side may be broad.
- We've taken  $\langle v_2^{\text{trig}} \rangle * \langle v_2^{\text{asso}} \rangle$ . This is OK even flow fluctuation is large, because  $v_2$  already include fluctuation. However, if  $v_2^{\text{asso}}$  is correlated with  $\phi_s$ , then  $\langle v_2^{\text{trig}} * v_2^{\text{asso}} \rangle \neq \langle v_2^{\text{trig}} \rangle * \langle v_2^{\text{asso}} \rangle$ .
- Actually, the inclusive correction  $\langle v_2^{\text{trig}} * v_2^{\text{asso}} \rangle$  should be larger than  $\langle v_2^{\text{trig}} \rangle * \langle v_2^{\text{asso}} \rangle_{\text{max}}$ ! The flow subtraction to inclusive dihadron correlation would be too small.

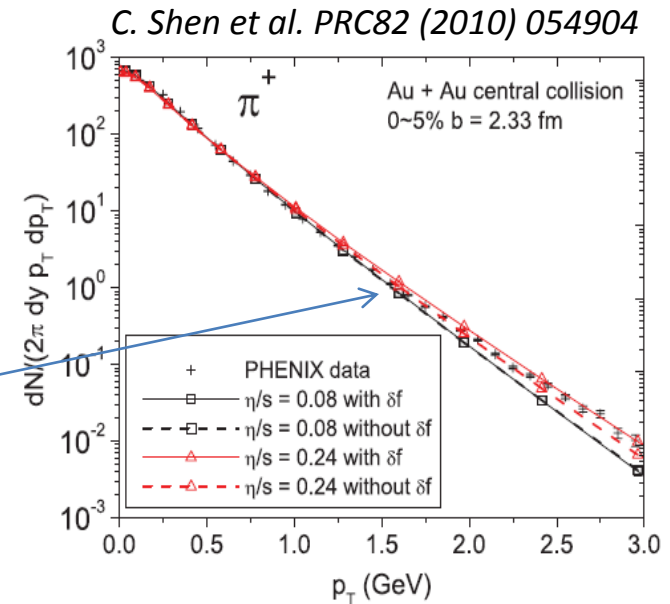


# Jet-like correlations

Trigger  $p_T=3-4$  GeV/c, Assoc  $p_T=1-2$  GeV/c

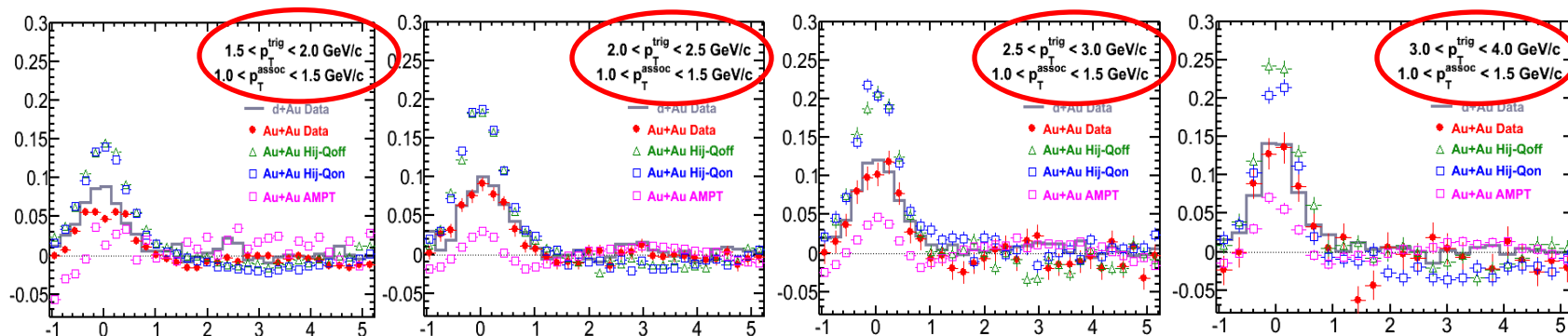


- Jet-like correlations with  $p_T > 3$  GeV/c trigger particles are invariant in-plane to out-of-plane, and from pp to central AA  $\rightarrow$  these high  $p_T$  particles are mainly from jets.
- Going to lower trigger  $p_T$ , expect hydro contribution to particle production  $\rightarrow$  Jet-like correlations will be reduced.
- Look at low  $p_T$  triggered dihadron correlations

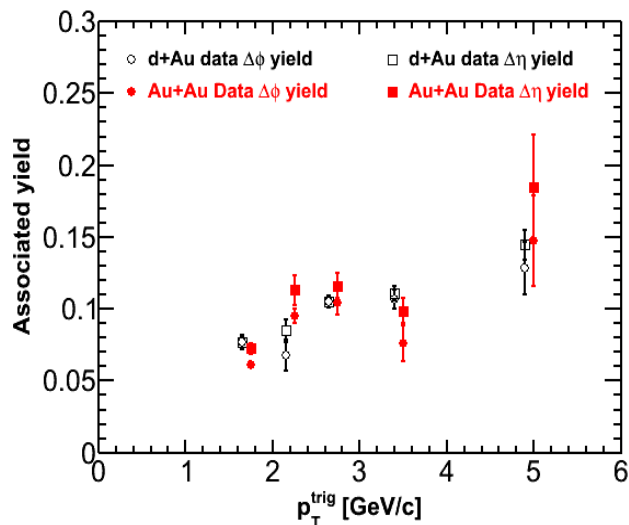


# Low- $p_T$ trigger jet-like correlations

Konzer (STAR) QM 2012



STAR Preliminary



- Surprisingly similar Au+Au and d+Au near-side jet-like signal strength
- No evidence of trigger dilution from background triggers
- Models do not reproduce data
- Challenges current understanding of particle production mechanisms



# Correlation and Flow:

## Do we understand them as well as we claim?

**NO.**

### Issues and open questions:

- Need to refrain a bit from simple Fourier components.
- $\eta$  gap to reduce nonflow may have undesired side effect.
- Is it possible to remove all nonflow?
- $\langle v_2^{\text{trig}} * v_2^{\text{asso}} \rangle \neq \langle v_2^{\text{trig}} \rangle * \langle v_2^{\text{asso}} \rangle$
- Can hydro particles have same jet-like correlations as in d-Au?
- Do we understand p-p and d-Au?