

# Accessing GPDs at LHeC

LHeC workshop , Chavannes-de-Bogis , June 2012

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based on work done with L Szymanowski, J Wagner, S. Wallon

Section 5.2.3 in LHeC report

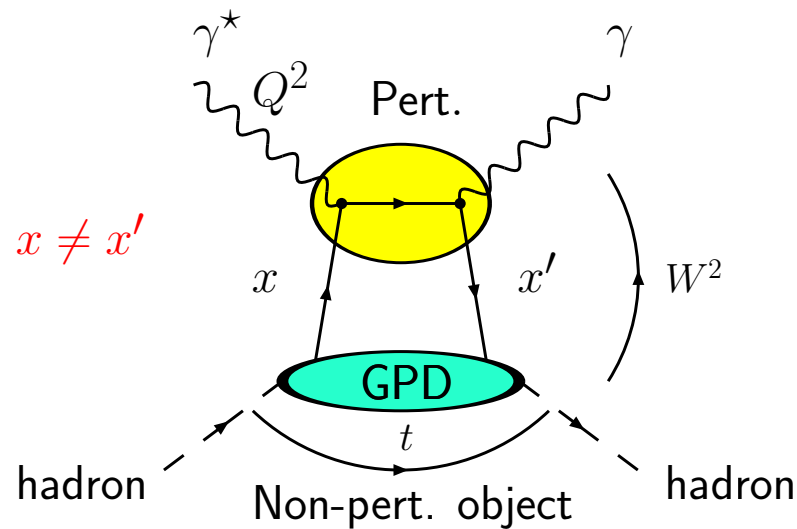
(Link to section 3 of EIC report : [arXiv :1108.1713 \[nucl-th\]](https://arxiv.org/abs/1108.1713) )

## Plan

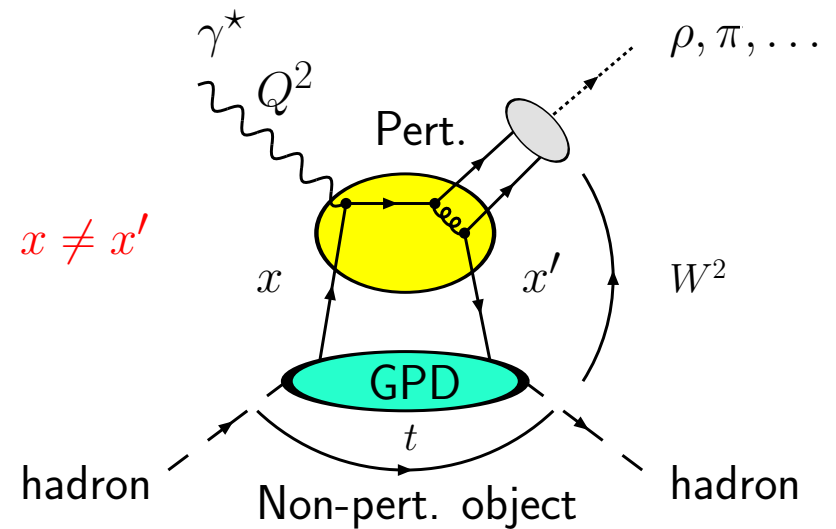
- ⇒ QCD (collinear) factorization for exclusive reactions
  - \* Success in DVCS : JLab, HERMES, HERA
  - GPD properties ; transverse imaging of the nucleon
- ⇒ DVCS and Timelike Compton scattering at very high energies
  - NLO corrections ; higher twist contributions
  - access in UPC at LHC ; access at LHeC
- ⇒ Exclusive meson electroproduction, spin physics and chiral-odd GPDs
- ⇒  $k_T$  vs collinear factorization ; searching for the Odderon

# QCD factorization in Exclusive processes

## DVCS



## Meson Production



⇒ **Factorisation** between a hard part (perturbatively calculable) and a soft part (non-perturbative) *Generalized Parton Distribution* **demonstrated** for

$$Q^2, W^2 \rightarrow \infty, x_B = \frac{Q^2}{Q^2 + W^2} \text{ fixed and } |t| \ll Q^2 \text{ fixed}$$

D. Mueller *et al.*, X. Ji, A. Radyushkin, J. Collins *et al.* , '94, '96, '98

## Some very good news

The simplest hard exclusive process :  $\gamma^*\gamma \rightarrow \pi^0$

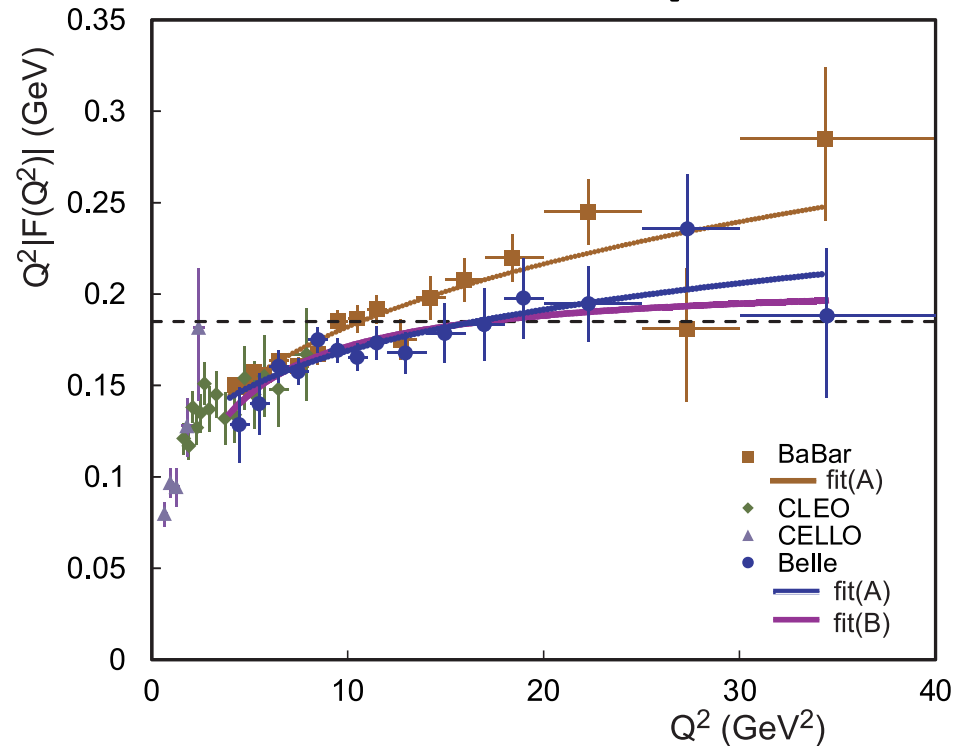
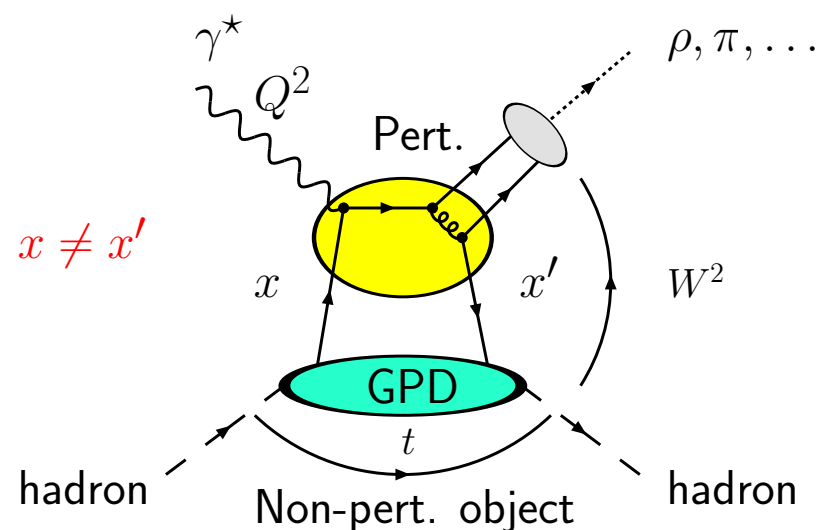
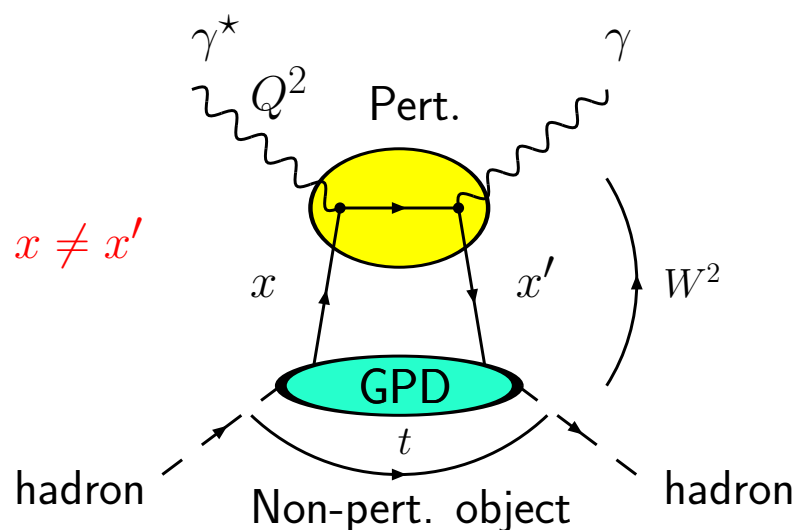


FIG. 24: Comparison of the results for the product  $Q^2|F(Q^2)|$  for the  $\pi^0$  from different experiments. The

**Belle confirms the applicability of perturbative QCD  
analysis to  $\pi^0\gamma$  transition form factor  
after BABAR data shaking the faith of theorists.**

# Generalised Parton Distributions

Non-Local operators (as in DIS) and **non diagonal** matrix elements  
= soft part of the amplitude for exclusive reactions



**GPD**  $(x, \xi, t)$  = Fourier Transform of matrix elements

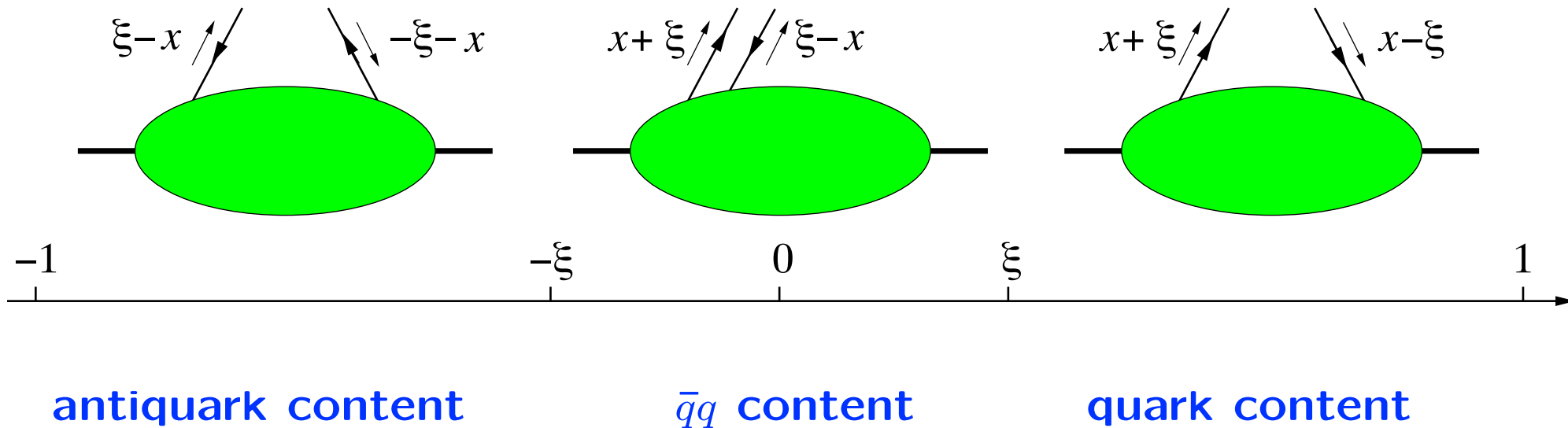
$$\langle N(p', \lambda') | \bar{\psi}(-z/2)_\alpha [-z/2; z/2] \psi(z/2)_\beta | N(p, \lambda) \rangle \Big|_{z^+=0, z_T=0}$$

**ON THE LIGHT CONE**  $z^2 = 0$

$$p' - p = \Delta \quad \Delta^2 = t \quad \Delta^+ = -\xi(p + p')^+ \quad x - x' = 2\xi$$

# Energy flow in GPDs

Three different regions : mean momentum fraction  $x$  vs skewness  $\xi$



Two different evolution equations

as  $\bar{q}(-x, Q^2)$

as  $\Phi^\pi(z, Q^2)$

as  $q(x, Q^2)$

DGLAP

ERBL

DGLAP

$\rightarrow \delta(-x)$

$\rightarrow \Phi_{as}^\pi(z, Q^2) = 6z\bar{z}$

$\rightarrow \delta(x)$

# Impact picture Representation

M. Burkhardt, JP Ralston and BP, M. Diehl

$t$  dependence of GPDs maps **transverse position**  $b_T$  of quarks.

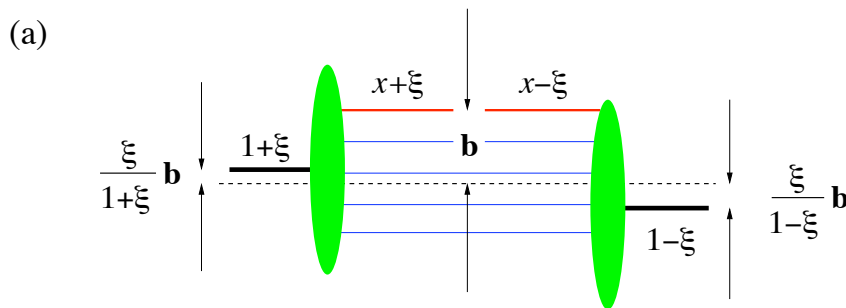
Fourier transform GPD at zero skewness

$$q(x, b_T) = (2\pi)^{-2} \int d^2 \Delta_T e^{i \Delta_T \cdot b_T} H(x, \xi = 0, t) \quad \text{probability}$$

Generalize at  $\xi \neq 0 \rightarrow$  **Quantum femtophotography**.

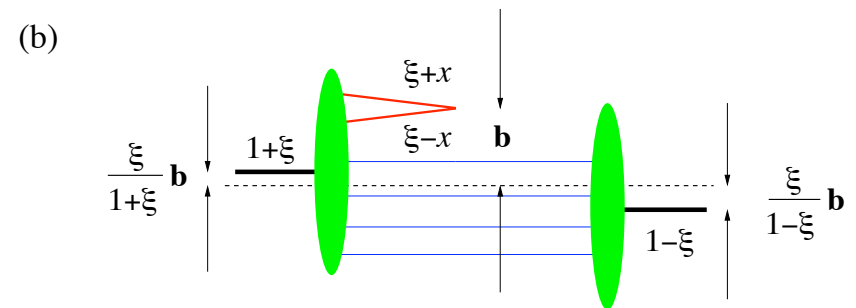
The  $t$ -dependence of dVCS localizes transversally in the proton the  $q$  and  $g$  (DGLAP) or the  $\bar{q}q$  and  $gg$  pairs of size  $\frac{1}{Q}$  (ERBL)

DGLAP region ( $x > \xi$ )



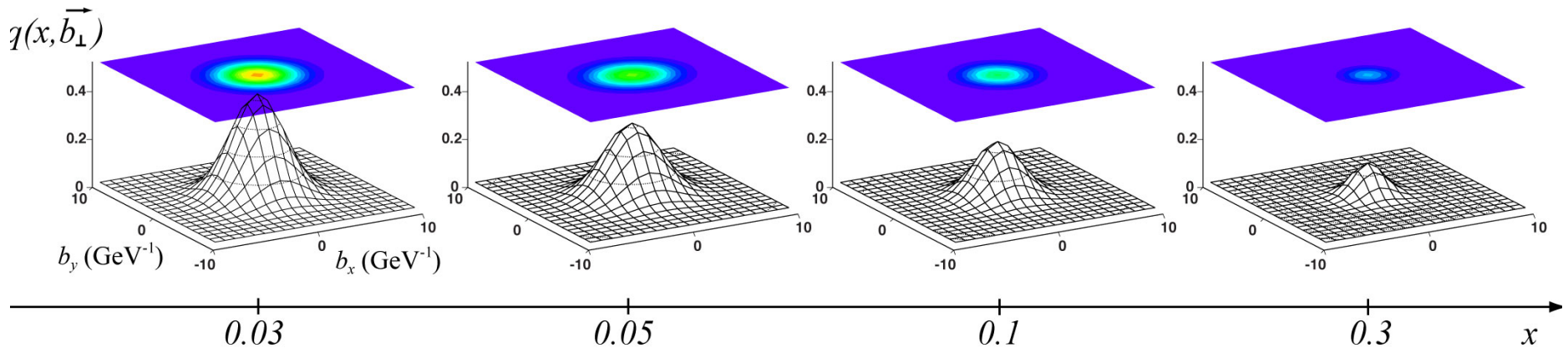
**Femtophotography of quark  
in the proton**

ERBL region ( $x < \xi$ )



**Femtophotography of quark-antiquark  
pair in the proton**

## Parton imaging with an EIC



for this dream to become textbook for students in the 2030's

one needs accurate measurement down to small values of  $\xi$

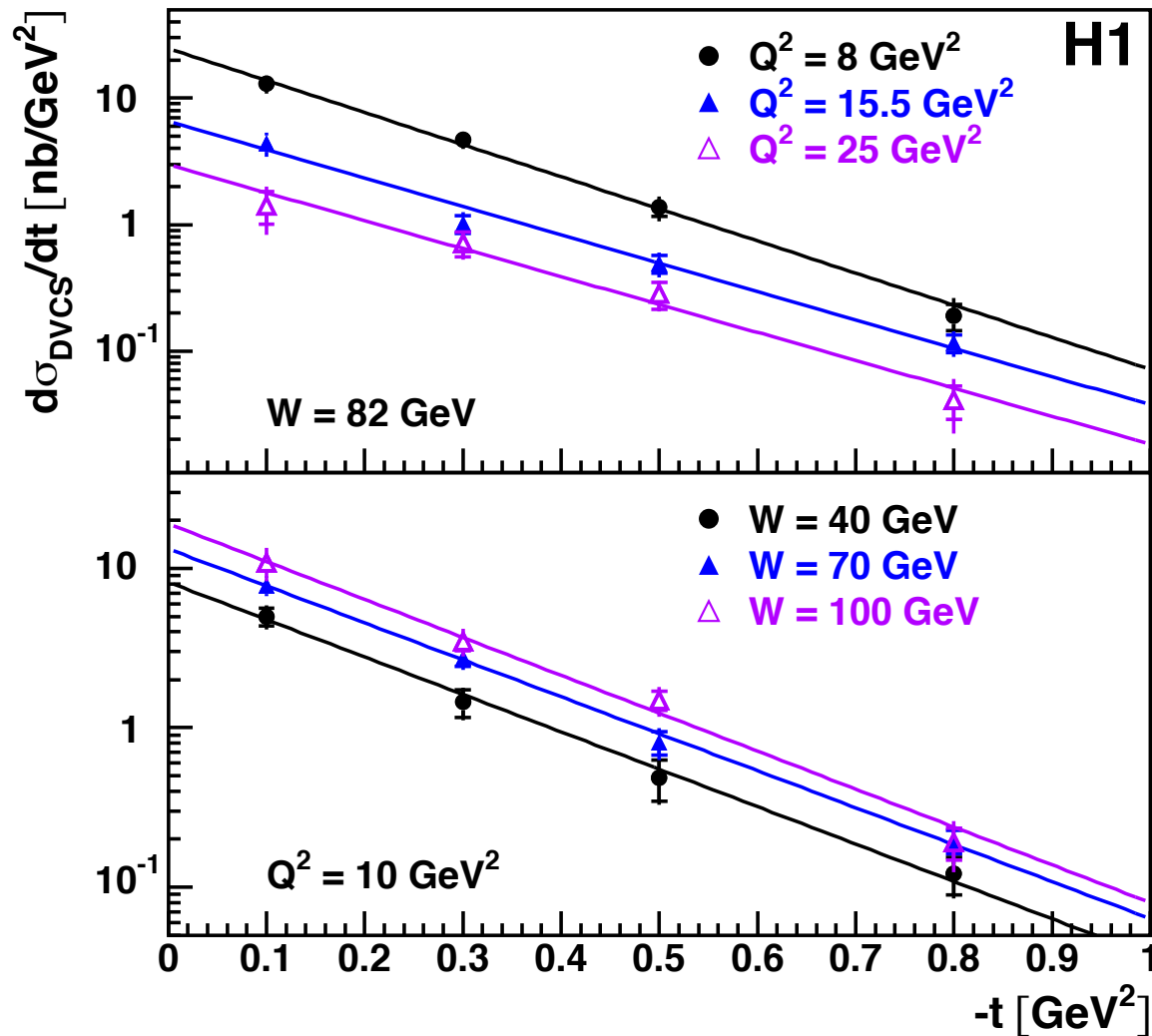
→ High energy, high luminosity electron nucleon colliders essential



## This is the reason I consider GPDs as a major breakthrough in QCD physics

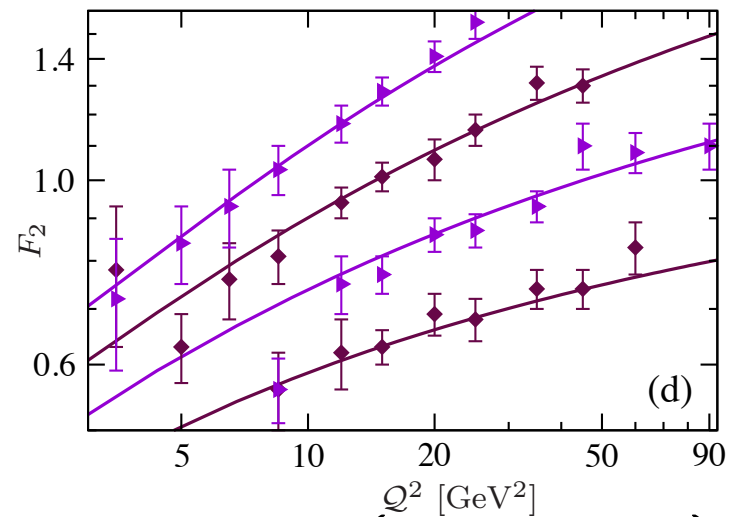
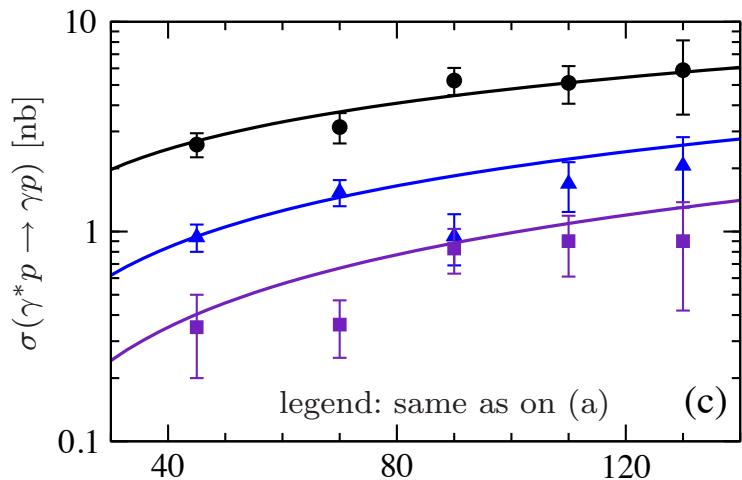
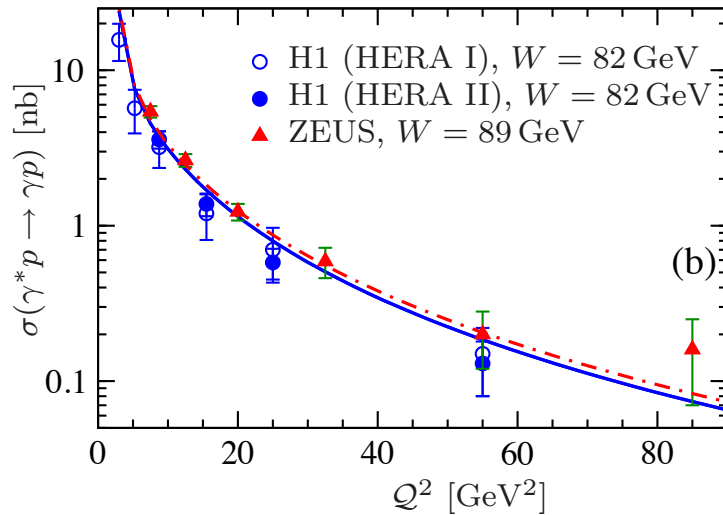
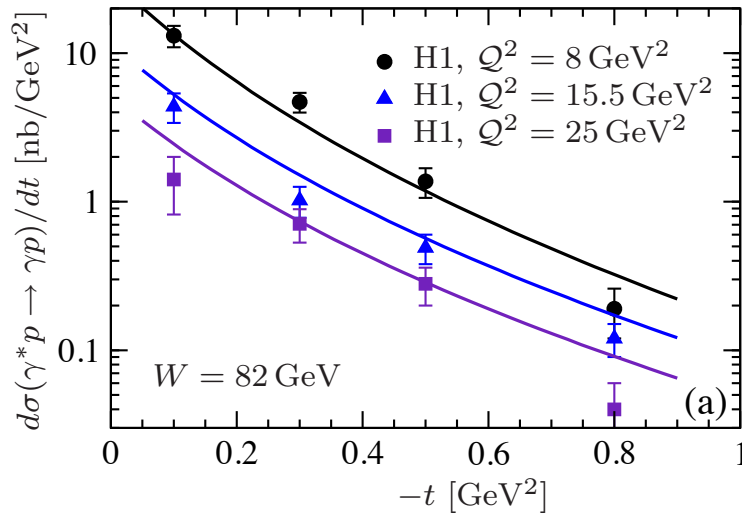
- ⇒ Beautiful progress in **forward** exclusive photon (DVCS) and meson (DVMP) experiments and analysis
- ⇒ Need to test universality of GPDs : **TCS vs DVCS** extractions
- ⇒ Need to better understand NLO and twist 3 contributions (  $\rightarrow \rho_T$  )
- ⇒ Need to resum soft gluon contributions (**Altinoluk et al 2012**)
- ⇒ Need to go to higher energies, smaller skewness (**EIC ; LHeC**)

# HERA results on DVCS



I forget on purpose the beautiful data at lower energies from  
**HERMES and JLab**

# Describing DVCS results with GPDs



Fitting DVCS (and DIS) with GPDs (and PDFs)

from K Kumericki and D. Mueller ArXiv 0904.0458

# DVCS simulation for LHeC

for  $\xi \approx x_B/2$  from  $2 \cdot 10^{-5}$  to  $6 \cdot 10^{-3}$

(for  $1 \text{ fb}^{-1}$  luminosity)

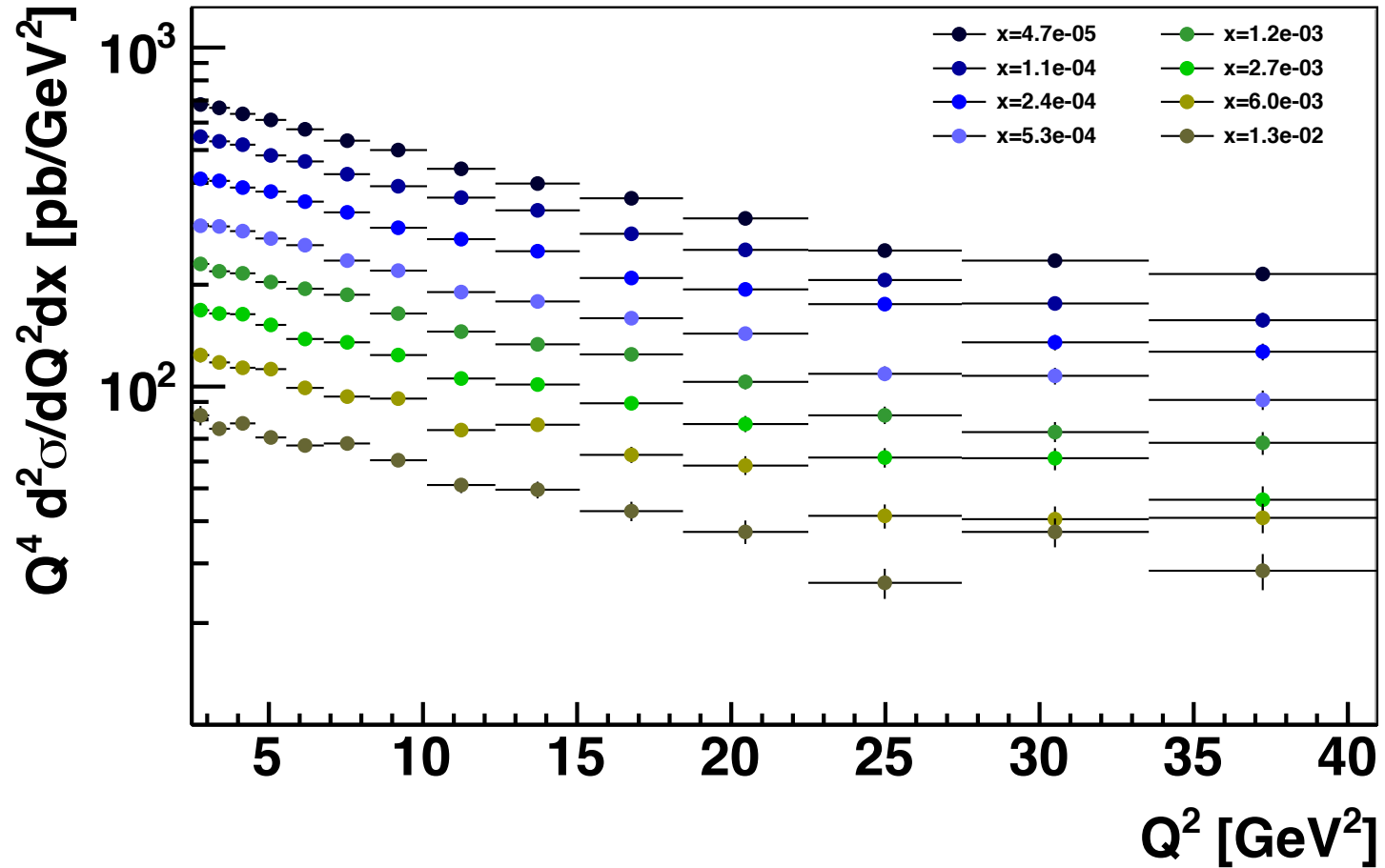


Figure 5.27: Simulated LHeC measurement of the DVCS cross section multiplied by  $Q^4$  for different  $x$  values for a luminosity of  $1 \text{ fb}^{-1}$ , with  $E_e = 50 \text{ GeV}$ , and electron and photon acceptance extending to within  $1^\circ$  of the beam pipe with a cut at  $P_T^\gamma = 2 \text{ GeV}$ . Only statistical uncertainties are considered.

and at very large  $Q^2$

$\xi$  from  $6 \cdot 10^{-5}$  to  $4 \cdot 10^{-3}$

(for  $100 \text{ fb}^{-1}$  luminosity)

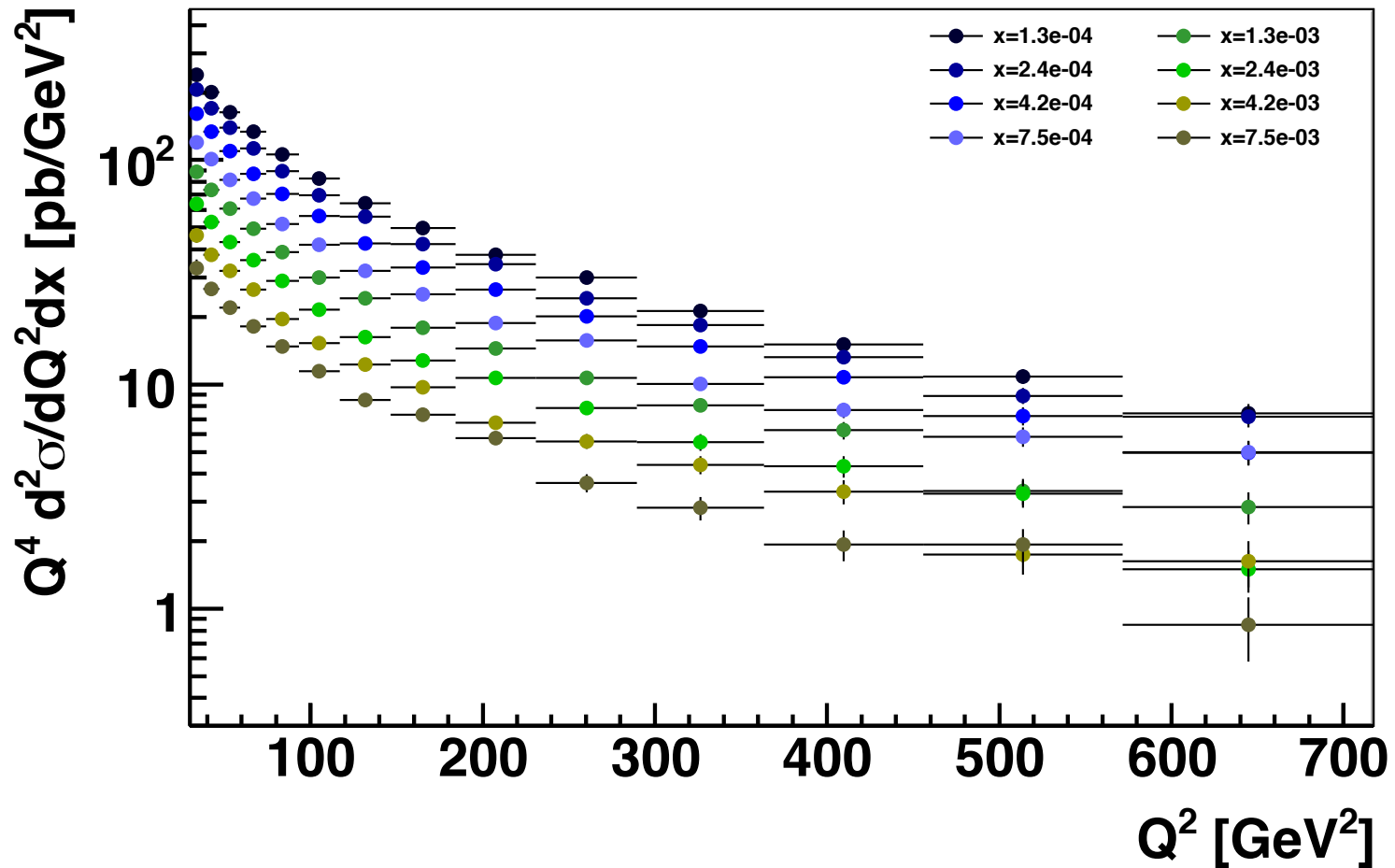
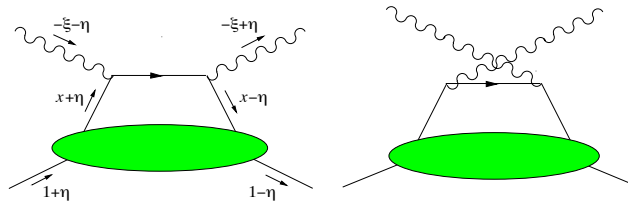


Figure 5.28: Simulated LHeC measurement of the DVCS cross section multiplied by  $Q^4$  for different  $x$  values for a luminosity of  $100 \text{ fb}^{-1}$ , with  $E_e = 50 \text{ GeV}$ , and electron and photon acceptance extending to within  $10^\circ$  of the beampipe with a cut at  $P_T^\gamma = 5 \text{ GeV}$ . Only statistical uncertainties are considered.

# On spacelike vs timelike probe

$$\gamma^*(q)N(p) \rightarrow \gamma^*(q')N'(p')$$

**DVCS vs TCS**



spacelike  $q^2 < 0$  ;  $q'^2 = 0$

vs

timelike  $q^2 = 0$  ;  $q'^2 > 0$

$$e N \rightarrow e' N \gamma$$

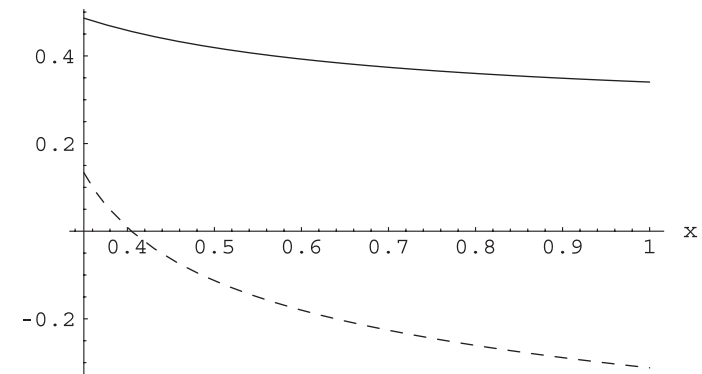
vs

$$\gamma N \rightarrow N \mu^+ \mu^-$$

$$\text{LO} : \mathcal{A}_{DVCS} = \mathcal{A}_{TCS}^*$$

$$R_{T-S}^q = \frac{C_{1(TCS)}^q - C_{1(DVCS)}^{q*}}{C_0^q}.$$

$$\text{NLO} : \mathcal{A}_{DVCS} \neq \mathcal{A}_{TCS}^*$$

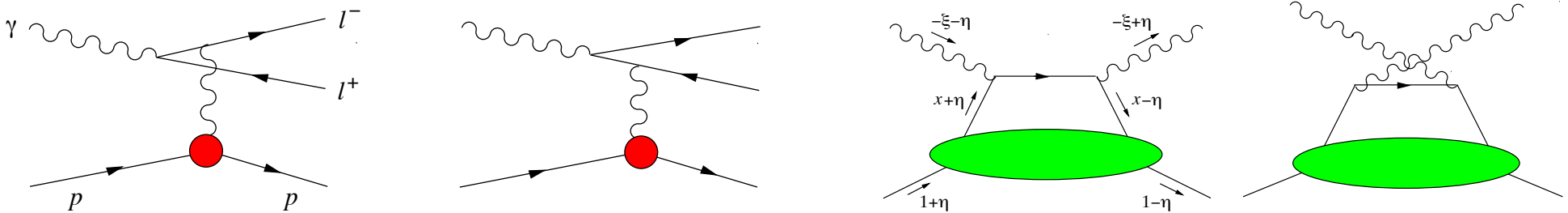


⇒ Both **timelike** and **spacelike** data useful to check NLO analysis !

# GPDs at LHC (and RHIC)

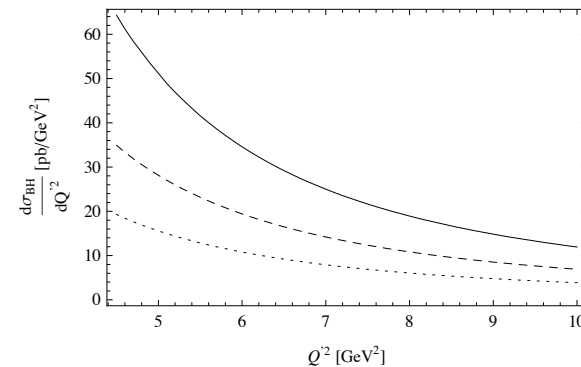
⇒ **Ultraperipheral Collisions : quasi real photons from proton beam**

$\mu^+ \mu^-$  pair production



QED dominates over TCS but in specific kinematics

→ **cutting out QED with angular cuts :**



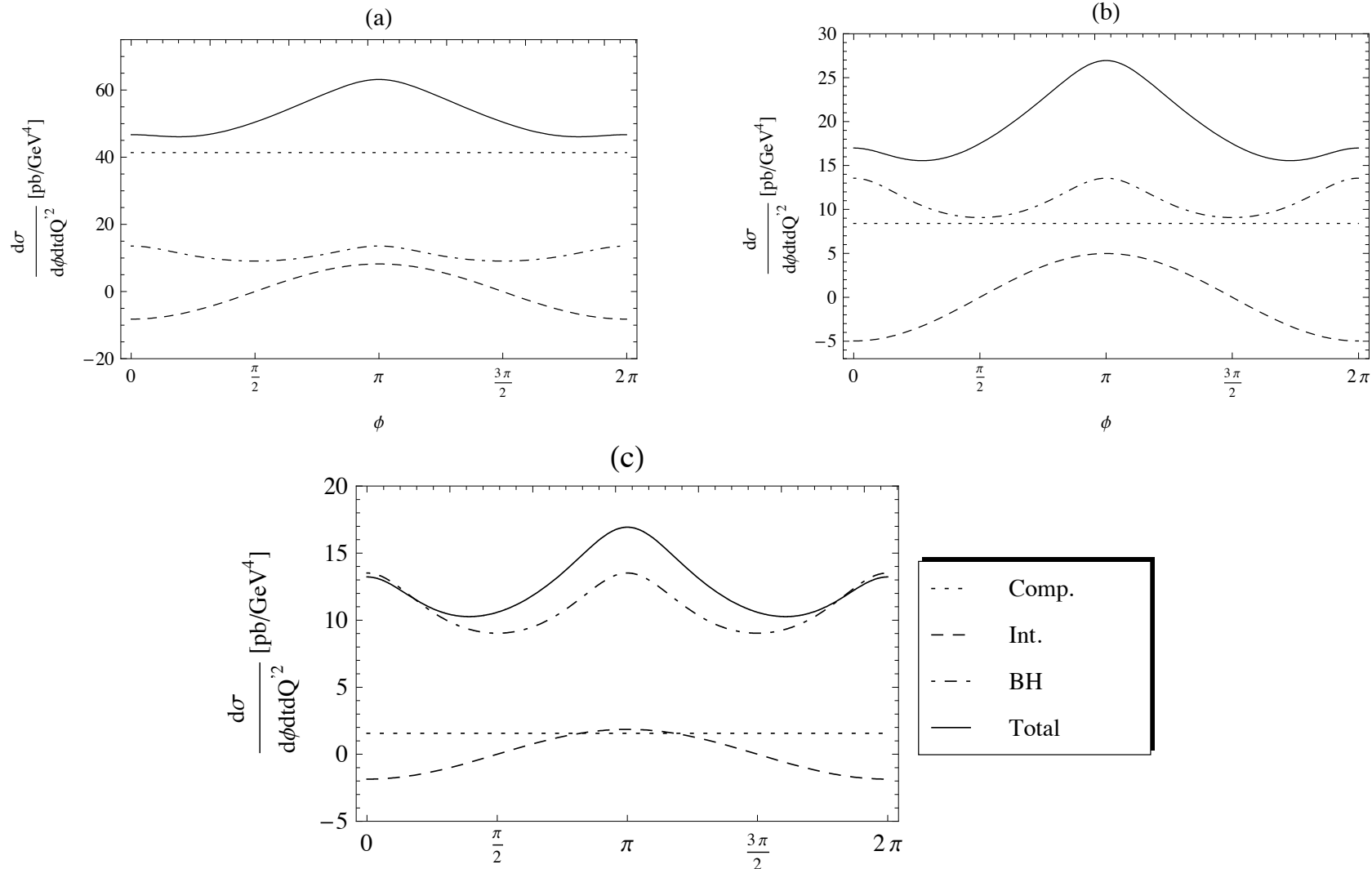
**GPDs are expected to be large at small  $x \approx \xi$**

$$\xi \approx Q^2 / s_{\gamma p}$$

⇒ **Probe of sea and gluon GPDs in small  $x$  regime**

# Observing TCS at LHC

➡ **Characteristic signal from interference** (charge conj. odd)

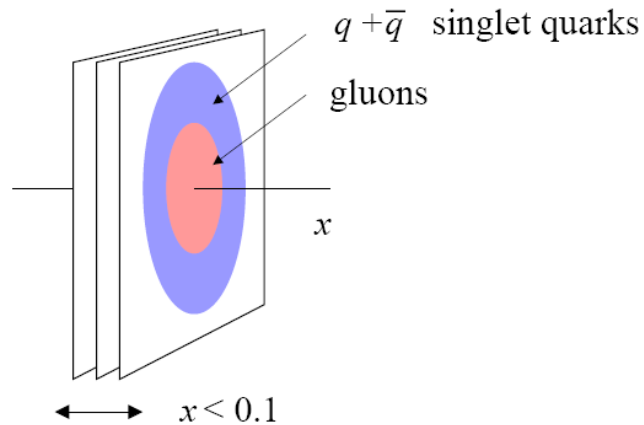


$$s_{\gamma p} = 10^7 \text{ GeV}^2 \quad s_{\gamma p} = 10^5 \text{ GeV}^2 \quad s_{\gamma p} = 10^3 \text{ GeV}^2 \quad Q'^2 = 5 \text{ GeV}^2$$

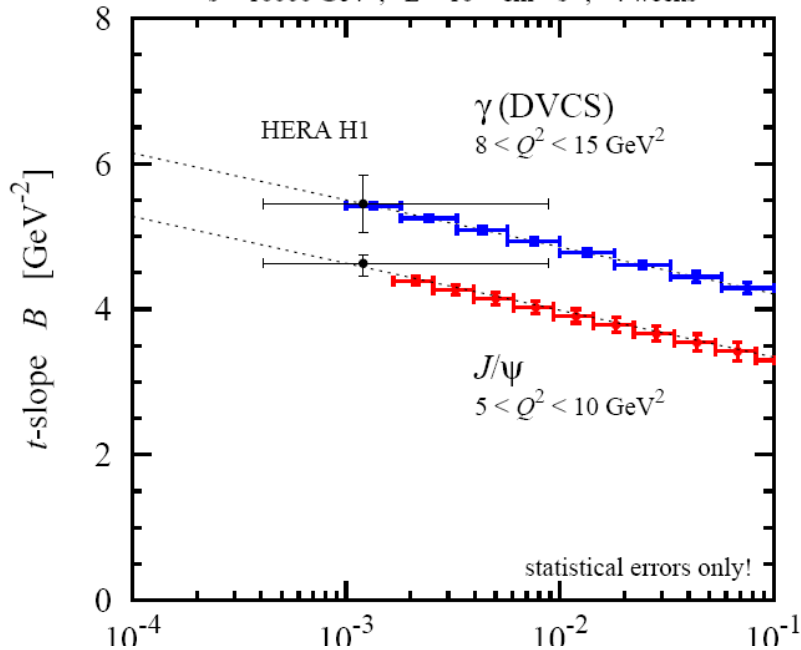


# Gluon imaging through DVCS

from T Horn, EIC study



$s = 10000 \text{ GeV}^2$ ,  $L = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ , 4 weeks



- Do singlet quarks and gluons have the same transverse distribution?

- Hints from HERA:  $\text{Area}(q + \bar{q}) > \text{Area}(g)$
- Dynamical models predict difference: pion cloud, constituent quark picture [Strikman, Weiss 09]
- No difference assumed in present pp MC generators for LHC!

- EIC: gluon size from  $J/\psi$ , singlet quark size from DVCS

- $x$ -dependence: quark vs. gluon diffusion in wave function
- Detailed analysis: LO NLO [Mueller et al.]

Detailed differential image of nucleon's partonic structure

# Exclusive Meson production and GPDs

**Vector meson production** the most dominant process

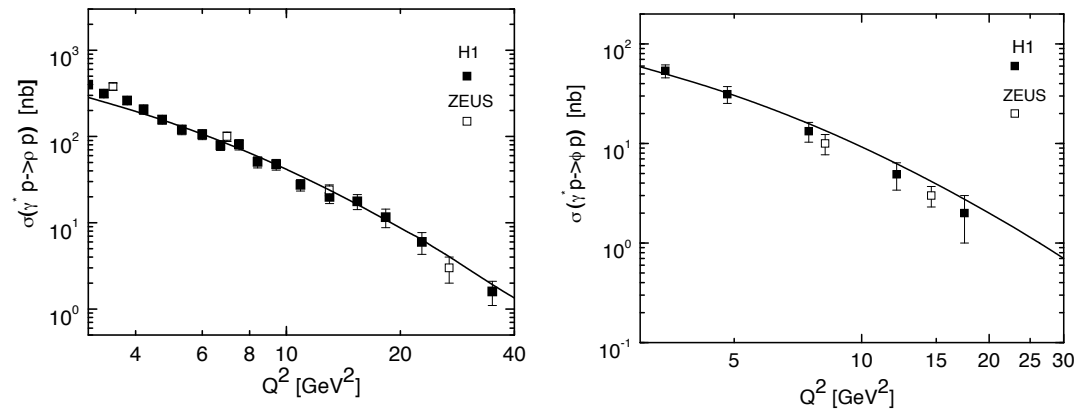


Figure 5: The integrated cross section for  $\gamma^* p \rightarrow \rho p$  (left) and  $\gamma^* p \rightarrow \phi p$  (right) versus  $Q^2$  at  $W \simeq 75$  GeV. Data taken from [12, 44] (filled squares) and [13, 45] (open squares) for  $\rho$  and  $\phi$  production, respectively. The solid lines represent our results.

Goloskokov Kroll hep-ph/0501242

$\rho_L$  production leading twist and dominant.

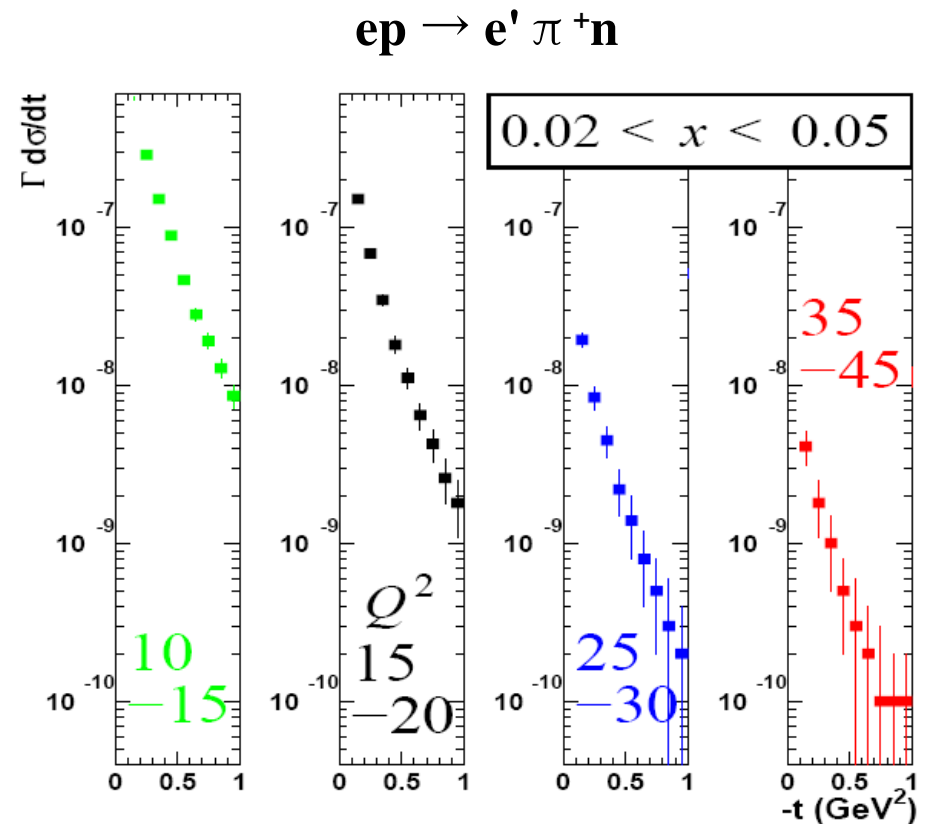
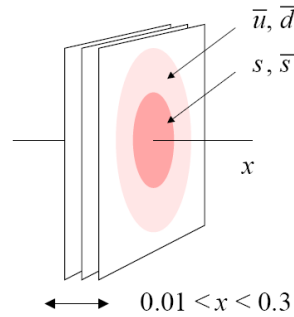
$\rho_L^0$  production mostly sensitive to gluon GPDs

probes  $H(x, \xi, t)$  and  $E(x, \xi, t)$  GPDs

# Sea quark imaging through $\pi$ electroproduction

from T Horn, EIC study

- New territory for collider!
- Spatial structure of *non-perturbative sea*
  - Closely related to JLab 12 GeV
    - Quark spin/flavor separations
    - Nucleon/meson structure
- Simulation for  $\pi^+$  production assuming 100 days at a luminosity of  $10^{34}$  with 5 on 50 GeV ( $s=1000 \text{ GeV}^2$ )
  - V. Guzey, C. Weiss: Regge model
  - T. Horn: empirical  $\pi^+$  parameterization
- Lower and more symmetric energies essential



[Tanja Horn, Antje Bruell, Christian Weiss]

Transverse spatial structure of non-perturbative sea quarks!

## Spin dependent GPDs

### A virtue of exclusive processes

→ spin physics without polarized beam/target !

$\pi$  or  $\eta$  production (twist 2) : sensitive to  $\tilde{H}(x, \xi, t)$  and  $\tilde{E}(x, \xi, t)$  GPDs

$$\text{recall : } \tilde{H}^q(x, 0, 0) = \Delta^q(x); \tilde{H}^g(x, 0, 0) = x\Delta^G(x)$$

because  $\gamma^5$  in meson DA selects  $\gamma^5$  in GPD operator

Vector meson production selects "helicity averaged" GPDs ;

Photon (DVCS or TCS) production mixes all 4 chiral even GPDs

What about chiral - odd GPDs ?

# Transversity GPDs

**Transverse spin** structure is very badly known !

⇒ Even at the PDF level ( $\Delta_{Tq}(x)$ ) :  
interesting but indirect (through TMDs) measurement of transversity  
⇒ Usefulness of direct measurements of chiral-odd GPDs.

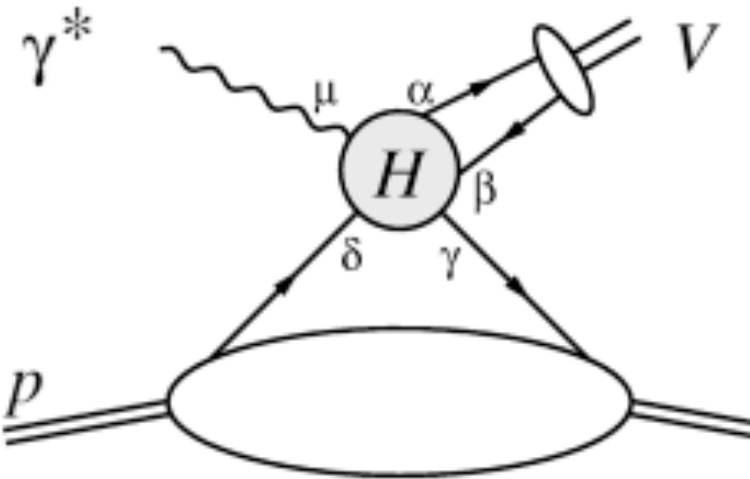
4 leading twist C-O GPDs :  $H_T^q(x, \xi, t)$ ,  $E_T^q(x, \xi, t)$ ,  $\tilde{H}_T^q(x, \xi, t)$ ,  $\tilde{E}_T^q(x, \xi, t)$

with  $H_T^q(x, 0, 0) = \Delta_{Tq}(x)$ .

Access to more features about the spin nucleon structure as :

- Momentum correlations of transversally polarized partons in a transversally polarized nucleon
- Transverse asymmetry of the angular momentum carried by quarks
- Density of polarized quarks in the impact parameter plane

# Exclusive electroproduction of one vector meson $V_T = \rho_T$



Diehl, Gousset, BP, Phys. Rev. D59, 034023  
Collins and Diehl, Phys. Rev. D61, 114015

**DA of  $V_T \propto \langle 0 | \bar{u}(0) \sigma^{\alpha\beta} u(x) | V \rangle$**

**C-O GPDs  $\propto \langle p' | \bar{u}(0) \sigma^{\delta\gamma} u(x) | p \rangle$**

**Basic Idea : C-O nature of  $\rho_T$  DA reveals the C-O GPDs**

**BUT** due to angular momentum and chirality conservation

$$\text{Tr}[H_{\alpha\beta\gamma\delta}^{\mu} \sigma^{\alpha\beta} \sigma^{\delta\gamma}] = 0$$

**at leading power in  $1/Q$  to all orders in the strong coupling**

**way out  $\rightarrow$  twist 3 contributions to  $\pi$  electroproduction ...or ...**

# Photo- or electroproduction of 2 vector mesons

Ivanov, BP, Szymanowski and Teryaev, Phys. Lett. B550, 65  
Enberg, BP and Szymanowski, Eur. Phys. J. C47, 87

Process at **high energy**, governed by the virtuality of the Pomeron

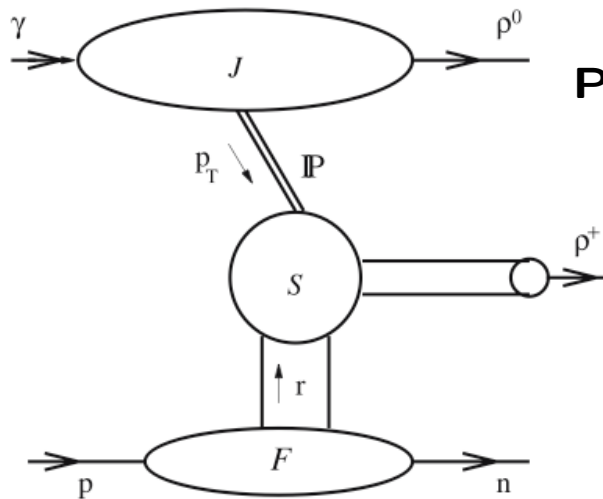


Photo- or electroproduction on a (polarized) nucleon target

Mesons ( $\rho_L^0$  and  $\rho_T^+$ ) separated by a large rapidity gap

**Hard scale** = virtuality of the Pomeron (or  $p_T$  of the  $\rho^0$ )

Estimate of rate and access to  $H_T$  in the ERBL domain

An hybrid description of exclusive processes

mixing  $k_T$  factorization (à la BFKL) and collinear factorization.

# Modeling Chiral-Odd GPDs

model for GPD based on tensor meson exchange :

$$H_T^{ud}(x, \xi, 0) = \frac{g_{b_1 NN} f_{b_1}^T \langle k_\perp^2 \rangle}{2\sqrt{2} M_N m_{b_1}^2} \frac{\phi_\perp^{b_1} \left( \frac{x+\xi}{2\xi} \right)}{2\xi}$$

+ some existing first Lattice QCD calculations

→ Born amplitude through  $\gamma \rightarrow \rho^0$  Impact factor  $J^{\gamma \rightarrow \rho^0}$  :

$$\begin{aligned} \mathcal{M}^{\gamma p \rightarrow \rho^0 \rho_T^+ n} &= \sin \theta \, 16\pi^2 W^2 \alpha_s f_\rho^T \xi \sqrt{\frac{1-\xi}{1+\xi}} \frac{C_F}{N_c (p_T^2)^2} \\ &\times \int_0^1 \frac{du \, \phi_\perp(u)}{u^2 \bar{u}^2} J^{\gamma \rightarrow \rho^0}(u p_T, \bar{u} p_T) \frac{H_T^{ud}(\xi(2u-1), \xi, t)}{\sqrt{2}} \end{aligned}$$

Nucleon spin orientation appears only through  $\sin(\theta)$

**unpolarized cross section non-zero !**



## Estimated cross sections

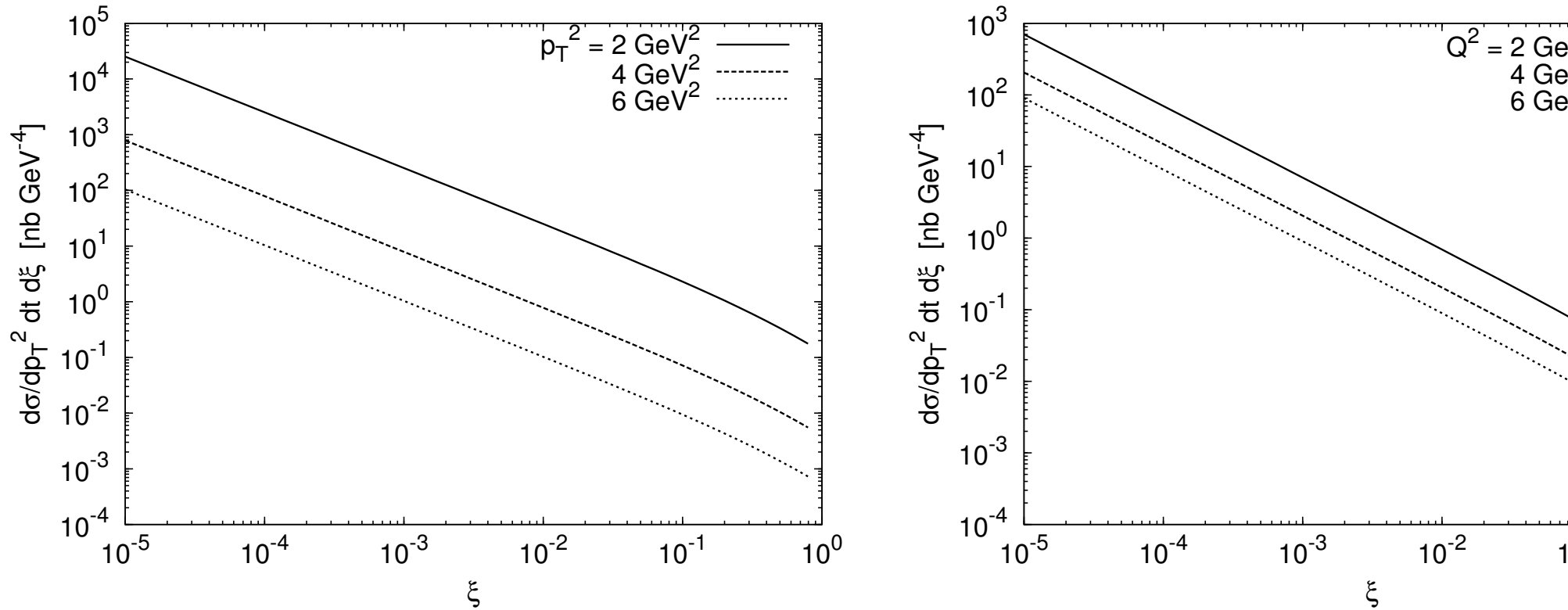


Figure 5.30: The differential cross section for the photoproduction (a) and electroproduction (b) of the meson pair  $\rho_T^0 \rho_T^+$  as a function of  $\xi$  for (a)  $p_T^2 = 2, 4$  and  $6 \text{ GeV}^2$  and for (b)  $p_T^2 = 2 \text{ GeV}^2$  and  $Q^2 = 2, 4$  and  $6 \text{ GeV}^2$ . The cross sections for the production of the meson pair  $\rho_T^0 \rho_T^0$  are two times smaller.

**Chiral-Odd GPDs should be measured at EIC / LHeC**

# Looking for the Odderon

(within  $k_T$  factorization - link to twist 3 gluon GPD ?)

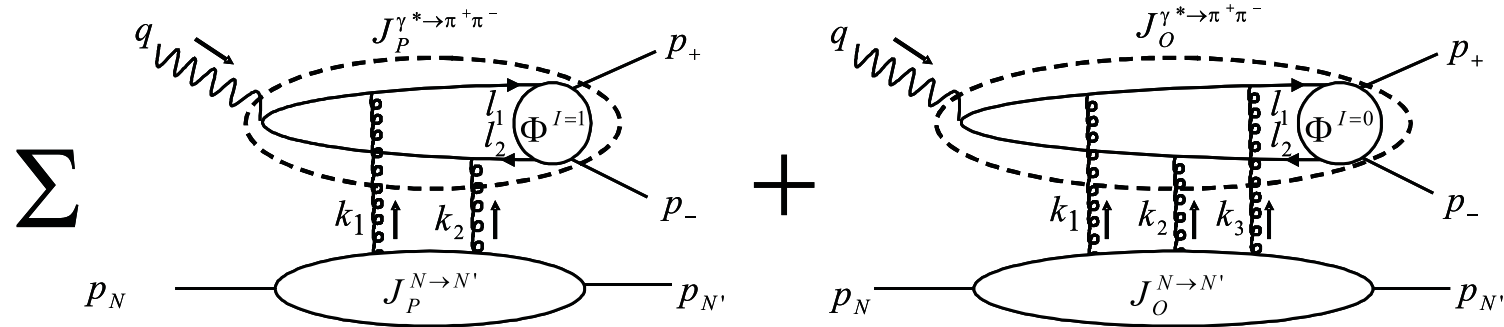


Figure 6.33: Feynman diagrams describing  $\pi^+\pi^-$  electroproduction in the Born approximation.

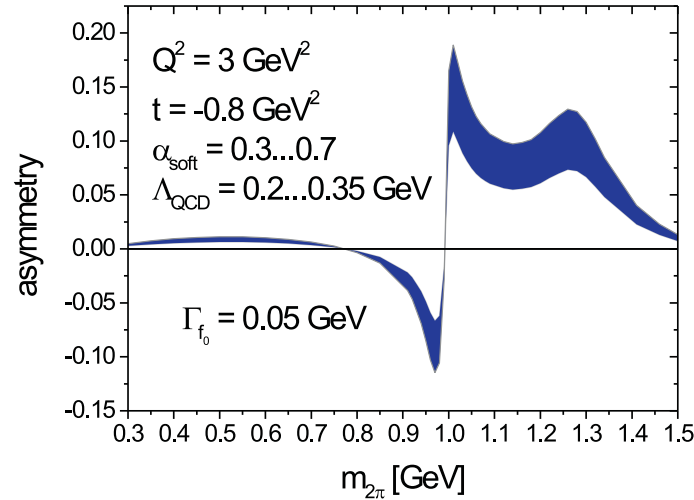


Figure 6.34: The charge asymmetry defined in Eq. (6.17) as a function of the  $\pi^+\pi^-$  invariant mass  $m_{2\pi}$ .

$$A(Q^2, t, m_{2\pi}^2) = \frac{\int \cos \theta d\sigma(s, Q^2, t, m_{2\pi}^2, \theta)}{\int d\sigma(s, Q^2, t, m_{2\pi}^2, \theta)} = \frac{\int_{-1}^1 \cos \theta d\cos \theta 2 \operatorname{Re} \left[ \mathcal{M}_P^{\gamma_L^*} (\mathcal{M}_O^{\gamma_L^*})^* \right]}{\int_{-1}^1 d\cos \theta \left[ |\mathcal{M}_P^{\gamma_L^*}|^2 + |\mathcal{M}_O^{\gamma_L^*}|^2 \right]}$$

## Conclusions

- ⇒ The future of GPD measurements is bright at medium energies : JLab 12, COMPASS 2
- ⇒ Much work began to uncover EIC possibilities in this domain.
- ⇒ Much remains to be done for LHeC, including detector requirements.

**NB : some uncovered items : nuclear GPDs, including coherent and break-up cases, saturation phenomena ...**

**Thank you**