



FIAS Frankfurt Institute
for Advanced Studies



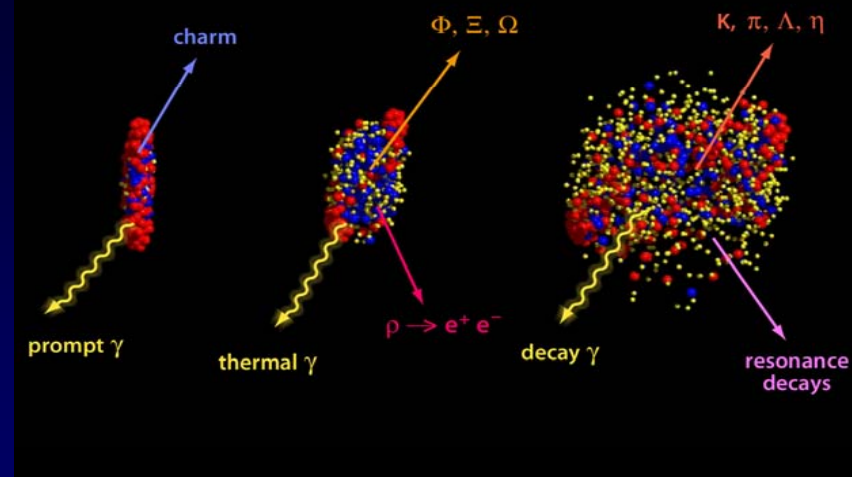
Heavy quarks and high p_T in nuclear collisions

Elena Bratkovskaya

13.09.2007 , Meeting of the Virtual Institute for Strongly Interacting
Matter (VI-SIM), "Heavy Ion Physics Perspectives",
Monbachtal e.v., Bad Liebenzell

Signals of the phase transition:

- Strangeness enhancement
- Multi-strange particle enhancement
- Charm suppression
- Collective flow (v_1, v_2)
- Thermal dileptons
- Jet quenching and angular correlations
- High p_T suppression of hadrons
- Nonstatistical event by event fluctuations and correlations
- ...



Experiment: measures final hadrons and leptons

How to learn about physics from data?

Compare with theory!

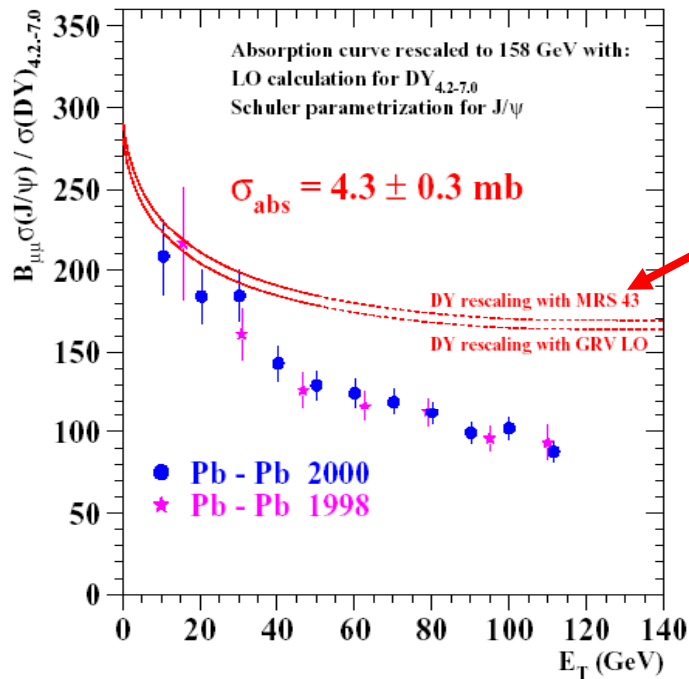
Microscopical transport models provide a unique dynamical description of nonequilibrium effects in heavy-ion collisions

Open and hidden charm

,Anomalous' J/Ψ suppression in A+A

Heavy flavor sector reflects the actual dynamics since heavy hadrons can **only be formed in the very early phase** of heavy-ion collisions !

■ **Anomalous J/Ψ suppression in A+A (NA38/NA50/NA60)**



J/Ψ ,normal' absorption
by nucleons
(Glauber model)

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Experimental observation:
extra suppression in A+A
collisions; increasing with
centrality

Scenarios for charmonium suppression in A+A

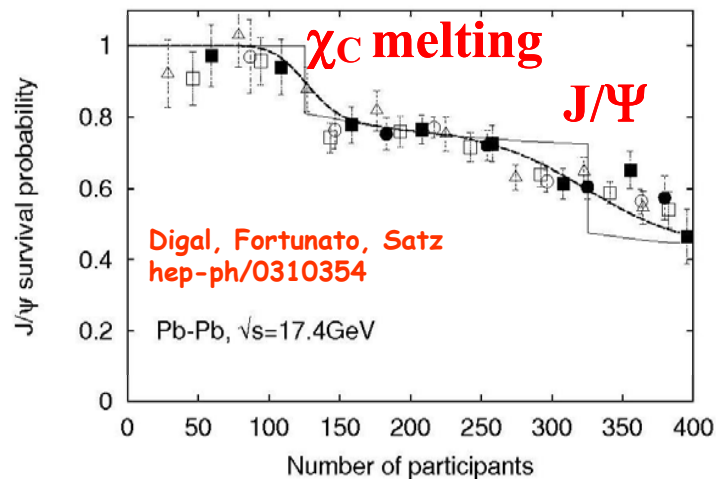
• QGP threshold melting

[Satz et al'03]

Quarkonium dissociation temperatures:

state	J/ψ(1S)	χ _c (1P)	ψ'(2S)	Υ(1S)	χ _b (1P)	Υ(2S)	χ _b (2P)	Υ(3S)
T _d /T _c	2.10	1.16	1.12	> 4.0	1.76	1.60	1.19	1.17

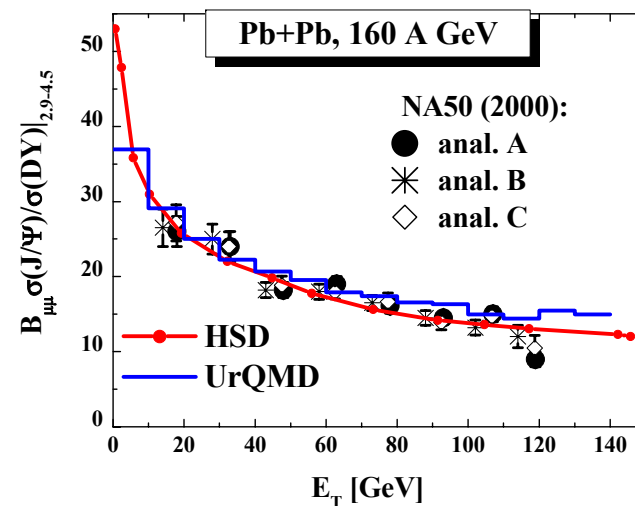
Dissociation energy density $\epsilon_d \sim 2(T_d/T_c)^4$



• Comover absorption

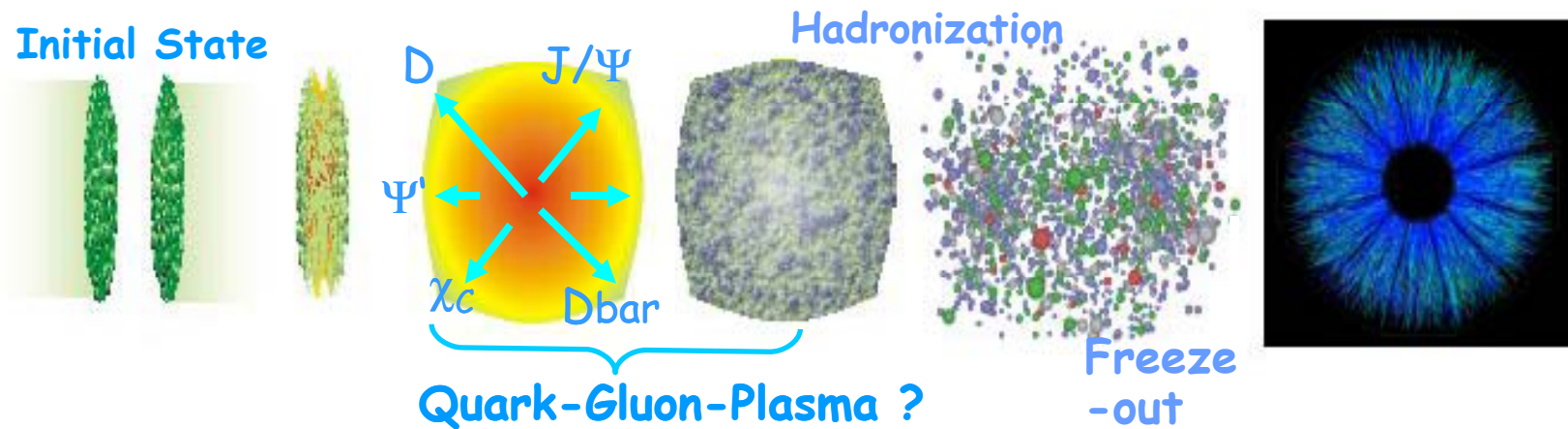
[Gavin & Vogt, Capella et al.'97]:

charmonium absorption by low energy inelastic scattering with ,comoving' mesons (m=π,η,ρ,...):



Charmonium dynamics -> HSD

Check scenarios for charmonium suppression in A+A
using **microscopic transport models**



Transport models

Microscopic transport models provide the dynamical description of **nonequilibrium** effects in heavy-ion collisions

HSD – **H**adron-**S**tring-**D**ynamics transport approach

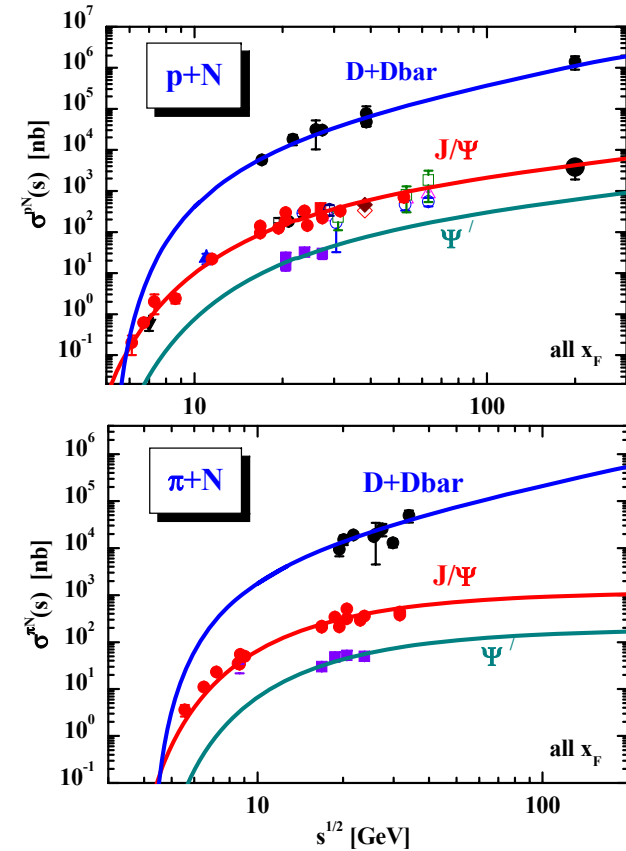
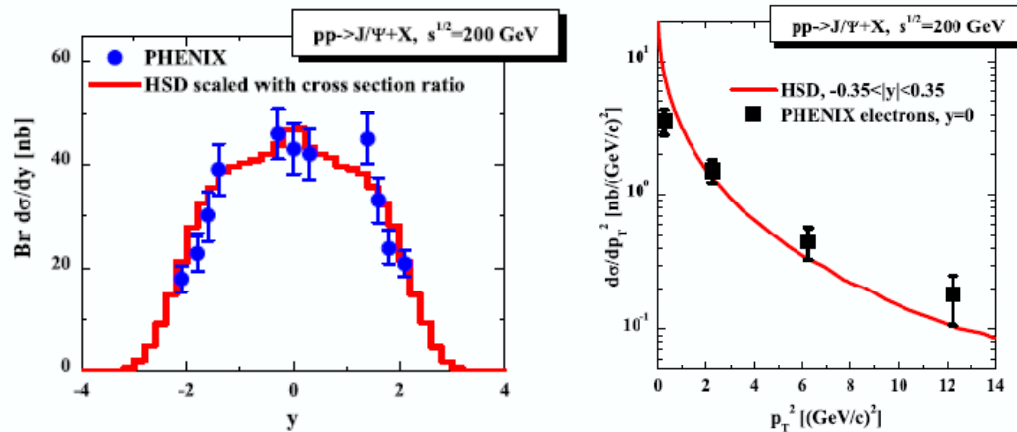
Charm and Charmonium production and absorption in HSD

- **Charmonium = hard probe**
 => **binary scaling!**
- **Production $\sigma(J/\Psi)$ and $\sigma(\Psi')$ in N+N and π +N collisions: parametrization of the available exp. data**

Coupled channel problem:

$$\sigma_{J/\Psi}^{exp} = \sigma_{J/\Psi} + B(\chi_c \rightarrow J/\Psi) \sigma_{\chi_c} + B(\Psi' \rightarrow J/\Psi) \sigma_{\Psi'}$$

- **y -, p_T - distributions of J/Ψ in pp at RHIC are controlled by the PHENIX data :**



Charmonia-baryon dissociation cross sections

$$J/\Psi (\chi_c, \Psi') + B \rightarrow D + \bar{D} + X$$

- Charmonia-baryon dissociation cross sections can be fixed from **p+A** data:

HSD-2003:

Pre-resonance c-cbar pairs (color-octet states):

$$\sigma_{cc\ B} = 6 \text{ mb} \quad (\tau_{cc} = 0.3 \text{ fm}/c)$$

Formed charmonium (color-singlet states):

$$\sigma_{J/\Psi\ B} = 4 \text{ mb}, \quad \sigma_{\chi\ B} = 5 \text{ mb}, \quad \sigma_{\Psi'\ B} = 8 \text{ mb}$$

HSD-2006/2007:

$$\sigma_{cc\ B} = \sigma_{J/\Psi\ B} = \sigma_{\chi\ B} = 4.18 \text{ mb}, \quad \sigma_{\Psi'\ B} = 7.6 \text{ mb}$$

adopting a new Glauber fit from NA50

Modelling of the comover scenario in HSD

1. Charmonia **dissociation** cross sections with π , ρ , K and K^* mesons $J/\Psi (\chi_c, \Psi') + \text{meson } (\pi, \rho, K, K^*) \leftrightarrow D+Dbar$

- **Phase-space model for charmonium + meson dissociation:**

$$\sigma_{1+2 \rightarrow 3+4}(s) = g_{\text{isospin}} 2^4 \frac{E_1 E_2 E_3 E_4}{s} |M_i|^2 \left(\frac{m_3 + m_4}{\sqrt{s}} \right)^6$$

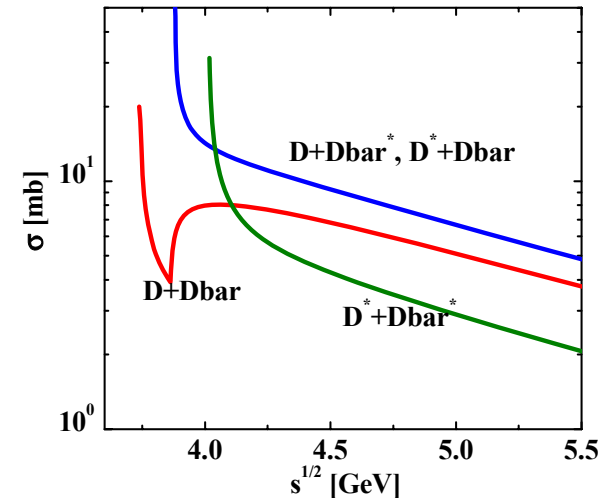
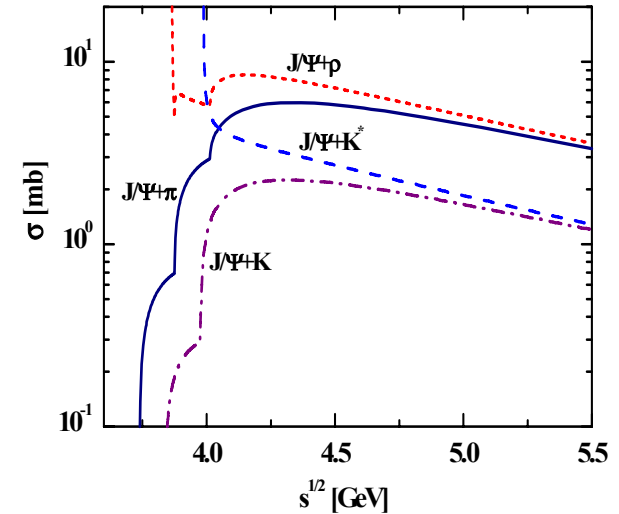
$$i = \chi_c, J/\Psi, \Psi'$$

$$|M_{J/\Psi}|^2 = |M_{\chi_c}|^2 = |M_{\Psi'}|^2 = |M_0|^2$$

constant matrix element

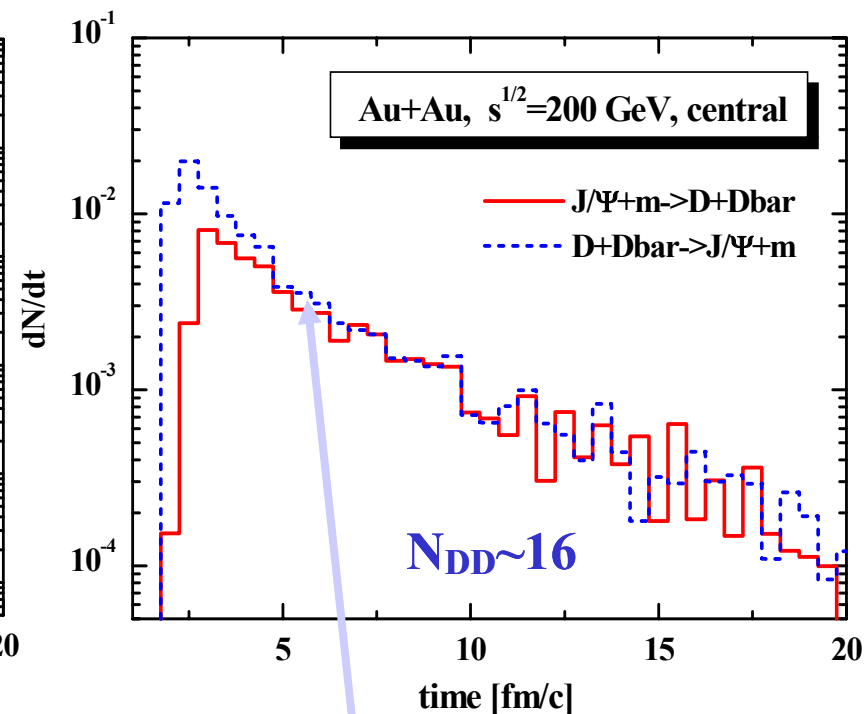
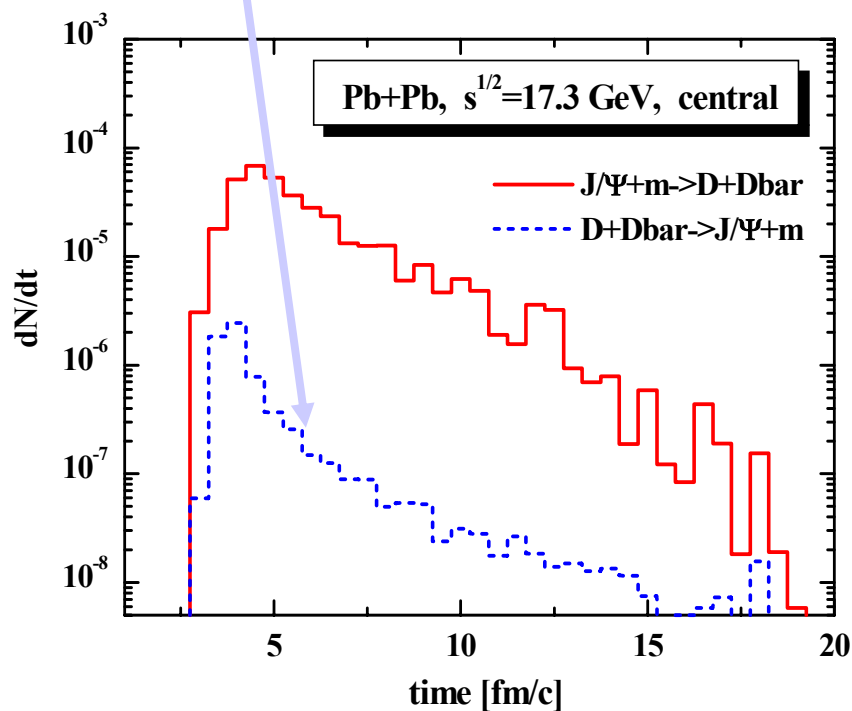
2. J/Ψ **recombination** cross sections by $D+Dbar$ annihilation:

$D+Dbar \rightarrow J/\Psi (\chi_c, \Psi') + \text{meson } (\pi, \rho, K, K^*)$
are determined by detailed balance!



Charmonium recombination by D-Dbar annihilation

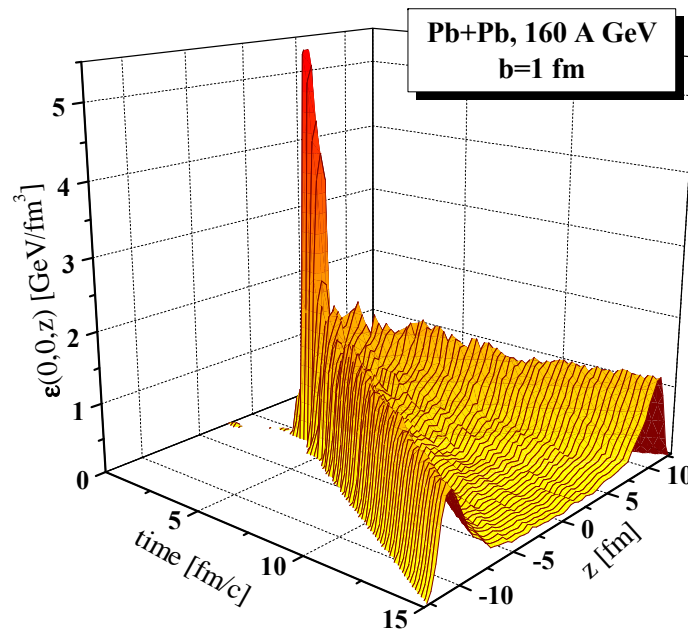
At **SPS** recreation of J/Ψ by $D+Dbar$ annihilation is **negligible**



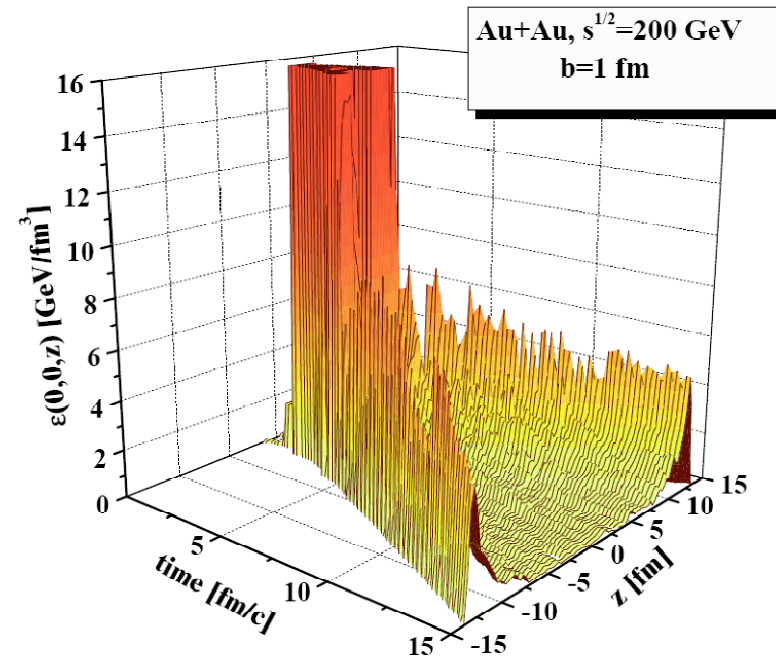
But at **RHIC** recreation of J/Ψ by $D+Dbar$ annihilation is **strong!**

Modelling of the QGP melting scenario in HSD

**Energy density $\epsilon(x=0,y=0,z;t)$ from HSD
for Pb+Pb collisions at 160 A GeV**



**Energy density $\epsilon(x=0,y=0,z;t)$ from HSD
for Au+Au collisions at 21300 A GeV**



Threshold energy densities:

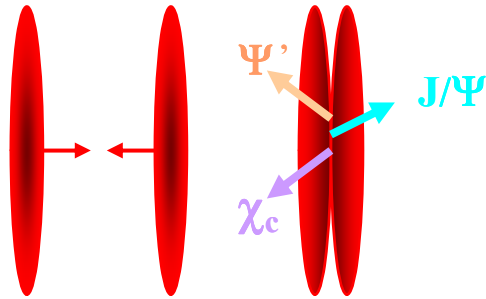
J/ Ψ melting: $\epsilon(\text{J}/\Psi)=16 \text{ GeV}/\text{fm}^3$

χ_c melting: $\epsilon(\chi_c)=2 \text{ GeV}/\text{fm}^3$

Ψ' melting: $\epsilon(\Psi')=2 \text{ GeV}/\text{fm}^3$

[Olena Linnyk et al.,
nucl-th/0612049, NPA 786 (2007) 183]

,Local' energy density ϵ versus Bjorken energy density ϵ_{Bj}



- transient time for central Au+Au at 200 GeV:
 $t_r \sim 2R_A/\gamma_{cm} \sim 0.13 \text{ fm}/c$
- c-cbar formation time:
 $\tau_C \sim 1/M_T \sim 1/4\text{GeV} \sim 0.05 \text{ fm}/c < t_r$
- c-cbar pairs are produced in the **initial hard NN collisions** in time period t_r

- Bjorken energy density:

$$\epsilon_{Bj} = \frac{1}{\tau} \cdot \frac{1}{A_{\perp}} \frac{dE_T}{dy}$$

A_T is the nuclei transverse overlap area
 τ is the formation time of the medium

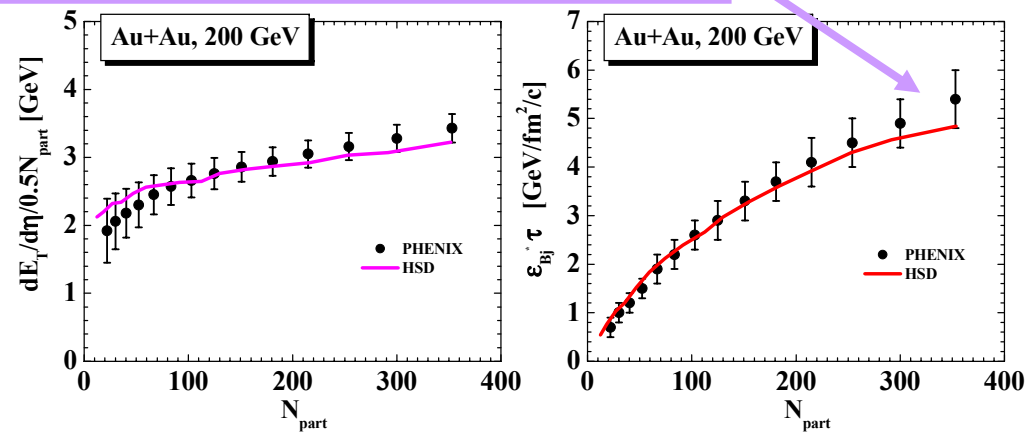
• at RHIC $\epsilon_{Bj} \tau \sim 5 \text{ GeV}/\text{fm}^2/c$

,Local' energy density ϵ during

transient time t_r :

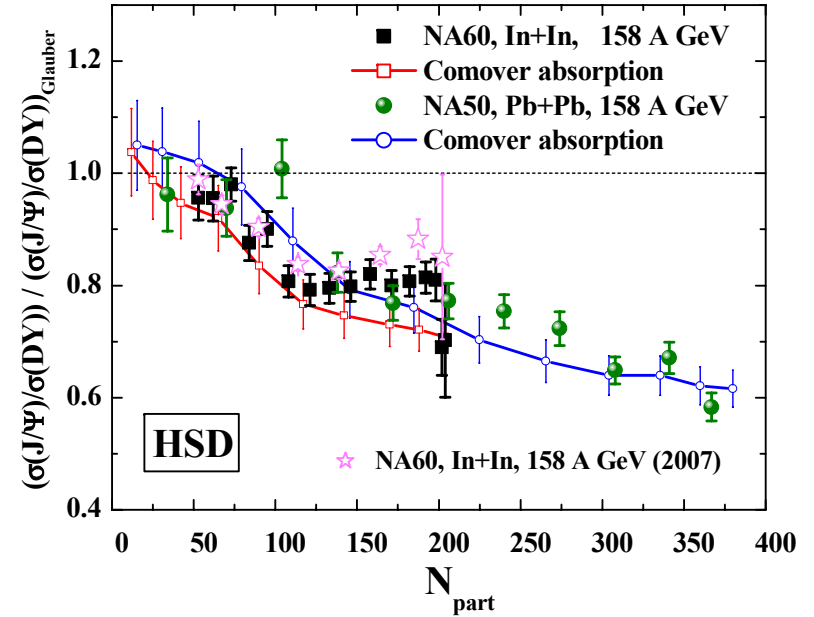
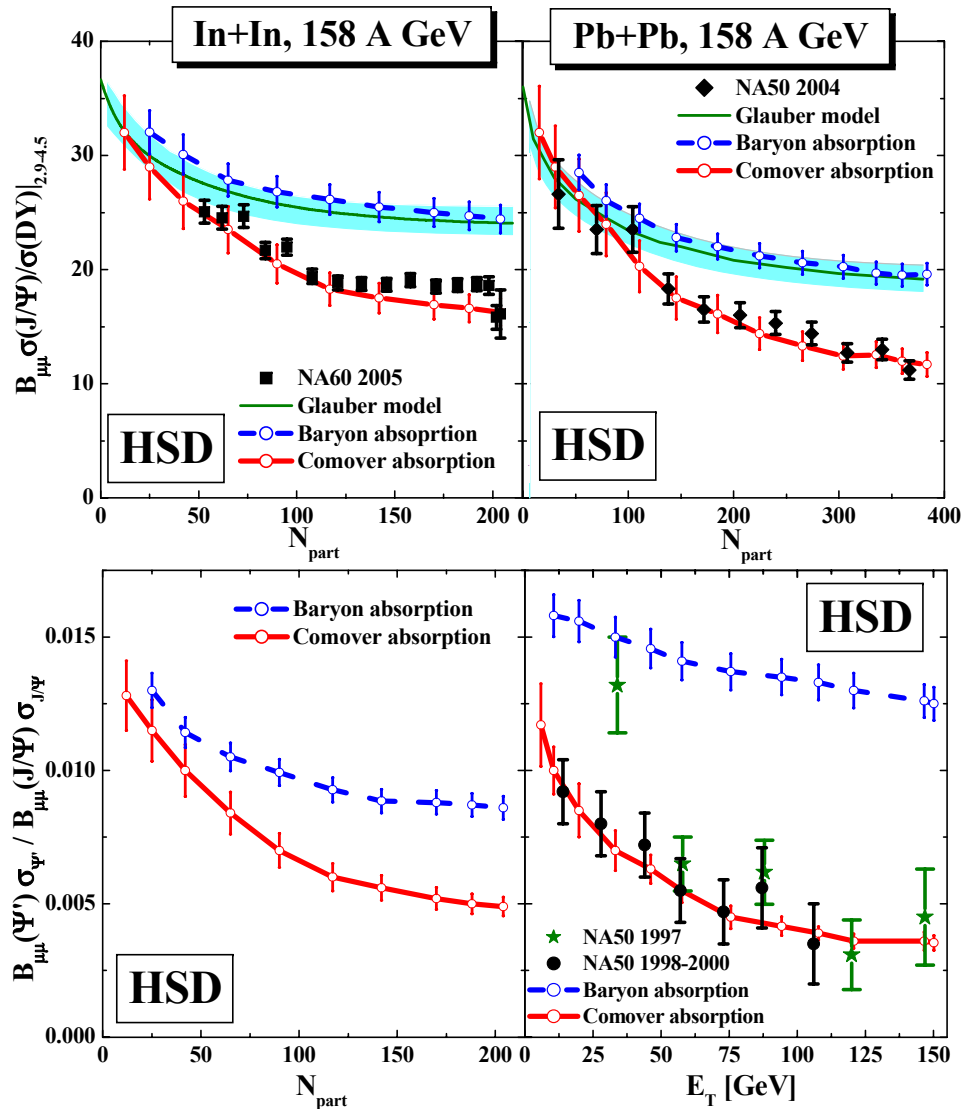
$$\epsilon \sim 5[\text{GeV}/\text{fm}^2/c] / [0.13 \text{ fm}/c] \\ \sim 30 \text{ GeV}/\text{fm}^3$$

accounting τ_C : $\epsilon \sim 28 \text{ GeV}/\text{fm}^3$



- ✓ HSD reproduces PHENIX data for Bjorken energy density very well
- ✓ HSD results are consistent with simple estimates for the energy density

J/Ψ and Ψ' suppression in In+In and Pb+Pb at SPS: Comover absorption

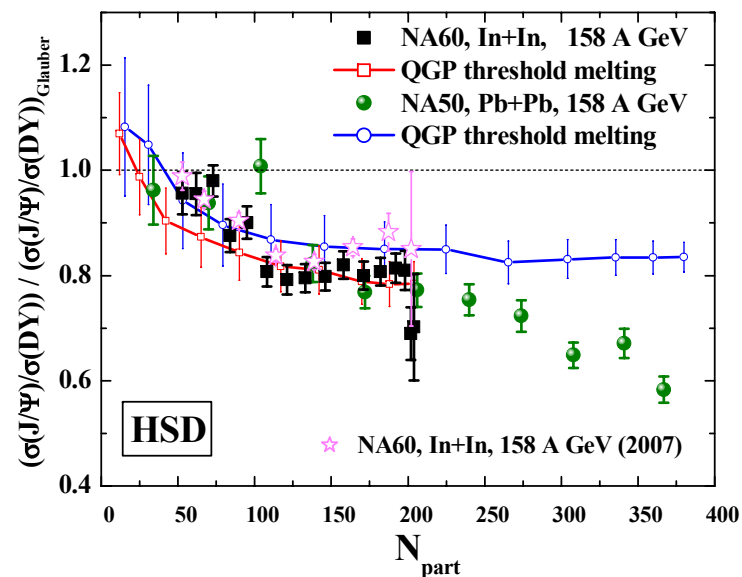
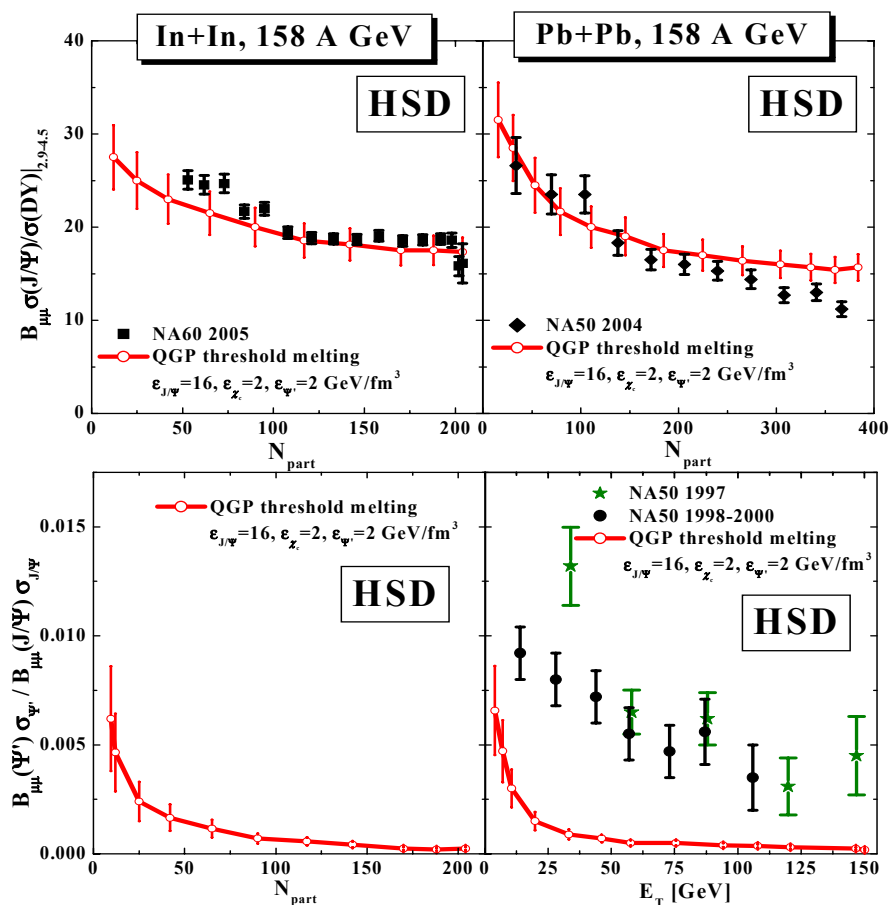


- Exp. data (NA50/NA60) for Pb+Pb and In+In at 160 A GeV are consistent with the **comover absorption model for the same set of parameters!**

[Olena Linnyk et al.,
nucl-th/0612049, NPA 786 (2007) 183]

J/Ψ and Ψ' suppression in In+In and Pb+Pb at SPS: QGP threshold scenario

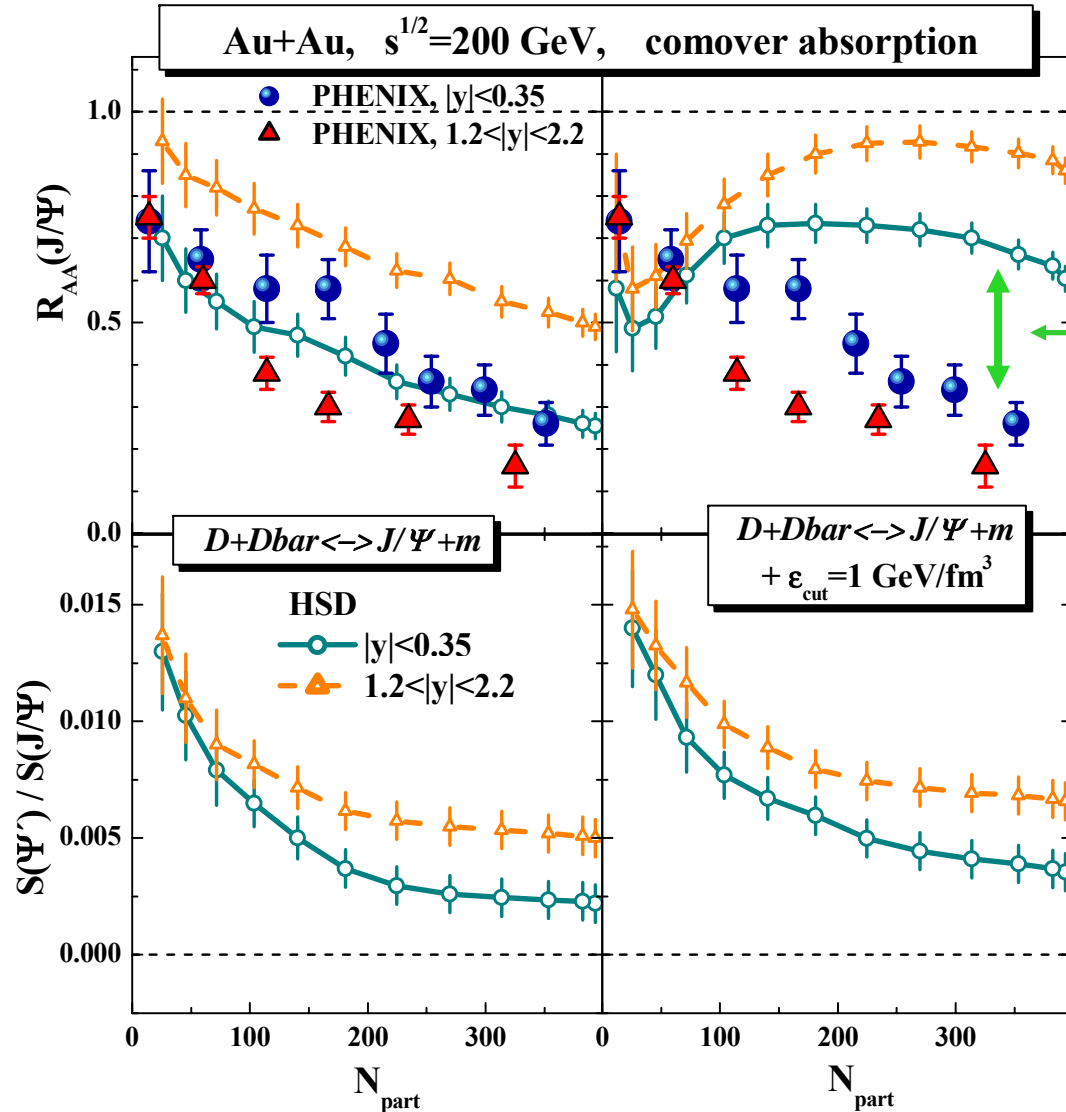
Set 1: $\epsilon(J/\Psi)=16 \text{ GeV/fm}^3$, $\epsilon(\chi_c)=2 \text{ GeV/fm}^3$, $\epsilon(\Psi')=2 \text{ GeV/fm}^3$



- **Set 1:** QGP threshold melting scenario with dissociation energy densities $\epsilon(J/\Psi)=16 \text{ GeV/fm}^3$, $\epsilon(\chi_c)=2 \text{ GeV/fm}^3$, $\epsilon(\Psi')=2 \text{ GeV/fm}^3$ shows too strong Ψ' absorption which contradicts to the NA50 data!

[Olena Linnyk et al.,
nucl-th/0612049, NPA 786 (2007) 183]

J/Ψ and Ψ' suppression in Au+Au at RHIC: Comover absorption

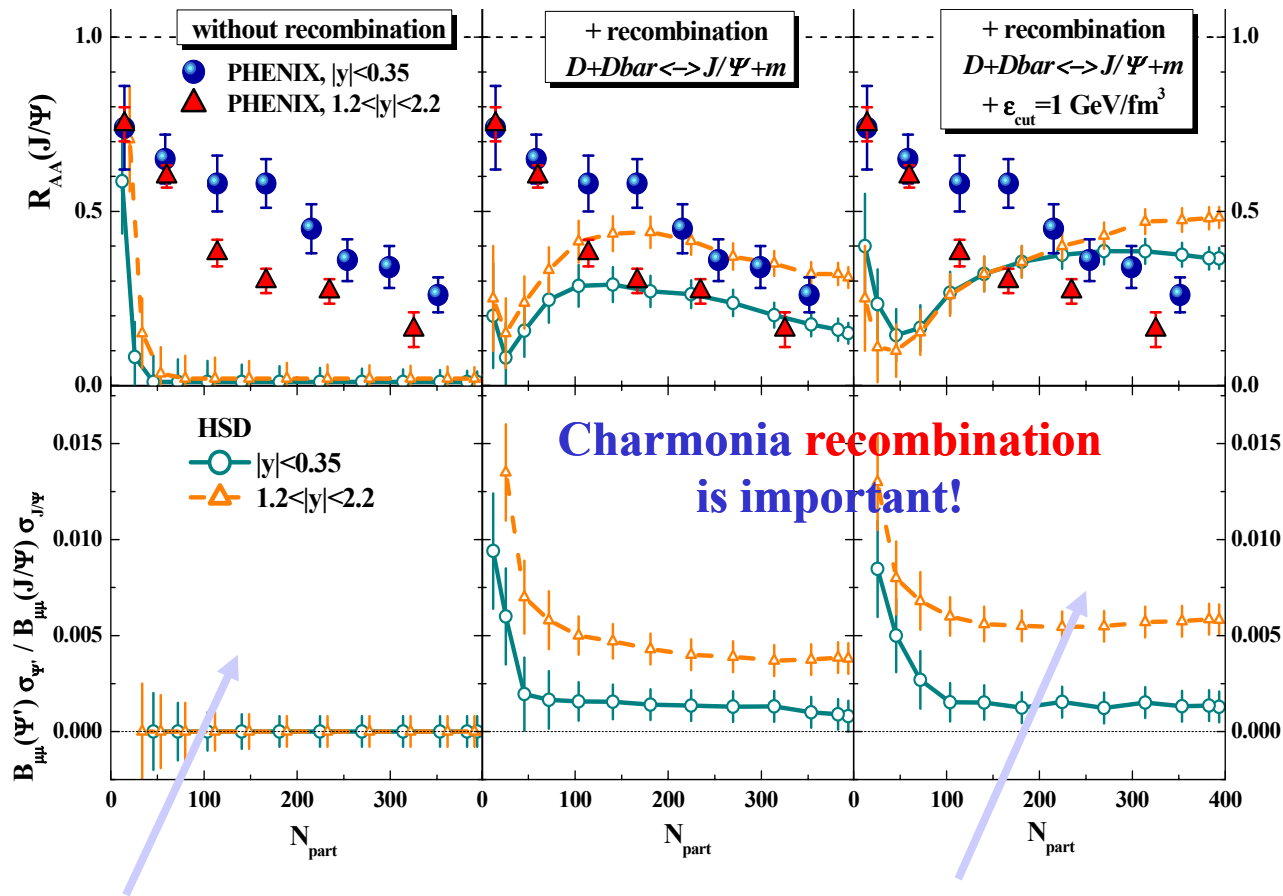


Energy density cut
 $\epsilon_{cut} = 1 \text{ GeV}/\text{fm}^3$ reduces the
 meson comover absorption

||
 Space for partonic effects

In the comover scenario
 the J/Ψ suppression at
mid-rapidity is stronger
than at forward rapidity,
 unlike the data!

J/Ψ and Ψ' suppression in Au+Au at RHIC: QGP threshold scenario

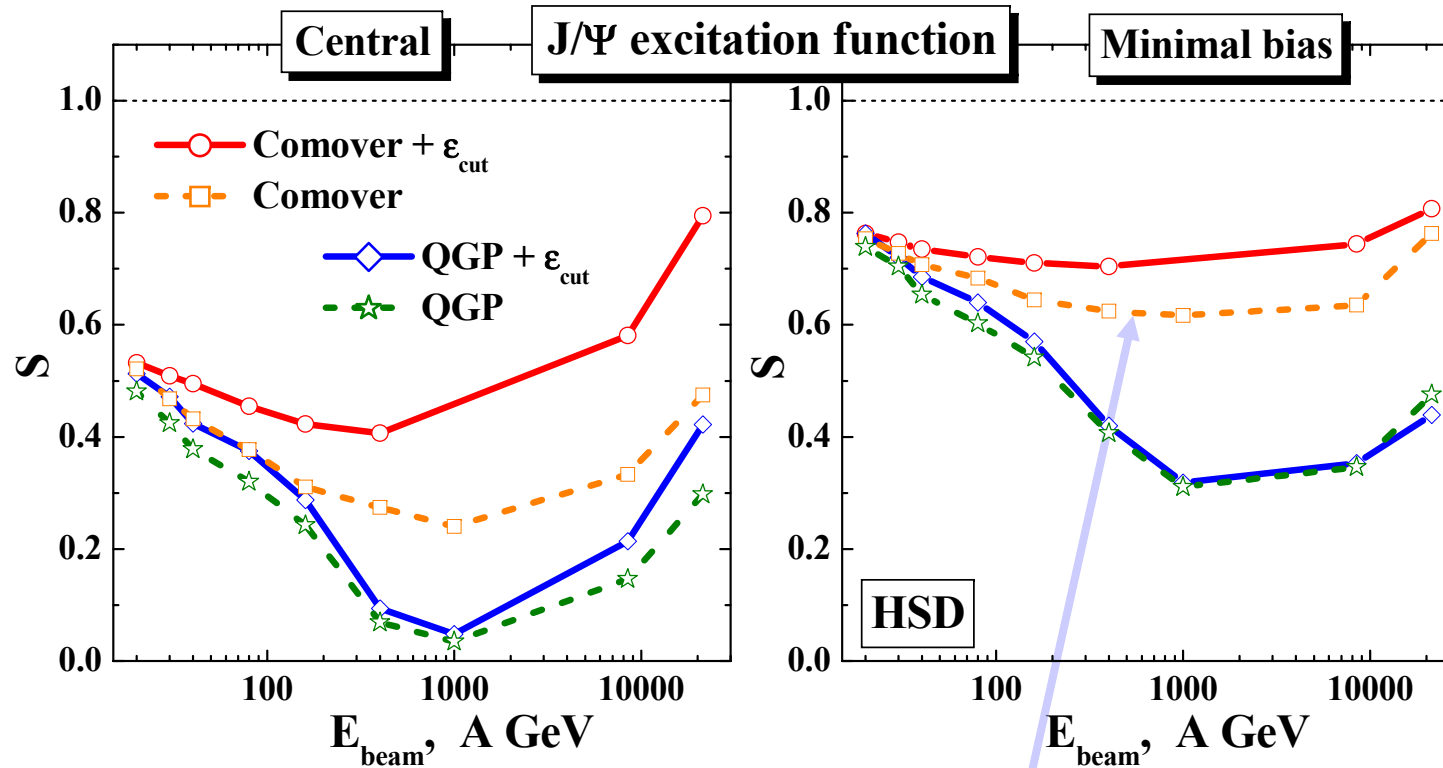


Satz's model: complete dissociation of initial J/Ψ and Ψ' due to the huge local energy densities !

Energy density cut $\epsilon_{\text{cut}}=1 \text{ GeV/fm}^3$ reduces the meson comover absorption, however, D+Dbar annihilation can not generate enough charmonia, especially for peripheral collisions!

QGP threshold melting scenario is ruled out by PHENIX data!

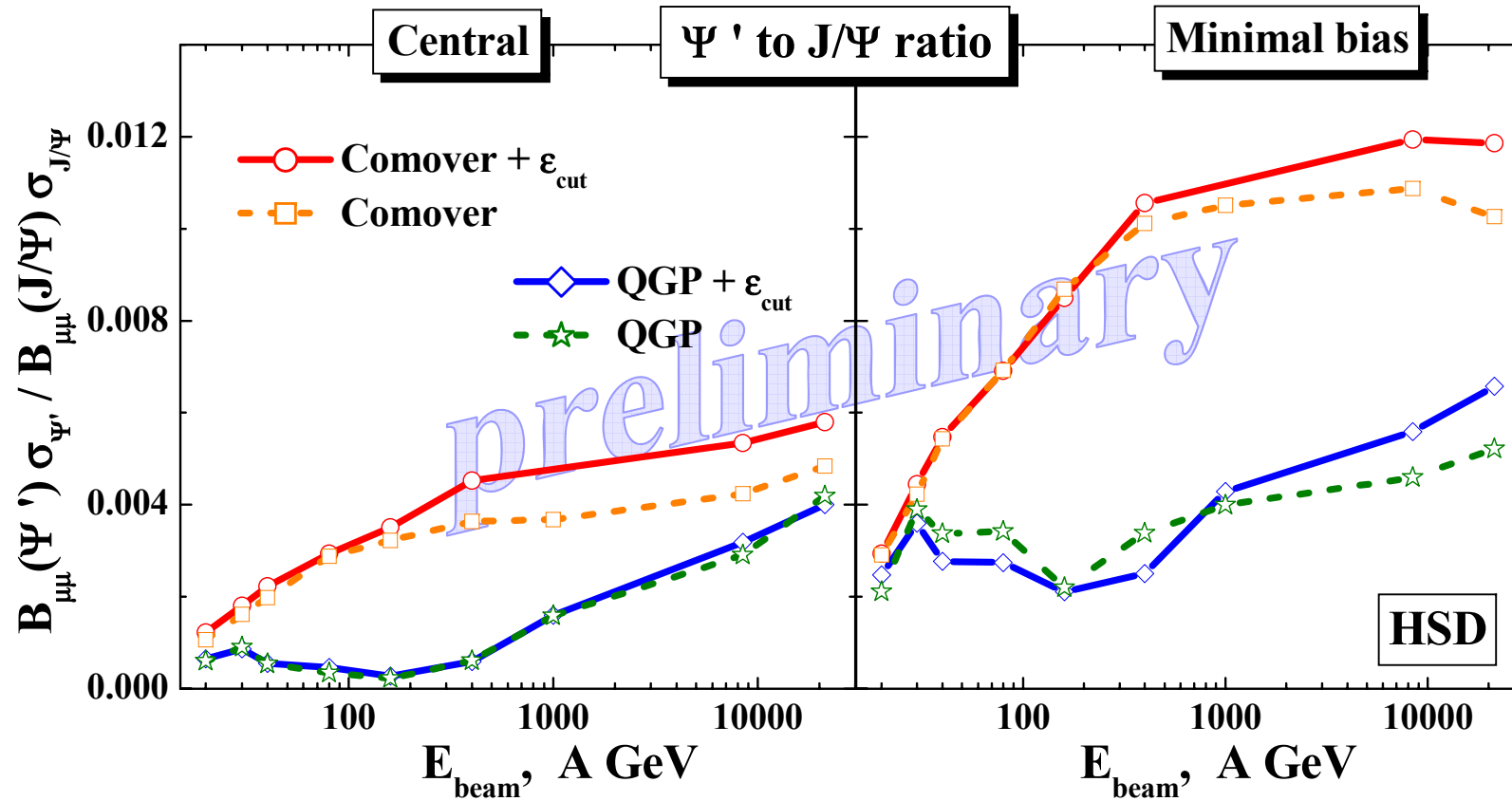
J/Ψ excitation function



- **Comover** reactions in the hadronic phase give almost a constant suppression; pre-hadronic reactions lead to a larger recreation of charmonia with E_{beam}

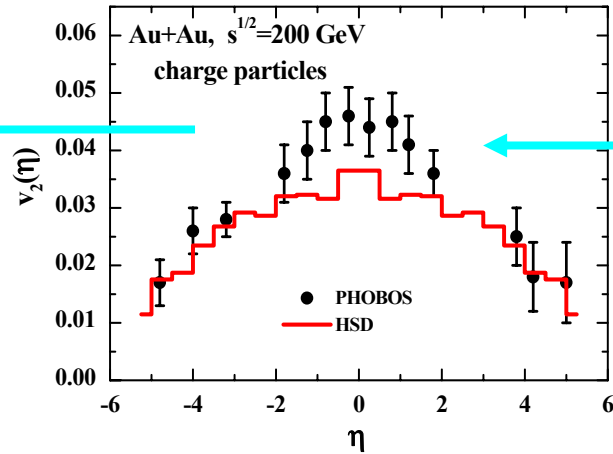
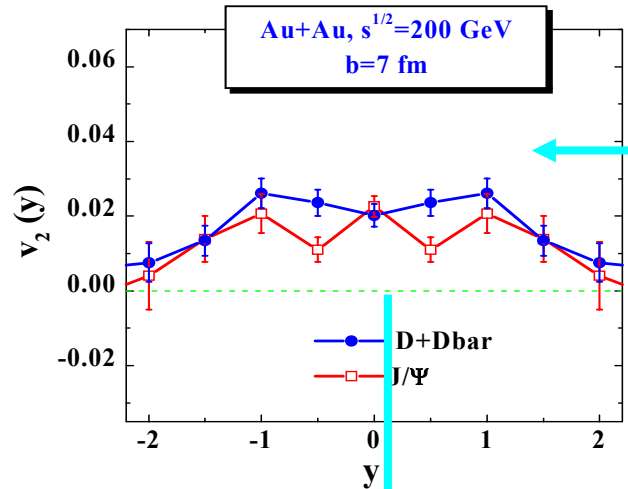
- The **J/Ψ melting scenario** with hadronic comover recreation shows a **maximum suppression** at $E_{\text{beam}} = 1$ A TeV; **exp. data ?**

Ψ' excitation function

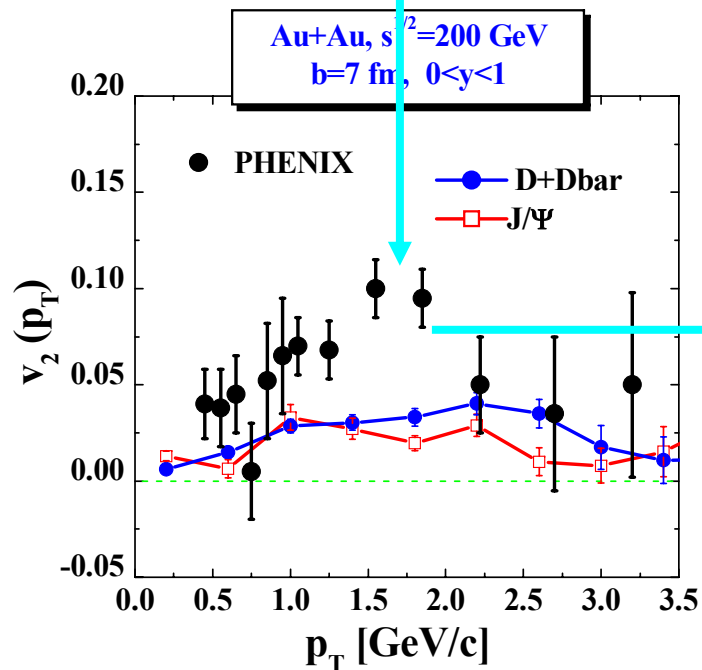


Ψ' suppression provides independent information
on absorption versus recreation mechanisms !

HSD: v_2 of D+Dbar and J/ Ψ from Au+Au versus p_T and y at RHIC



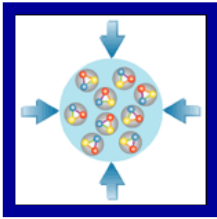
Collective flow from **hadronic interactions** is too low at midrapidity !



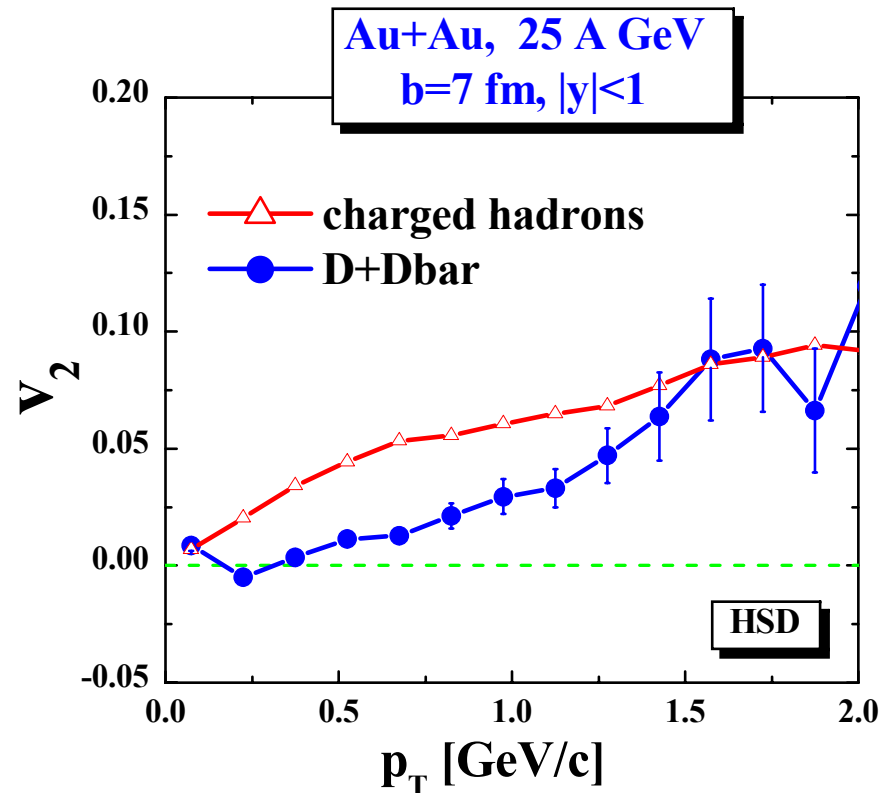
• **HSD: D-mesons and J/ Ψ follow the charged particle flow** => **small $v_2 < 3\%$**

• **Exp. data at RHIC show large collective flow of D-mesons up to $v_2 \sim 10\%$!**

=> **strong initial flow of non-hadronic nature!** (cf. AMPT by Che-Ming Ko)



HSD predictions for CBM - elliptic flow at 25 A GeV



Challenge for CBM!

• **HSD**: D-mesons and J/ Ψ follow the charged particle flow
 \Rightarrow **small v_2**

AMPT with string melting shows much **stronger v_2** !
(Che-Ming Ko)

Possible observation at CBM:
strong initial flow of D-mesons and J/ Ψ due to **partonic interactions!**

Summary I.

- **J/Ψ probes early stages** of fireball and HSD is the tool to model it.
- **Comover absorption and threshold melting** both reproduce J/Ψ survival in Pb+Pb as well as in In+In at SPS, while **Ψ'/J/Ψ data** appear to be in conflict with the melting scenario.
- **Comover absorption and threshold melting** fail to describe the RHIC data at $s^{1/2}=200$ GeV for Au+Au at mid- and forward-rapidities simultaneously.
- STAR data on **v₂ of high p_T** charged hadrons and charm D mesons are NOT reproduced in the hadron-string picture => evidence for a huge plasma pressure ?!

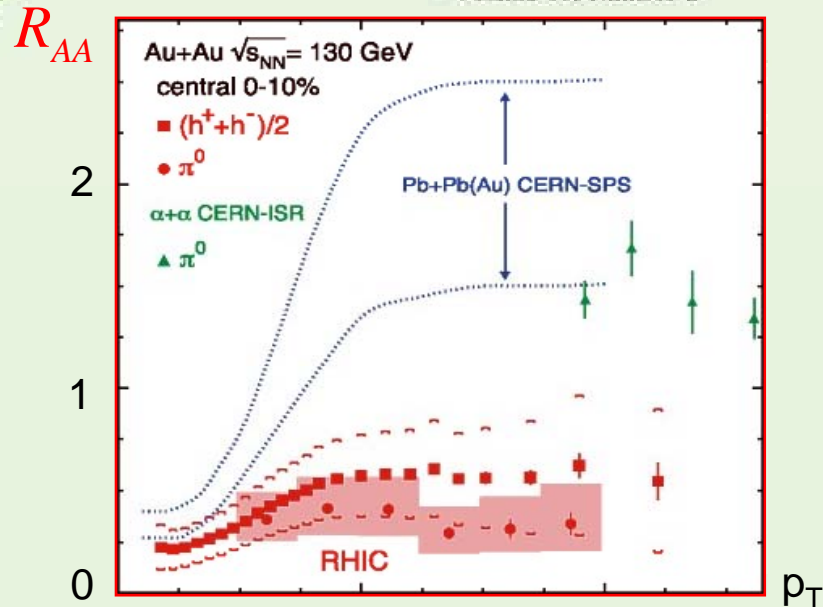
High p_T suppression of hadrons

PHYSICAL
REVIEW
LETTERS

$$R_{AA} = \frac{d\sigma^{AA}}{dyd^2p_T} / \frac{N^{bin} d\sigma^{pp}}{dyd^2p_T}$$

14 January 2002

Volume 88, Number 2

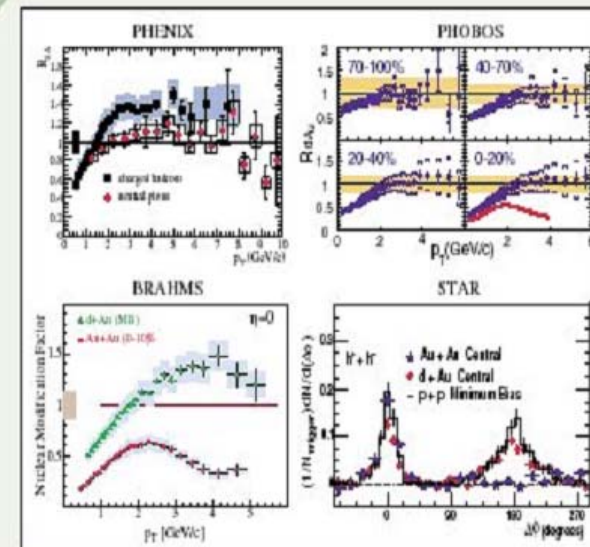


PHYSICAL
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Articles published week ending

15 AUGUST 2003

Volume 91, Number 7



High p_T suppression of hadrons: exp. observables

Nuclear Modification Factor:

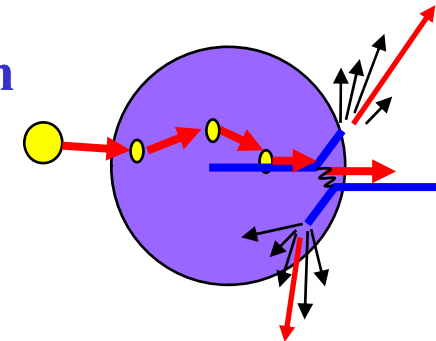
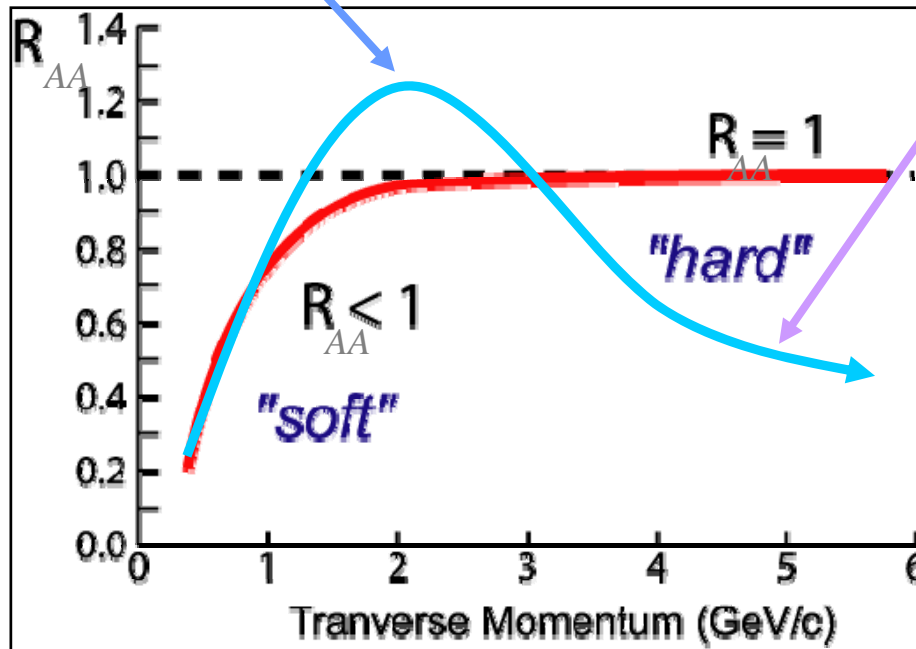
$$R_{AA}(p_T) \propto \frac{\sigma_{AA}}{\langle N_{coll} \rangle \cdot \sigma_{pp}}$$

Probe: initial + final state effects

Cronin effect

“jet quenching”

hadronic + partonic suppression



$R_{AA} < 1$ in regime of **soft physics**

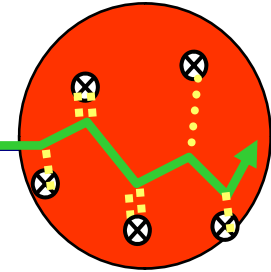
$R_{AA} = 1$ at high- p_T where hard scattering dominates

$R_{AA} > 1$ - Cronin Effect

Suppression:

$R_{AA} < 1$ at high- p_T “jet quenching”

Cronin effect at RHIC (HSD)



Cronin effect: initial state **semi-hard gluon radiation** increases p_T spectra already in p+A or d+A

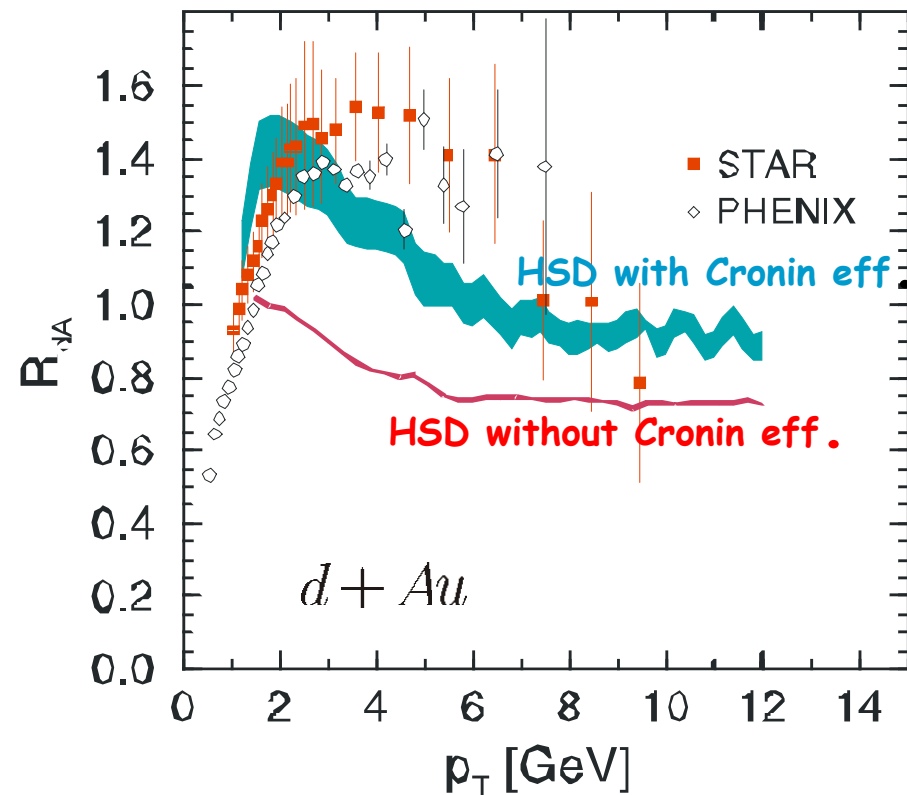
Dynamical modeling of the Cronin effect in HSD:

$$R_{dA}(p_T) = \frac{1/N_{dA}^{\text{event}} \cdot d^2N_{dA}/dydp_T}{\langle N_{\text{coll}} \rangle / \sigma_{pp}^{\text{inelas}} \cdot d\sigma_{pp}/dydp_T}$$

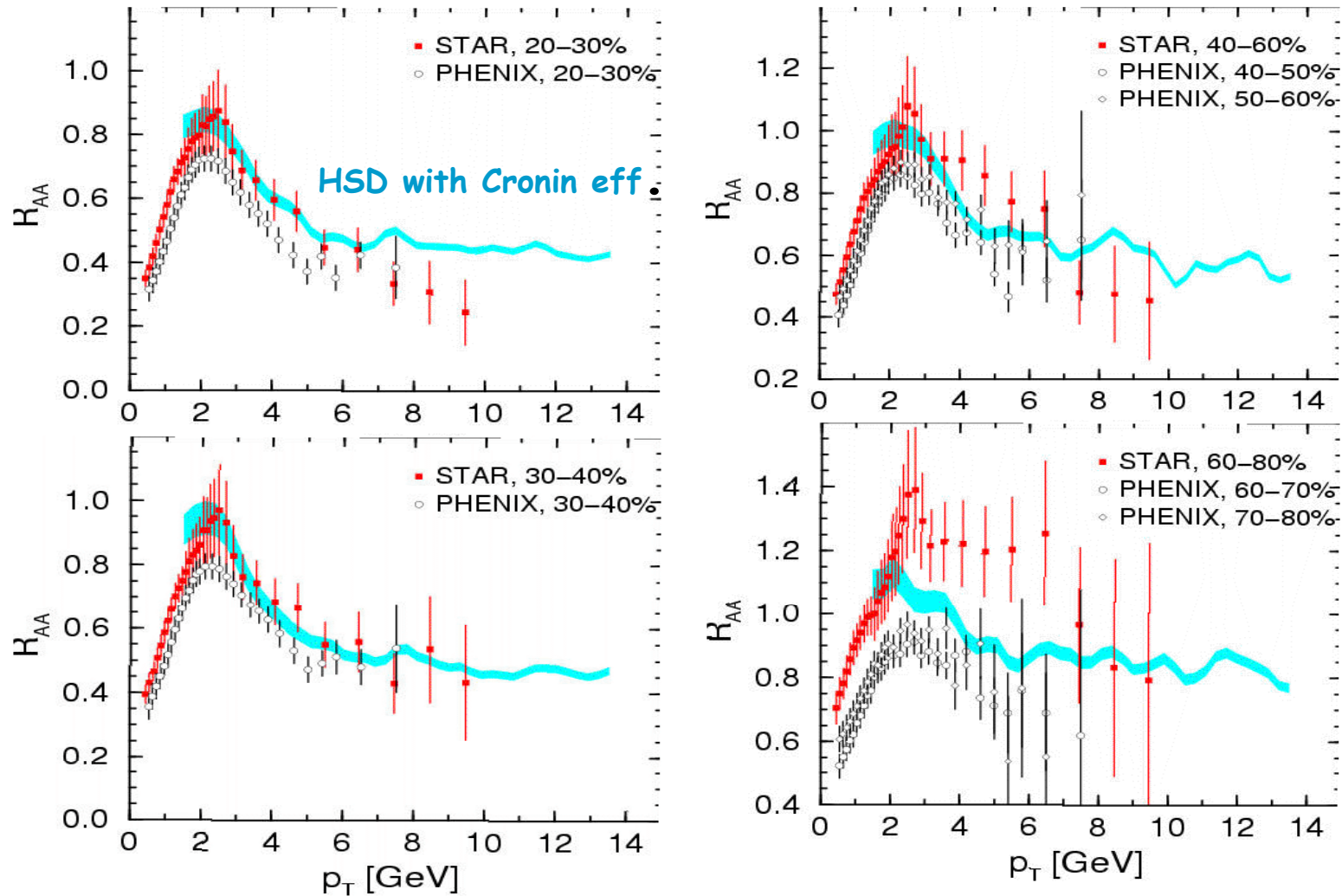
$$\langle k_T^2 \rangle_{AA} = \langle k_T^2 \rangle_{PP} (1 + a N_{\text{Prev}})$$

N_{Prev} = number of previous collisions (dynamically calculated for each hadron!)

parameter $a = 0.25 - 0.4$



High p_T suppression in non-central Au+Au (HSD)



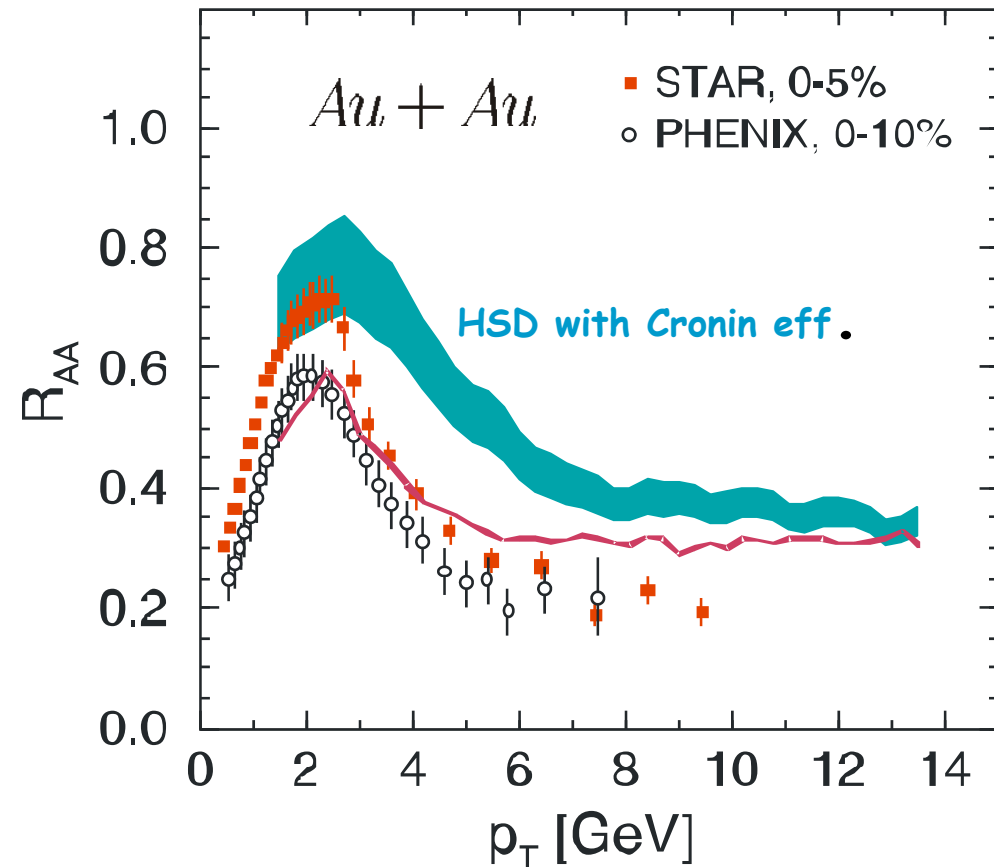
- Hadron-string model with Cronin effect provides ~ enough high p_T suppression for non-central Au+Au

Cronin effect in central Au+Au at RHIC (HSD)

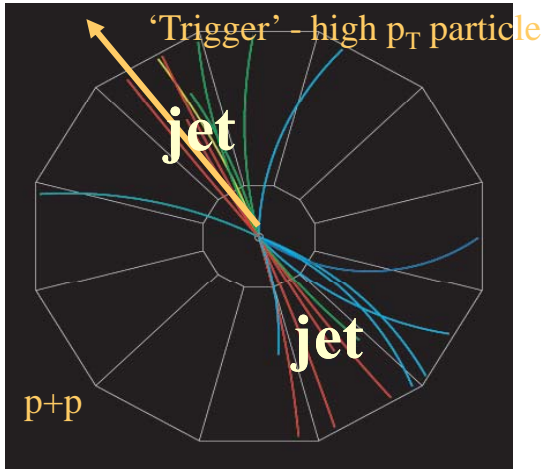
$$R_{AA}(p_T) = \frac{1/N_{AA}^{\text{event}} \cdot d^2N_{AA}/dydp_T}{\langle N_{\text{coll}} \rangle / \sigma_{pp}^{\text{inelas}} \cdot d\sigma_{pp}/dydp_T}$$

- Hadron-string model **doesn't** provide enough high p_T suppression for central Au+Au

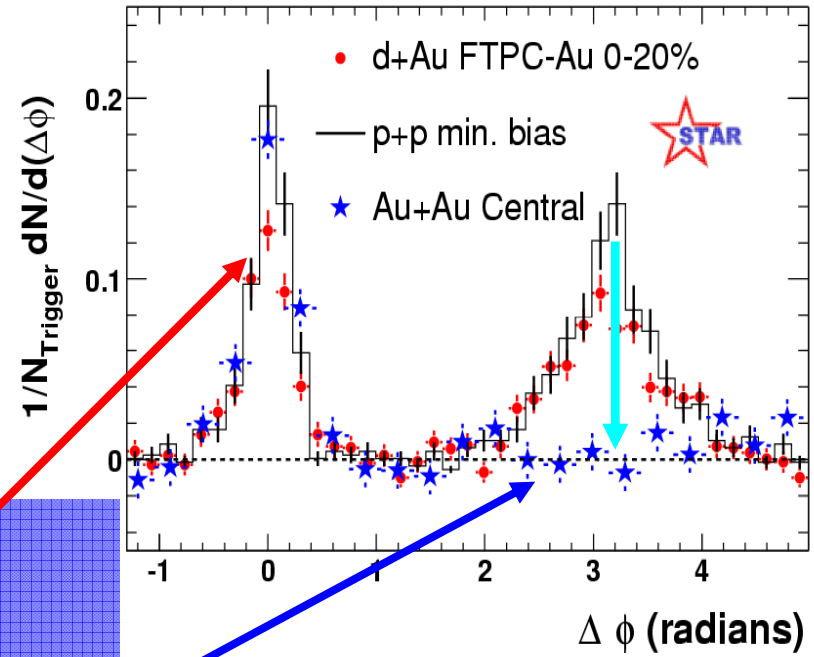
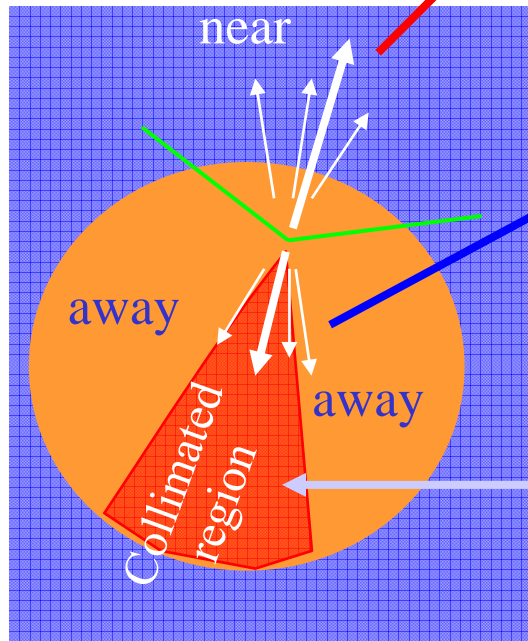
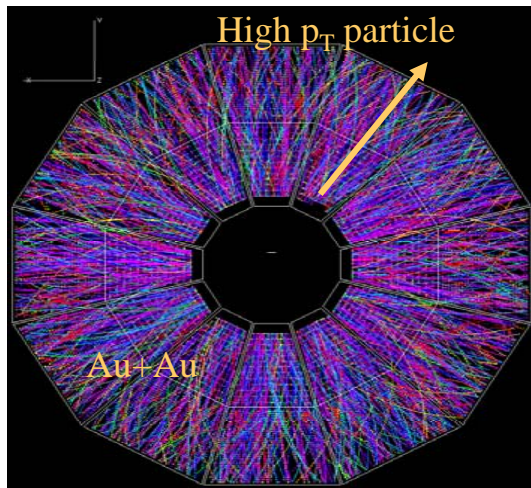
→ **Extra suppression :**
from early partonic phase ?!



Jet suppression: $dN/d\phi$



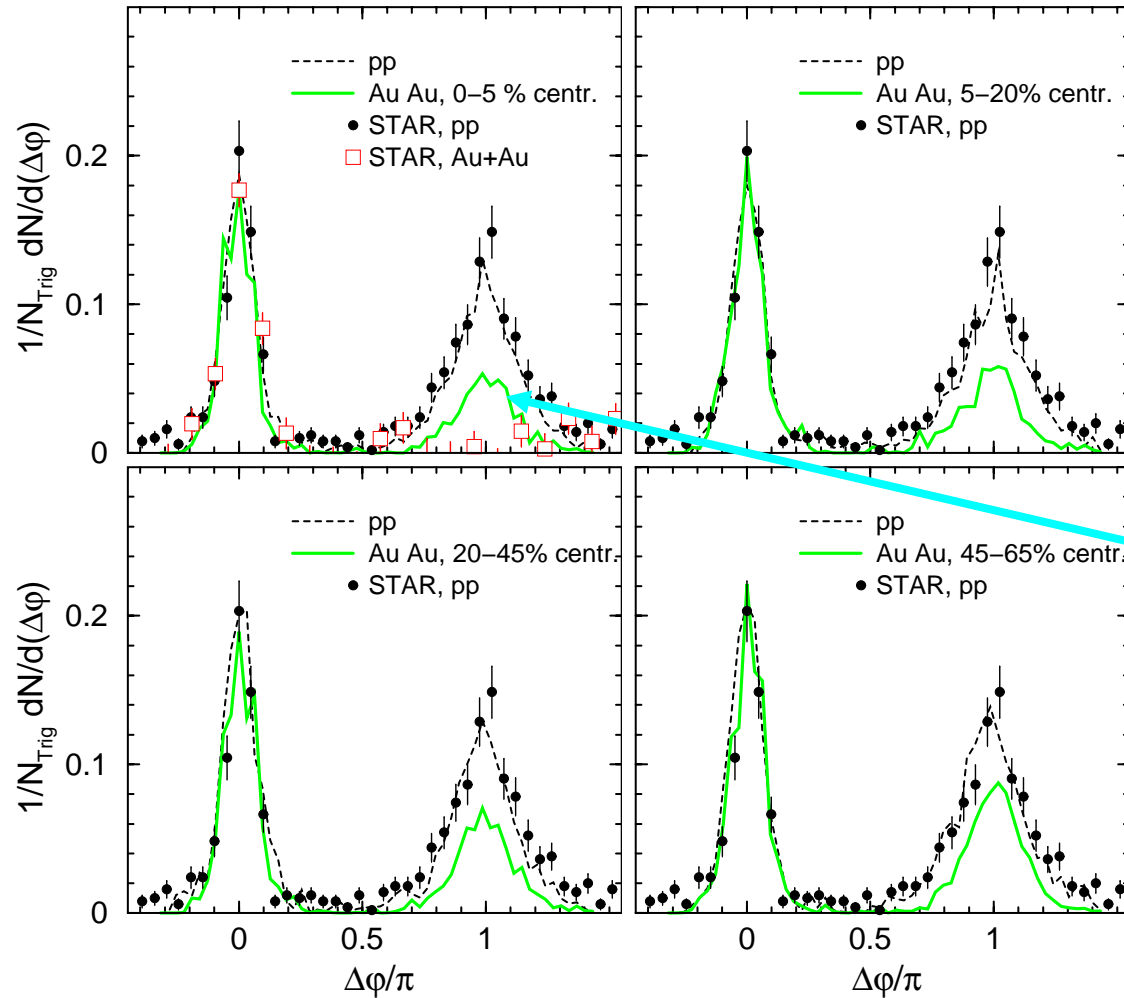
$4 < p_T(\text{trigger}) < 6 \text{ GeV}/c$
 $p_T(\text{assoc.}) > 2 \text{ GeV}/c$



QGP suppression ?!

jet energy loss in the medium

Jet suppression: $dN/d\phi$



- The jet angular correlations for **pp** are fine!
- Au+Au: The near-side jet angular correlation is well described, but the **suppression of the far-side jet is too low!**

Summary II.

- **Cronin effect** (initial state semi-hard gluon radiation) increases p_T spectra already in p+A or d+A
 - The **attenuation of high p_T -hadrons (R_{AA})** is well reproduced in the hadron-string approach for non-central Au+Au collisions at top RHIC energies, however, underestimated in **central** Au+Au collisions.
 - The jet angular correlations for **pp** are fine, however, the **suppression of the ,far-side‘ jet** is underestimated in central Au+Au collisions at $s^{1/2} = 200$ GeV
- ➔ **System interacts more strongly in the early phase than hadron-like matter !**

Outlook

- A deconfined phase is clearly reached at RHIC, but a **theory** having the **relevant/proper degrees of freedom (quarks and gluons)** in this regime is **needed** to study its transport properties.

HSD → **Parton-Hadron-String-Dynamics**
(PHSD)

PHSD –
transport description of the **partonic** and **hadronic** phase

Thanks

Olena Linnyk (FIAS)

Wolfgang Cassing (Giessen Uni.)

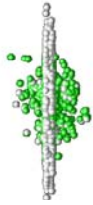
Volodymyr Konchakovski (Frankfurt Uni.)

Andrej Kostyuk (FIAS)

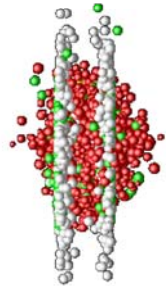
Laura Tolos (FIAS)

+ Horst Stöcker (GSI, FIAS, Frankfurt Uni.)

Outlook: PHSD - basic concepts



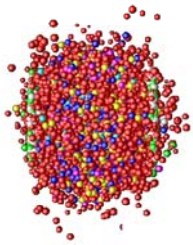
Initial A+A collisions – HSD: string formation and decay to pre-hadrons



Fragmentation of pre-hadrons into quarks: using the quark spectral functions from the Dynamical QuasiParticle Model (DQPM) approximation to QCD

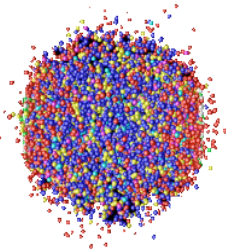
DQPM: Peshier, Cassing, PRL 94 (2005) 172301;
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

Partonic phase: quarks and gluons (= ,dynamical quasiparticles‘) with off-shell spectral functions (width, mass) defined by DQPM



elastic and inelastic parton-parton interactions: using the effective cross sections from the DQPM

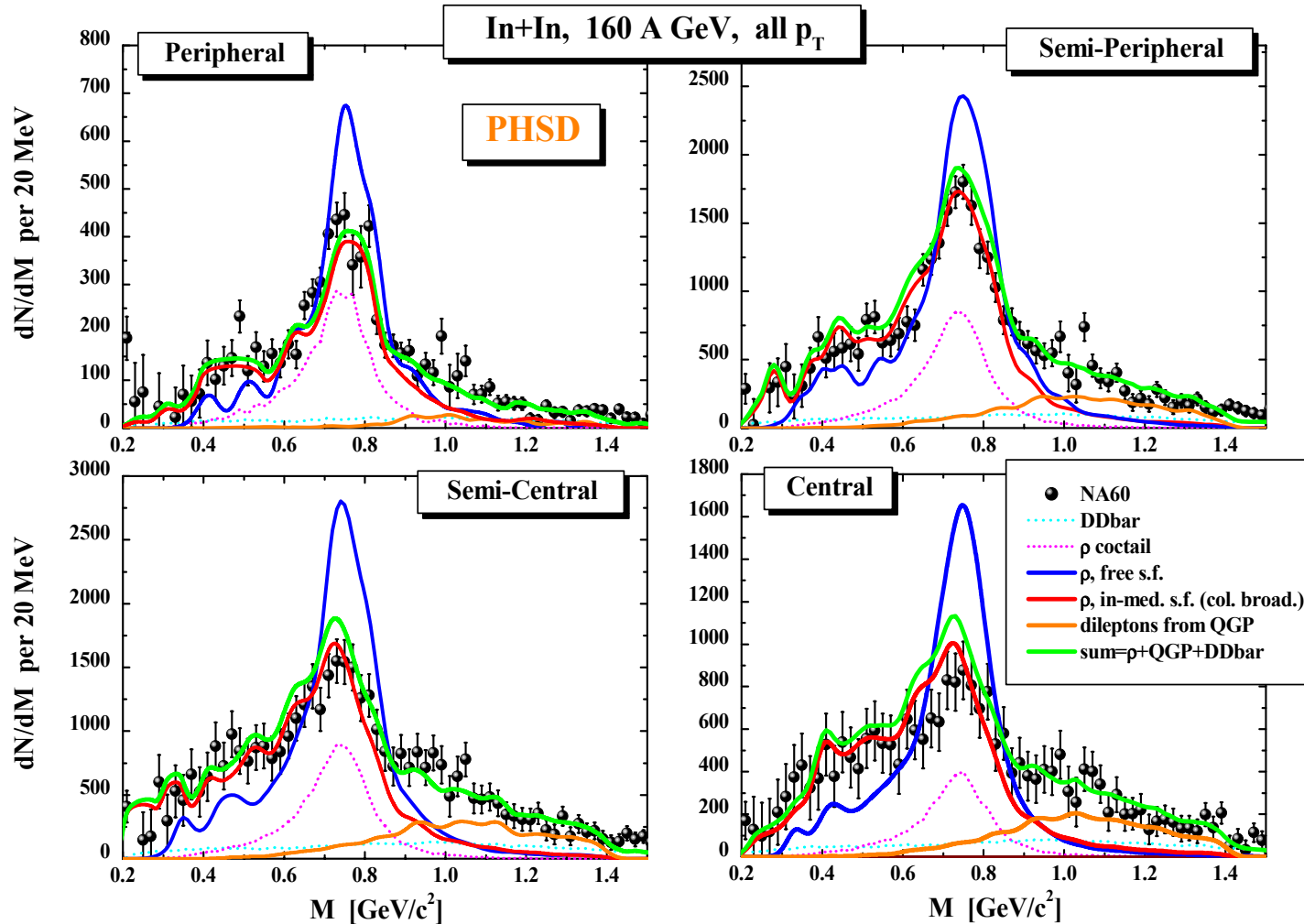
- ✓ $q + q\bar{q}$ (flavor neutral) \Leftrightarrow gluon (colored)
- ✓ gluon + gluon \Leftrightarrow gluon (possible due to large spectral width)
- ✓ $q + q\bar{q}$ (color neutral) \Leftrightarrow hadron resonances



Hadronization: based on DQPM - massive, off-shell quarks and gluons with broad spectral functions hadronize to off-shell mesons and baryons:
gluons \rightarrow $q + q\bar{q}$; $q + q\bar{q} \rightarrow$ meson; $q + q + q \rightarrow$ baryon

Hadronic phase: hadron-string interactions – off-shell HSD

Dilepton radiation from the sQGP – NA60



Conjecture: the sQGP shows up already at SPS energies !