



ALICE

Measurement of jet spectra in Pb-Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV with ALICE

Salvatore Aiola,
on behalf of the ALICE collaboration
INFN Catania and LBNL

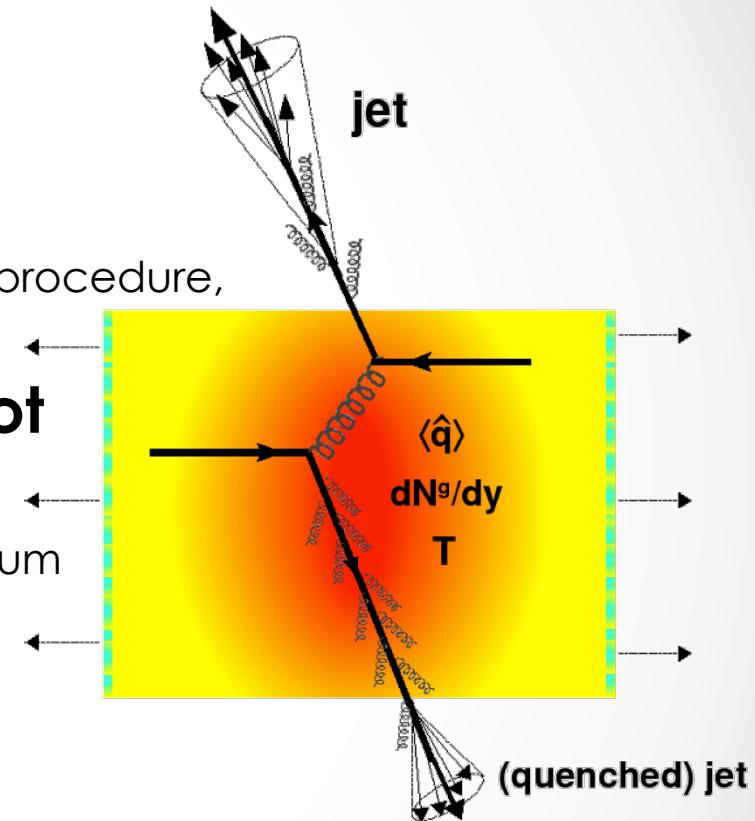


Outline

- Introduction
- Analysis overview
- Jet reconstruction
- Background estimation
- Unfolding
- Jet spectra
- Conclusions & Outlook

Jets in heavy-ion collisions

- Jet: **collimated spray of hadrons**
 - QCD branching of a high p_T parton
 - Subsequent hadronization of fragments
 - Experimentally grouped according to given procedure, *jet algorithm*
- Jets can be used to **probe the hot QCD medium**
 - Observable properties modified by the medium
 - p_T distribution
 - ...and many more



Jets in heavy-ion collisions

- Jet: **collimated spray of hadrons**

- QCD branching of a high p_T parton
- Subsequent hadronization of fragments
- Experimentally grouped according to given procedure,
jet algorithm

- Jets can be used to **probe the hot QCD medium**

- Observable properties modified by the medium
 - p_T distribution
 - ...and many more

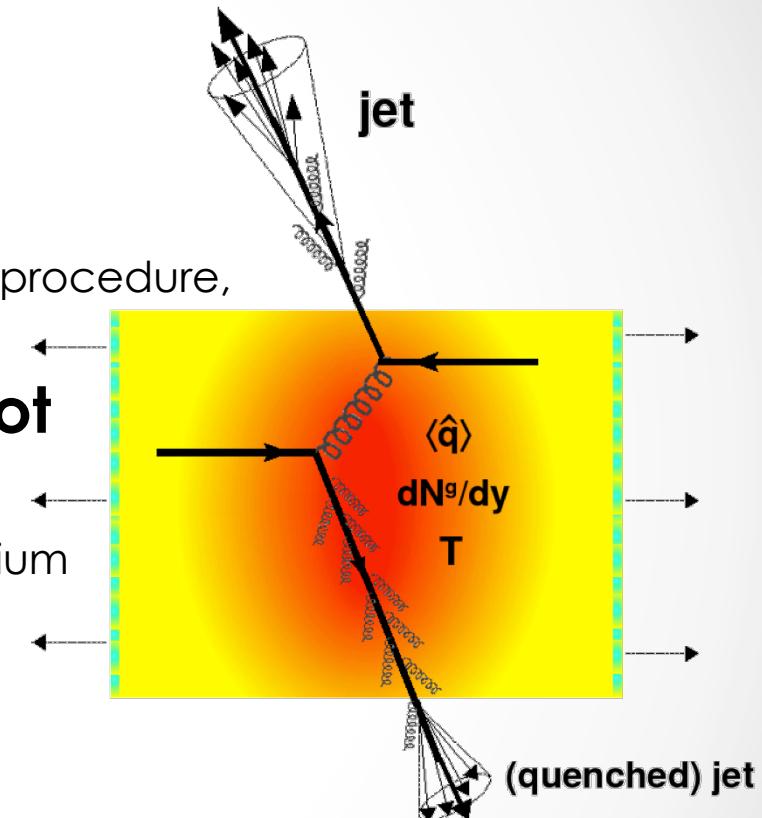
- **Experimentally challenging**

- Huge background given by the underlying event (UE)

1. **Average background**

2. **Combinatorial jets**

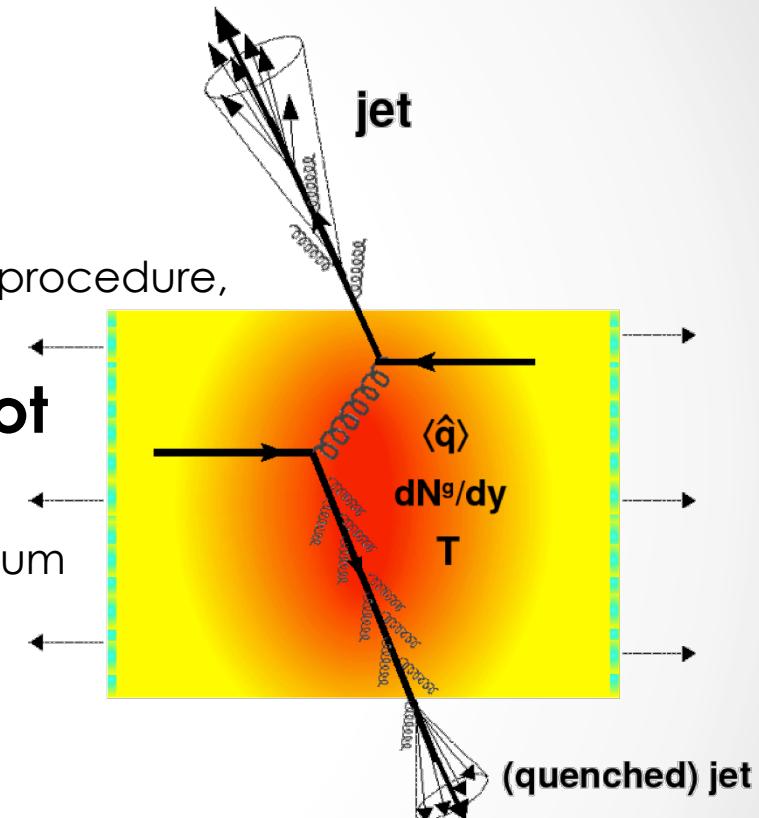
3. **p_T smearing due to background fluctuations**



Jet-by-jet and
event-by-event correction

Jets in heavy-ion collisions

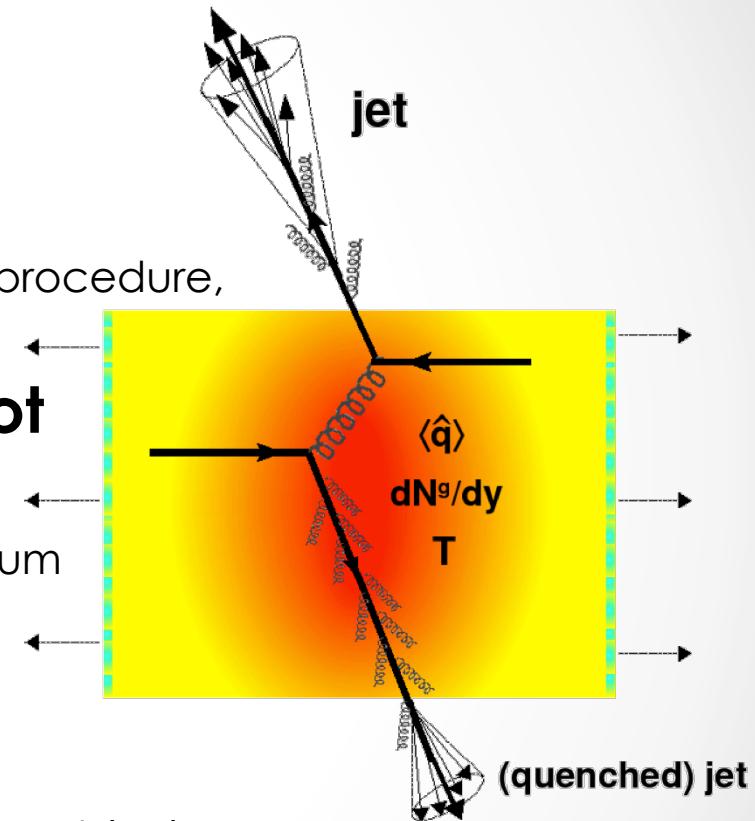
- Jet: **collimated spray of hadrons**
 - QCD branching of a high p_T parton
 - Subsequent hadronization of fragments
 - Experimentally grouped according to given procedure, *jet algorithm*
- Jets can be used to **probe the hot QCD medium**
 - Observable properties modified by the medium
 - p_T distribution
 - ...and many more
- **Experimentally challenging**
 - Huge background given by the underlying event (UE)
 1. **Average background**
 2. **Combinatorial jets**
 3. **p_T smearing due to background fluctuations**



Efficiently removed
by a leading hadron
 p_T requirement

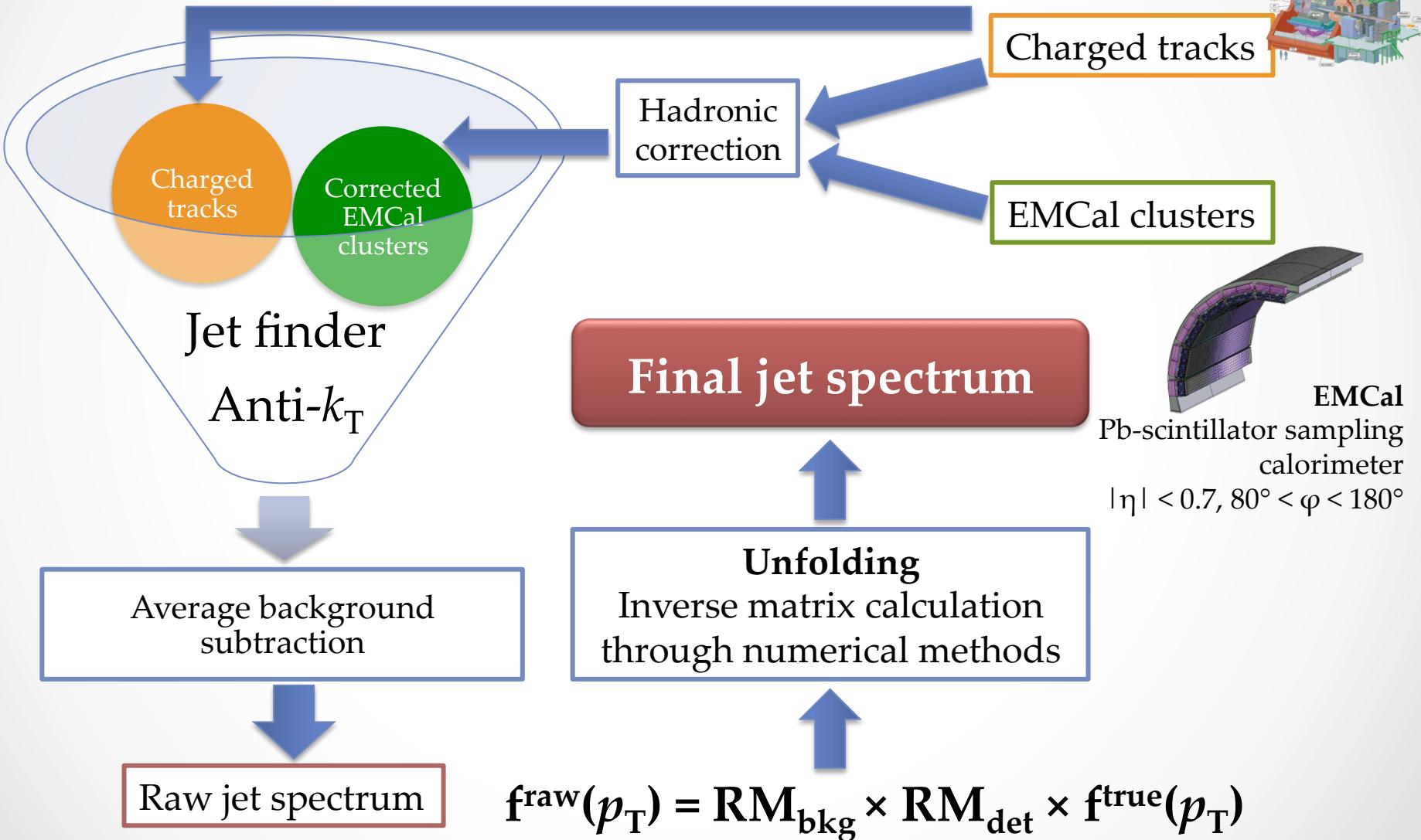
Jets in heavy-ion collisions

- Jet: **collimated spray of hadrons**
 - QCD branching of a high p_T parton
 - Subsequent hadronization of fragments
 - Experimentally grouped according to given procedure, *jet algorithm*
- Jets can be used to **probe the hot QCD medium**
 - Observable properties modified by the medium
 - p_T distribution
 - ...and many more
- **Experimentally challenging**
 - Huge background given by the underlying event (UE)
 1. **Average background**
 2. **Combinatorial jets**
 3. **p_T smearing due to background fluctuations**



Statistical
correction based on
 δp_T distributions

Analysis overview



Jet reconstruction

- Inputs to the jet finder
 - Assumed to be massless
 - **Charged tracks** with $p_T > 150 \text{ MeV}/c$
 - **EMCal clusters** with $E_T > 300 \text{ MeV}/c$ after charged particle correction
- Jet reconstructed using FastJet* package
 - **Infrared- and Collinear-Safe** algorithms
 - Good for comparison with theory
 - **Radii** $R = 0.2, 0.3$
 - **Area cut** $A > 0.6 * \pi R^2$ removes extremes
 - Fiducial cut selects jets **fully contained in the EMCal** acceptance

Sequential recombination

$$d_{ij} = \min(k_{Ti}^{2p}, k_{Tj}^{2p}) \frac{\Delta_{ij}^2}{R^2} \quad d_{iB} = k_{Ti}^{2p}$$

$p = 1$ **k_T algorithm**

Clustering starts from the low p_T particles

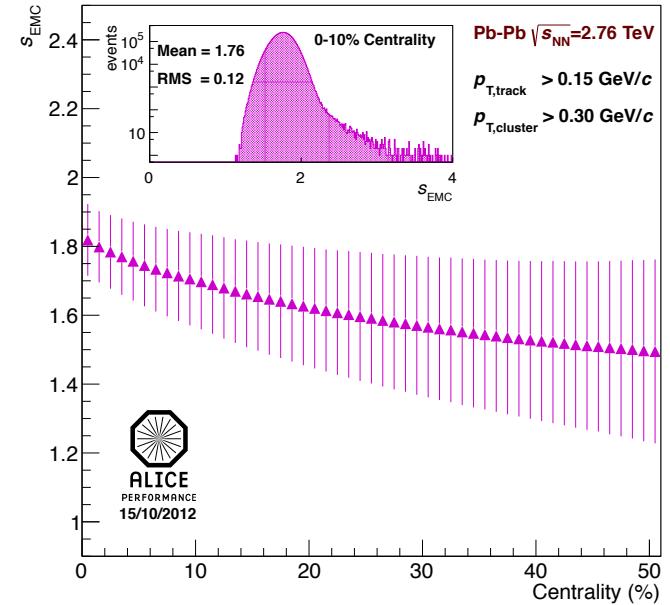
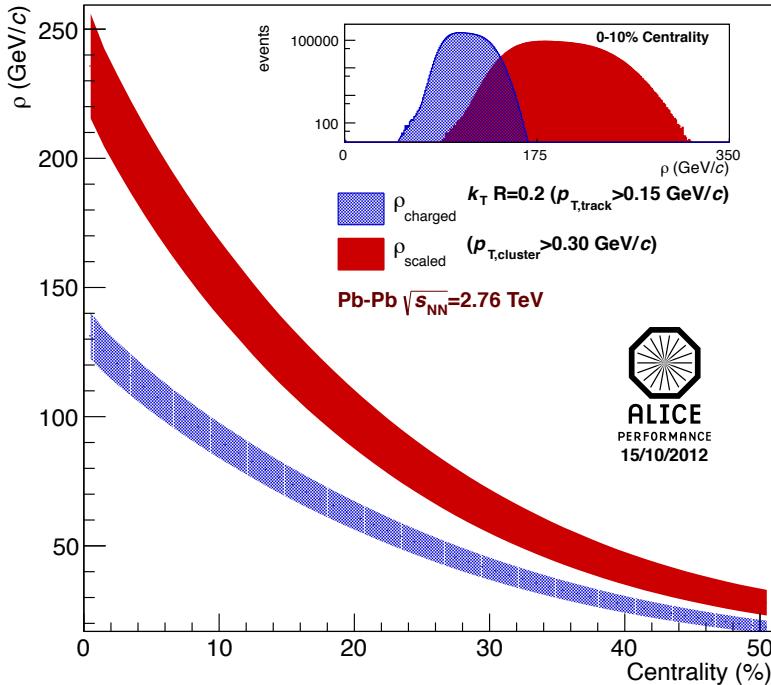
Used for Q calculation

$p = -1$ **Anti- k_T algorithm**

Clustering starts from the highest p_T particle, regular cone-like jet shapes

Used for signal jets

1. Average background

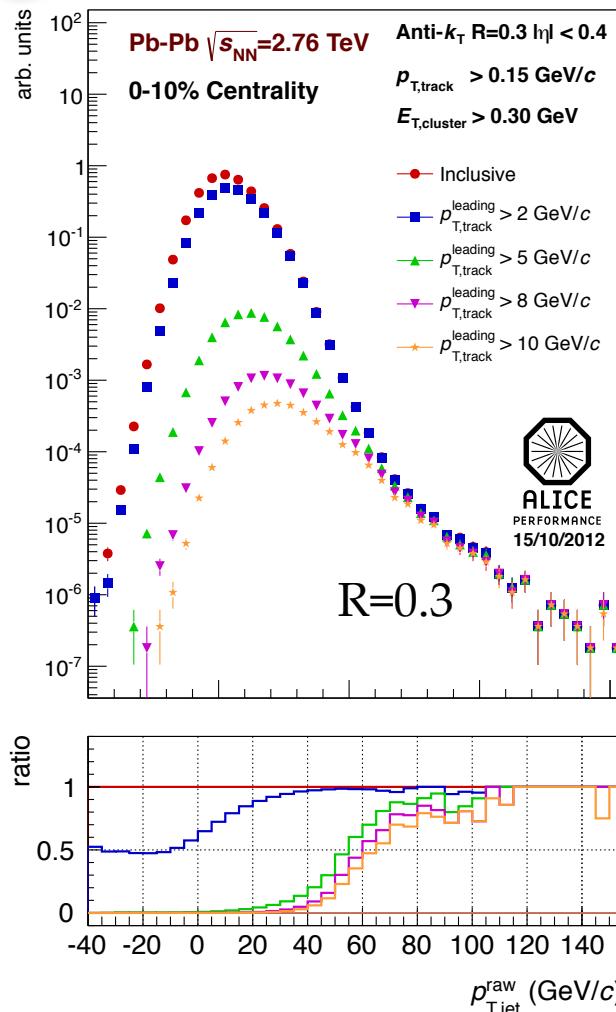
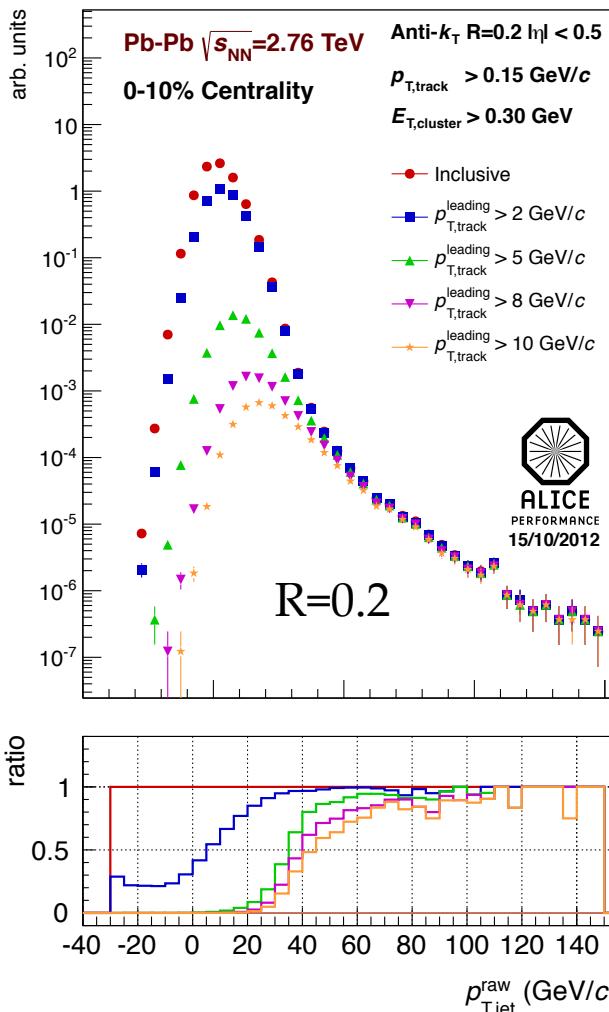


Charged+neutral momentum
density / charged only
momentum density

- Calculated on an **event-by-event basis**
- The median of the density of k_T charged jets gives $\rho_{\text{ch}} = \text{median} \left[\left\{ \frac{p_{Tj}^{k_T}}{A_j^{k_T}} \right\} \right]$
- A parametrization of the scale s_{EMC} factor as a function of centrality is used to obtain $\rho_{\text{ch+ne}} = \rho_{\text{ch}} \times s_{\text{EMC}}$

$$\rho_{\text{ch}} = \text{median} \left[\left\{ \frac{p_{Tj}^{k_T}}{A_j^{k_T}} \right\} \right]$$

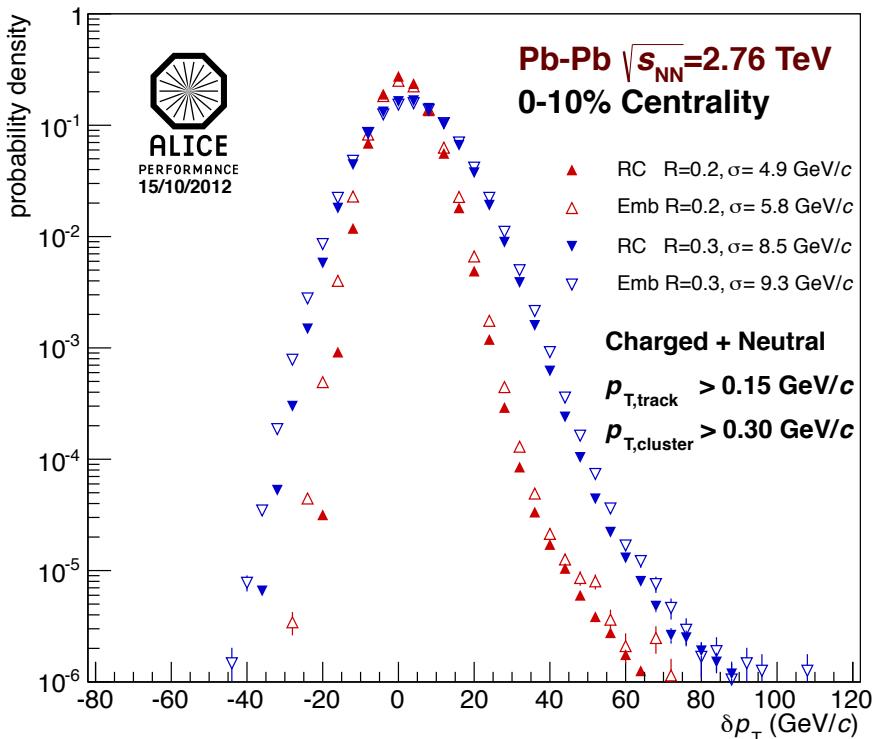
2. Leading hadron trigger



5 GeV/c is our default choice for R=0.2 jet spectra

- **Combinatorial jets:** jet finder clusters together also soft particles from the UE
- Efficiently remove the fake jets by **requiring a high p_T constituent**
- Bigger effect when changing the p_T threshold from $0 \rightarrow 5 \text{ GeV}/c$ than $5 \rightarrow 10 \text{ GeV}/c$
- Bias still effective up to 100 GeV/c

3. Background fluctuations



Random cones

1. Throw a random direction in the (η, ϕ) plane
2. Sum up the momenta of all particles in the cone of radius R
3. Calculate

$$\delta p_T^{RC} = p_T^{RC} - \pi R^2 \rho$$

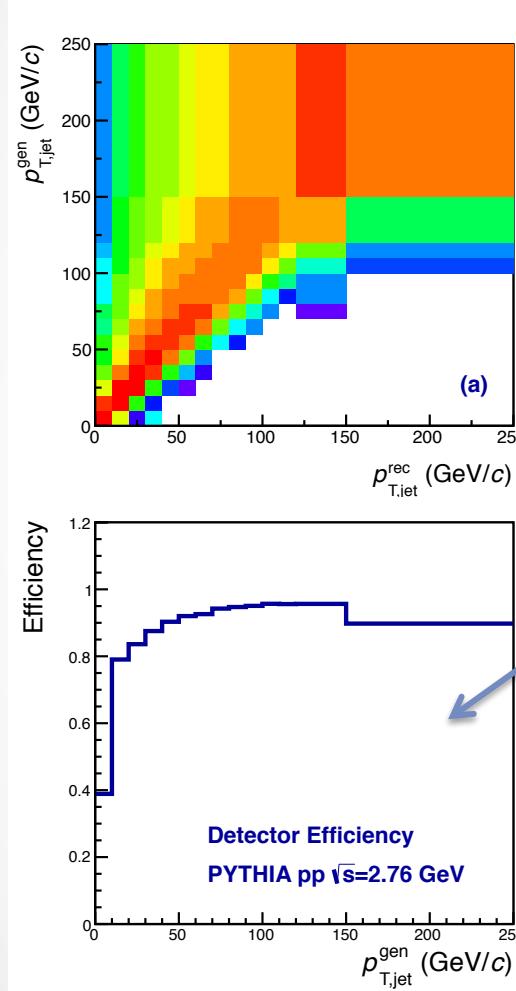
- Region-by-region background fluctuations are estimated through **random cones** (solid symbols) and **single particle embedding** (open symbols)
- Fluctuations limit **jet energy resolution**
 - **Smearing** of the p_T raw jet spectrum

Single particle embedding

1. Embed a high p_T particle
2. Run the Anti- k_T jet finder
3. Pick up the jet which contain the embedded particle and calculate

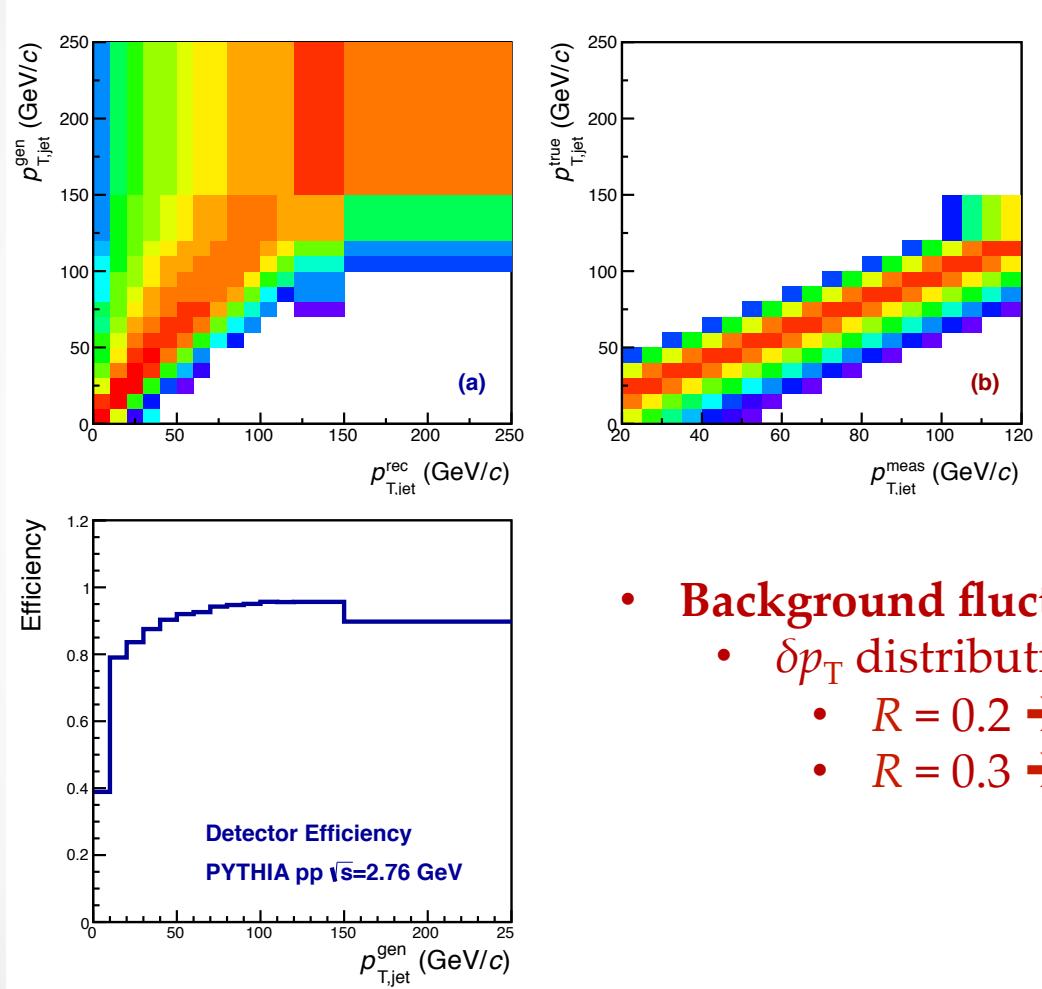
$$\delta p_T^{emb} = p_T^{jet} - p_T^{probe} - A^{jet} \rho$$

Response matrices



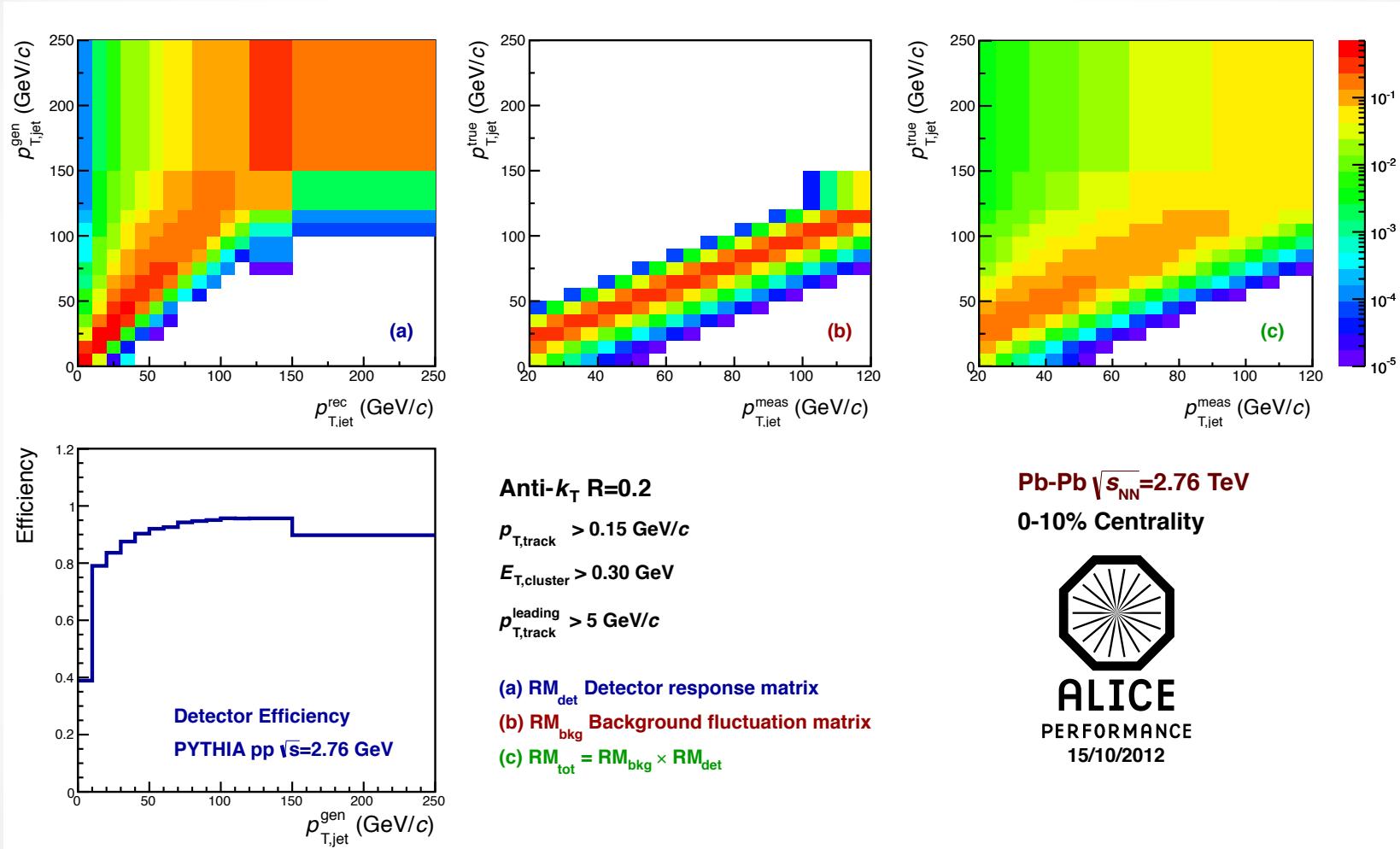
- **Detector effects**
 - PYTHIA+GEANT3 simulation
 - Assumes jet detector response is the same as in **vacuum fragmentation**
 - **Shift jet energy scale** $\sim 20\text{-}25\%$
 - Unreconstructed neutrons and K_L^0
 - Tracking inefficiency
 - **Jet energy resolution** $\sim 18\%$
 - **Jet reconstruction efficiency** $\sim 95\%$ at 80 GeV/c

Response matrices



- **Background fluctuations**
 - δp_T distribution (RC)
 - $R = 0.2 \rightarrow \sigma = 4.9 \text{ GeV}/c$
 - $R = 0.3 \rightarrow \sigma = 8.5 \text{ GeV}/c$

Response matrices



Why unfolding

- Remove detector and analysis specific features from the observables to be compared with theory
- Another simpler but less elegant approach: use the response matrix to include these distortions in the theoretical predictions
- **Advantages of unfolding**
 - Comparison between experiments is possible
 - Final results are more “fundamental” and easier to read

Why unfolding

- Remove detector and analysis specific features from the observables to be compared with theory
- Another simpler but less elegant approach: use the response matrix to include these distortions in the theoretical predictions
- **Advantages of unfolding**
 - Comparison between experiments is possible
 - Final results are more “fundamental” and easier to read
- Different unfolding methods
 - **χ^2 minimization**

$$\chi^2 = \sum_{refolded} \left(\frac{y_{refolded} - y_{measured}}{\sigma_{measured}} \right)^2$$

Why unfolding

- Remove detector and analysis specific features from the observables to be compared with theory
- Another simpler but less elegant approach: use the response matrix to include these distortions in the theoretical predictions
- **Advantages of unfolding**
 - Comparison between experiments is possible
 - Final results are more “fundamental” and easier to read
- Different unfolding methods
 - **χ^2 minimization**
 - **Bayesian**
 - Uses the Bayes’ theorem to iteratively update the estimators’ probabilities, starting from a “prior” guess

Bayes’ theorem

$$P(C_i | E) = \frac{P(E | C_i)P(C_i)}{\sum_{l=1}^{n_c} P(E | C_l)P(C_l)}$$

Why unfolding

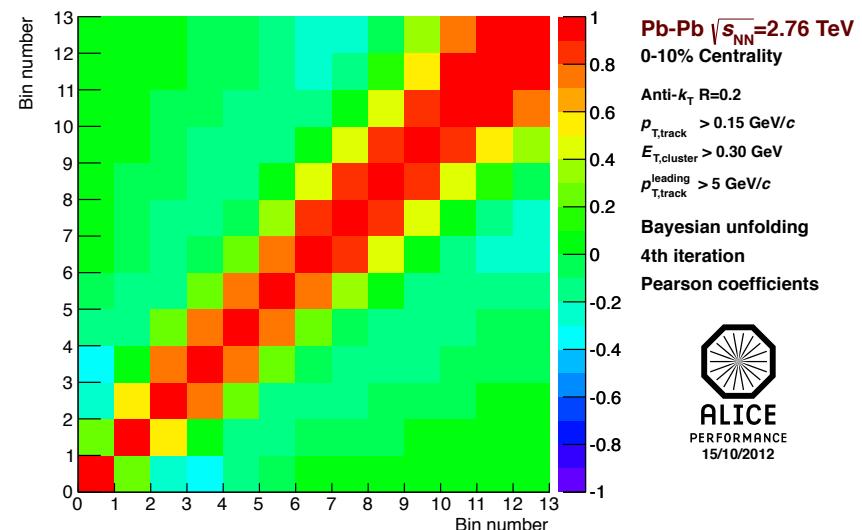
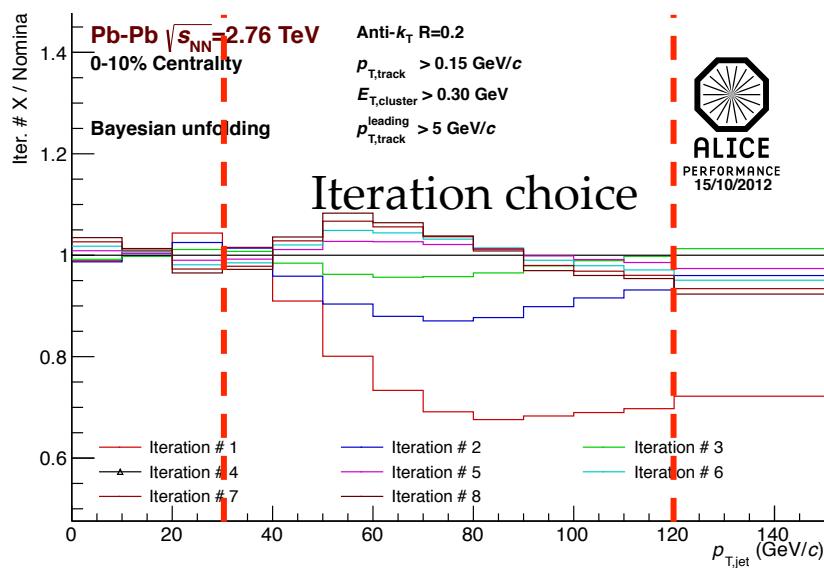
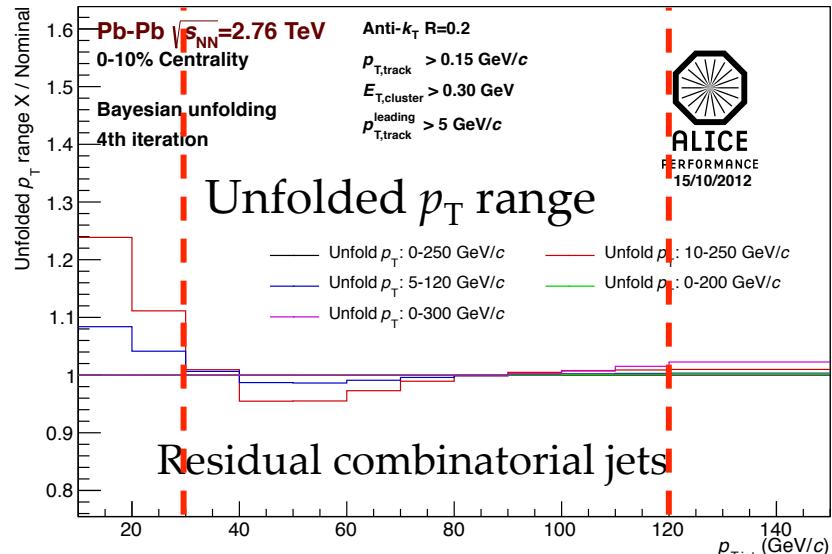
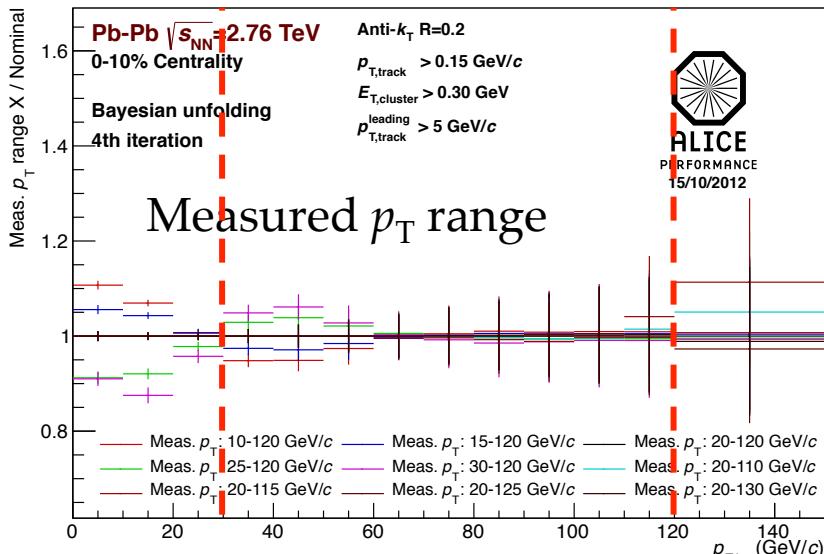
- Remove detector and analysis specific features from the observables to be compared with theory
- Another simpler but less elegant approach: use the response matrix to include these distortions in the theoretical predictions
- **Advantages of unfolding**
 - Comparison between experiments is possible
 - Final results are more “fundamental” and easier to read
- Different unfolding methods
 - **χ^2 minimization**
$$\chi^2 = \sum_{refolded} \left(\frac{y_{refolded} - y_{measured}}{\sigma_{measured}} \right)^2$$
 - **Bayesian**
 - Uses the Bayes’ theorem to iteratively update the estimators’ probabilities, starting from a “prior” guess
 - Infinite iterations would give the best mathematical solution, but highly fluctuating, not useful for physics

Why unfolding

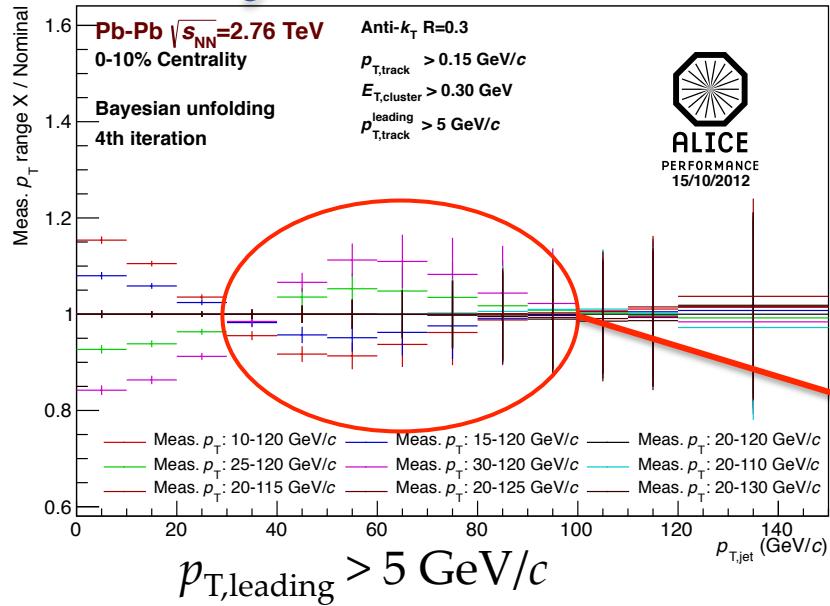
- Remove detector and analysis specific features from the observables to be compared with theory
- Another simpler but less elegant approach: use the response matrix to include these distortions in the theoretical predictions
- **Advantages of unfolding**
 - Comparison between experiments is possible
 - Final results are more “fundamental” and easier to read
- Different unfolding methods
 - **X² minimization**
$$\chi^2 = \sum_{refolded} \left(\frac{y_{refolded} - y_{measured}}{\sigma_{measured}} \right)^2 + \beta \sum_{unfolded} \left(\frac{d^2 \log y_{unfolded}}{d \log p_T^2} \right)$$
 - **Bayesian**
 - Uses the Bayes' theorem to iteratively update the estimators' probabilities, starting from a “prior” guess (**finite # of iterations**)
 - Infinite iterations would give the best mathematical solution, but highly fluctuating, not useful for physics
- **Regularization parameter:** *a priori* knowledge about spectrum shape
 - Avoids wildly fluctuating results
 - Introduce a bias in the result

G. D'Agostini, Nuclear Instr. Meth., 362 (1995) 487-498
G. Cowan, A Survey of Unfolding Methods in Part. Phys.,
Proc. Adv. Stat. Tech. in Part. Phys., Durham (2002) ([link](#))

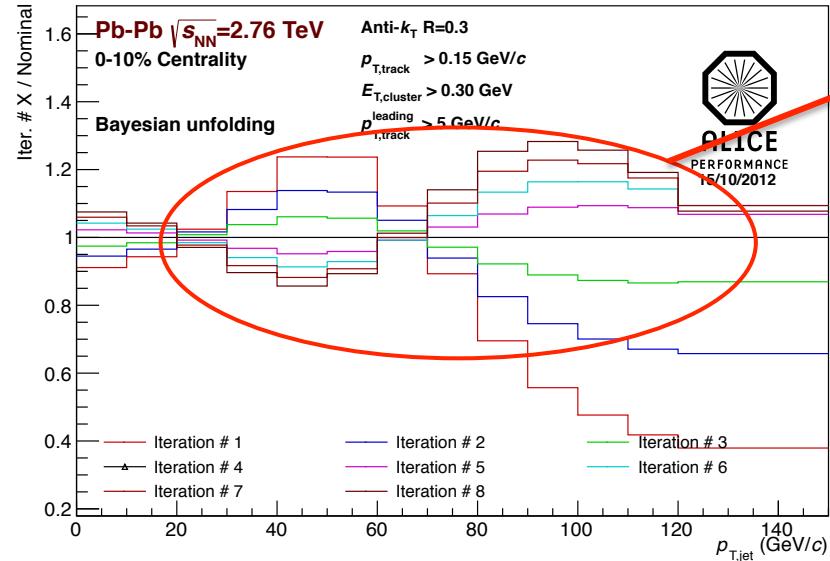
Bayesian unfolding results – R=0.2



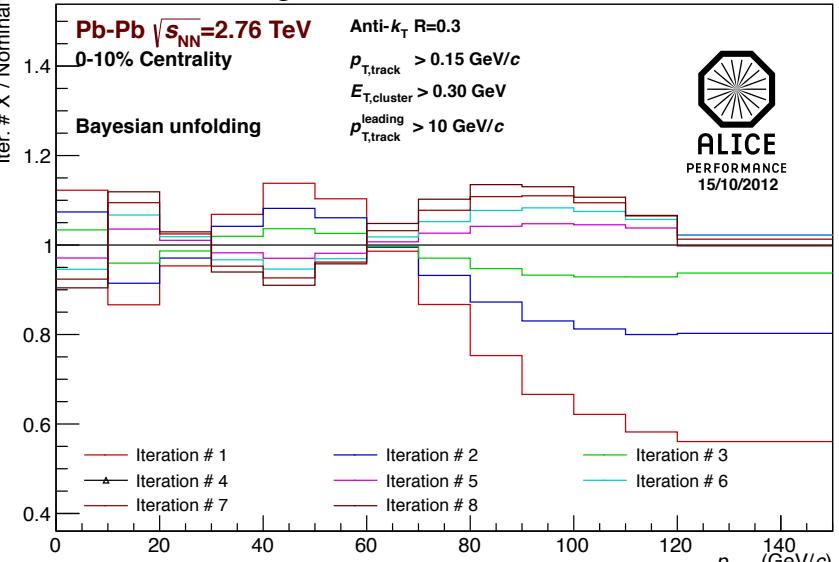
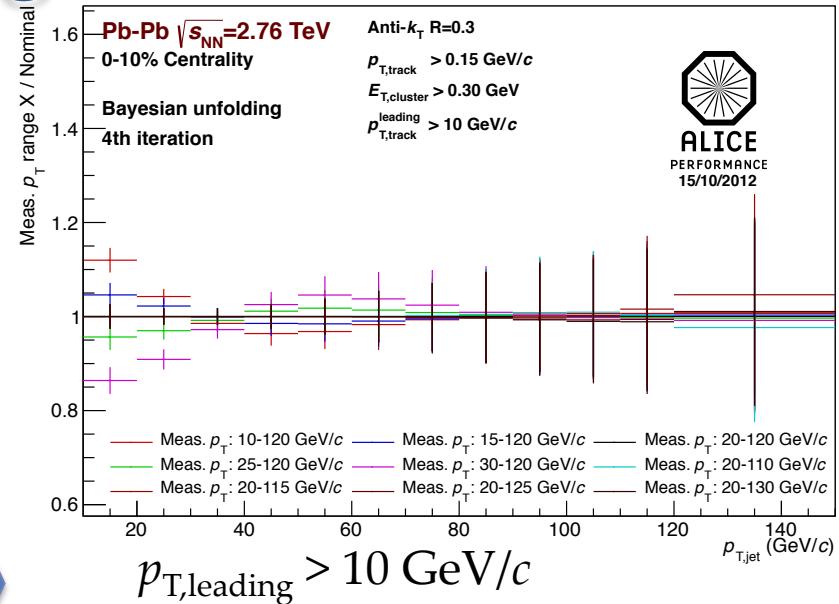
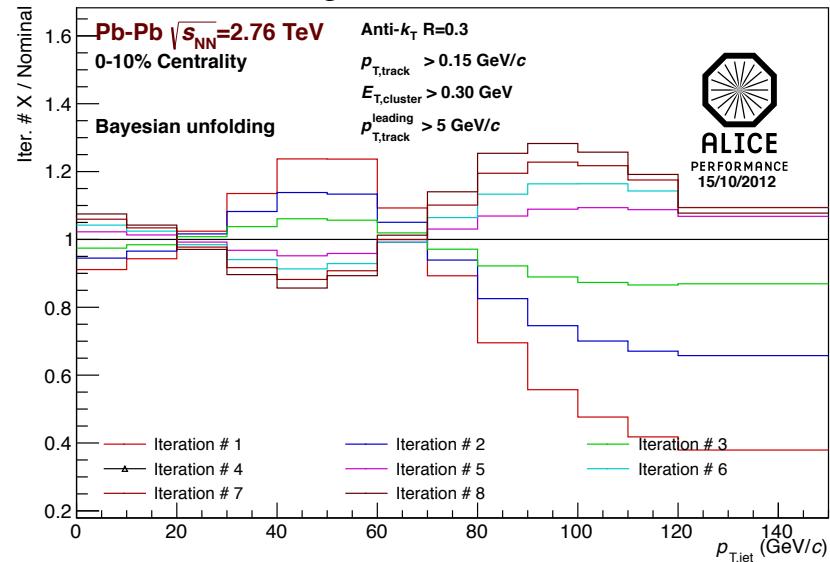
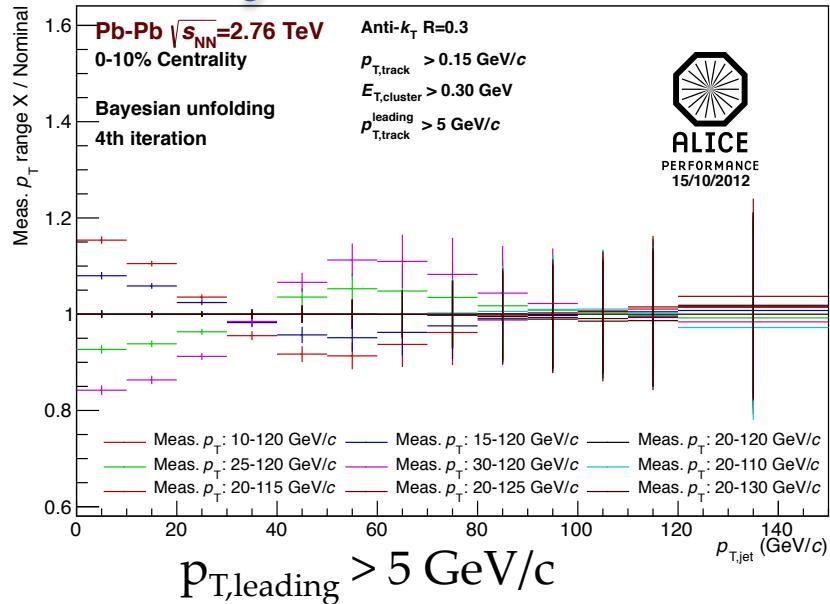
Bayesian unfolding results – R=0.3



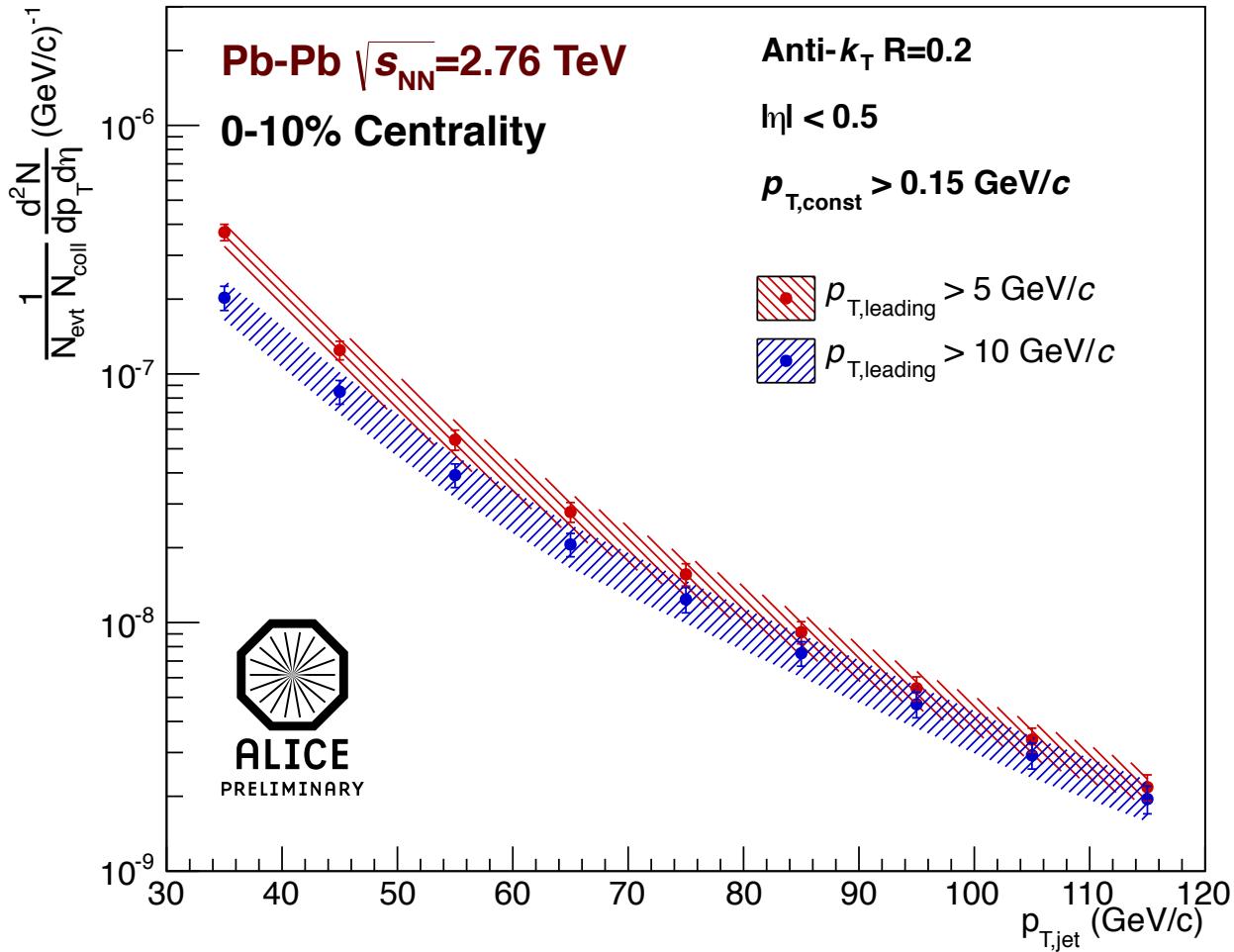
Unfolding instabilities...



Bayesian unfolding results – R=0.3

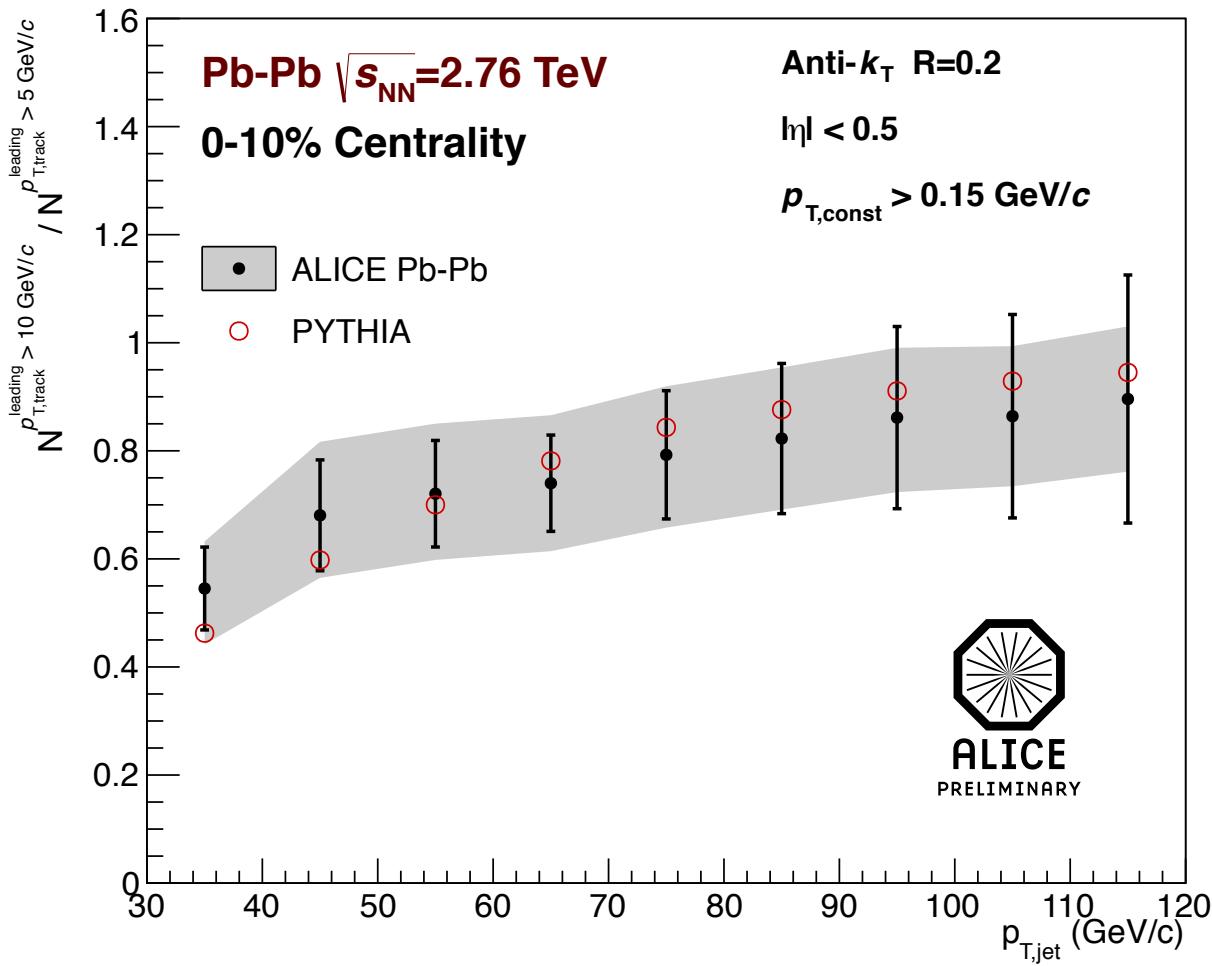


Unfolded biased spectra



- Fully corrected p_T jet spectra
 - $R=0.2$
 - $p_{T,\text{const}} > 0.15 \text{ GeV}/c$
- Two different leading hadron triggers
 - trigger only on charged tracks to avoid difficulties in the theoretical predictions
 - $p_{T,\text{track}} > 5 \text{ GeV}/c$
 - $p_{T,\text{track}} > 10 \text{ GeV}/c$

Effect of the leading hadron requirement



- Statistics and systematics
 - Take into account correlated uncertainties
- Agreement with PYTHIA is consistent with vacuum-like fragmentation of the jet core

Conclusions & Outlook

- **Background** in Pb-Pb central collisions has been studied in order to be able to extract a “clean” signal
- **Bayesian unfolding** used to correct spectra for **background fluctuations** and **detector-induced effects**
 - Studies are ongoing to unfold the **R=0.3 spectra**
- Effect of the **leading hadron requirement** studied for two different thresholds, 5 GeV/c and 10 GeV/c
 - Agreement with PYTHIA is **consistent with vacuum-like fragmentation** of the jet core

Backup

Systematic uncertainties

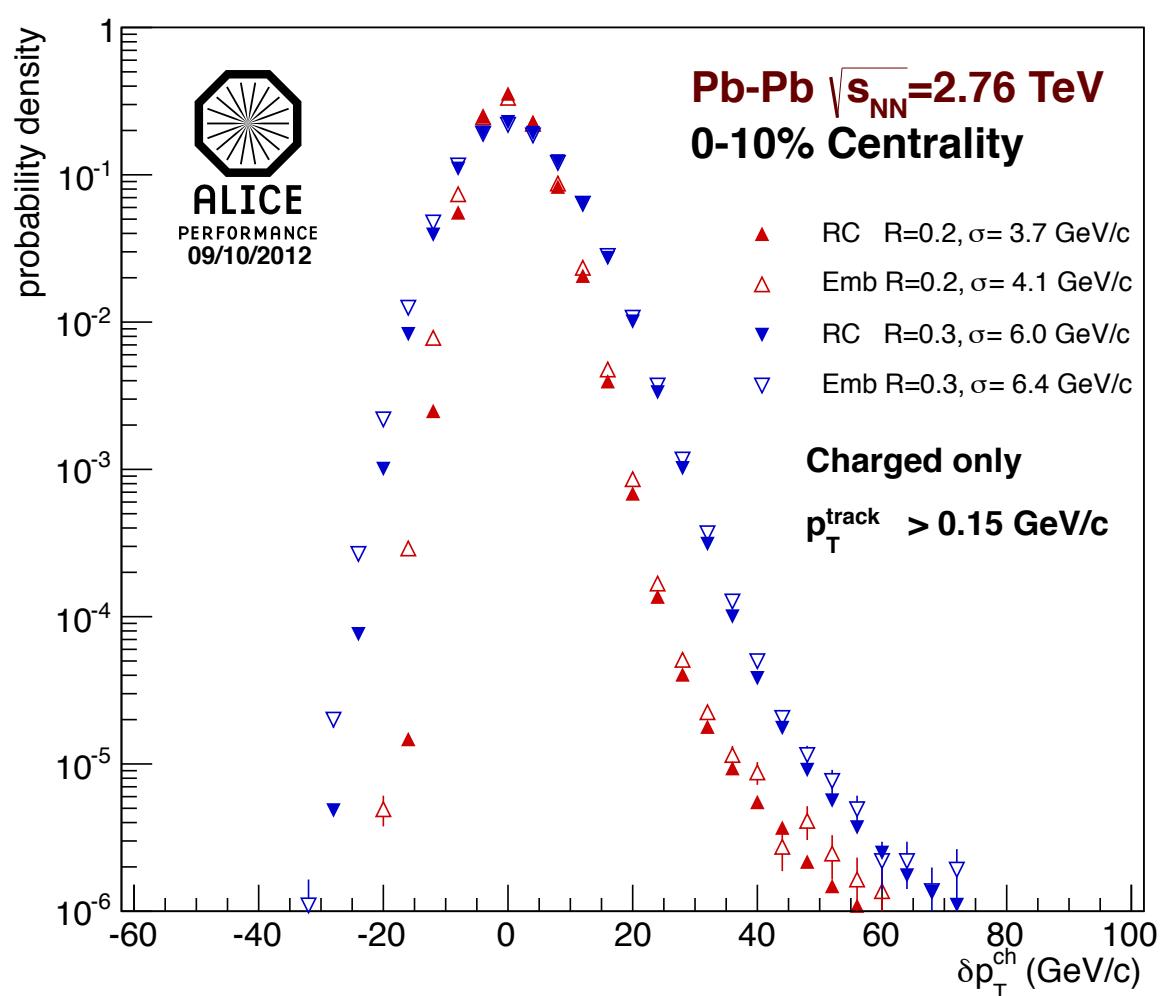
Uncertainty	% on jet spectrum
EMCal (resolution, energy scale, clusterizer, non-linearity)	5.3
Tracking Efficiency	10
Hadronic Correction	10
Scale Factor	2
δp_T RC vs. Embedding	5
Unfolding	14
Flow bias	< 1%
Total	19

Unfolding uncertainty	% on jet spectrum
Iteration	5
Measured p_T range	9
Unfolded p_T range	5
Prior choice	10
Total unfolding	14

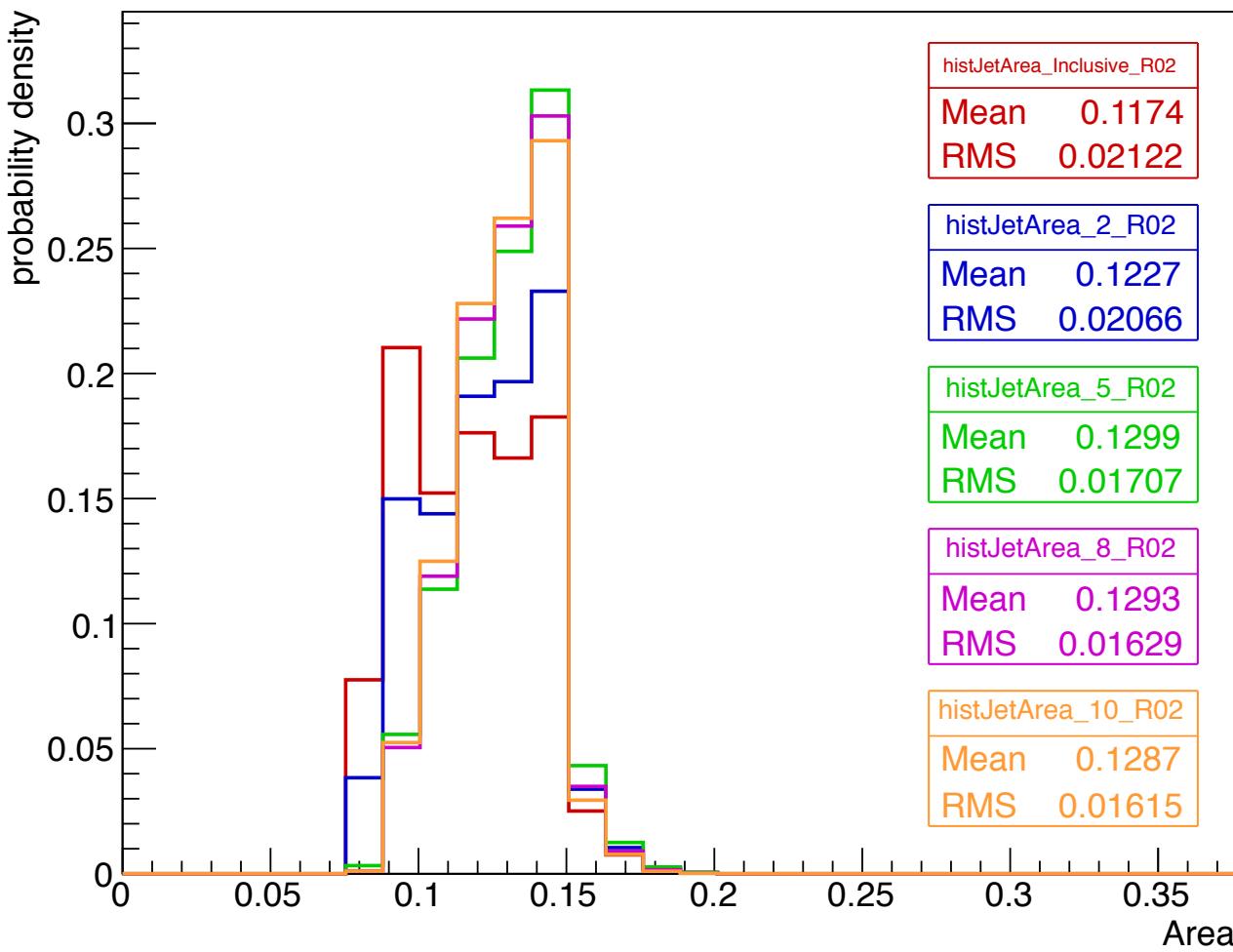
Some of the uncertainties are asymmetric and p_T dependent.
Here only the maximum value is shown.

Systematic uncertainties for 10 GeV/c bias are very similar

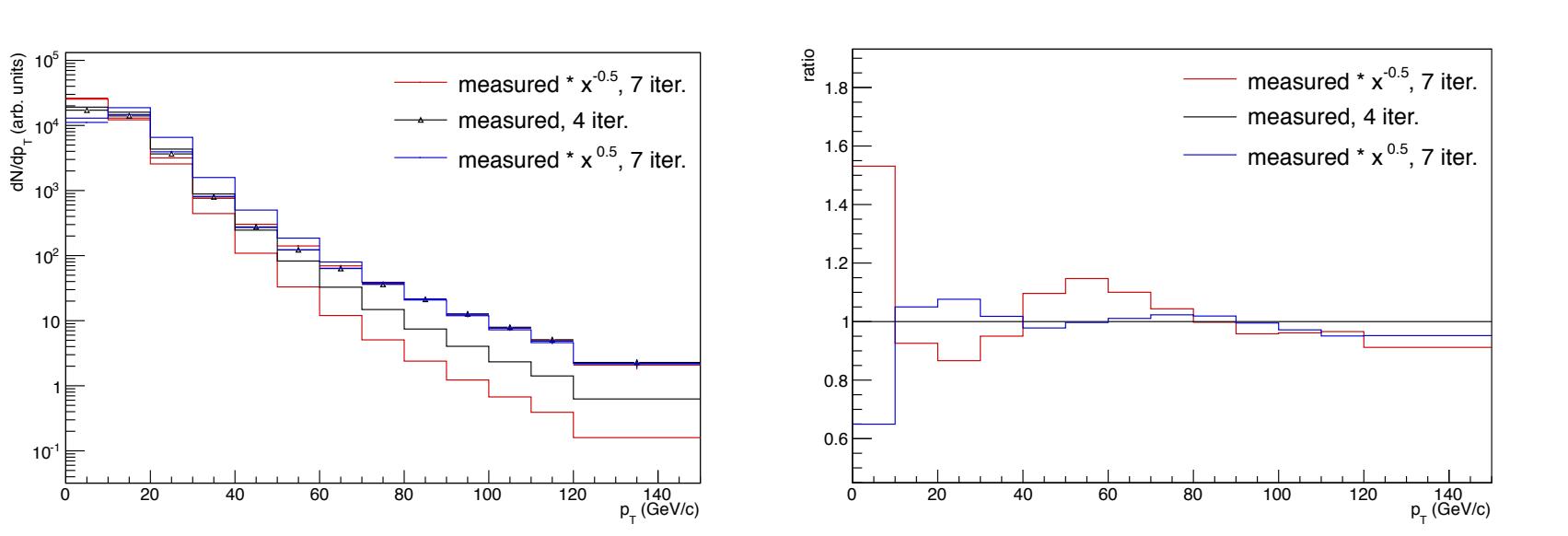
Charged δp_T



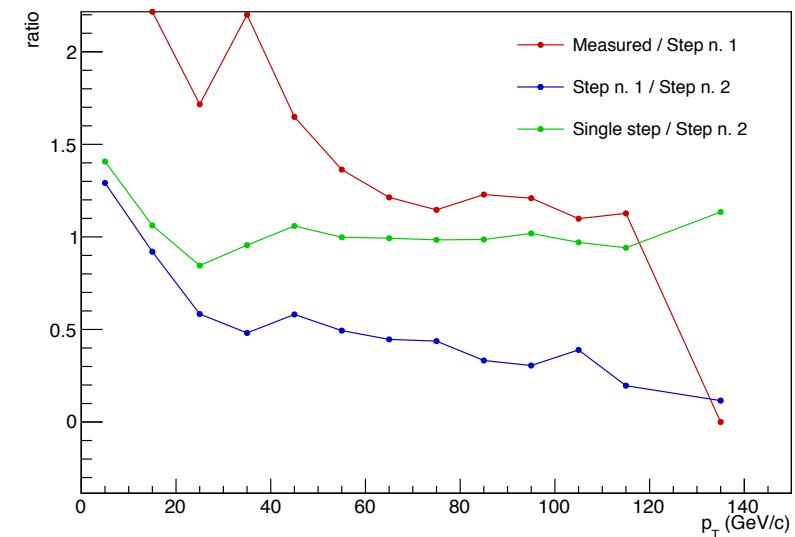
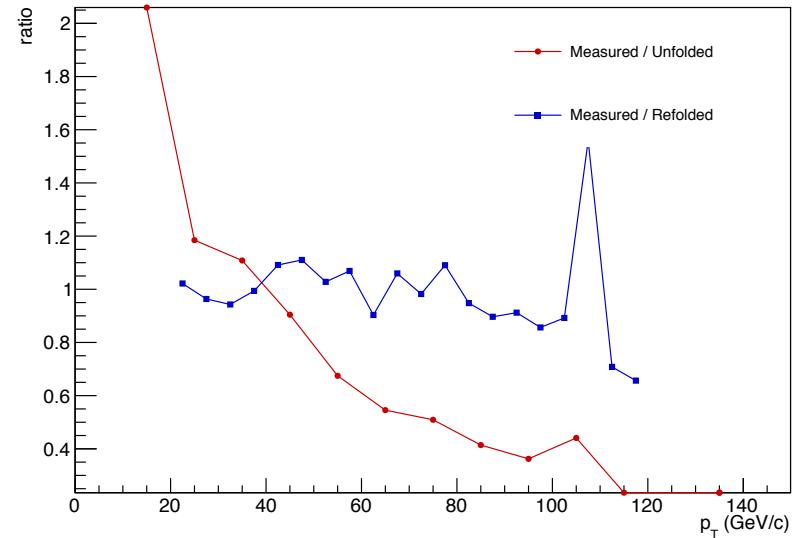
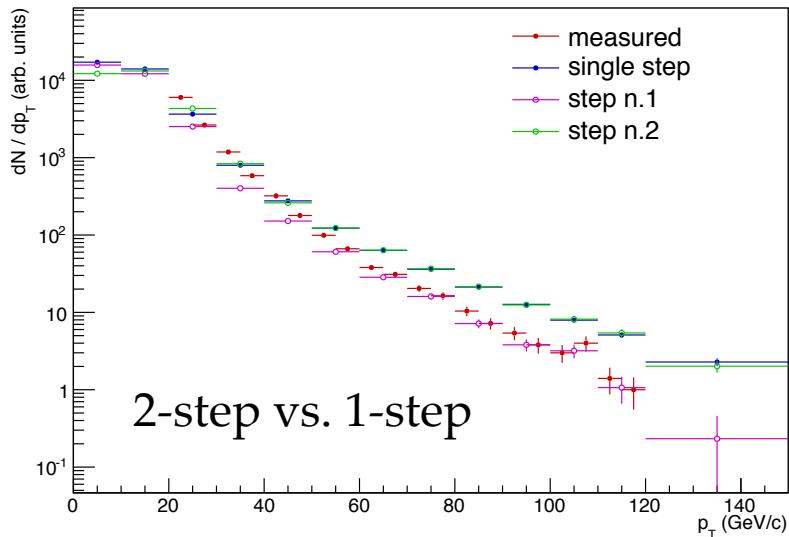
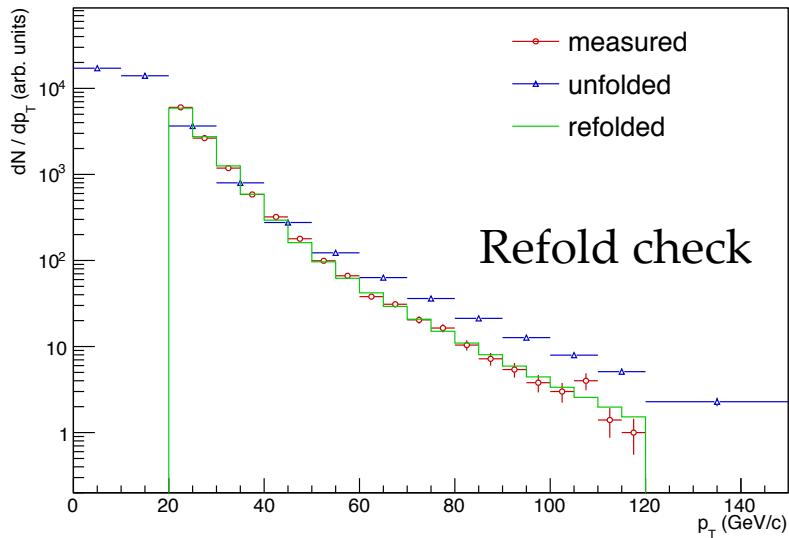
Jet area distribution



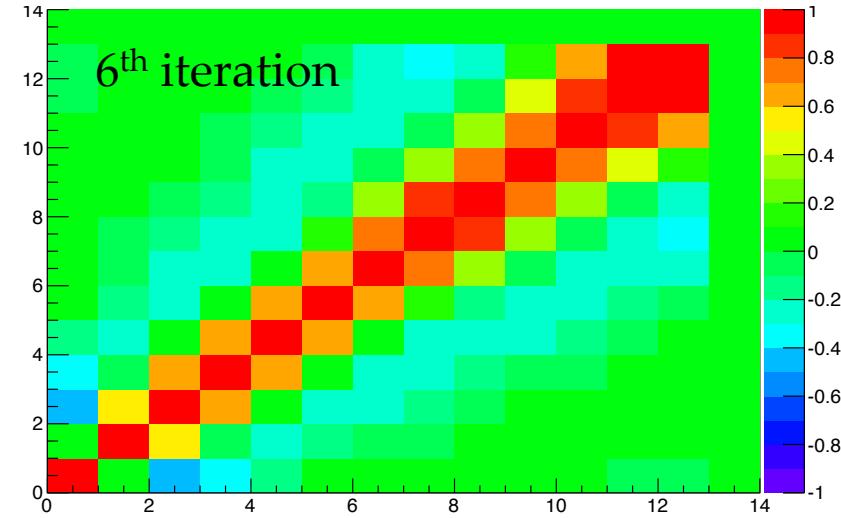
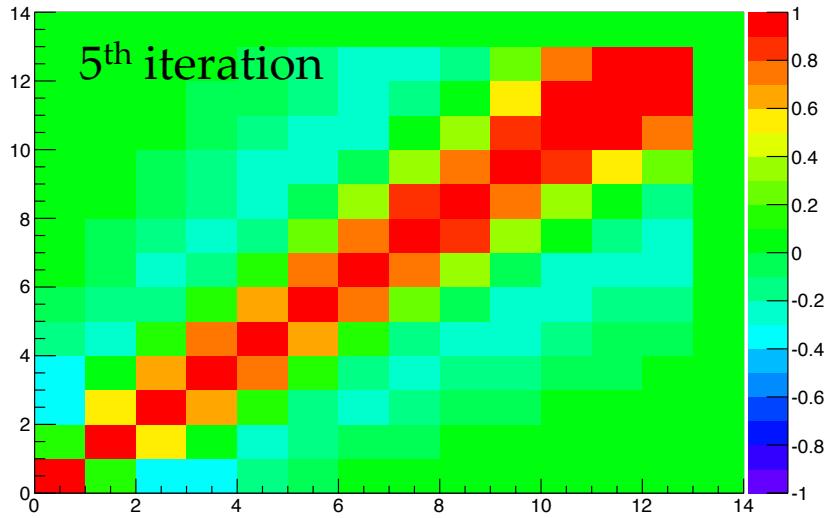
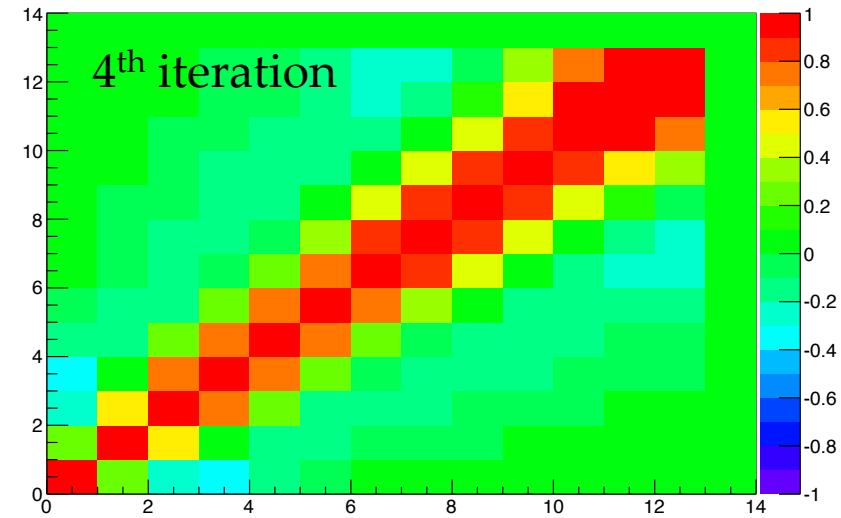
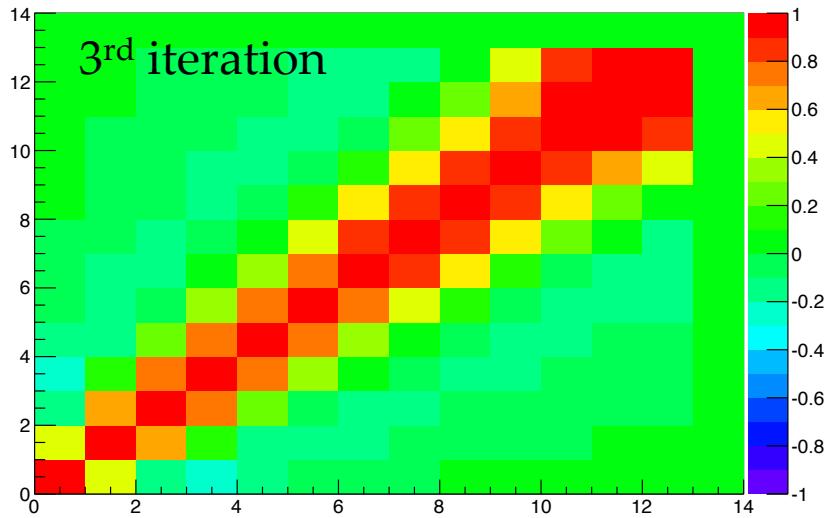
Unfolding prior choice



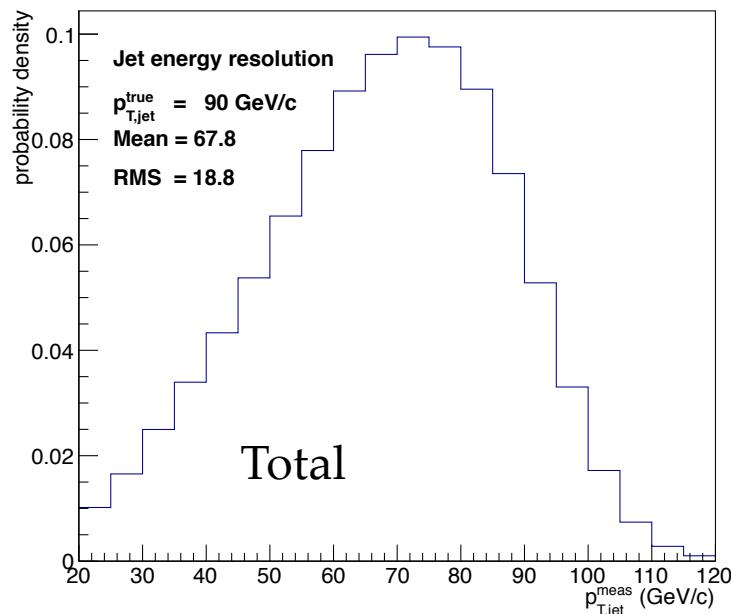
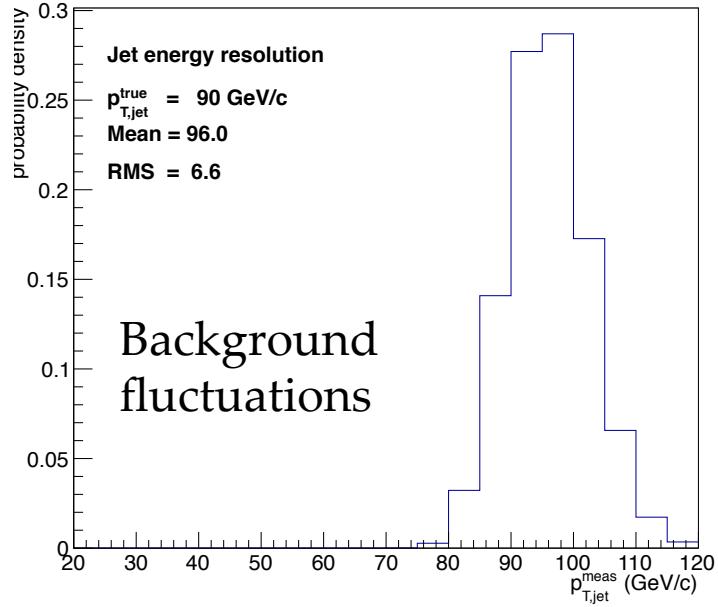
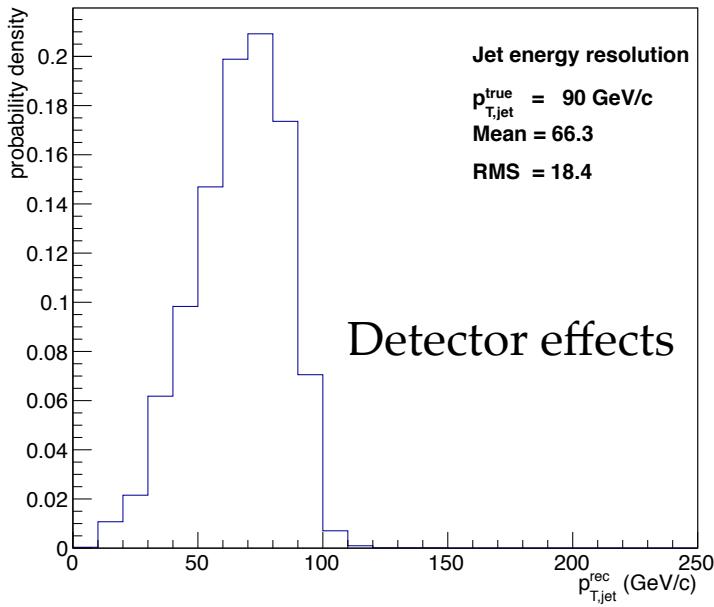
Refold, 2-step



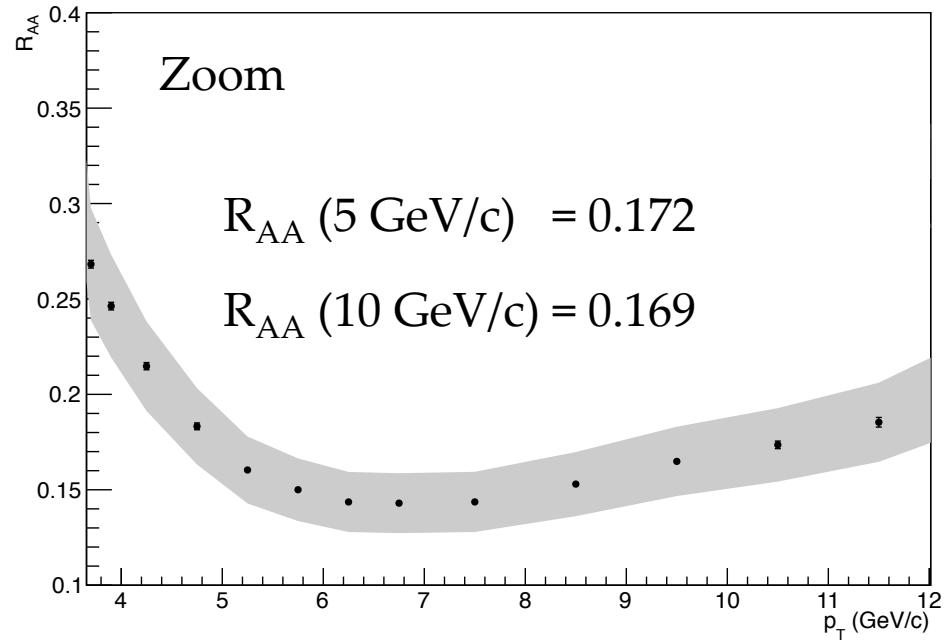
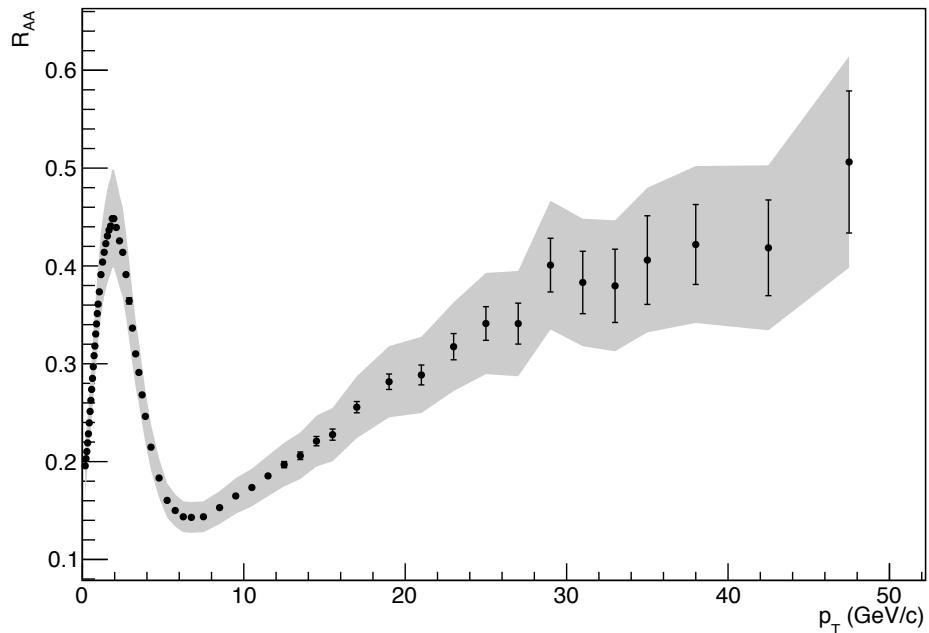
Pearson coefficients



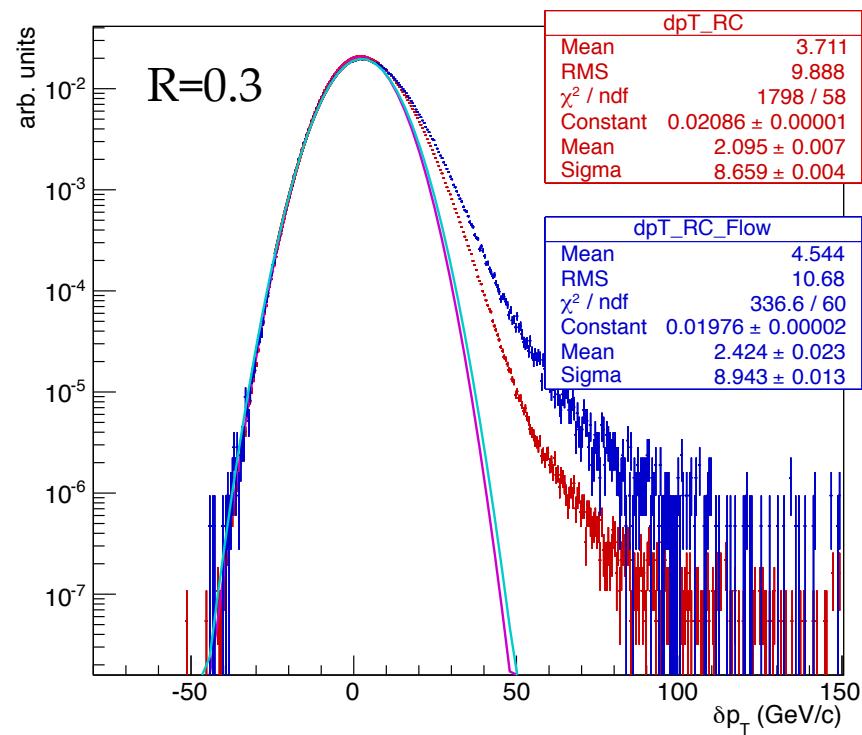
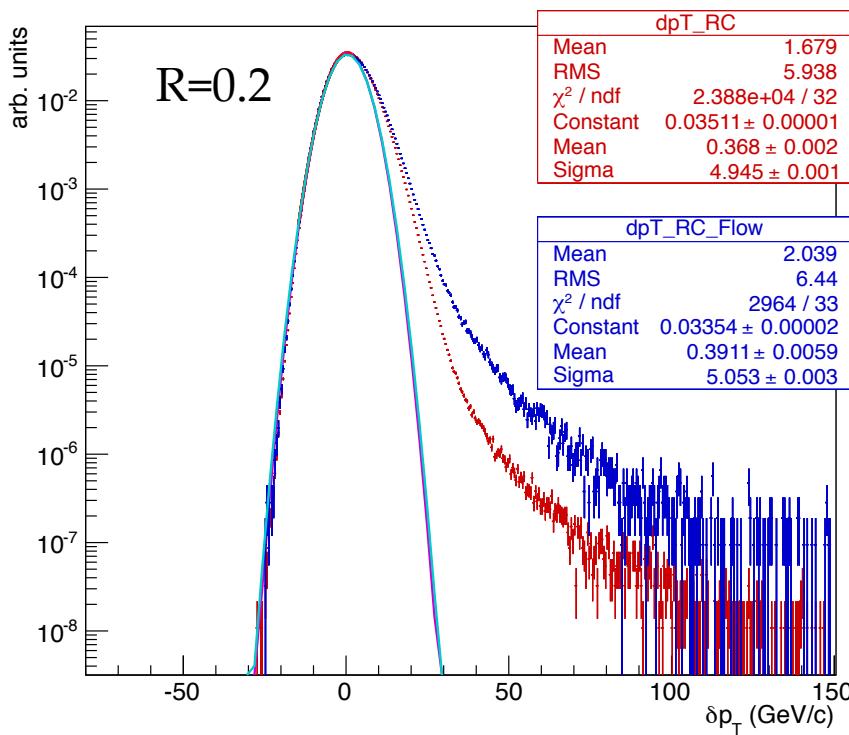
Jet energy resolution



Charged particles R_{AA}



Flow bias



Blue = random cones δp_T distributions obtained requiring a 5 GeV/c charged track in the calorimeter

Red = regular random cones δp_T distributions

The shift measured with the mean of the Gaussian fit on the LHS is < 100 MeV/c for R=0.2 and 300 MeV/c for R=0.3