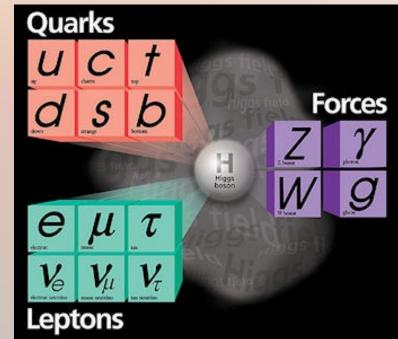
Combined search for the SM Higgs Boson at DØ

James Kraus University of Mississippi On behalf of the DØ Collaboration



The Higgs Boson in the Standard Model

- Higgs mechanism responsible for electroweak symmetry breaking
- Interactions with Higgs field gives particles their masses
- Standard Model does not predict the Higgs mass



• CMS and ATLAS announced the discovery of a Higgs-like boson in the H \rightarrow ZZ and H $\rightarrow \gamma\gamma$ channels in July 2012.

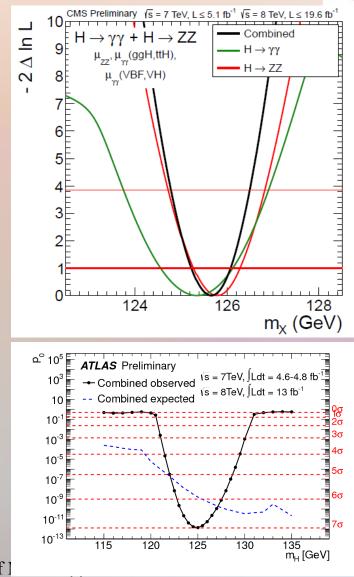




3

Standard Model Higgs Boson at LHC

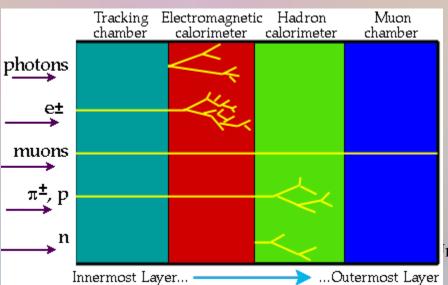
- The ATLAS and CMS collaborations announced the discovery of a new, Higgs-like particle at a mass of ~125 GeV in July 2012
- So far all properties are consistent with the SM Higgs boson.

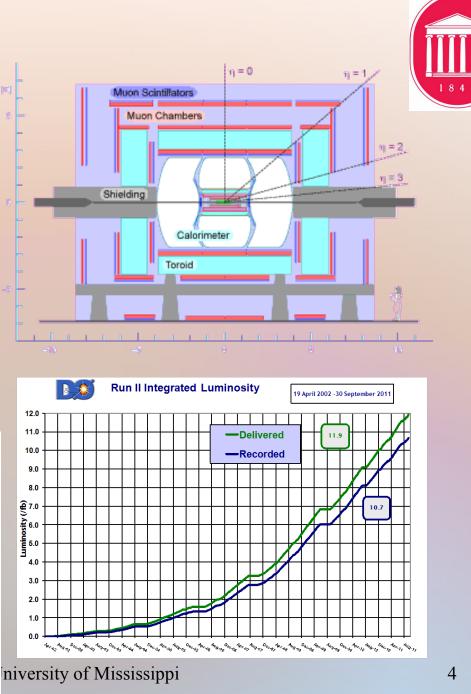




The DØ Experiment

- A multipurpose particle detector
- Innermost detectors are the trackers, followed by calorimetry and muon chambers

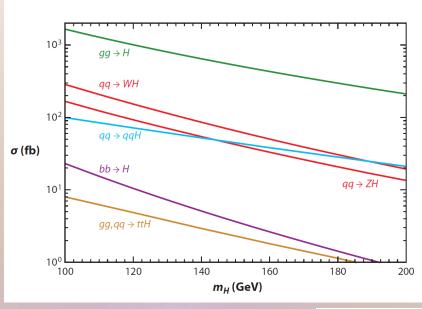




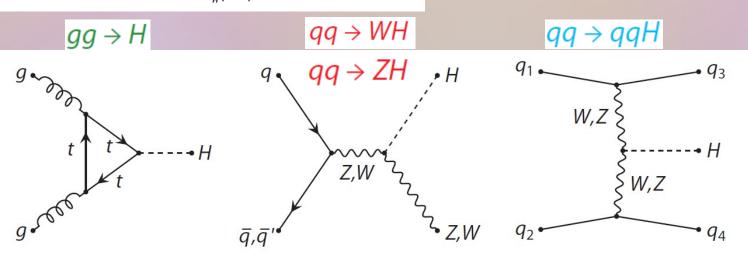




Higgs Production at the Tevatron



The Higgs at Tevatron produced by gluon fusion $(gg \rightarrow H)$, associated production $(qq \rightarrow W/Z+H)$, and vector boson fusion $(qq \rightarrow qqH)$

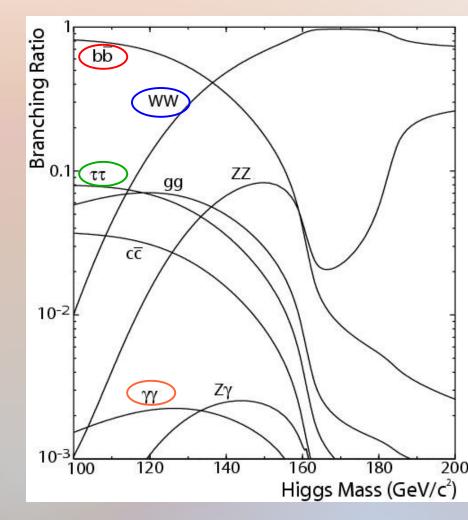






Higgs Decay Modes

- The Higgs decays primarily to bottom quarks at low mass, WW at higher mass
- The searches for the Higgs at the Tevatron focus on H→bb at low mass, H→WW at high mass.
- Also use $H \rightarrow \tau\tau$, $H \rightarrow \gamma\gamma$ due to low backgrounds

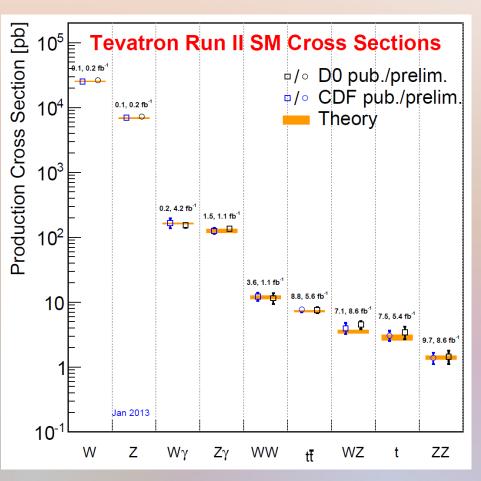






Higgs Backgrounds

- The rate of Higgs production is small compared to backgrounds
 - Bottom decays swamped by QCD background except for associated ZH/WH production
 - WW cleaner, but still large backgrounds
- These backgrounds are modeled with Alpgen+Pythia, Pythia, and CompHEP
 - Multijet and some other backgrounds modeled with data

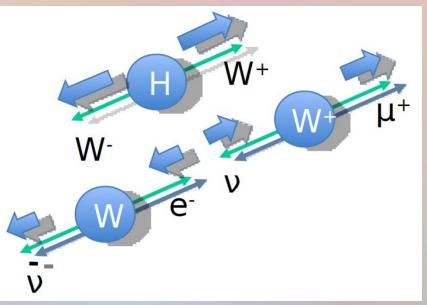






Signal vs. Background

- Once we have a good background model, we look for differences between signal and background
- Example from $H \rightarrow WW \Delta \phi$ between leptons



- The Higgs is a spin-0 particle; W is spin-1
- \Rightarrow small opening angle in leptons from the W decays
- Leptons from WW isotropic; from Z back-toback

James Kraus, University of Mississippi





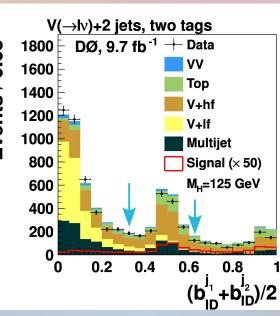
Signal vs. Background

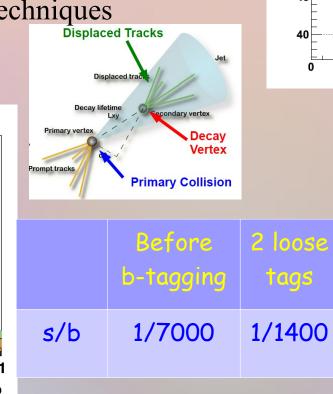
- Once we have a good background model, we look for differences between signal and background
- Example $\Delta \phi$ between leptons
- The signal to background ratio is so small that any one variable will not be sufficient to isolate the Higgs
- Instead, we look at many variables associated with each event at once multivariate analysis

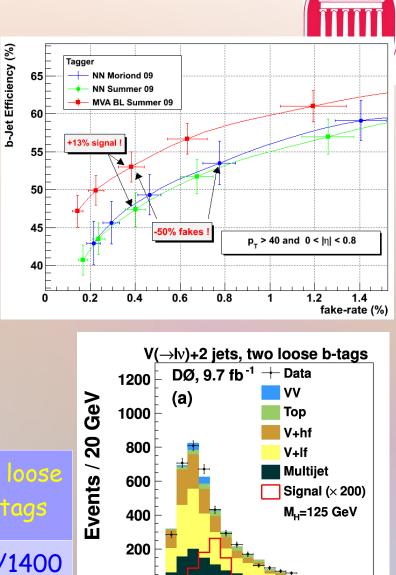
S vs. B in VH \rightarrow Vbb

- Important:
 - Jet energy resolution $\Delta m/m \sim 15\%$
 - b-tagging
 - Multivariate techniques









0

50 100 150 200 250 300 350 400 **Dijet mass [GeV]**

James Kraus, University of Mississippi

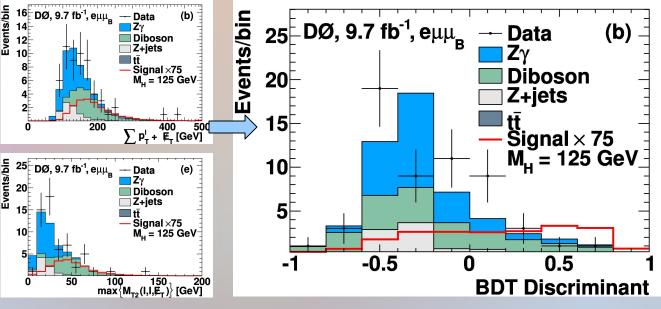




Multivariate Analysis

- Take multiple variables, each of which has some separating power
- Combine them into one variable that separates the signal and background

14 variables +

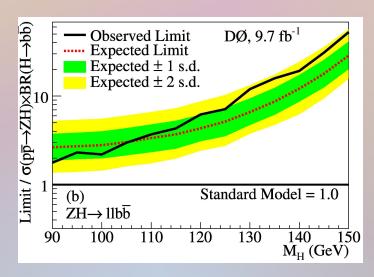


James Kraus, University of Mississippi





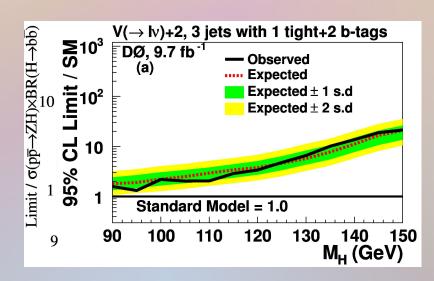
We have limited sensitivity in any one final state
 – We therefore look at ZH→llbb, (link)







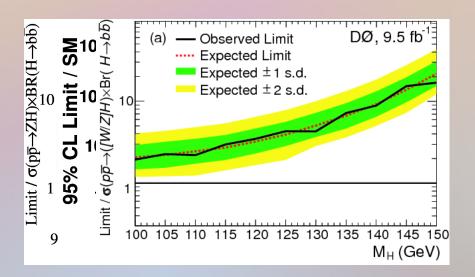
We have limited sensitivity in any one final state
 – We therefore look at ZH→llbb, WH→lvbb, (link)







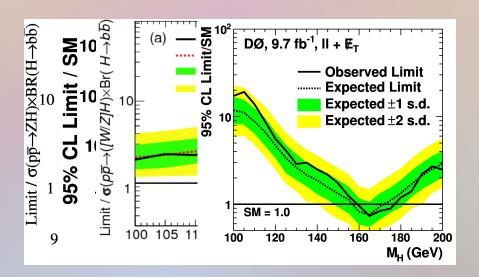
- We have limited sensitivity in any one final state
 - We therefore look at ZH→llbb, WH→lvbb, ZH→vvbb,
 (link)







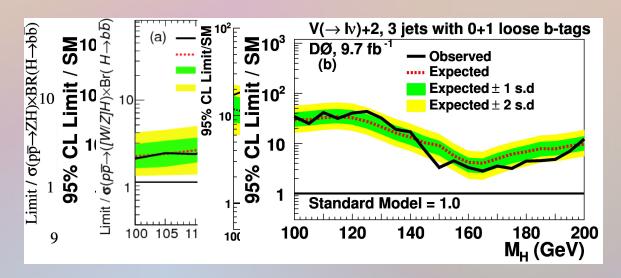
- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, (<u>link</u>)







- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', (<u>link</u>)



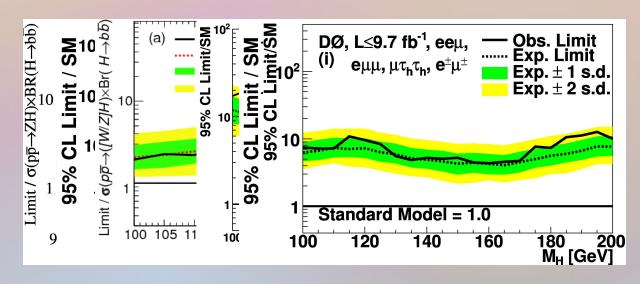
James Kraus, University of Mississippi







- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', VH \rightarrow trileptons and e⁺µ⁺, (link)



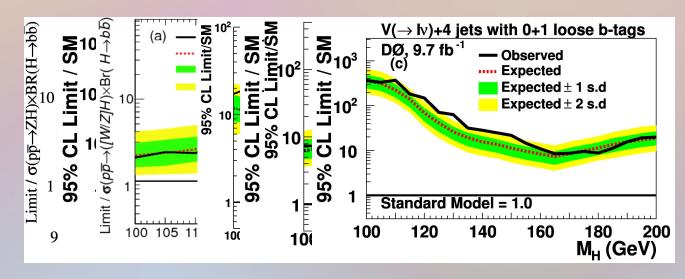
James Kraus, University of Mississippi







- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', VH \rightarrow trileptons and e⁺µ⁺, VH \rightarrow lvqqqq', (<u>link</u>)

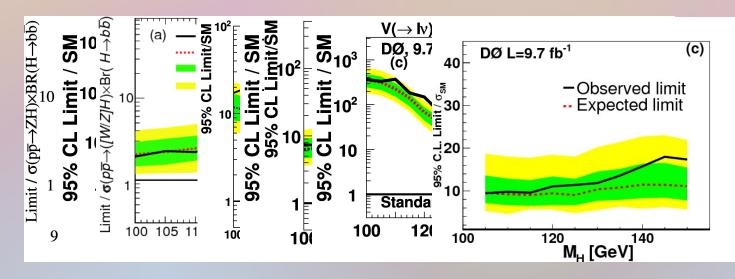








- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', VH \rightarrow trileptons and e⁺µ⁺, VH \rightarrow lvqqqq', H \rightarrow ττ,(link)



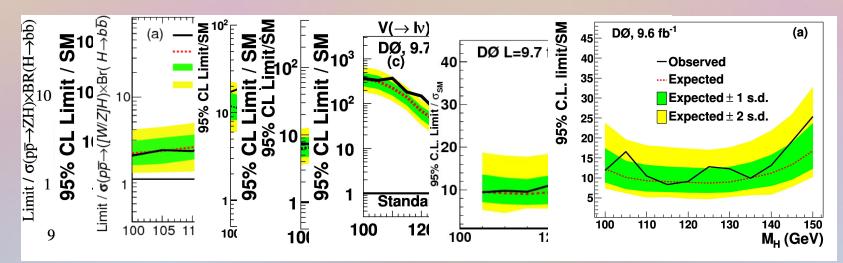
James Kraus, University of Mississippi







- We have limited sensitivity in any one final state
 - We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', VH \rightarrow trileptons and e⁺µ⁺, VH \rightarrow lvqqqq', H \rightarrow t\tau, and H $\rightarrow\gamma\gamma$ (link)









• We have limited sensitivity in any one final state

- We therefore look at ZH \rightarrow llbb, WH \rightarrow lvbb,ZH \rightarrow vvbb, H \rightarrow WW \rightarrow lvlv, H \rightarrow WW \rightarrow lvqq', VH \rightarrow trileptons and e⁺µ⁺, VH \rightarrow lvqqqq', H \rightarrow ττ, and H $\rightarrow\gamma\gamma$

Channel $(V = W, Z \text{ and } \ell = e, \mu)$		Luminosity (fb^{-1})	$M_H \; (\text{GeV})$
$WH \to \ell \nu bb$	$H o b \bar{b}$	9.7	90-150
$ZH \to \ell\ell b \bar{b}$		9.7	90 - 150
$ZH o \nu \bar{\nu} b \bar{b}$		9.5	100 - 150
$H \to W^+ W^- \to \ell^+ \nu \ell^- \bar{\nu}$	$H \to W^+ W^-$	9.7	100-200
$H + X \to W^+ W^- \to \mu^{\pm} \tau_h^{\mp} + \leq 1$ jet		7.3	155 - 200
$H \to W^+ W^- \to \ell \nu q' \bar{q}$		9.7	100 - 200
$VH \rightarrow ee\mu/\mu\mu e + X$		9.7	100 - 200
$VH \to e^{\pm}\mu^{\pm} + X$		9.7	100 - 200
$VH o \ell \nu q' \bar{q} q' \bar{q}$		9.7	100 - 200
$VH \to \tau_h \tau_h \mu + X$	$H \to \tau^+ \tau^-$	8.6	100 - 150
$H + X \to \ell \tau_h j j$		9.7	105 - 150
$H \to \gamma \gamma$		9.7	100-150



DØ, L_{int} ≤ 9.7 fb⁻¹

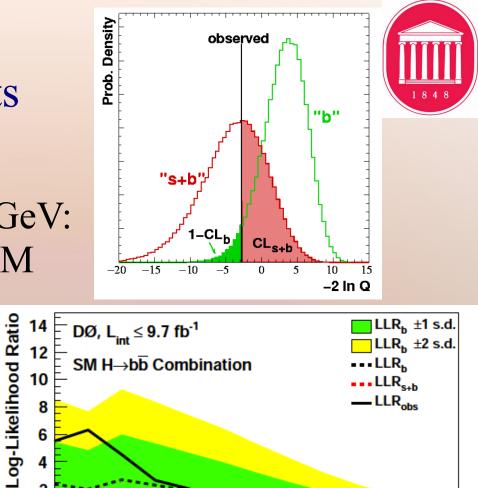
SM H→bb Combination

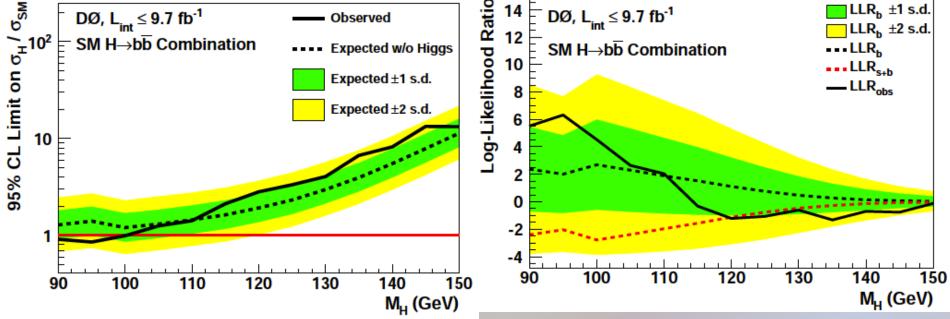
$VH \rightarrow Vbb$ Results



Observed

Expected w/o Higgs







H→WW Result



Dilepton only:

Exclude 159–176 GeV For 125 GeV, exp. 3.4×SM, obs. 4.1×SM

Full H→WW:

Exclude 157–-178 GeV For 125 GeV, exp. 2.9×SM, obs. 4.6×SM

15% improvement over dilepton alone 95% CL Limit/SM 01 01 DØ, 9.7 fb⁻¹, II + ∉_T 0 SM DØ, $L_{int} \le 9.7 \text{ fb}^{-1}$ Observed SM H→W⁺W⁻ Combination Expected w/o Higgs ь **Observed Limit** 95% CL Limit on Expected Limit Expected ±1 s.d. Expected ±1 s.d. 10 Expected ±2 s.d. Expected ± 2 s.d. SM = 1.0180 120 140 160 200 100 mes Kraus, Un M_н (GeV) M_L (GeV)





$H \rightarrow \tau \tau$ and $H \rightarrow \gamma \gamma$ Results

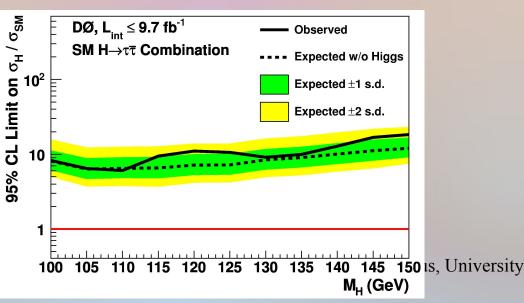
 $H \rightarrow \tau \tau$:

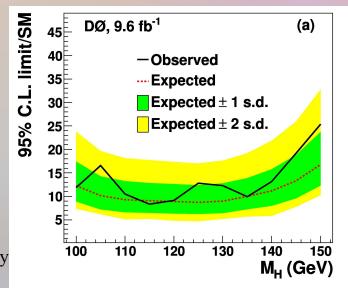
$H \rightarrow \gamma \gamma$:

For 125 GeV, exp 7.25×SM but small obs 10.4×SM For 125 Ge

Expect a narrow resonance, but small branching ratio

For 125 GeV, exp 8.7×SM obs 12.8×SM



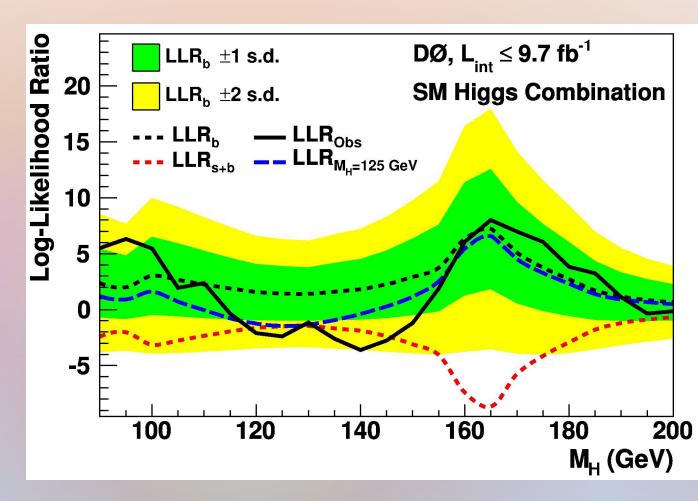


24



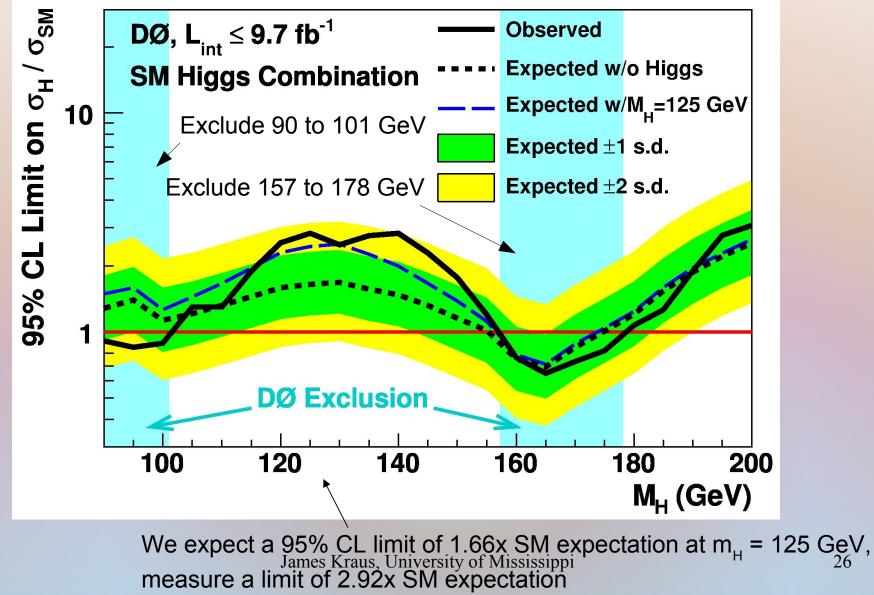


The blue dashed line is the expected LLR for background + Higgs boson signal having a mass of 125 GeV with expected SM cross section.







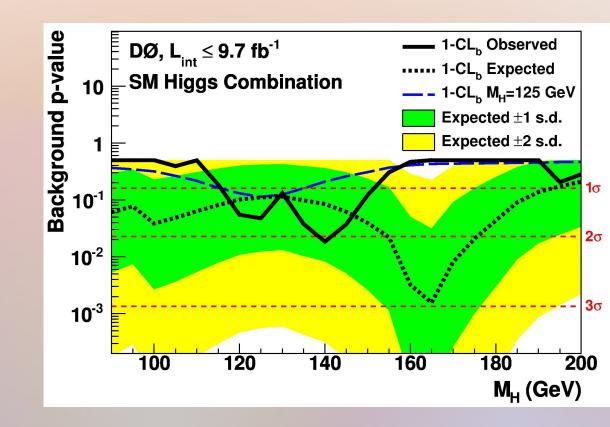




Magnitude of Excess



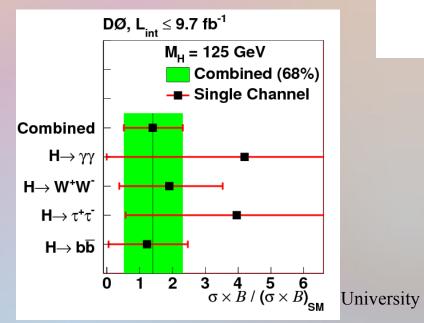
Between M_H
 120–145 GeV,
 up to a 2σ excess

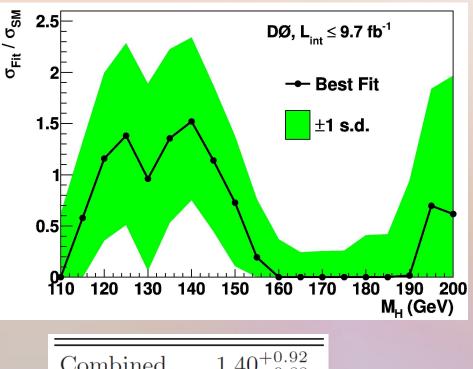




Favored Cross Sections

- With M_H=125 GeV, best fit cross-section is 1.4× SM cross-section
 - Excess in all 4 main subchannels





Combined	$1.40_{-0.88}^{+0.02}$
$H\to\gamma\gamma$	$4.20^{+4.60}_{-4.20}$
$H \to W^+ W^-$	$1.90^{+1.63}_{-1.52}$
$H \to \tau^+ \tau^-$	$3.96^{+4.11}_{-3.38}$
$H \to b \bar{b}$	$1.23^{+1.24}_{-1.17}$



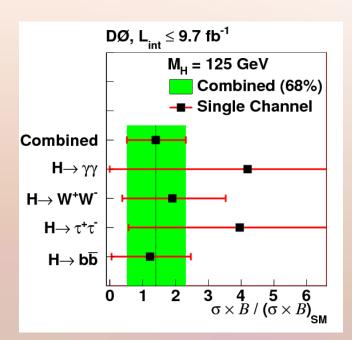


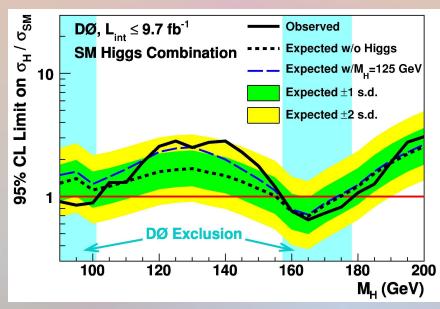
Conclusions

• Have set limits on SM Higgs boson production at DØ

Sub. to PRD (arXiv:1303.0823)

- With M_{H} =125 GeV, best fit cross-section is $1.4 \times SM$ cross-section
 - Excess seen in all four main subchannels
- DØ excludes Higgs boson masses between 90–101 GeV and 157 – 178 GeV @ 95% CL









Backup Slides

James Kraus, University of Mississippi

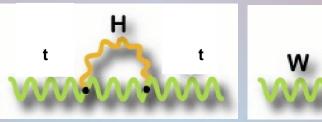


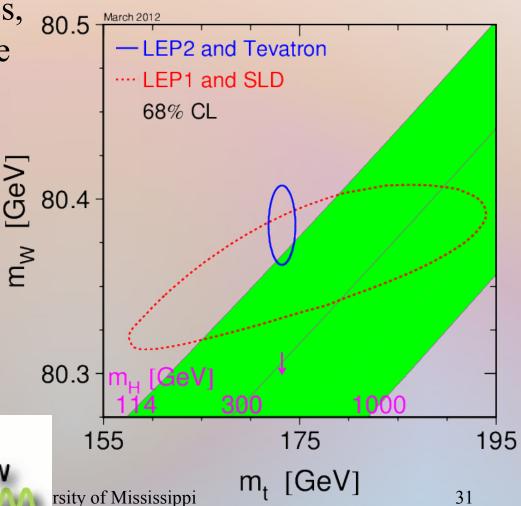


Indirect Measurements

- Because of quantum mechanics, the Higgs bosons can influence other particles without being directly detected
- More mass ⇒ more interactions with virtual Higgs bosons
- Measuring the mass of the W boson and top quark can tell us about the Higgs boson

н



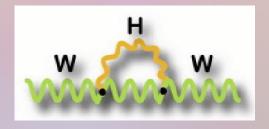


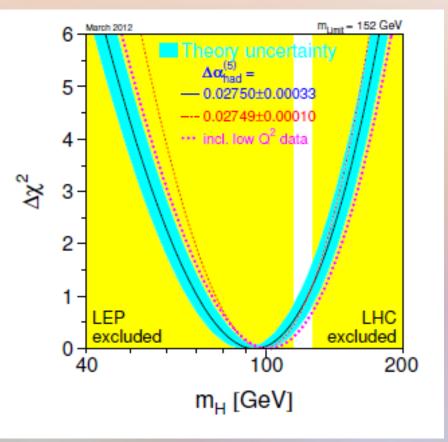


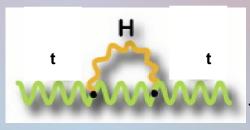


Indirect Measurements

- Fit to the precision EW data Prefers a low mass Higgs $\Rightarrow M_{\rm H} \approx 94 \, {\rm GeV}$
- \Rightarrow M_H \leq 152 GeV at 95% CL





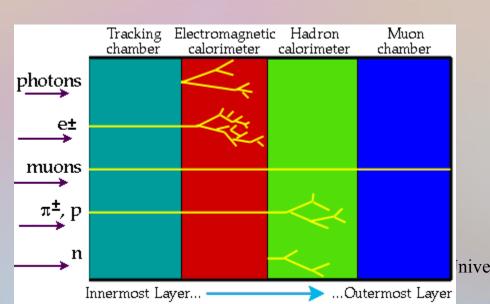


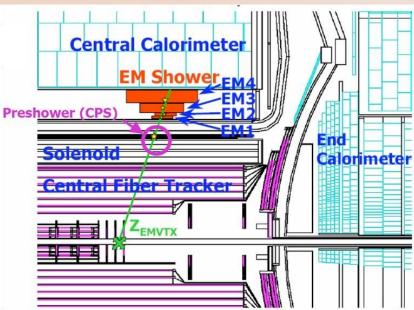




Electrons, Photons

- Lighter particles (e, γ) are stopped in the first few layers of the the calorimeter
- Hadrons make it farther into the detector





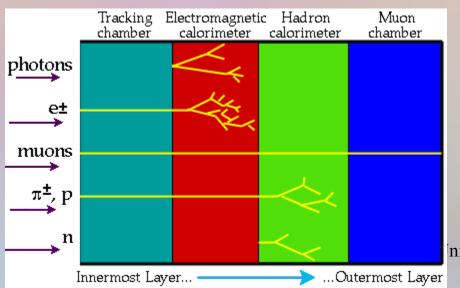
- Only charge particles leave tracks
- Track matching allow us to distinguish e from γ niversity of Mississippi

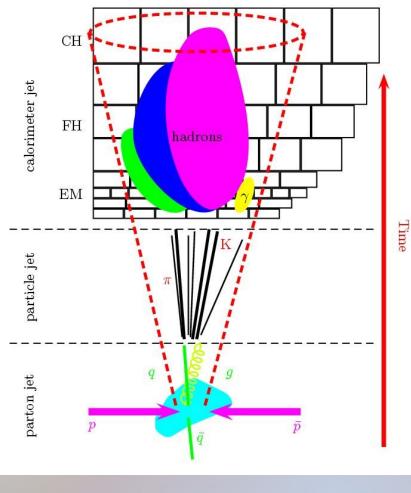


Hadronic Jets



• When a high momentum quark or gluon is produced in a collision, results in a jet of hadronic particles into the detector





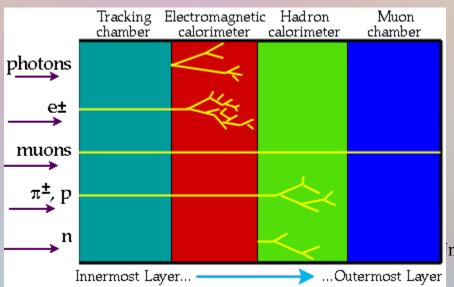
niversity of Mississippi

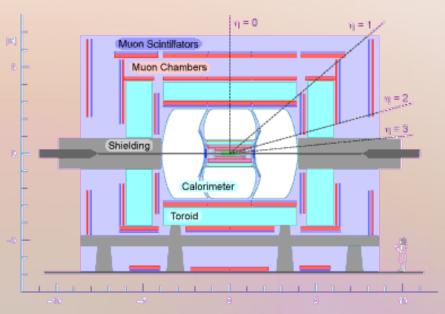




Muons

- Muons can pass through our entire detector
 - μ are massive enough not to be stopped like e or γ in EM calorimeter
 - Don't feel strong force, so hard interactions are rare





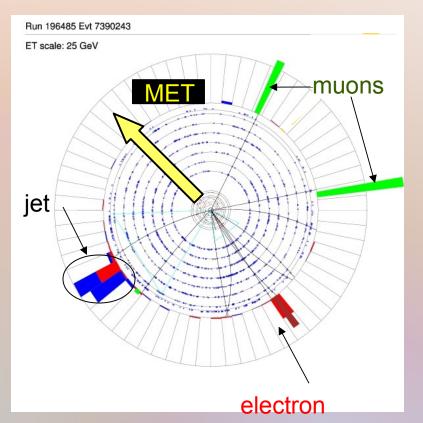
- Muons are charged, so they leave tracks
- Place trackers outside the calorimeter to detect passage, and match to a niversity of Mississippin inner tracker





Missing Transverse Energy

- Neutrinos do not interact with our detector
 - Infer their presence through missing transverse energy
 - Transverse = \perp to beamline
- To conserve momentum,
- vector sum <u>SpT</u> = 0
 - If non-zero, then either
 - Energy Mis-measurement



• Missed one or more particles (usually neutrinos)





Prob. Density

Based on log-likelihood ratio

 $-2\ln Q = LLR = -2\ln\left(\frac{L(\text{data} \mid s + b, \hat{\theta})}{L(\text{data} \mid b, \hat{\theta})}\right)$

- -L is the Poisson likelihood that the s+b (b) correctly models the data
- -20 -15 -10 -5 $-\theta$ represents the systematic uncertainties on the measurements (luminosity, energy scale, etc.)
- Calculate probability density of LLR based on models of signal and background

1–CL_b



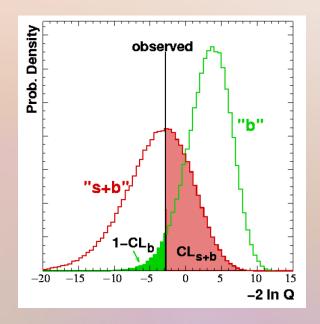


How We Set Limits

• Based on log-likelihood ratio

 $-2\ln Q = LLR = -2\ln\left(\frac{L(\text{data} \mid s + b, \hat{\theta})}{L(\text{data} \mid b, \hat{\theta})}\right)$

CL_b and CL_{s+b} are given by the integrals of the *b* and *s+b* LLR distributions above observed LLR



- Represents how well the given model agrees with data
- We vary the signal content of the s+b model to find where $CL_s = CL_{s+b}/CL_b < 0.05$
 - We exclude those points at the 95% Confidence Level (CL) James Kraus, University of Mississippi 38

