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Recent Developments in Vincia

W. Giele, D.A. Kosower, P. Skands

A. Gehrmann-De Ridder, L. Hartgring, E. Laenen,
A. Larkoski, J.J. Lopez-Villarejo, MR

Overview

introduction

extension of Vincia to hadron collisions

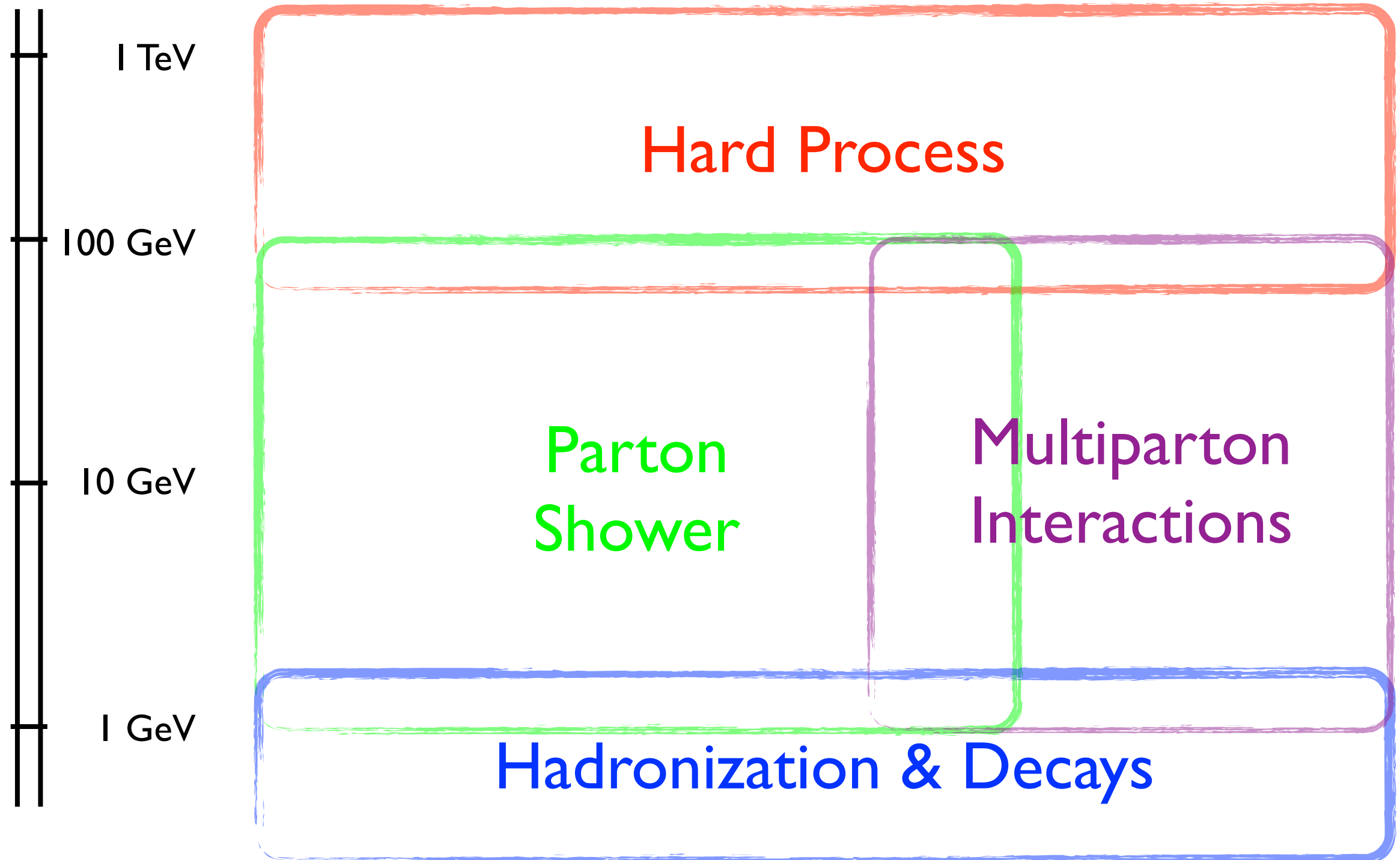
uncertainties

matching

summary and outlook

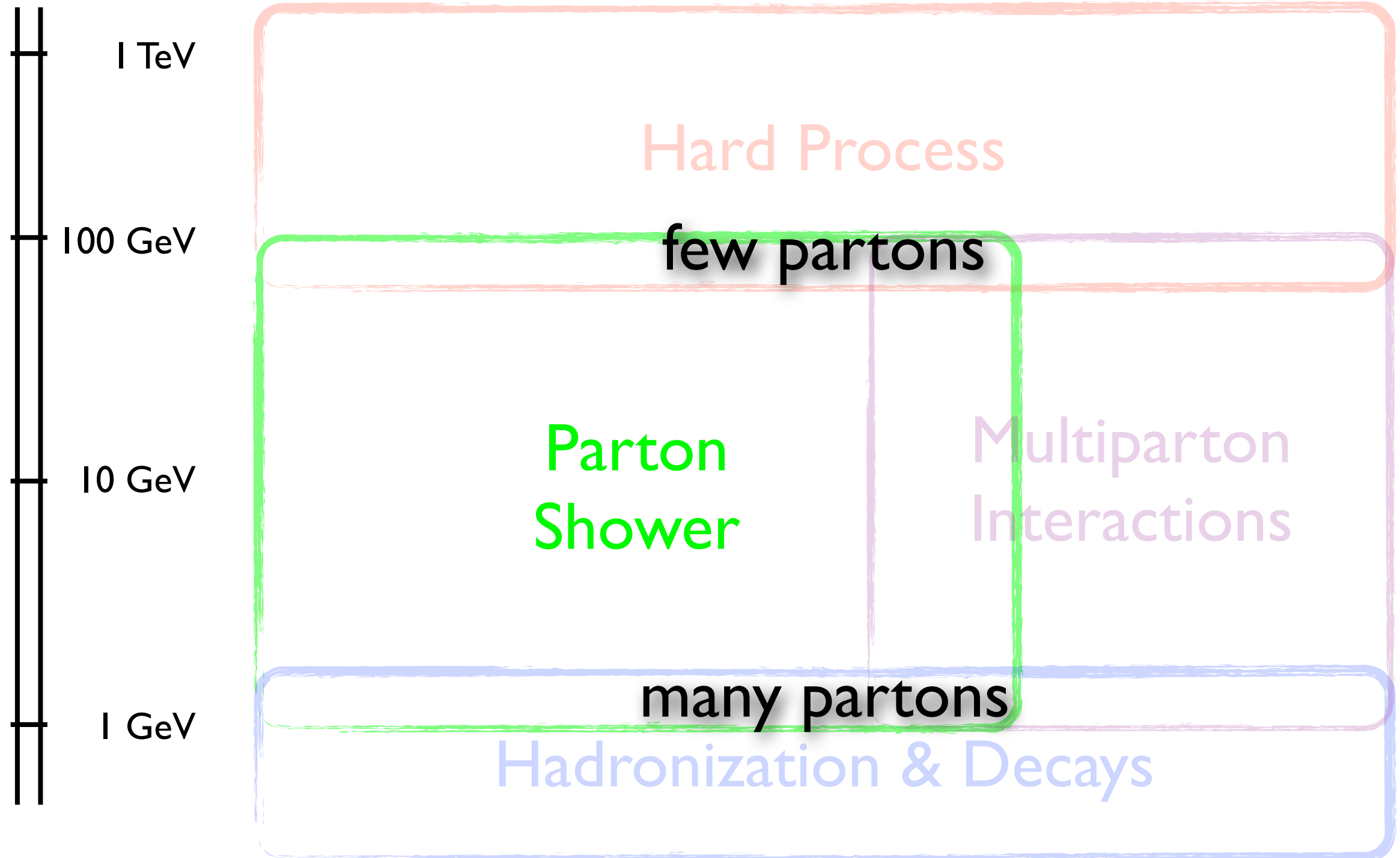
Event Generators - Cartoon

scale



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scale



The Vincia Parton Shower

W. Giele, D.A. Kosower, P. Skands
0707.3652, 1102.2126

estimates its uncertainty

facilitates matching to fixed-order

integrates into Pythia 8^[1] as a plugin

[1] T. Sjöstrand, S. Mrenna, P. Skands
0710.3820

Formalism

use approximate factorization of matrix element

$$|M_n| \approx \sum_{(ijk)} a(i, j, k) |M_{n-1}(\dots, l, K, \dots)|^2$$

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and exact factorization of phase space

$$d\Phi_n = d\Phi_{\text{ant}}(i, j, k) d\Phi_{n-1}(\dots, l, K, \dots)$$

to approximate σ_n from σ_{n-1}

Extension to Hadron Collisions

D.A. Kosower, P. Skands, MR

1210.6345

the **convolution** factorizes, not the phase space

$$\int \frac{dx_a}{x_a} \frac{dx_b}{x_b} f_a(x_a) f_b(x_b) d\Phi_{2 \rightarrow n} =$$
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ratios of parton distribution functions enter the branching probabilities

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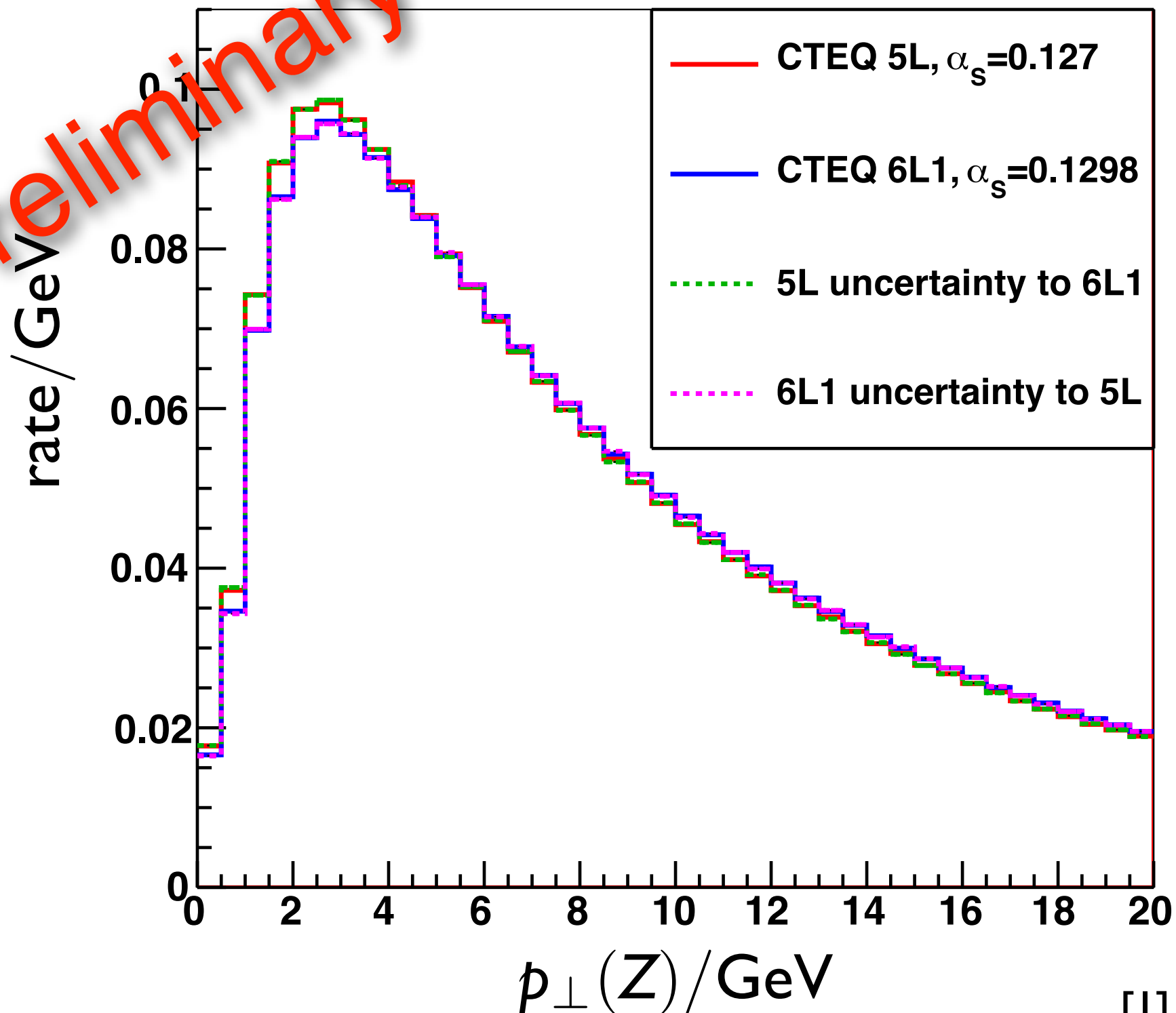
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⇒ Vincia can generate uncertainty bands

Test of Uncertainty Variations

Drell-Yan^[1] in Pythia 8.176 + Vincia

Preliminary



[1] pure parton shower

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Vincia uses a veto algorithm, generating trial branchings according to simple trial antennae

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what the shower would do without matching

Matching

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note: events still unweighted, $P_{\text{ME}} \xrightarrow{\text{unresolved}} 1$

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[1] J.J. Lopez-Villarejo, P. Skands
1109.3608

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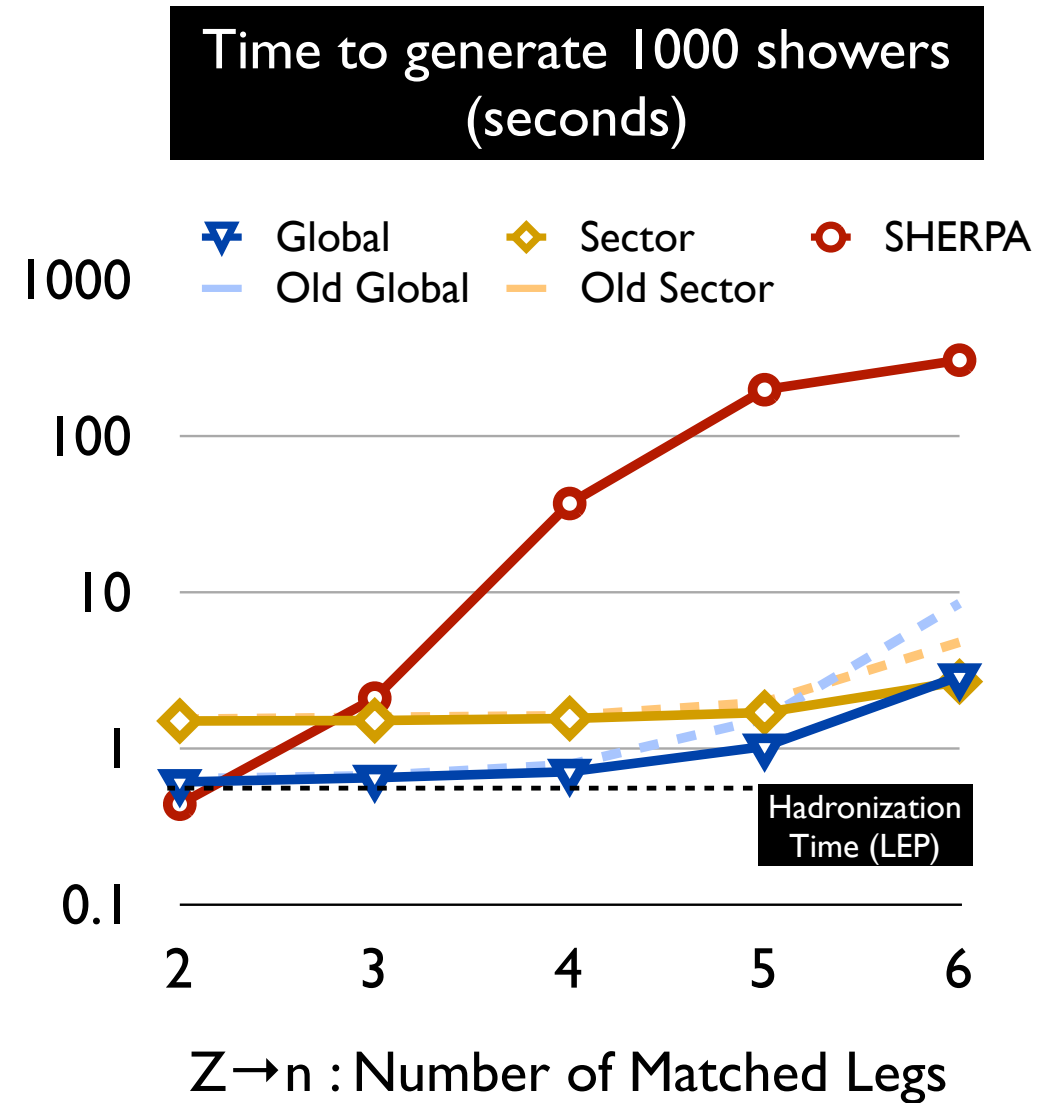
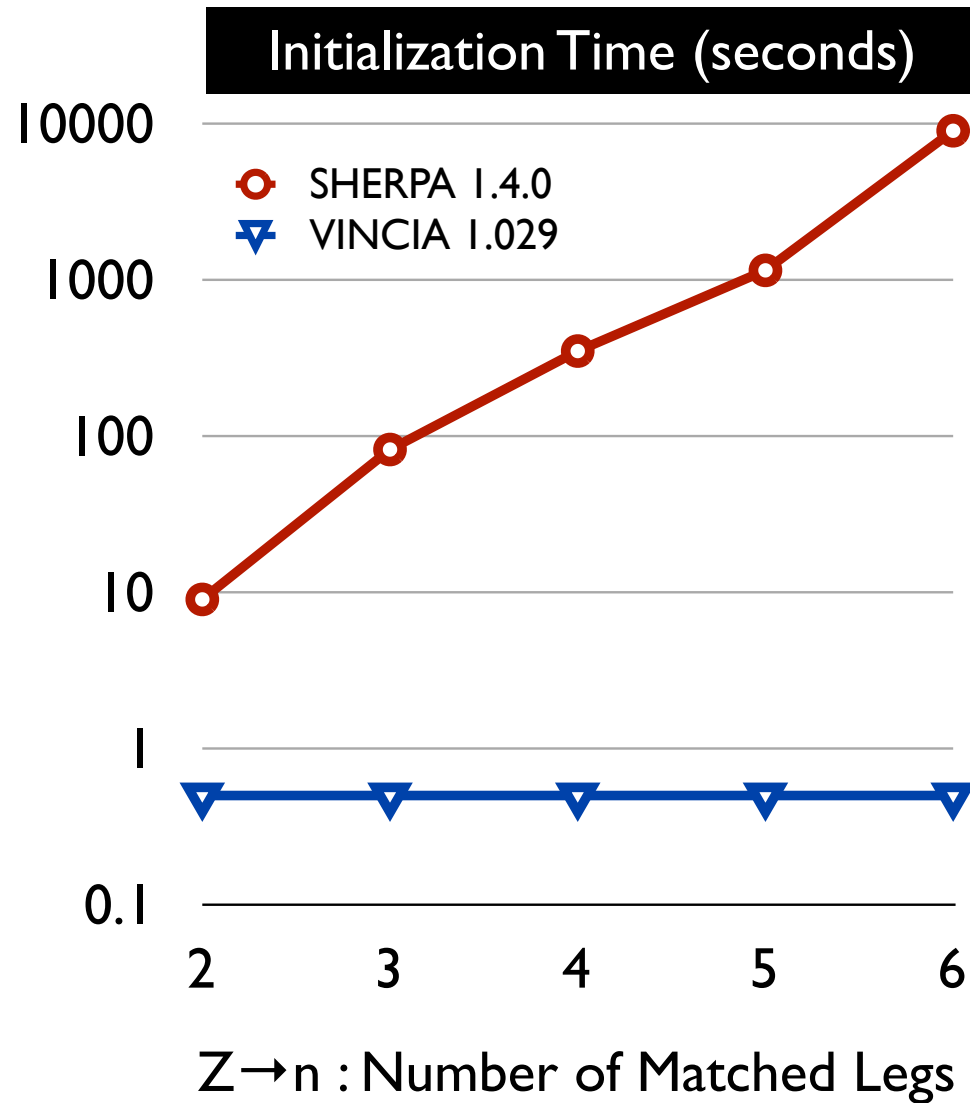
helicity-identified shower^[2]: M much cheaper

[1] J.J. Lopez-Villarejo, P. Skands
1109.3608

[2] A. Larkoski, J.J. Lopez-Villarejo, P. Skands
1301.0933

Matching

speed of tree-level matching in Z decay



One-loop matching

L. Hartgring, E. Laenen, P. Skands

1303.4974

multiplicative matching like at tree-level

events remain unweighted

tree-level matching at higher multiplicities stays

tuning of shower matched to $Z \rightarrow 3$ at one loop to event shapes (two-loop running & CMW scheme):

$$\alpha_s^{\text{NLO}}(m_Z) \approx 0.122 \quad \text{compared to} \quad \alpha_s^{\text{LO}}(m_Z) \approx 0.139$$

Summary and Outlook

Vincia has been extended to hadron collisions

unitary and efficient matching at tree- and one-loop level has been demonstrated for $e^+ e^-$ collisions

next: matching for hadron collisions

Thanks for your attention

Backup

One-loop matching

L. Hartgring, E. Laenen, P. Skands

1303.4974

analogous to tree-level matching:

$$\text{matched} = (I + V_3) \text{ (shower approximation)}$$

rate for exactly three resolved partons at Q_{had} :

$$\text{matched} = |M_3^0|^2 + 2\Re(M_3^0 M_3^{1*}) + \int_0^{Q_{\text{had}}^2} \frac{d\Phi_4}{d\Phi_3} |M_4^0|^2$$

$$\text{(s. approx.)} = (I + V_2) |M_3^0|^2 \Delta_2(m_Z^2, Q_3^2) \Delta_3(Q_3^2, Q_{\text{had}}^2)$$

One-loop matching

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$$\begin{aligned} V_3 = & -V_2 + \frac{2\Re(M_3^0 M_3^{1*})}{|M_3^0|^2} \\ & + \int_{Q_3^2}^{m_2^2} d\Phi_{\text{ant}} a(2 \rightarrow 3) + \sum_{3 \rightarrow 4} \int_0^{m_3^2} d\Phi_{\text{ant}} a(3 \rightarrow 4) \\ & + \int_0^{Q_{\text{had}}^2} \frac{d\Phi_4}{d\Phi_3} \frac{|M_4^0|^2}{|M_3^0|^2} - \sum_{3 \rightarrow 4} \int_0^{Q_{\text{had}}^2} d\Phi_{\text{ant}} a(3 \rightarrow 4) \end{aligned}$$

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shower is matched at tree-level

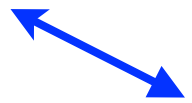
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singularities cancel


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integral over $3 \rightarrow 4$ antennae has two pieces:

$$\sum_{3 \rightarrow 4} \int_0^{m_3^2} d\Phi_{\text{ant}} [a^{\text{std}}(3 \rightarrow 4) + \delta a(3 \rightarrow 4)]$$

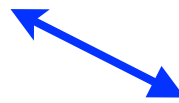
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integrals $2 \rightarrow 3$ and $3 \rightarrow 4$ are **two pieces:**

universal, divergent, integrated analytically



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$$\sum_{3 \rightarrow 4} \int_0^{m_3^2} d\Phi_{\text{ant}} [a^{\text{std}}(3 \rightarrow 4) + \delta a(3 \rightarrow 4)]$$

process-dependent, finite, integrated numerically