

Vector boson + jets from k_t - dependent parton showers

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- power - suppressed terms in k_t dependent shower
- W+n jet production using
 - new small x improved TMD gluon density including uncertainties
 - comparison with W+2 jet measurements
 - observables from double parton scattering

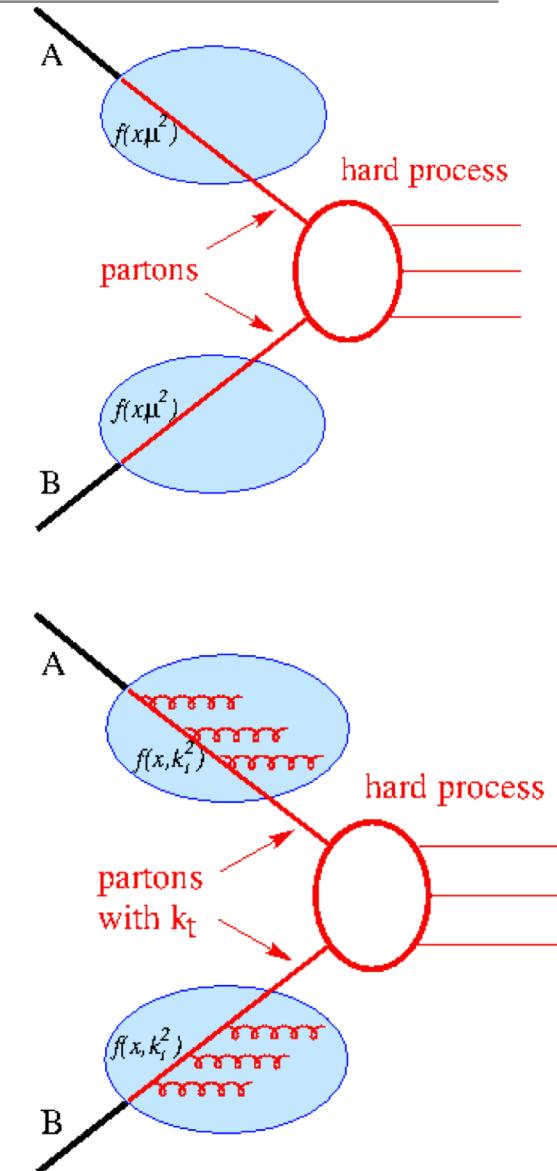
The past ...

- The standard ansatz: collinear factorization

$$\sigma = f_i^A(x_1, \mu^2) \hat{\sigma}(i + j \rightarrow X) f_j^B(x_2, \mu^2)$$

- even at NLO kinematic constraints lead to inconsistencies (see talk by S. Dooling) especially away from central region.

- The extension: transverse momentum dependent factorization
 - factorization of TMDs proven by J. Collins (CUP 2011)
 - gluon TMD at small x determined, but it is still a challenge to include quarks



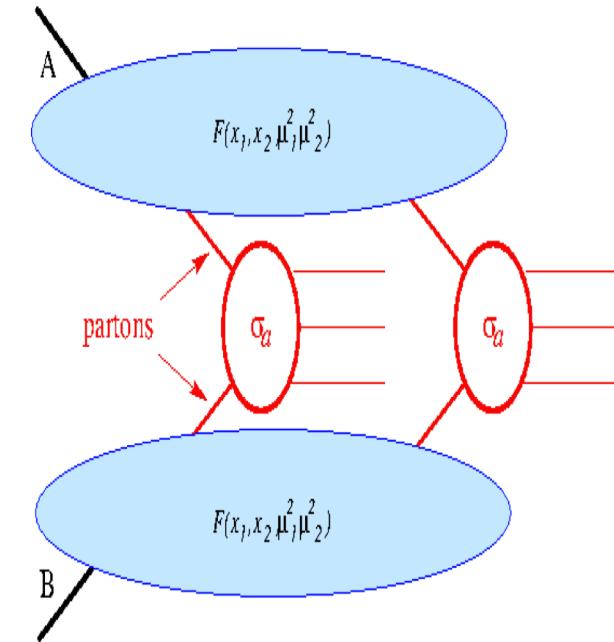
The past ...

- **The limitations:**

$$\sigma = c_1 f_g f_g + c_2 f_q f_q + c_3 f_g F + c_4 F F$$

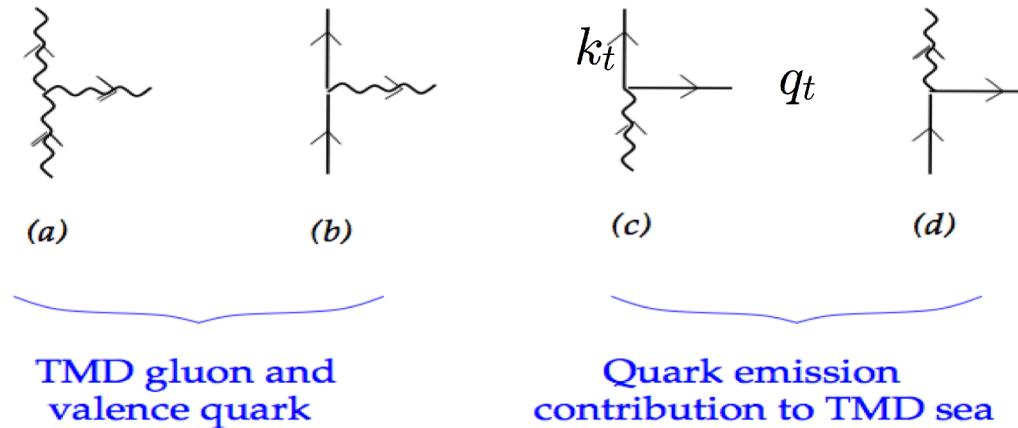
- last terms are suppressed by powers of μ
- ansatz of collinear factorization
- inclusion of power suppressed terms leads to double parton scattering

- **Can one do differently ?**



Treatment of power suppressed terms

- un-integrated sea-quark distribution (see talk by M. Hentschinski *TMD quark distributions at small x* in Structure function session and Hautmann, Hentschinski, Jung, Nucl.Phys. B865 (2012) 54-66, arXiv 1205.6358)



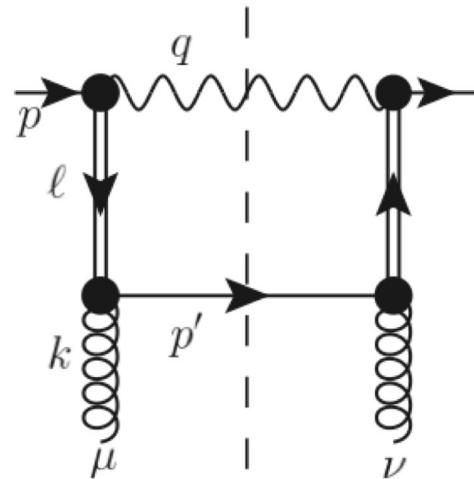
- quark sea: flavor singlet evolution coupled to gluons at small x via

$$\mathcal{P}_{g \rightarrow q}(z, q_t, k_t) = P_{qg, DGLAP}(z) \left(1 + \sum_{n=0}^{\infty} b_n(z) \left(\frac{k_t^2}{q_t^2} \right)^n \right)$$

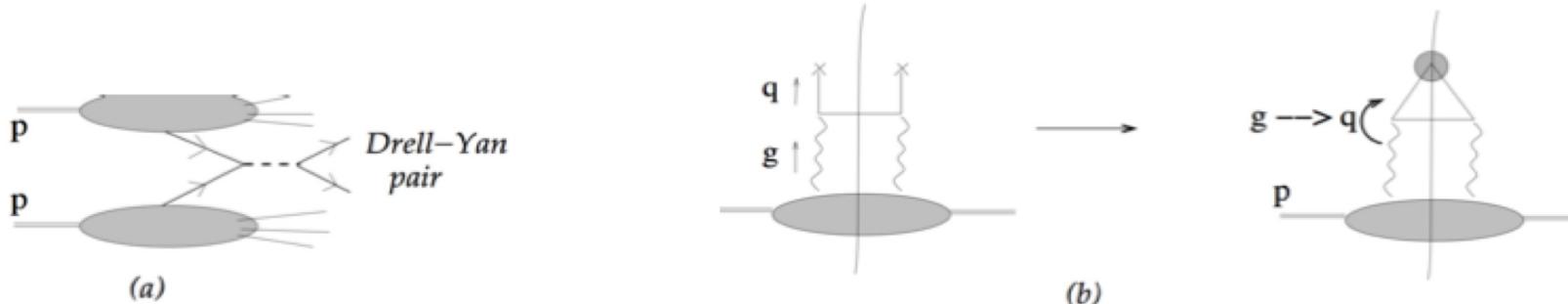
- all b_n known: $\mathcal{P}_{g \rightarrow q}$ computed in closed form (pos definite) in Catani, Hautmann 1994, Ciafaloni et al 2005-2006 in small x factorization
- Note: power suppressed terms only in small x improved TMDs (not covered by J. Collins factorization.)**

Vector bosons + jets at high energies

- Partonic matrix elements (gauge invariant, despite off-shell parton)
(see Ball & Marzani NPB814 (2009) 246, Hautmann, Hentschinski, Jung NPB865 (2012) 54)



- DY production and $g \rightarrow q$ splitting contribution to sea quark distribution



CCFM evolution and k_t – dependent shower

- Color coherence requires angular ordering instead of p_t ordering ...

$$q_i > z_{i-1} q_{i-1} \quad \text{with} \quad q_i = \frac{p_{ti}}{1 - z_i}$$

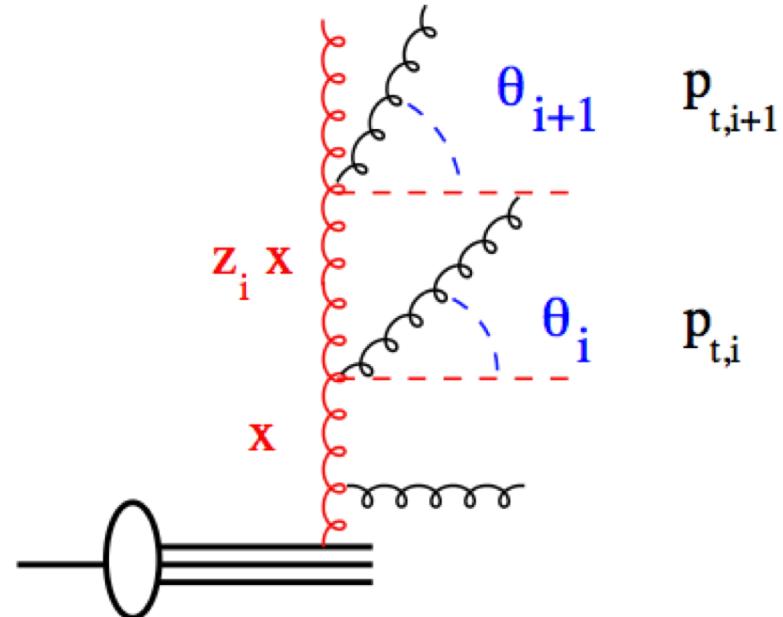
→ recover DGLAP with q ordering

at medium and large x

→ HERWIG uses: $q_i > q_{i-1}$

→ at small x , no restriction on q

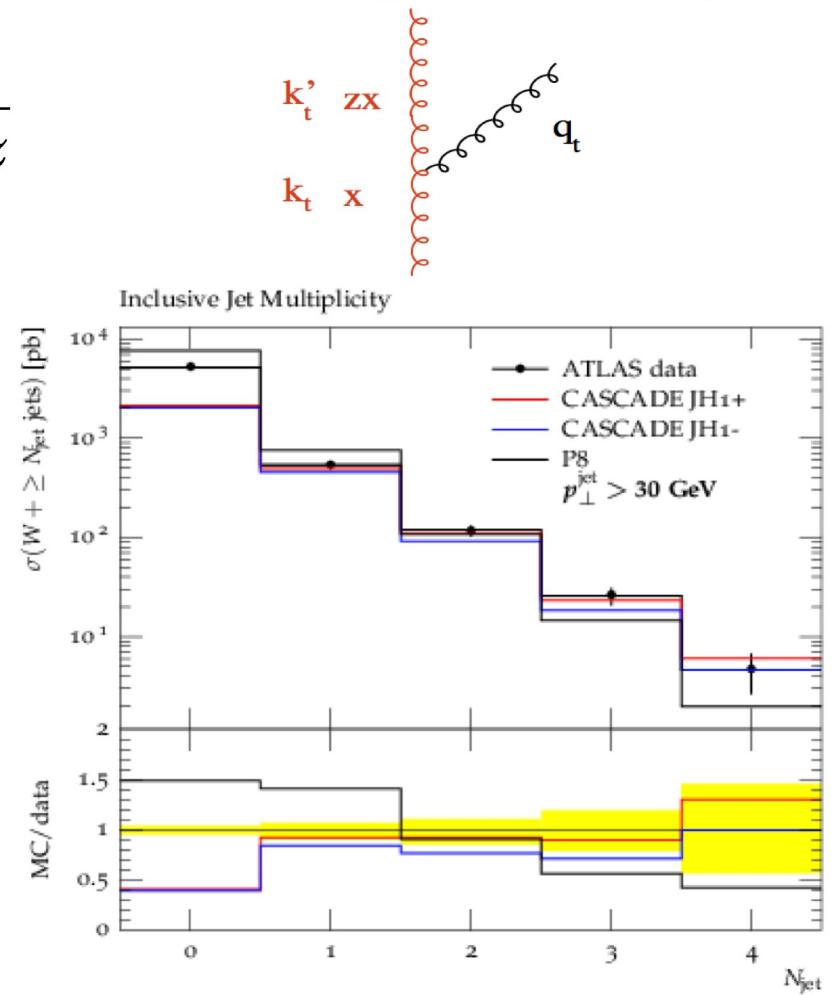
p_{ti} can perform a random walk



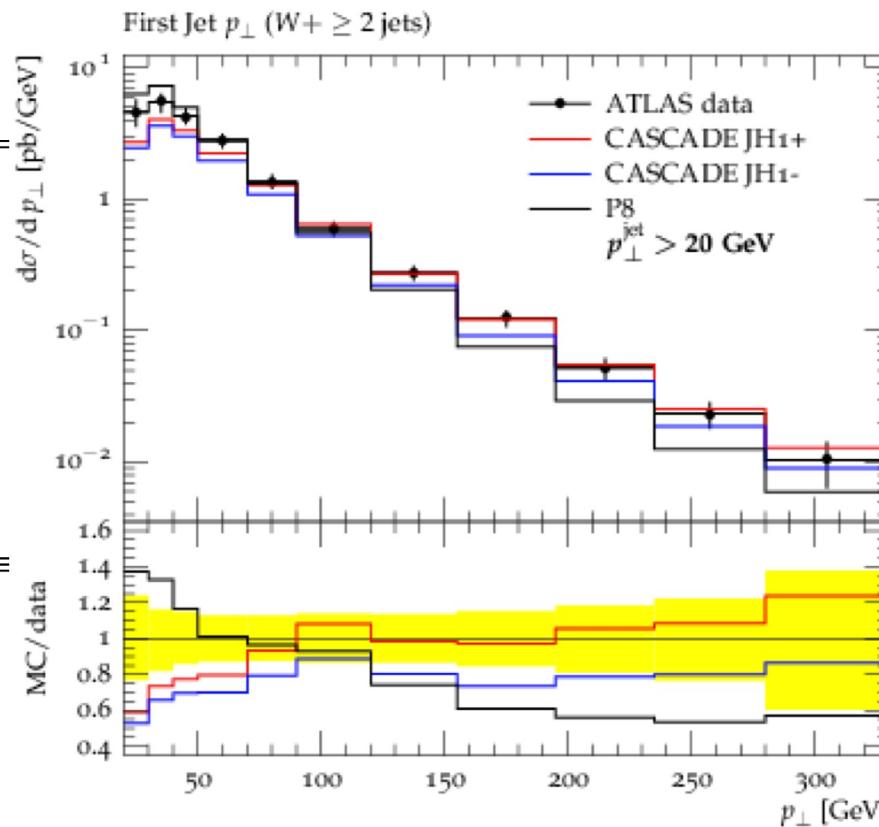
→ Catani Ciafaloni Fiorani Marchesini evolution forms a bridge between DGLAP and BFKL evolution

W + jets

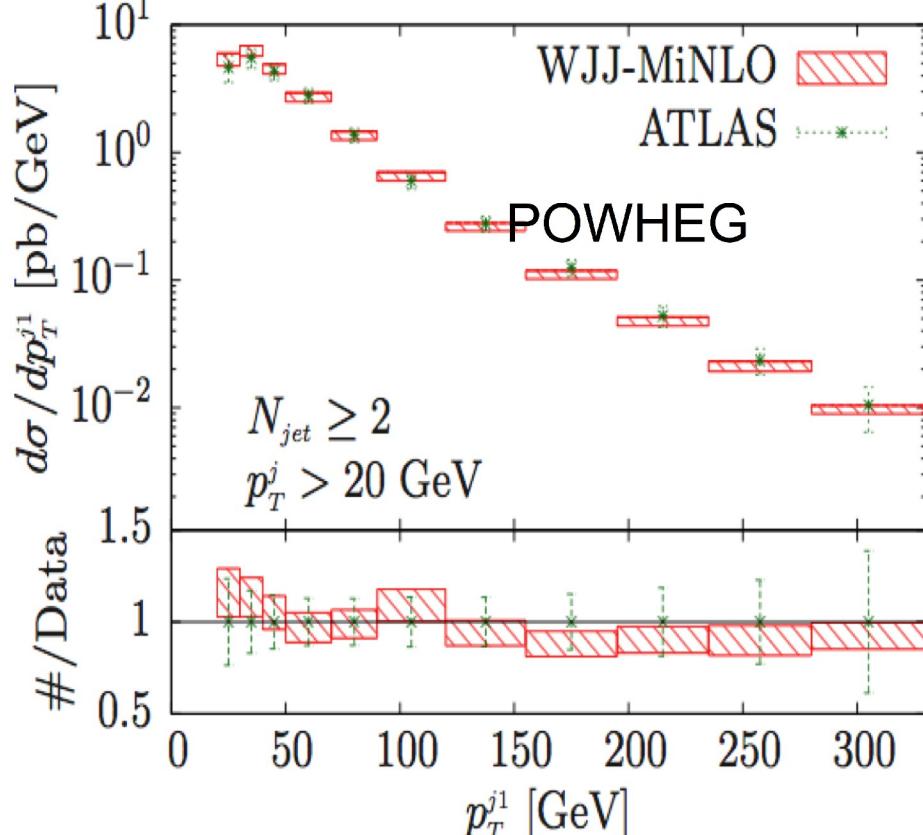
- using CCFM gluon TMD, determined from HERA F_2 including experimental and model uncertainties
- initial parton shower according to CCFM evolution in angular ordered phase space
 - $q_i > z_{i-1} q_{i-1}$ with $q_i = \frac{p_{ti}}{1-z}$
 - no p_t constraint at small x
 - jets can have large p_t
- Jet multiplicities are reproduced
 - 1 jet → from ME
 - 2-4 jets from shower
- **Note:** PYTHIA with p_t -ordered shower cannot predict higher jet multiplicities



$W + n\text{-jets}$: k_t shower vrs NLO

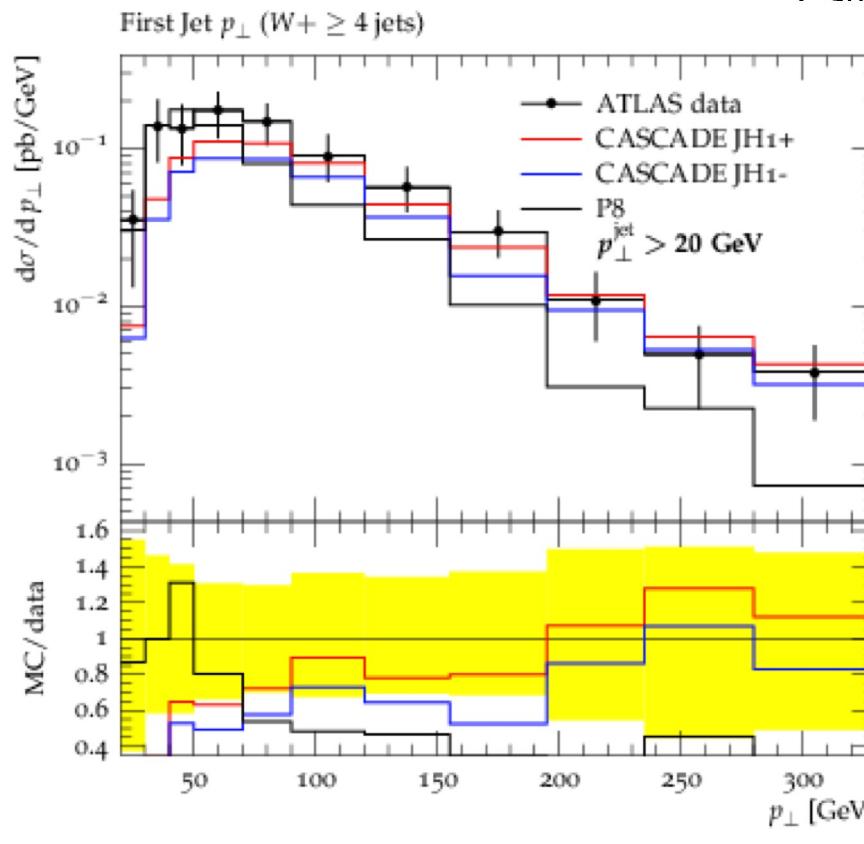
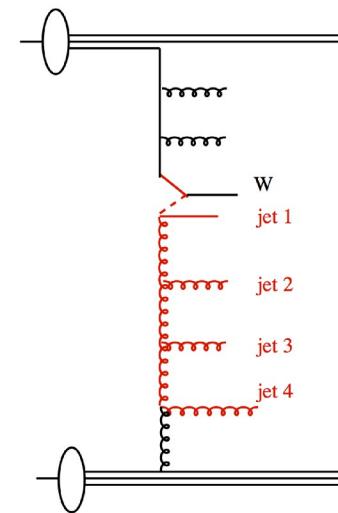


Campbell, Ellis, Nason, Zanderighi, arXiv 1305.5447

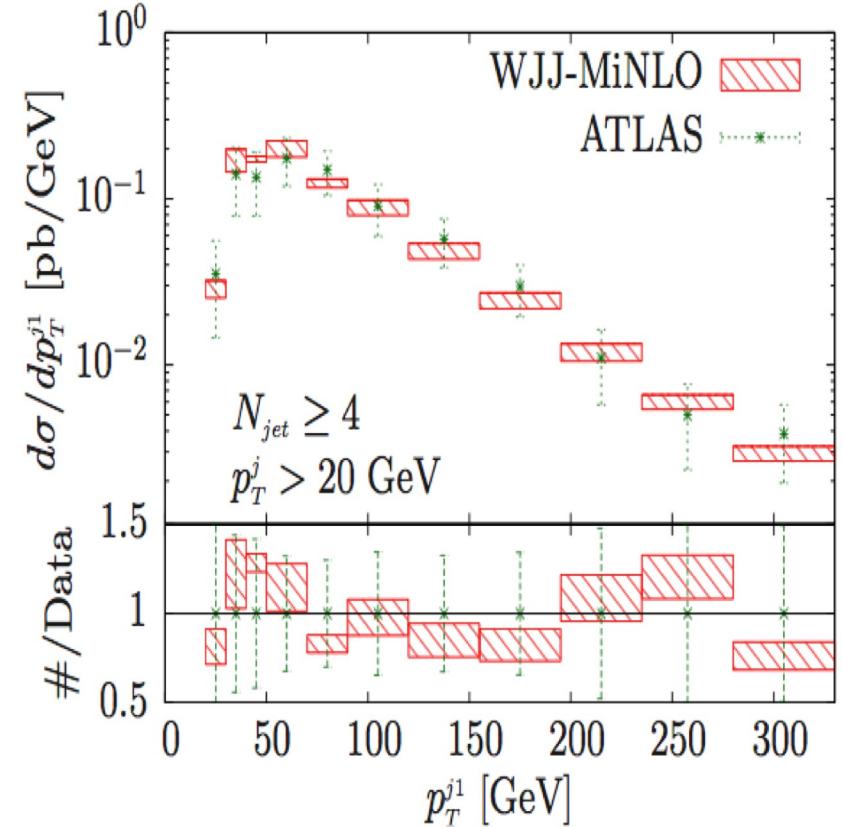


- 1st jet in $W + 2$ jet events:
 - off shell ME + CCFM k_t - shower (CASCADE) comparable with NLO $W + 2$ jet (POWHEG)
 - uncertainties studied in CASCADE: pdf and scale uncertainties
 - PYTHIA + shower starts to fail at large p_T

$W + n\text{-jets}$: k_T shower vrs NLO

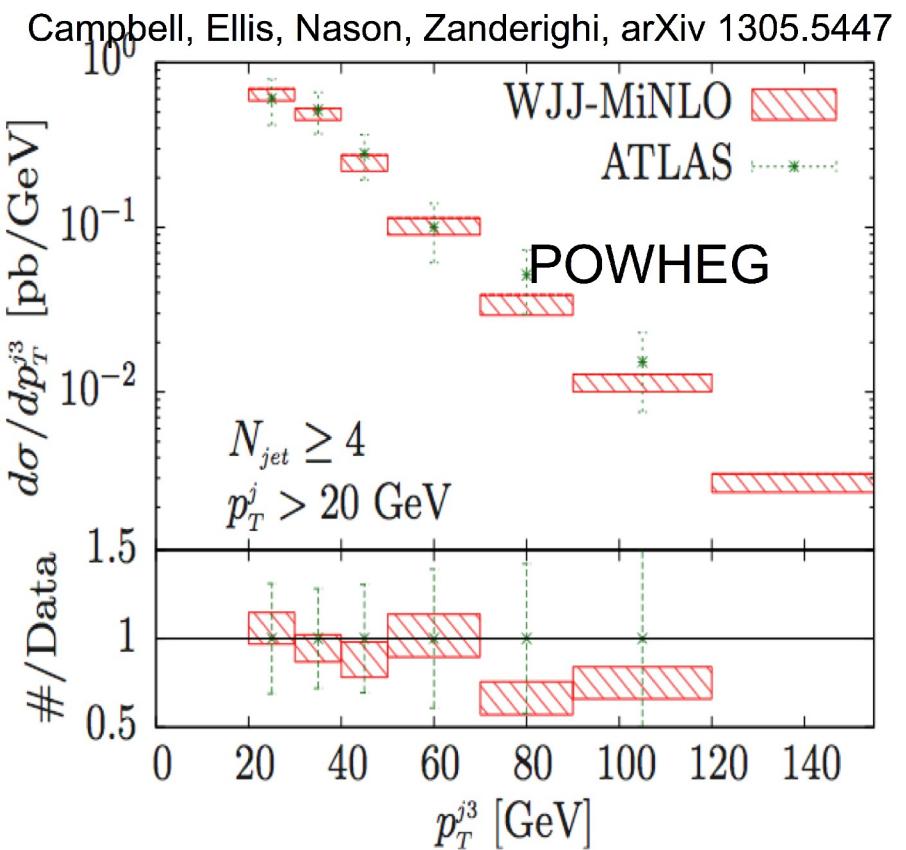
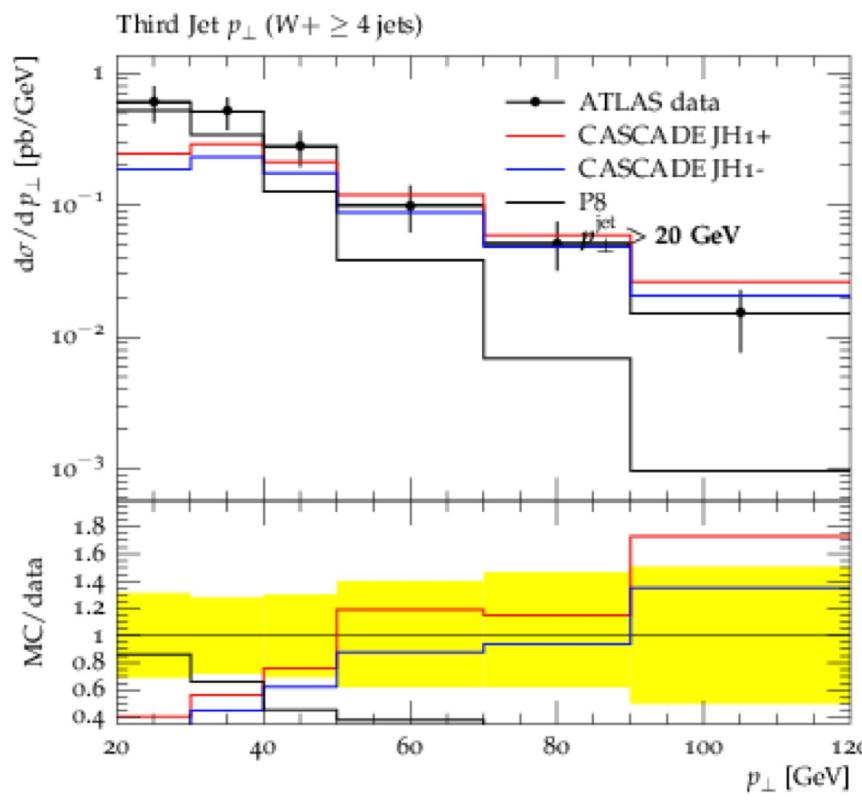
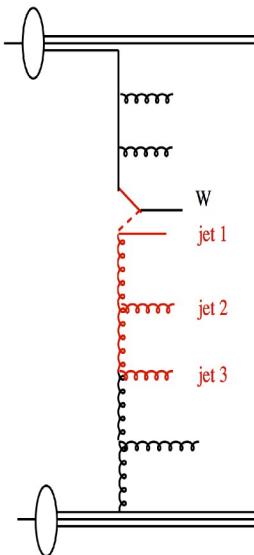


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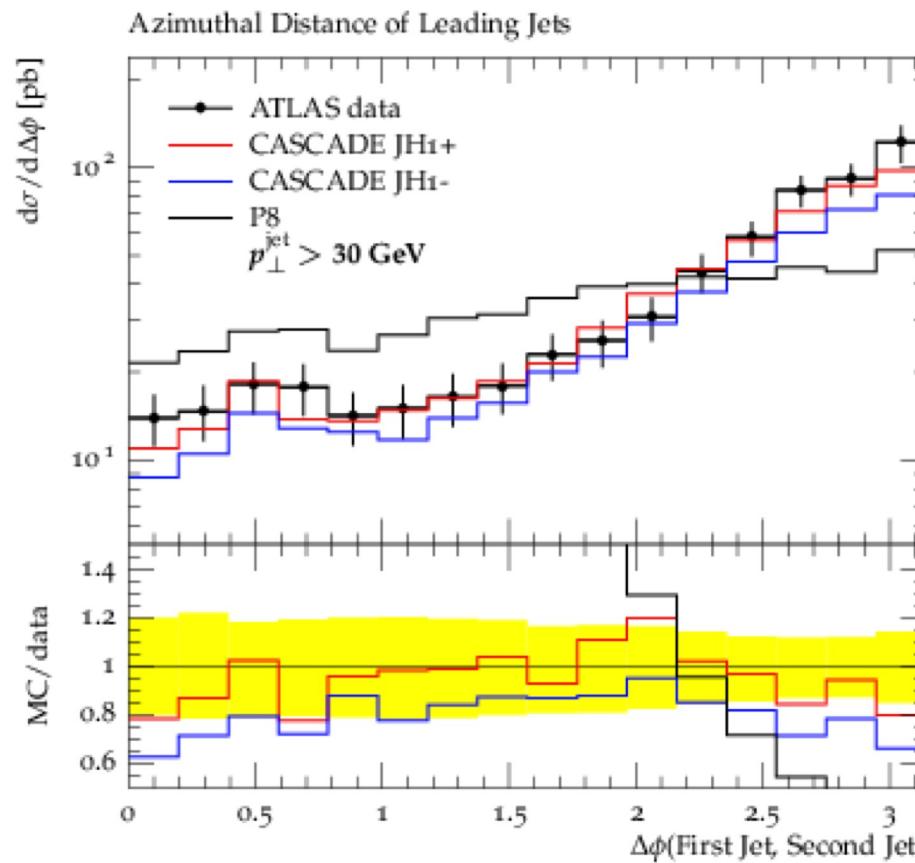
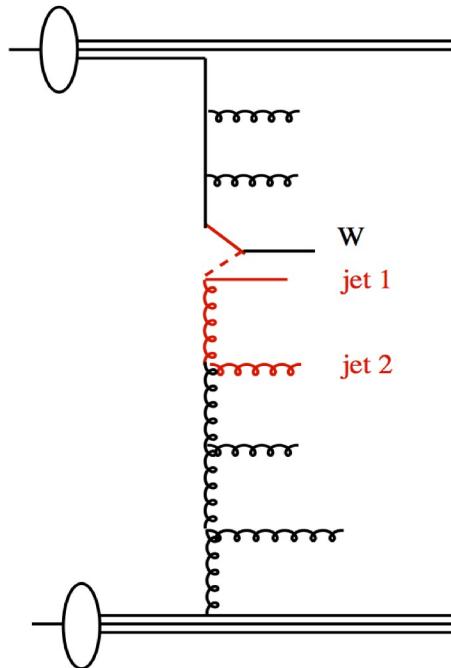
- off-shell ME + CCFM k_T -shower (CASCADE) comparable with NLO $W+4\text{jet}$
- first jet comes from hard process, other jets partially from shower
 - CCFM k_T -shower works fine even for high pt
 - P8 shower cannot describe shape

$W + n\text{-jets}$: k_T shower vrs NLO



- off-shell ME + CCFM k_T - shower predicts correct x-section and shape for 3rd jet (similar to NLO $W+2\text{jet}$) !
 - 3rd jet comes from CCFM k_T - shower
 - collinear shower pt ordered shower (PYTHIA) fails to describe shape

$W + n\text{-jets}$: k_t shower vrs NLO

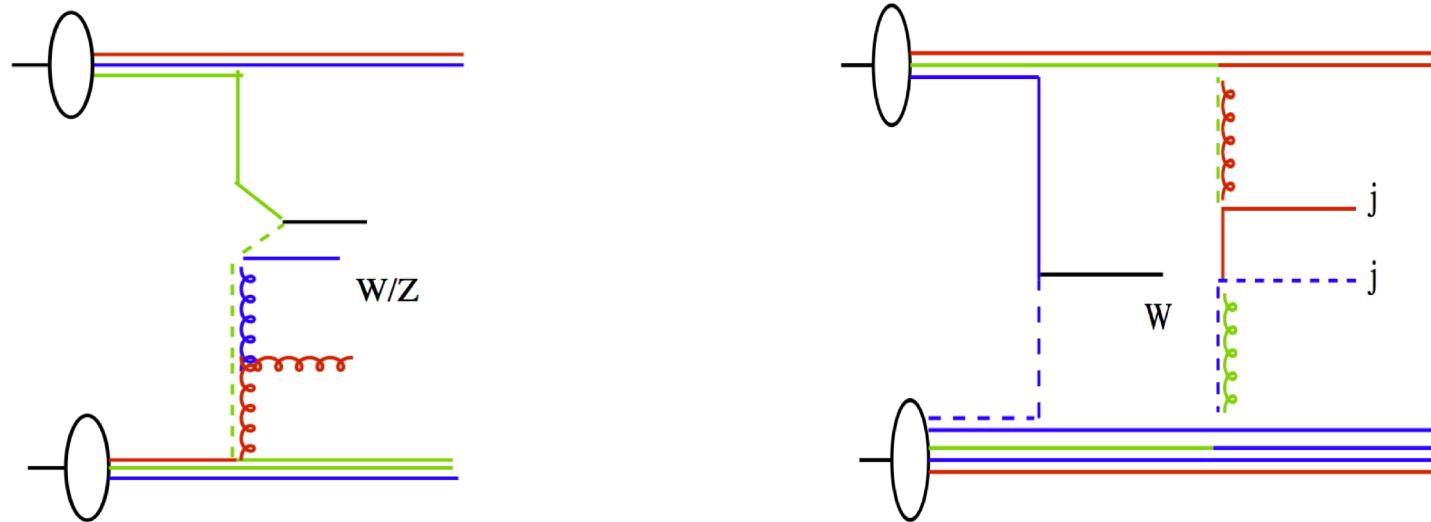


- off-shell ME + CCFM k_t -shower for x-section and shape for $\Delta\phi$ between first 2 jets agrees with measurements within uncertainties:
 - sensitive probe of shower:
 - back to back region and decorrelation region well reproduced !
 - not described by collinear shower pt ordered shower (PYTHIA)

What is the gain ?

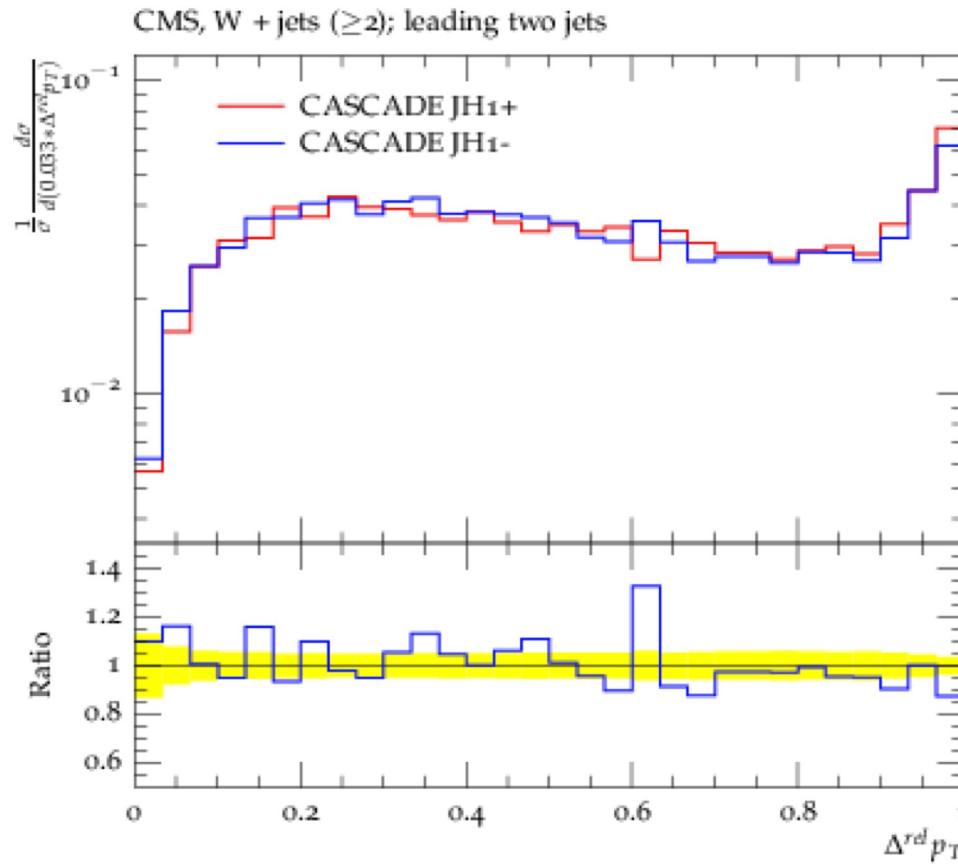
- CCFM gluon TMD and k_t dependent shower with off shell ME give similar results as NLO matched with collinear shower
- calculation arranged in a very efficiency way → fast calculation
- jet production from TMD and k_t dependent shower **extendable to any number of jets** without further adjustment and tuning
 - CCFM + k_t dependent shower describes well high pt jet production
- Advantage of CCFM+ k_t dependent shower:
 - matching with $2 \rightarrow n$ off-shell parton calculation (automated method, see *talk by A. van Hameren, small x session*)
 - opens possibility for full LHC phenomenology of QCD, EWK and BSM processes

$W + 2 \text{ jet}$: signal for double-parton scattering ?

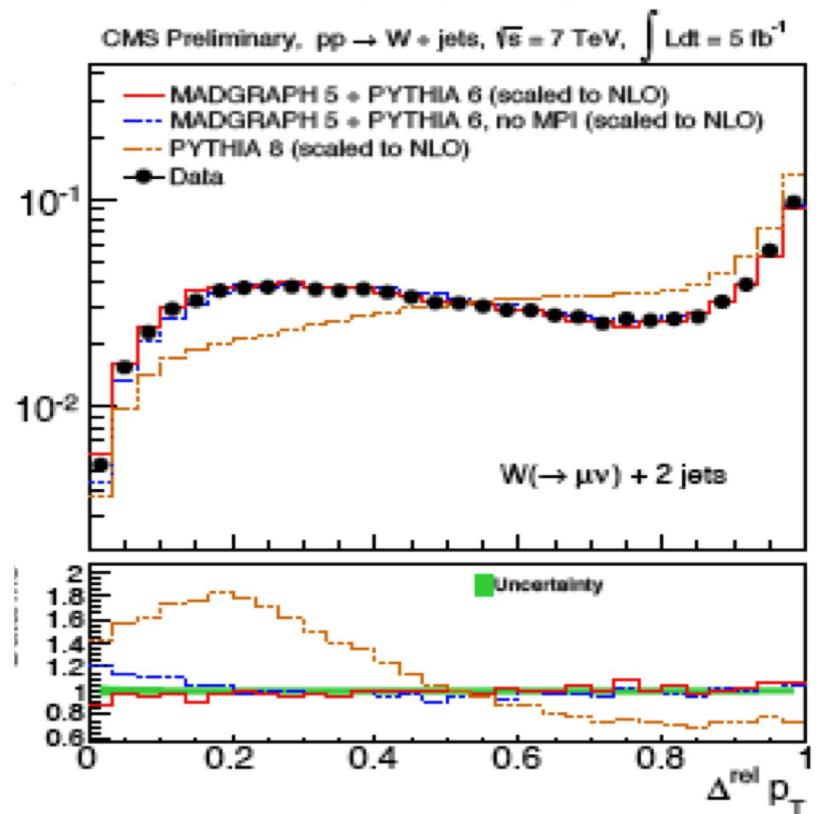


- DPS signal: de-correlated jets compared to W
 - what is the contribution from single chains ?
 - are jets just power-suppressed terms from evolution or do they come from independent scatterings ?

W+2 jet: signal for double-parton scattering ?

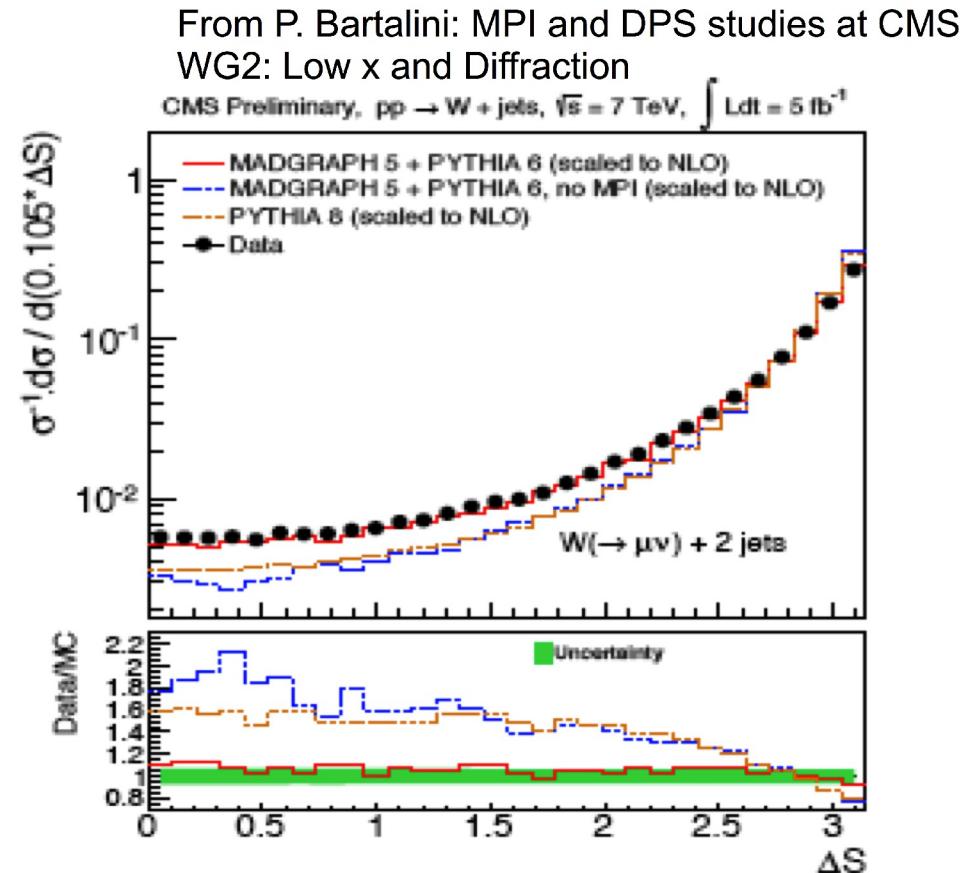
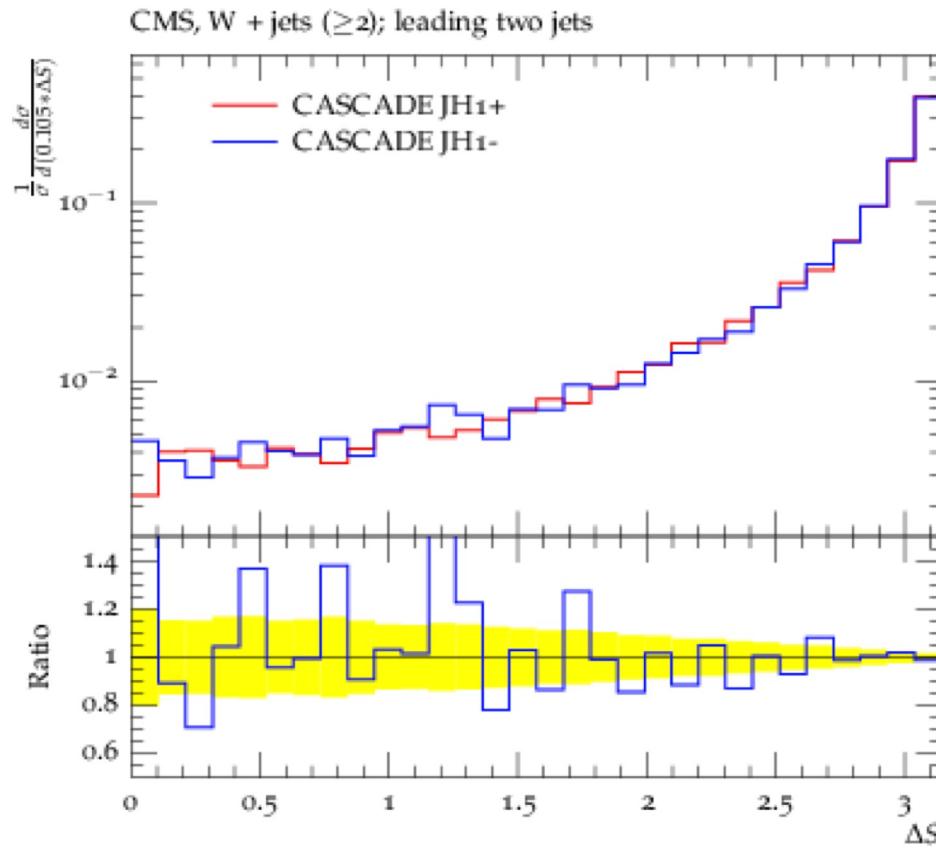


From P. Bartalini: MPI and DPS studies at CMS
WG2: Low x and Diffraction



- off-shell ME & CCFM + k_t shower predict a similar shape as seen in latest CMS measurement

W+2 jet: signal for double-parton scattering ?

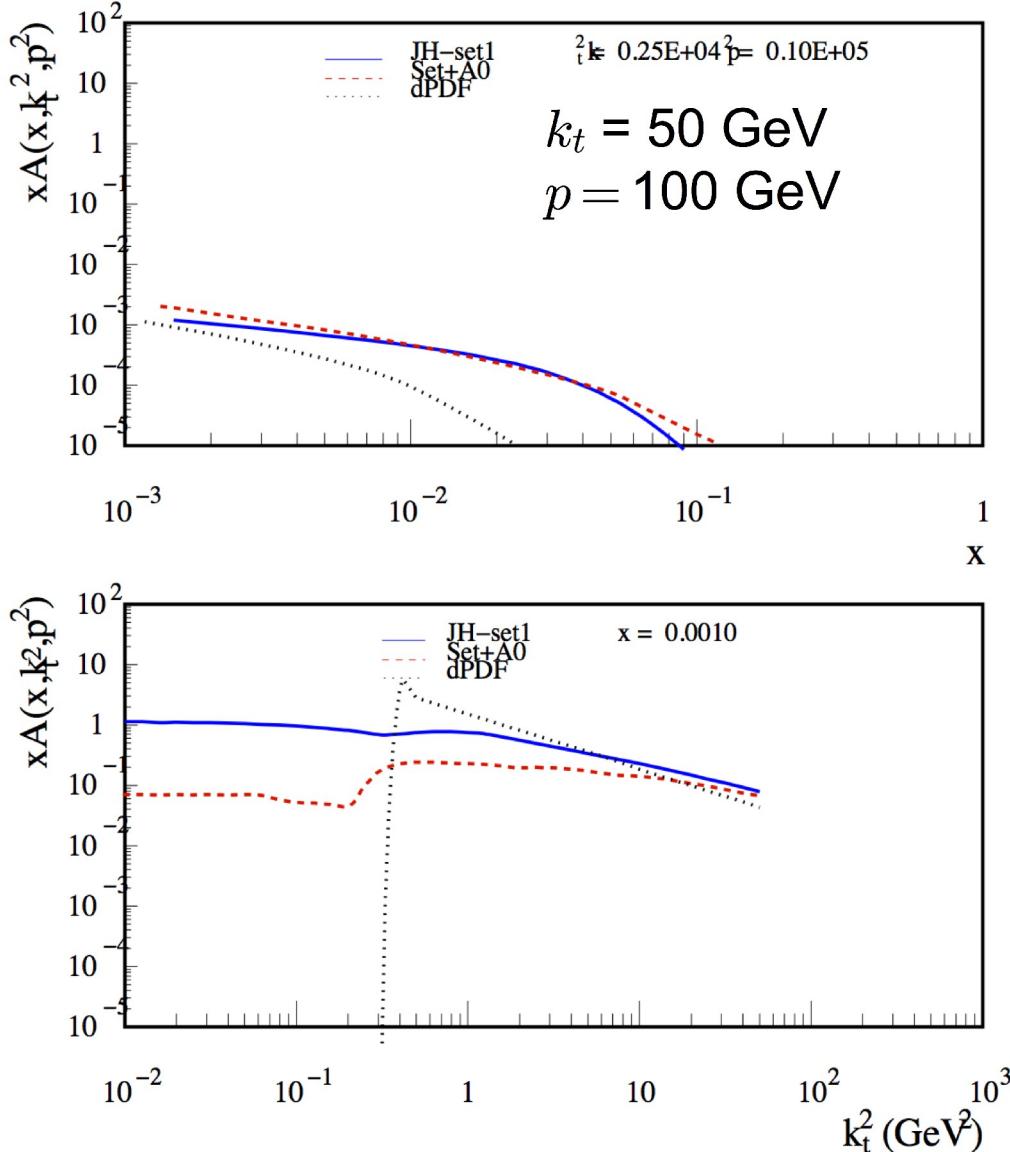


- off-shell ME & CCFM + k_t shower predict a similar shape as seen in latest CMS measurement.
 - how much room for DPS is left in frame of high-energy factorization ?

Conclusion

- terms which are neglected in collinear factorization can be important:
 - some are included in high energy factorization and small x improved TMD gluon densities
- using latest gluon TMD including uncertainties
 - reasonable description of $W+n$ jet observables are obtained
 - description is similar to $W+2$ jet NLO
- k_t dependent CCFM shower is appropriate for multi-jet kinematics
 - shape and $p_t x$ -sections of 4 hard jets (most coming from shower) are well reproduced
- off-shell ME + CCFM TMD with k_t dependent CCFM shower predicts shape of double-parton scattering variables close to latest measurements

CCFM gluon from F_2 fit

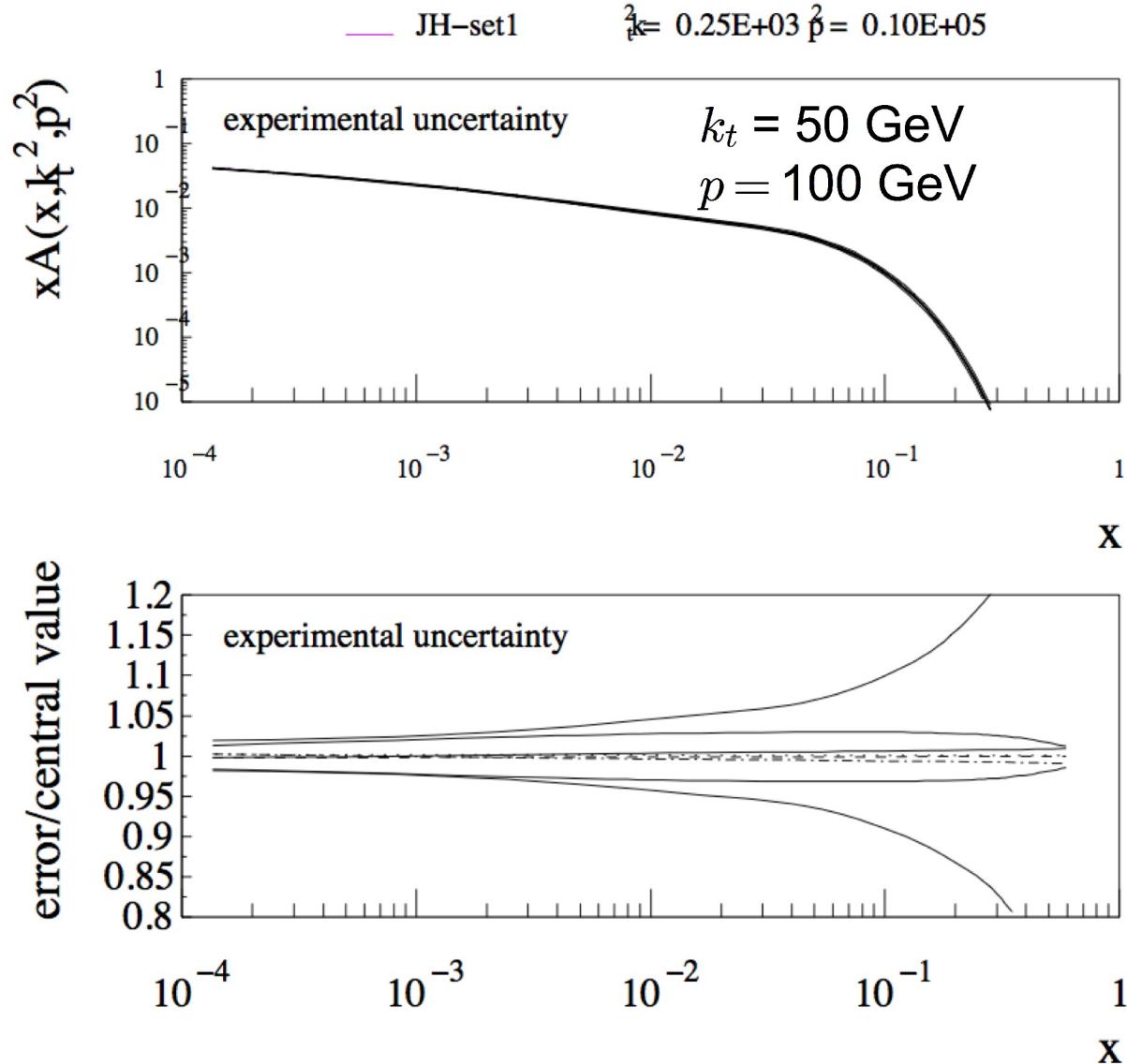


- Fit function:

$$\begin{aligned} \mathcal{A}_0(x) &= N_g x^{-B_g} (1-x)^{C_g} \\ &\times (1 - D_g x \\ &+ E_g \sqrt{x} + F_g x^2) \end{aligned}$$

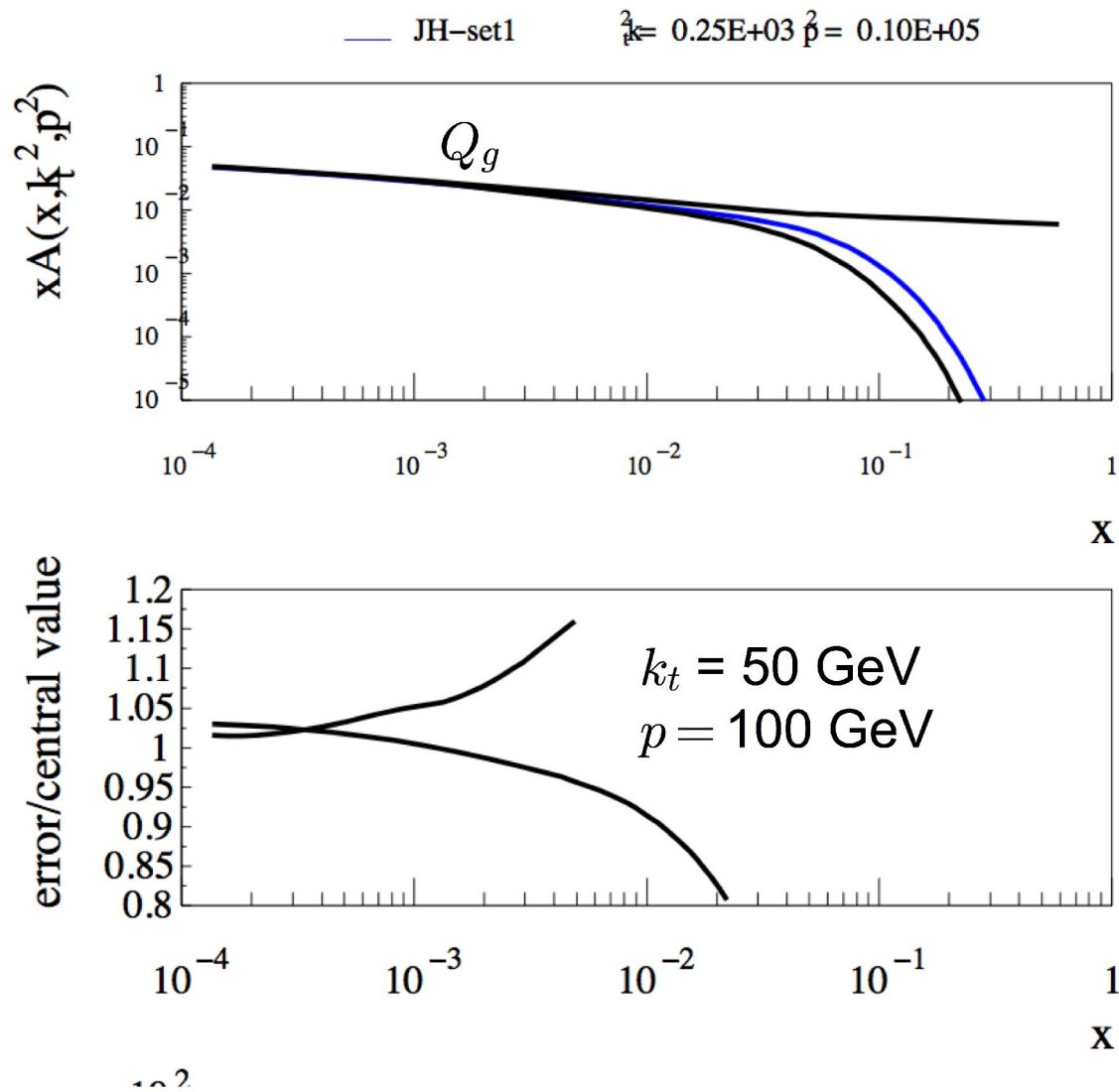
- 2-loop α_s
- gluon splitting function with non-singular terms
- $Q^2 > 5 \text{ GeV}, x < 0.005$
 - new fit gives $\chi^2/ndf \sim 1.2$
 - depending on number of parameters
- details are different from previous uPDF set A_0 and from derivative of collinear gluon with CTEQ6

uncertainties of CCFM gluon



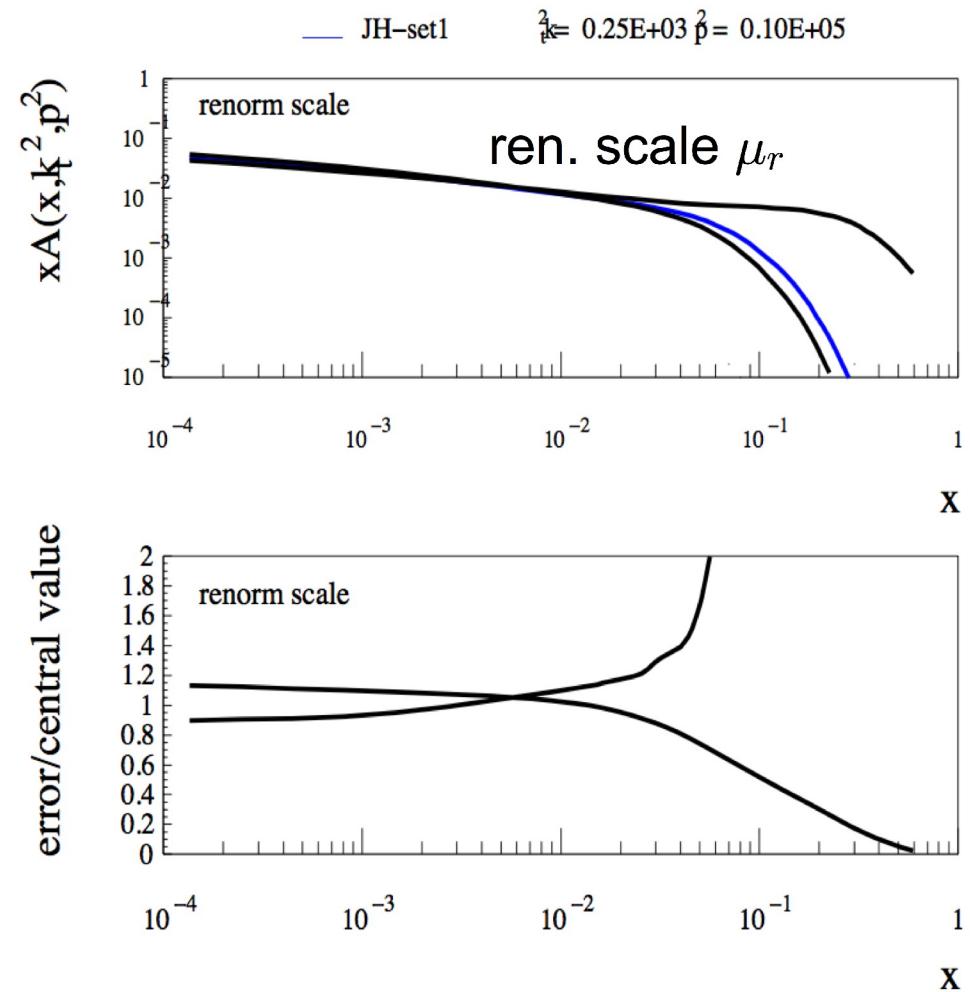
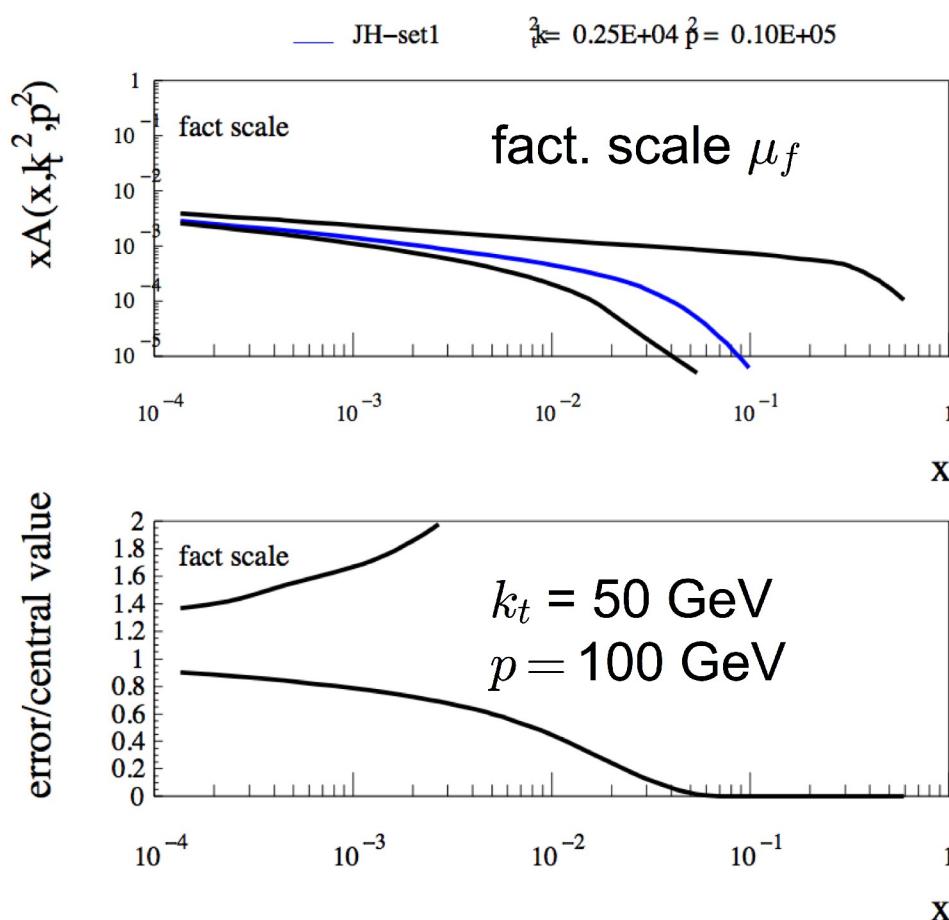
- experimental uncertainties result in 10-20 % for gluon uncertainty at medium and large x
- uncertainties at small x very small

uncertainties of CCFM gluon



- for the first time:
estimation theory
uncertainties in TMDs
- uncertainties from starting scale Q_g
 - small at small x
 - explode at large x ,
because of no
constraint from data:
 $Q^2 > 5 \text{ GeV}, x < 0.005$

model uncertainties of CCFM gluon



- factorization scale and renormalization scale varied by factor 2
 - what matters is x-section (convolution of σ with pdf !)
 - at large x no constraint (since $x < 0.005$)