BFKL Evolution as a Communicator between Small and Large Energy Scales

H. Kowalski, L.N. Lipatov, D.A. Ross

Outline:

Gluon density - analyzed by BFKL equation with running α_s

▶ discrete solution

Gluon density becomes a system of quasi-bound states

(in contrast to the DGLAP evolution)

Application to HERA data, F_2

Future application to LHC data: DY processes

Physics motivation

Sensitivity to BSM effects

Pomeron-Graviton Correspondence

H. Kowalski, Marseille 24th of April 2013

the talk is based on 3 papers

The Green Function for BFKL Pomeron and the Transition to DGLAP Evolution.

H. Kowalski, L.N. Lipatov, D.A. Ross, in preparation

BFKL Evolution as a Communicator Between Small and Large Energy Scales

H. Kowalski, L.N. Lipatov, D.A. Ross, arXiv:1205.6713 and 1109.0432

Using HERA data to determine the infrared behaviour of the BFKL amplitude

H. Kowalski, L.N. Lipatov, D.A. Ross and G. Watt, EPJC 70: 983, 2010

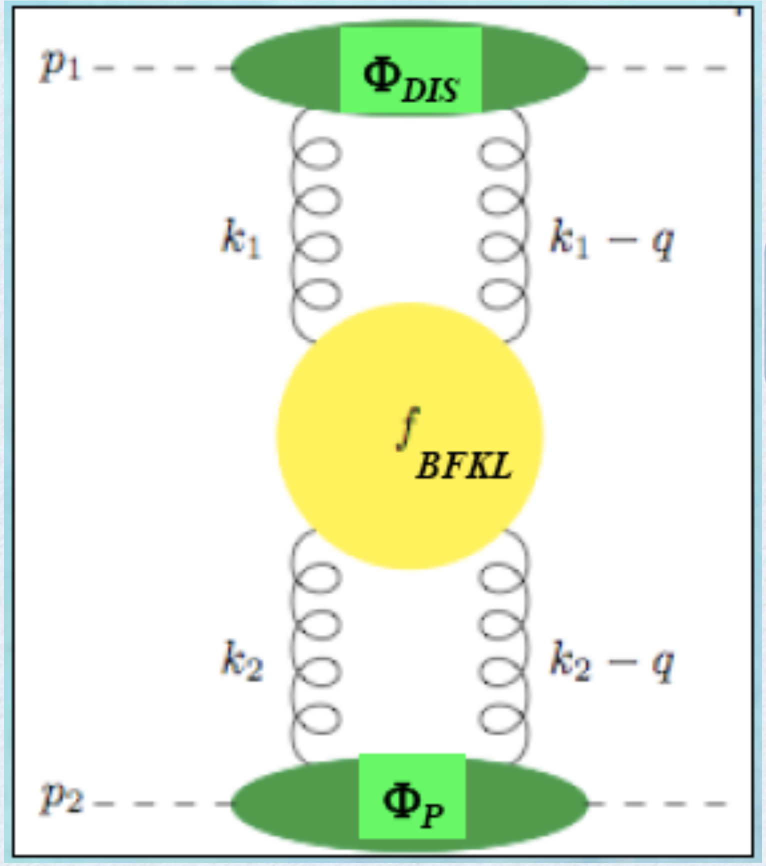
Evidence for the discrete asymptotically-free BFKL Pomeron from HERA data

J. Ellis, H. Kowalski, D.A. Ross

Physics Letters B 668 (2008) 51–56

The dynamics of Gluon Density at low x is determined by the amplitude for the scattering of a gluon on a gluon, described by the BFKL equation

$$\frac{\partial}{\partial \ln s}\mathcal{A}(s,\mathbf{k},\mathbf{k}') = \delta(k^2 - k'^2) + \int dq^2 \mathcal{K}(\mathbf{k},\mathbf{q})\mathcal{A}(s,\mathbf{q},\mathbf{k}')$$



solved by the Green function method, in terms of the eigenfunctions of the kernel

$$\int dk'^2 \mathcal{K}(\mathbf{k},\mathbf{k}') f_\omega(\mathbf{k}') = \omega f_\omega(\mathbf{k})$$

in LO, with fixed α_s

$$f_\omega(\mathrm{k}) = \exp(i\nu \ln k^2)/k$$

$$\omega = \alpha_s \chi_0(\nu)$$

Green f. method - preserves the scaling (conformal) invariance of BFKL

⇒ most consistent solution of BFKL

a possible bridge to Pomeron-Graviton?

Properties of the BFKL Kernel

Quasi-locality

$$\mathcal{K}(\mathbf{k},\mathbf{k}') = \frac{1}{kk'}\sum_{n=0}^{\infty} c_n \delta^{(n)}\left(\ln(\mathbf{k}^2/\mathbf{k}'^2)\right)$$

$$c_n = \int_0^\infty dk'^2 \mathcal{K}(\mathbf{k},\mathbf{k}') \frac{k}{k'} \frac{1}{n!} \left(\ln(\mathbf{k}^2/\mathbf{k}'^2)\right)^n$$

Similarity to the Schroedinger equation

$$k\int dk'^2 \mathcal{K}(\mathbf{k},\mathbf{k}') f_\omega(\mathbf{k}') = \sum_{n=0}^{\infty} c_n \left(\frac{d}{d\ln(\mathbf{k}^2)}\right)^n \bar{f}_\omega(\mathbf{k}) = \omega \bar{f}_\omega(\mathbf{k})$$

Characteristic function

$$k\int dk'^2 \mathcal{K}(\mathbf{k},\mathbf{k}') f_\omega(\mathbf{k}') = \chi\left(-i\frac{d}{d\ln k^2}, \alpha_s(k^2)\right) \bar{f}_\omega(k) = \omega \bar{f}_\omega(k)$$

with running α_s, BFKL frequency ν becomes k-dependent, $\nu(k)$

$$\alpha_s(k^2)\chi_0(\nu(\mathbf{k})) + \alpha_s^2(k^2)\chi_1(\nu(\mathbf{k})) \;=\; \omega \qquad \text{NLO}$$

ν has to become a function of k because ω is a constant

GS resummation applied

evaluation in diffusion ($\nu \approx 0$) or semiclassical approximation ($\nu > 0$)

For sufficiently large k, there is no longer a real solution for ν.
The transition from real to imaginary $\nu(k)$ singles out a special value of

$k = k_{crit}$, with $\nu(k_{crit}) = 0$.

The solutions below and above this critical momentum k_{crit} have to match. This fixes the phase of ef's.

Near $k{=}k_{crit}$, the BFKL eq. becomes the Airy eq. which is solved by the Airy eigenfunctions (to a very good approximation)

$$k\, f_\omega(k) = \bar{f}_\omega(k) = \mathrm{Ai}\left(-(\frac{3}{2}\phi_\omega(k))^{\frac{2}{3}}\right)$$

with

$$\phi_\omega(k) = 2\int_k^{k_{\mathrm{crit}}} \frac{d\,k'}{k'}|\nu_\omega(k')|$$

instead of

$f_\omega(\mathrm{k}) = \exp(i\nu \ln k^2)/k$

for $k{<<}k_{crit}$ the Airy function has the asymptotic behaviour

$$k\, f_\omega(k) \sim \sin\left(\phi_\omega(k) + \frac{\pi}{4}\right)$$

The two fixed phases at $k{=}k_{crit}$ and at $k{=}k_0$ (near Λ_{QCD}) lead to the quantization condition

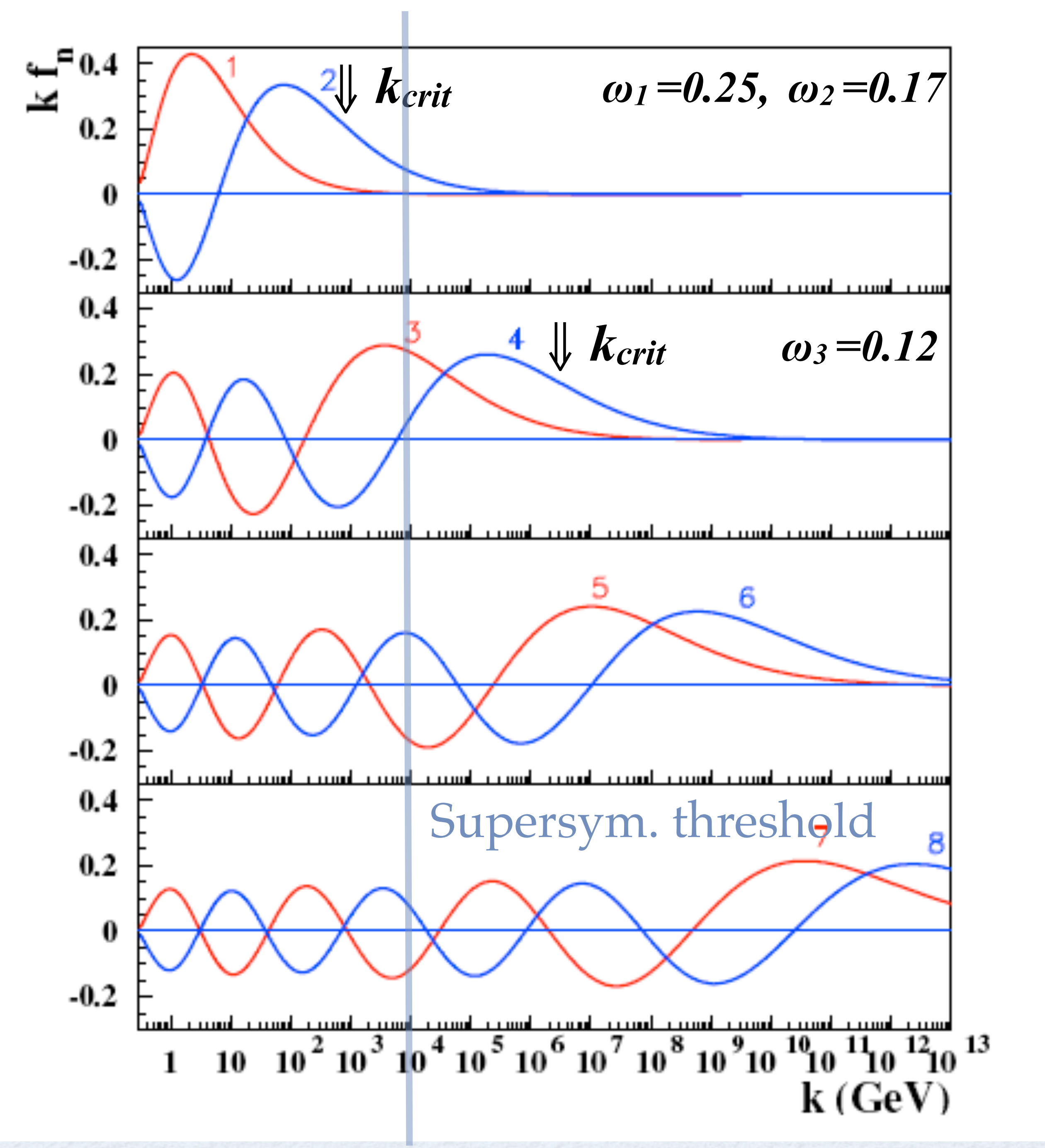
$$\phi_\omega(k_0) = \left(n - \frac{1}{4}\right)\pi + \eta\,\pi$$

Discrete Pomeron Solution of the BFKL eq

The first eight eigenfunctions determined at $\eta=0$

$k_{crit} \simeq c \exp(4n)$
$c \simeq \Lambda_{QCD}$

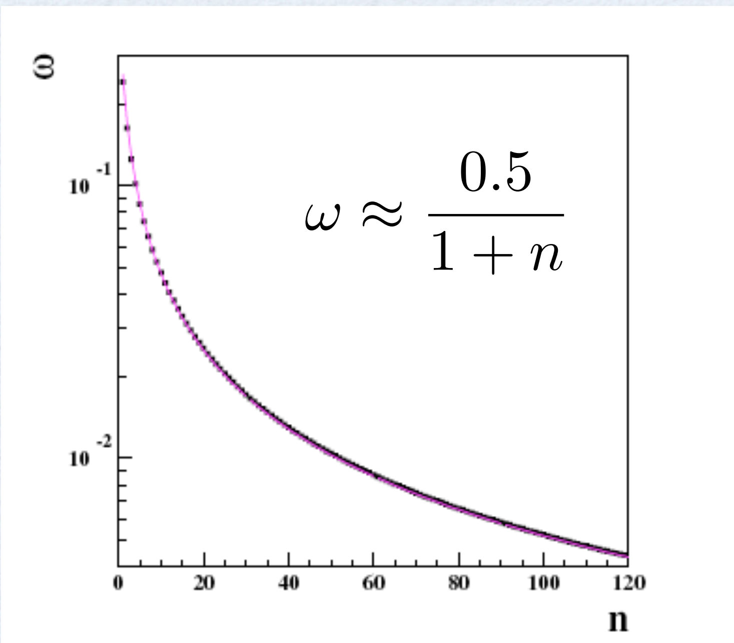
Similarity to WKB solutions of the Schrödinger eq for the potential well



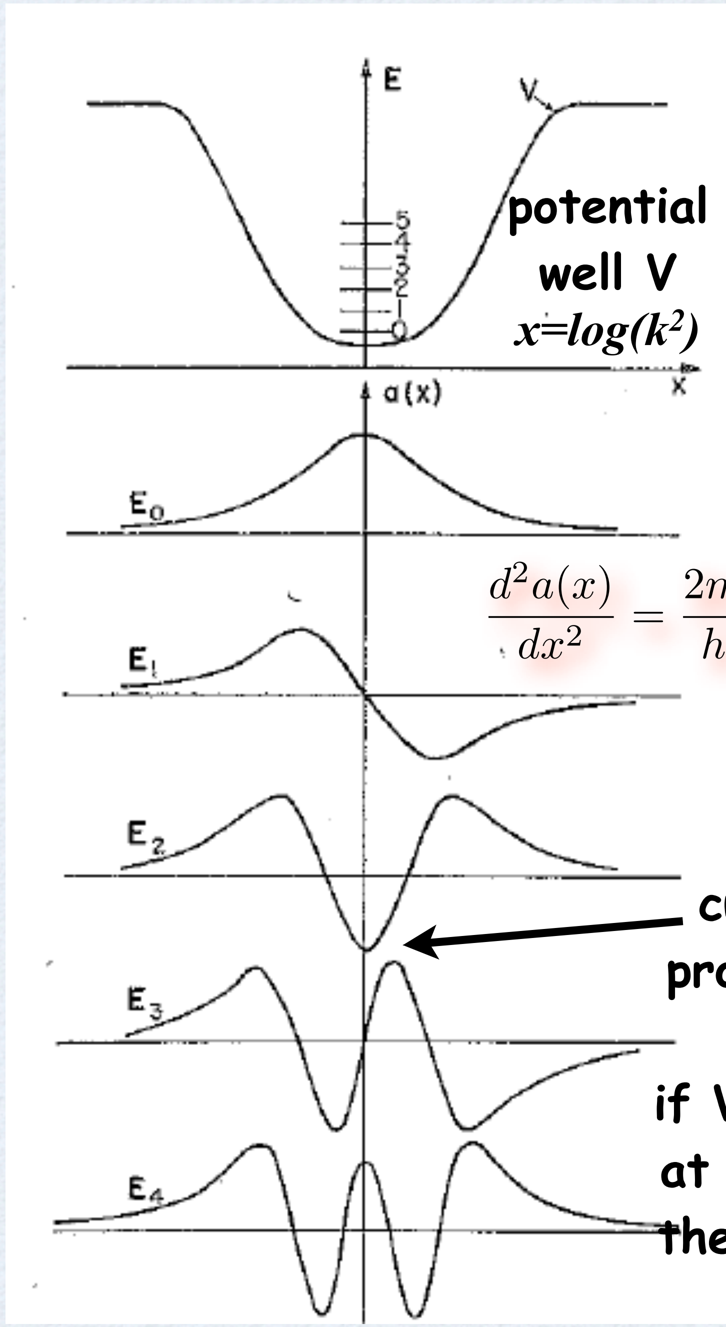
Similarity with the Schroedinder eq. for the potential well

Feynman Lecture III

BFKL eq is similar to S. eq for the potential well with the dynamically increasing width



$$\omega \approx \frac{0.5}{1+n}$$



potential well V

$x=log(k^2)$

$$\frac{d^2a(x)}{dx^2} = \frac{2m}{h}[V - E]a(x)$$

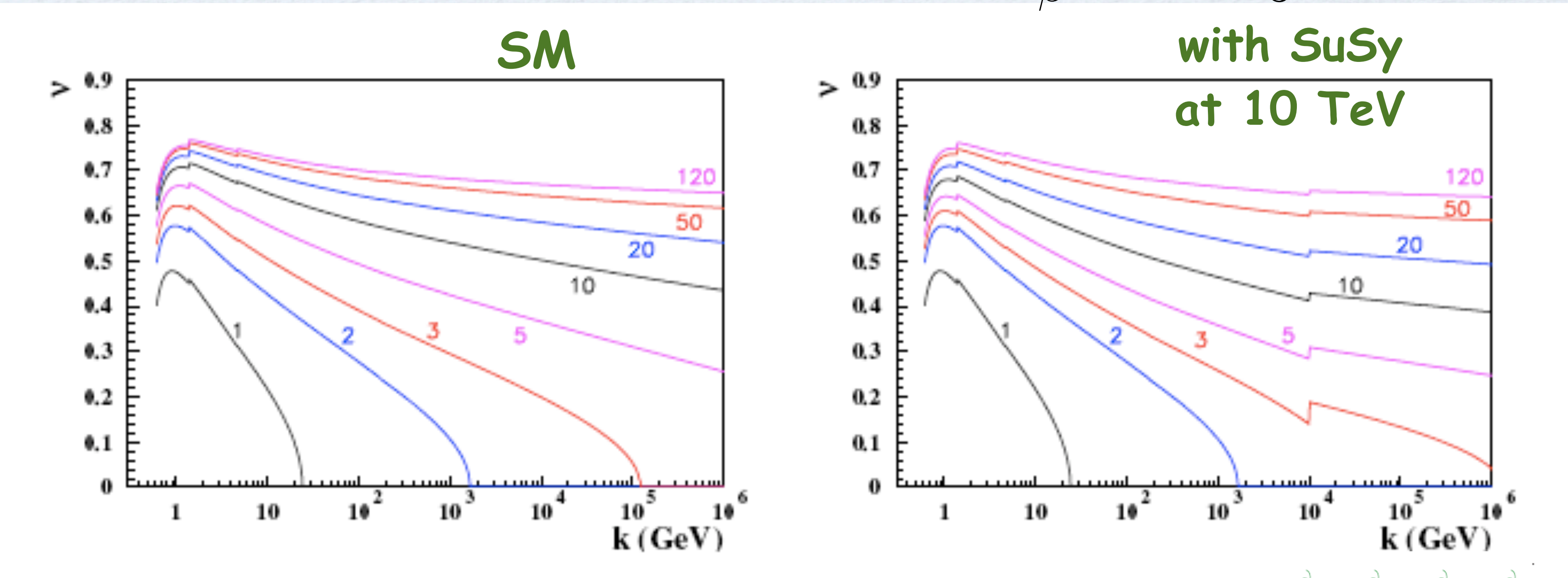
analogy worked out with J.Bartels

curvature is proportional to (V-E)

if V is modified at some x then the whole wf is changed

Sensitivity of the frequencies $\nu(k)$ to thresholds

change of β function in α_s (LO) $$\frac{\beta^{SM}}{\beta^{SUSY}} = \frac{7}{3}$$

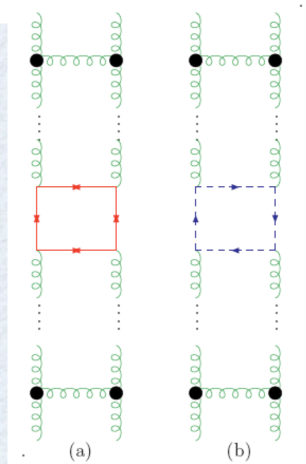


$$\delta_f \chi_1(\nu) = \frac{\pi^2}{32} \frac{\sinh(\pi\nu)}{\nu(1+\nu^2)\cosh^2(\pi\nu)} \left(\frac{11}{4} + 3\nu^2\right)$$

for gluinos

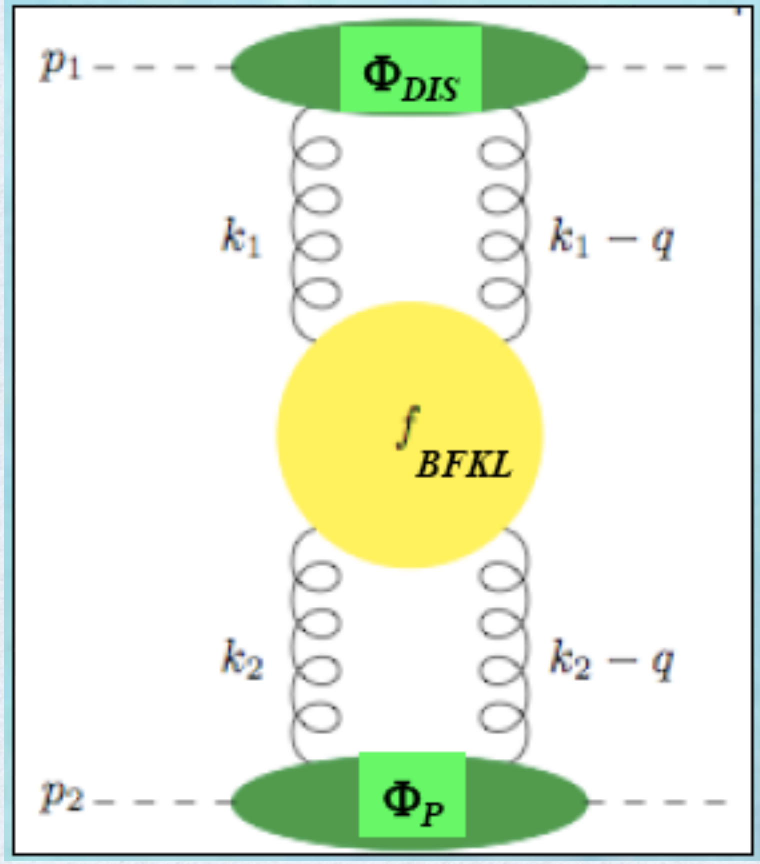
$$\delta_s \chi_1(\nu) = -\frac{\pi^2}{32} \frac{n_f}{C_A^3} \frac{\sinh(\pi\nu)}{\nu(1+\nu^2)\cosh^2(\pi\nu)} \left(\frac{5}{4} + \nu^2\right)$$

for squarks



Kotikov, Lipatov 2003

Comparison with HERA data



Discreet Pomeron Green function

$$\mathcal{A}(\mathbf{k},\mathbf{k}') = \sum_{m,n} f_m(\mathbf{k}) \mathcal{N}^{-1}_{mn} f_n(\mathbf{k}') \left(\frac{s}{kk'}\right)^{\omega_n}.$$

Integrate with the photon and proton impact factors

$$\mathcal{A}^{(U)}_n \equiv \int_x^1 \frac{d\xi}{\xi} \int \frac{dk}{k} \Phi_{\mathrm{DIS}}(Q^2,k,\xi) \left(\frac{\xi k}{x}\right)^{\omega_n} f_n(\mathbf{k})$$

$$\mathcal{A}^{(D)}_m \equiv \int \frac{dk'}{k'} \Phi_p(k') \left(\frac{1}{k'}\right)^{\omega_m} f_m(\mathbf{k}').$$

$$F_2(x,Q^2) = \sum_{m,n} \mathcal{A}^{(U)}_n \mathcal{N}^{-1}_{nm} \mathcal{A}^{(D)}_m$$

the infrared boundary condition

Proton impact factor

$$\Phi_p(\mathbf{k}) \;=\; A\,k^2 e^{-bk^2}$$

The fit is not sensitive to the particular form of the impact factor.

The support of the proton impact factor is much smaller than the oscillation period of f_n and because the frequencies ν have a limited range

➤ many eigenfunctions have to contribute and η has to be a function of n. Phase condition at $\tilde{k}_0$, (close to Λ_{QCD})

$$\eta \;= \eta_0 \left(\frac{n-1}{n_{\max}-1}\right)^{\kappa}$$

additional parameter k_0 which should be in the perturbative region but close to Λ_{QCD}

$$\phi_n(\tilde{k}_0) \;=\; \phi_n(k_0) - 2\nu_n^0 \ln\left(\frac{k_0}{\tilde{k}_0}\right),$$

Fits to F_2, $Q^2 > 8$ GeV^2, $x > 0.01$ $N = 108$, (two loop α_s)

SUSY Scale (TeV)	χ^2	κ	$\tilde{k}_0$ (GeV)	η_0	A	b
3	125.7	0.555	0.288	-0.87	201.2	10.6
6	114.1	0.575	0.279	-0.880	464.8	15.0
10	109.9	0.565	0.275	-0.860	720.1	17.7
15	110.1	0.555	0.279	-0.860	882.2	18.6
30	117.8	0.582	0.278	-0.870	561.6	16.2
50	114.9	0.580	0.279	-0.870	627.4	16.8
90	114.8	0.580	0.279	-0.870	700.2	17.5
∞	122.5	0.600	0.274	-0.800	813.1	17.5

$\chi 2/N=$ 110/108=1.02

Table 1: Fits for N=1 SUSY at different scales. The bottom row corresponds to the Standard Model. All fits are performed with $n_{max} = 100$.

Note: we are partially absorbing the SUSY effects into the free parameters of the boundary conditions: e.g best SuSy fit with η_0, κ of SM gives $\chi 2$ ~ 400

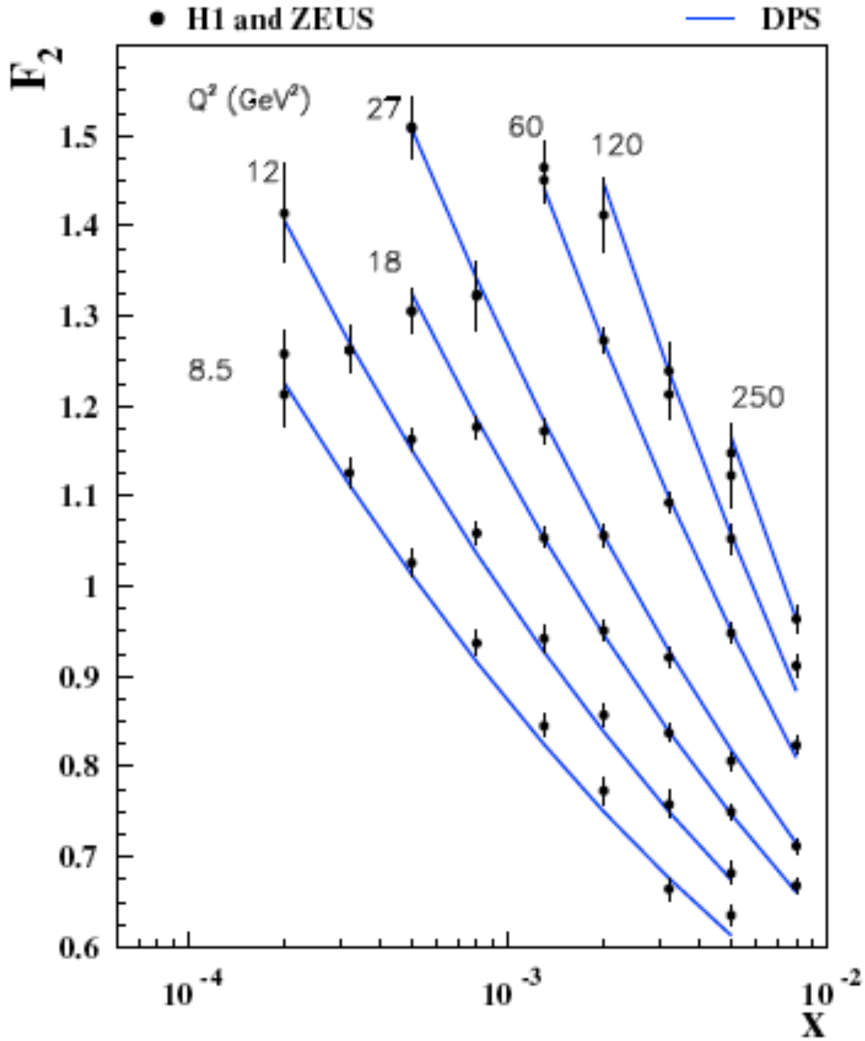
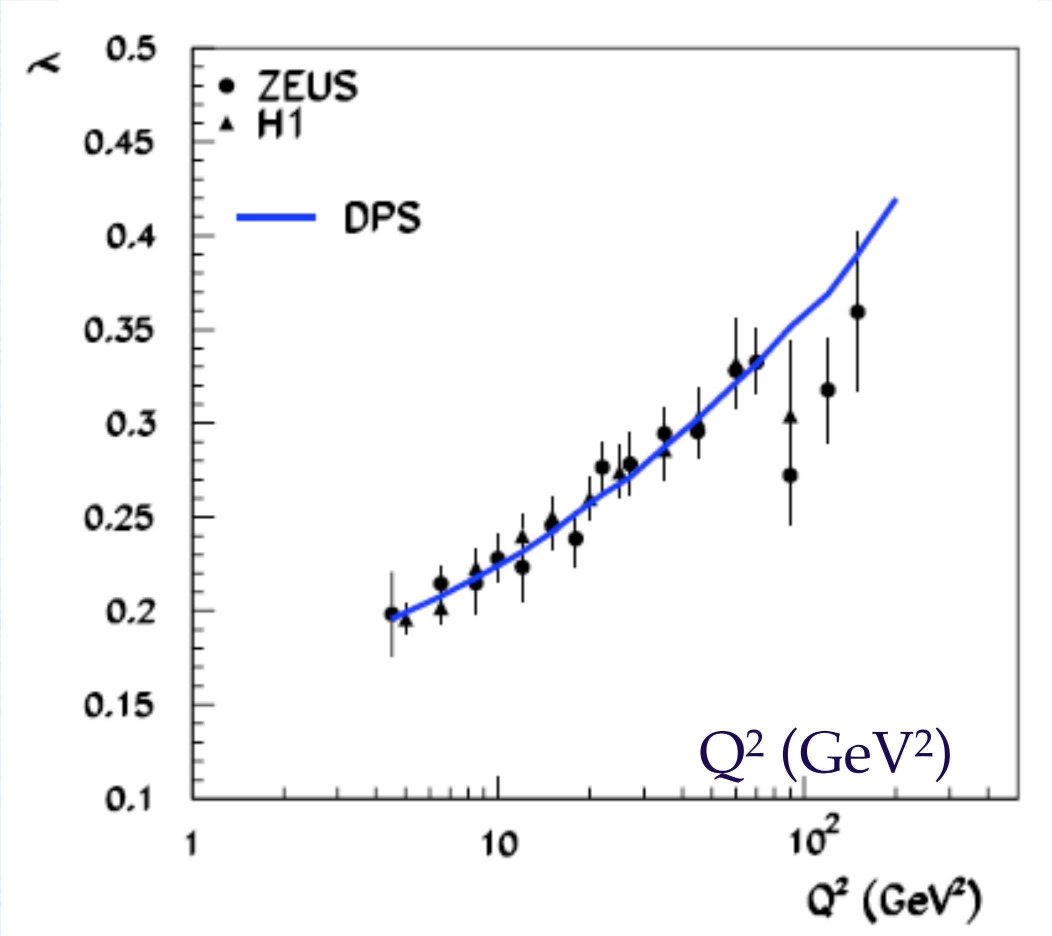


Figure 7: Comparison of the DPS fit with $M_{SUSY} = 10$ TeV with HERA data.

The rate of rise λ

$F_2 \sim (1/x)^\lambda$



The first successful pure BFKL description of the λ plot.

For many years it was claimed that BFKL analysis was not applicable to HERA data because of the observed substantial variation of λ with Q^2

The qualities of fits for various numbers of eigenfunctions, $Q^2 > 4$ GeV^2 (one loop α_s)

$n_{\max}$	χ^2/N_{df}	κ	A	b
1	10811 /125	—	146	30.0
5	350.0 /125	3.78	$3.1 \cdot 10^6$	78.0
20	286.5 /125	0.96	632	15.8
40	193.3 /125	0.84	2315	23.2
60	163.3 /125	0.78	3647	25.6
80	156.5 /125	0.73	3081	24.4
100	149.1 /125	0.69	2414	22.8
120	143.7 /125	0.66	2041	21.8

➤ new data are crucial for finding the right solution

the differences in the fit qualities would be negligible if the errors where more than 2-times larger

Discrete BFKL-Pomeron

Why so many eigenfunctions?

the contribution of large n ef's is only weakly suppressed,
enhancement by $(1/x)^{\omega}$ is not very large because
$\omega_1 \approx 0.25$, $\omega_5 \approx 0.1$, $\omega_{10} \approx 0.05$

suppression of large n contribution only by

the normalization condition for eigenfunctions $\sim 1/\sqrt{n}$

alternating signs of the proton overlap $(-1)^n$

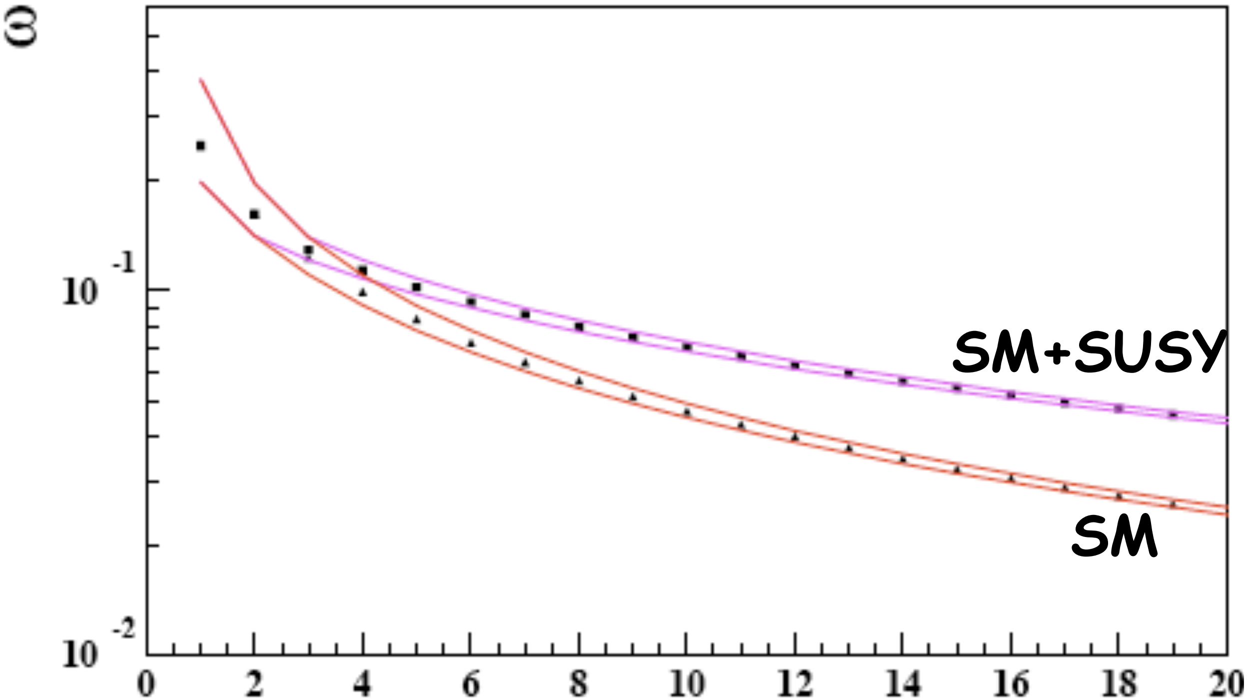
Eigenvalues of the Discrete BFKL-Pomeron

LO evaluation

$$\omega_n = \frac{0.96}{\pi\beta} \cdot \frac{1}{\eta + \mathbf{n} - 1/4}$$

$$\frac{\beta^{SM}}{\beta^{SUSY}} = \frac{7}{3}$$

NLO numerical evaluation

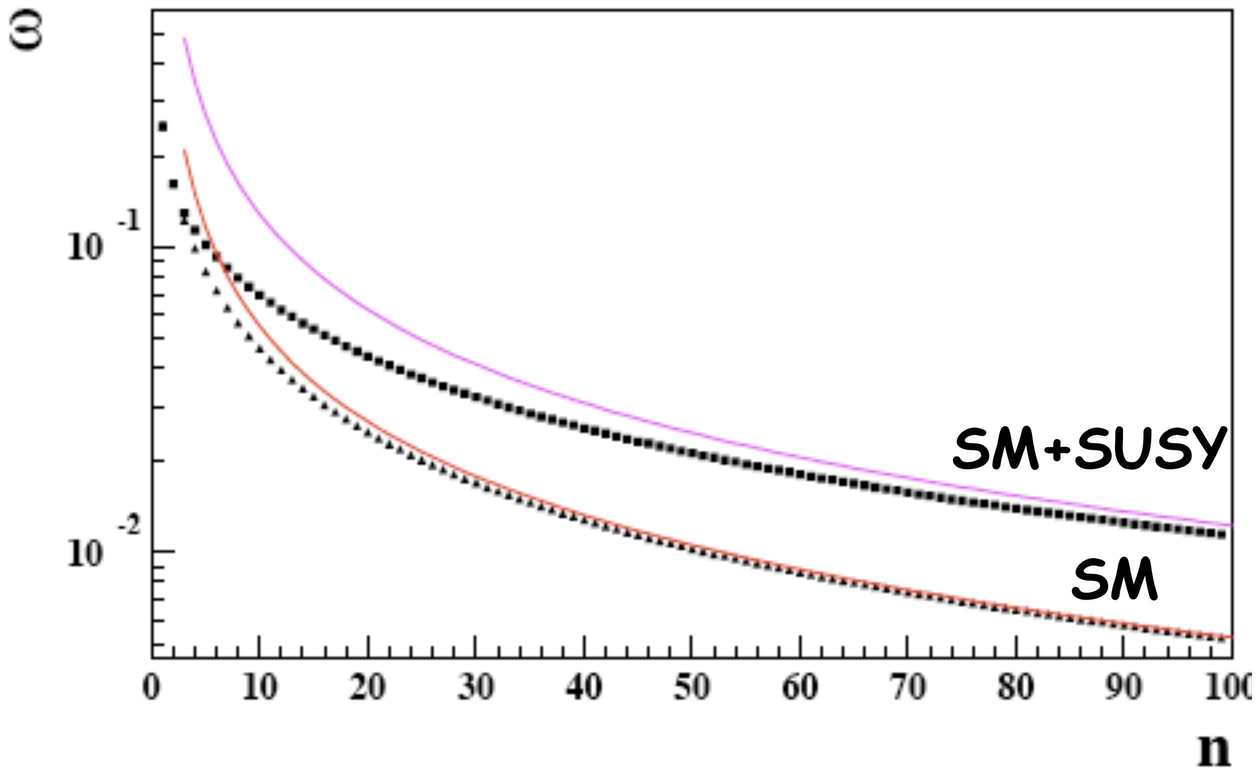


lines indicate the uncertainty of the phase
(η can only vary between 0 and π)

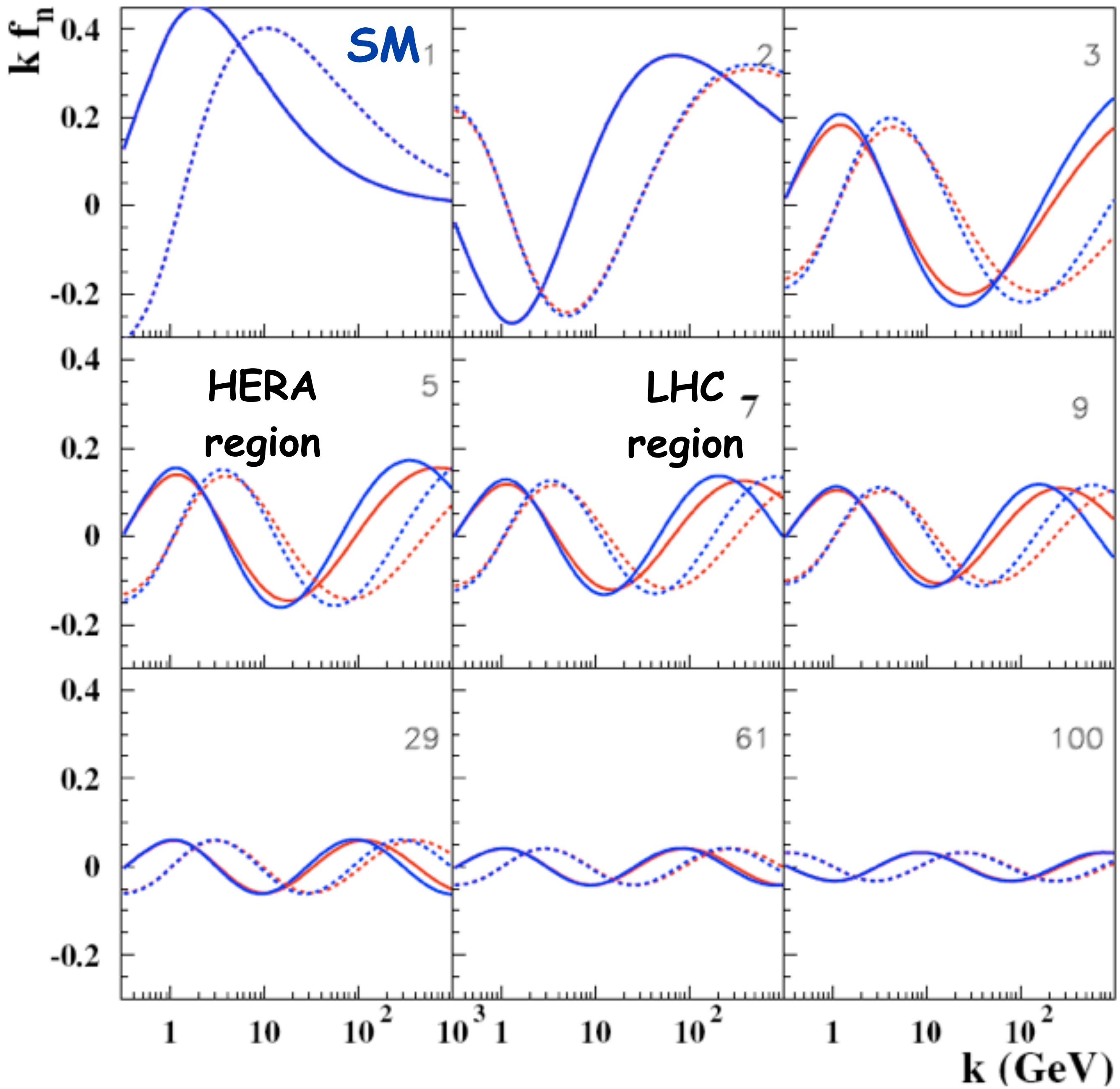
difference between SM and SM+SUSY is substantially larger than the uncertainty of the phase

Eigenvalues of the Discrete BFKL-Pomeron

Comparison of the LO analytical (lines) and the NLO numerical evaluation (symbols)



difference between SM and SM+SUSY is substantially larger than the difference between LO and NLO

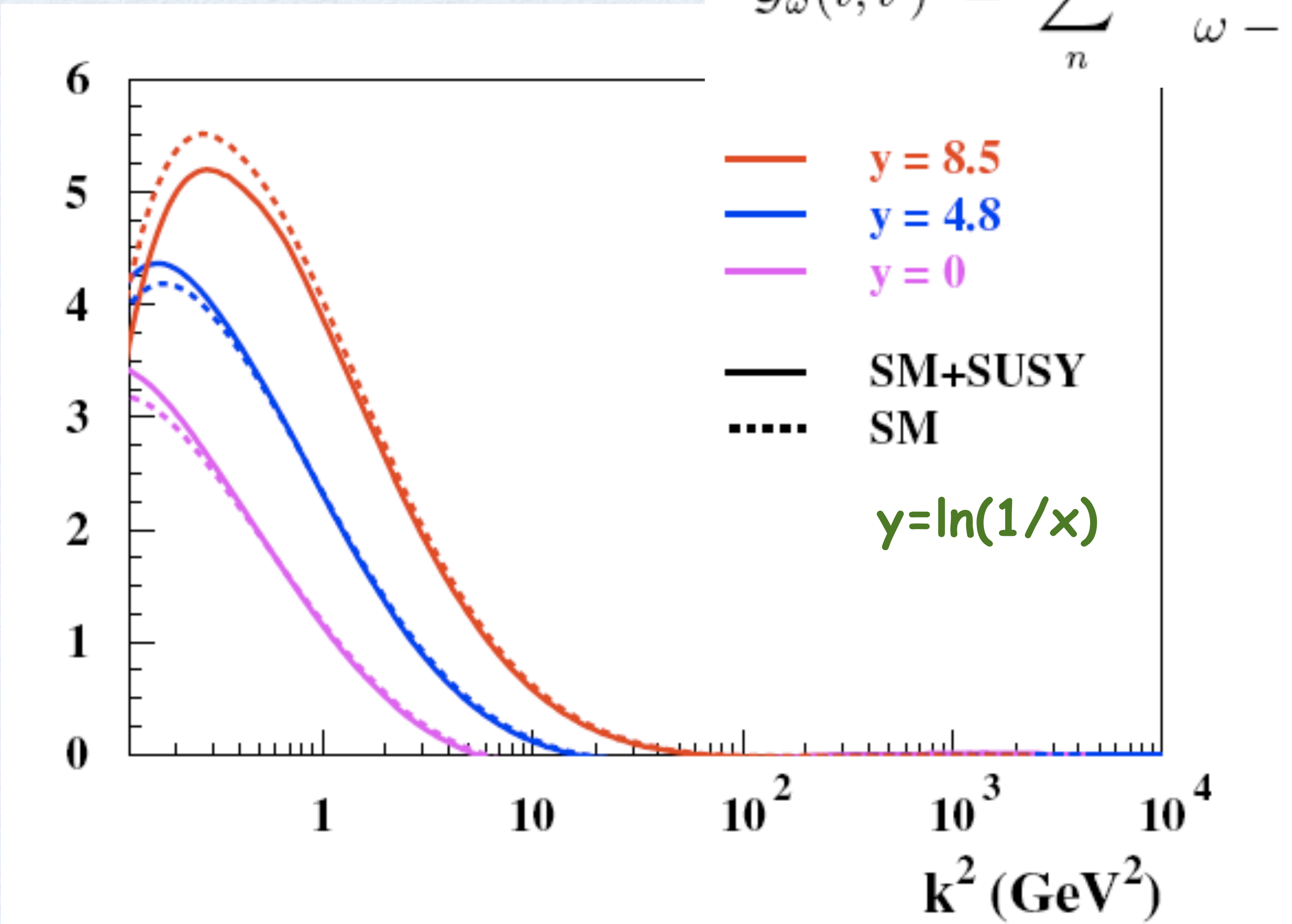


k f_n
SM
1
2
3
HERA
region
5
LHC
region
7
9
29
61
100
0.4
0.2
0
-0.2
1
10
10^2
10^3
k (GeV)

Evolution of the wave packet in DPS

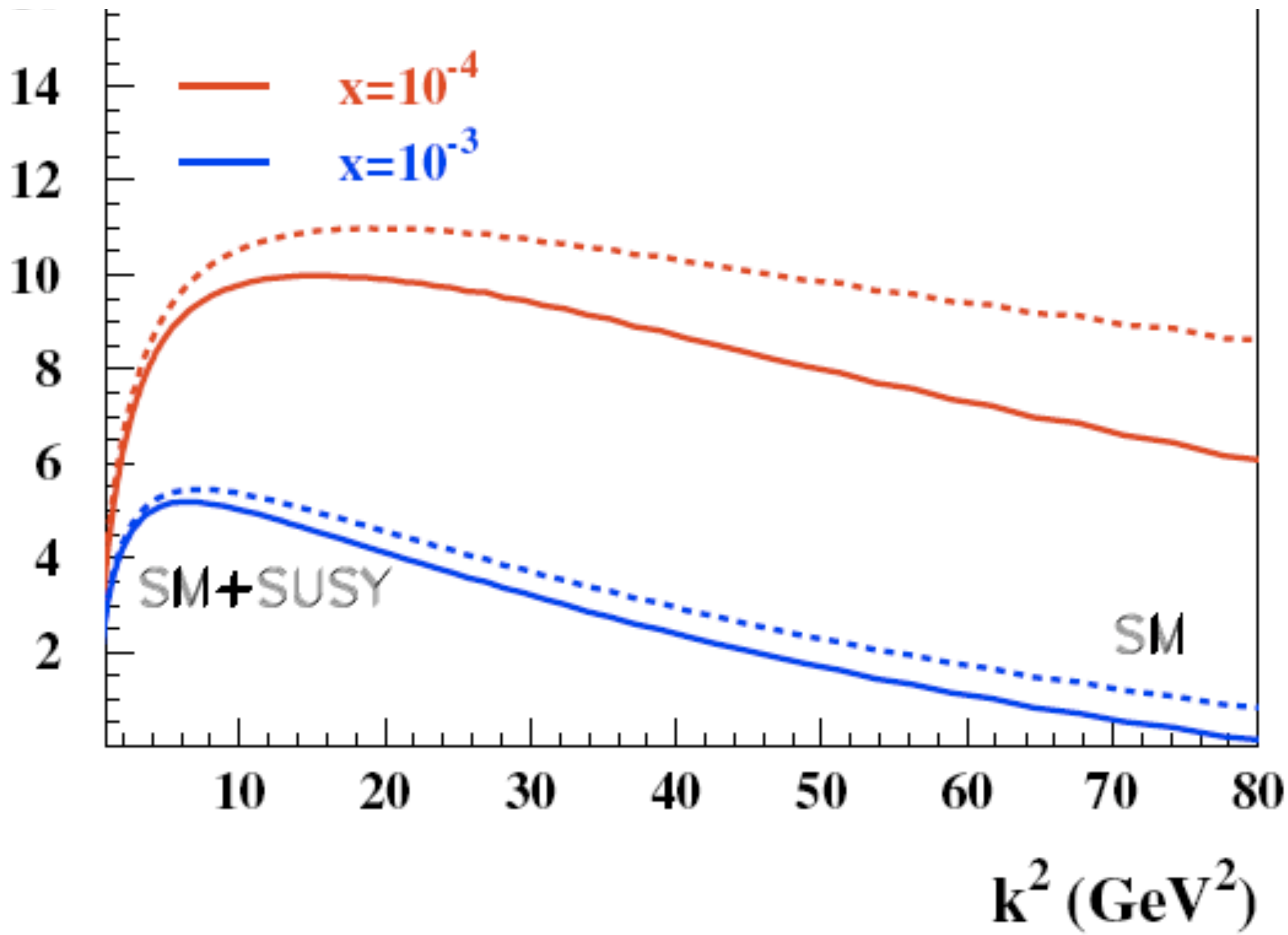
$$\mathcal{WP}(y,k^2) \;=\; \int d\omega \int \frac{dk'^{\,2}}{k'^{\,2}}\, e^{\omega y}\, \hat{\mathcal{G}}_\omega(k,k')\Phi_p(k'),$$

$$\tilde{\mathcal{G}}_\omega(t,t') \;=\; \sum_n \frac{f^*_{\omega_n}(t')f_{\omega_n}(t)}{\omega-\omega_n}\;.$$



Evolution of the gluon density in DPS

$$x\dot{g}(x,k^2) = k^2 \int d\omega \int \frac{dk'^2}{k'^2} \left(\frac{kx}{k'}\right)^{-\omega} \hat{\mathcal{G}}_\omega(k,k')\Phi_p(k'),$$



Next steps necessary for description of the low-x and high Q^2 processes (DY at LHC)

perform a full evaluation of the Green function:

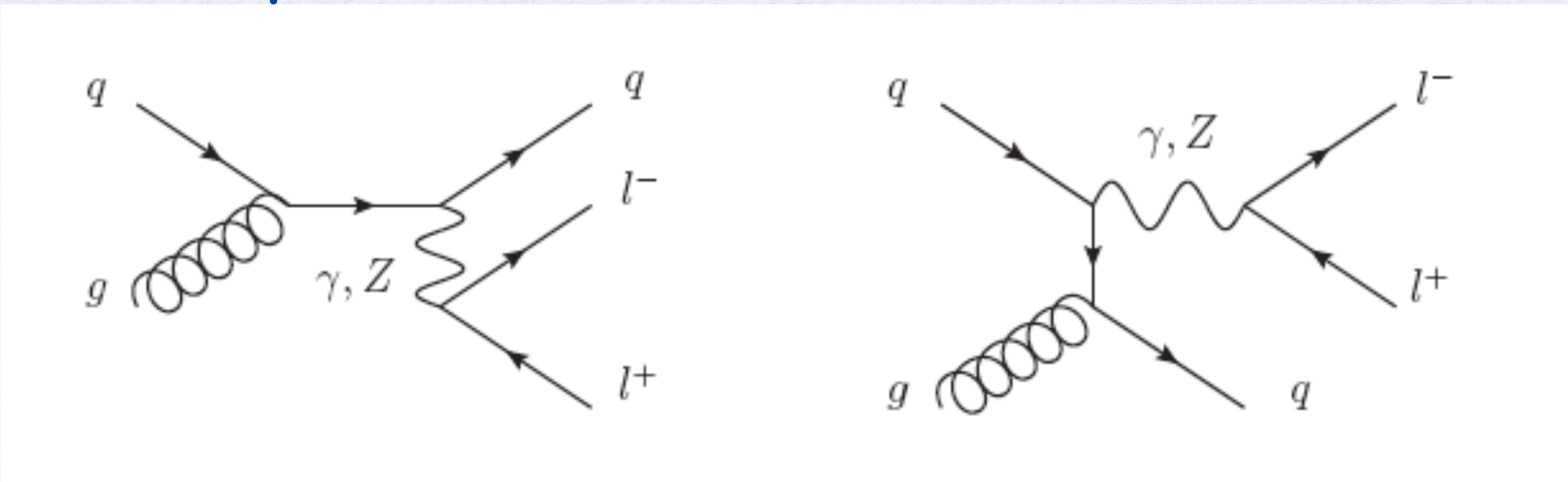
- sum over n to infinity (instead to n=O(100))
- evaluate a possible contribution of negative ω's

$$\tilde{\mathcal{G}}_\omega(t,t') = \sum_n \frac{f^*_{\omega_n}(t')f_{\omega_n}(t)}{\omega-\omega_n} + \frac{1}{2\pi i}\int_{-\infty}^{0} d\omega' \frac{f^*_{\omega'}(t')f_{\omega'}(t)}{\omega-\omega'+i\epsilon}$$

the large n and negative ω's contributions can be evaluated in LO only (presumably)

Drell-Yan processes at LHC

Dominant process at LHC



Additional requirement: add valence quarks contribution, i.e; gluon and sea-quark contribution like in DPS and valence quarks like in DGLAP

necessary requirement: obtain DGLAP from DP-BFKL paper in progress

The Green Function for the Discrete BFKL Pomeron and the Transition to DGLAP Evolution.

H. Kowalski, L.N. Lipatov, D.A. Ross + ...

Obtain the BFKL Green Function

$$\left(\omega - \hat{\Omega}(\omega, t, \hat{\nu})\right) \mathcal{G}_\omega(t, t') \;=\; \delta(t - t'),$$

$$t \equiv \ln\left(k^2/\Lambda^2_{QCD}\right)$$

from the generalized Airy operator
(valid in diffusion and semiclassical approximation)

$$\left(\omega - \hat{\Omega}\left(\omega, t, -i\frac{\partial}{\partial t}\right)\right) \;=\; \frac{1}{N_\omega(t)}\left(\dot{z}z - \frac{\partial}{\partial t}\frac{1}{\dot{z}}\frac{\partial}{\partial t}\right)\frac{1}{N_\omega(t)},$$

$$s_\omega(t) \;=\; \int_t^{t_c} dt'\, \nu_\omega(t').$$

$$z(t) = -\left(\frac{3}{2}s_\omega(t)\right)^{\frac{2}{3}}$$

Generalized Airy Green Function

$$\mathcal{G}_\omega(t,t') = \pi N_\omega(t) N_\omega(t') \left(\overline{B_i}(z(t)) A_i(z(t'))\theta(t'-t) + A_i(z(t))\overline{B_i}(z(t'))\theta(t-t')\right)$$

with $\overline{B_i}(z) = B_i(z) + c(\omega) A_i(z).$ $c(\omega) = \cot\left(\phi(\omega)\right)$

$$\phi(\omega) = \eta_{np}(\omega, t_0) - \frac{\pi}{4} - s_\omega(t_0).$$

leads to a similar pole term contribution, $\Delta y = \ln(1/x)$

$$\mathcal{G}_\omega^{\mathrm{pole}}(t,t') = \sum_n \pi N_{\omega_n}(t) N_{\omega_n}(t') \frac{A_i\left(z(t)\right) A_i\left(z(t')\right)}{\phi'(\omega_n)(\omega - \omega_n)},$$

+ a possible contribution of the cut at negative ω

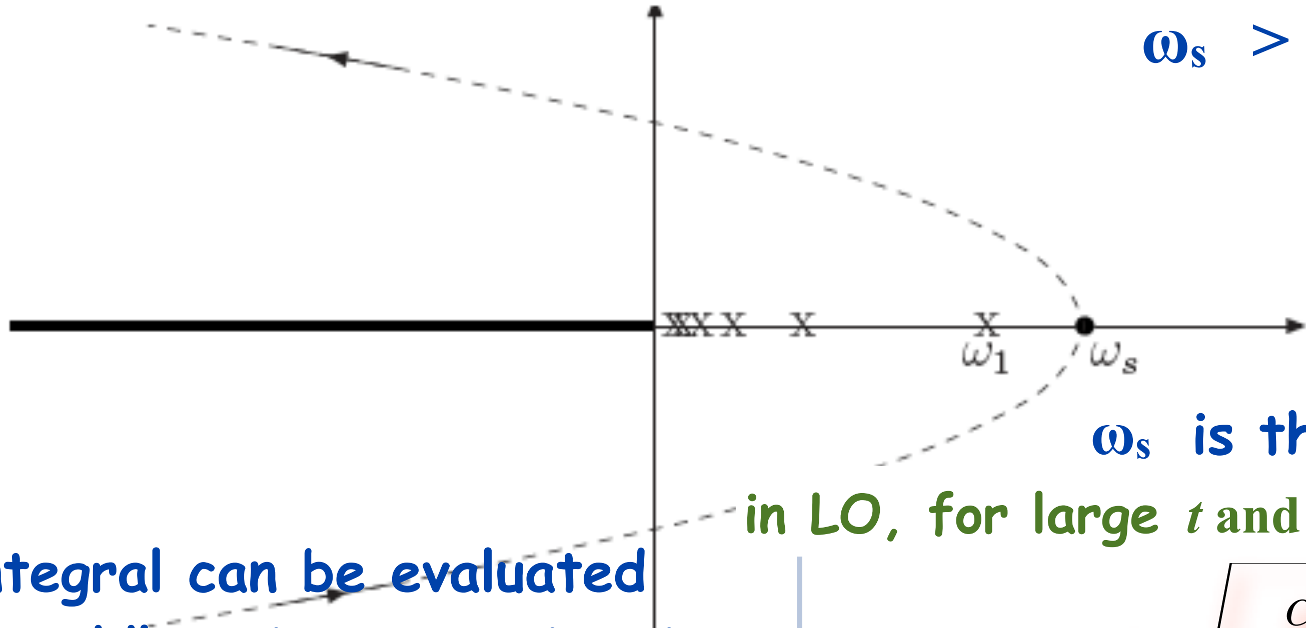
$$N_\omega(t) = \frac{(-z(t))^{1/4}}{\sqrt{\frac{1}{2}\Omega'(\omega, t, \nu_\omega)}},$$

Unintegrated gluon density

$$\dot{g}(x,t) = \frac{1}{2\pi i}\int_C d\omega x^{-\omega} \int dt' \mathcal{G}_\omega(t,t')\Phi_P(t'),$$

Integration over contour C

$\omega_s > \omega_1$



ω_s is the saddle point

in LO, for large t and $\ln(1/x)$

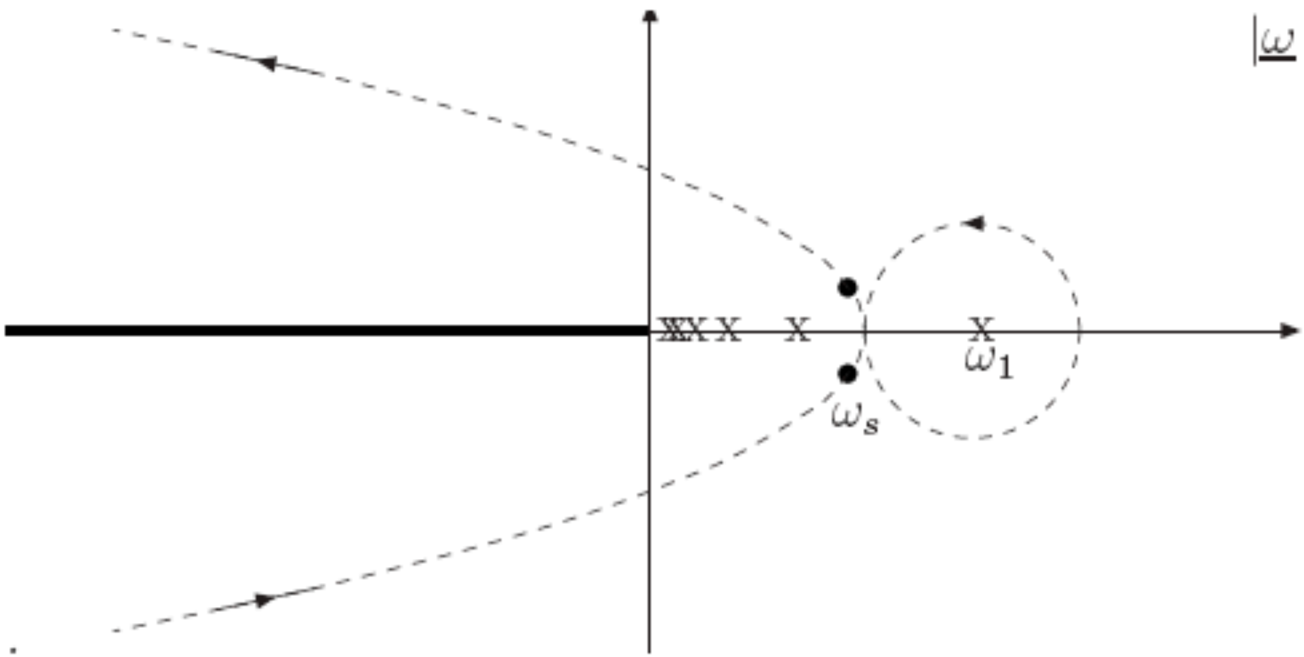
$$\omega_s \approx \sqrt{\frac{\alpha_s\, t}{\ln(1/x)}}$$

Integral can be evaluated in saddle point approximation

it agrees with the DLL limit of DGLAP

for $t \ll \ln(1/x)$
$\mathrm{Re}(\omega_s) < \omega_m$

$\mathrm{Re}(\omega_s) < \omega_1$



Conclusions

The Discrete-Pomeron solution of BFKL provides a very good description of HERA data

The DP-BFKL has a genuine sensitivity to BSM effects. The BSM effects are affecting the eigenvalues and eigenfunctions at larger n which are almost independent on the higher order QCD effects and the lack of knowledge of the Infrared Boundary Condition (IBC).

The data evaluation depends on IBC. IBC is a physical quantity, its understanding can be substantially improved:

by analyzing different physics reactions,

e.g.: F_2 together with LHC Drell-Yan

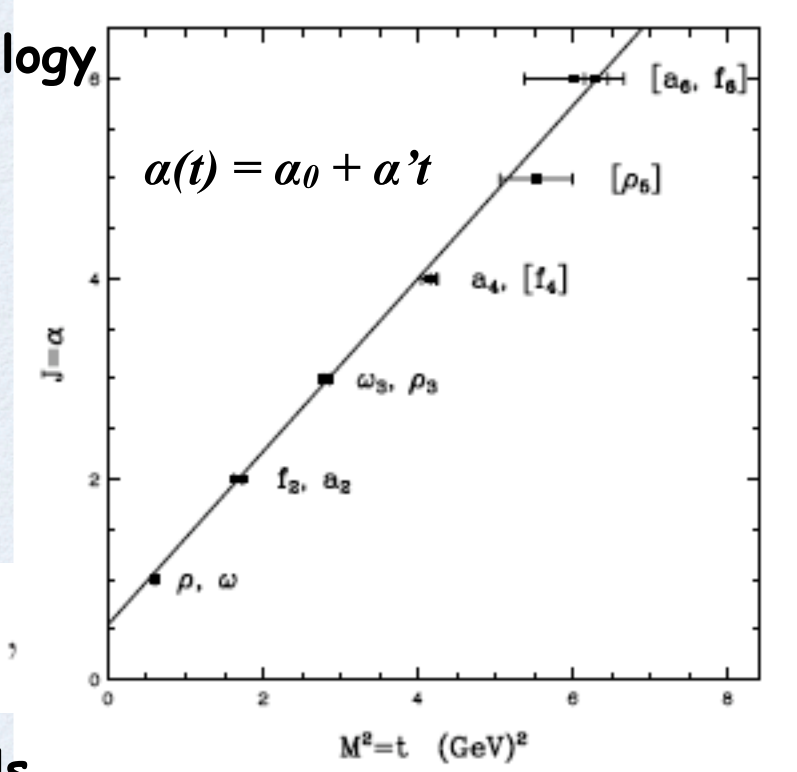
+ diffractive processes...

by involving more sophisticated theoretical methods

Back up slides

Pomeron - Graviton Correspondence

String theory emerged out of phenomenology of hadron-hadron scattering - Dolan-Horn-Schmid duality



$$\sum_r \frac{g_r^2(t)}{s-(M_r-i\Gamma_r)^2} \simeq \beta(t)(-\alpha' s)^{\alpha(t)}$$

▶ Veneziano amplitude

$$A_{\pi^+\pi^-\to\pi^+\pi^-}(s,t) = g_o^2\,\frac{\Gamma[1-\alpha_\rho(t)]\Gamma[1-\alpha_\rho(s)]}{\Gamma[1-\alpha_\rho(s)-\alpha_\rho(t)]},$$

▶ generalization to dual resonance models, Veneziano amplitude for the pomeron trajectory has a pole for $s=t=0$ *with* $J=2$

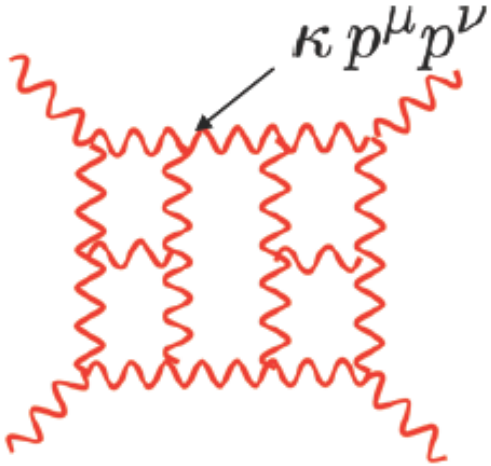
▶ starting point for a theory of quantum gravity

Maldacena Conjecture: (N=4 SUSY QCD) = (CFT in $ADS_5 \times S^5$)

from a talk by ZVI BERN

Is a UV finite theory of gravity possible?

$\kappa = \sqrt{32\pi G_N}$ ← Dimensionful coupling



Gravity: $\displaystyle \int \prod_{i=1}^{L} \frac{dp_i^D}{(2\pi)^D} \frac{(\kappa\, p_j^\mu p_j^\nu)\cdots}{\text{propagators}}$

Gauge theory: $\displaystyle \int \prod_{i=1}^{L} \frac{d^D p_i}{(2\pi)^D} \frac{(g\, p_j^\nu)\cdots}{\text{propagators}}$

Extra powers of loop momenta in numerator means integrals are badly behaved in the UV

Focus on N=8 supergravity and N=4 SUSY YM

High degree of symmetry => technical simplicity

new methods developed:

Modern Unitarity, symbology, BDS ...

focus on order by order finiteness - now up to 6 loops

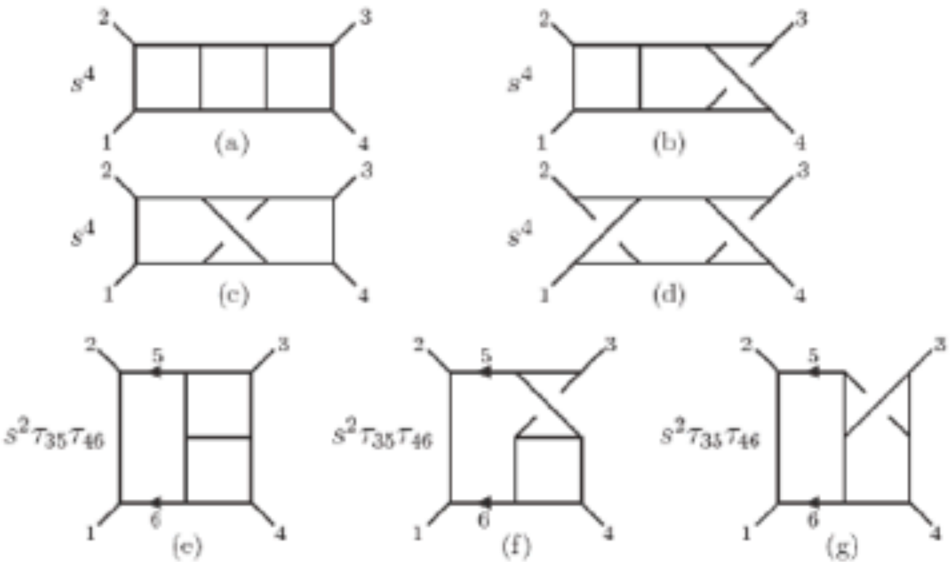
Infinite loop calculation could be possible in the Multi-Regge limit

Complete Three Loop Result

ZB, Carrasco, Dixon, Johansson, Kosower, Roiban; hep-th/0702112
ZB, Carrasco, Dixon, Johansson, Roiban arXiv:0808.4112 [hep-th]

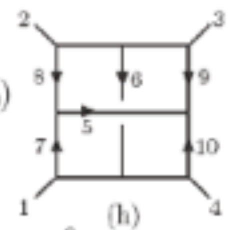
Obtained via maximal cut method:

from a talk by ZVI BERN

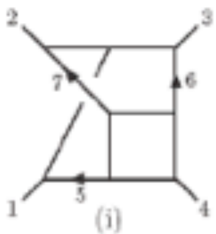


$$\tau_{ij} = 2k_i \cdot k_j$$

$$(s(\tau_{26}+\tau_{36})+t(\tau_{15}+\tau_{25})+st)^2$$
$$+(s^2(\tau_{26}+\tau_{36})-t^2(\tau_{15}+\tau_{25}))(\tau_{17}+\tau_{28}+\tau_{39}+\tau_{4,10})$$
$$+s^2(\tau_{17}\tau_{28}+\tau_{39}\tau_{4,10})+t^2(\tau_{28}\tau_{39}+\tau_{17}\tau_{4,10})$$
$$+u^2(\tau_{17}\tau_{39}+\tau_{28}\tau_{4,10})$$



$$(s\tau_{45}-t\tau_{46})^2-\tau_{27}(s^2\tau_{45}+t^2\tau_{46})$$
$$-\tau_{15}(s^2\tau_{47}+u^2\tau_{46})-\tau_{36}(t^2\tau_{47}+u^2\tau_{45})$$
$$+l_5^2s^2t+l_6^2st^2-\tfrac{1}{3}l_7^2stu$$



Three-loop is not only ultraviolet finite it is "superfinite"—cancellations beyond those needed for finiteness!

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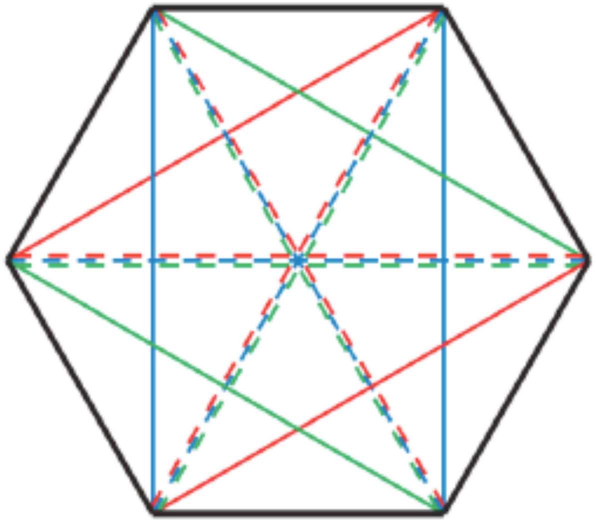
Scattering in Planar N=4 Super-Yang-Mills Theory and the Multi-Regge-Limit



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ICHEP Melbourne, Australia July 5, 2012

from the Summary:

- Multi-Regge limit of 6-gluon amplitude may well be first case solved to all orders