

# *Single bottom production and the heavy quark impact factor at NLO*

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# *Why single bottom production?*

## **Theoretical motivation:**

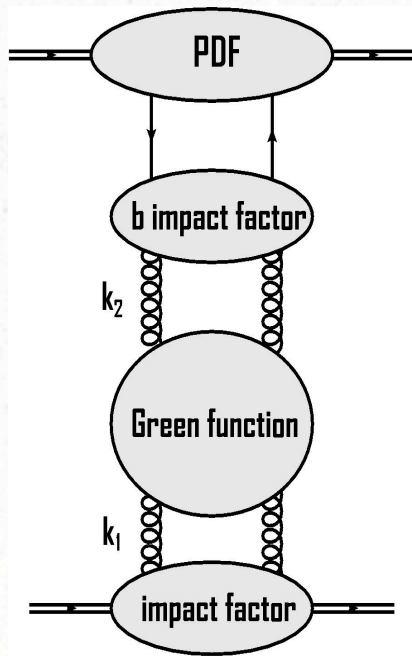
- Interesting for small-x physics
- Open questions of phenomenology

## **Experimental motivation:**

- A tagged b-quark
- Large rapidity coverage

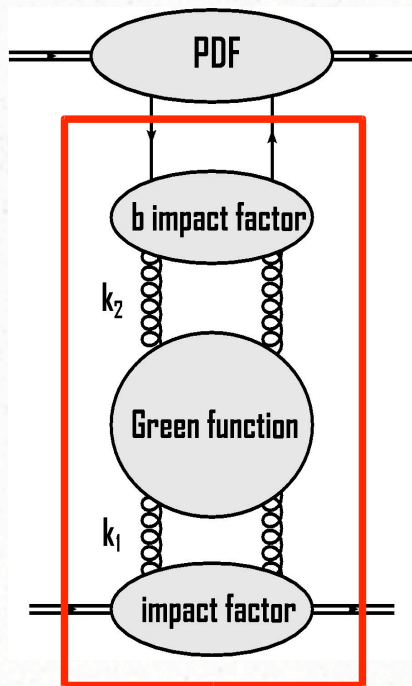
# *Framework*

- Diagrams and cross sections



# Framework

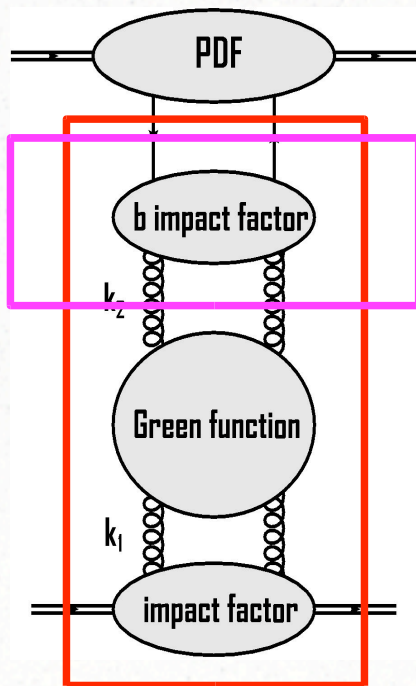
- Diagrams and cross sections



$$\frac{d\sigma_{ab}}{d[\mathbf{k}_1] d[\mathbf{k}_2]} = \int \frac{d\omega}{2\pi i\omega} h_a(\mathbf{k}_1) \mathcal{G}_\omega(\mathbf{k}_1, \mathbf{k}_2) \times \\ \times h_b(\mathbf{k}_2) \left( \frac{s}{s_0(\mathbf{k}_1, \mathbf{k}_2)} \right)^\omega$$

# Framework

- Diagrams and cross sections



$$\frac{d\sigma_{ab}}{d[\mathbf{k}_1] d[\mathbf{k}_2]} = \int \frac{d\omega}{2\pi i\omega} h_a(\mathbf{k}_1) \mathcal{G}_\omega(\mathbf{k}_1, \mathbf{k}_2) \times \\ \times h_b(\mathbf{k}_2) \left( \frac{s}{s_0(\mathbf{k}_1, \mathbf{k}_2)} \right)^\omega$$

*Heavy quark impact factor*

# Momentum space

- Single heavy quark production matrix element at NLO

$$F_q(z_1, \mathbf{k}_1, \mathbf{k}_2) = F_q^{m=0}(z_1, \mathbf{k}_1, \mathbf{k}_2) + \Delta F_q(z_1, \mathbf{k}_1, \mathbf{k}_2)$$

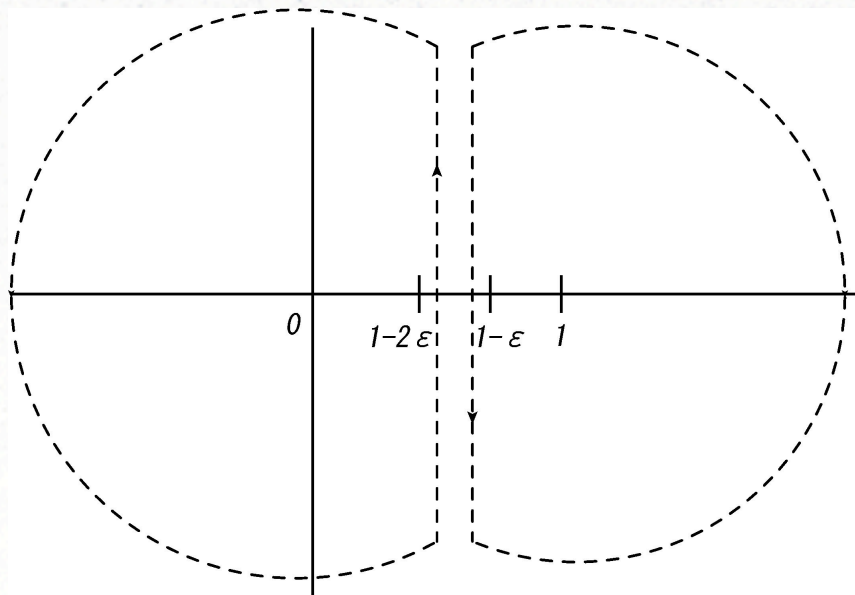
$$\begin{aligned} \Delta F_q(\mathbf{k}_2) &= \Delta F_{q,\text{real}}(\mathbf{k}_2) + \Delta F_{q,\text{virt}}(\mathbf{k}_2) \\ &= A_\varepsilon \left[ \frac{\Gamma(-\varepsilon)}{2(1+2\varepsilon)} \frac{(m^2)^\varepsilon}{\mathbf{k}_2^2} + \frac{\Gamma(1-\varepsilon)}{2} \left\{ \int_0^1 \int_0^1 dz_1 dx \left( \frac{1-z_1}{z_1} + \frac{1+\varepsilon}{2} z_1 \right) \times \right. \right. \\ &\quad \times \left[ \frac{1}{[x(1-x)\mathbf{k}_2^2 + m^2 z_1^2]^{1-\varepsilon}} - \frac{1}{[x(1-x)\mathbf{k}_2^2]^{1-\varepsilon}} \right] + \\ &\quad \left. \left. + \frac{2m^2}{\mathbf{k}_2^2} \int_0^1 \int_0^1 \frac{z_1(1-z_1) dz_1 dx}{[x(1-x)\mathbf{k}_2^2 + m^2 z_1^2]^{1-\varepsilon}} \right\} \right] \end{aligned}$$

- Integrals over  $x$  and  $z$  - finite

## *Some technical aspects*

- Sum of residues – picture

$$\Delta F_q(\mathbf{k}_2) = \frac{1}{m^2} \int_{1-2\varepsilon < \text{Re}\gamma < 1-\varepsilon} \frac{d\gamma}{2\pi i} \left( \frac{\mathbf{k}_2^2}{m^2} \right)^{-\gamma-\varepsilon} \Delta \tilde{F}_q(\gamma, \varepsilon)$$

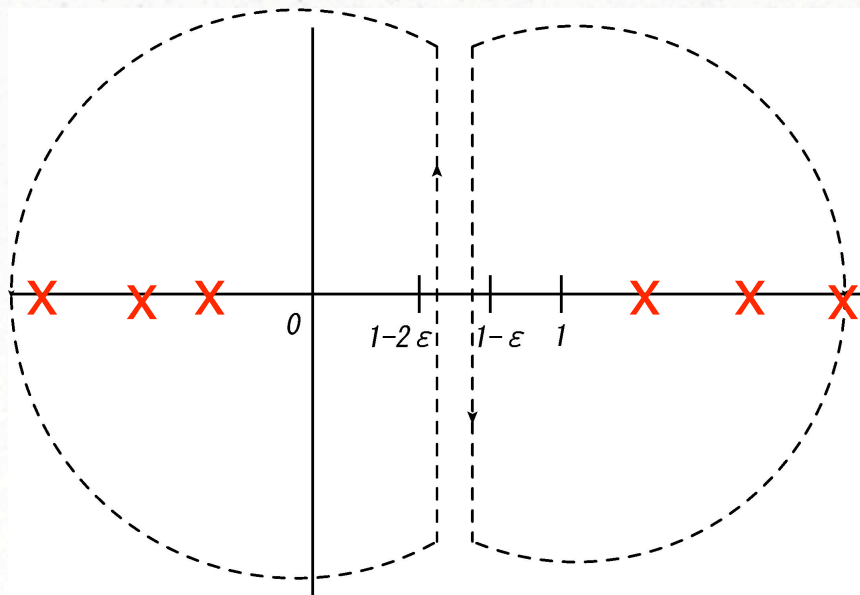




# *Some technical aspects*

- Sum of residues – picture

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# *Some technical aspects*

- Sum of residua
- $R = \mathbf{k}_2^2 / m^2$

- point:  $1 - 2\epsilon$

$$\text{contrib1} = -\frac{1}{\epsilon^2} + \frac{1}{2\epsilon} + \frac{1}{2} (3 + \log(R) + \log(R)^2)$$

- point:  $\frac{1}{2} - 2\epsilon$

$$\text{contrib2} = -\frac{3}{8}\pi^2\sqrt{R}$$

- point:  $-\epsilon$

$$\text{contrib3} = \frac{R}{4\epsilon p} + \frac{7R}{12}$$

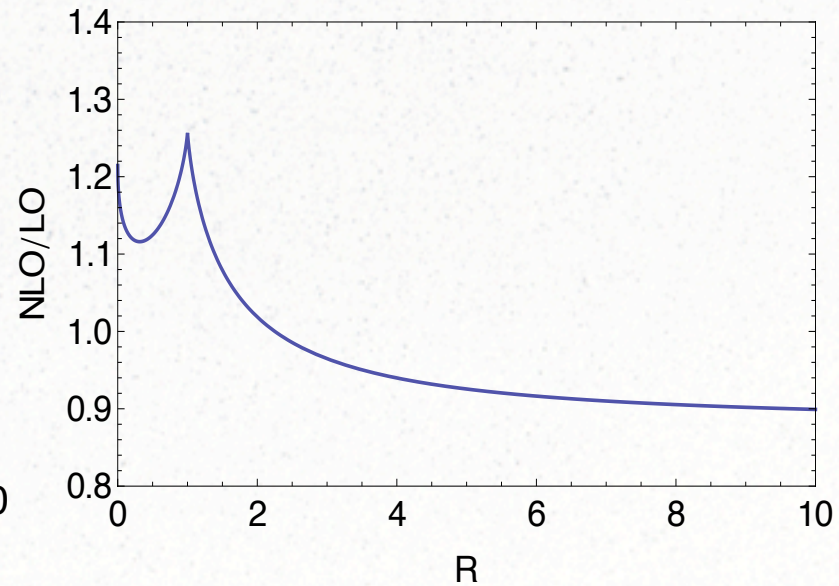
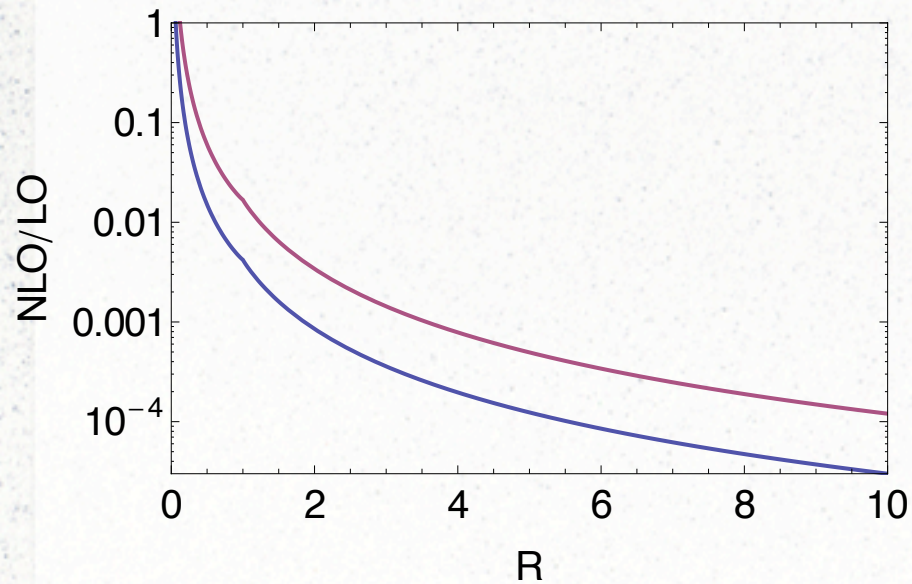
- point:  $-2\epsilon$

$$\text{contrib4} = \frac{1}{4}(R - R\log(R)) - \frac{R}{4\epsilon}$$

$$\begin{aligned} \sum_{n=1}^{n=\infty} \text{contrib5}(n) = & -1 - \frac{5R}{6} + \frac{3\pi^2\sqrt{R}}{8} \\ & + \frac{R^{3/2}\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right)}{2\sqrt{R+4}} \\ & - 2\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right)^2 \\ & + \frac{3\sqrt{R}\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right)}{\sqrt{R+4}} \\ & + \frac{4\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right)}{\sqrt{R}\sqrt{R+4}} \\ & - \frac{3}{4}\sqrt{R}\log(R)\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right) \\ & + \frac{3}{2}\sqrt{R}\log(\sqrt{R+4}-2)\text{csch}^{-1}\left(\frac{2}{\sqrt{R}}\right) \\ & - 3\sqrt{R}\text{Li}_2\left(\frac{2}{\sqrt{R}+\sqrt{R+4}}\right) + \frac{3\sqrt{R}}{4}\text{Li}_2\left(\frac{2}{R+\sqrt{R+4}\sqrt{R+2}}\right) \end{aligned}$$

# *Impact factor results*

- 2 masses of the heavy quark: 5 and 10 GeV
- Ratio



## *Summary and Outlook*

- Numerical results for heavy quark impact factor at NLO obtained

### **Next steps:**

- Running coupling
- Convolution with the gluon Green function and heavy quark PDFs
- Cross section for single bottom at large rapidities