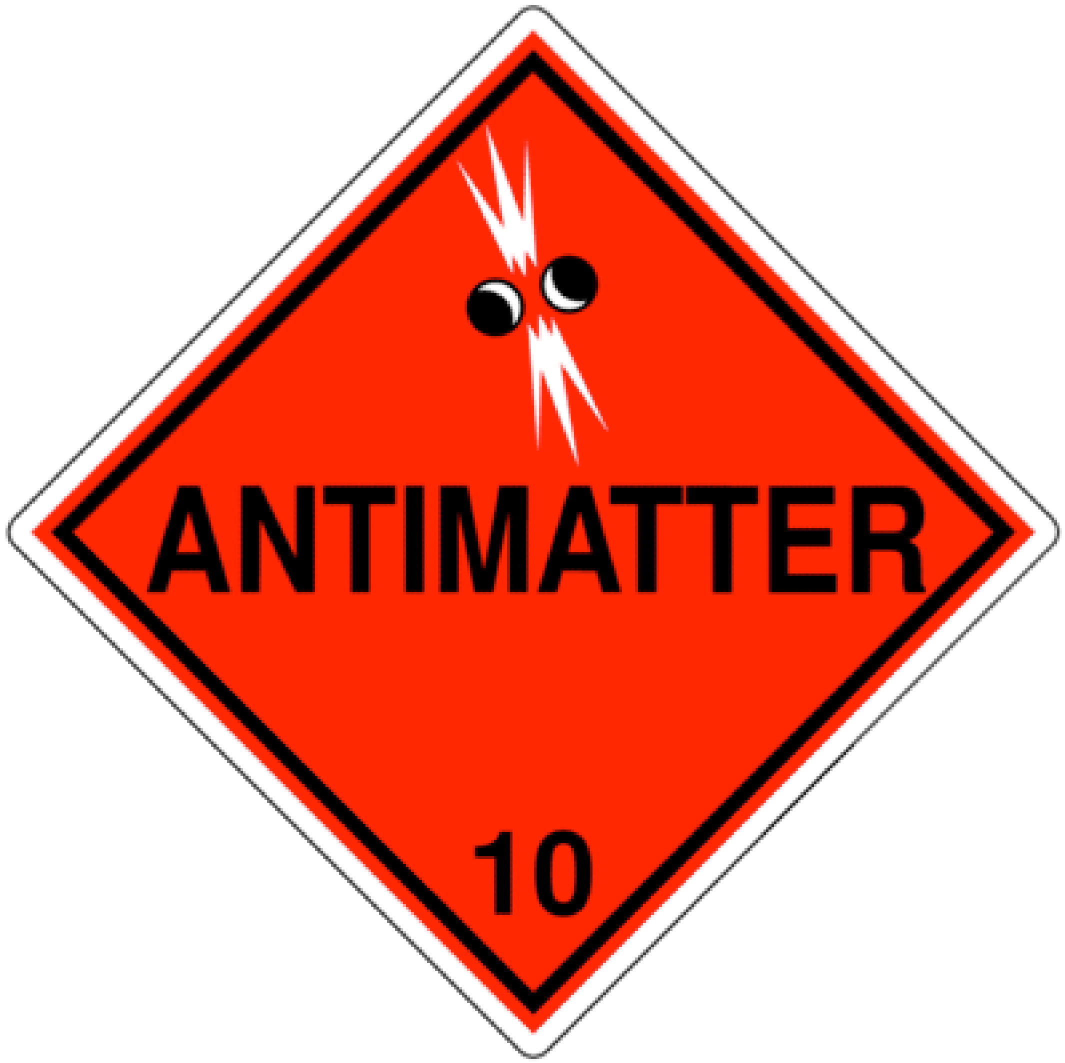
Lectures on Antimatter



Michael Doser / CERN

This is what it's all about:

This is what it's all about:

"abbiamo l'antimateria!"

(in the film Angels and Demons)

Overview:

1. Introduction and overview
2. Antimatter at high energies (SppS, LEP, Fermilab)
3. Meson spectroscopy (antimatter as QCD probe)

4. Astroparticle physics and cosmology
5. CP and CPT violation tests
6. Precision tests with Antimatter

7. Precision tests with Antihydrogen
8. Applications of antimatter

Acknowledgement:

These lectures contain a wide range of material, from many sources. I have endeavored to provide links to publications in many places. Some of the sources, from which slides, graphs, drawings or thoughts were liberally appropriated are in addition presentations, lectures or publications by:

Gerald Gabrielse, Eberhard Widmann, Rolf Landua, Michael Holzscheiter, and many resources from the internet, specifically those dealing with the astroparticle-physics and cosmological aspects of antimatter.

Overview:

1. Introduction and overview
2. Antimatter at high energies (SppS, LEP, Fermilab)
3. Meson spectroscopy (antimatter as QCD probe)

4. Astroparticle physics and cosmology
5. CP and CPT violation tests
6. Precision tests with Antimatter

7. Precision tests with Antihydrogen
8. Applications of antimatter

Introduction and overview

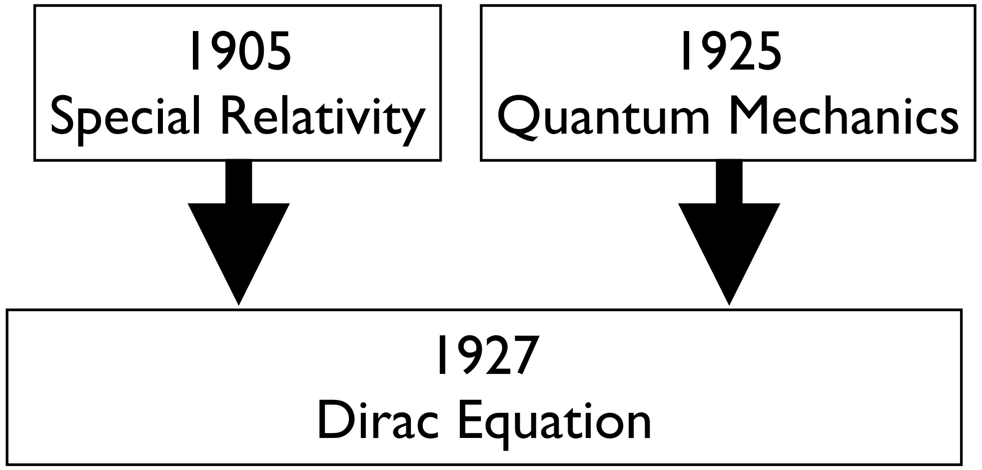
1. A bit of theory

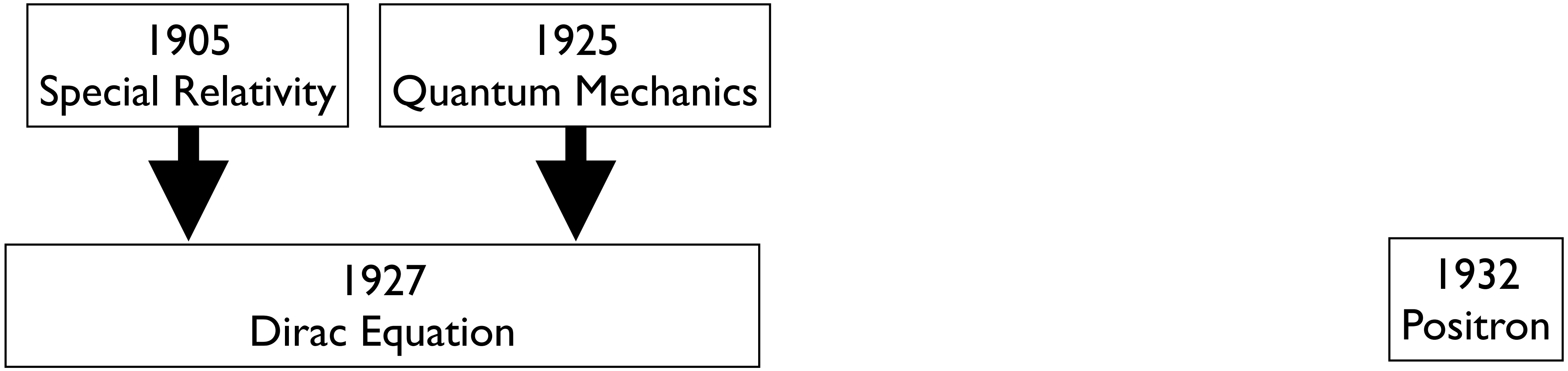
2. A bit of history

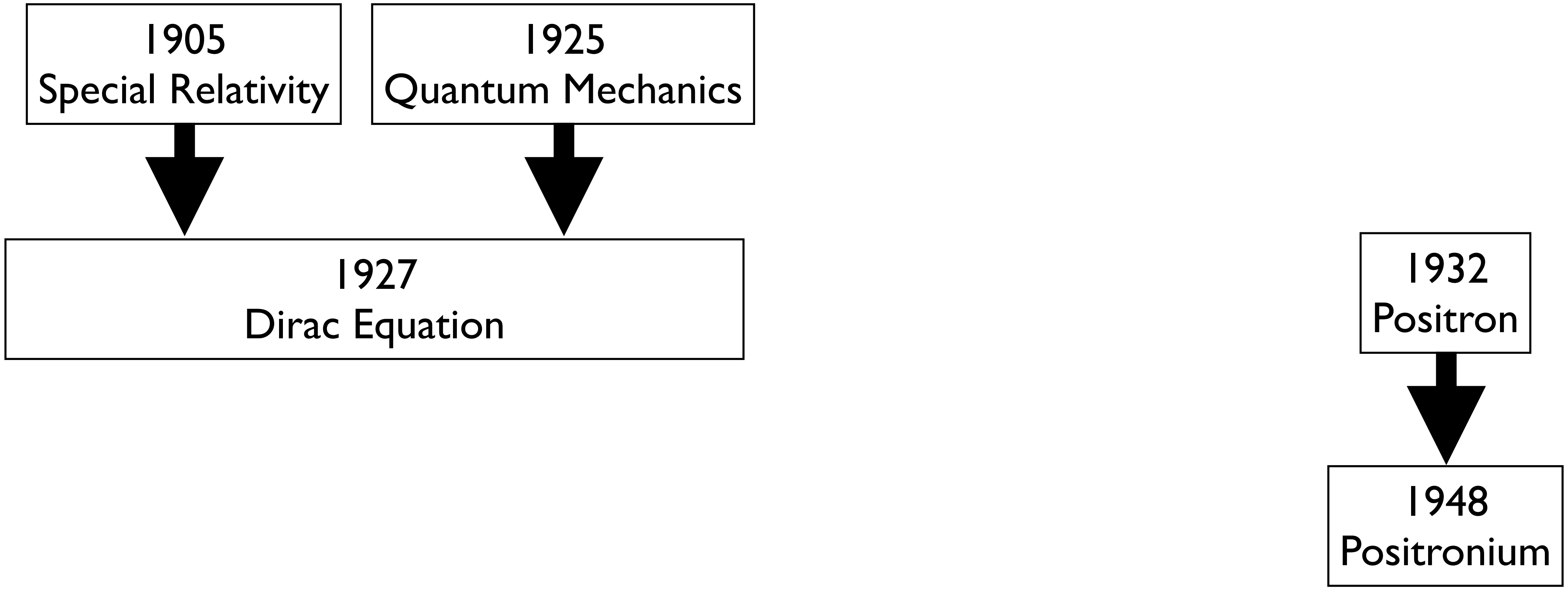
3. The making of...

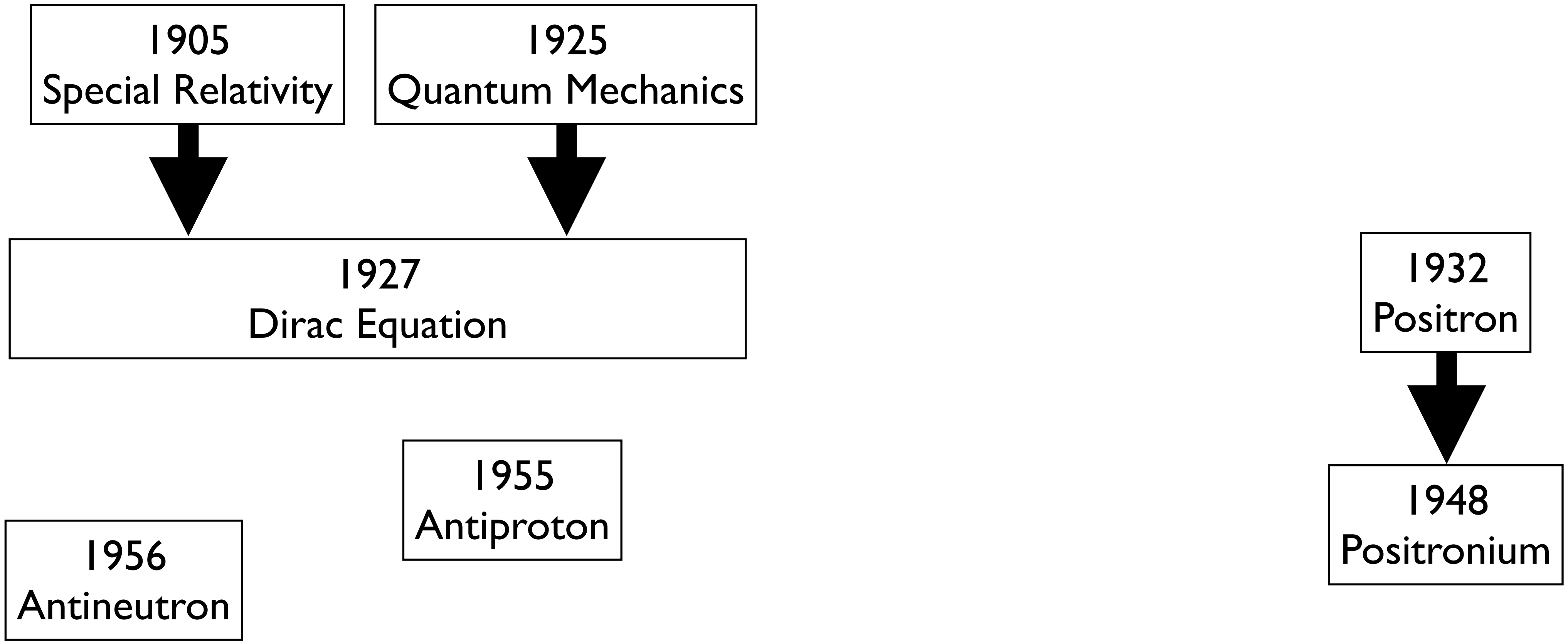
1905
Special Relativity

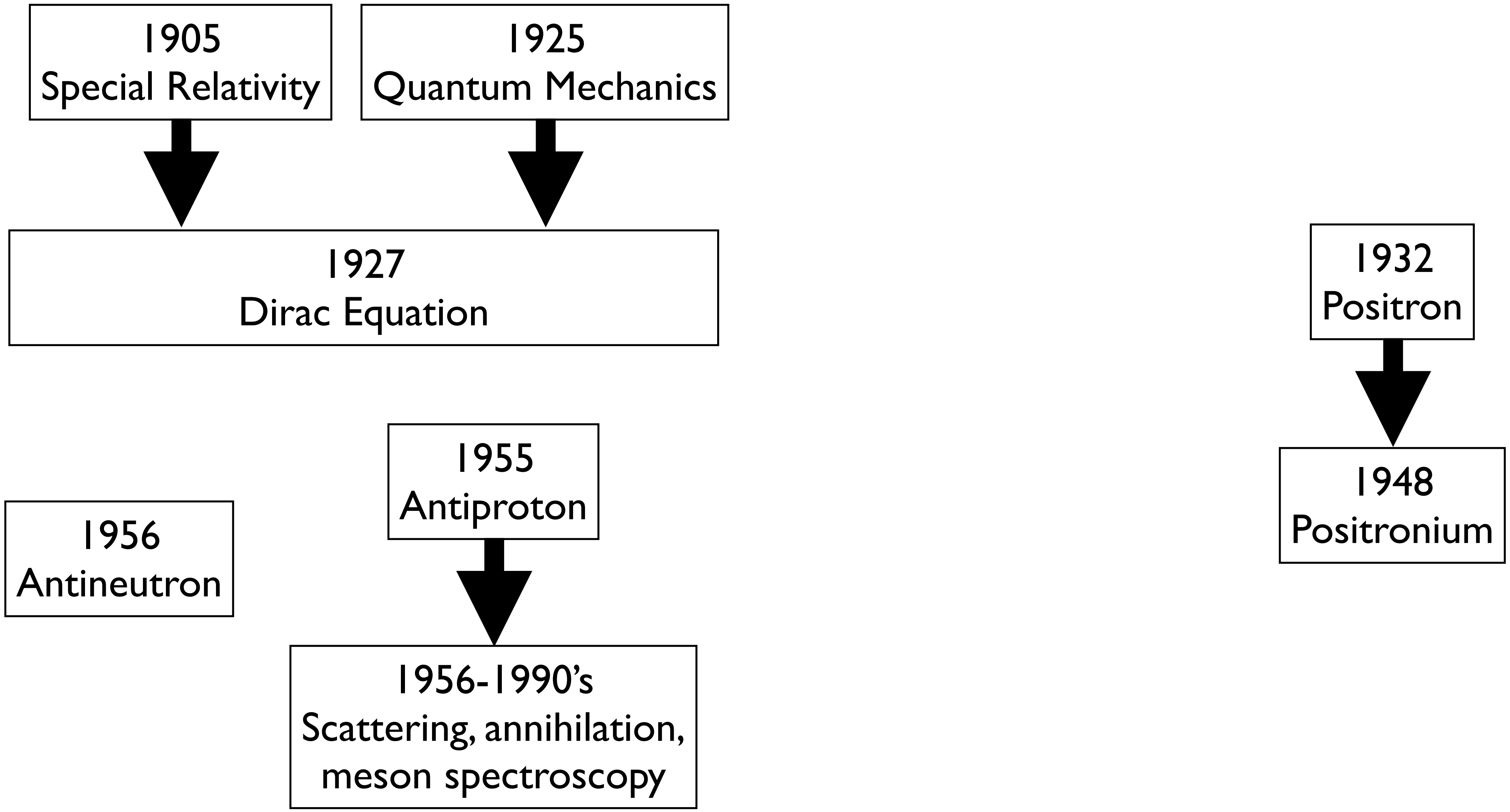
1905 Special Relativity	1925 Quantum Mechanics

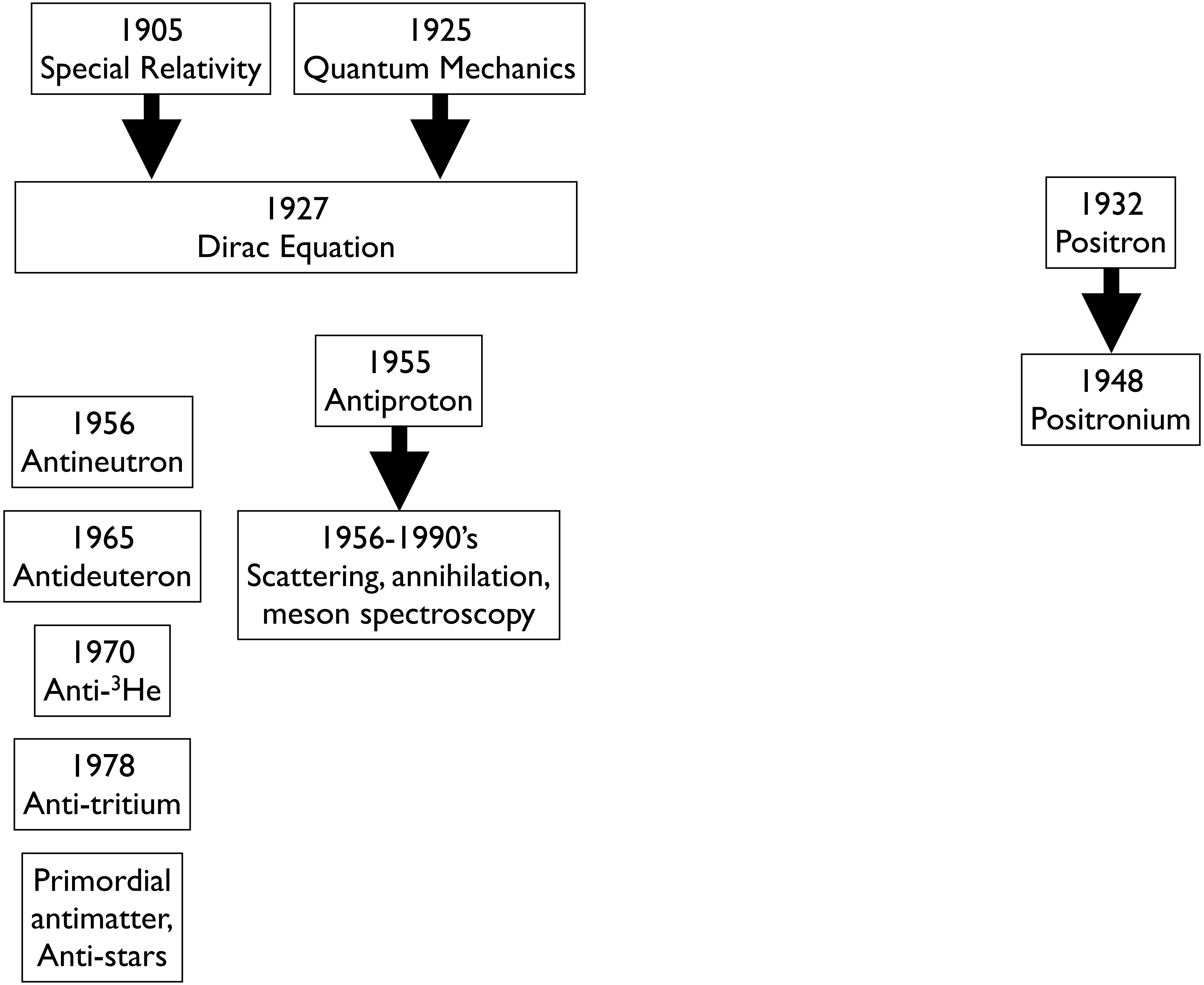

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation

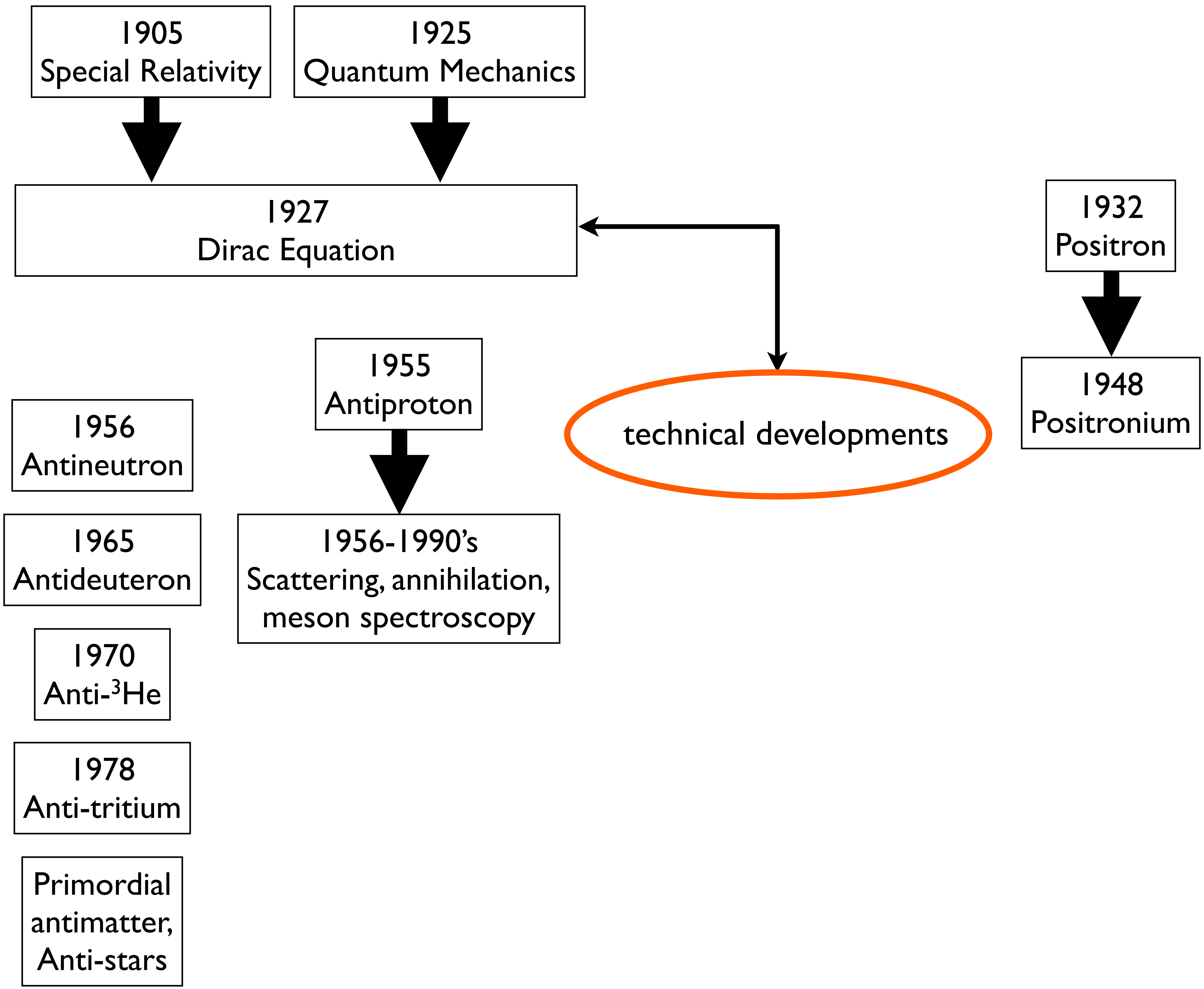

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron

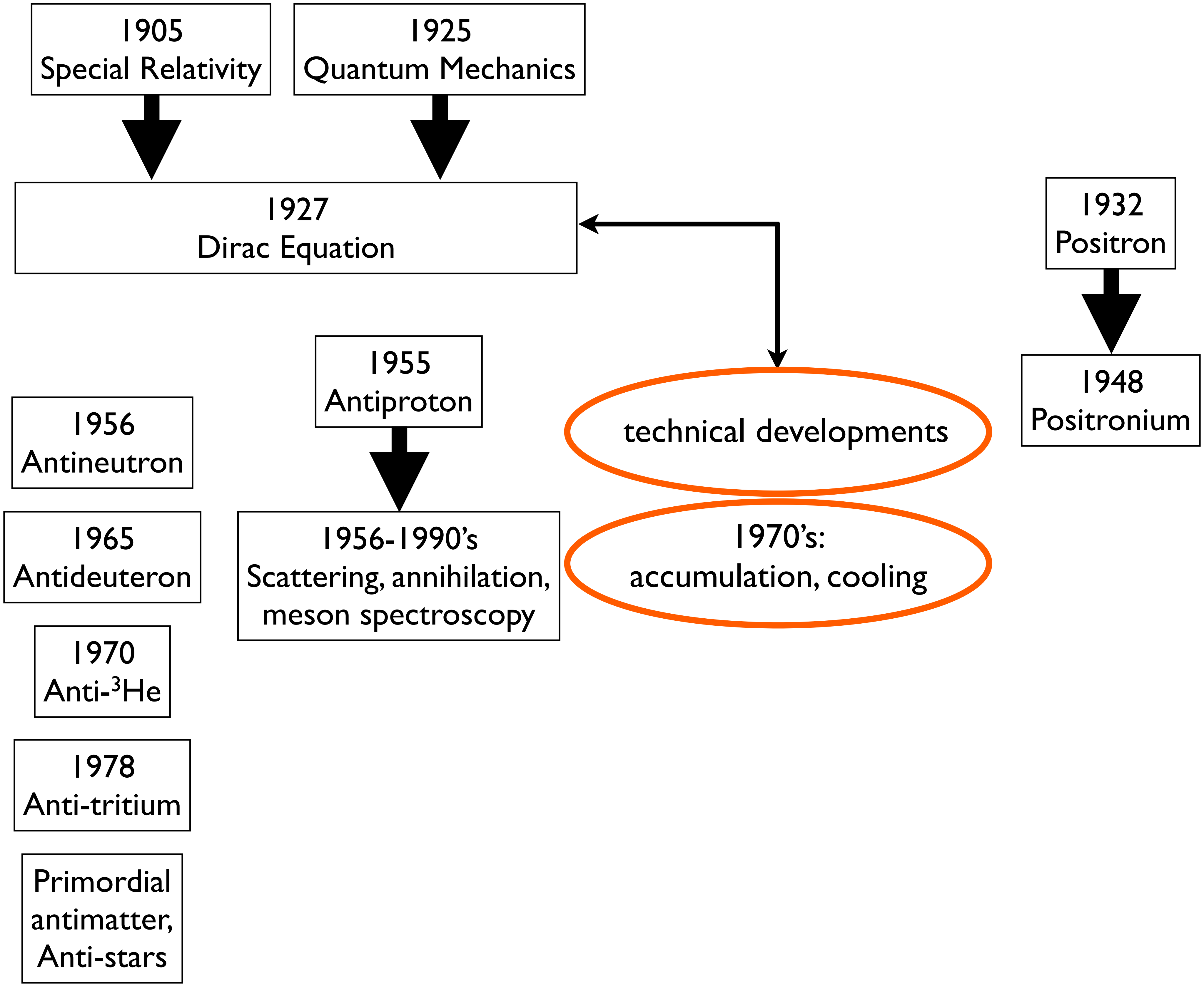

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956
Antineutron

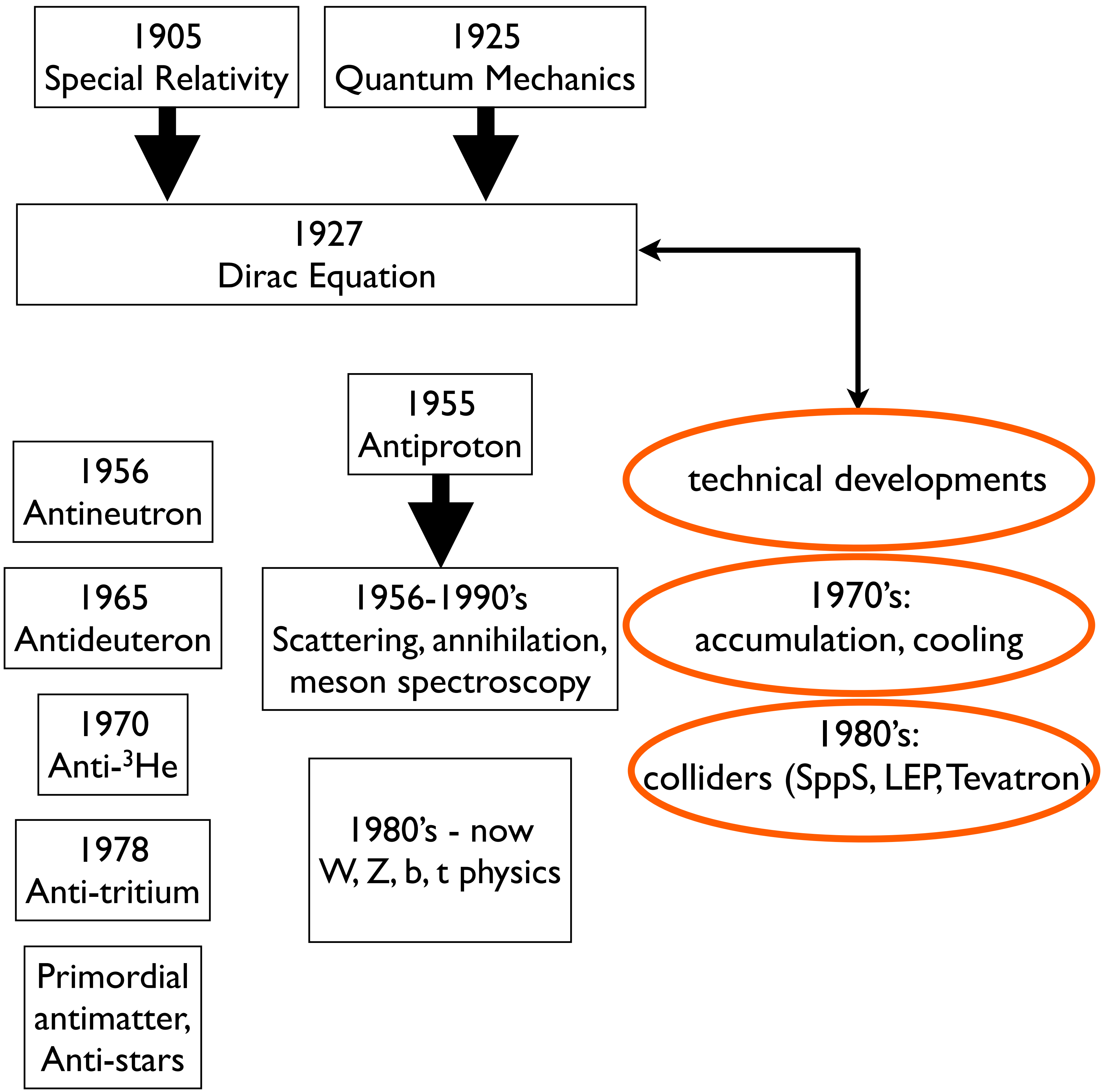

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956
Antineutron
1956-1990's
Scattering, annihilation,
meson spectroscopy

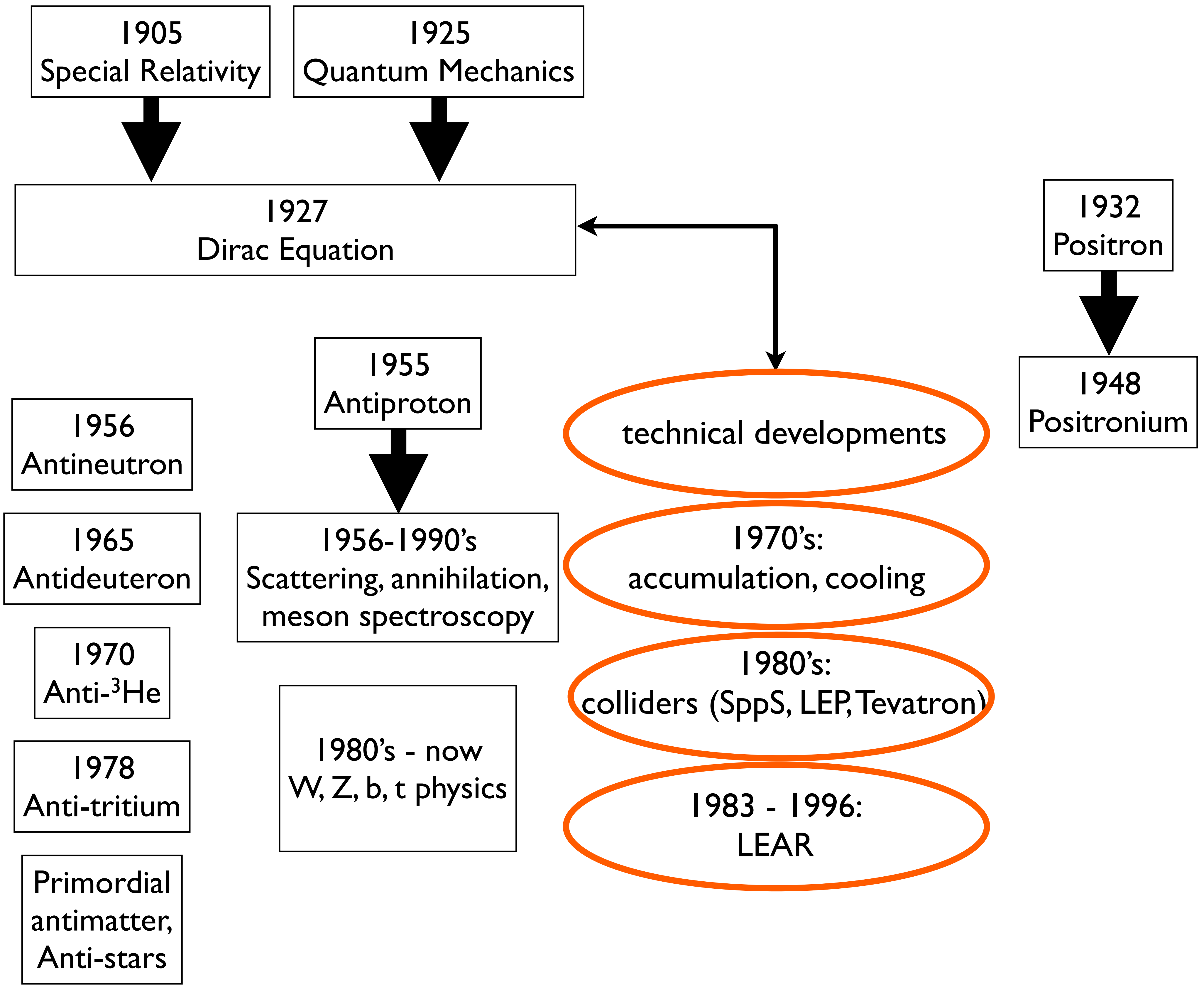

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956-1990's
Scattering, annihilation, meson spectroscopy
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial antimatter, Anti-stars

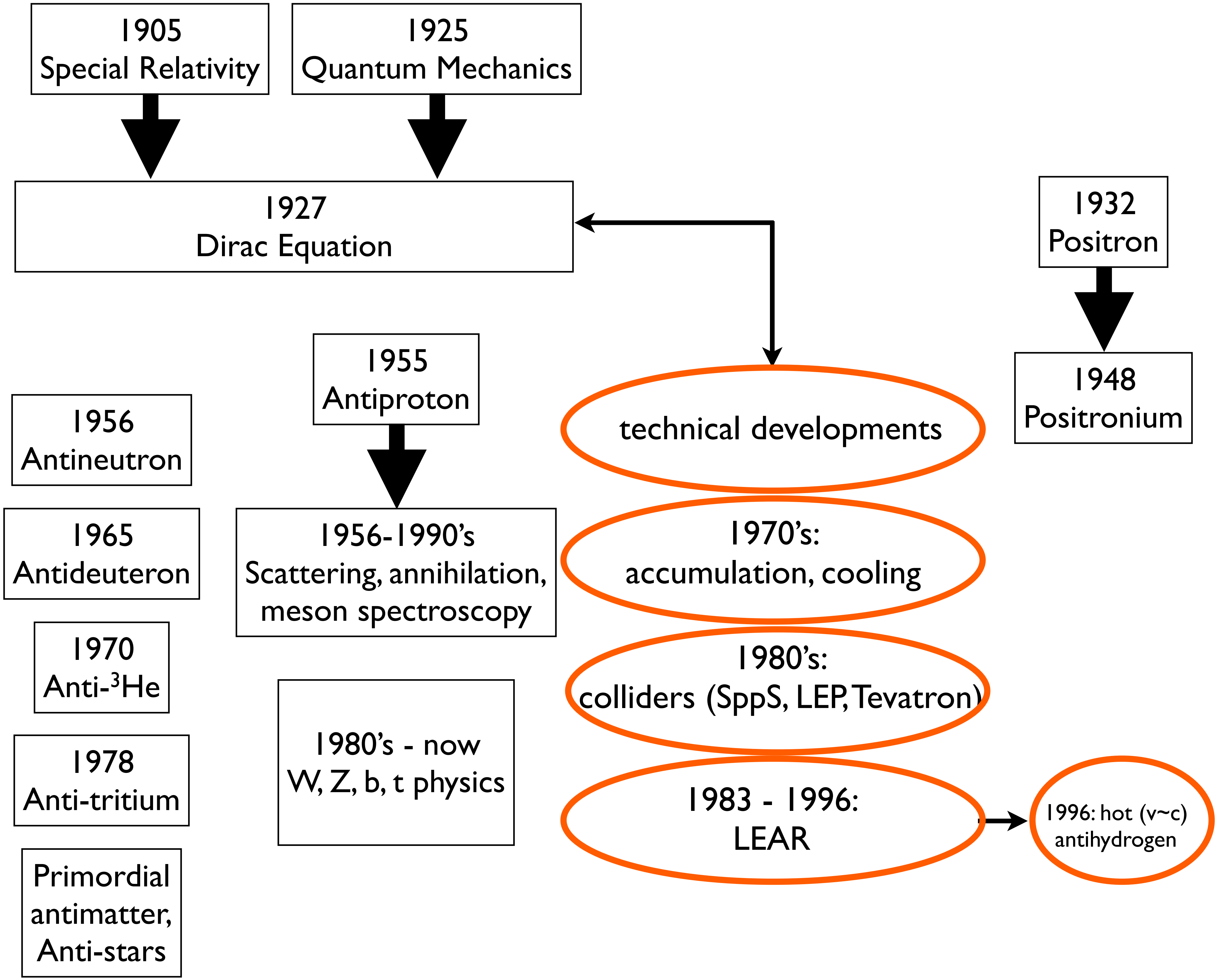

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
technical developments
1932
Positron
1948
Positronium
1955
Antiproton
1956-1990's
Scattering, annihilation,
meson spectroscopy
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial
antimatter,
Anti-stars

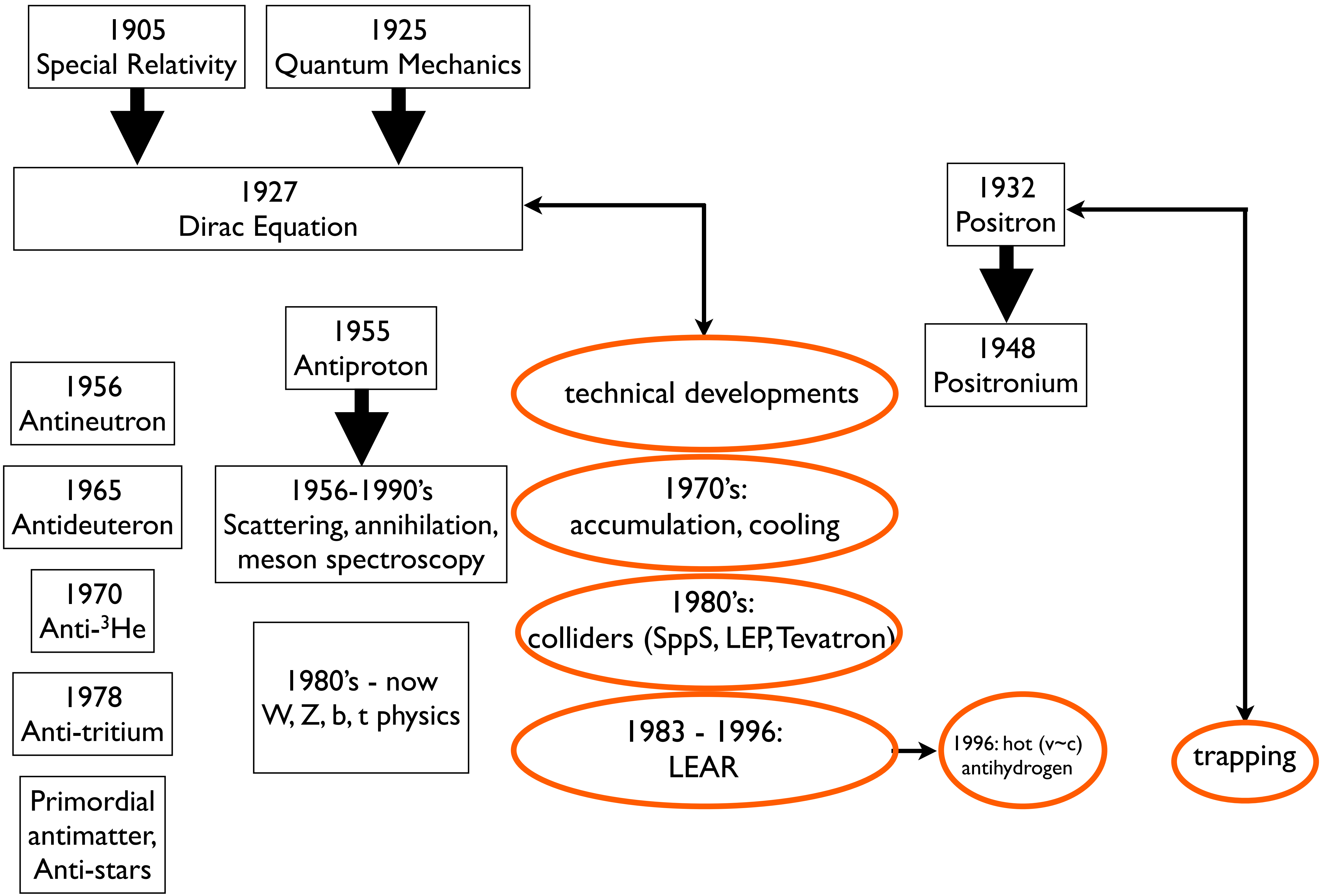


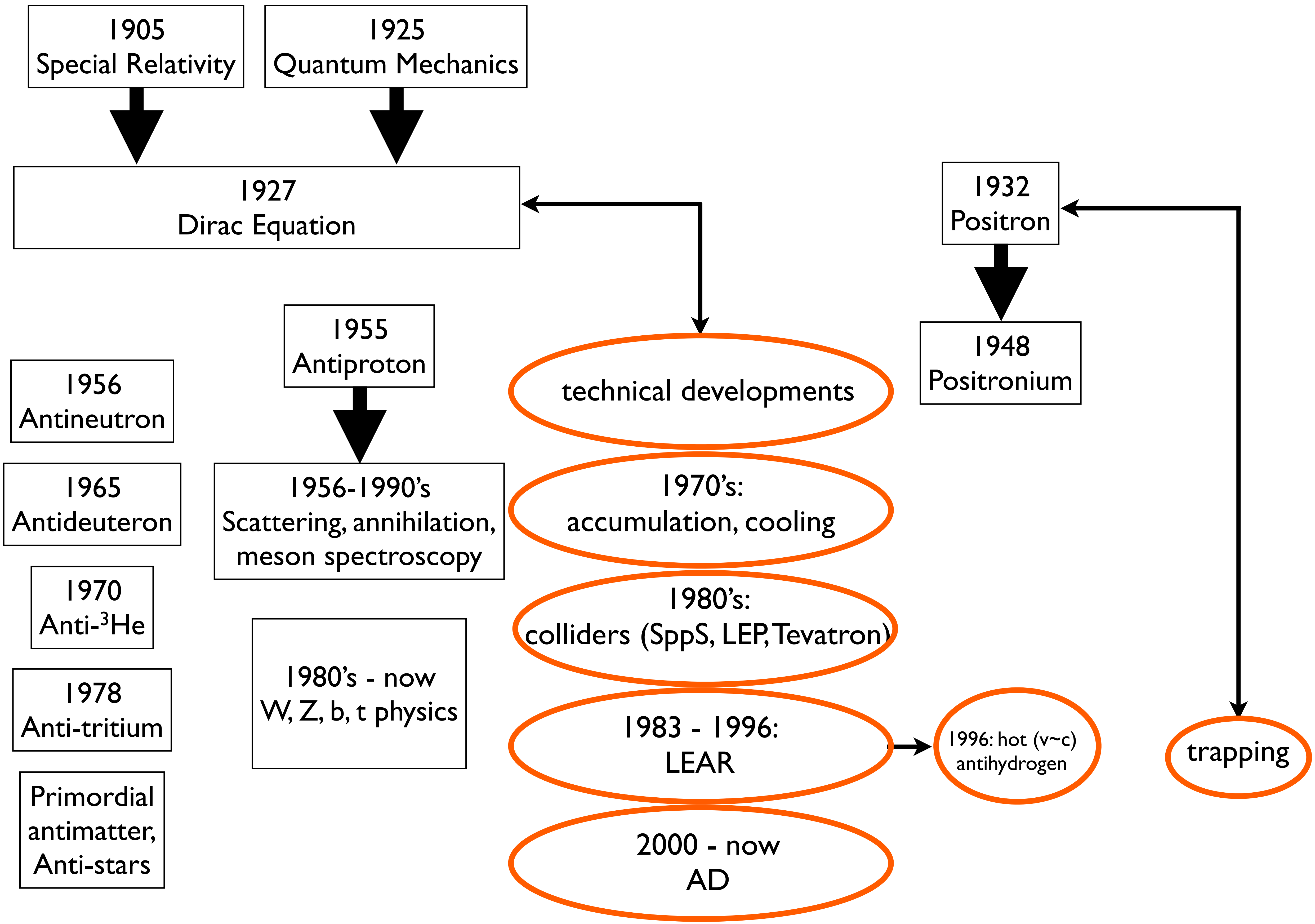
1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
technical developments
1956
Antineutron
1965
Antideuteron
1956-1990's
Scattering, annihilation,
meson spectroscopy
1970's:
accumulation, cooling
1970
Anti-^{3}He
1978
Anti-tritium
Primordial
antimatter,
Anti-stars

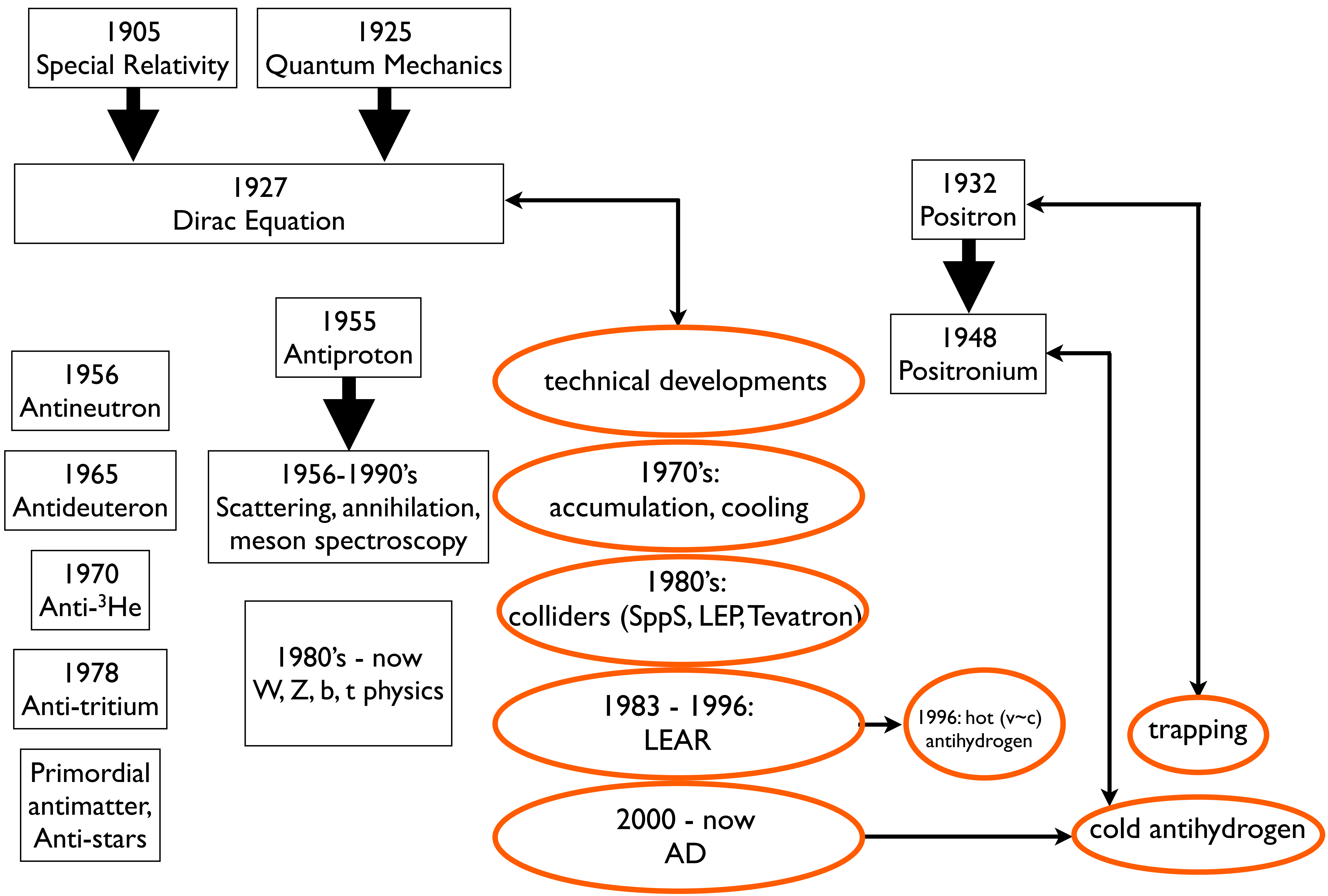

1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
technical developments
1970's:
accumulation, cooling
1980's:
colliders (SppS, LEP, Tevatron)
1955
Antiproton
1956-1990's
Scattering, annihilation, meson spectroscopy
1980's - now
W, Z, b, t physics
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial antimatter, Anti-stars


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
technical developments
1956
Antineutron
1956-1990's
Scattering, annihilation,
meson spectroscopy
1970's:
accumulation, cooling
1965
Antideuteron
1970
Anti-^{3}He
1980's:
colliders (SppS, LEP, Tevatron)
1978
Anti-tritium
1980's - now
W, Z, b, t physics
1983 - 1996:
LEAR
Primordial
antimatter,
Anti-stars


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial antimatter, Anti-stars
1956-1990's
Scattering, annihilation, meson spectroscopy
1980's - now
W, Z, b, t physics
technical developments
1970's:
accumulation, cooling
1980's:
colliders (SppS, LEP, Tevatron)
1983 - 1996:
LEAR
1996: hot (v~c) antihydrogen


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956-1990's
Scattering, annihilation, meson spectroscopy
technical developments
1970's:
accumulation, cooling
1980's:
colliders (SppS, LEP, Tevatron)
1983 - 1996:
LEAR
1996: hot (v~c) antihydrogen
trapping
1956
Antineutron
1965
Antideuteron
1970
Anti-3He
1978
Anti-tritium
Primordial antimatter, Anti-stars
1980's - now
W, Z, b, t physics


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial antimatter, Anti-stars
1956-1990's
Scattering, annihilation, meson spectroscopy
1980's - now
W, Z, b, t physics
technical developments
1970's:
accumulation, cooling
1980's:
colliders (SppS, LEP, Tevatron)
1983 - 1996:
LEAR
1996: hot (v~c) antihydrogen
2000 - now
AD
trapping


1905
Special Relativity
1925
Quantum Mechanics
1927
Dirac Equation
1932
Positron
1948
Positronium
1955
Antiproton
1956
Antineutron
1965
Antideuteron
1970
Anti-^{3}He
1978
Anti-tritium
Primordial
antimatter,
Anti-stars
1956-1990's
Scattering, annihilation,
meson spectroscopy
1980's - now
W, Z, b, t physics
technical developments
1970's:
accumulation, cooling
1980's:
colliders (SppS, LEP, Tevatron)
1983 - 1996:
LEAR
1996: hot (v~c)
antihydrogen
2000 - now
AD
trapping
cold antihydrogen

Schrödinger:

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

$$E = \frac{p^2}{2m} \rightarrow i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

Schrödinger:

$$E = \frac{p^2}{2m} \rightarrow \quad i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

relativistic energy-momentum relation

Klein-Gordon:

$$E^2 = p^2 + m^2 \rightarrow \quad -\hbar^2\frac{\partial^2}{\partial t^2}\psi = -\hbar^2\nabla^2\psi + m^2\psi$$

relativistic, spin 0
(number of particles not conserved)

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

Schrödinger:

$$E = \frac{p^2}{2m} \rightarrow \quad i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

relativistic energy-momentum relation

Klein-Gordon:

$$E^2 = p^2 + m^2 \rightarrow \quad -\hbar^2\frac{\partial^2}{\partial t^2}\psi = -\hbar^2\nabla^2\psi + m^2\psi$$

relativistic, spin 0
(number of particles not conserved)

energy Eigenvalues (free particle)
$E = +/- (\mathbf{p}^2+m^2)^{1/2}$

negative energy solutions with
negative probability density

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

Schrödinger:

$$E = \frac{p^2}{2m} \rightarrow \quad i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

relativistic energy-momentum relation

Klein-Gordon:

$$E^2 = p^2 + m^2 \rightarrow \quad -\hbar^2\frac{\partial^2}{\partial t^2}\psi = -\hbar^2\nabla^2\psi + m^2\psi$$

relativistic, spin 0
(number of particles not conserved)

negative energy solutions with negative probability density

energy Eigenvalues (free particle)
E = +/- ($\mathbf{p}^2$+m^2)$^{1/2}$

Dirac:

linear in $\frac{\partial}{\partial t}$ and ∇

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

Schrödinger:

$$E = \frac{p^2}{2m} \rightarrow i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

relativistic energy-momentum relation

Klein-Gordon:

$$E^2 = p^2 + m^2 \rightarrow -\hbar^2\frac{\partial^2}{\partial t^2}\psi = -\hbar^2\nabla^2\psi + m^2\psi$$

relativistic, spin 0
(number of particles not conserved)

energy Eigenvalues (free particle)
$E = +/- (\mathbf{p}^2+m^2)^{1/2}$

negative energy solutions with
negative probability density

Dirac:

linear in $\frac{\partial}{\partial t}$ and ∇

general form: $H\psi = (\boldsymbol{\alpha}\cdot\mathbf{P} + \beta m)\psi$

classical energy-momentum relation

acts on wavefuntion $\psi(\mathbf{x}, t)$

Schrödinger:

$$E = \frac{p^2}{2m} \rightarrow i\hbar\frac{\partial}{\partial t}\psi = -\frac{\hbar^2}{2m}\nabla^2\psi$$

non-relativistic

$E \rightarrow i\hbar\frac{\partial}{\partial t}$

$p \rightarrow -i\hbar\nabla$

differential operators

relativistic energy-momentum relation

Klein-Gordon:

$$E^2 = p^2 + m^2 \rightarrow -\hbar^2\frac{\partial^2}{\partial t^2}\psi = -\hbar^2\nabla^2\psi + m^2\psi$$

relativistic, spin 0
(number of particles not conserved)

energy Eigenvalues (free particle)
$E = +/- (\mathbf{p}^2+m^2)^{1/2}$

negative energy solutions with
negative probability density

Dirac:

linear in $\frac{\partial}{\partial t}$ and ∇

general form: $$H\psi = (\boldsymbol{\alpha}\cdot\mathbf{P} + \beta m)\psi$$

energy-momentum relationship: $$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

$$= (\alpha_i^2 P_i^2 + (\alpha_i \alpha_j + \alpha_j \alpha_i) P_i P_j + (\alpha_i \beta + \beta \alpha_i) P_i m + \beta^2 m^2)\psi$$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

$$= (\alpha_i^2 P_i^2 + (\alpha_i \alpha_j + \alpha_j \alpha_i) P_i P_j + (\alpha_i \beta + \beta \alpha_i) P_i m + \beta^2 m^2)\psi$$

!

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_i P_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

!

$$= (\underline{\alpha_i^2} P_i^2 + \underline{(\alpha_i \alpha_j + \alpha_j \alpha_i)} P_i P_j + (\alpha_i \beta + \beta \alpha_i) P_i m + \beta^2 m^2)\psi$$

1 0

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$= (\underline{\alpha_i^2} P_i^2 + \underline{(\alpha_i\alpha_j + \alpha_j\alpha_i)} P_i P_j + \underline{(\alpha_i\beta + \beta\alpha_i)} P_i m + \beta^2 m^2)\psi$$

1 0 0

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_iP_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$\boxed{H^2\psi = (\mathbf{P^2} + m^2)\psi}$$

$$= (\underline{\alpha_i^2} P_i^2 + \underline{(\alpha_i \alpha_j + \alpha_j \alpha_i)} P_i P_j + \underline{(\alpha_i \beta + \beta \alpha_i)} P_i m + \underline{\beta^2} m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_iP_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$$

lowest dim. matrices: 4x4 ; Pauli-Dirac representation

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$= (\alpha_i^2 P_i^2 + (\alpha_i \alpha_j + \alpha_j \alpha_i) P_i P_j + (\alpha_i \beta + \beta \alpha_i) P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$

lowest dim. matrices: 4x4 ; Pauli-Dirac representation

ψ : 4-component column vector (Dirac spinor)

(E>0,+1/2);(E>0,-1/2);(E<0,+1/2);(E<0,-1/2)

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

!

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_iP_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$$

lowest dim. matrices: 4x4 ; Pauli-Dirac representation

ψ : 4-component column vector (Dirac spinor)
(E>0,+1/2);(E>0,-1/2);(E<0,+1/2);(E<0,-1/2)

$\times\beta$

$$i\beta\frac{\partial\psi}{\partial t} = -i\beta\alpha\nabla\psi + m\psi$$

$$H\psi = (\boldsymbol{\alpha} \cdot \mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_iP_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$$

lowest dim. matrices: 4x4 ; Pauli-Dirac representation

ψ: 4-component column vector (Dirac spinor)
(E>0,+1/2);(E>0,-1/2);(E<0,+1/2);(E<0,-1/2)

$\times\beta$

$$i\beta\frac{\partial\psi}{\partial t} = -i\beta\alpha\nabla\psi + m\psi$$

$$(i\gamma^\mu\partial_\mu - m)\psi = 0$$

$$H\psi = (\boldsymbol{\alpha}\cdot\mathbf{P} + \beta m)\psi$$

$$H^2\psi = (\mathbf{P^2} + m^2)\psi$$

$$H^2\psi = (\alpha_i P_i + \beta m)(\alpha_j P_j + \beta m)\psi$$

!

$$= (\alpha_i^2 P_i^2 + (\alpha_i\alpha_j + \alpha_j\alpha_i)P_iP_j + (\alpha_i\beta + \beta\alpha_i)P_i m + \beta^2 m^2)\psi$$

1 0 0 1

$\alpha_1, \alpha_2, \alpha_3, \beta$ anticommute with each other

$\alpha_1^2 = \alpha_2^2 = \alpha_3^2 = \beta^2 = 1$

lowest dim. matrices: 4x4 ; Pauli-Dirac representation

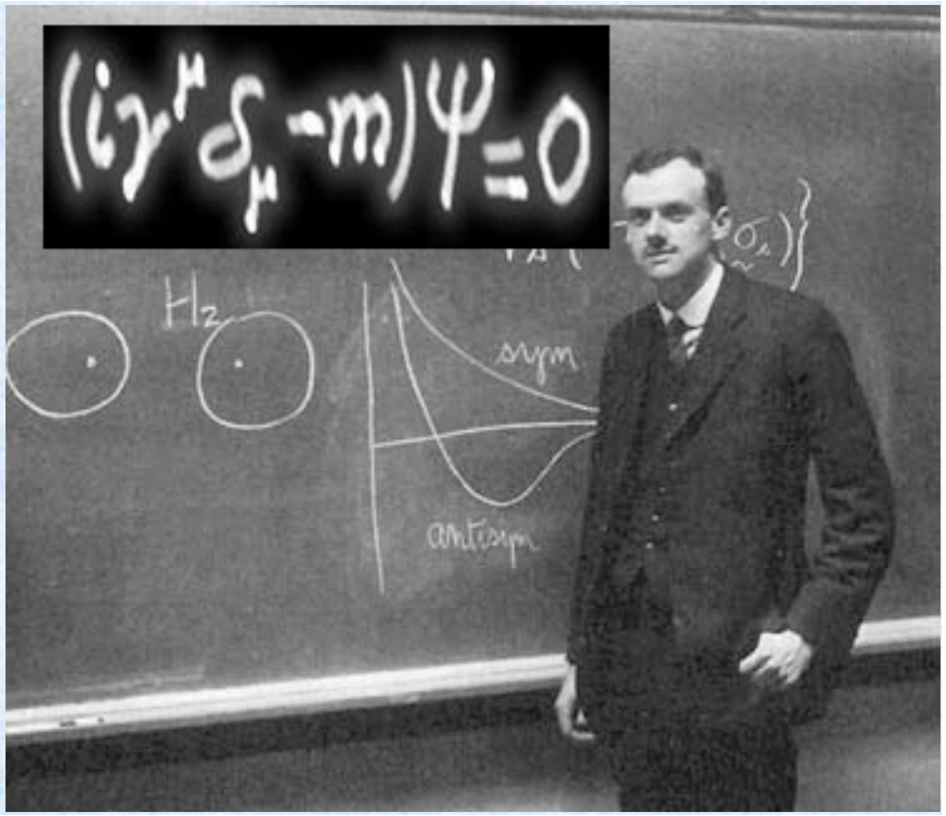
ψ : 4-component column vector (Dirac spinor)
(E>0,+1/2);(E>0,-1/2);(E<0,+1/2);(E<0,-1/2)

$\times\beta$

$$i\beta\frac{\partial\psi}{\partial t} = -i\beta\alpha\nabla\psi + m\psi$$

$$(i\gamma^\mu\partial_\mu - m)\psi = 0$$

relativistic, spin 1/2
(number of particles conserved)



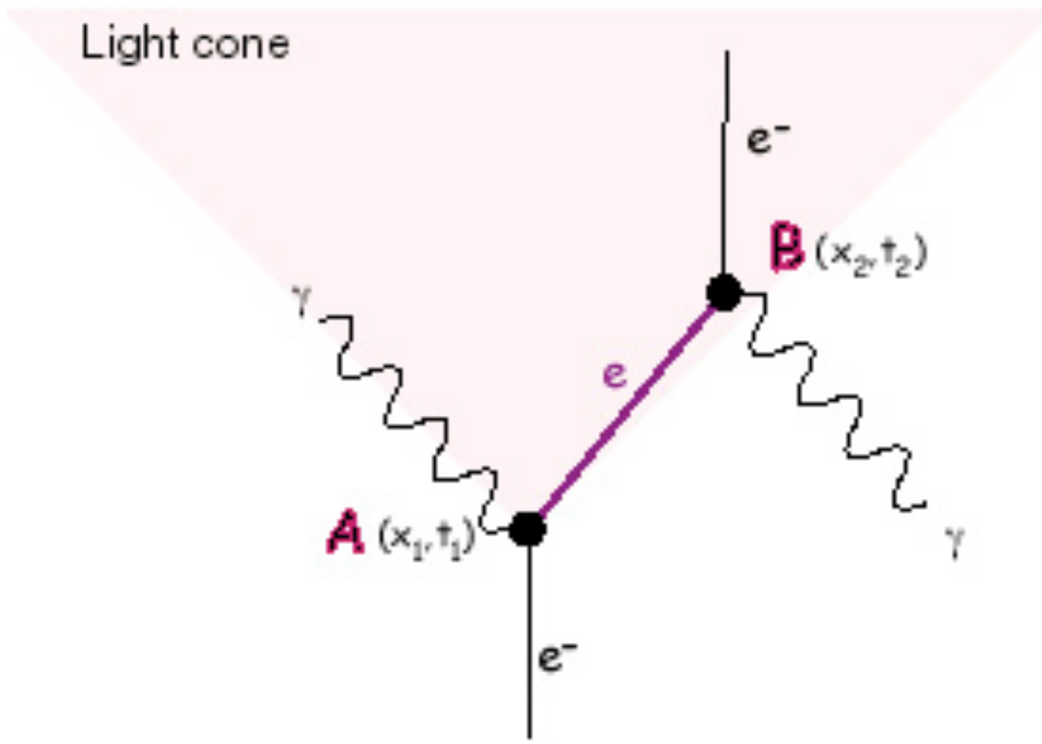
Benefit of hindsight: Quantum Field Theory

Benefit of hindsight: Quantum Field Theory

The electron (field) is no longer described by a wave function but an operator that creates and destroys particles. All energies are positive.

Benefit of hindsight: Quantum Field Theory

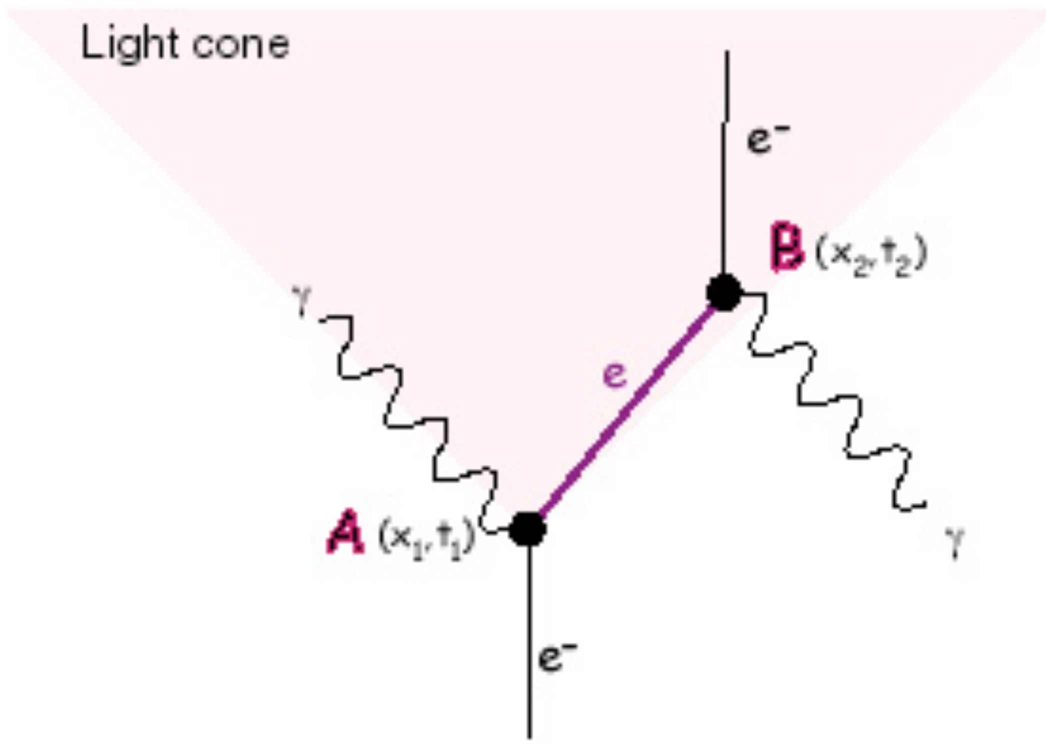
The electron (field) is no longer described by a wave function but an operator that creates and destroys particles. All energies are positive.



Observer #1 : A happens before B

Benefit of hindsight: Quantum Field Theory

The electron (field) is no longer described by a wave function but an operator that creates and destroys particles. All energies are positive.

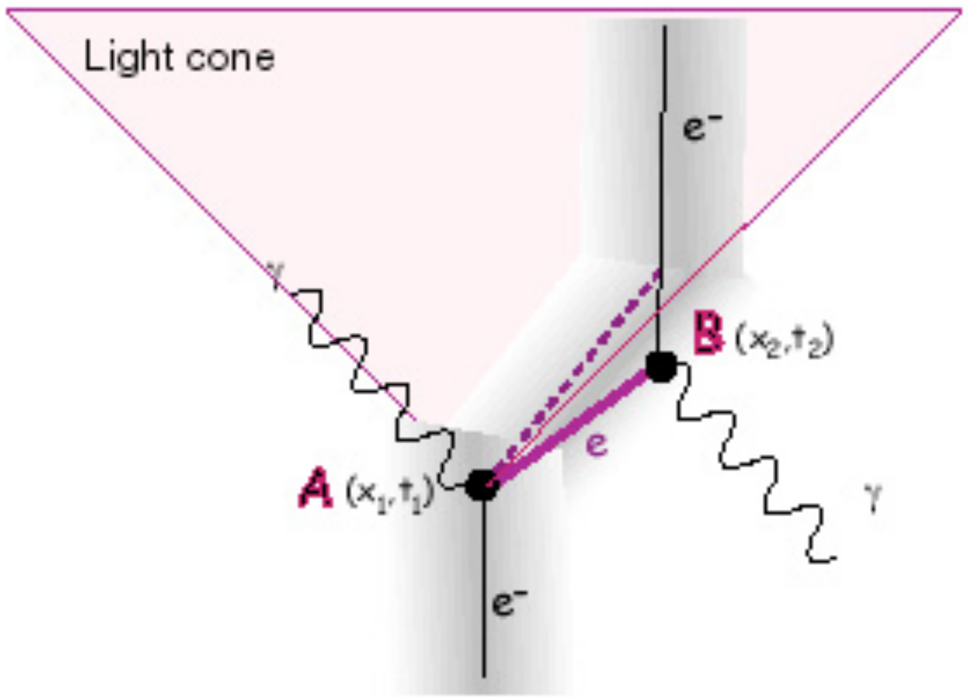


Observer #1 : A happens before B

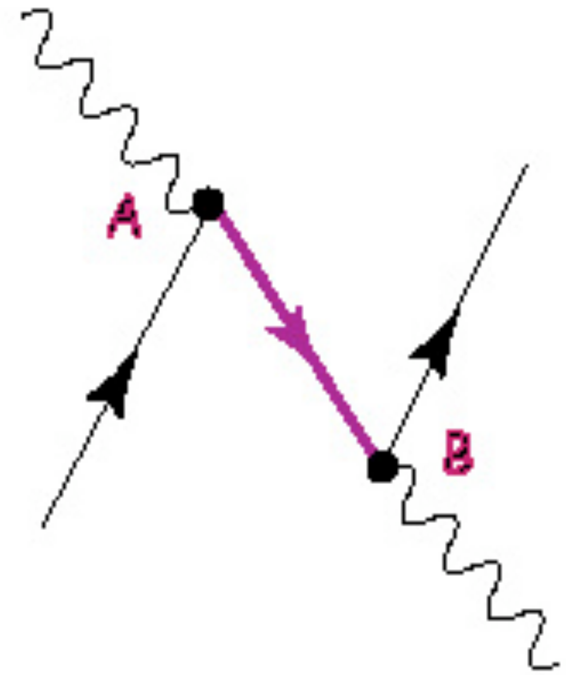
An electron can emit a photon at A, propagate a certain distance, and then absorb another photon at B.

Wave function only localized within Compton wave length (λ ~ 1/m).

Wave function only localized within Compton wave length ($\lambda \sim 1/m$).

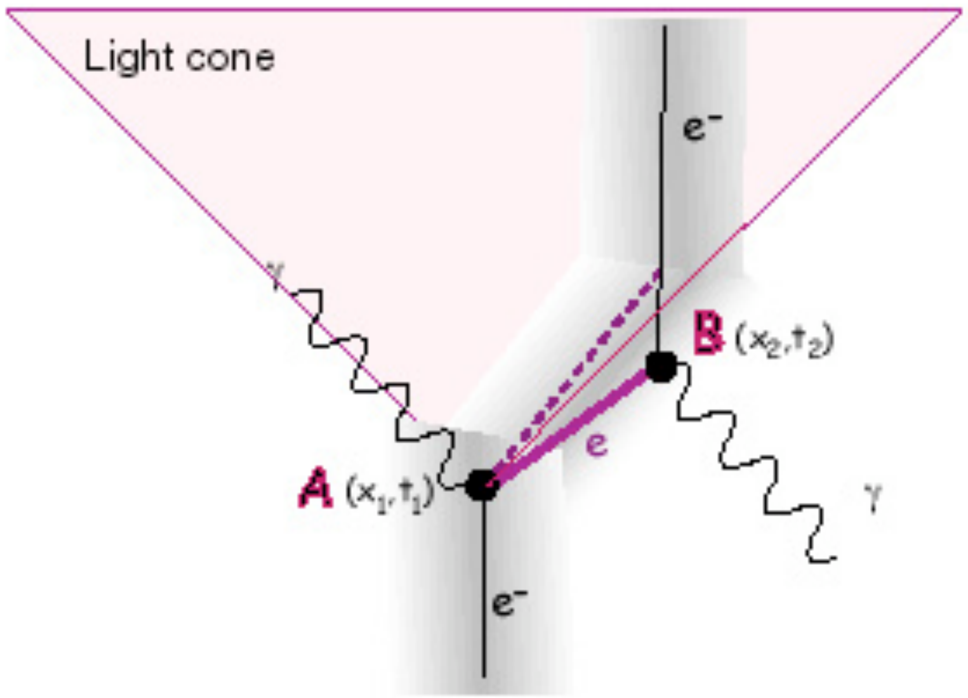


Quantum relativity: electron wave function can be **outside the light cone** (Compton wave length $l = h/m_e c$)

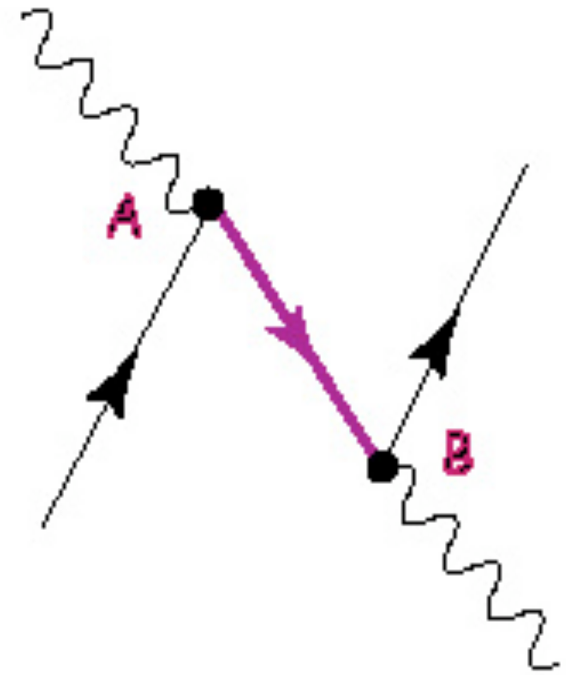


For a moving observer, event B can therefore happen before event A. The process at B is then interpreted as 'pair creation'.

Wave function only localized within Compton wave length ($\lambda \sim 1/m$).



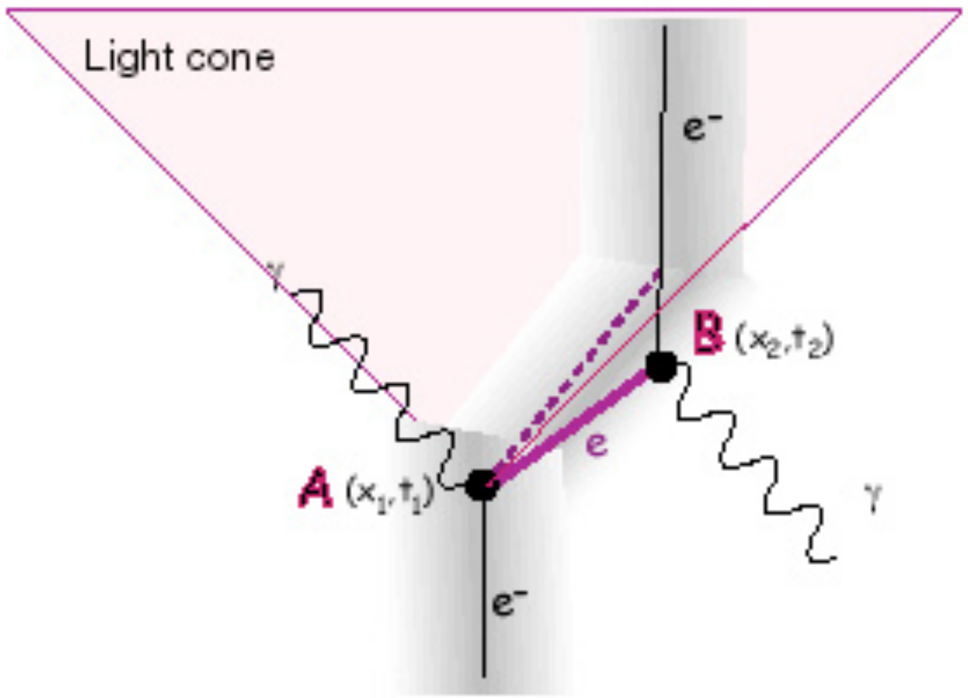
Quantum relativity: electron wave function can be **outside the light cone** (Compton wave length $l = h/m_e c$)



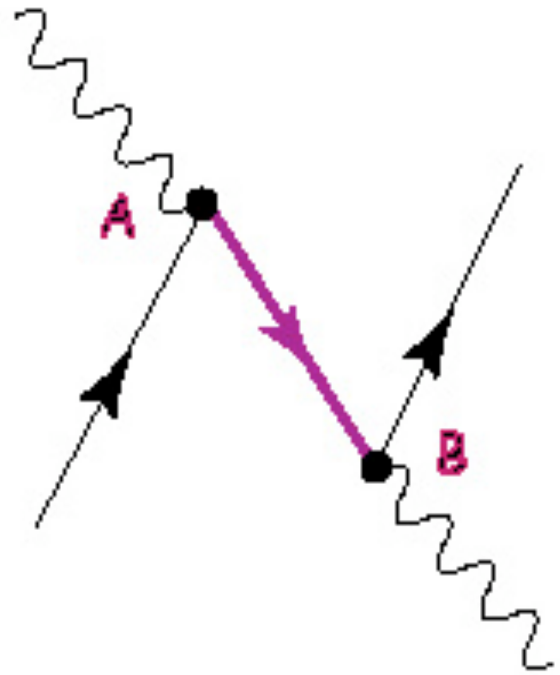
For a moving observer, event B can therefore happen before event A. The process at B is then interpreted as 'pair creation'.

"One observer's electron is the other observer's positron"

Wave function only localized within Compton wave length ($\lambda \sim 1/m$).



Quantum relativity: electron wave function can be **outside the light cone** (Compton wave length $l = h/m_e c$)



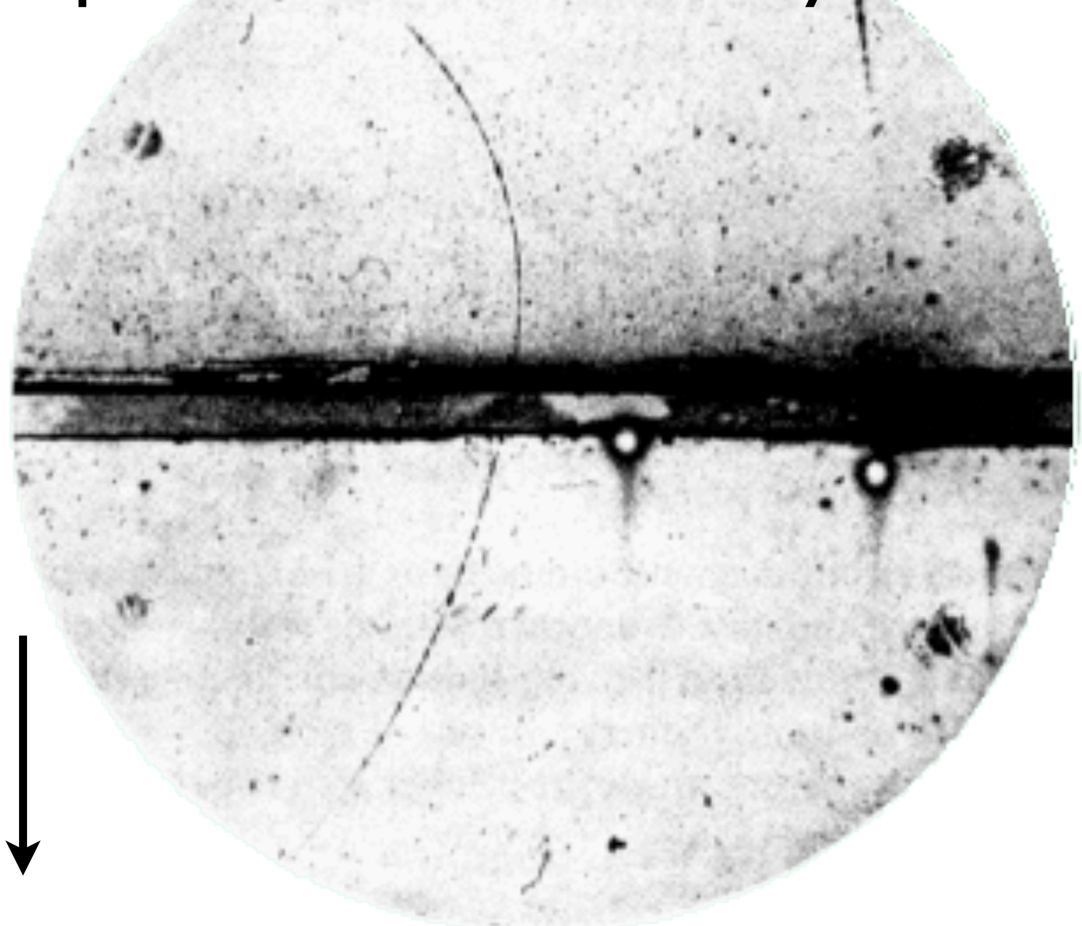
For a moving observer, event B can therefore happen before event A. The process at B is then interpreted as 'pair creation'.

"One observer's electron is the other observer's positron"

Causality requires antiparticles to exist

Antimatter:

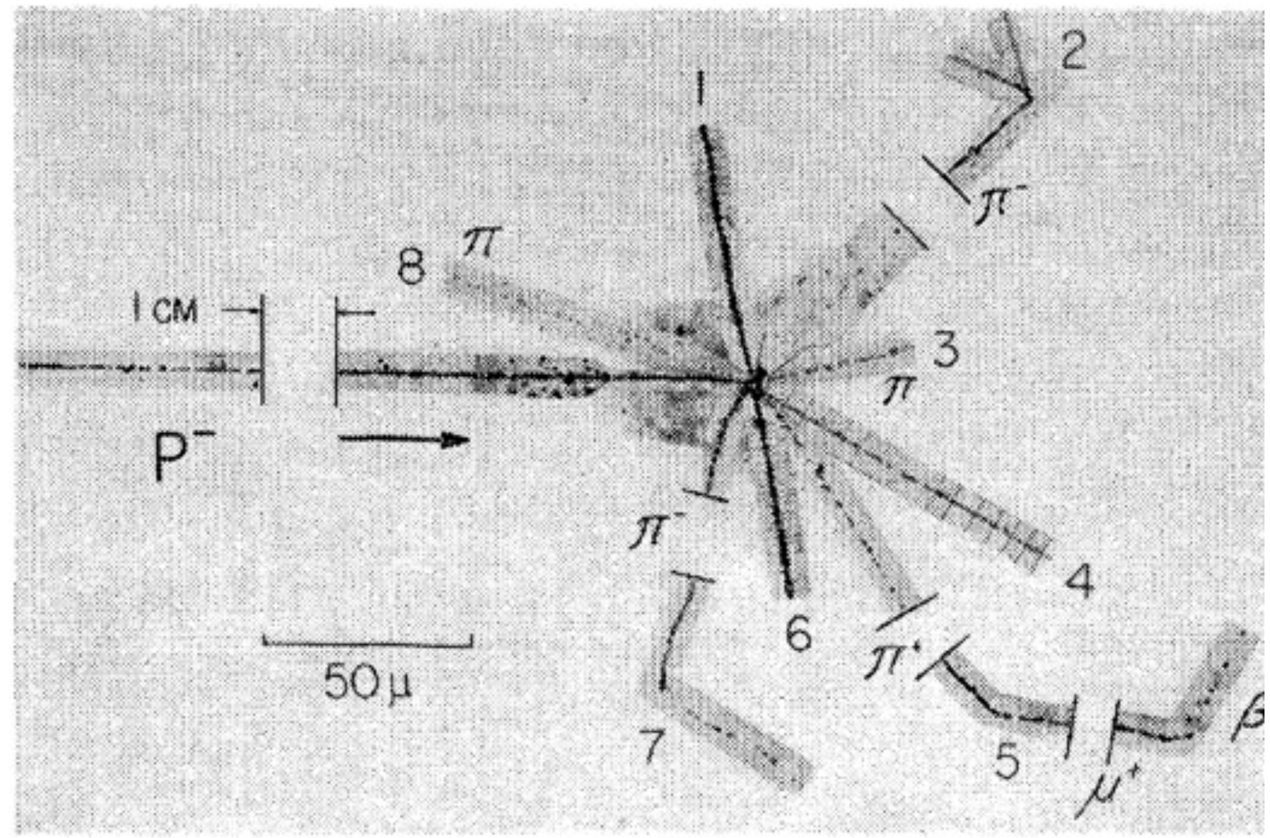
1932 - Anderson discovers the positron in cosmic rays



this way up ↓

Cloud chamber photograph by Andersen
Phys. Rev. 43, 491 (1933)
Nobel prize 1936

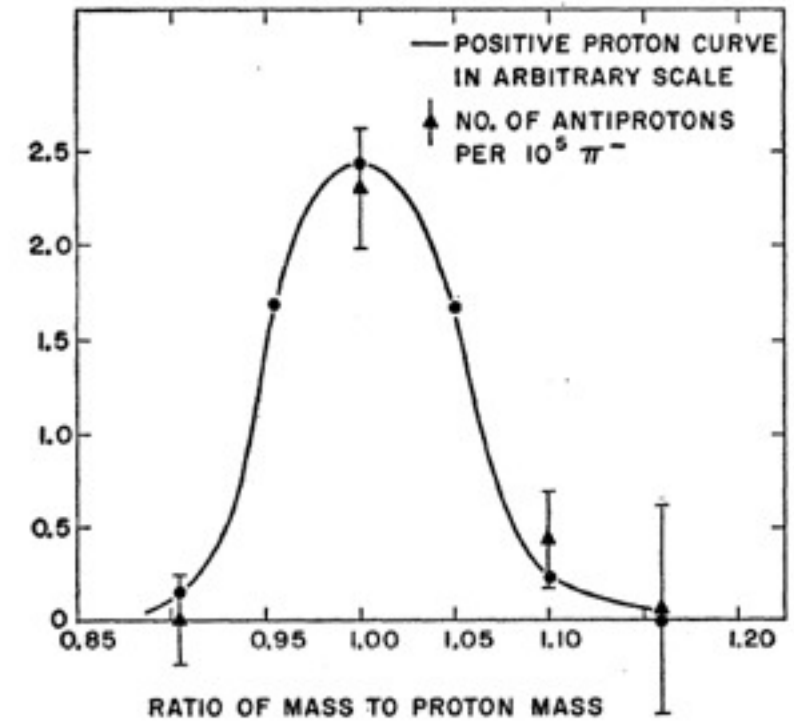
1955 - intentional production of antiprotons in an accelerator

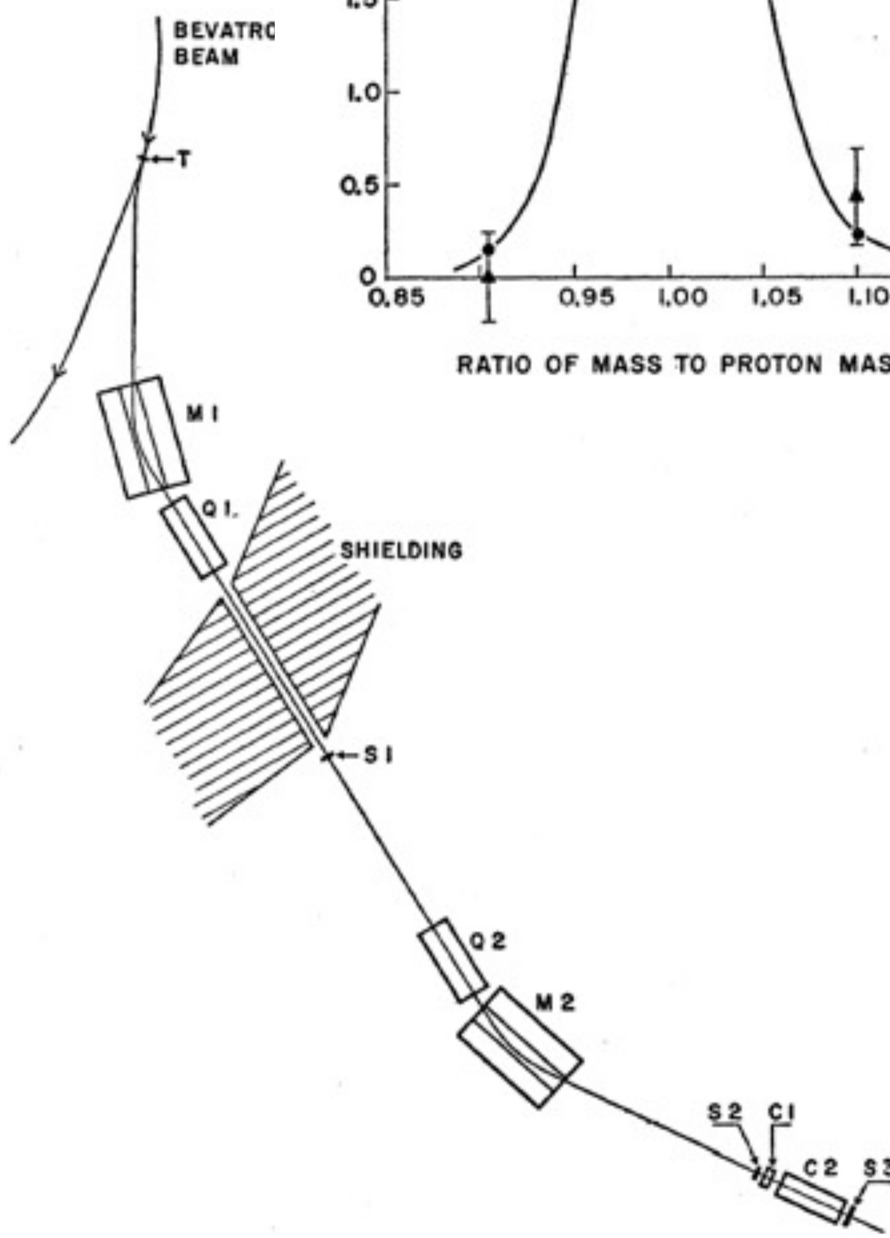


- Energy release 1350 ± 50 MeV > m_p
- Total 35 annihilations!
  - Chamberlain et al., Phys. Rev. 102, 902 (1956)
- final proof of antimatter character

Discovery of the Antiproton

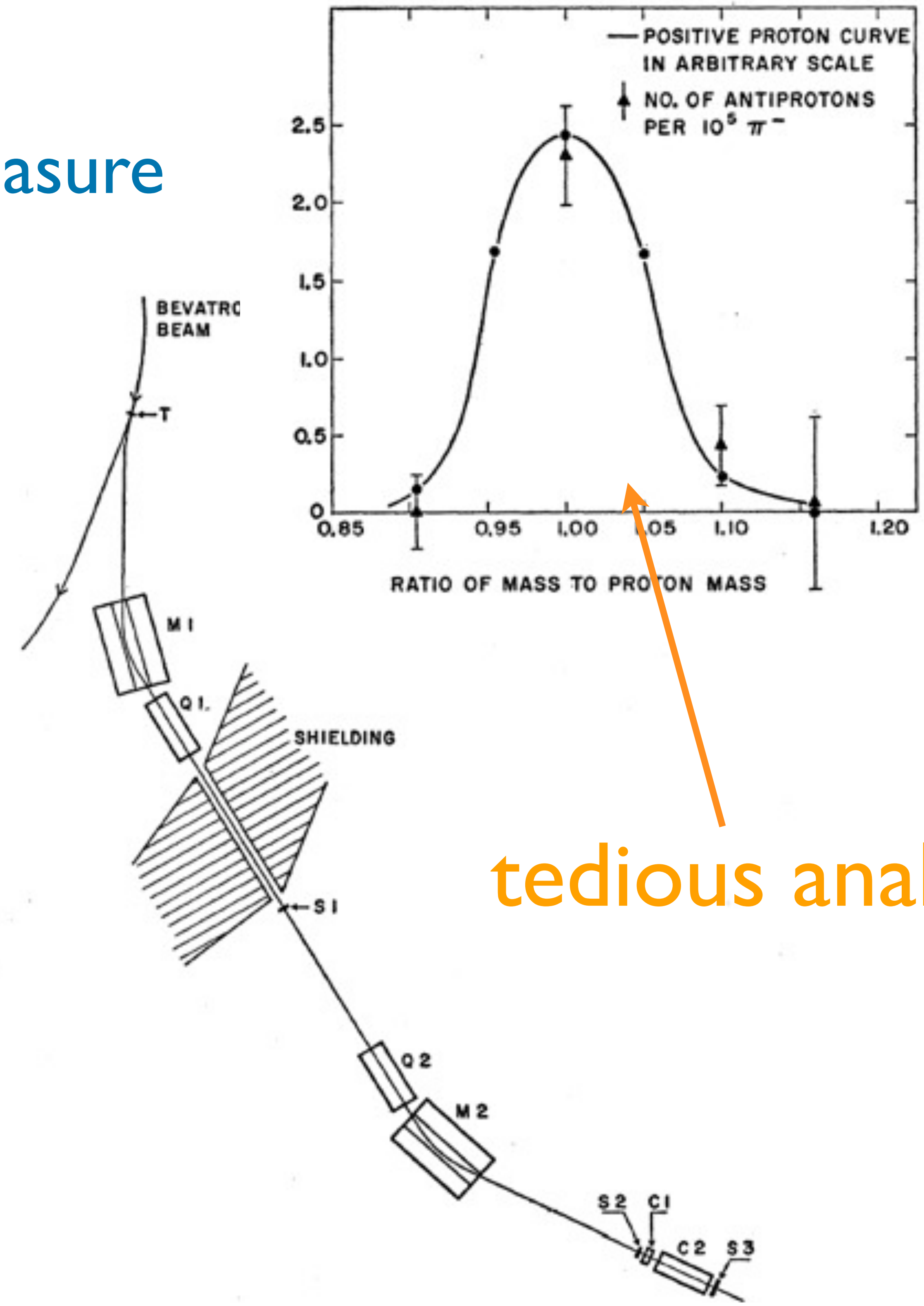
- Bevatron 5.6 GeV
  - Just at threshold!
- Discrimination against π^-: measure
  - Momentum
    - Magnets: 1.19 GeV
  - Velocity
    - TOF 51 vs. 40 ns
  - Cherenkov counter veto
- 60 events in 1955
- $\Delta m/m_p \sim 5\%$
- O. Chamberlain, E. Segre, C. Wiegand, T. Ypsilantis, Phys. Rev. 100, 947 (1955)
- Nobelprize Chamberlain & Segre 1959





Discovery of the Antiproton

- Bevatron 5.6 GeV
  - Just at threshold!
- Discrimination against π^-: measure
  - Momentum
    - Magnets: 1.19 GeV
  - Velocity
    - TOF 51 vs. 40 ns
  - Cherenkov counter veto
- 60 events in 1955
- $\Delta m/m_p$ ~ 5%
- O. Chamberlain, E. Segre, C. Wiegand, T. Ypsilantis, Phys. Rev. 100, 947 (1955)
- Nobelprize Chamberlain & Segre 1959

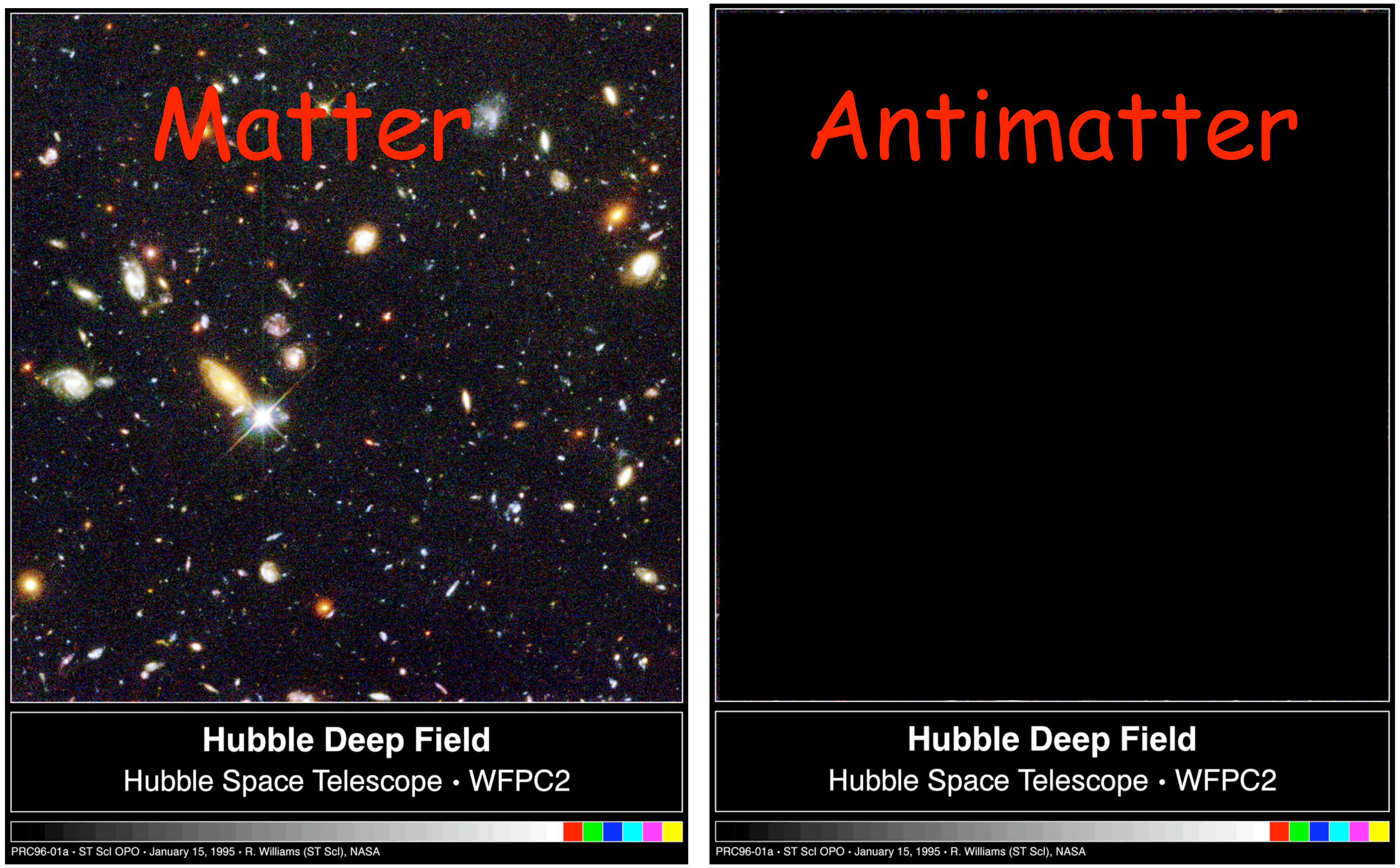


tedious analysis!

Study antimatter

Baryon asymmetry

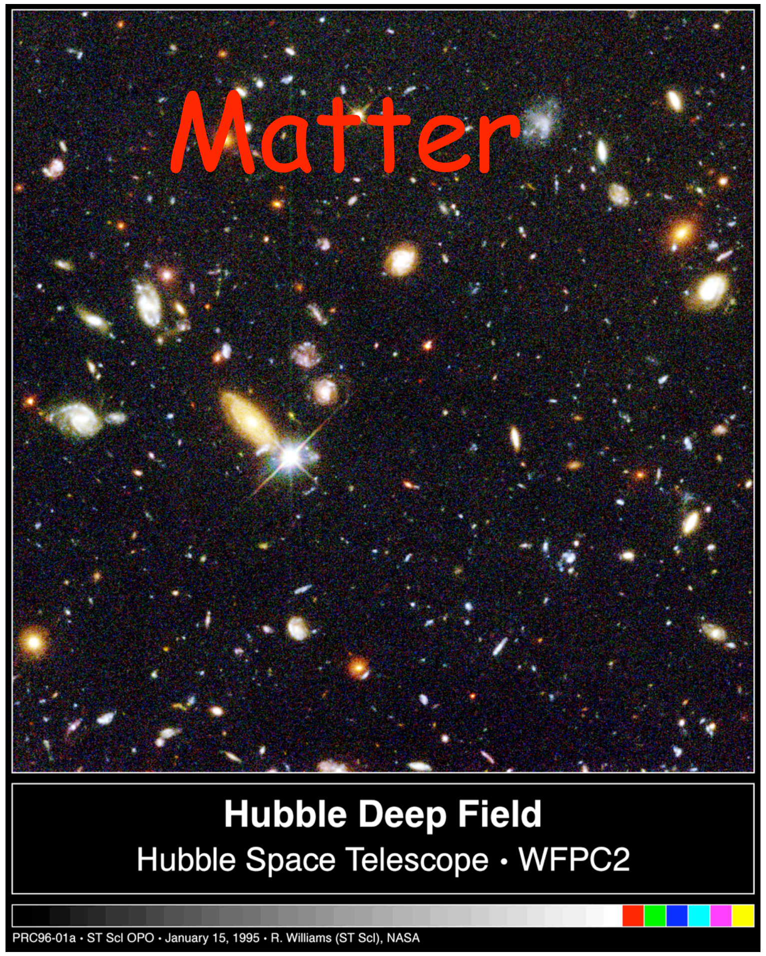
Investigate symmetries

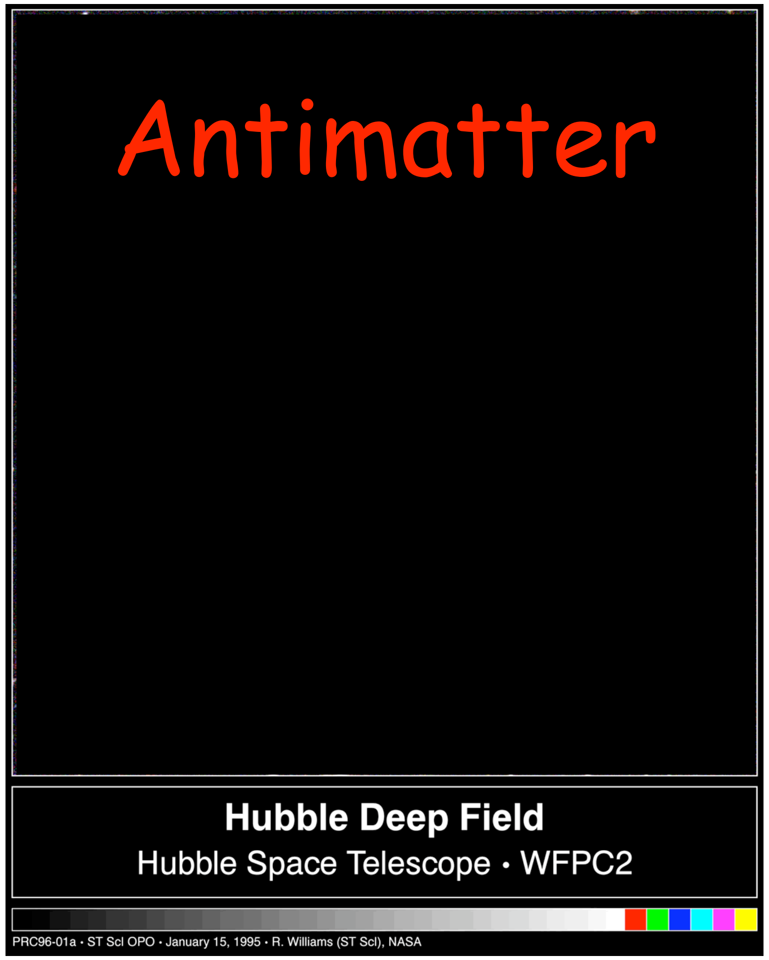


Study antimatter

Baryon asymmetry

Investigate symmetries





Use antimatter as tool

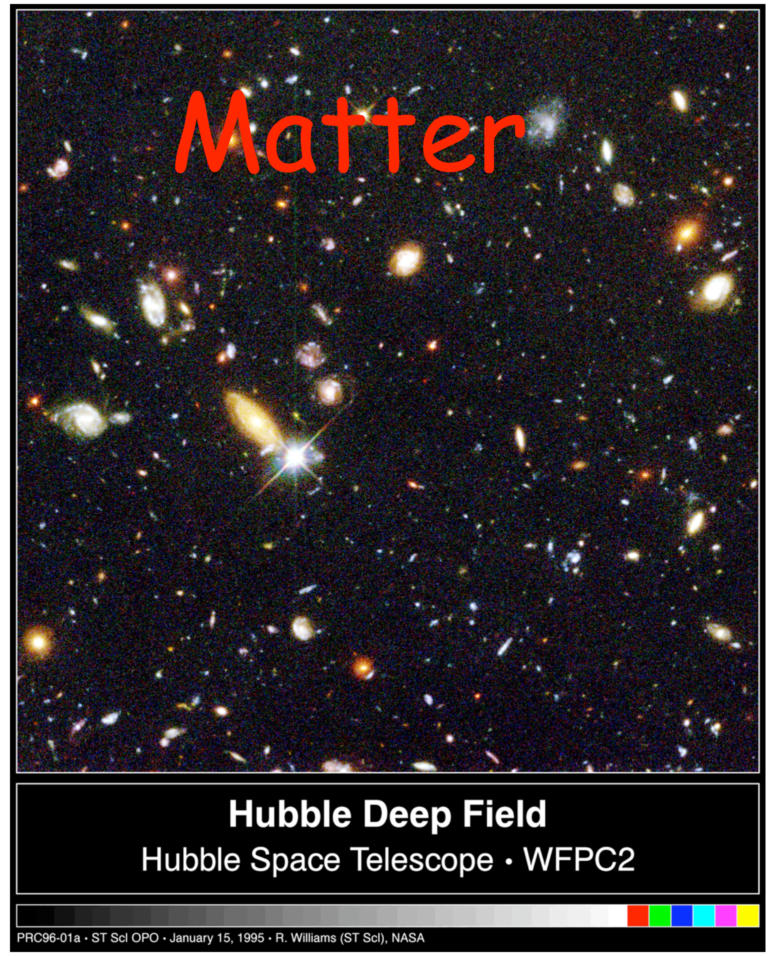
Matter-antimatter annihilation: source of new particles

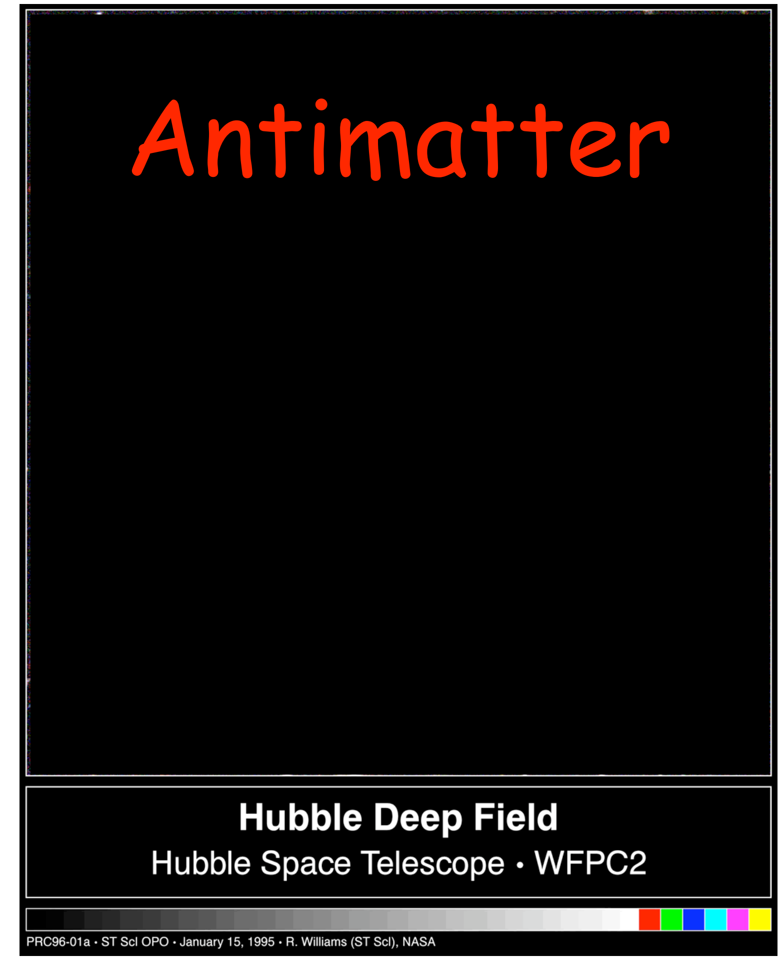
Investigate symmetries

Study antimatter

Baryon asymmetry

Investigate symmetries





Use antimatter as tool

Matter-antimatter annihilation: source of new particles

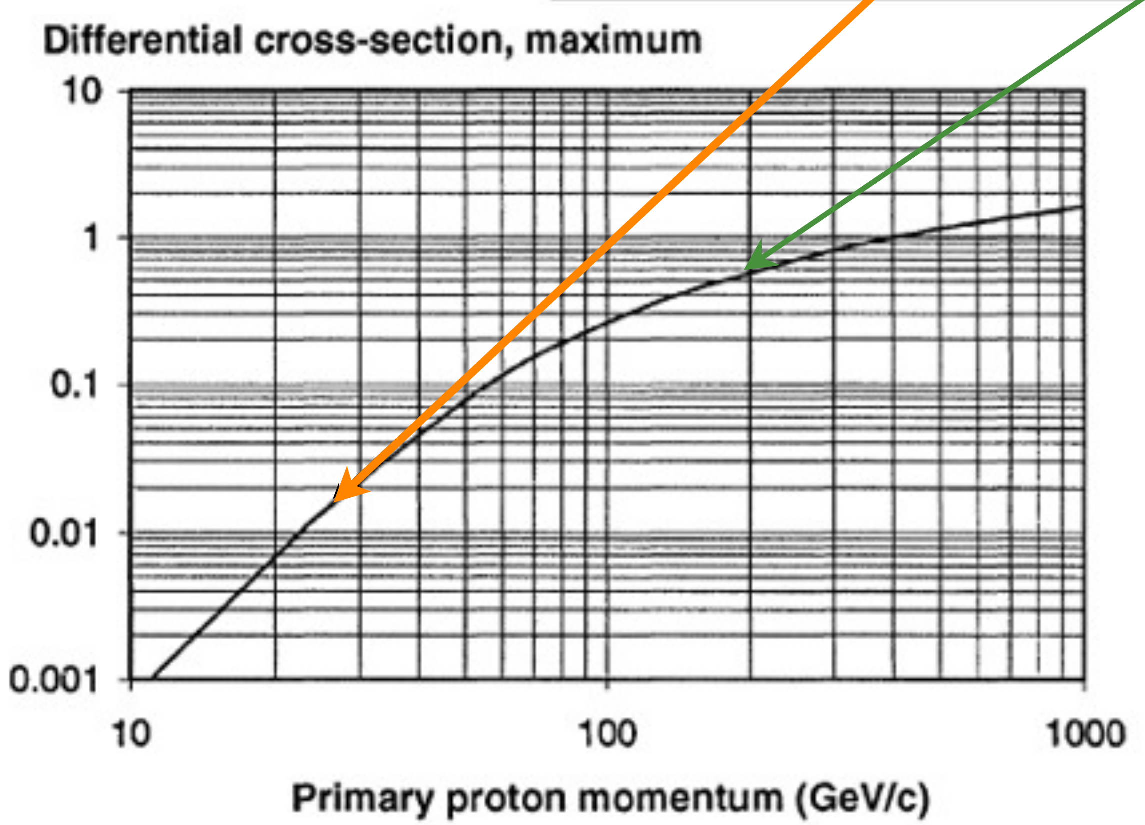
Investigate symmetries

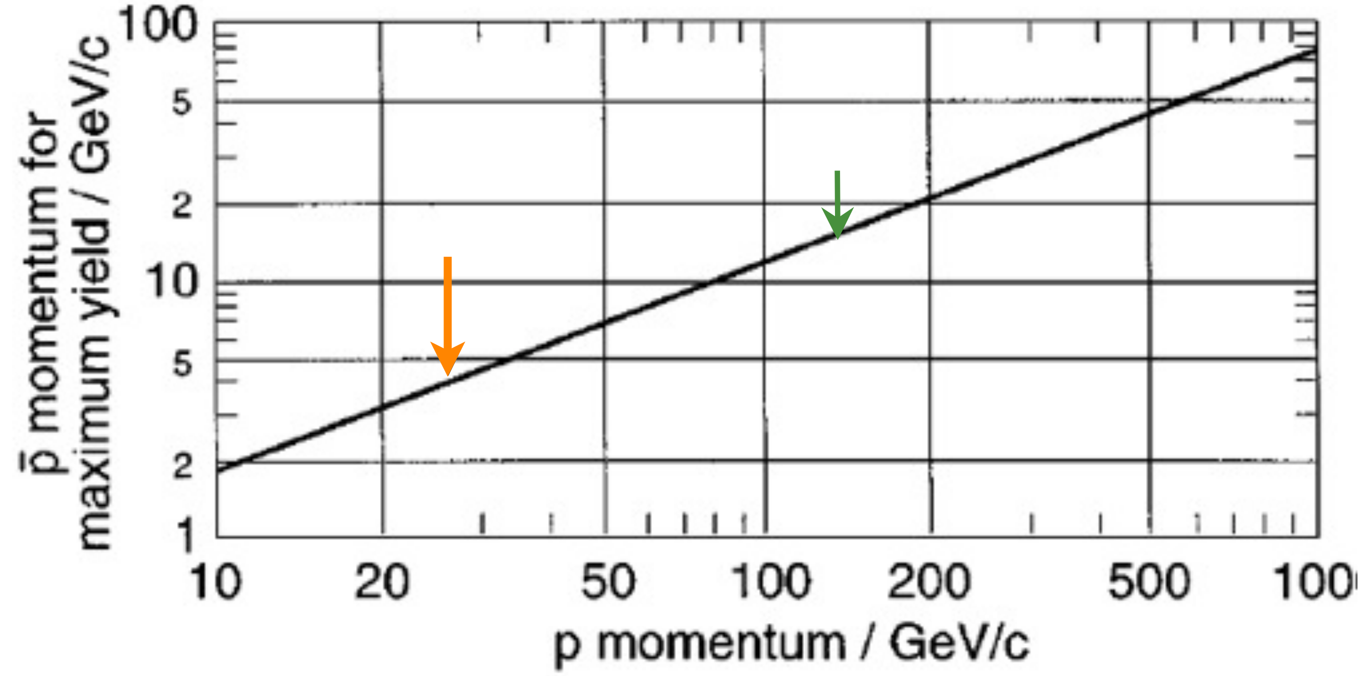
need to make it, though...

Production Energy $pN \rightarrow pXp\bar{p}$

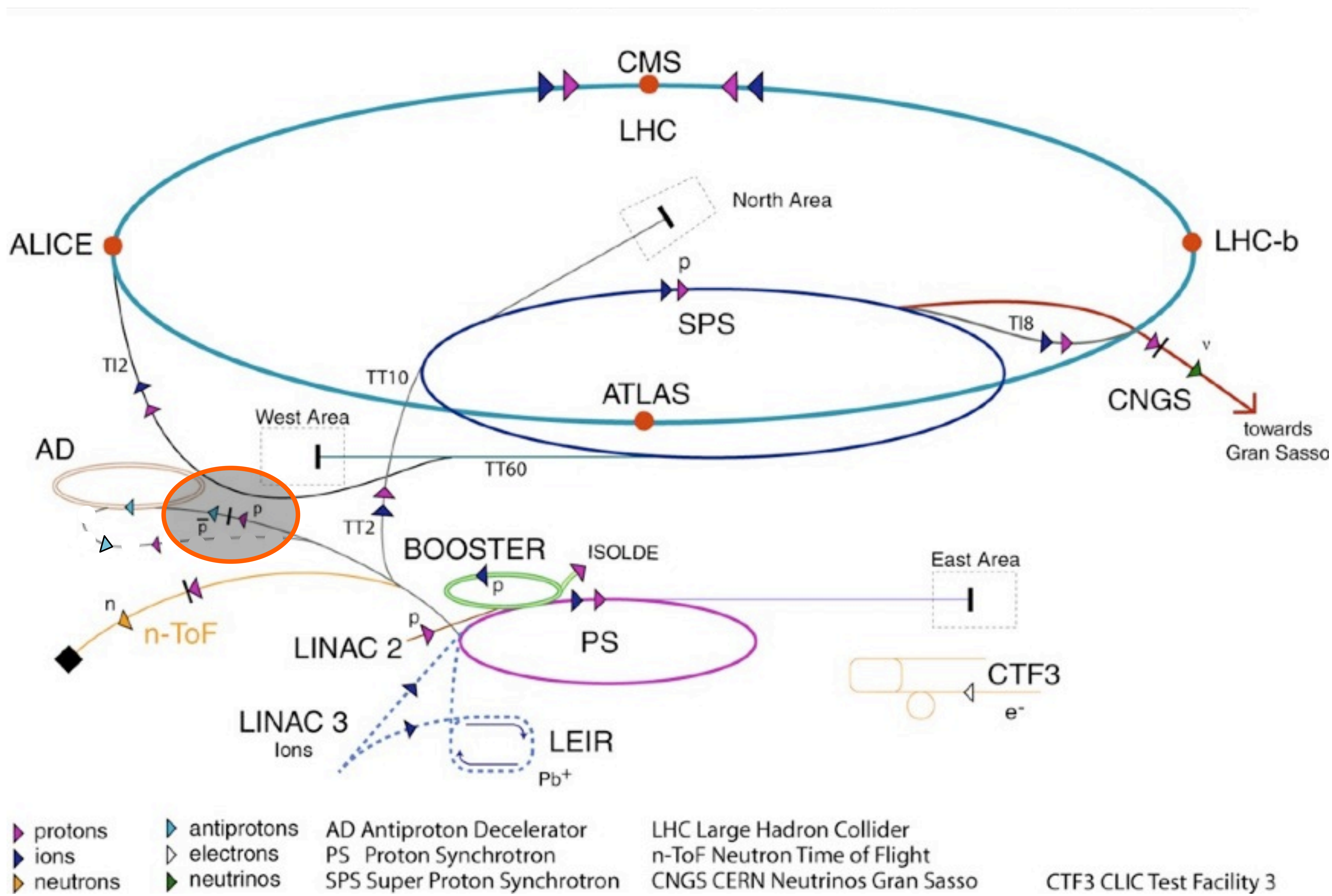
TABLE II. Comparison of CERN and Fermilab antiproton sources: for Fermilab the upgrading program quoted in Church and Marriner (1993) has been anticipated; for CERN the measured yield with magnetic horn has been used.

Machine		CERN Antiproton Collector	Fermilab debuncher
Production momentum (GeV/c)		26	120
Collection momentum (GeV/c)		3.5	9
$\bar{p}$/sr/GeV/c/Interacting p		0.013	0.25
Acceptances	A_h (π mm mrad)	200	25
	A_v (π mm mrad)	200	25
	$\Delta p/p \times 10^{-3}$	60	40
$\sqrt{A_h A_v} \times \Delta p/p$ (π mm mrad$\times 10^{-3}$)		12×10^3	10^3
Yield ($\bar{p}/p$)		3.5×10^{-6}	14×10^{-6}
Protons per pulse		1.5×10^{13}	0.5×10^{13}
Antiprotons per pulse		5×10^7	7×10^7





CERN Accelerator Complex



Overview:

1. Introduction and overview
2. Antimatter at high energies (SppS, LEP, Fermilab)
3. Meson spectroscopy (antimatter as QCD probe)

4. Astroparticle physics and cosmology
5. CP and CPT violation tests
6. Precision tests with Antimatter

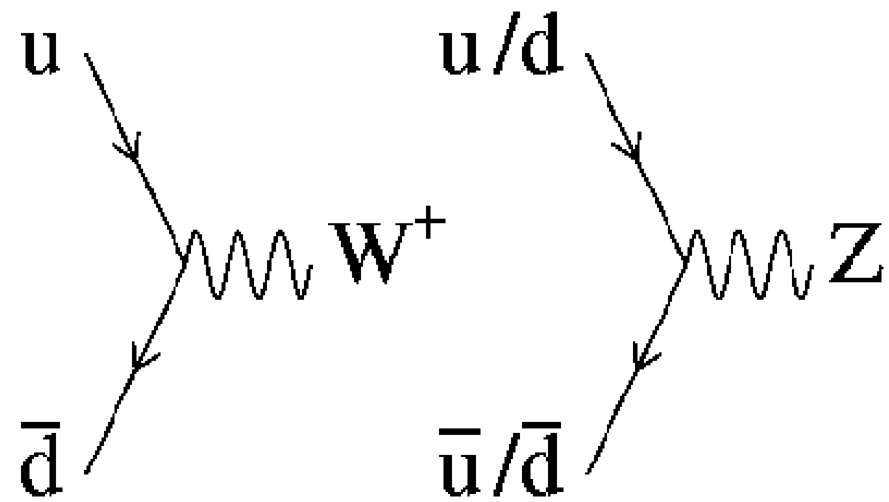
7. Precision tests with Antihydrogen
8. Applications of antimatter

Use matter and antimatter to study high energy interactions, and establish the standard model

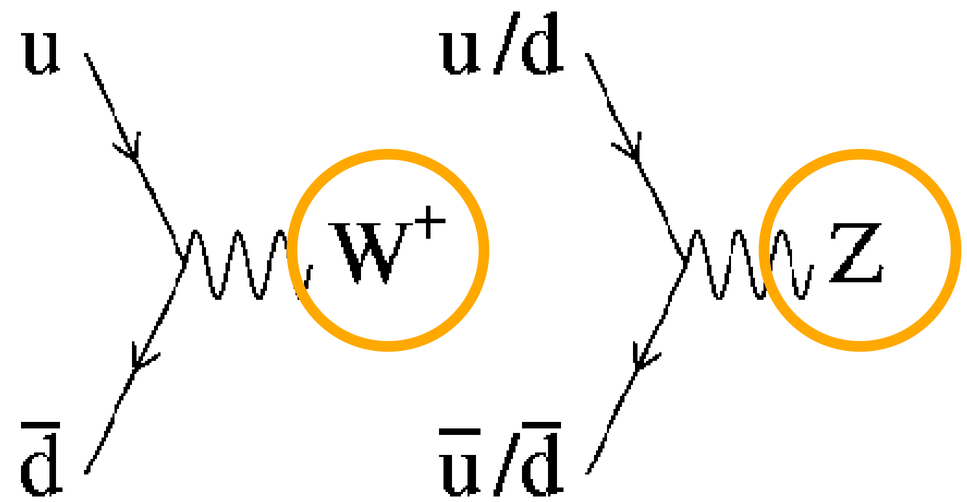
1. Proton-antiproton collisions at $Sp\bar{p}S$
2. Positron-electron interactions (at KEK, SLC, LEP)
3. Proton-antiproton interactions at Fermilab
4. Proton-antiproton for meson spectroscopy

Antimatter (+matter) is a tool to produce new particles, but it also allows to study the couplings between different particle types.

Electroweak interactions (1970's)

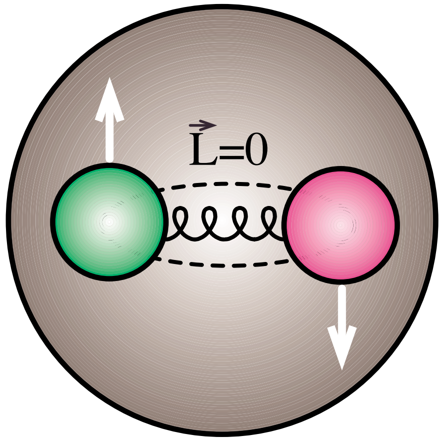


Electroweak interactions (1970's)

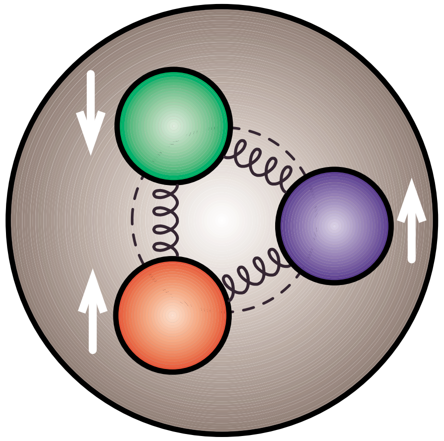


Where do we get the antiquarks from?

QCD

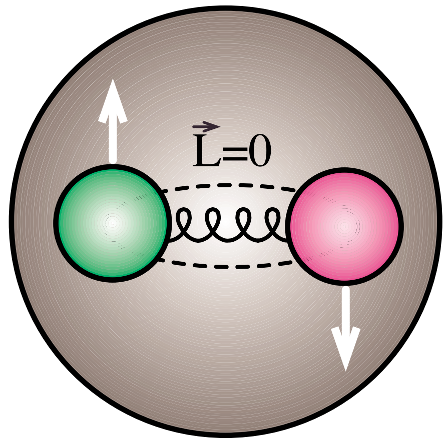


Meson ($q\bar{q}$)

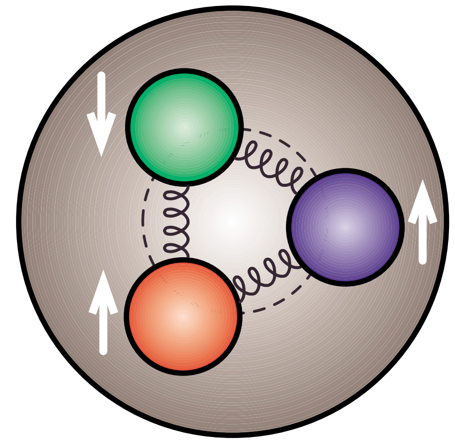


Baryon (qqq)

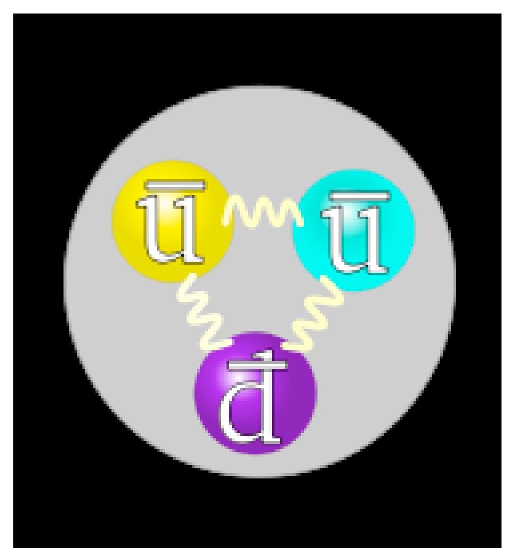
QCD



Meson ($q\bar{q}$)



Baryon (qqq)



Antibaryon ($\bar{q}\bar{q}\bar{q}$)

Collisional energy Q in parton-parton center-of-mass frame:

$$Q^2 = x_1 x_2 E^2_{cm}$$

The probability of a proton containing a parton of type i at the appropriate values of x_1 and Q^2 is given by a 'parton distribution function' (PDF), $f_i(x_1, Q^2)$

(must be measured, i.e. at H1/Zeus @ HERA)

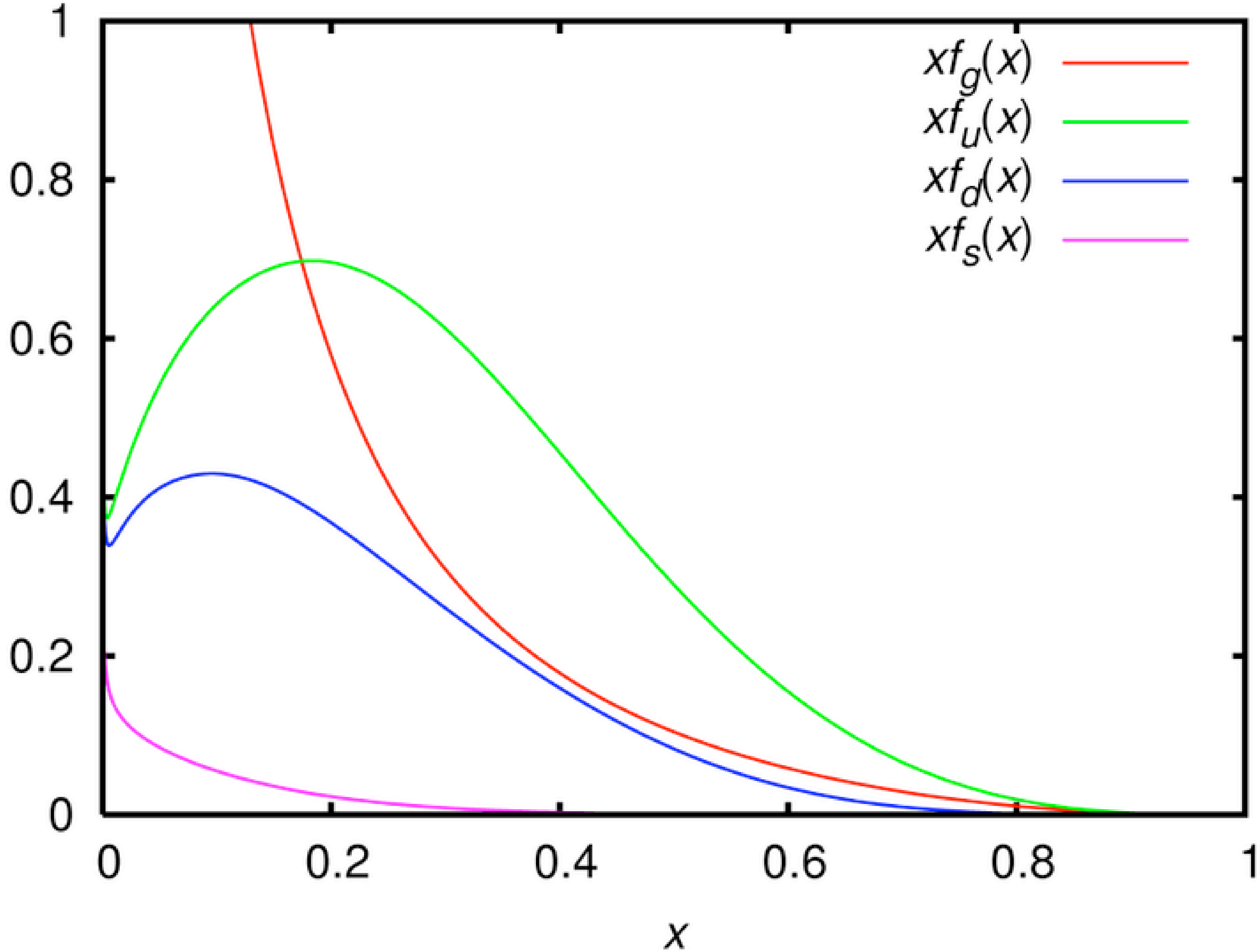
Sum over all possible combinations of incoming partons and integrate over the momentum fractions x_1 and x_2

$$\sigma = \sum_{i,j=q,\bar{q},g} \int dx_1 dx_2 \, f_i(x_1, Q^2) \cdot \bar{f}_j(x_2, Q^2) \cdot \hat{\sigma}(Q^2)$$

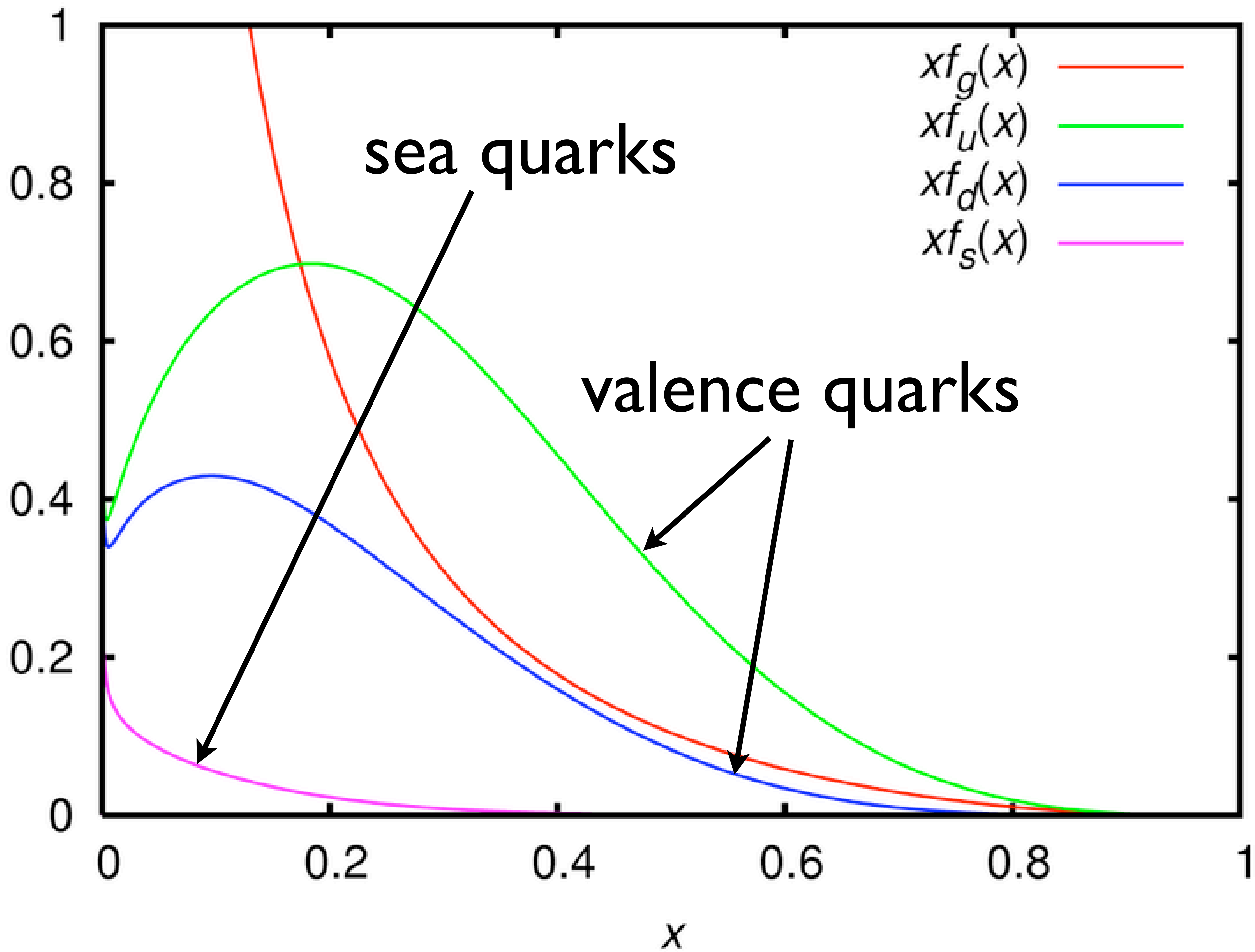
(anti)proton beam = broadband beam of (anti)partons

(initial-state partons have a high probability of radiating gluons before they collide, so not even the nominal energy is available)

Fraction of momentum carried by ...



Fraction of momentum carried by ...



The use of antiproton-proton collisions allows for
a higher average energy of collisions
between quarks and antiquarks
than would be possible in proton-proton collisions.

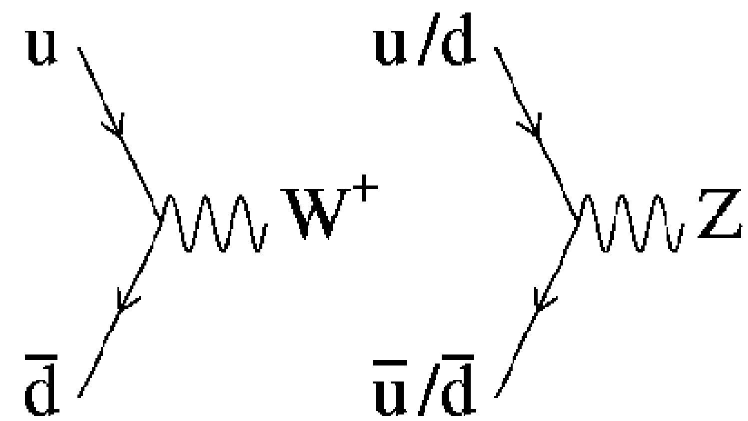
This is because the valence quarks in the proton,
and the valence antiquarks in the antiproton,
tend to carry the largest fraction of
the proton or antiproton's momentum.

The use of antiproton-proton collisions allows for a higher average energy of collisions between quarks and antiquarks than would be possible in proton-proton collisions.

This is because the valence quarks in the proton, and the valence antiquarks in the antiproton, tend to carry the largest fraction of the proton or antiproton's momentum.

= poor man’s high-energy collider

$Sp\bar{p}S$ (1980's)



valence quarks

sea quarks

requires antiprotons

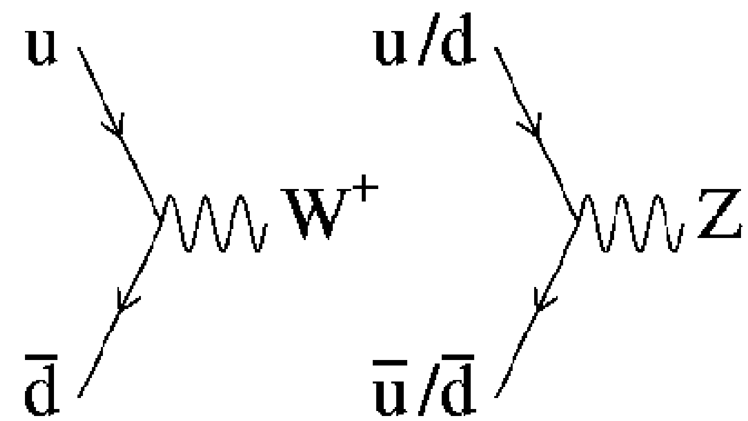
requires significantly higher energy

$$\sigma(p\bar{p} \to W^{\pm} \to e^{\pm} + \nu) \simeq 0.4 \times 10^{-33}\ k\ \mathrm{cm}^2$$

$$\sqrt{s} = 540\ \mathrm{GeV}$$

Design luminosity: 10^{30} $cm^{-2}s^{-1}$

Sp$\bar{\text{p}}$S (1980's)



valence quarks

sea quarks

requires antiprotons

requires significantly higher energy

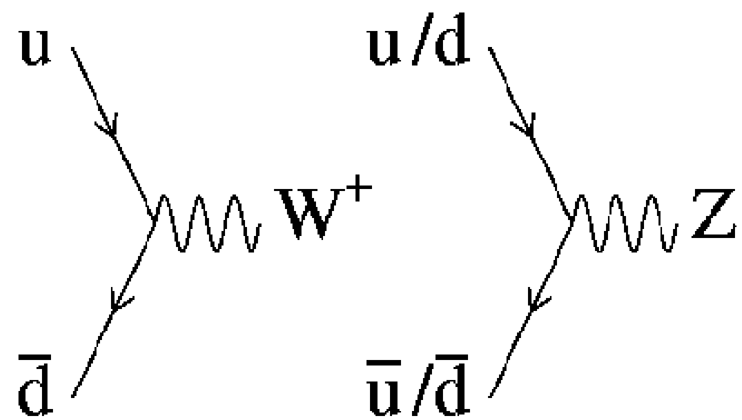
$$\sigma(p\bar{p} \to W^{\pm} \to e^{\pm} + \nu) \simeq 0.4 \times 10^{-33}\ k\ \text{cm}^2$$

$\sqrt{s}$ = 540 GeV

Design luminosity: 10^{30} $\text{cm}^{-2}\text{s}^{-1}$

We can now report successful storage of protons and antiprotons at 270 GeV with lifetimes of several hours. Typically two bunches of 5×10^{10} protons each were colliding against one bunch of about 10^9 antiprotons, giving an initial luminosity of 2×10^{25} $\text{cm}^{-2}\text{s}^{-1}$ per interaction point in these first runs.

$Sp\bar{p}S$ (1980's)



valence quarks

sea quarks

requires antiprotons

requires significantly higher energy

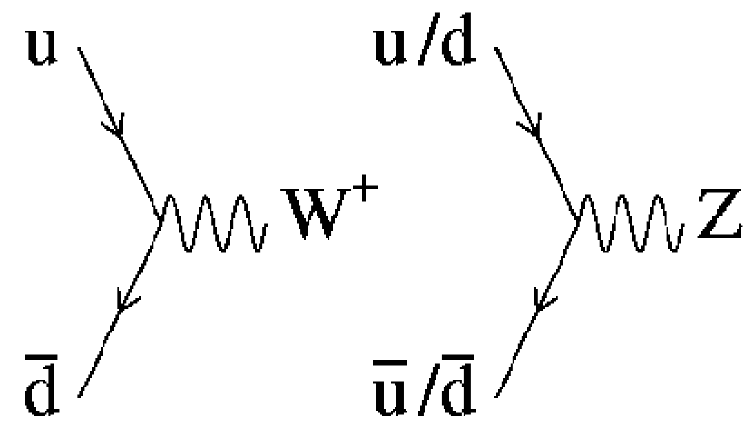
$$\sigma(p\bar{p} \to W^{\pm} \to e^{\pm} + \nu) \simeq 0.4 \times 10^{-33}\ k\ \mathrm{cm}^2$$

$$\sqrt{s} = 540\ \mathrm{GeV}$$

Design luminosity: 10^{30} $cm^{-2}s^{-1}$

We can now report successful storage of protons and antiprotons at 270 GeV with lifetimes of several hours. Typically two bunches of 5×10^{10} protons each were colliding against one bunch of about 10^9 antiprotons, giving an initial luminosity of 2×10^{25} $cm^{-2}s^{-1}$ per interaction point in these first runs.

$Sp\bar{p}S$ (1980's)



valence quarks

sea quarks

requires antiprotons

requires significantly higher energy

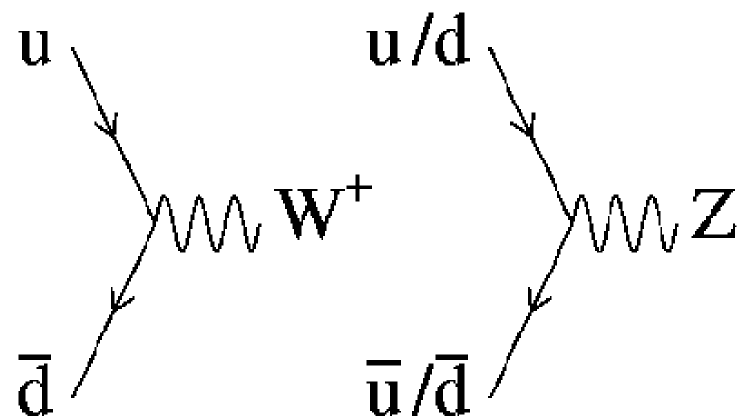
$$\sigma(p\bar{p} \rightarrow W^{\pm} \rightarrow e^{\pm} + \nu) \simeq 0.4 \times 10^{-33}\ k\ \mathrm{cm}^2$$

$\sqrt{s} = 540$ GeV

Design luminosity: 10^{30} $cm^{-2}s^{-1}$

We can now report successful storage of protons and antiprotons at 270 GeV with lifetimes of several hours. Typically two bunches of 5×10^{10} protons each were colliding against one bunch of about 10^9 antiprotons, giving an initial luminosity of 2×10^{25} $\mathrm{cm}^{-2}\mathrm{s}^{-1}$ per interaction point in these first runs.

$Sp\bar{p}S$ (1980's)



valence quarks

sea quarks

requires antiprotons

requires significantly higher energy

$$\sigma(p\bar{p} \to W^{\pm} \to e^{\pm} + \nu) \simeq 0.4 \times 10^{-33}\ k\ \mathrm{cm}^2$$

$$\sqrt{s} = 540\ \mathrm{GeV}$$

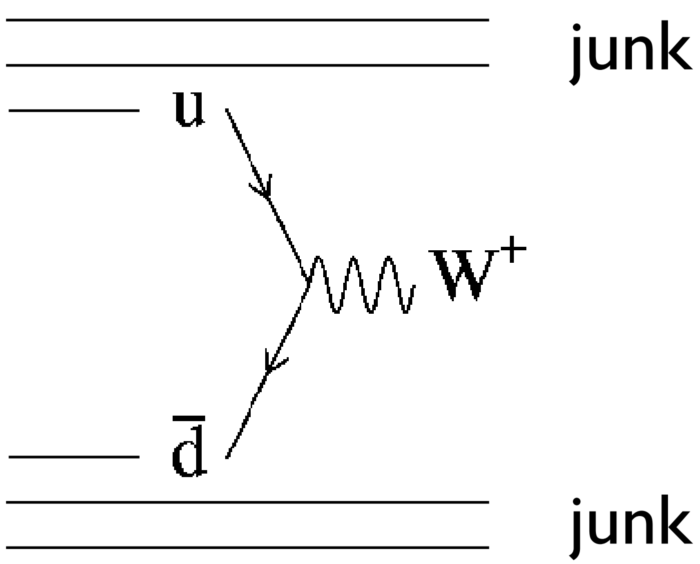
Design luminosity: 10^{30} cm^{-2}s^{-1}

We can now report successful storage of protons and antiprotons at 270 GeV with lifetimes of several hours. Typically two bunches of 5×10^{10} protons each were colliding against one bunch of about 10^{9} antiprotons, giving an initial luminosity of 2×10^{25} cm^{-2}s^{-1} per interaction point in these first runs.

UA1 results
1) how to detect W

$$\bar{p} + p \rightarrow W^{\pm} + X, \; W \rightarrow e^{\pm} + \nu;$$

- isolated large E_T electrons
- isolated large E_T neutrinos

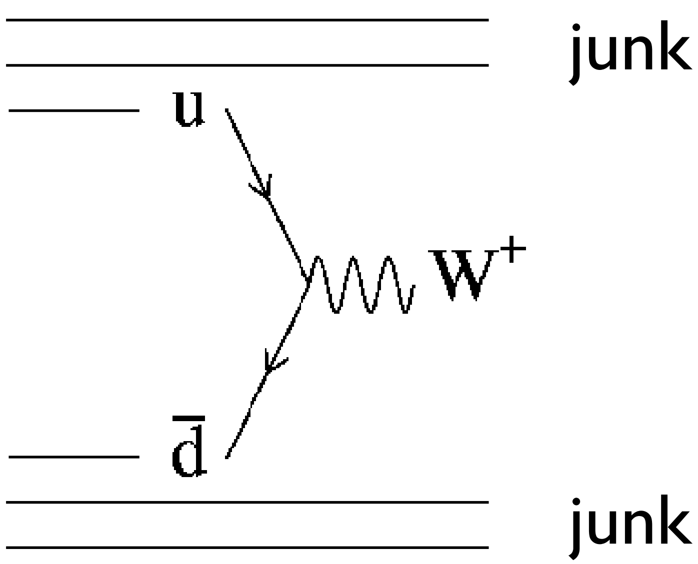


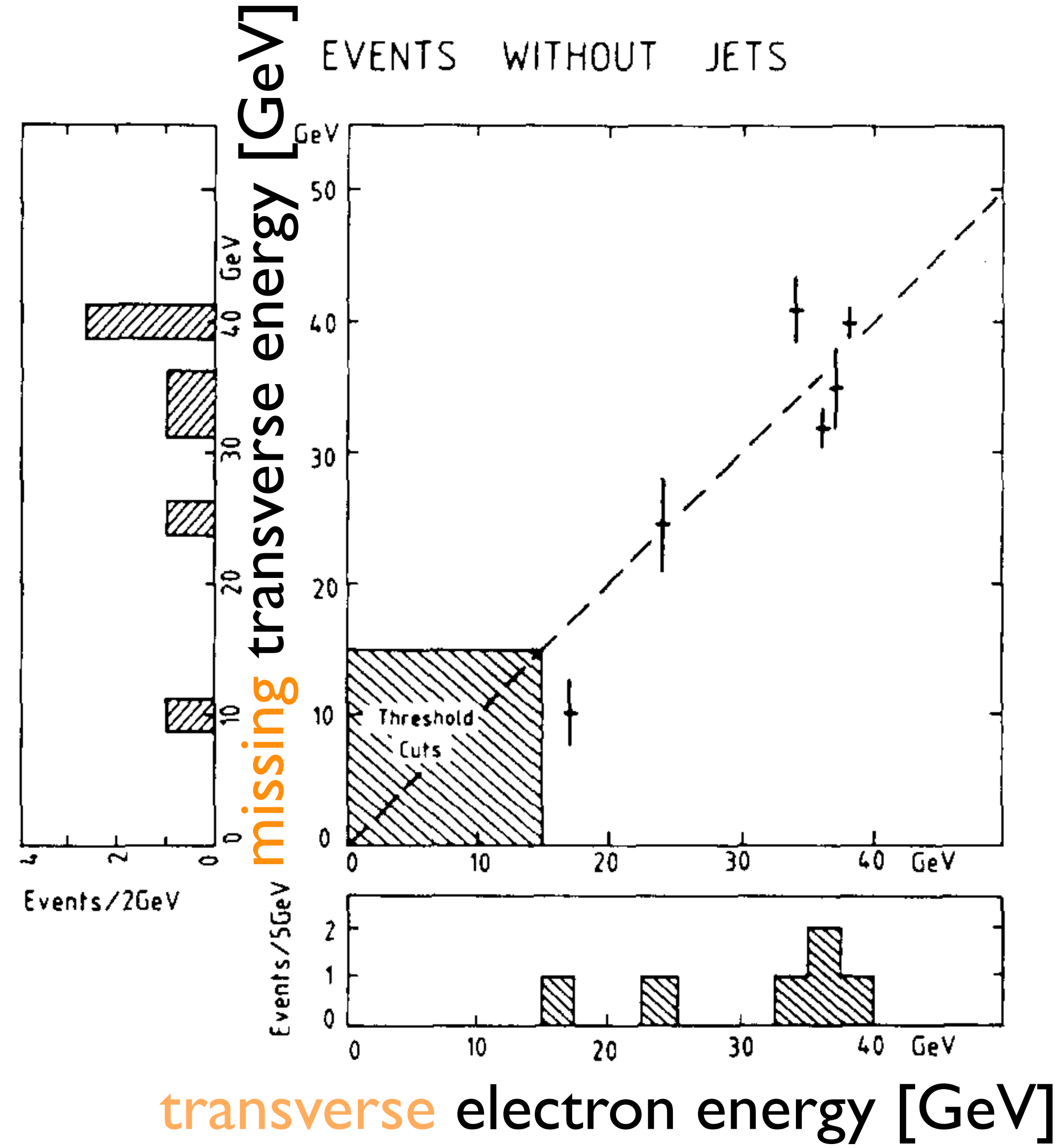
Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of isolated large transverse energy electrons with associated missing energy at s = 540 GeV. *Phys. Lett. B* **122**, 103–116 (1983)

UA1 results
1) how to detect W

$$\bar{p} + p \rightarrow W^{\pm} + X, \; W \rightarrow e^{\pm} + \nu;$$

- isolated large E_T electrons
- isolated large E_T neutrinos



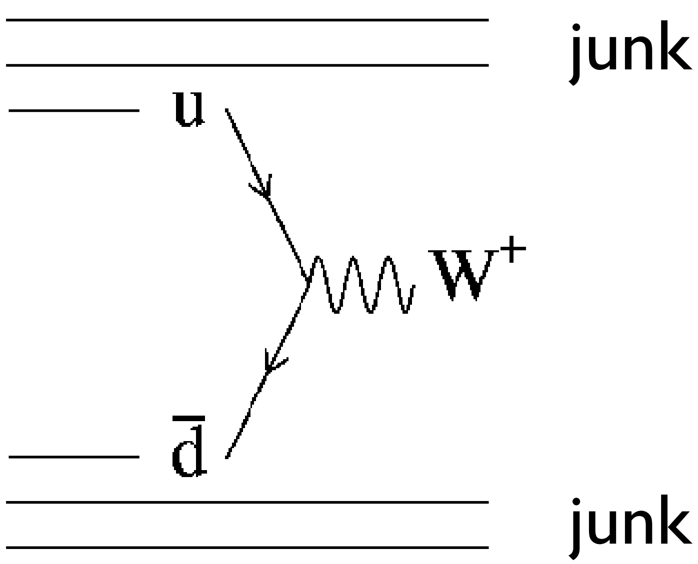


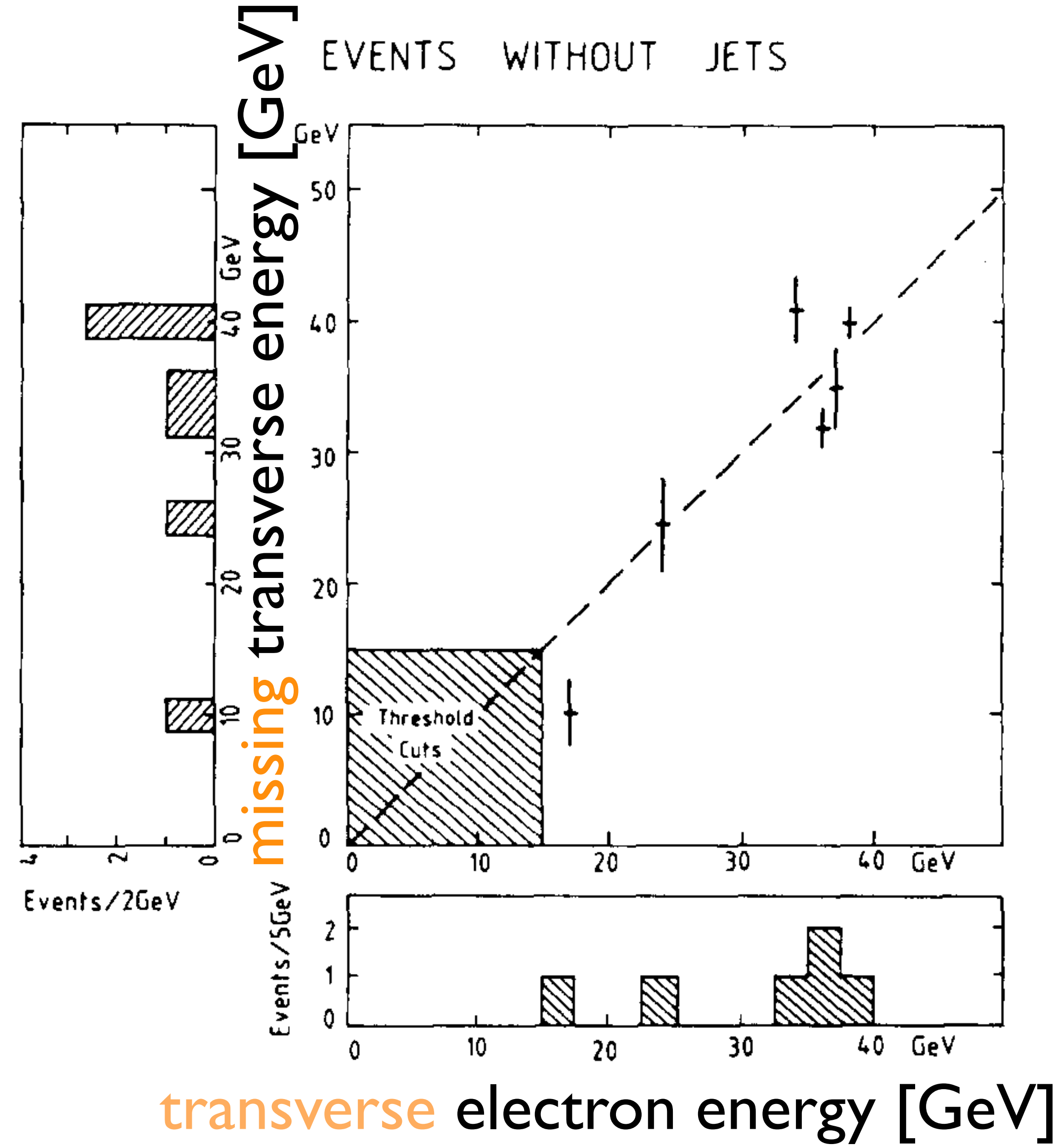
Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of isolated large transverse energy electrons with associated missing energy at s = 540 GeV. *Phys. Lett. B* **122**, 103–116 (1983)

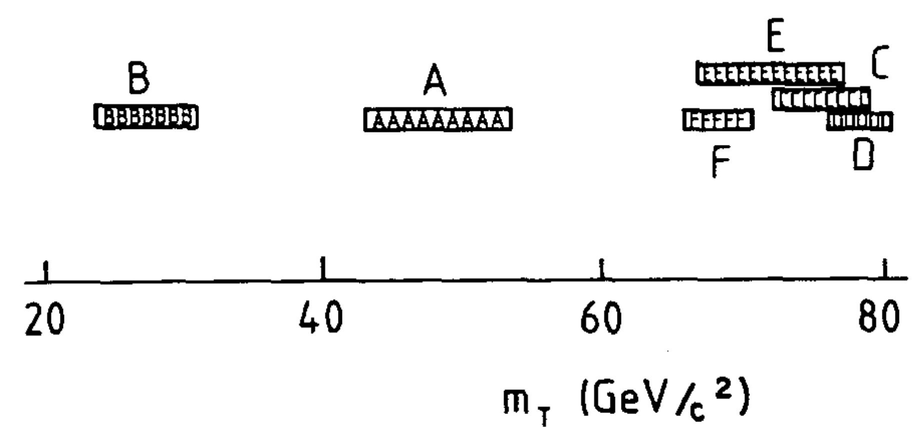
UA1 results
1) how to detect W

$$\bar{p} + p \rightarrow W^{\pm} + X, W \rightarrow e^{\pm} + \nu;$$

- isolated large E_T electrons
- isolated large E_T neutrinos





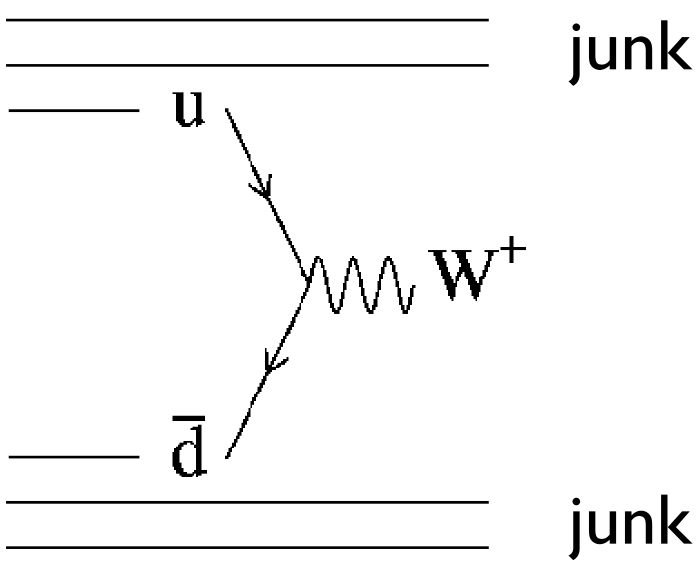


Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of isolated large transverse energy electrons with associated missing energy at s = 540 GeV. *Phys. Lett. B* **122**, 103–116 (1983)

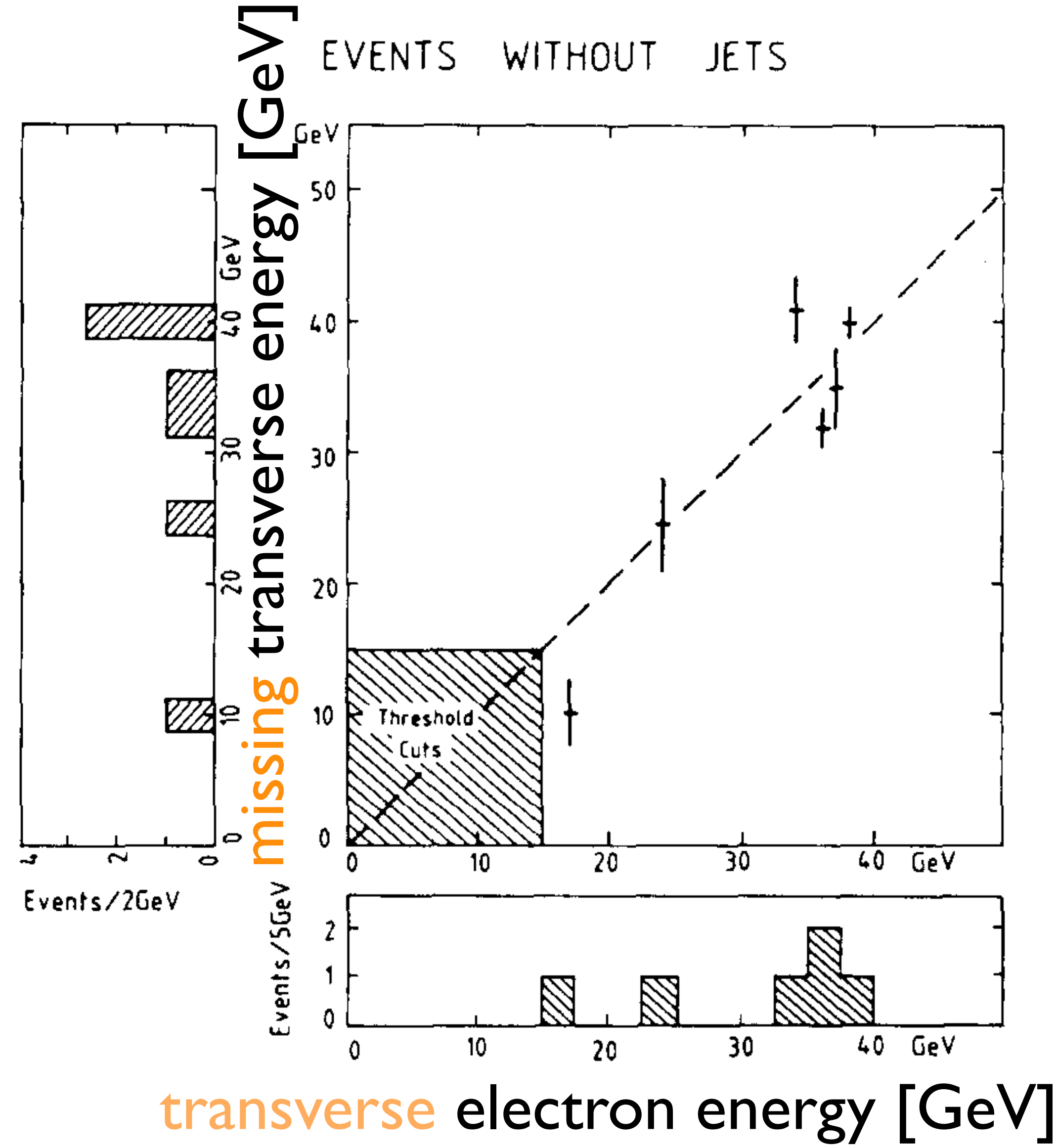
UA1 results
1) how to detect W

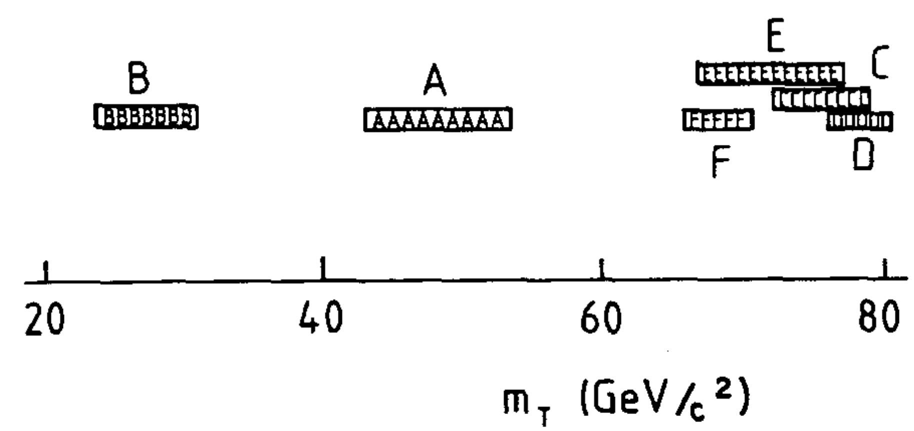
$$\bar{p} + p \to \mathbf{W^{\pm}} + \mathbf{X},\ W \to e^{\pm} + \nu;$$

- isolated large E_T electrons
- isolated large E_T neutrinos



Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of isolated large transverse energy electrons with associated missing energy at s = 540 GeV. *Phys. Lett. B* **122**, 103–116 (1983)





$$m_W = (81 {}^{+5}_{-5})\ \text{GeV}/c^2$$

Monday, July 30, 2012

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\quad \hookrightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^\pm \rightarrow e^\pm \nu$ events have been recorded [5] ‡2. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] ‡3 $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^\pm$ leptonic decay channels [9] ‡4.

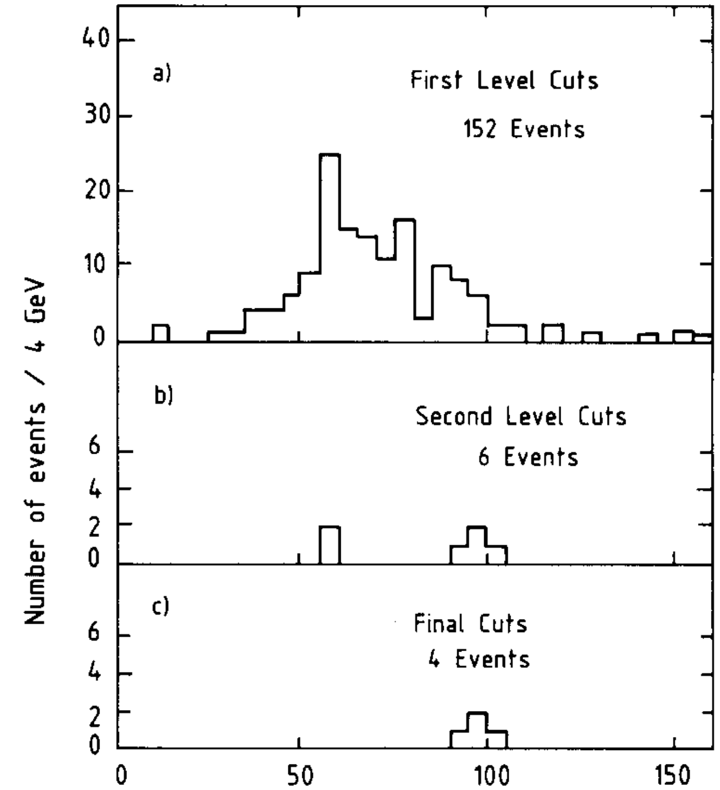
- two isolated electrons
- two isolated muons

Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/c^2 at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\hookrightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^\pm \rightarrow e^\pm \nu$ events have been recorded [5] ‡2. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] ‡3 $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^\pm$ leptonic decay channels [9] ‡4.



all

at the end of a track

"isolation"

invariant mass of two EM clusters [GeV]

- two isolated electrons
- two isolated muons

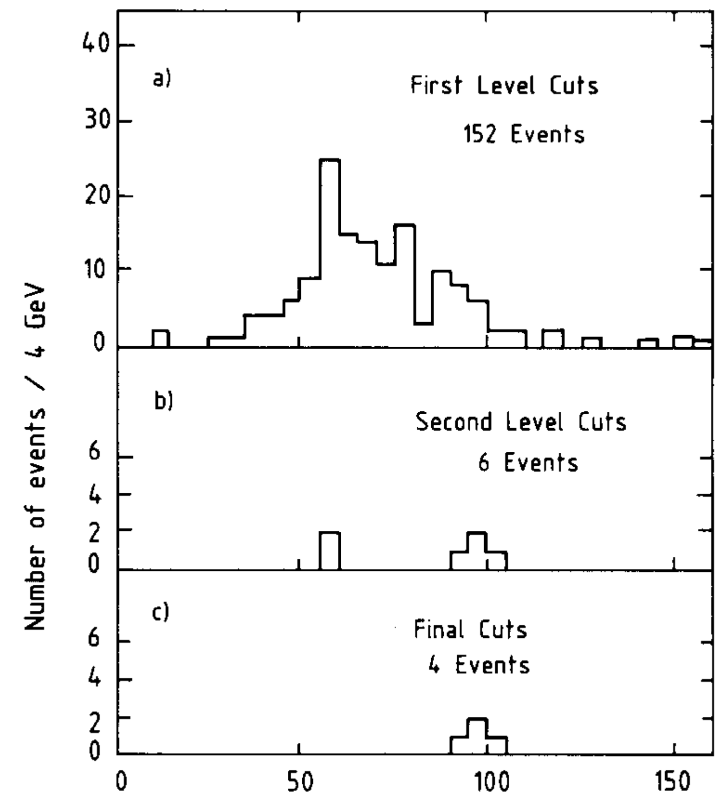
Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/*c*² at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\quad\quad\quad \hookrightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^{\pm} \rightarrow e^{\pm}\nu$ events have been recorded [5] ‡2. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] ‡3 $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^{\pm}$ leptonic decay channels [9] ‡4.

- two isolated electrons
- two isolated muons

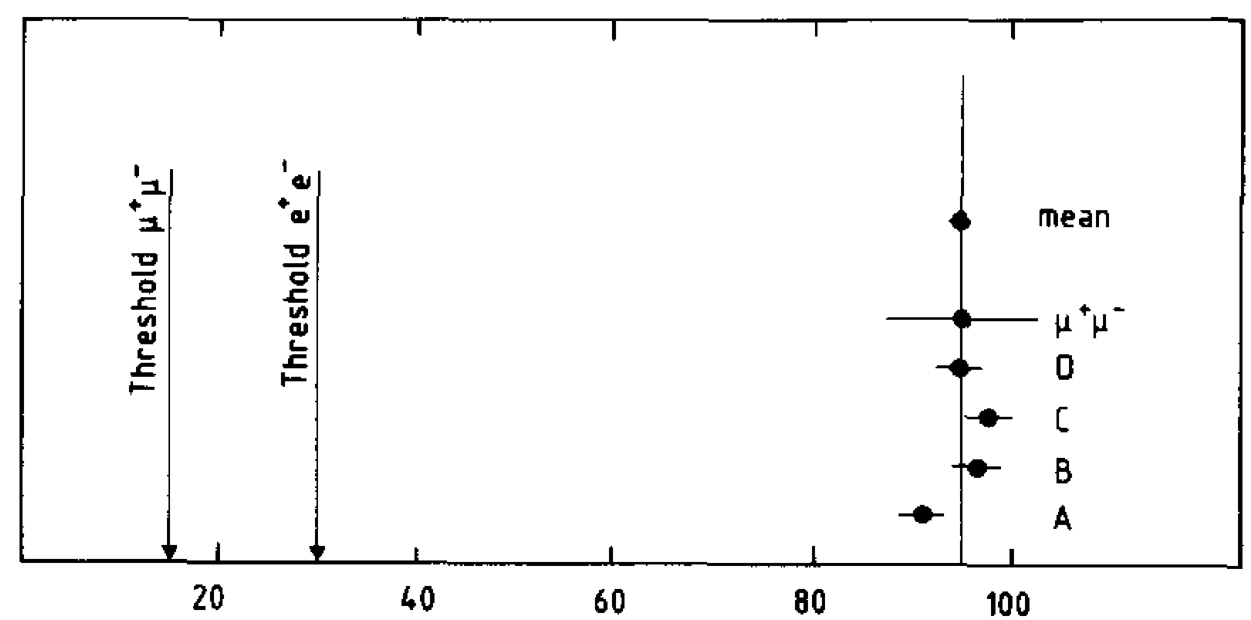


all

at the end of a track

"isolation"

invariant mass of two EM clusters [GeV]

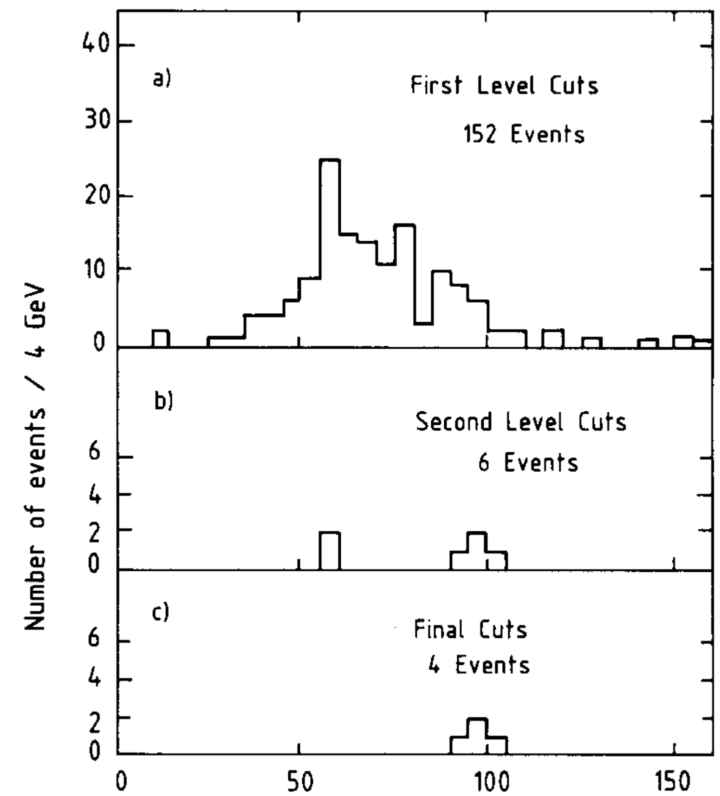


invariant mass of lepton pairs [GeV]

Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/c^2 at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\quad\quad Z^0 \rightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$



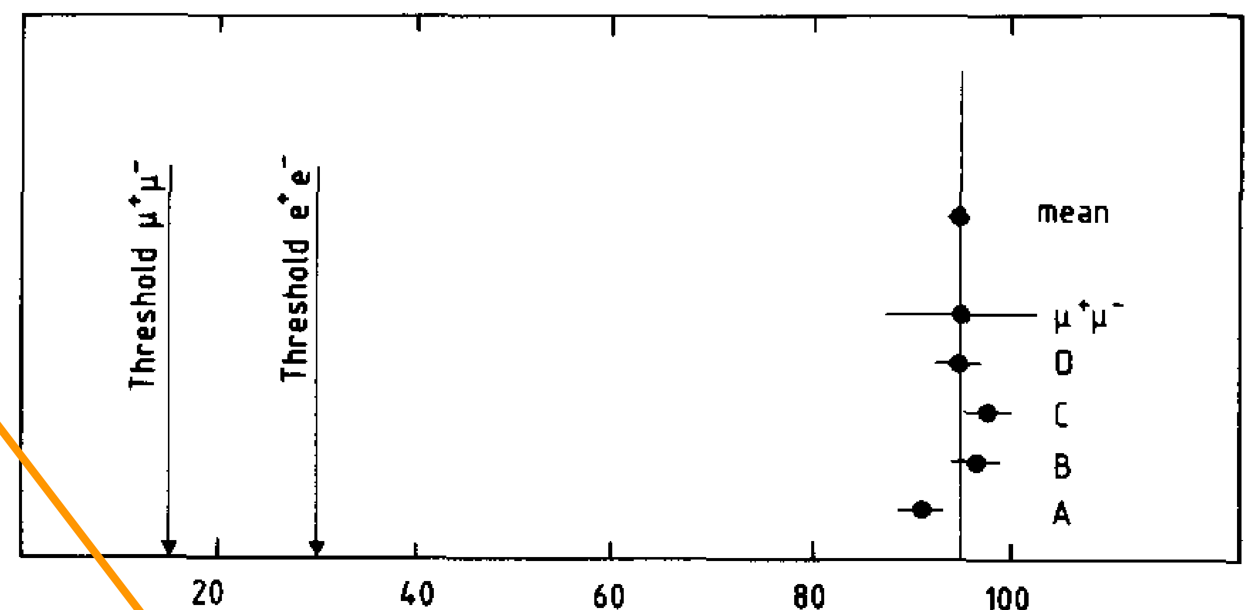
invariant mass of two EM clusters [GeV]

all

at the end of a track

“isolation”

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^\pm \rightarrow e^\pm \nu$ events have been recorded [5] [‡2]. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] [‡3] $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^\pm$ leptonic decay channels [9] [‡4].



invariant mass of lepton pairs [GeV]

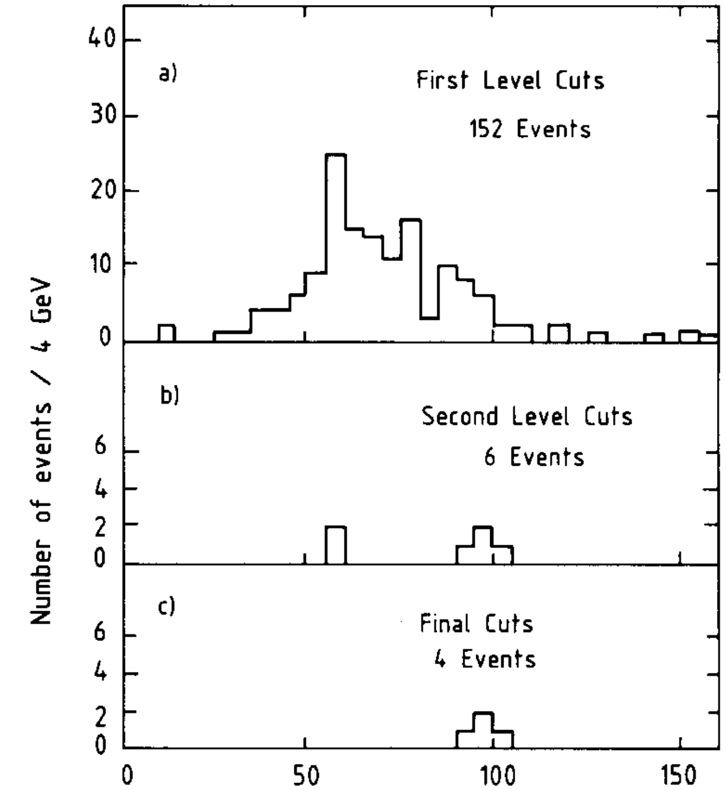
$$m_{Z^0} = (95.2 \pm 2.5)\ \mathrm{GeV}/c^2$$

- two isolated electrons
- two isolated muons

Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/*c*² at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\quad \hookrightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$



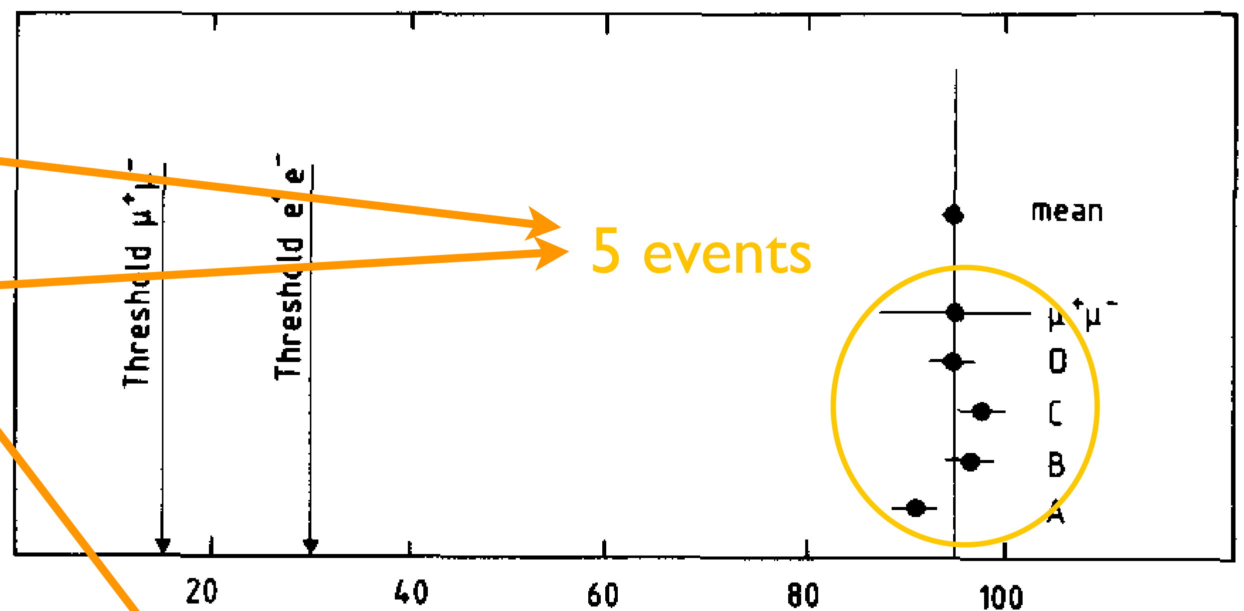
invariant mass of two EM clusters [GeV]

all

at the end of a track

"isolation"

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^\pm \rightarrow e^\pm \nu$ events have been recorded [5] [‡2]. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] [‡3] $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^\pm$ leptonic decay channels [9] [‡4].



5 events

invariant mass of lepton pairs [GeV]

$$m_{Z^0} = (95.2 \pm 2.5)\ \text{GeV}/c^2$$

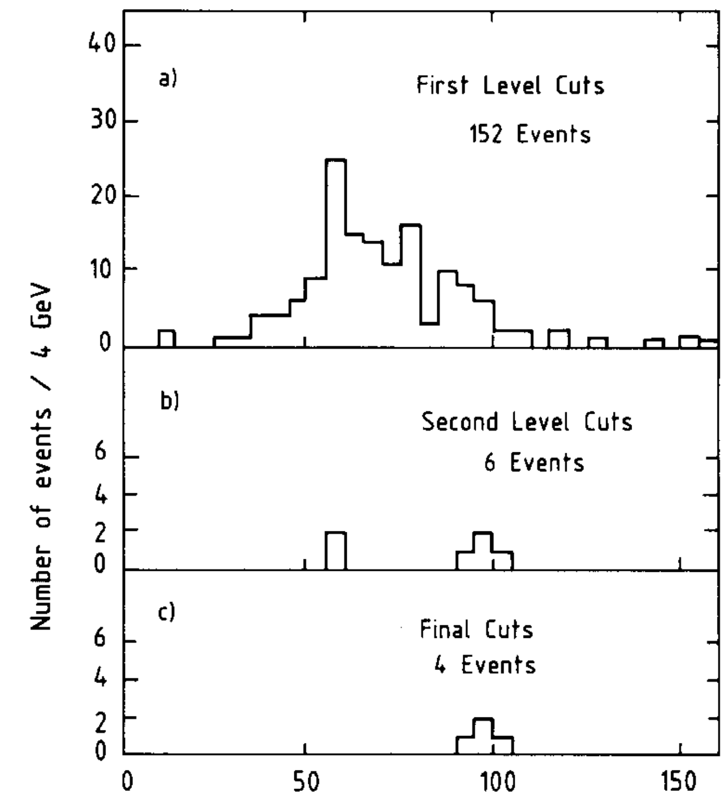
- two isolated electrons
- two isolated muons

Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/c^2 at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

UA1 results
2) how to detect Z

$$\bar{p} + p \rightarrow Z^0 + X$$
$$\hookrightarrow e^+ + e^- \quad \text{or} \quad \mu^+ + \mu^-$$

4 events — 1 event



all

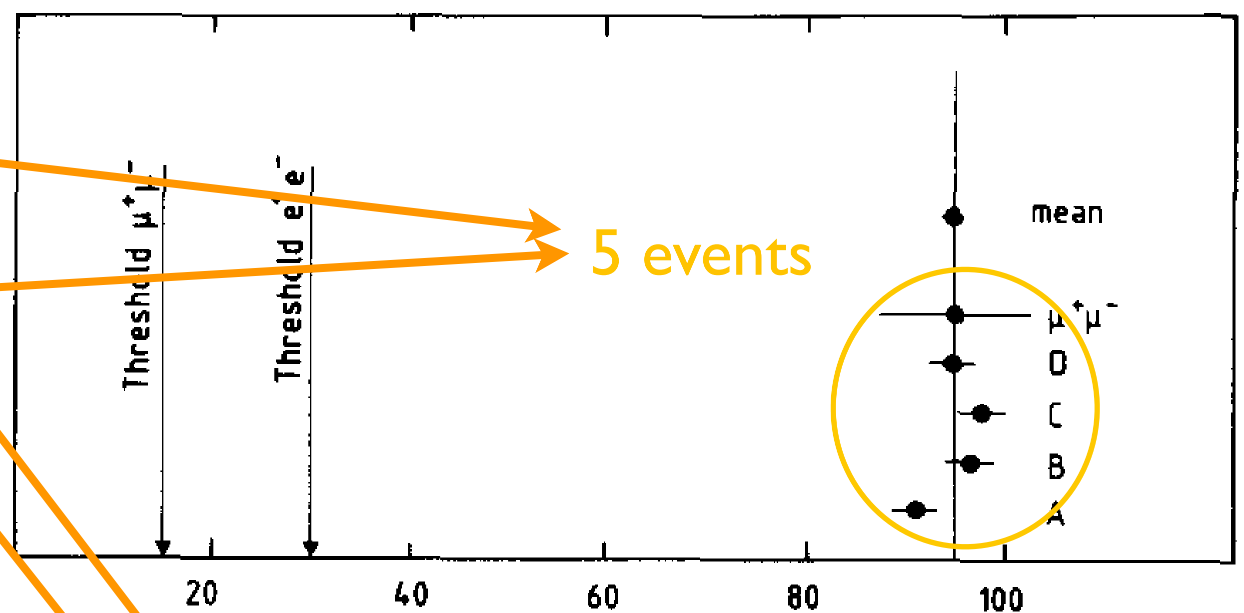
at the end of a track

"isolation"

invariant mass of two EM clusters [GeV]

The paper is based on an early analysis of a sample of collisions with an integrated luminosity of 55 nb^{-1}. In this event sample, 27 $W^\pm \rightarrow e^\pm \nu$ events have been recorded [5] ‡2. According to minimal SU(2) × U(1), the Z^0 mass is predicted to be [6] ‡3 $m_{Z^0} = 94 \pm 2.5$ GeV/c^2. The reaction (1) is then approximately a factor of 10 less frequent than the corresponding $W^\pm$ leptonic decay channels [9] ‡4.

5 events



invariant mass of lepton pairs [GeV]

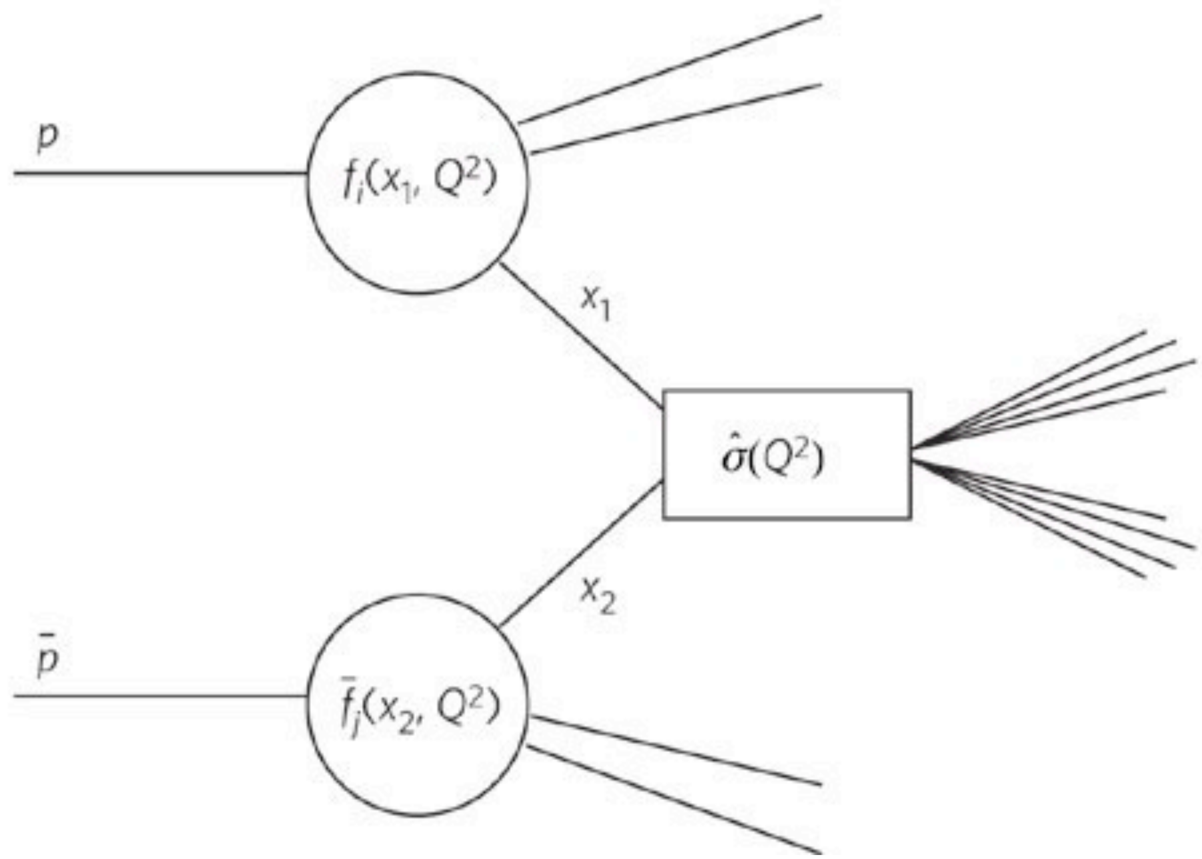
$$m_{Z^0} = (95.2 \pm 2.5)\ \text{GeV}/c^2$$

lepton universality

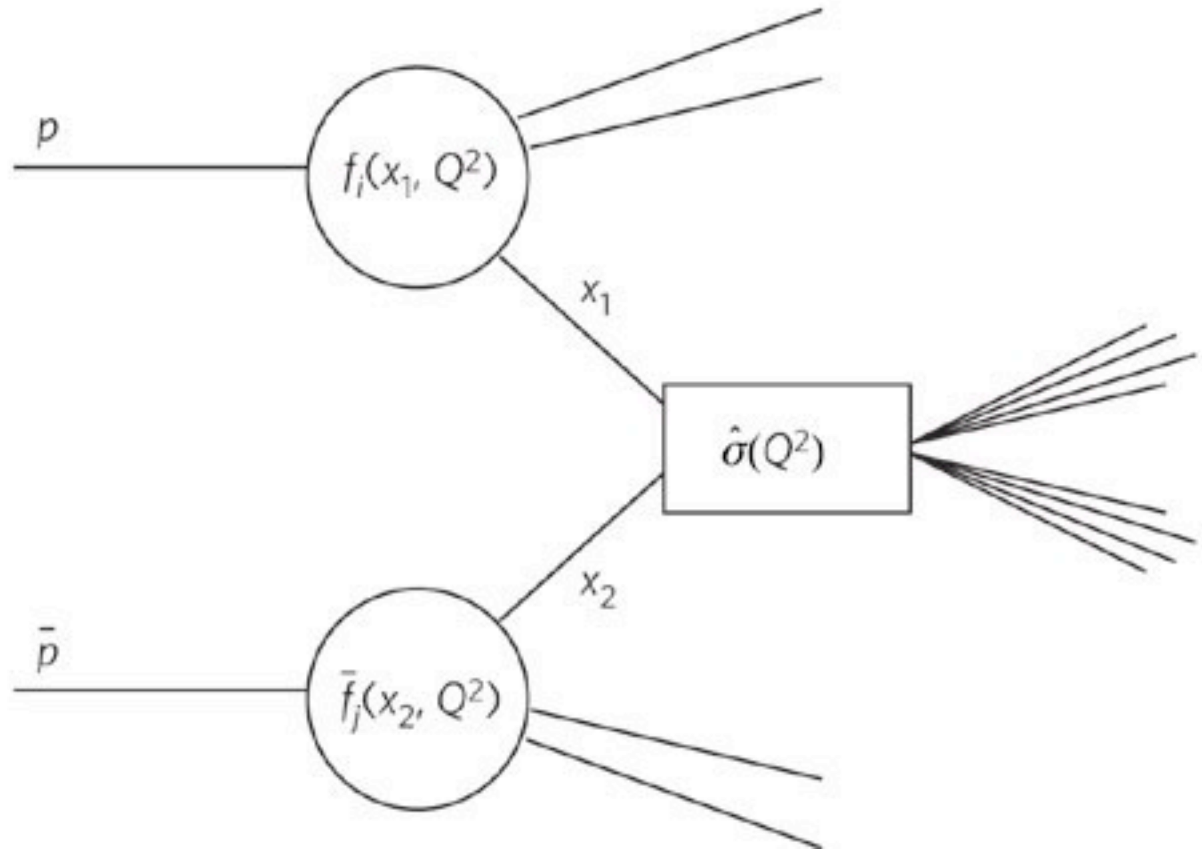
- two isolated electrons
- two isolated muons

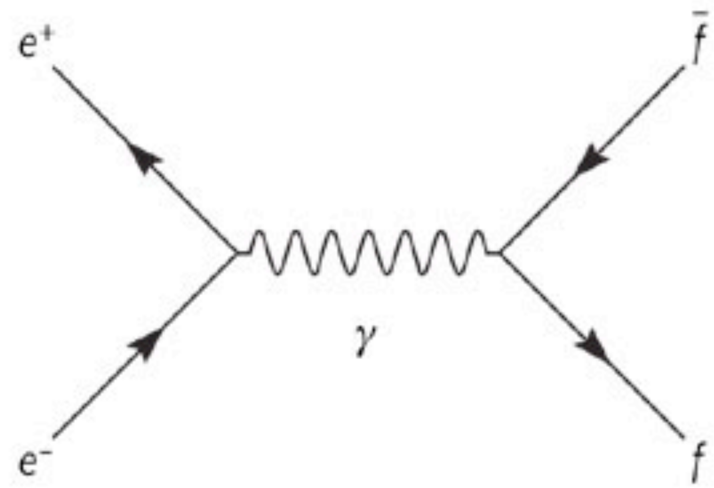
Arnison, G. *et al*. (UA1 Collaboration). Experimental observation of lepton pairs of invariant mass around 95 GeV/c^2 at the CERN SPS collider. *Phys. Lett. B* **126**, 398–410 (1983).

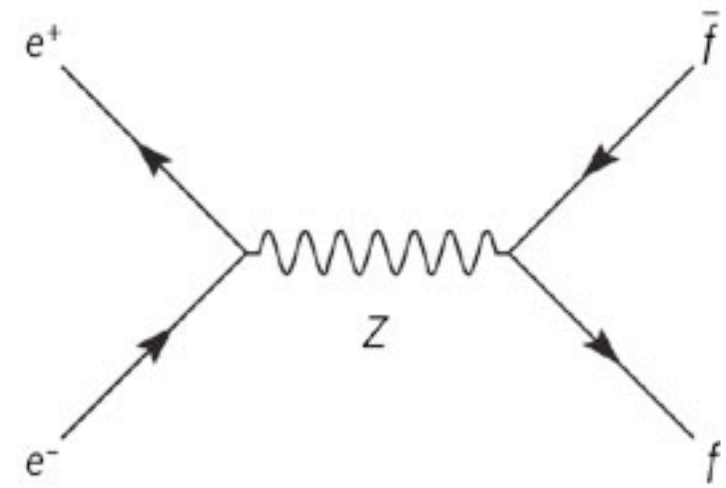
Comparing $\bar{p}p$ with e^+e^-



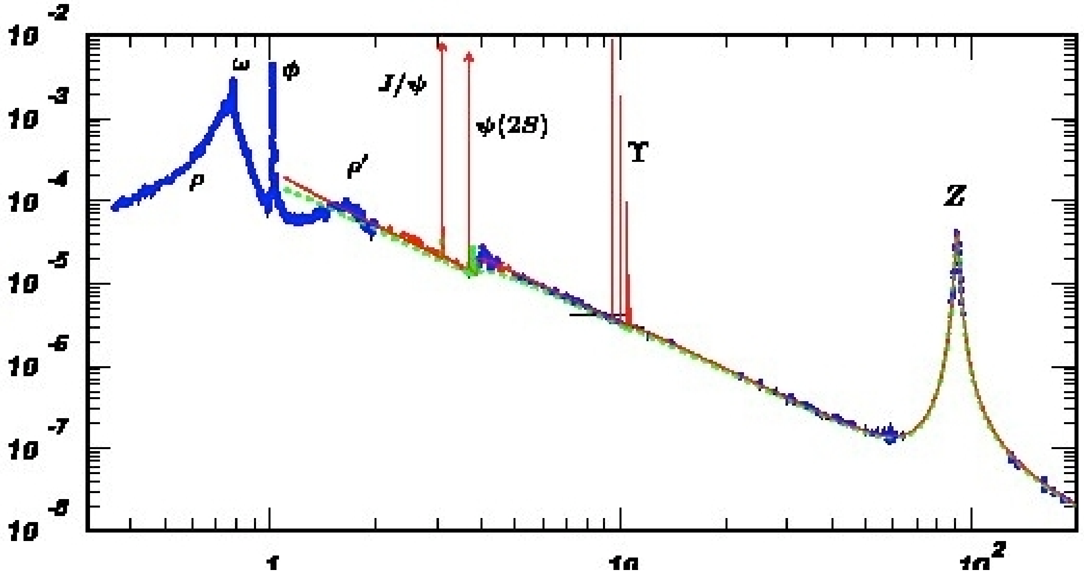
Comparing $\bar{p}p$ with e^+e^-



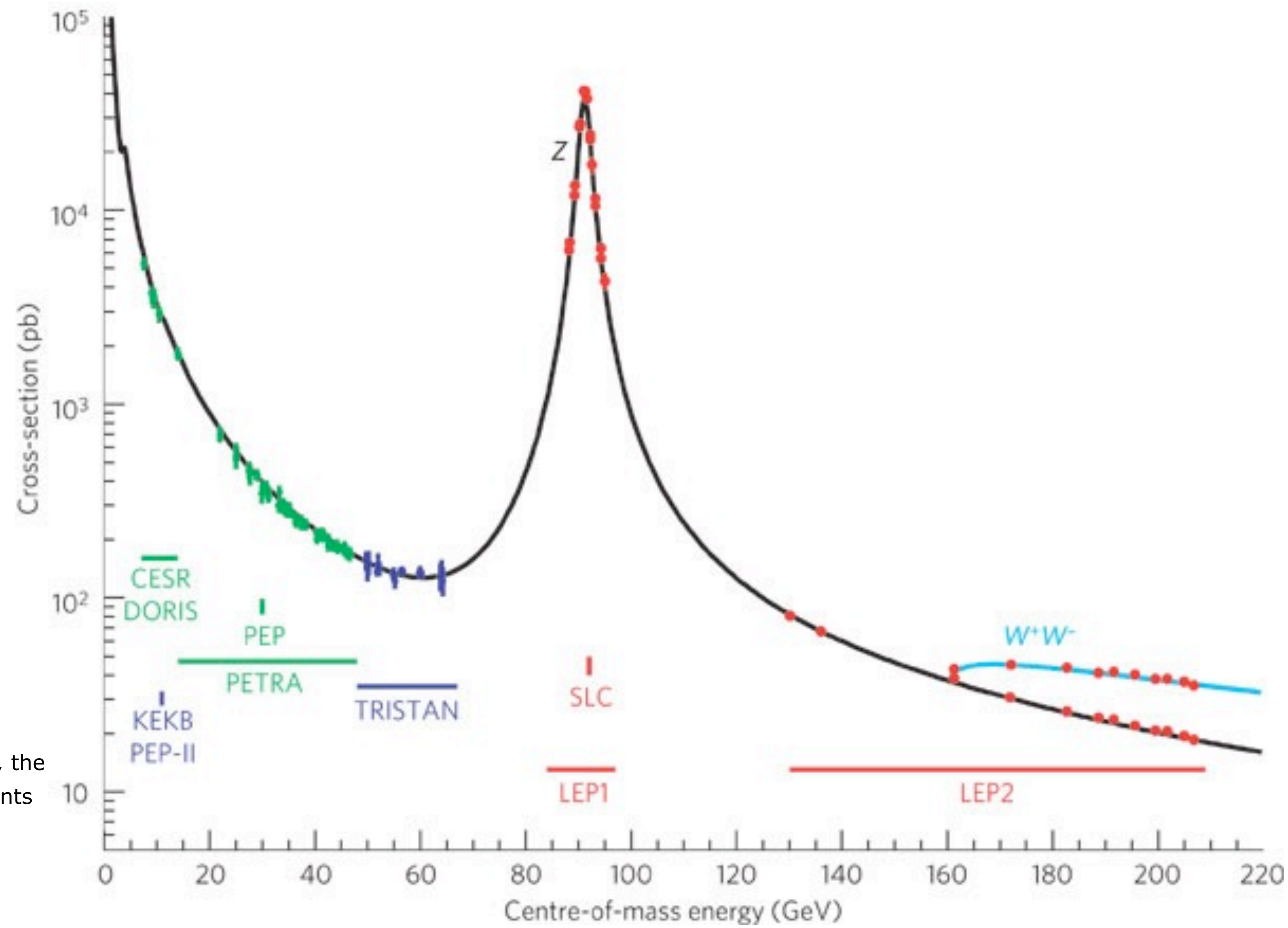




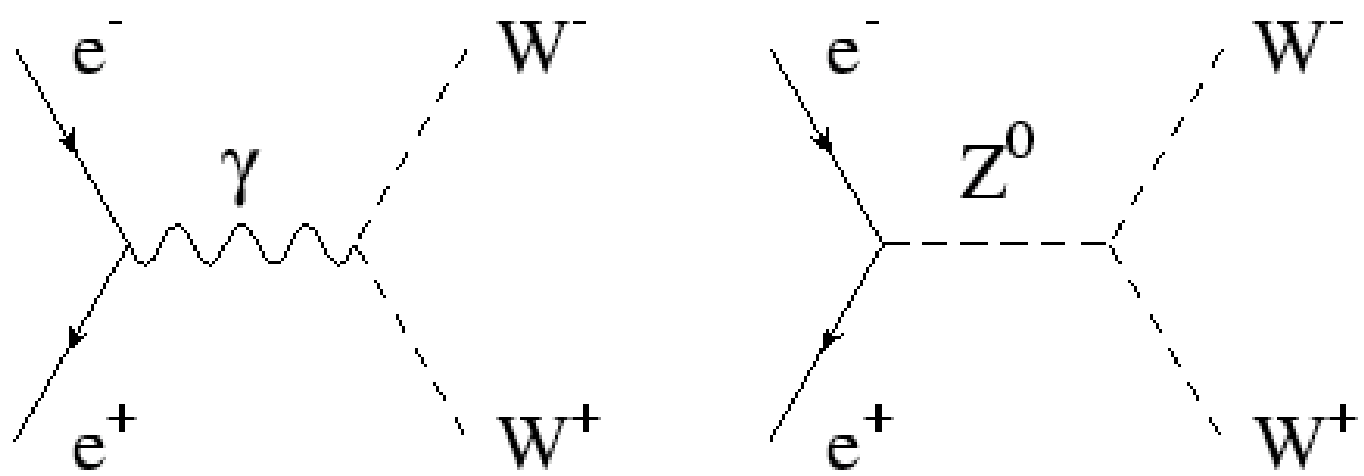
e^+e^- colliders up to LEP

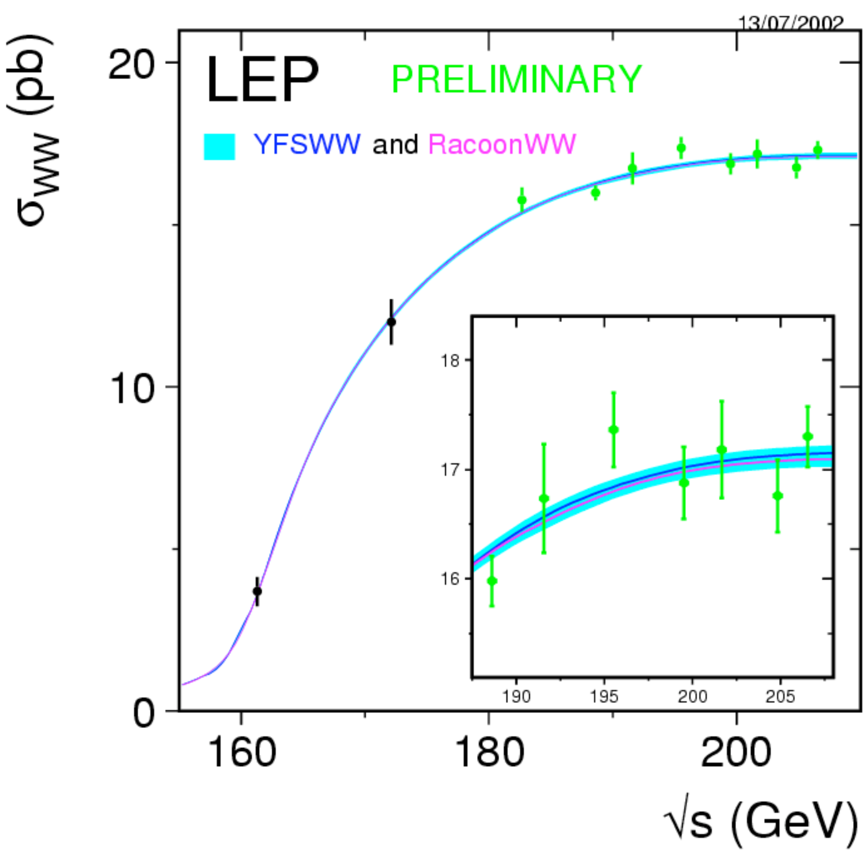


ALEPH, DELPHI, L3, OPAL, SLD collaborations, LEP Electroweak Working Group, the SLD Electroweak and Heavy Flavour Groups. Precision electroweak measurements on the *Z* resonance. *Phys. Rep.* **427**, 257–454 (2006)



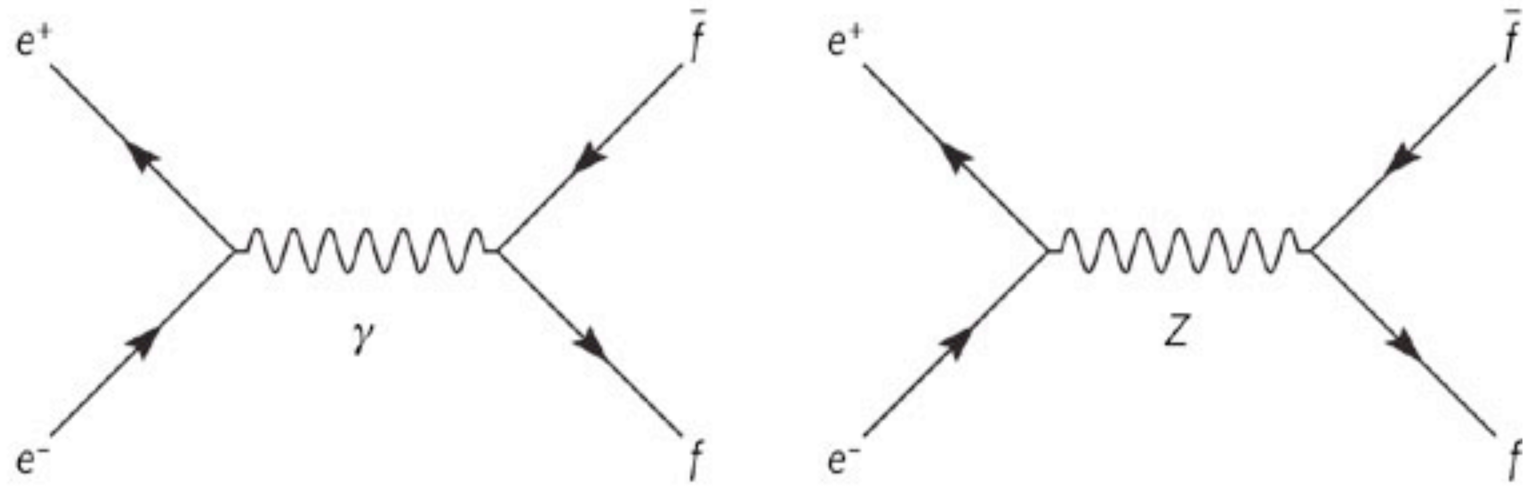
W pair production (LEP2)



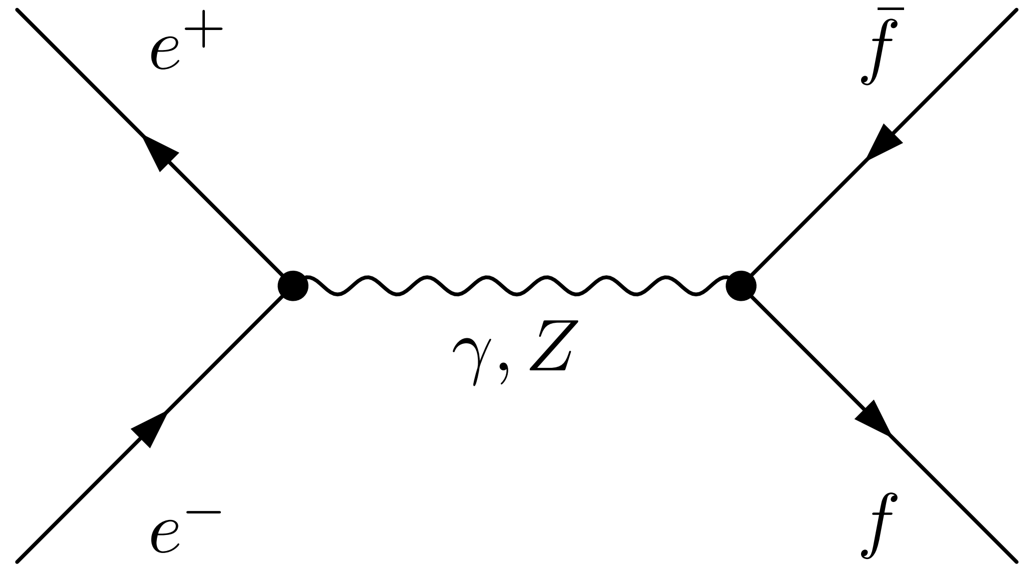


many (confirming) results.....
but the t was is still missing....

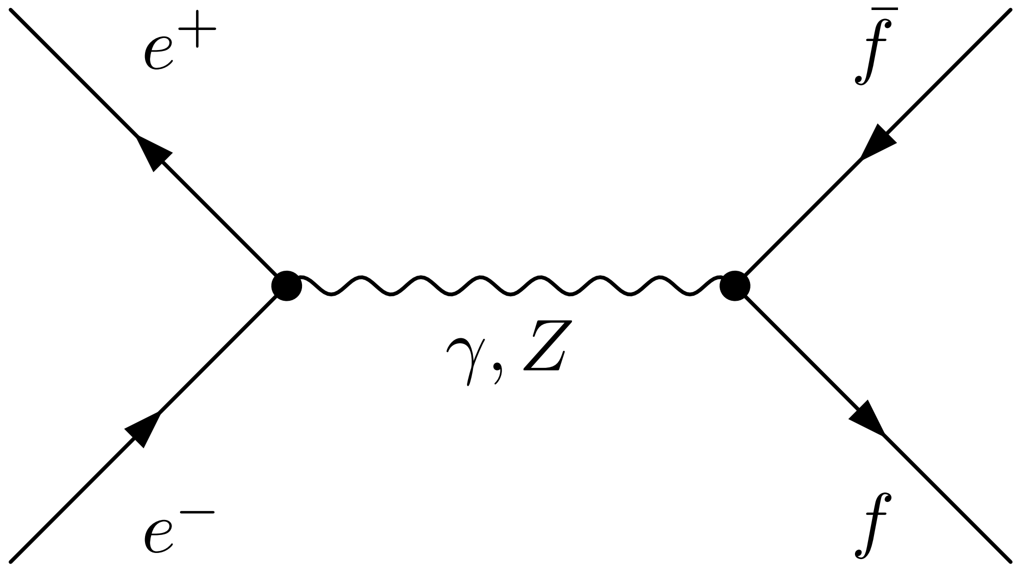
Interference effects in $e^+e^- \rightarrow f\bar{f}$



Interference effects in $e^+e^- \rightarrow f\bar{f}$



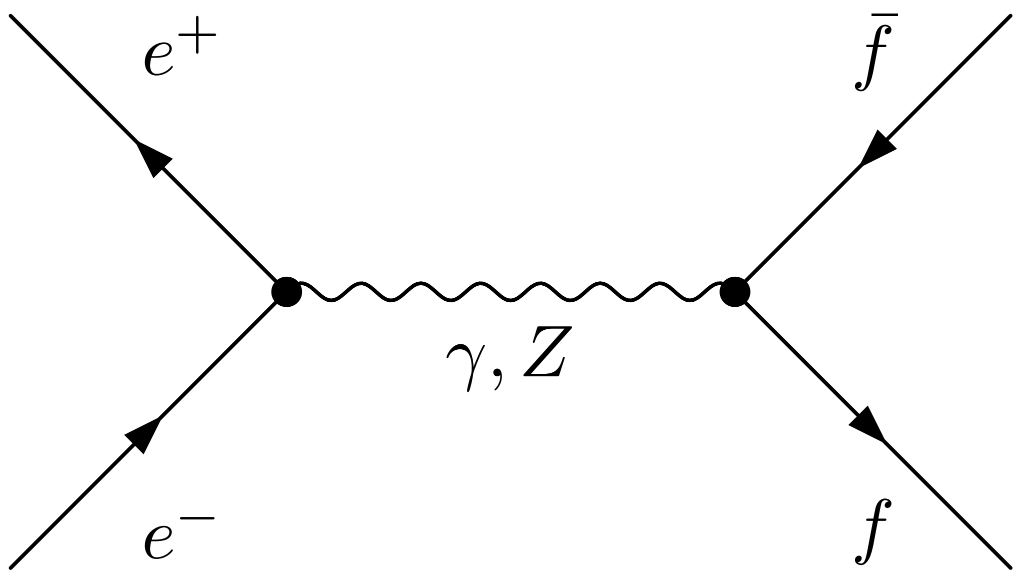
Interference effects in $e^+e^- \rightarrow f\bar{f}$



interference from presence of axial+vector couplings of leptons, quarks to Z

$$\frac{d\sigma_{f\bar{f}}}{d\cos\theta} = \frac{3}{8}\sigma_{f\bar{f}}(1+\cos^2\theta + \frac{8}{3}A^f_{FB}\cos\theta),$$

Interference effects in $e^+e^- \rightarrow f\bar{f}$



interference from presence of axial+vector couplings of leptons, quarks to Z

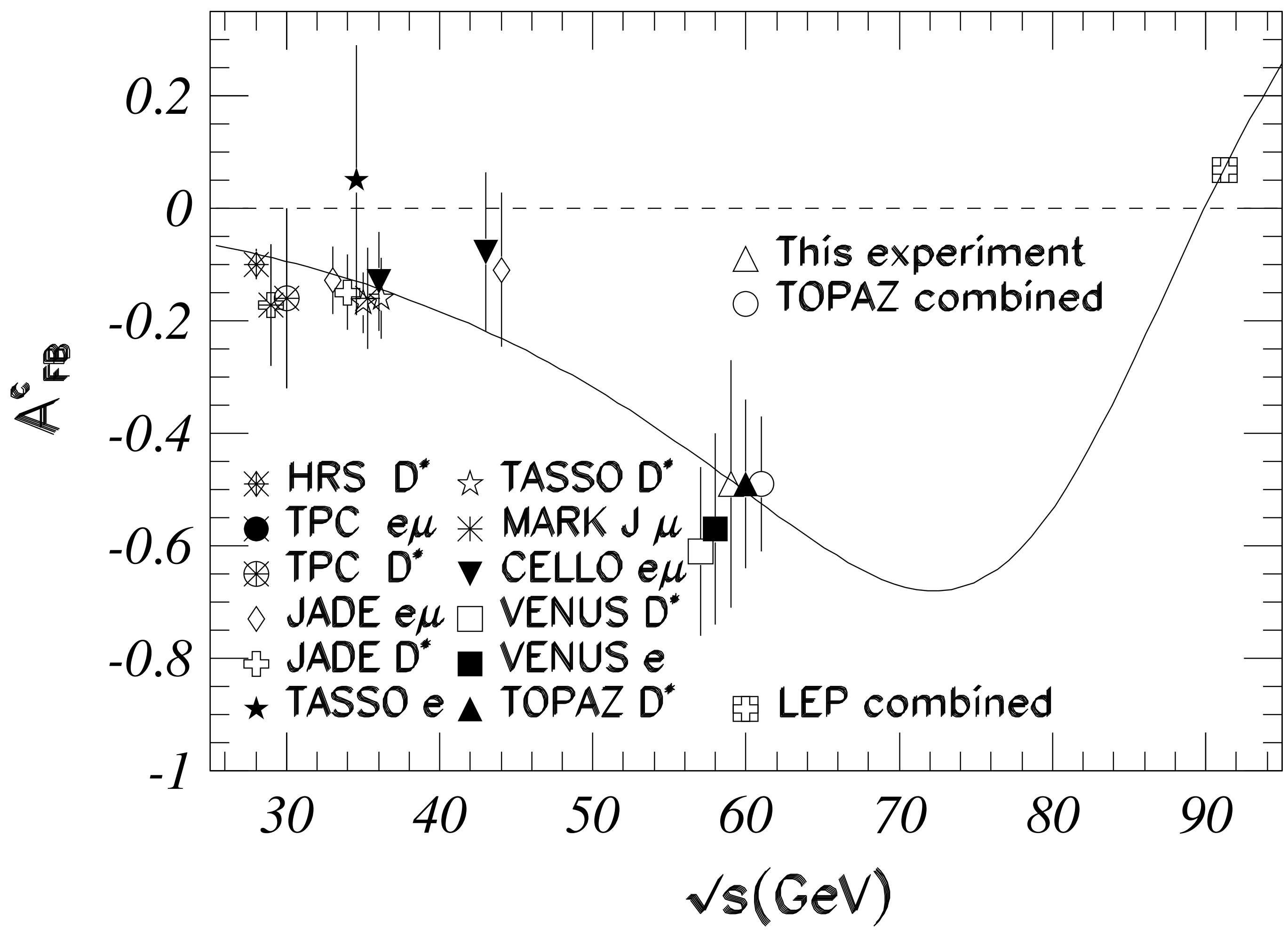
$$\frac{d\sigma_{f\bar{f}}}{d\cos\theta} = \frac{3}{8}\sigma_{f\bar{f}}(1 + \cos^2\theta + \frac{8}{3}A^f_{FB}\cos\theta),$$

Effects small and swamped by huge Z exchange cross section on Z pole

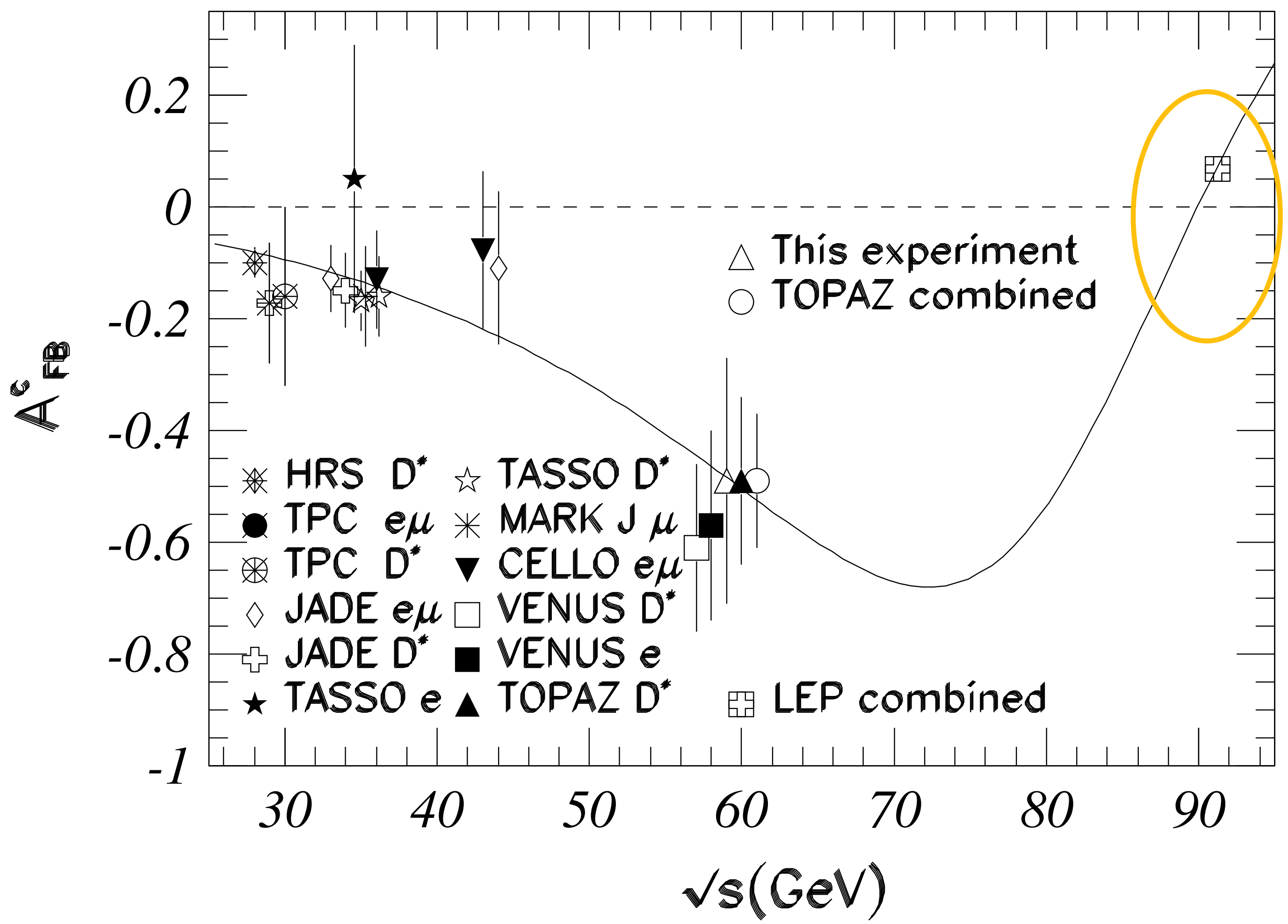
A_{FB} depends on weak isospin, charge of quarks. At TRISTAN (60 GeV):

$$A^c_{FB} = -0.47,$$
$$A^b_{FB} = -0.59$$

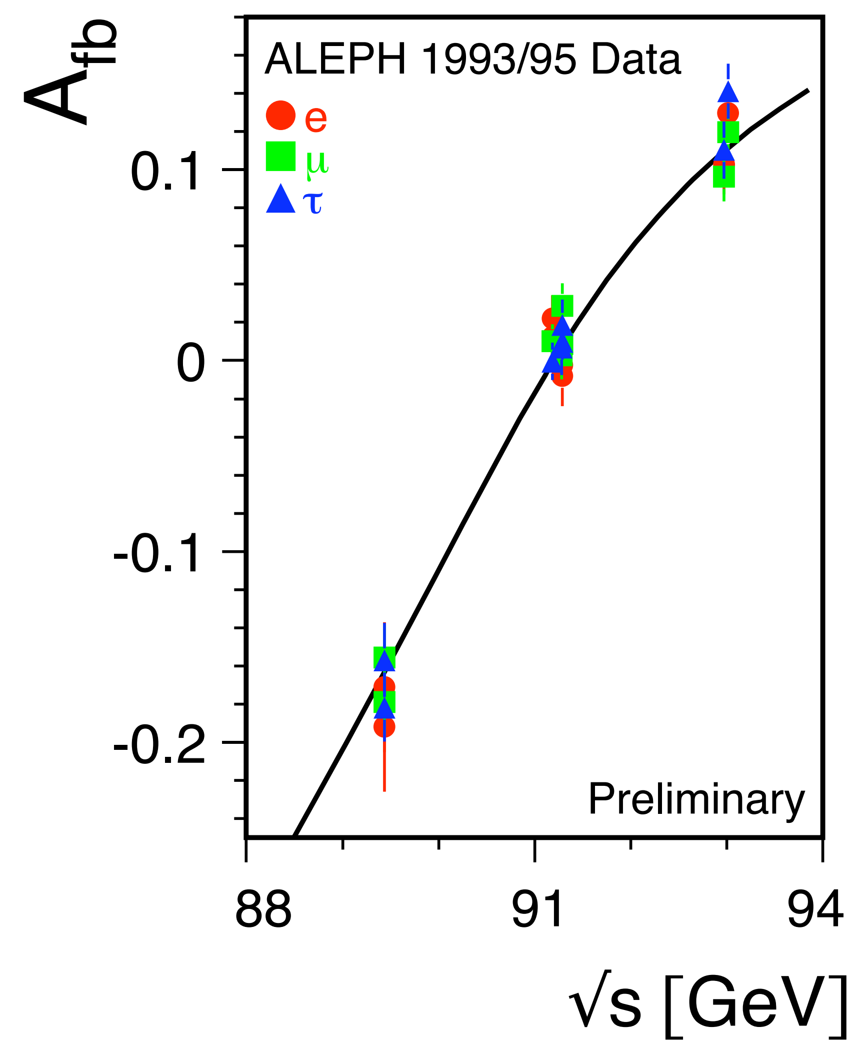
TRISTAN at KEK (60 GeV)



TRISTAN at KEK (60 GeV)



ALEPH at LEP (90 GeV)



LEP and SLD

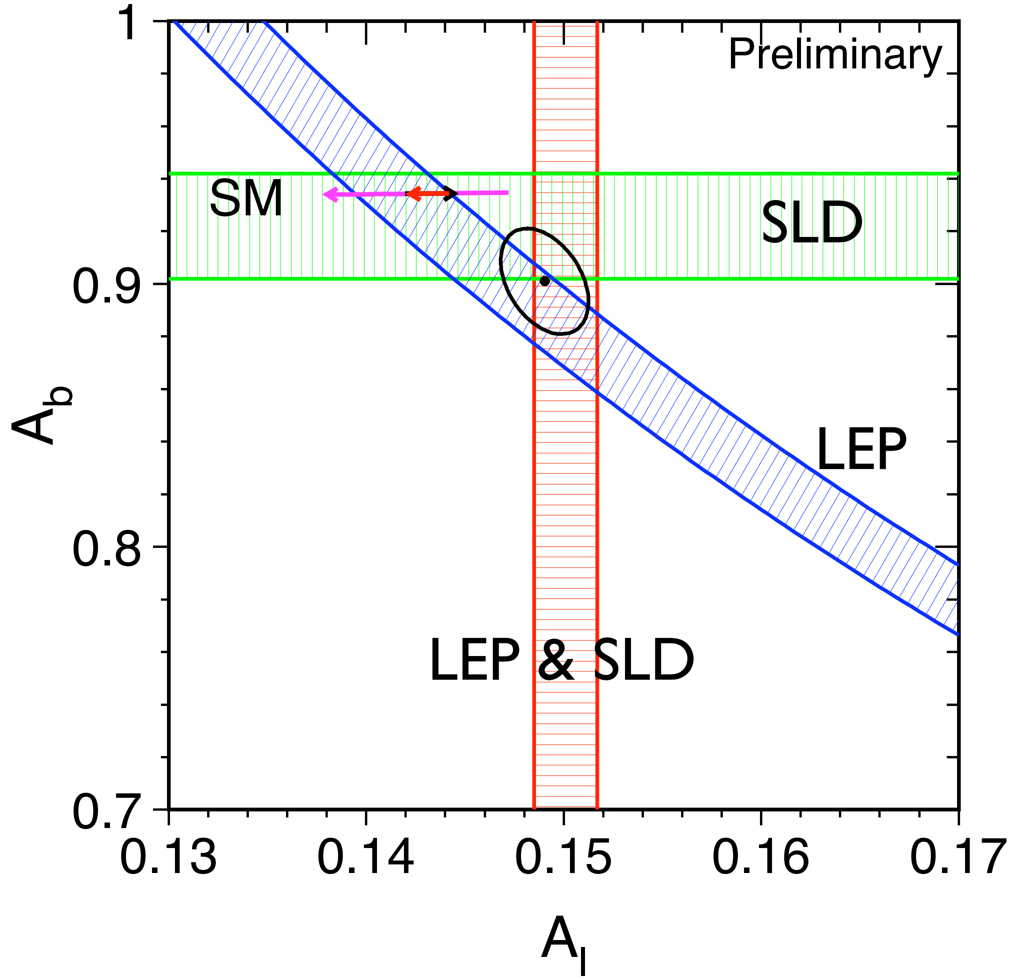
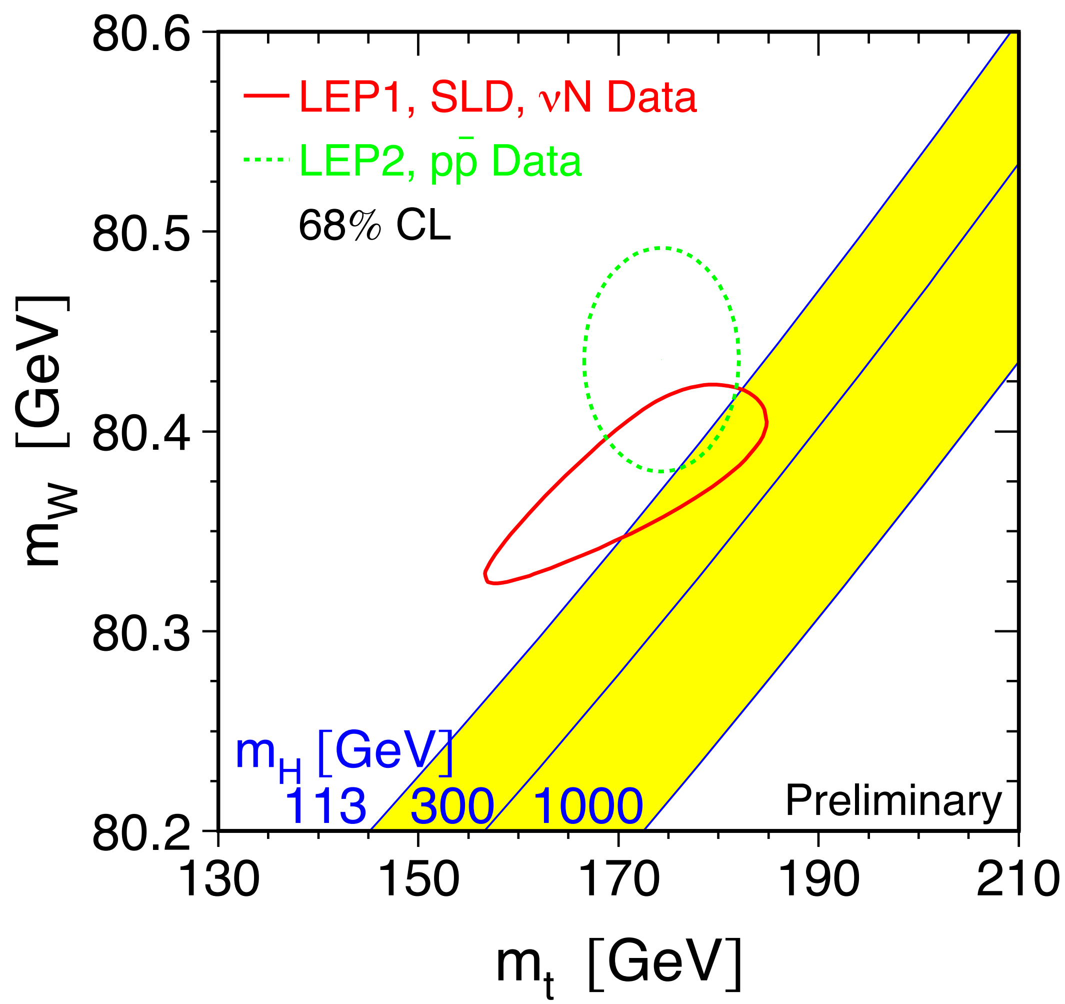


Figure 5. The measurements of the combined LEP+SLD $\mathcal{A}_l$ (vertical band), SLD $\mathcal{A}_b$ (horizontal band) and LEP $A_{\mathrm{FB}}^{b,0}$ (diagonal band), compared to the Standard Model expectation (arrow).



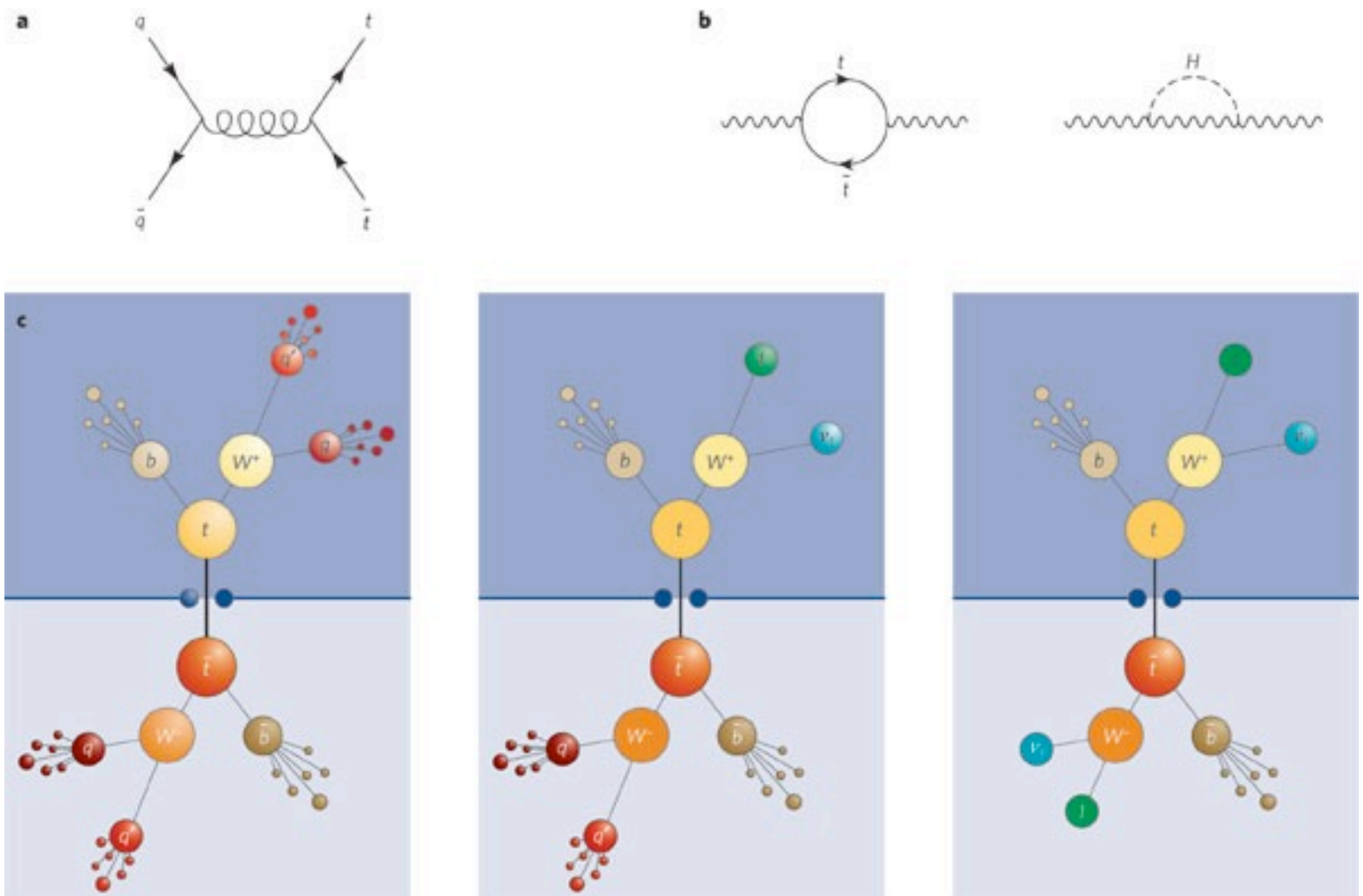
precision measurements were sensitive to m_t before top was discovered (and also sensitive to m_H)

Tevatron: top physics
W physics
search for Higgs

Tevatron: top physics

W physics

search for Higgs

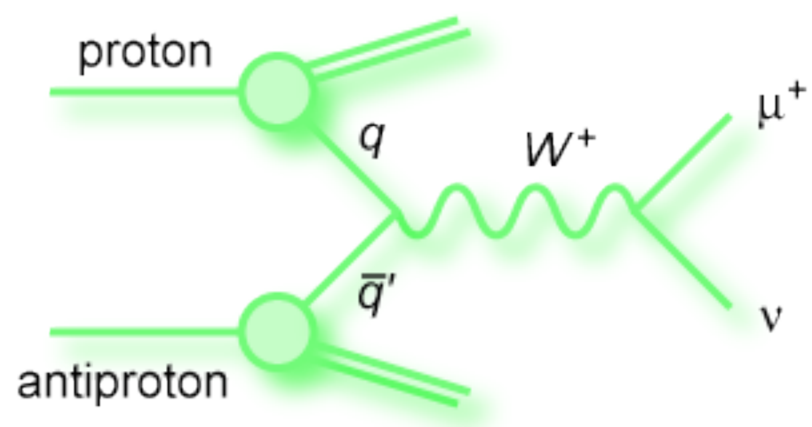


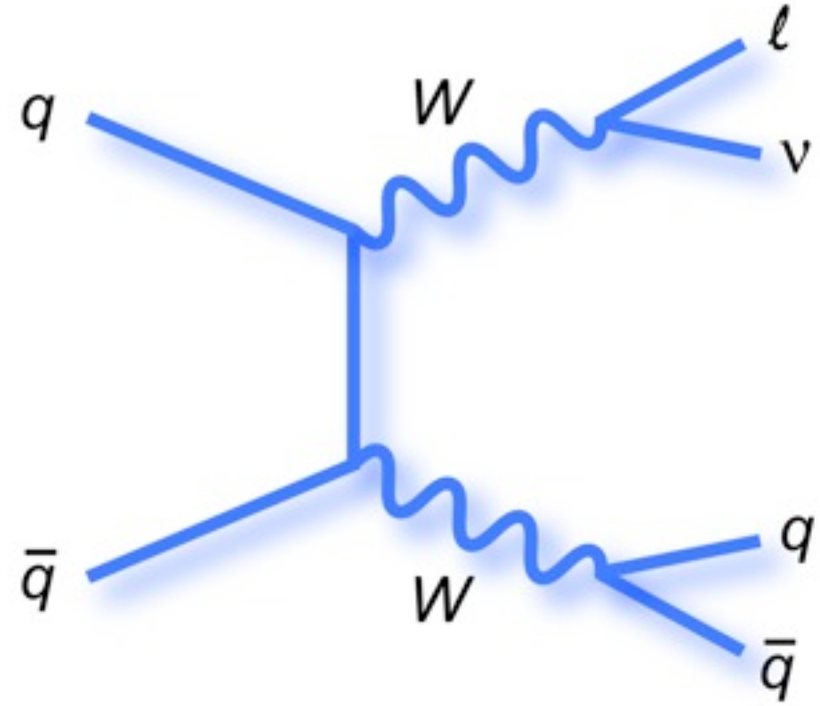
Tevatron:

top physics

W physics

search for Higgs





Tevatron:

top physics

W physics

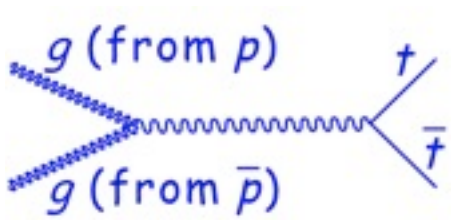
search for Higgs

Tevatron: Discovery of the Top-Quark

FNAL: 1995

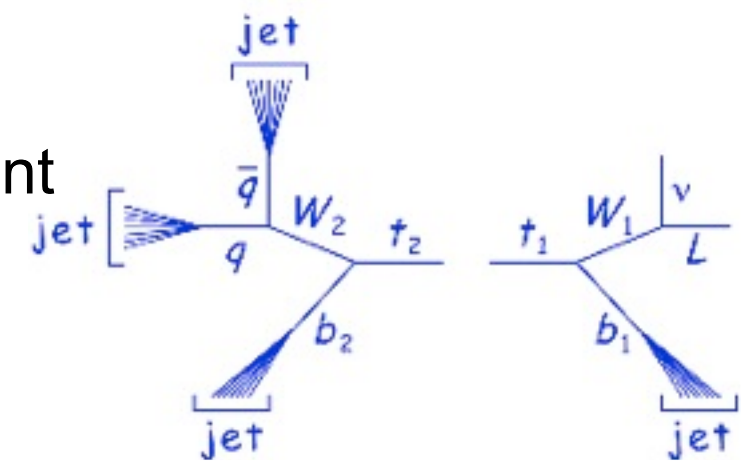
Tevatron : $\sqrt{s}$ = 1.8 TeV
Detectors : CDF, D∅

Dominant Production

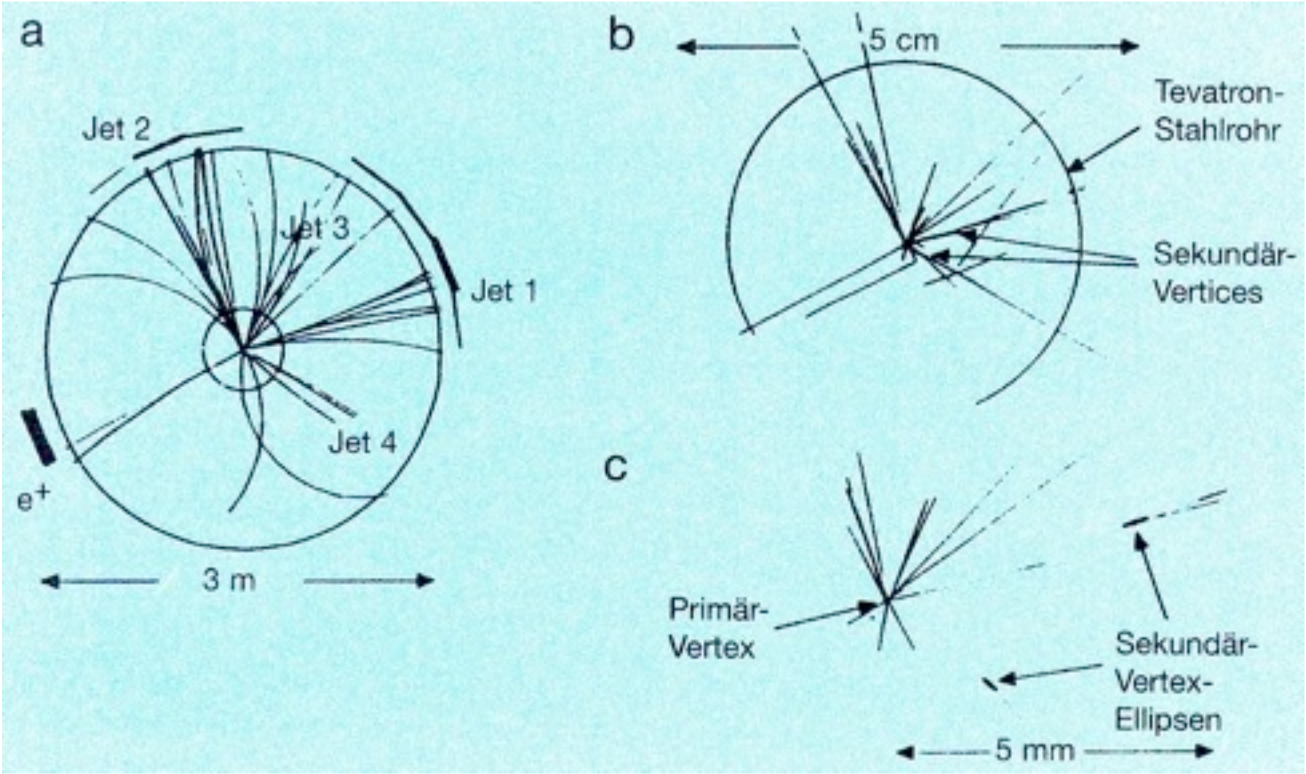


σ ≈ 4 pb (High p⊥)
(σTot ≈ 60 mb (10 o.m. bigger) (< p⊥ > ≈ 0.5 GeV))

Dominant Decay



Trigger on high p⊥ and secondary (b) vertex



⇒ m_t = (174.3 ± 5.1) GeV/c2

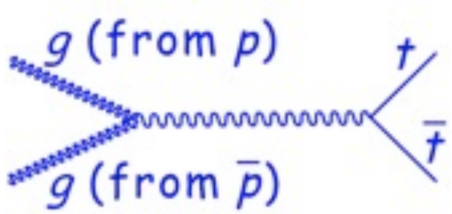
σ(pp̄) → tt̄ in agreement with SM-predictions

Tevatron: Discovery of the Top-Quark

FNAL: 1995

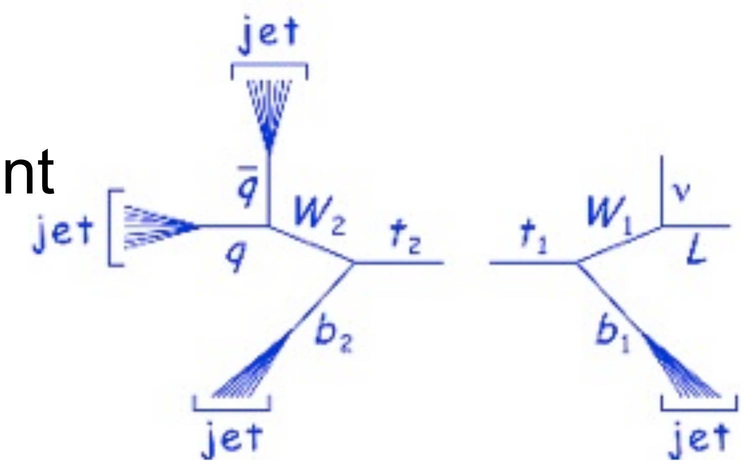
Tevatron : $\sqrt{s}$ = 1.8 TeV
Detectors : CDF, D∅

Dominant Production

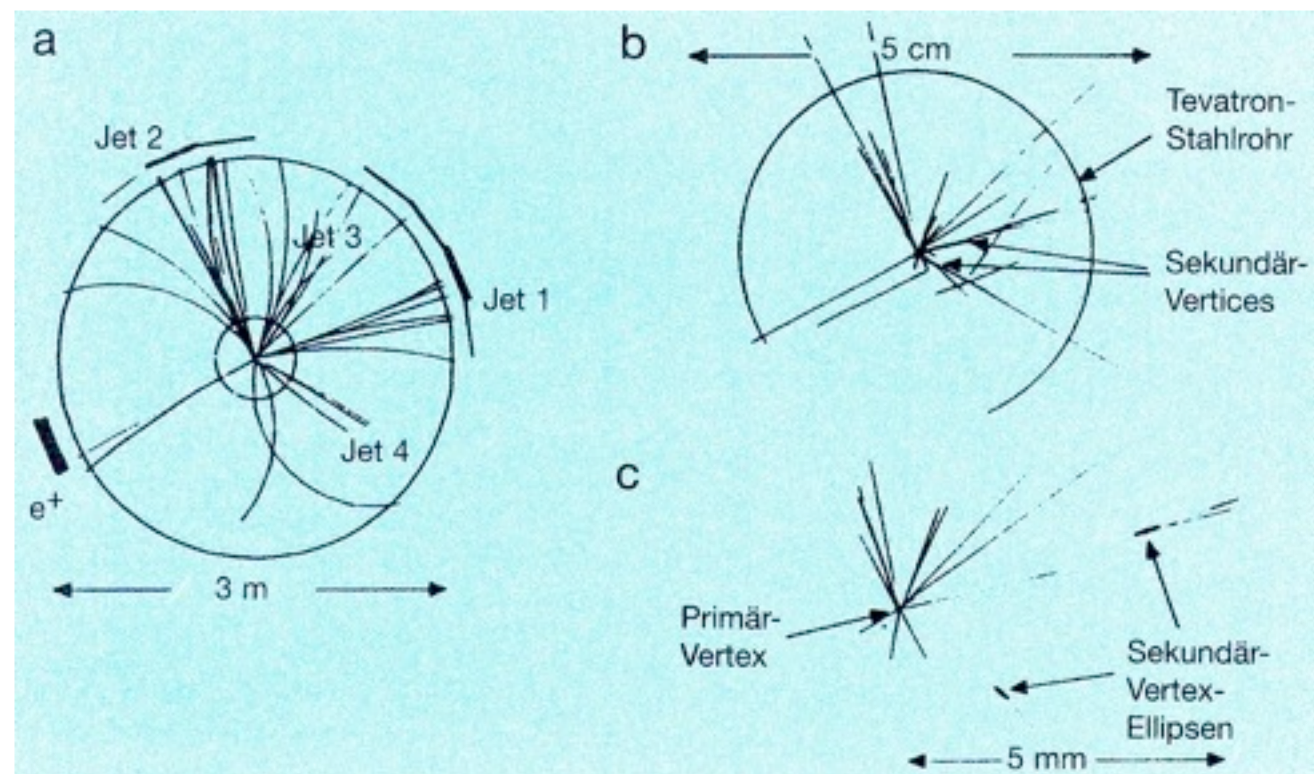


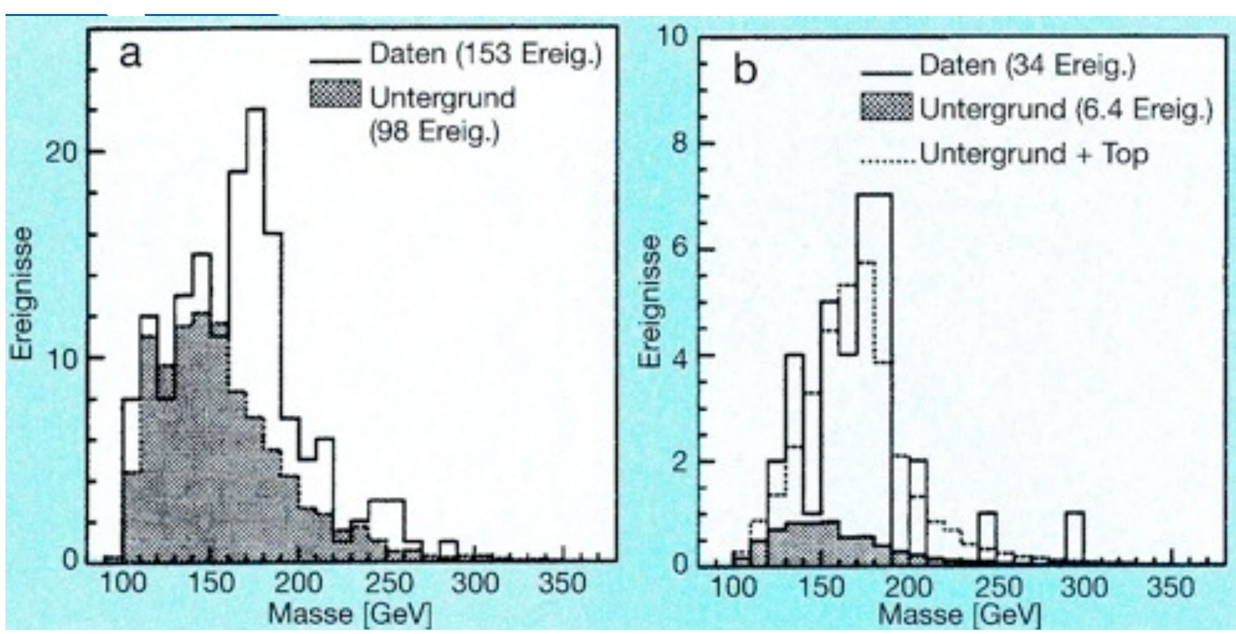
$\sigma \approx 4$ pb (High p⊥)
(σTot ≈ 60 mb (10 o.m. bigger) (< p⊥ > ≈ 0.5 GeV))

Dominant Decay



Trigger on high p⊥ and secondary (b) vertex

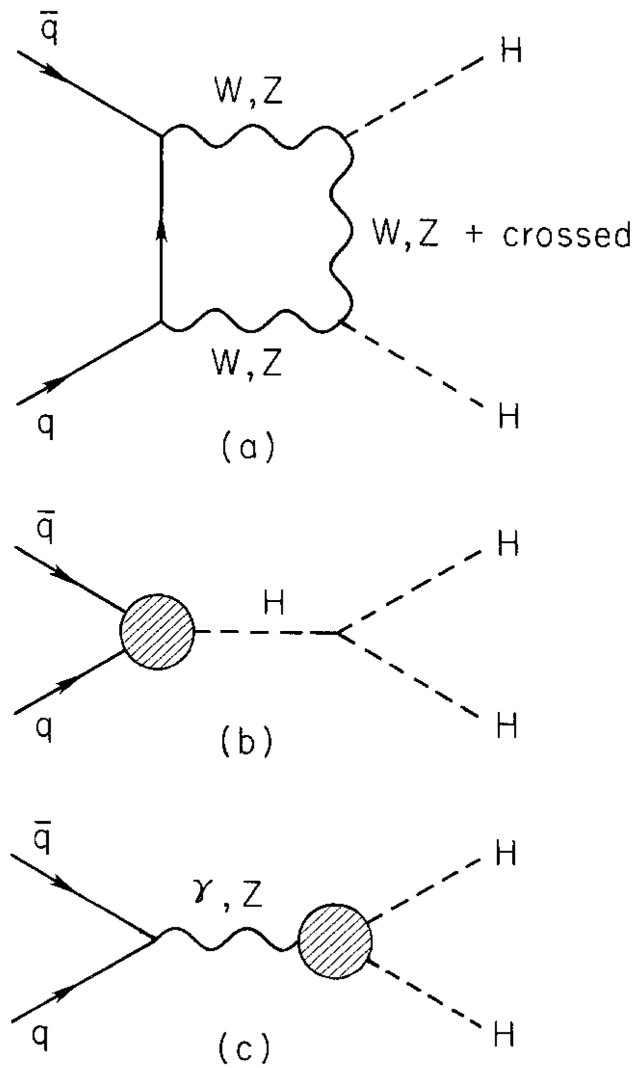


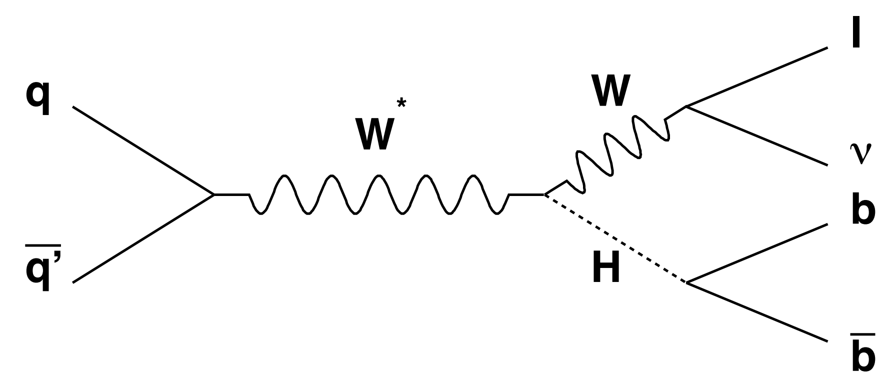


⇒ m_t = (174.3 ± 5.1) GeV/c2

$\sigma(p\bar{p}) \rightarrow t\bar{t}$ in agreement with SM-predictions

on to the Higgs; why not $\bar{p}p$?

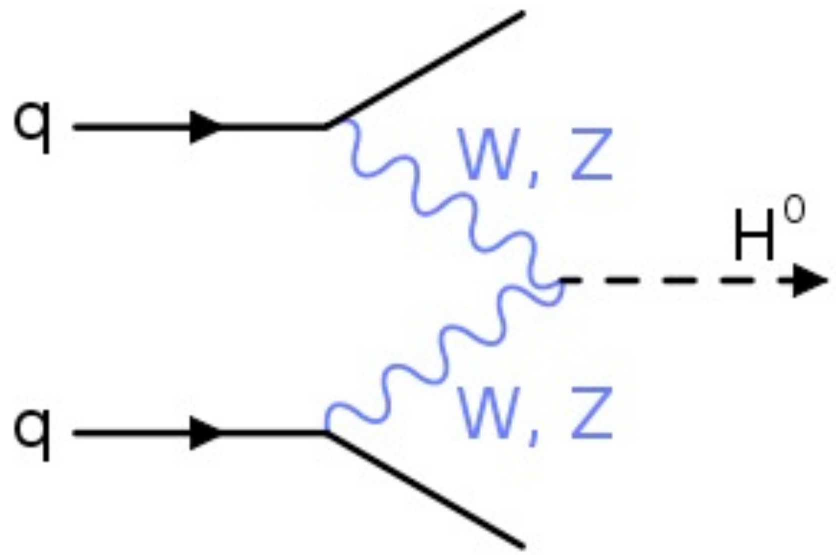




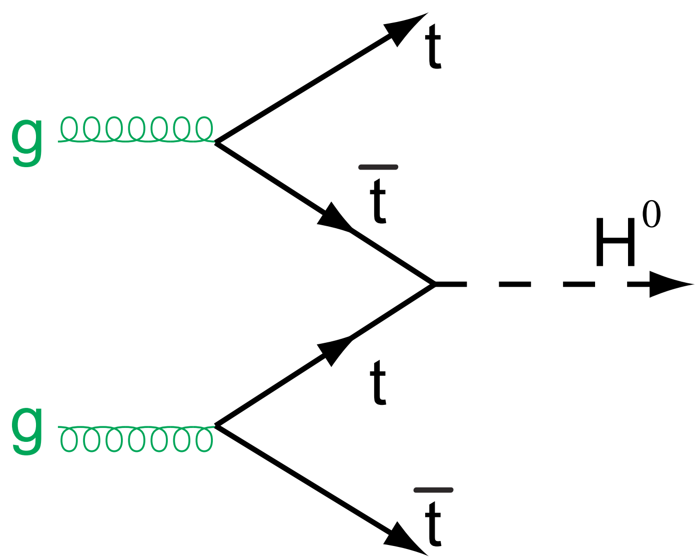
all perfectly respectable
production mechanisms, but ...

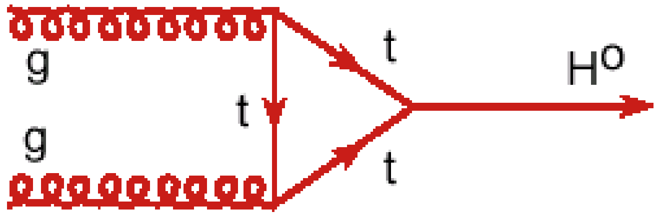
on to the Higgs; why not $\bar{p}p$?

all perfectly respectable
production mechanisms, but ...

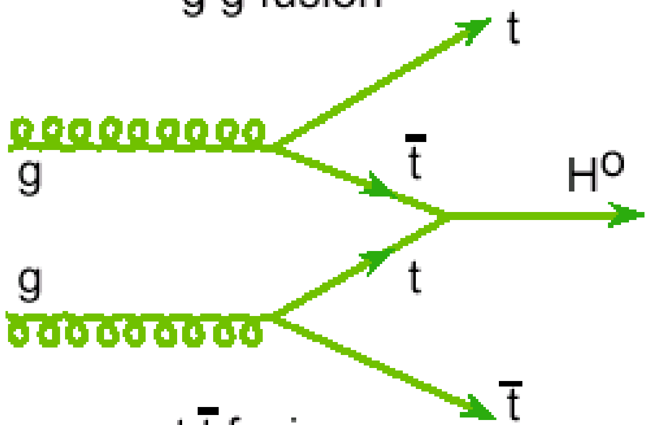


&

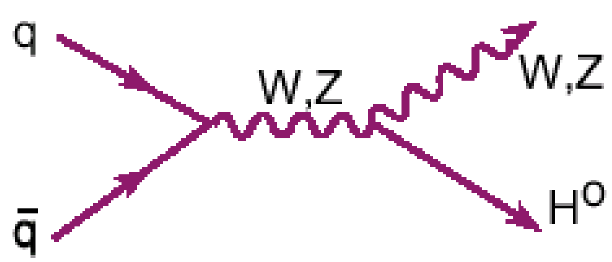




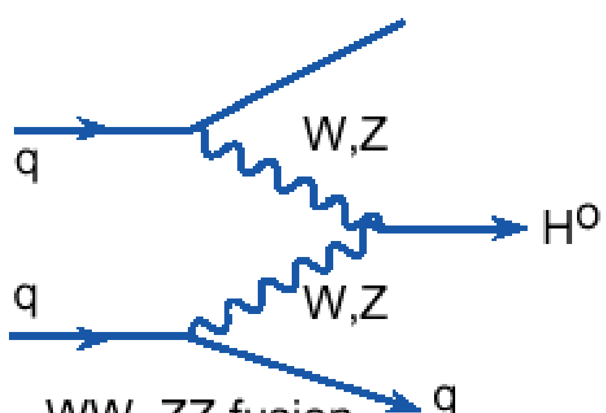
g g fusion



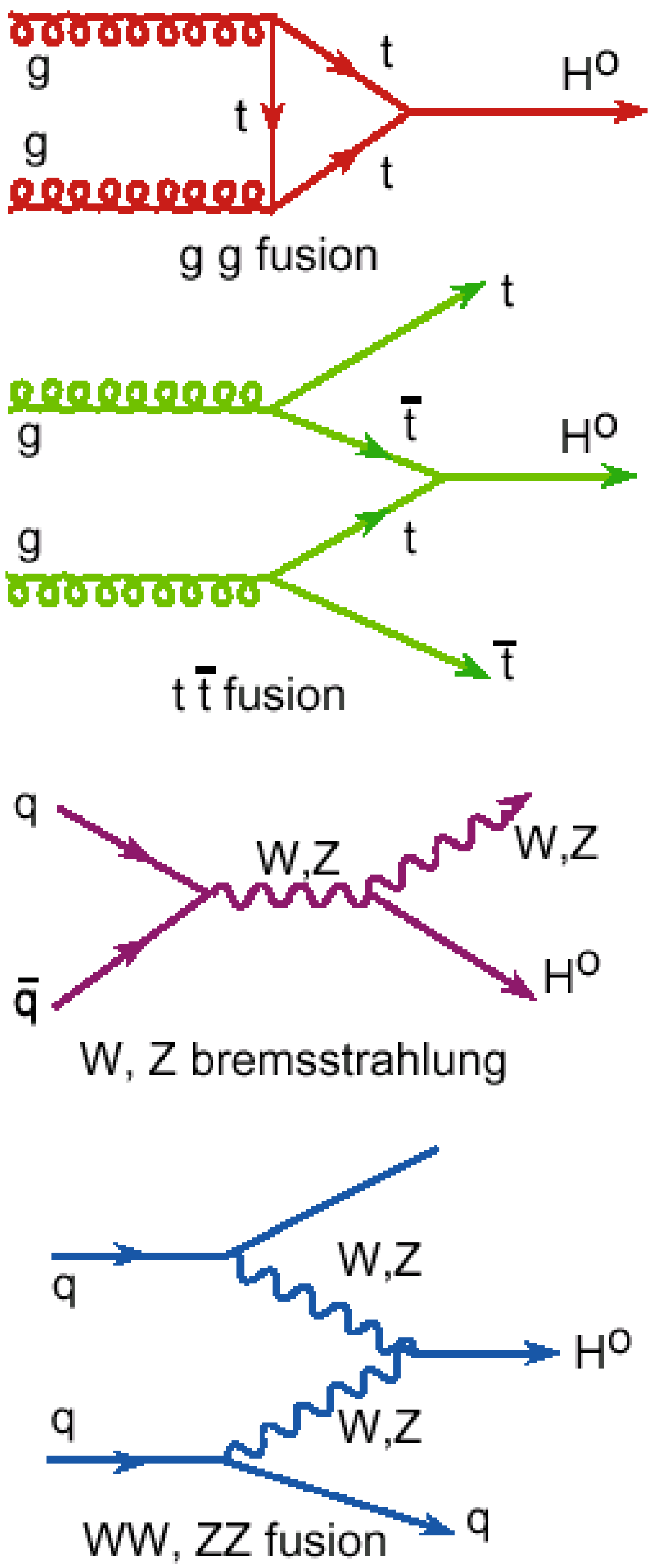
t t̄ fusion



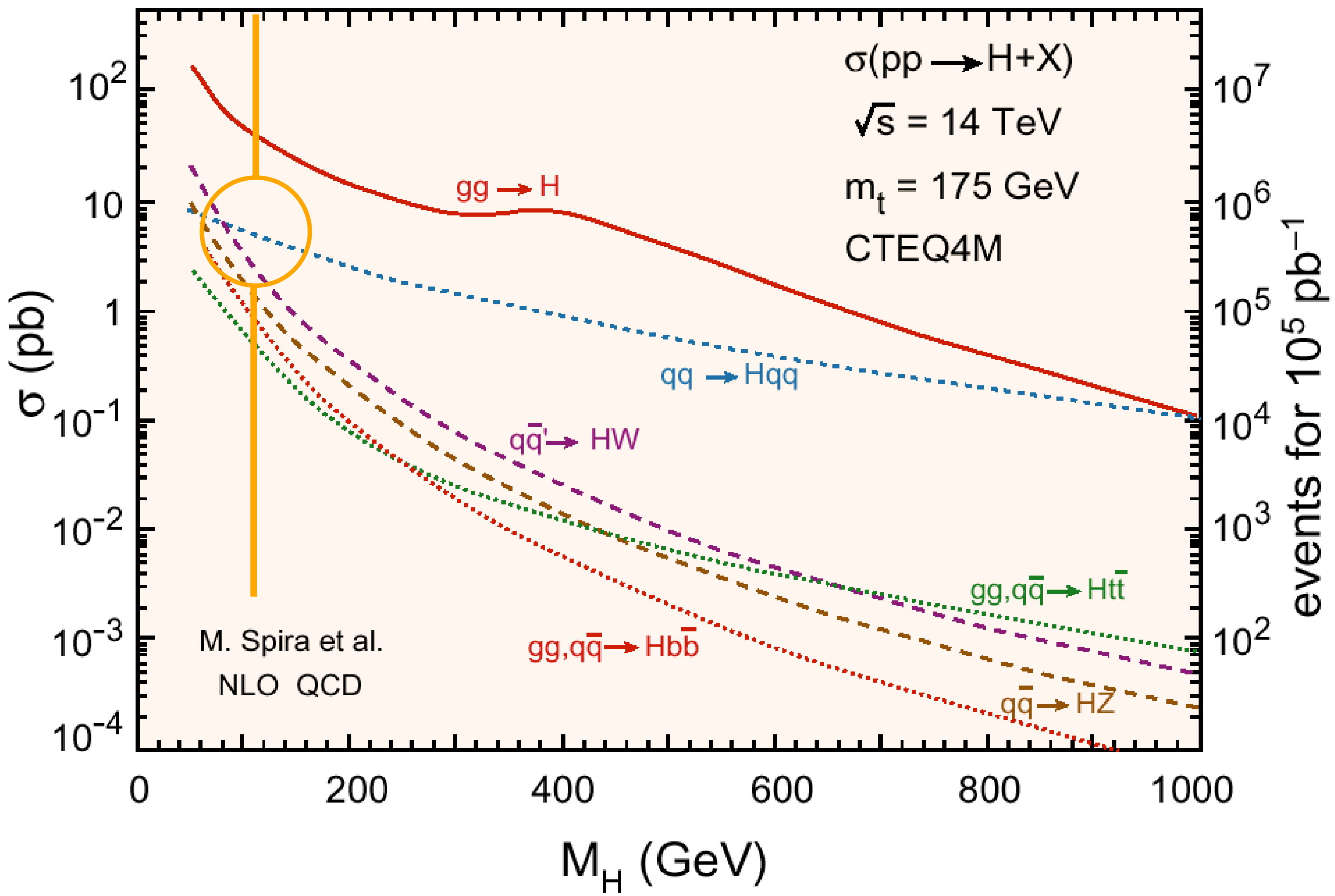
W, Z bremsstrahlung



WW, ZZ fusion



g
g
t
t
t
Ho
g g fusion
g
g
t
t̄
t
t̄
Ho
t t̄ fusion
q
q̄
W,Z
W,Z
Ho
W, Z bremsstrahlung
q
q
W,Z
W,Z
Ho
q
WW, ZZ fusion



σ(pp → H+X)
√s = 14 TeV
m_t = 175 GeV
CTEQ4M
gg → H
qq → Hqq
qq̄ → HW
gg,qq̄ → Htt̄
gg,qq̄ → Hbb̄
qq̄ → HZ
M. Spira et al.
NLO QCD
σ (pb)
events for 10^5 pb^-1
M_H (GeV)

Advantages of $\bar{p}$-p vs. p-p

Advantages of p-p vs. p-$\bar{p}$

Advantages of $\bar{p}$-p vs. p-p

higher reaction rates at low (~ 1TeV) energies for specific processes

Advantages of p-p vs. p-$\bar{p}$

higher reaction rates at high (~ 10 TeV) energies

Advantages of $\bar{p}$-p vs. p-p

higher reaction rates at low (~ 1TeV) energies for specific processes

quark-antiquark fusion dominant at low energies

Advantages of p-p vs. p-$\bar{p}$

higher reaction rates at high (~ 10 TeV) energies

gluon fusion is dominant process in any hadronic machine at high energies

Advantages of $\bar{p}$-p vs. p-p

higher reaction rates at low (~ 1 TeV) energies for specific processes

quark-antiquark fusion dominant at low energies

Advantages of p-p vs. p-$\bar{p}$

higher reaction rates at high (~ 10 TeV) energies

gluon fusion is dominant process in any hadronic machine at high energies

at high energies, gluon fusion is the dominant process, and the gluon pdf's are the same for p as for $\bar{p}$

Advantages of $\bar{p}$-p vs. p-p

higher reaction rates at low (~ 1TeV) energies for specific processes

quark-antiquark fusion dominant at low energies

one single set of magnet rings (counter-propagating beams, same charges)

Advantages of p-p vs. p-$\bar{p}$

higher reaction rates at high (~ 10 TeV) energies

gluon fusion is dominant process in any hadronic machine at high energies

two magnet rings required (counter-propagating beams, opposite charges)

at high energies, gluon fusion is the dominant process, and the gluon pdf's are the same for p as for $\bar{p}$

Advantages of $\bar{p}$-p vs. p-p

higher reaction rates at low (~ 1TeV) energies for specific processes

quark-antiquark fusion dominant at low energies

one single set of magnet rings (counter-propagating beams, same charges)

Advantages of p-p vs. p-$\bar{p}$

higher reaction rates at high (~ 10 TeV) energies

gluon fusion is dominant process in any hadronic machine at high energies

two magnet rings required (counter-propagating beams, opposite charges)

far easier production of projectiles (antiproton production and cooling is still very difficult and inefficient)

at high energies, gluon fusion is the dominant process, and the gluon pdf's are the same for p as for $\bar{p}$

Overview:

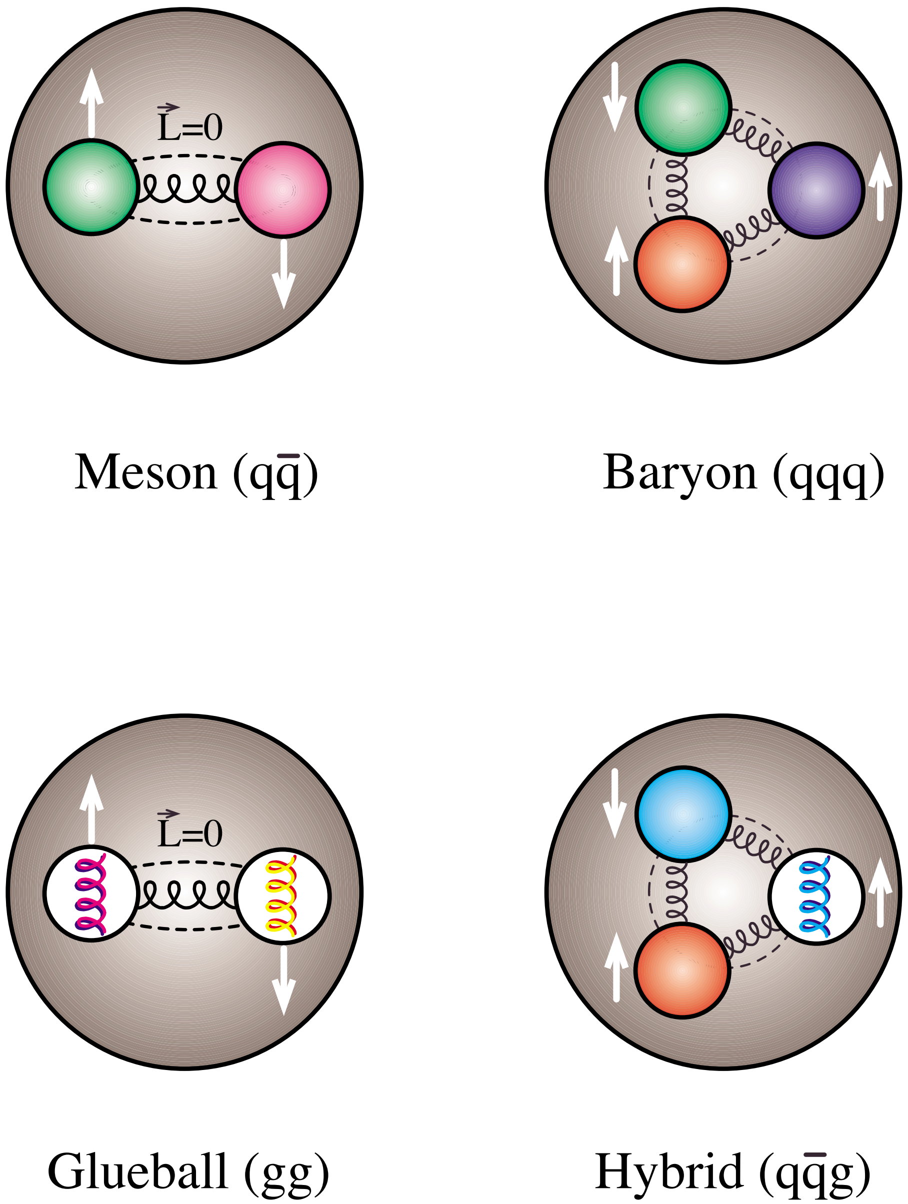
1. Introduction and overview
2. Antimatter at high energies (SppS, LEP, Fermilab)
3. Meson spectroscopy (antimatter as QCD probe)

4. Astroparticle physics and cosmology
5. CP and CPT violation tests
6. Precision tests with Antimatter

7. Precision tests with Antihydrogen
8. Applications of antimatter

Testing QCD with antimatter

QCD

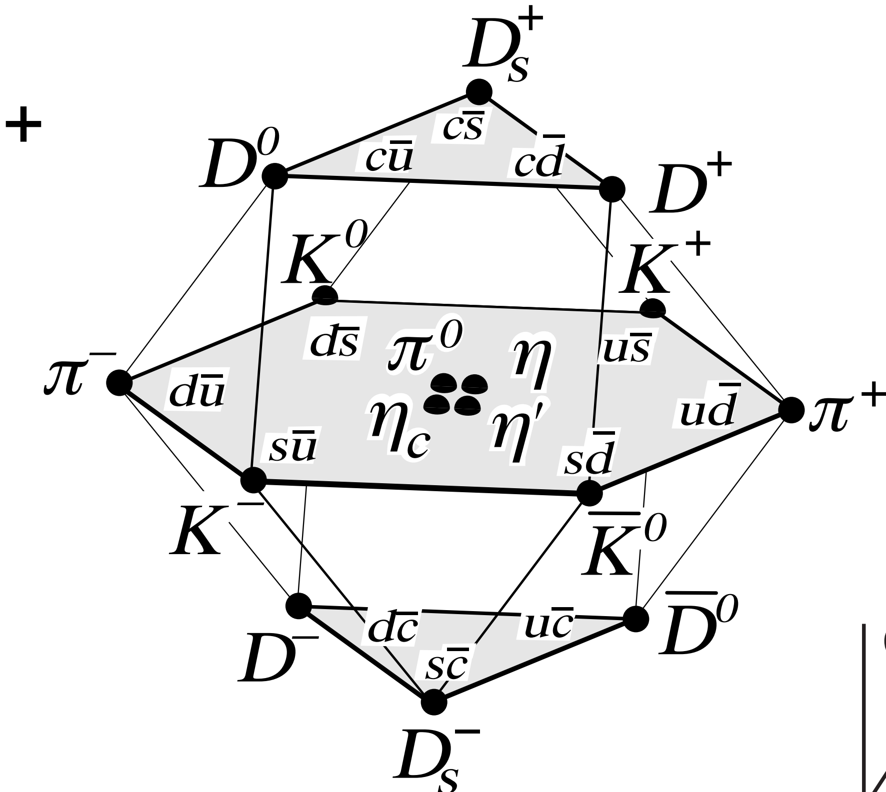


$q\bar{q}$ states

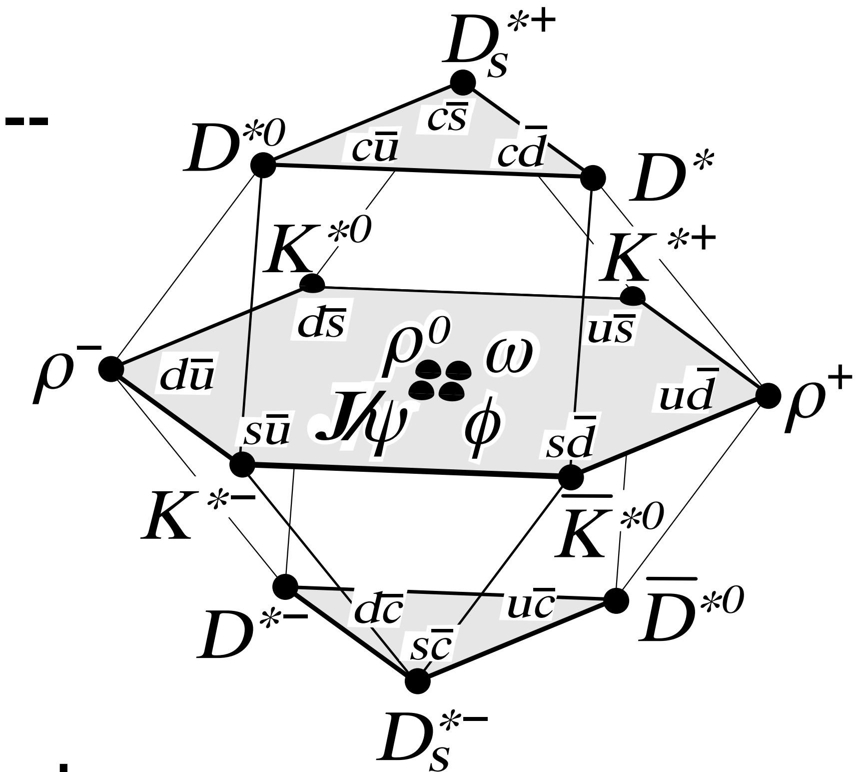
Classification scheme: multiplets

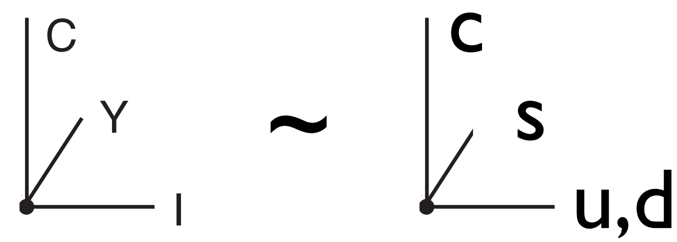
$P(\bar{q}q) = (-1)^L$
$C(\bar{q}q) = (-1)^{L+S}$

0^{-+}



1^{--}





3 quarks: SU(3) $3 \otimes 3 = 8 \oplus 1$ symmetry breaking through quark mass difference

But of course, there are gluons, virtual quark-antiquark pairs, leading to a whole cryptozoology of exotics (glueballs, hybdrids, pentaquarks, ...)

Testing the quark model = search for non-$\bar{q}q$ states

fermionic system

$P(\bar{q}q) = (-1)^{L}$

$C(\bar{q}q) = (-1)^{L+S}$

mesons

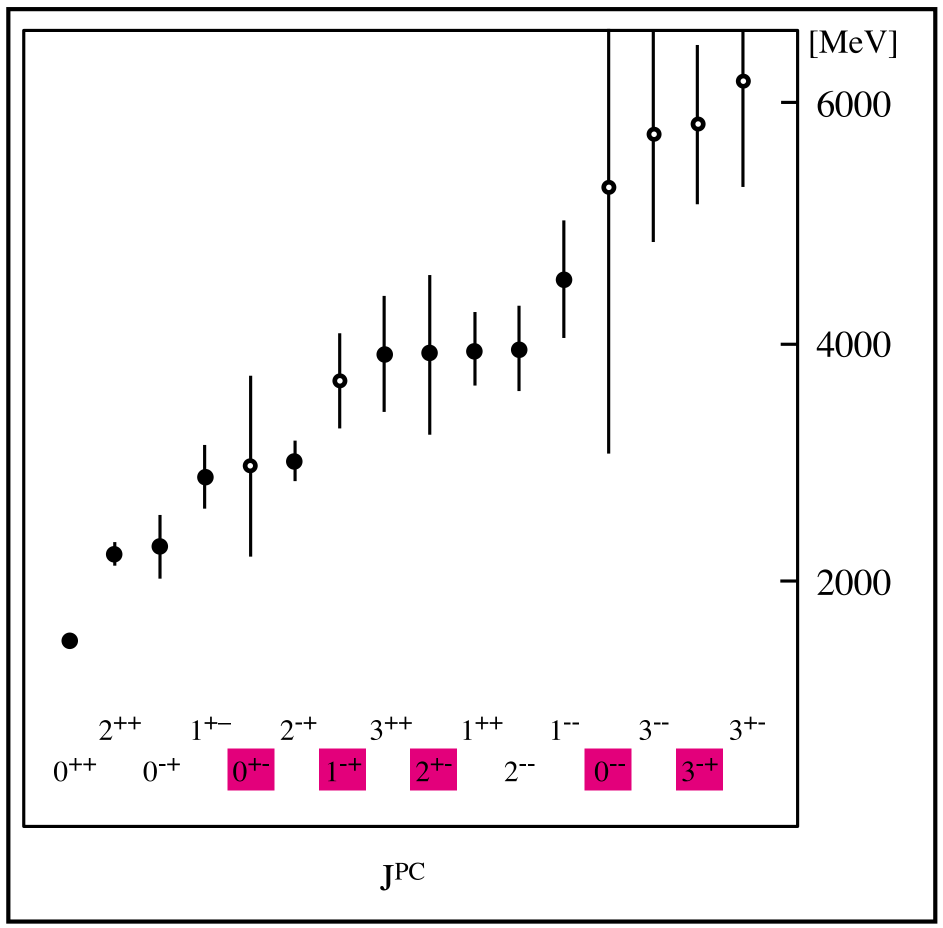
bosonic system

$P = (-1)^{L+1}$

$C = (-1)^{L+S}$

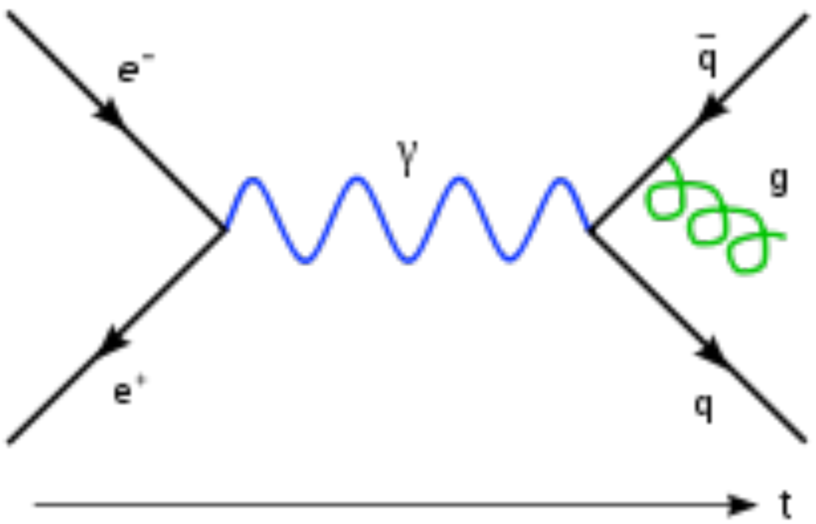
glueballs

color charge:
gluons couple to
other gluons and
can form
bound states



The glueball spectrum predicted by lattice calculations [10]. Exotic quantum numbers are marked as boxes.

Evidence for gluons: e^+e^- annihilation



The idea of searching for gluon jets had actually been proposed by John Ellis, Mary Gaillard and Graham Ross in a seminal paper that appeared in 1976. Under the apparently imperative title "Search for Gluons in e^+-e^- Annihilation", the authors suggested the existence of "hard-gluon bremsstrahlung", which should give rise to events with three jets in the final state. According to the laws of field theory, the outgoing quarks can radiate field quanta of the strong interaction, i.e. gluons, which should in turn fragment into hadrons and thus create a third hadron jet forming a plane with the other two (see figure 1). At the particle energies of up to 15GeV per beam delivered by DESY's newly built PETRA electron-positron storage ring, the probability for such hard-gluon bremsstrahlung processes to occur might amount to a few percent.

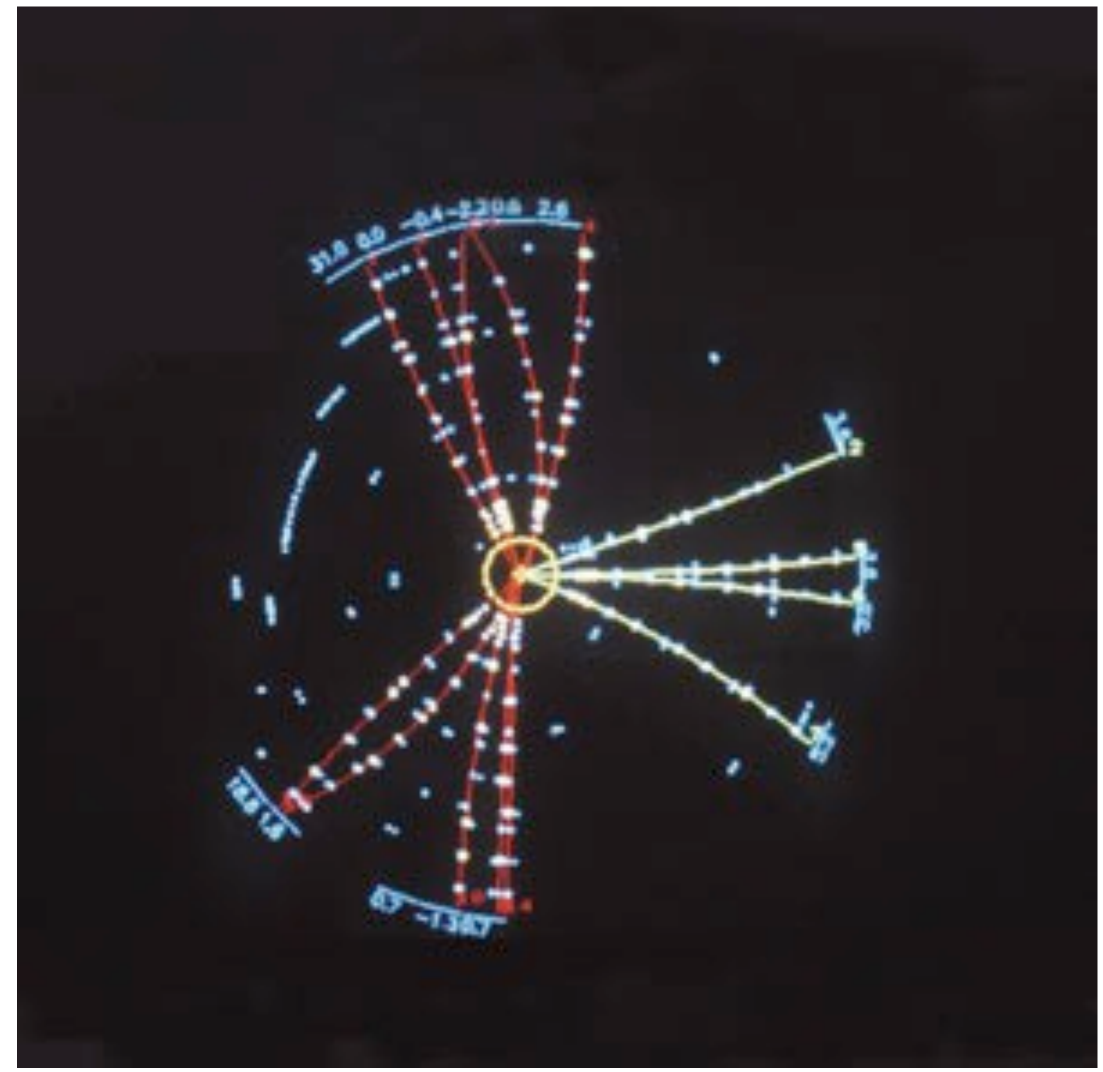


Fig. 10.19 The same as Fig. 10.17 except that this event is one of the rare, separated, three jet events. The total energy is 35.16 GeV.

TASSO experiment at DESY (PETRA, 1978)

Antiproton-proton annihilation (at rest)

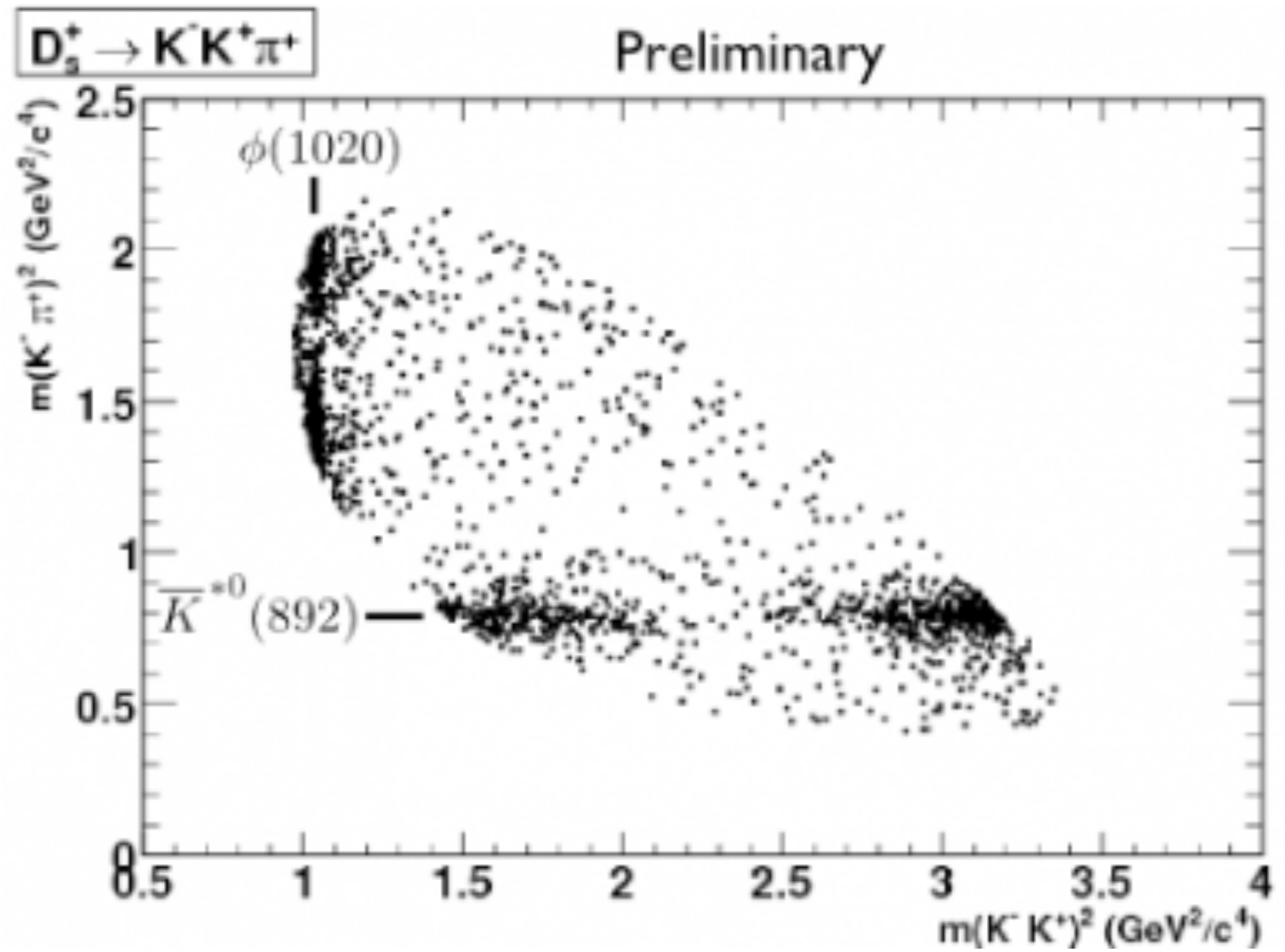
Available energy = 2 m_p <annihilation> ~ 3π

Dalitz plot (any 3-body final state)

m^2 is relativistically invariant;
plot m^2_{12} vs. m^2_{23}

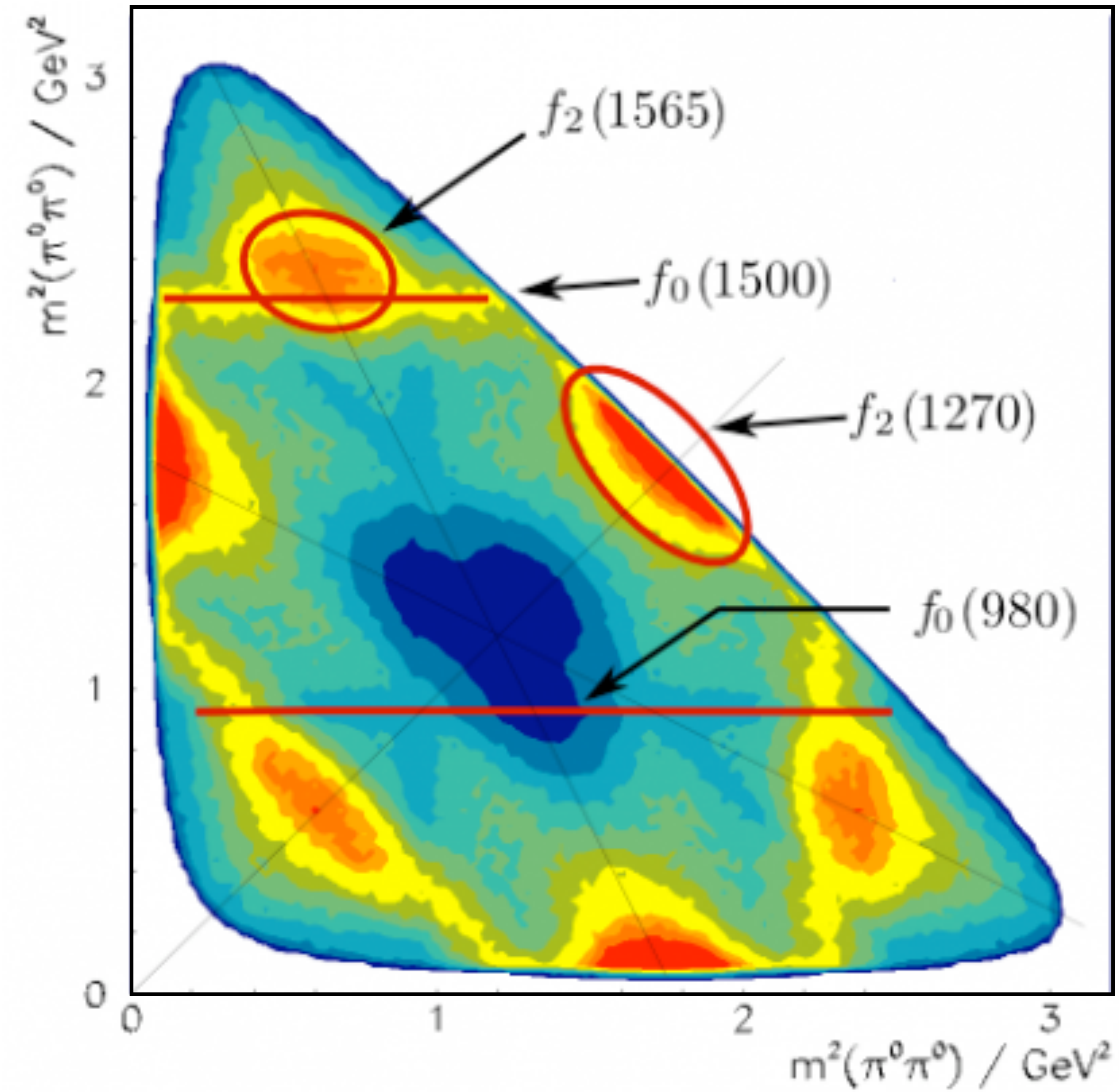
energy-momentum conservation
= limits of contour

no resonances = uniform population
intermediate states = structures

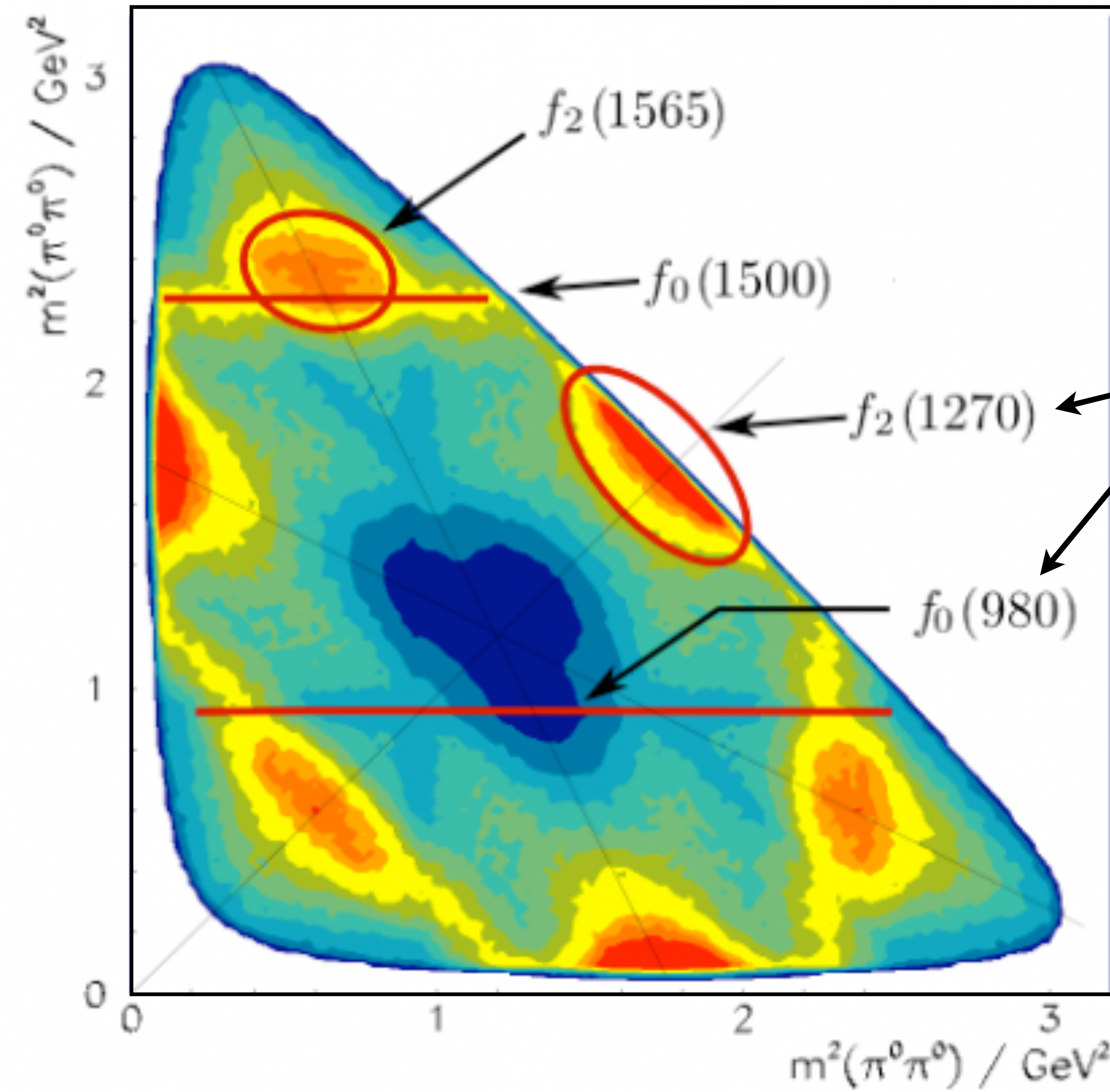


http://superweak.wordpress.com/2006/07/31/dalitz-plots/

$p\bar{p} \rightarrow 3\pi^0$

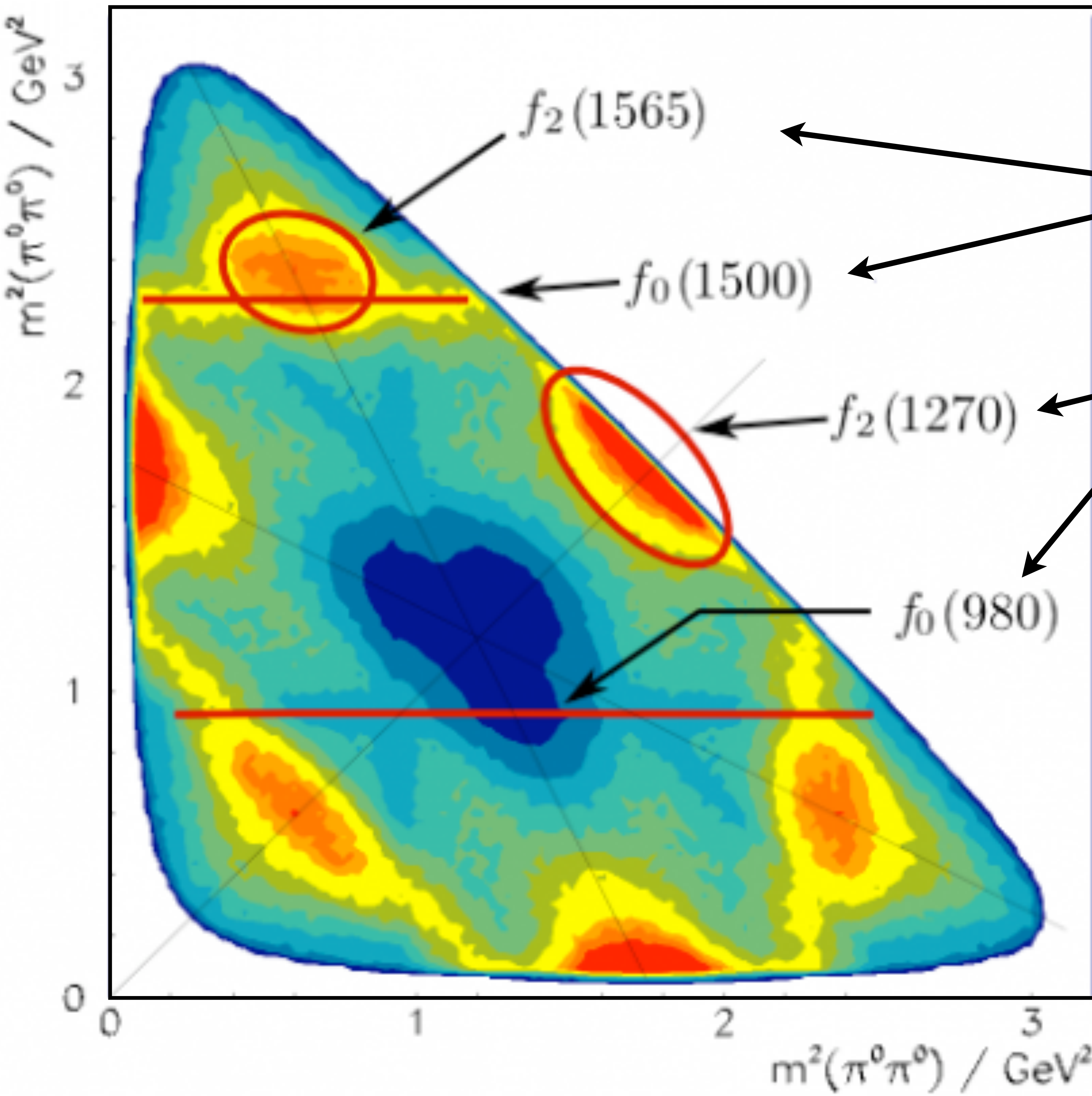


$p\bar{p} \rightarrow 3\pi^0$



standard (known) mesons

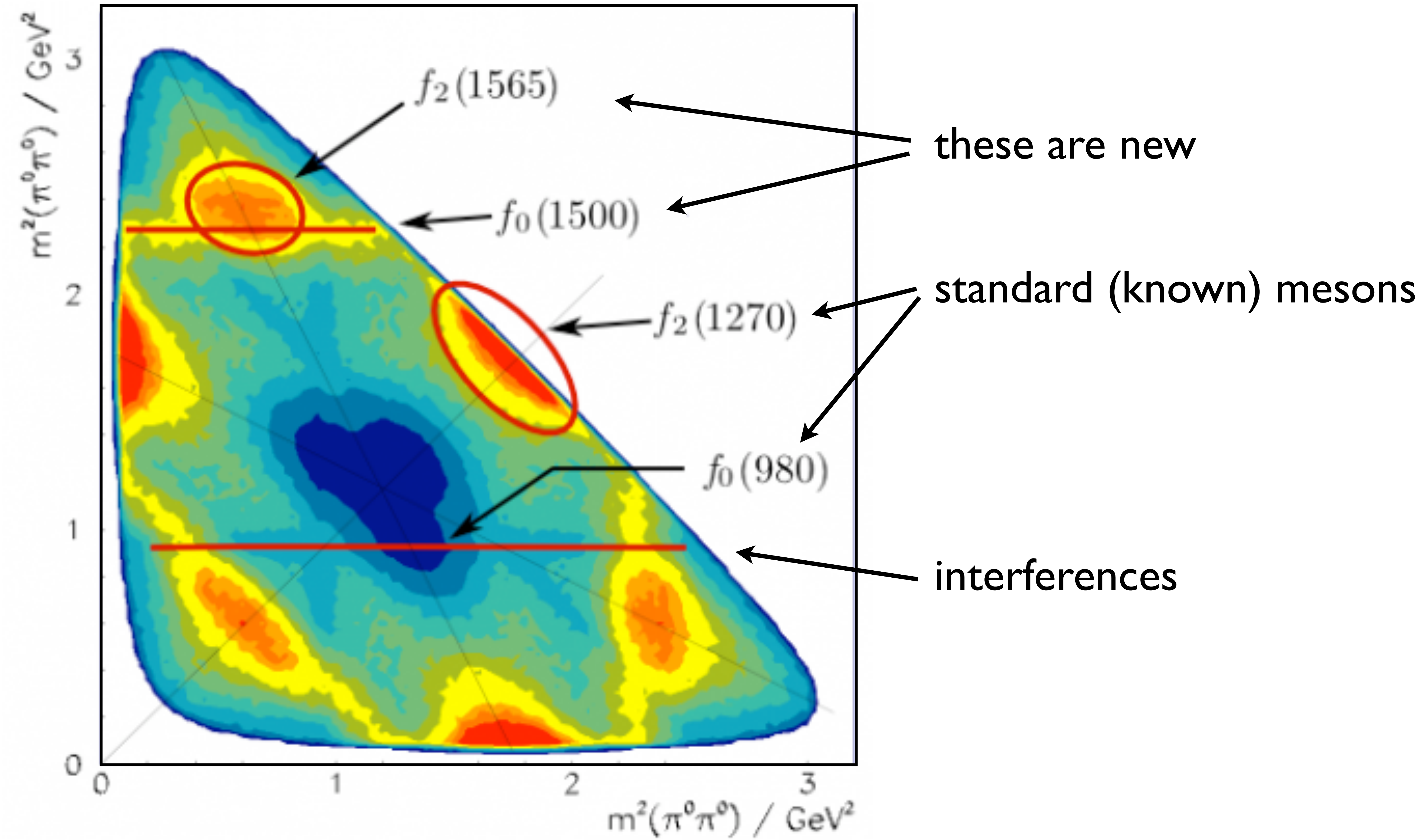
$p\bar{p} \rightarrow 3\pi^0$



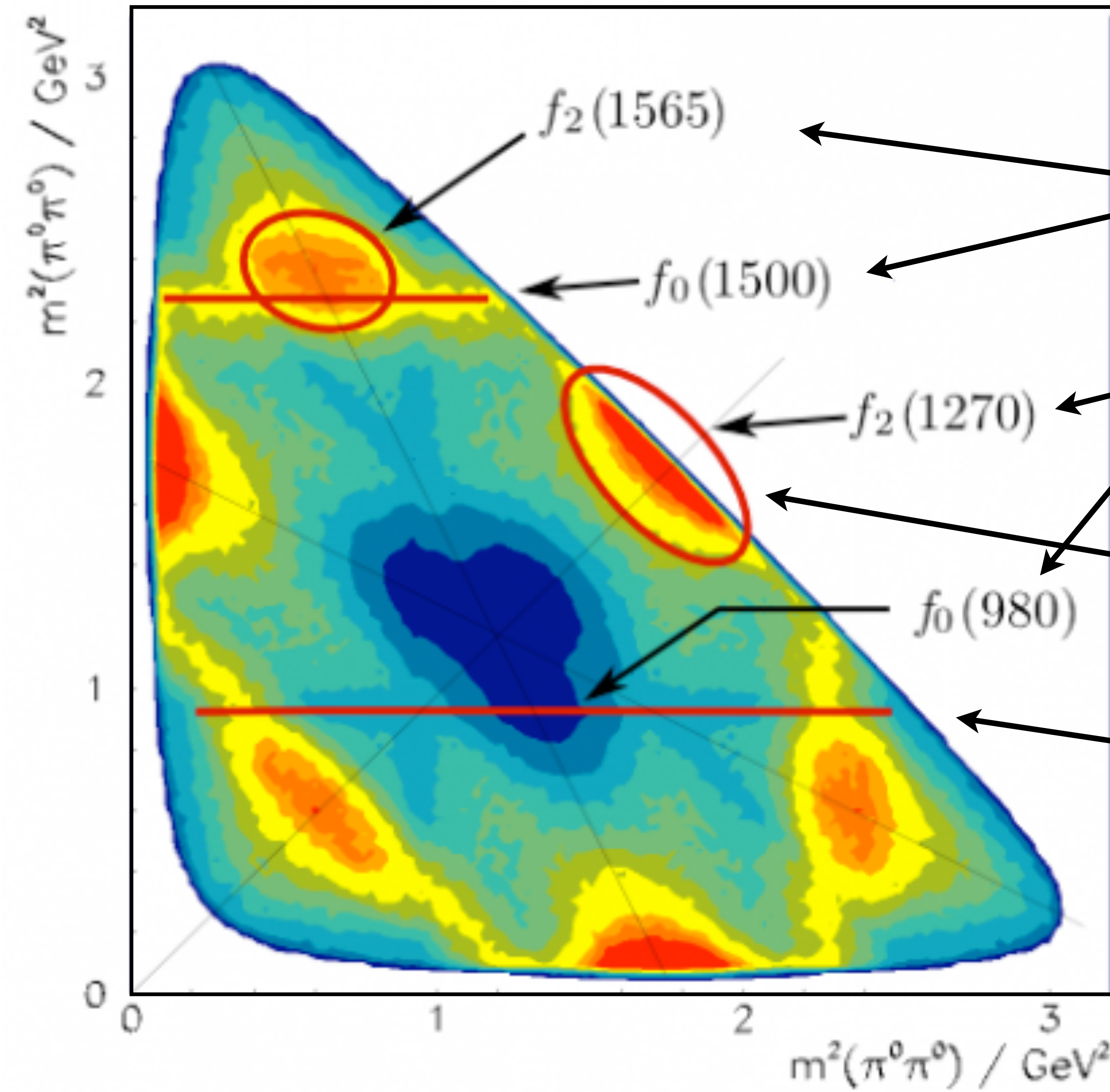
these are new

standard (known) mesons

$p\bar{p} \rightarrow 3\pi^0$



$p\bar{p} \rightarrow 3\pi^0$



these are new

standard (known) mesons

angular distribution ~ J^{PC}

interferences

Dalitz plot formalism

3-body decay of a spin 0 particle into pseudoscalars

$$\Gamma = \frac{1}{(2\pi)^3 32\sqrt{s^3}} \left|\mathcal{M}\right|^2 dm_{ab}^2\, dm_{bc}^2 \,,$$

kinematic factors (the first factor) dynamics ($|\mathcal{M}|^2$)

$|\mathcal{M}|^2$ constant = uniform population

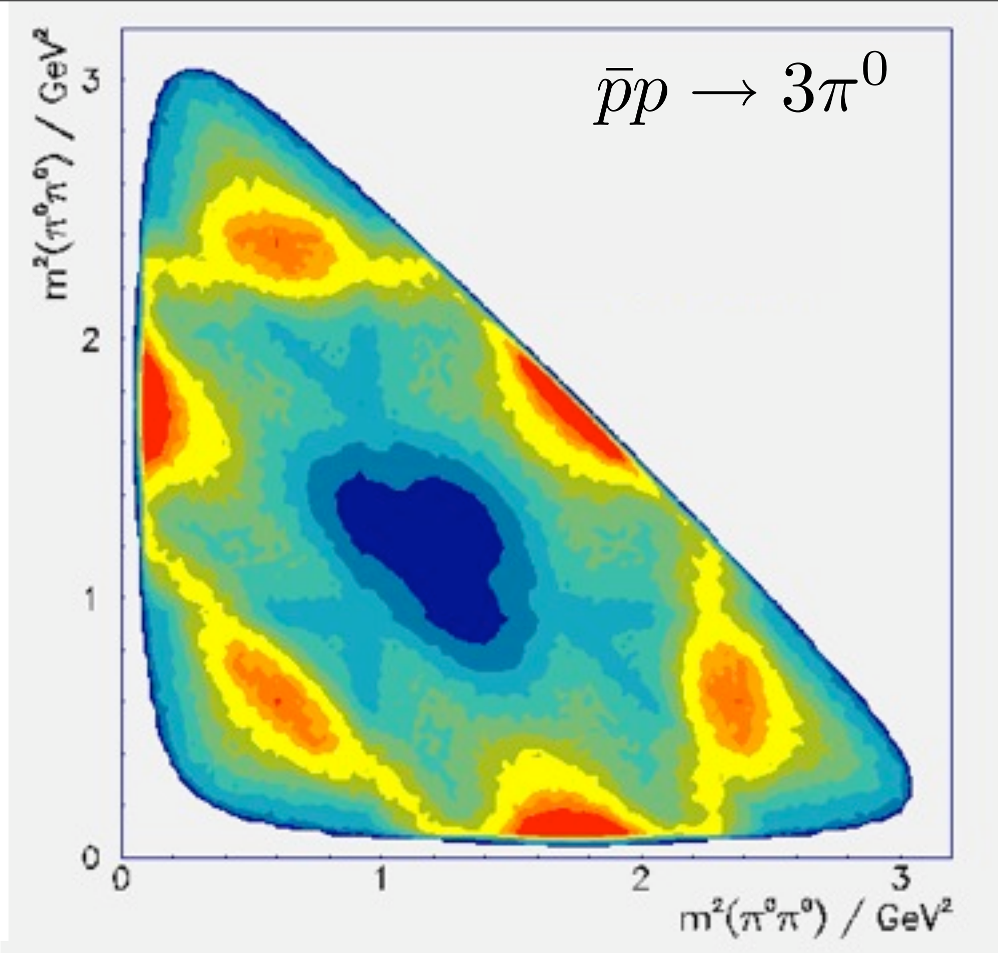
non-uniform population = dynamics

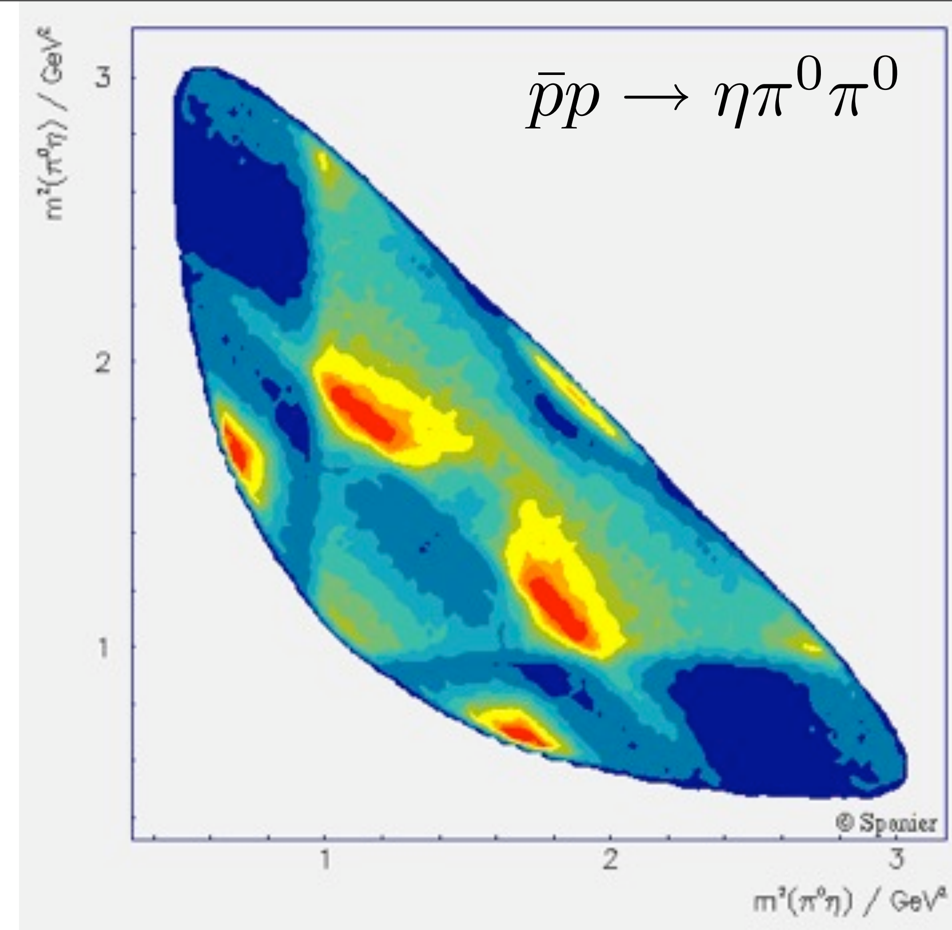
$R \to rc, r \to ab$

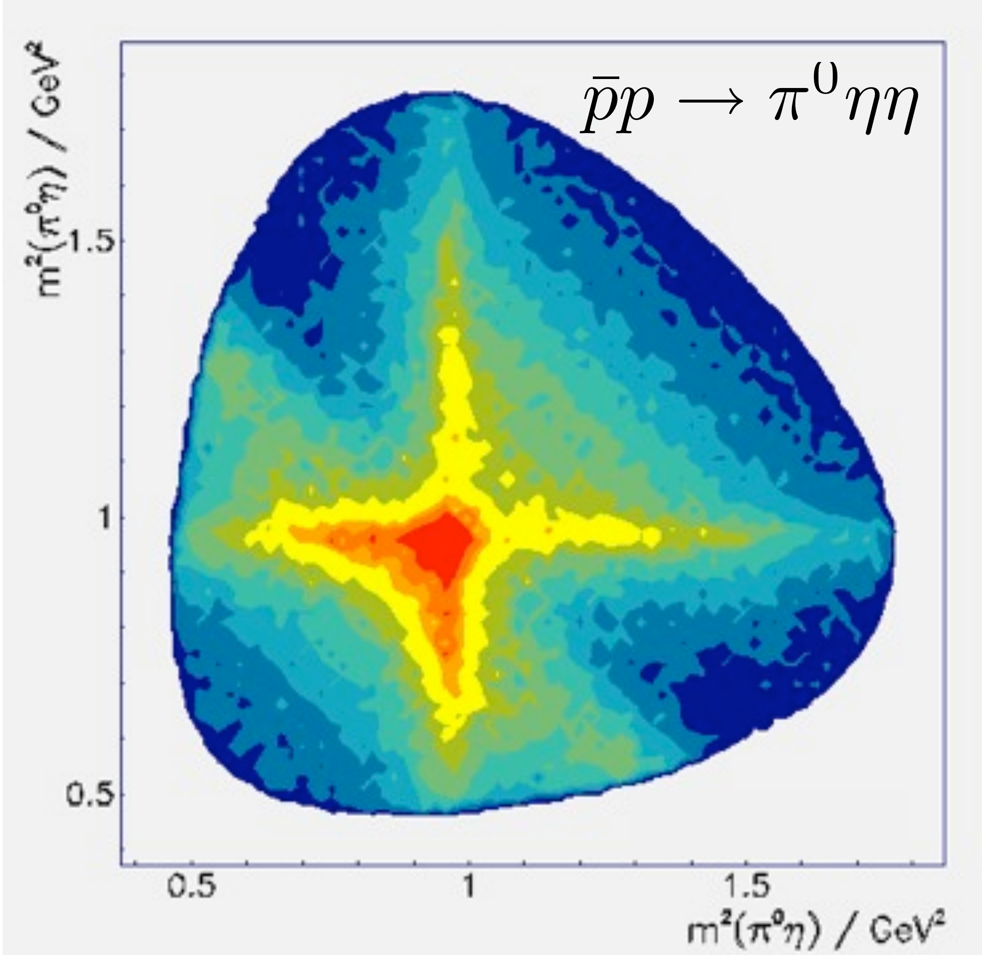
$$\mathcal{M}_r(J,L,l,m_{ab},m_{bc}) = \sum_\lambda \langle ab|r_\lambda\rangle\, T_r(m_{ab})\, \langle cr_\lambda|R_J\rangle$$

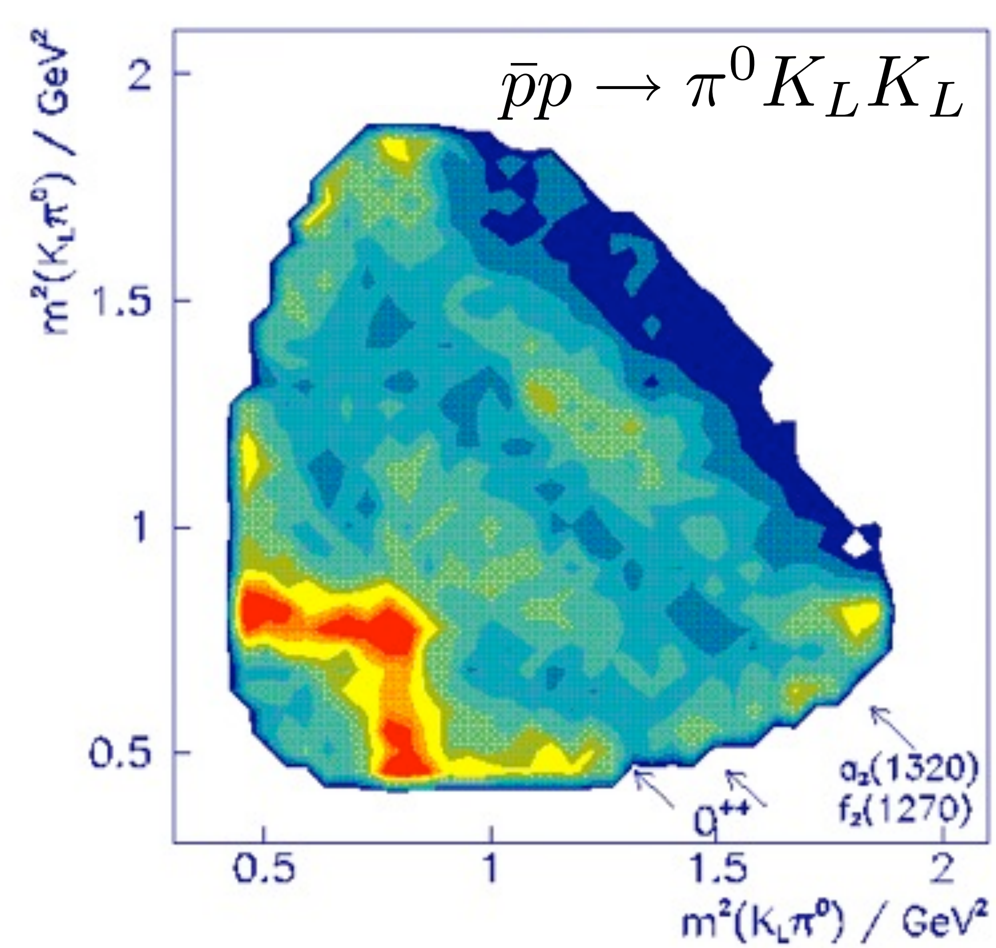
$$= Z(J,L,l,\vec{p},\vec{q}) B_L^R(|\vec{p}|) B_L^r(|\vec{q}|) T_r(m_{ab})\,.$$

- helicity states: λ (the sum over λ)
- angular distribution: Z
- momenta in r rest frame: $\vec{p}, \vec{q}$
- barrier factors: $B_L^R(|\vec{p}|)$, $B_L^r(|\vec{q}|)$
- dynamical function descr. resonance = Breit-Wigner or K-matrix or ...: $T_r(m_{ab})$


$\bar{p}p \rightarrow 3\pi^0$
$m^2(\pi^0\pi^0)$ / GeV^2
$m^2(\pi^0\pi^0)$ / GeV^2


$\bar{p}p \rightarrow \eta\pi^0\pi^0$
$m^2(\pi^0\eta)$ / GeV^2
$m^2(\pi^0\eta)$ / GeV^2
© Spanier


$\bar{p}p \rightarrow \pi^0\eta\eta$
$m^2(\pi^0\eta)$ / GeV^2
$m^2(\pi^0\eta)$ / GeV^2


$\bar{p}p \rightarrow \pi^0 K_L K_L$
$m^2(K_L\pi^0)$ / GeV^2
$m^2(K_L\pi^0)$ / GeV^2
0^{++}
$a_2(1320)$
$f_2(1270)$

Review of Particle Physics 2000

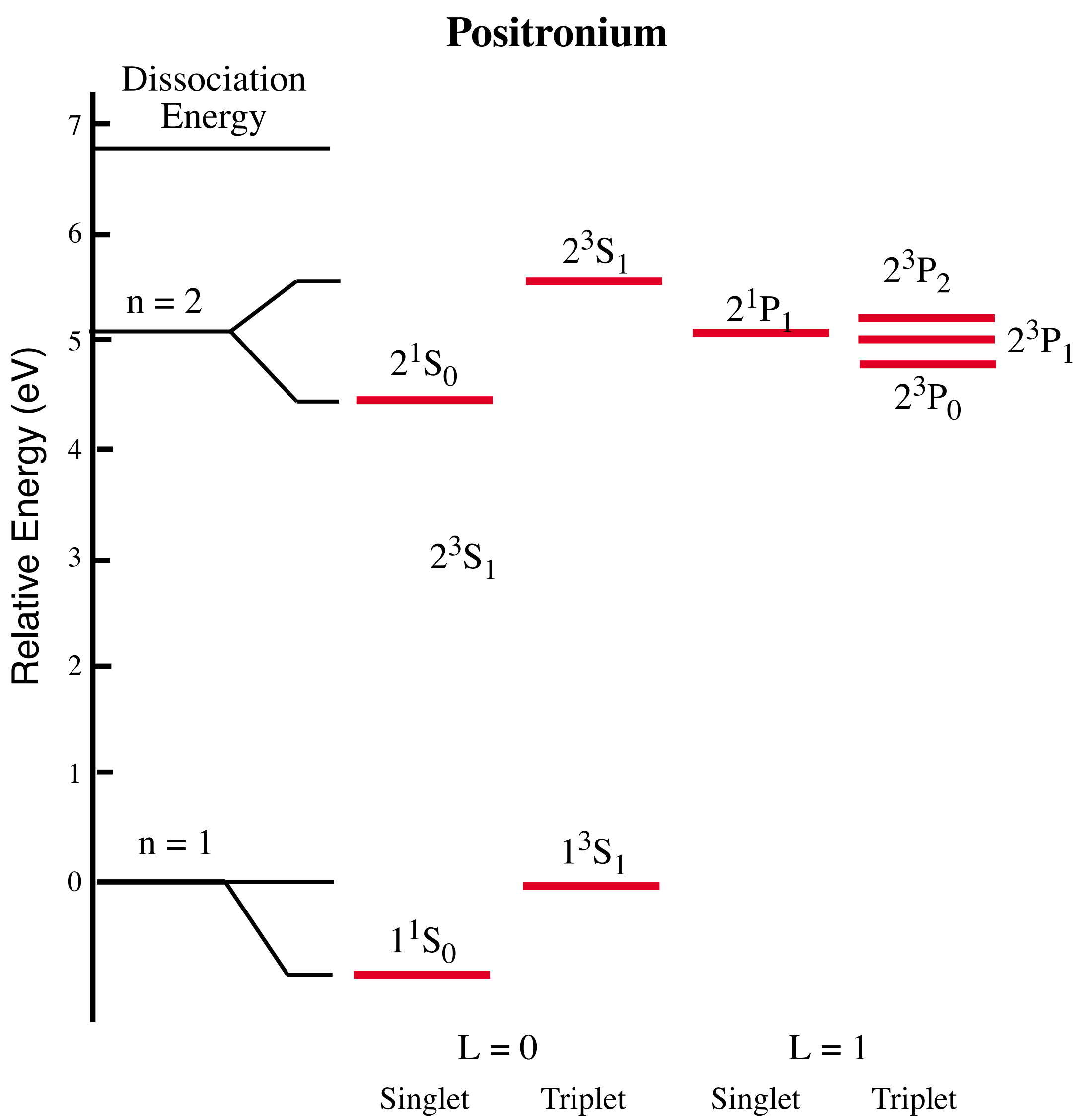
$N\ ^{2S+1}L_J$	J^{PC}	$u\bar{d}, u\bar{u}, d\bar{d}$ $I = 1$	$u\bar{u}, d\bar{d}, s\bar{s}$ $I = 0$	$\bar{s}u, \bar{s}d$ $I = 1/2$
$1\ ^1S_0$	0^{-+}	π	η, η'	**K**
$1\ ^3S_1$	1^{--}	ρ	ω, ϕ	**K*(892)**
$1\ ^1P_1$	1^{+-}	**b_1(1235)**	**h_1(1170)**, h_1(1380)	**K_{1B}**†
$1\ ^3P_0$	0^{++}	**a_0(1450)***	**f_0(1370)***, **f_0(1710)***	**K_0*(1430)**
$1\ ^3P_1$	1^{++}	**a_1(1260)**	**f_1(1285)**, **f_1(1420)**	**K_{1A}**†
$1\ ^3P_2$	2^{++}	**a_2(1320)**	**f_2(1270)**, **f_2'(1525)**	**K_2*(1430)**
$1\ ^1D_2$	2^{-+}	**π_2(1670)**	η_2(1645), η_2(1870)	**K_2(1770)**
$1\ ^3D_1$	1^{--}	**ρ(1700)**	**ω(1650)**	**K*(1680)**‡
$1\ ^3D_2$	2^{--}			**K_2(1820)**
$1\ ^3D_3$	3^{--}	**ρ_3(1690)**	**ω_3(1670)**, **ϕ_3(1850)**	K_3*(1780)
$1\ ^3F_4$	4^{++}	**a_4(2040)**	**f_4(2050)**, f_4(2220)	**K_4*(2045)**
$2\ ^1S_0$	0^{-+}	**π(1300)**	**η(1295)**, **η(1440)**	K(1460)
$2\ ^3S_1$	1^{--}	**ρ(1450)**	**ω(1420)**, **ϕ(1680)**	**K*(1410)**‡
$2\ ^3P_2$	2^{++}		f_2(1810), **f_2(2010)**	K_2*(1980)
$3\ ^1S_0$	0^{-+}	**π(1800)**	η(1760)	K(1830)

(shaded) contributions from LEAR experiments

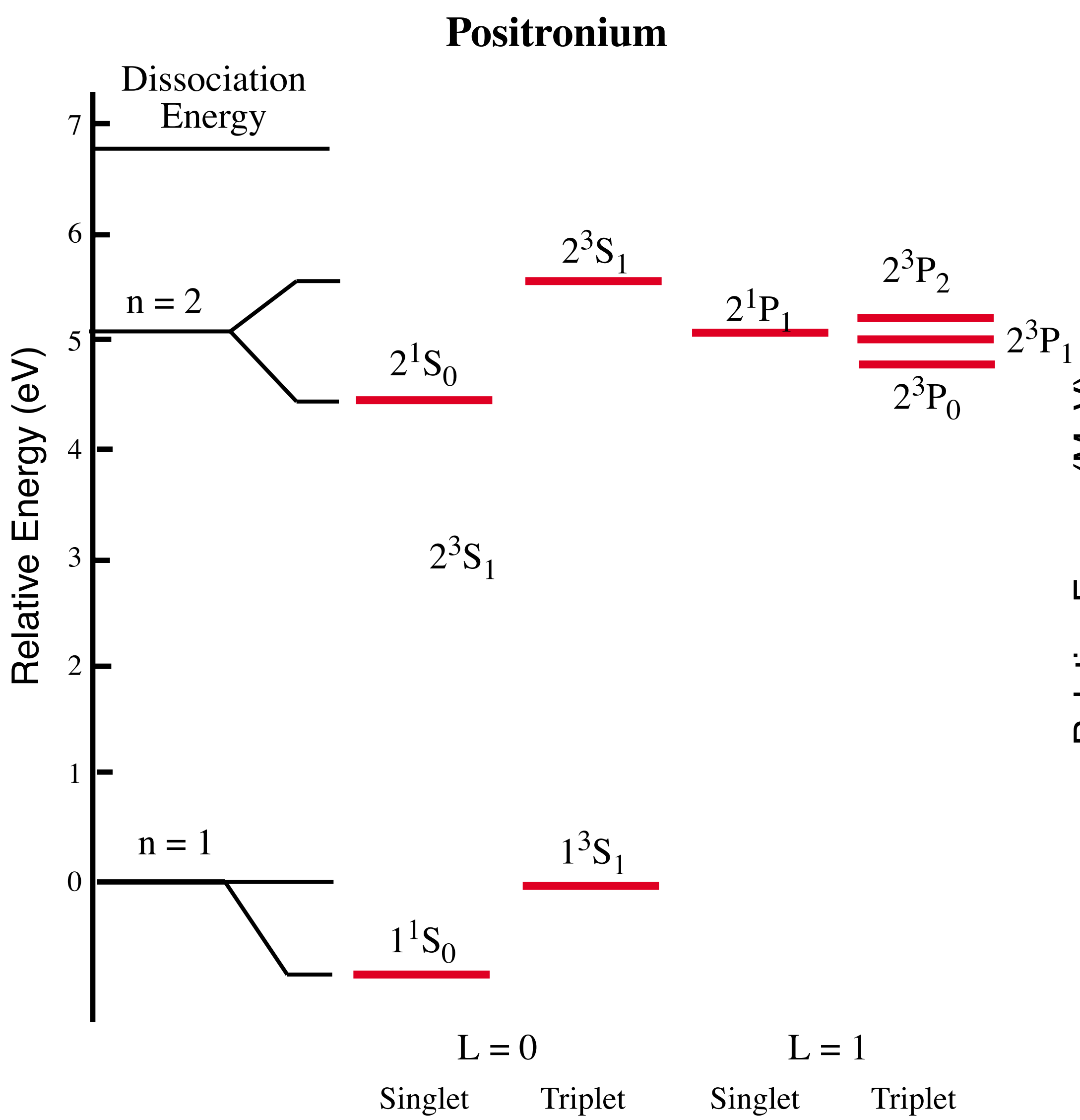
significant contributions, but:

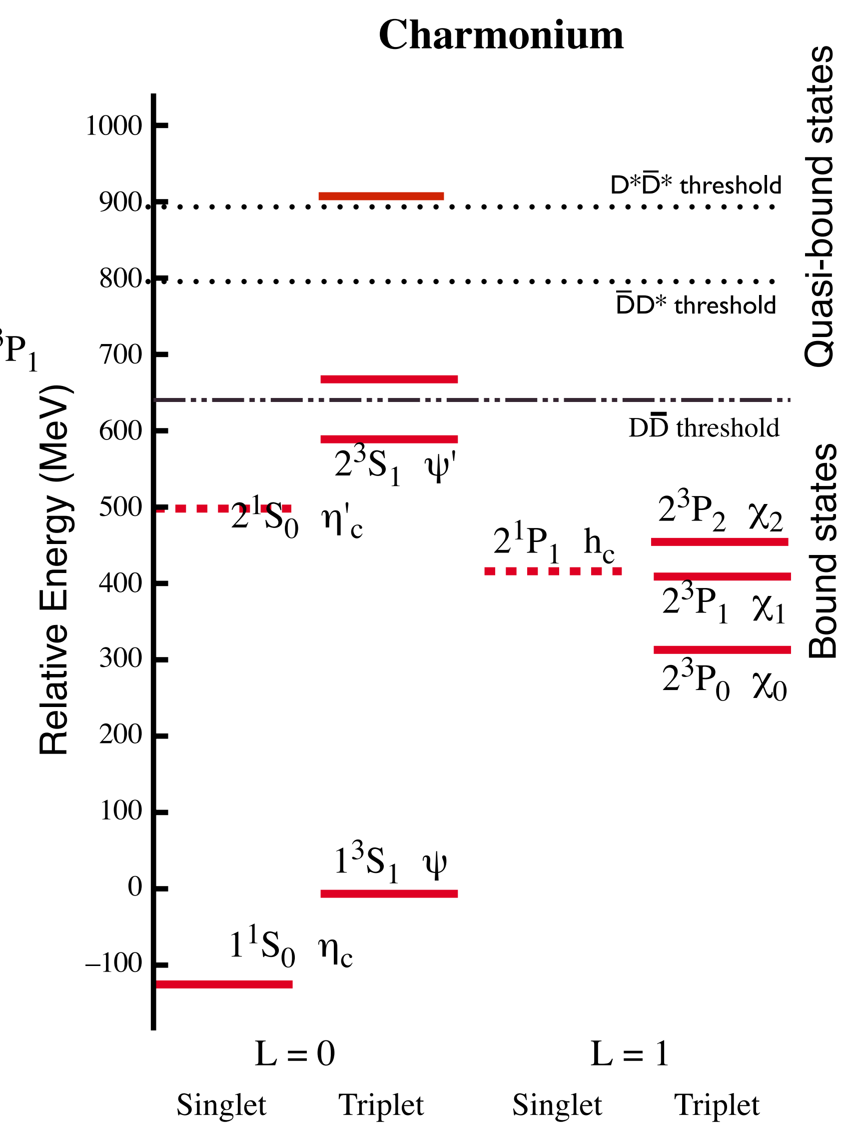
- mass range limited
- states are broad
- no good theory predictions
- need input from other production mechanisms

“cleaner” systems

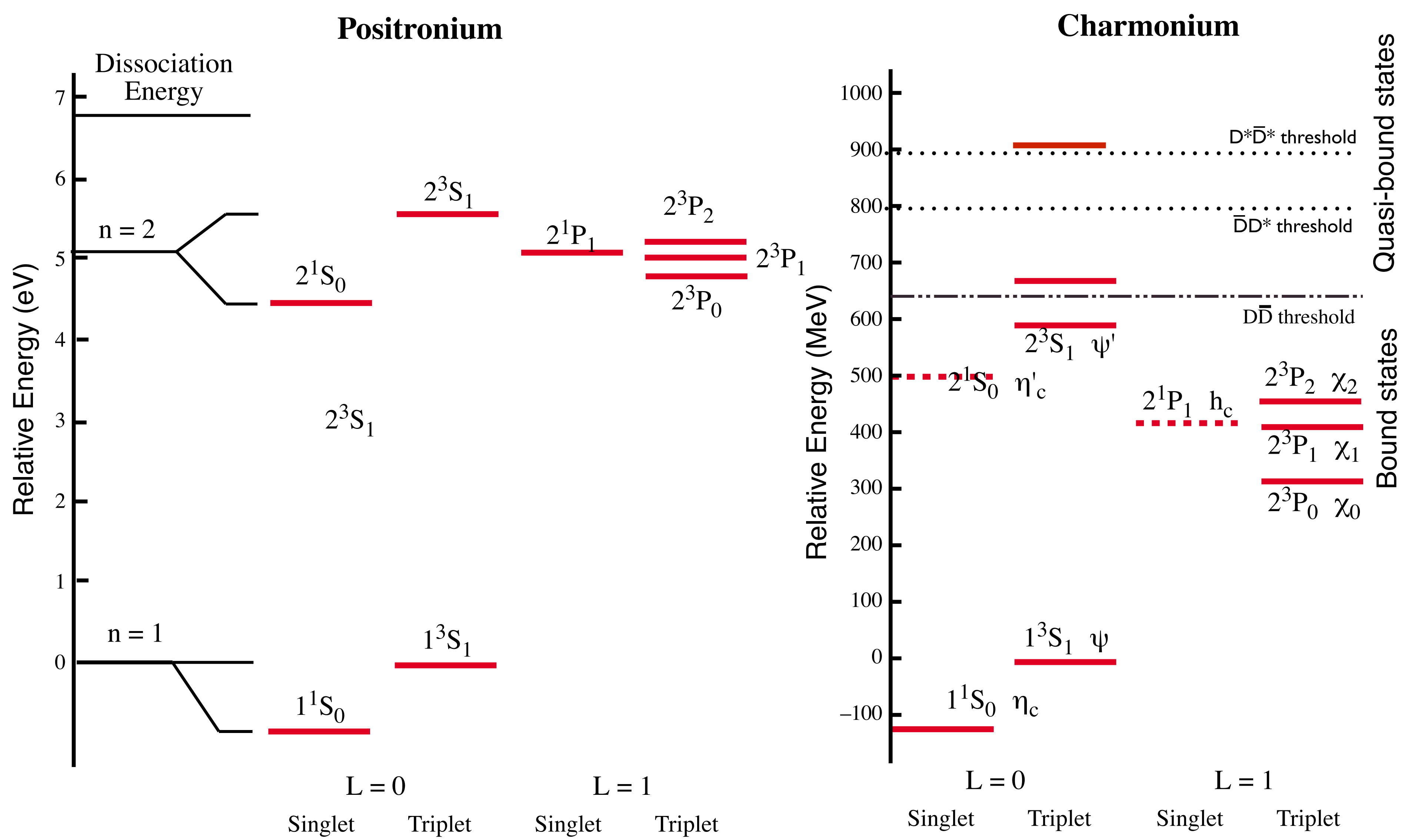


"cleaner" systems





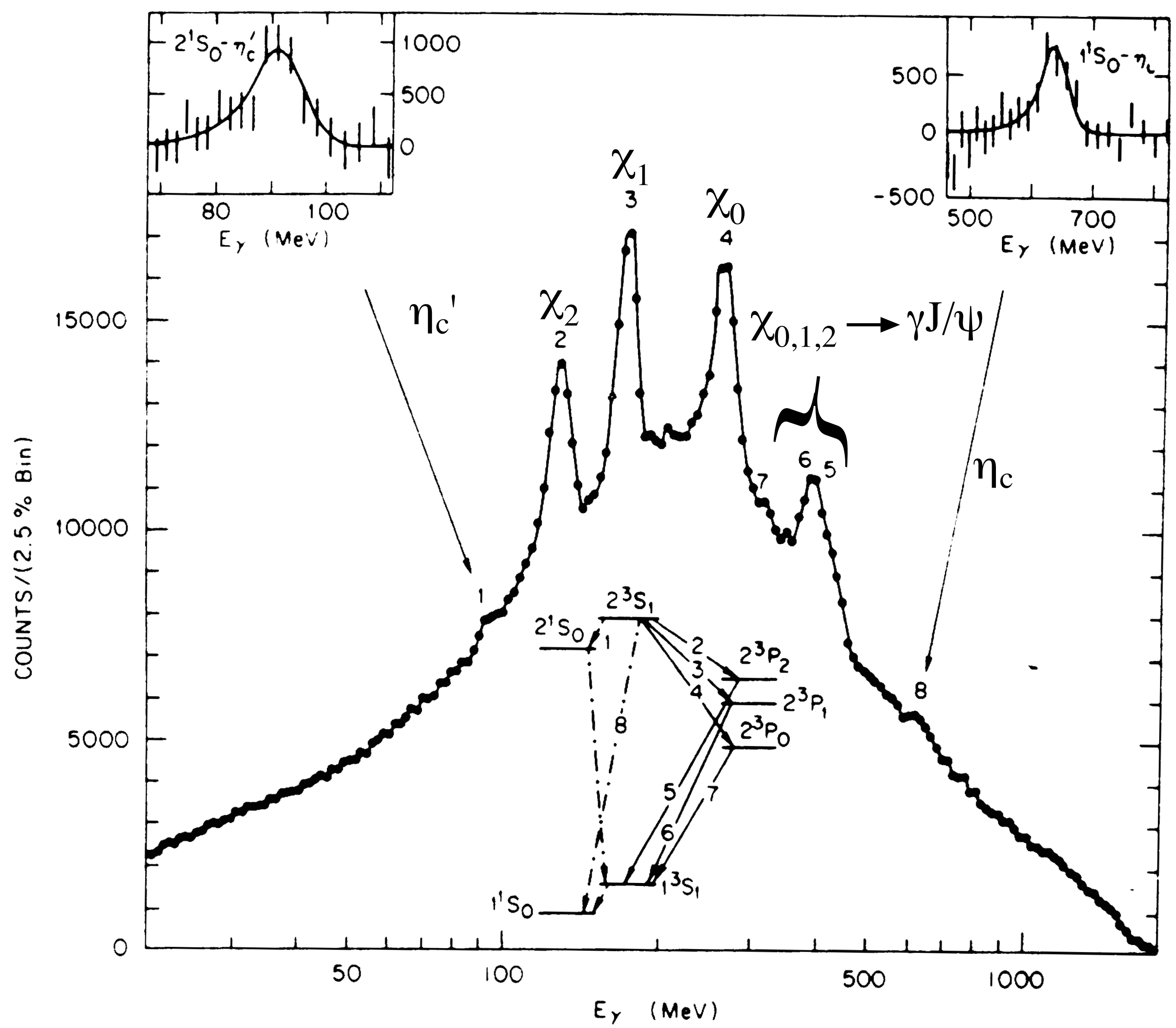
“cleaner” systems



charmonium is the positronium of QCD

Charmonium Spectrum

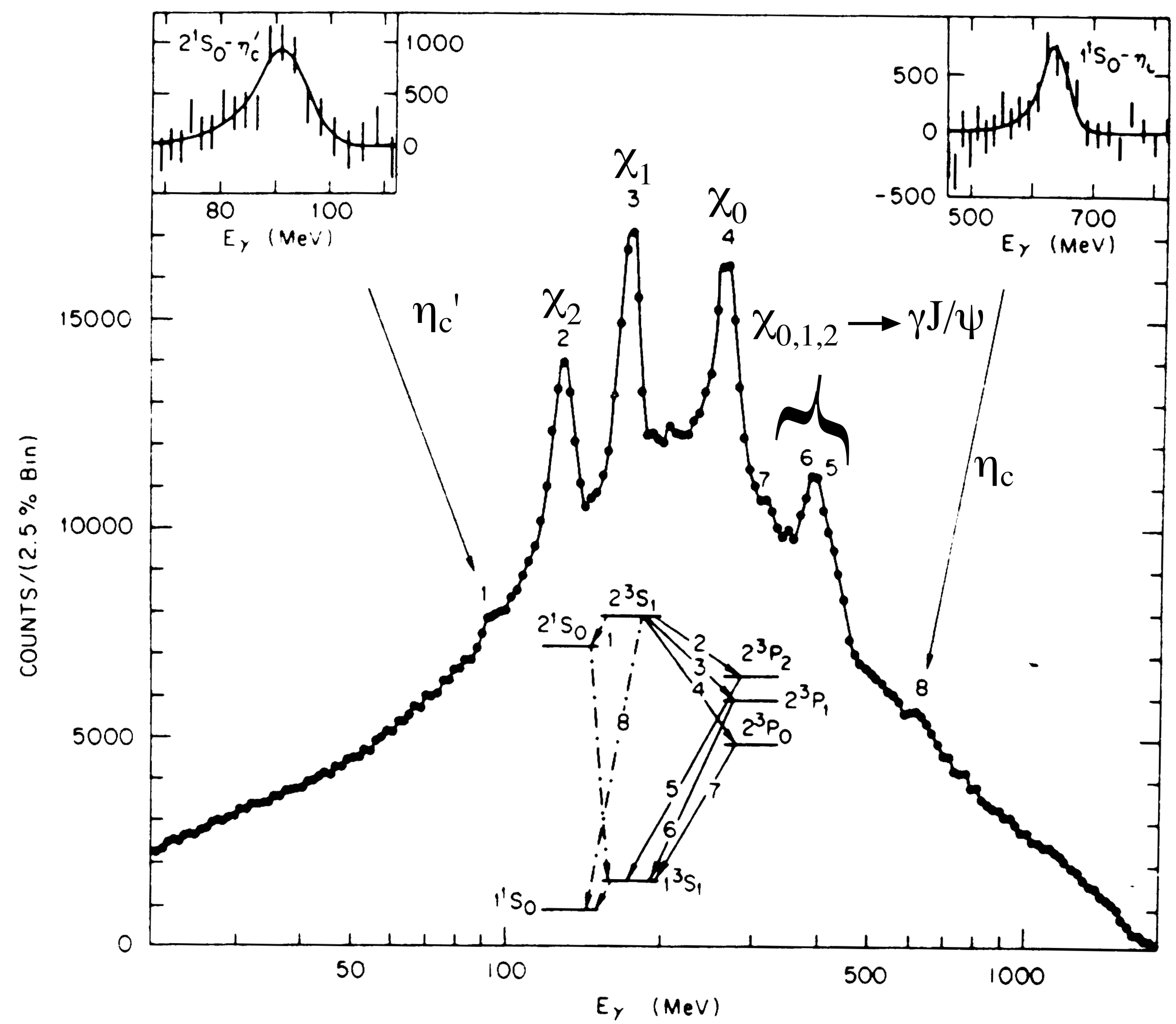
Crystal Ball (e^+e^- collisions)



Charmonium Spectrum

"atomic" spectroscopy of $c\bar{c}$ system

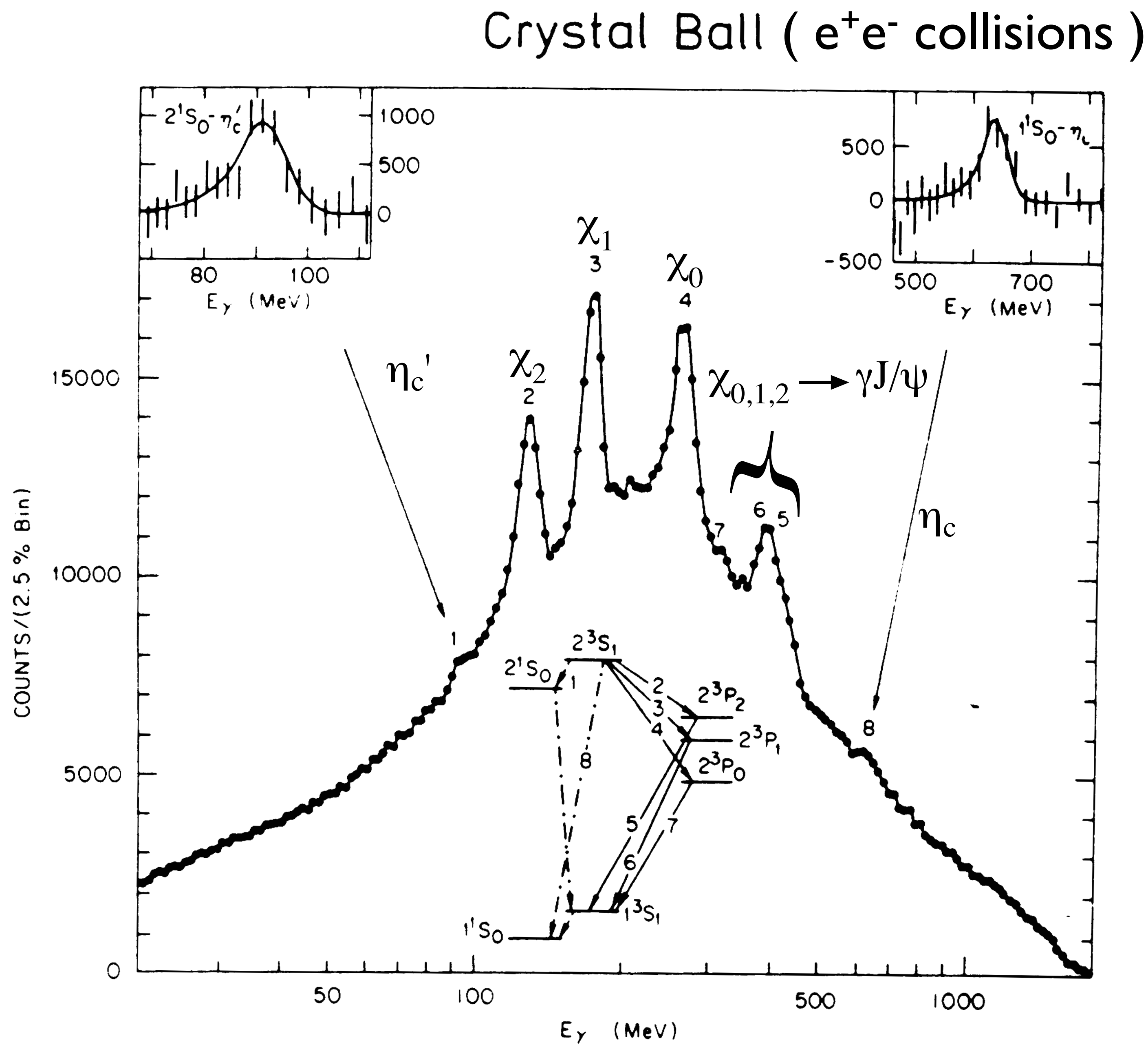
Crystal Ball (e^+e^- collisions)



Charmonium Spectrum

"atomic" spectroscopy of $c\bar{c}$ system

clean data but... picture is incomplete



Production:

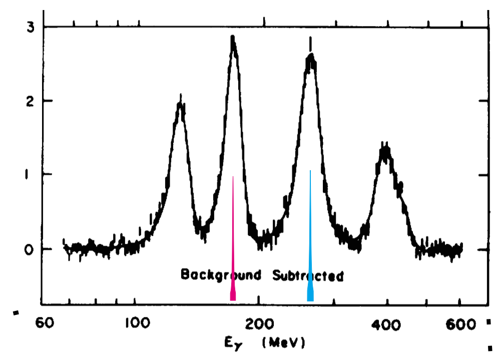
$e^+e^- \rightarrow \psi'$

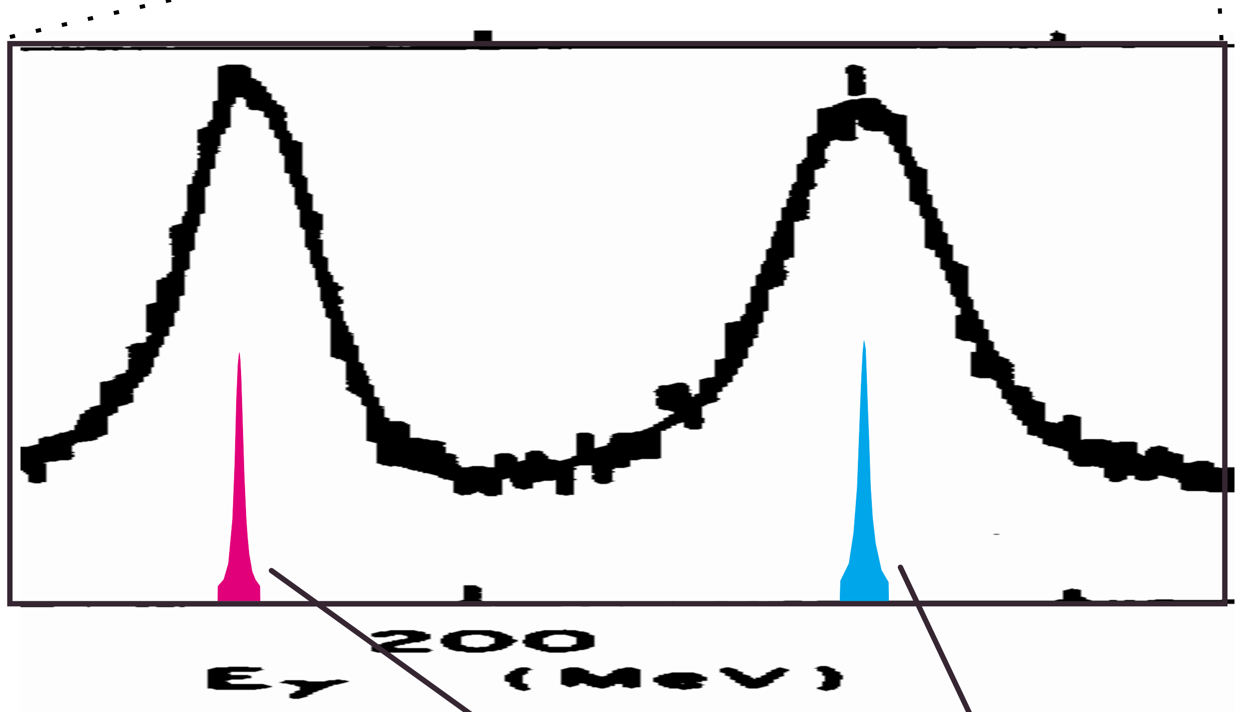
$\hookrightarrow \gamma\,\chi_{1,2}$

$\hookrightarrow \gamma\,\gamma\,J/\psi$

$\hookrightarrow \gamma\,\gamma\,e^+e^-$

Crystal Ball





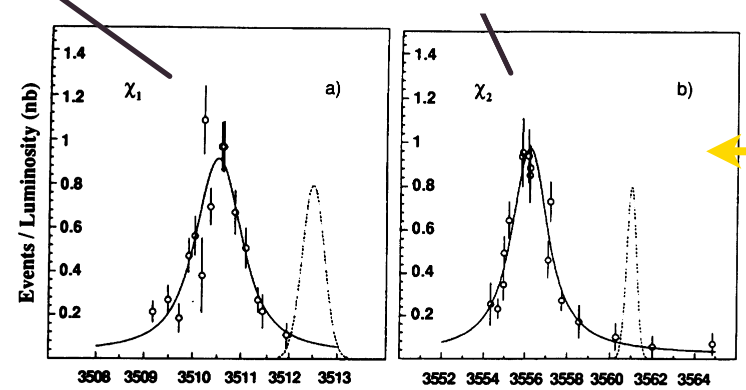
E 760 (Fermilab)

Formation:

$\bar{p}p \rightarrow \chi_{1,2}$

$\hookrightarrow \gamma\,J/\psi$

$\hookrightarrow \gamma\,e^+e^-$



σ_m (beam) = 0.5 MeV

Monday, July 30, 2012

Production:

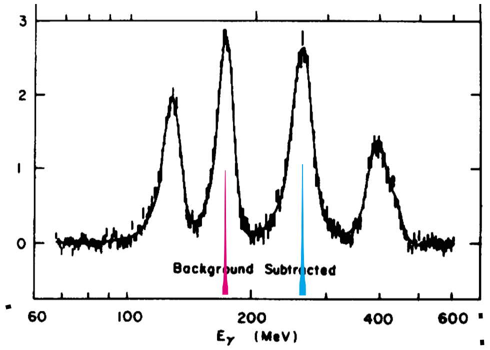
$e^+e^- \rightarrow \psi'$

$\quad \hookrightarrow \gamma\, \chi_{1,2}$

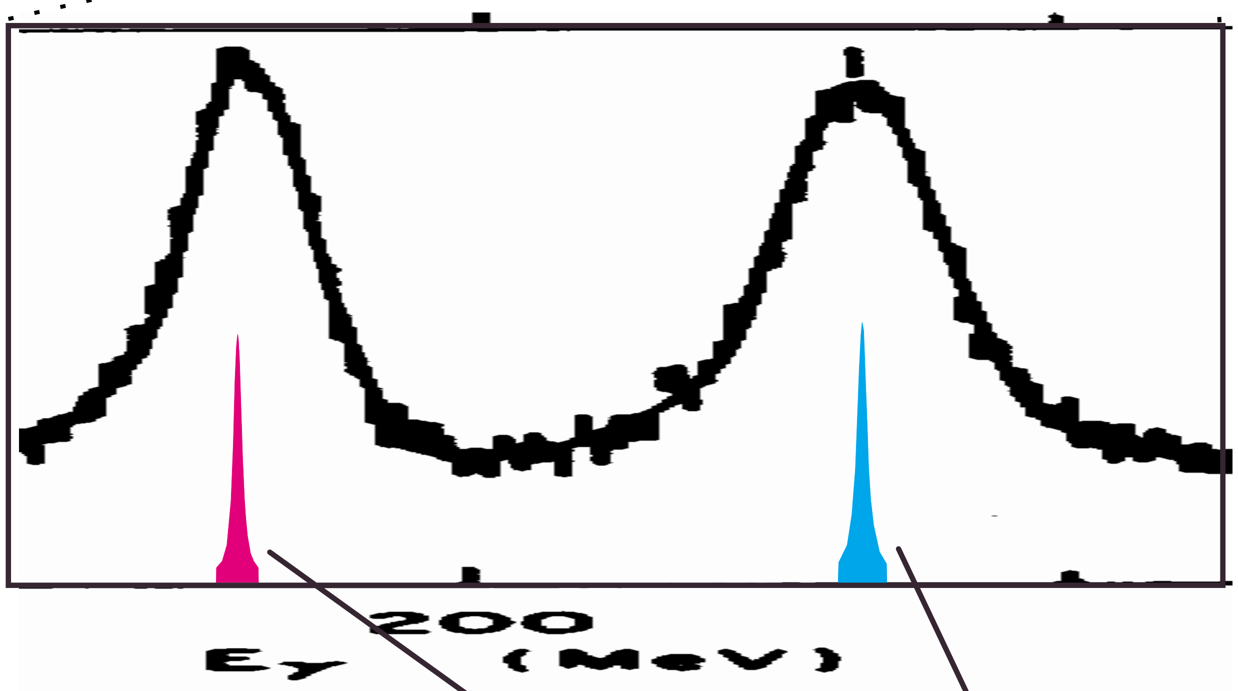
$\qquad \hookrightarrow \gamma\, \gamma\, J/\psi$

$\qquad\quad \hookrightarrow \gamma\, \gamma\, e^+e^-$

Crystal Ball



← resolution limited by detector



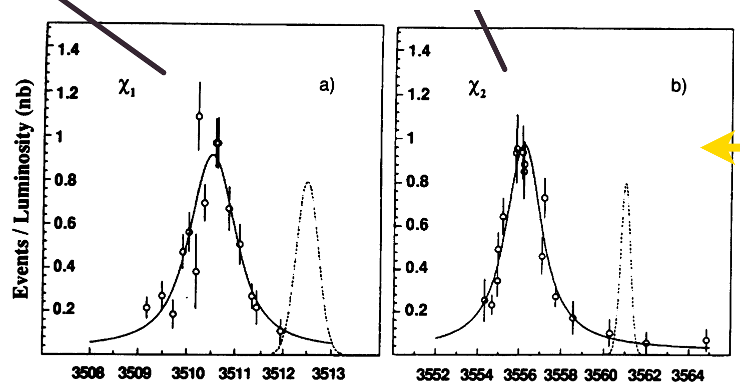
E 760 (Fermilab)

Formation:

$\bar{p}p \rightarrow \chi_{1,2}$

$\quad \hookrightarrow \gamma\, J/\psi$

$\qquad \hookrightarrow \gamma\, e^+e^-$



← resolution limited by knowledge of accelerator frequency

σ_m (beam) = 0.5 MeV

... in spite of many years of efforts, no clean understanding of low energy QCD. It is still a field with many open questions...

HEP however has mostly moved on ...

... in spite of many years of efforts, no clean understanding of low energy QCD. It is still a field with many open questions...

HEP however has mostly moved on ...

The end

(Actually, not really. Rather, the beginning: tomorrow, we go back to the Big Bang)