## Discovery and measurement of Supersymmetry at LHC

- Discovery of Supersymmetry
- Parameter measurements
- +connection to the cosmology
- Fulll reconstructions
- jets at LHC


## DM and collider signature



- "SUSY signature" "Models with new colored particles decaying into a stable neutral particle--LSP"
- "New physics" are migrated into SUSY category.
- Universal extra dimension lightest of first level KK is stable. .
- Little Higgs model with T parity. T parity in the model, T odd sector has stable particle ( $\mathrm{BH}_{\mathrm{H}}$ )
- Signal: assume mass difference is large

$\mathrm{E}_{\text {Tmiss }}>\max \left(100 \mathrm{GeV}, 0.2 \mathrm{M}_{\text {eff }}\right)$


## Background and discovery

CMS

- The typical number of SUSY events are $10^{5}$ for $10 \mathrm{fb}^{-1}$, while BG rate is $10^{9-8}$ for $\mathrm{W}, \mathrm{Z}$ and ttbar productions. $10^{-4}$ rejection of SM process is required.
- Understanding of the distribution is the key issue. For discovery
- $P_{T}$ distribution of the jets, $M_{\text {eff }}$ distribution. (theoretical complexities)

- $E_{\text {tmiss }}$ distributions (Experimental complexities)


## signal and background separation

top partner pair


( $\mathrm{T} \overline{\mathrm{T}}$ production)
-Two particles produced at threshold
-Two DMs can escape same direction

- Higher ETmiss/Meff

Typical signal structure when 2DM's is in a event !!!!

$i^{v}$ jets SM particles are boosted for large Meff

- neutrinos are back to
back
- lower Etmiss /Meff

From Kanay’s Slide in SUSY06

## Discovery Potential

5 -sigma discovery potential on $m_{0}-m_{1 / 2}$ plane
 $\mathrm{m}^{1 / 2}$


Only statistical error is included.

- Backgound is estimated by Alpgen.

Fast simulation result
Signal : Isawig/Jimmy
Background : Alpgen

- O-lepton mode : More statistics is available.
- 1-lepton mode : Relatively smaller background uncertainty. Major background is tt(+njets) is comparatively predictable.


## The "discovery reach" depends on "MSUGRA" assumptions.

choice for large $\mu$

- Ex. KKLT (or MMAM mixed modulas anomaly mediation model ) If both volume modulas $T$ and compensator C contribute to the SUSY breaking.

$$
M_{a}=\left(\frac{l_{a}}{R}+\frac{b_{a g} g_{G T}}{16 \pi^{2}}\right) m_{3 / 2}
$$

- mass spectrum can be quite degenerated. Change FT/FC, MSUGRA $\rightarrow$ UED like $\rightarrow$ AM
- The most degenerate spectrum corresponds to mixed dark matter. Consistent to $\Omega \mathrm{M}$



## SUSY at LHC in degenerated limit

- degenerate $S U S Y=$ lower $P_{T}$ jets, small $M_{\text {eff, }}$ Small missing energy. Discovery gets difficult (no chance if all masses are same of couse )
- Need to take into account the background seriously. $\mathrm{S} / \mathrm{N}<1$, discovery is in ?? because of the background uncertainty

- Recent preliminary ATLAS simulation (by S. Okada +et al Kobe group including QCD, W, Z, ttbar background with ME collections) confirms the same tendency. Discovery is extremely tough if $m$ (LSP)>0.7 M(squark.



## Significance (0-lepton mode)


S. Okada et al (Kobe group very preliminary

We need to study more on degenerate cases

Measurement of basic parameters

## SUSY parameter measurement A brief history

- early 1990
- JLC study: define LC as the machine to measure SUSY parameters, spin, and interaction. check GUT relation $M_{1}: M_{2}$
- $\mathrm{LHC}=$ a discovery machine.


Tsukamoto et al '93

- Snowmass 1996:
- Trying to establish US participation at LHC, "(ex-) Theorists"(Hinchliffe, Paige, ...) took LC concepts. Techniques for mass reconstruction were established at that time.
- ILC: SUSY coupling measurements ('96 Nojiri et al , Feng et al....): physics point of LC over LHC


Nojiri et al '96

## measurement at LHC a check list

- In the past ILC study, emphases were on the measurements on Supersymmetric parameters. Let check how much LHC can do.
- mass
- MSSM parameters
- spin
- SUSY couplings
possible using end point method and transverse parameters (such as MT2

Some guess using branching ratios
charge asymmetry and cross section
not yet

## SUSY parameter measurements


$\mathrm{m}(\mathrm{jll})$ with $\mathrm{mll}>0.5 \mathrm{mll}(\max )$


LSP
 ee $+\mu \mu-e \mu$ subtraction is effective to select single channel

Hinchliffe et al (97)


## Spin Effect (fermion ino + chiral interaction )



## Best Case at LHC



- Evidence of multiple ino give our access to ino masses, thus $\mu$ and $M_{-} i$
- squarks and gluino work as the source of EW sector of Supersymmetry.


# Summary in SPS1a (most lucky case) 

from LHC/LC study

| particle | mass | error(low) | error(high) |  |
| :---: | :---: | :---: | :---: | :---: |
| gluino | 595 | 16.3 | 8.0 | bbll |
| squark(L) | 540 | 21.2 | 8.7 | jll |
| squark(R) | 520 | 17.7 | 11.8 | $\mathrm{M}_{\text {T2 }} 10 \mathrm{GeV}$ sys |
| $\tilde{\chi}_{4}^{0}$ | 378 | 14.6 | 5.1 |  |
| $\tilde{\chi}_{2}^{0}$ | 177 | 13.4 | 4.7 |  |
| $\tilde{\chi}_{1}^{0}$ | 96 | 13.2 | 4.7 |  |

- LSP mass error is large, but mass differences are known precisely. There are correlated overall error to the mass scale.
- Access to 2 or 3 neutralino mass, information on 2 of 4 LSP parameters
- selectron and smuon mass error is about same to that of LSP


## Connection to Cosmos.

At LHC
Collider physics meets Astrophysics

Discovery of DM at LHC, Parameter measurements

Thermal relic density
prediction of $\sigma(\mathrm{XN} \rightarrow \mathrm{XN})$ or $\sigma(\mathrm{XX} \rightarrow \mathrm{SM})---r, \mathrm{e}^{+}$flux

if it is there......


Direct detection of DM, and cosmic ray measurements


PAMERA, GLAST, CDMSII etc..

Astro/cosmo observations

## DM density control parameters

1)bulk: LSP=Bino like.

Slepton exchange
$\Omega h^{2} \propto m_{\tilde{l}}^{4} / m_{\tilde{\chi}}^{2}$ too large mass density
2)Higgs pole effect $\mathrm{m}_{\boldsymbol{H}}=2 \mathrm{~m}_{x}$
3) coannihilation $\tilde{\tau} \tilde{\chi}$
4)focus point region:
higgsino-gaugino mixing


Gaugino mass

Dominant uncertainty of ino mixing angle comes from $\tan \beta$


## Importance to measure tau mode.


tau -fake tau

- If tau coannihilation is on, the mass density is very sensitive to the stau LSP mass difference.
- Due to the left right mixing, 2nd lighest neutralino tend to go to tau mode. finding edges exclude the possibility of stau LSP coannihilation.
- tau jets are identified as narrow jets. Tau is experimentally difficult because QCD fake tau + particles converted at inner trackers. $41 \%$ tau tagging efficiency with significant background.
- we took 62 GeV \pm 5 GeV for this point.


## Importance to get lower limit on MA

Assumes all SUSY particle is heavy

The limit from non -SUSY production for $\tan \beta=10$ is around 250 GeV

We also have heavy Higgs from SUSY production with significant cross section

- $m(A / H) \leq 315 \mathrm{GeV} \quad$ for $\tilde{\chi}_{4(3)}^{0} \rightarrow \tilde{\chi}_{1}^{0} A / H$
- $m(A / H) \leq 230 \mathrm{GeV} \quad$ for $\tilde{\chi}_{4(3)}^{0} \rightarrow \tilde{\chi}_{2}^{0} A / H$ and $\tilde{\chi}_{2}^{ \pm} \rightarrow \tilde{\chi}_{1}^{ \pm} A / H$.

- We may have an access to heavy higgs sector up to 300 GeV also for tanbeta $=10$ for this point .


## Trying to pin down Dark matter nature

- DM density: for SPS1a
- roughly $20 \%$ at LHC. The plots are based on reconstruction of O (10000) different experiments-Giacomo likes this....
- LSP-N cross section. almost no bound for LHC. Disregard fake peak....
- It is certainly not general, no conclusions.. at Les Houches, we have selected several points to cover NUHM case (mixed DM)

MN, Polesello and Tovey hep-ph/0512204


Probability density

$\mathrm{P}-\chi$ scattering cross section(pb)
Baltz et al (2006)

## Full reconstruction at LHC

## Solving missing momentum



If we know all masses, there are 5 mass constraints for LSP momentum, therefore event can be fully solved.

If we do not know any of those masses, each event gives you 4 dim hyper surface in the 5 dim
 mass space. (one constraint. )

5 events=> all masses
For presentation, assume we know mass of

$$
\chi_{1}^{0}, \chi_{2}^{0}, \tilde{l}
$$

## One event $\Rightarrow$

## probability density for true masses(logL)

 from expected $b$ jet smearing
## Gluino mass






Gluino-sbottom
$\log L(1)+\log L(2) \quad+\log L(3)+\quad \log L(4)+\ldots$
$=\Sigma \log L\left(\sim \Delta \chi^{2}\right)$
(a) $\tan \beta=10$, all


## Other application

(Tovey, Polesello, very very preliminary )

idea
End point method does not have much constraint to the overall mass scale

Look for 4 leptons events, with two golden cascades, when ETmiss can be compared with solved ETmiss for the assumed masses.

This may be more sensitive to the overall mass scale.

- Error of LSP improves from $8.0 \%$-> $6.5 \%$


## jets! at LHC

What hell we can if your signals are all jetsExample: study of "top-partners".

- Lepton + jets are easy because we study the events with relatively small number of jets and jets are isolated. however, there are important process with many jets in the same direction.
- in MSUGRA, stops are lighter than other squarks want to establish gluino-> stop cascade decays
- in Little Higgs model, top partner is again the lightest among quark partners. The cross section is high because it is fermion. Want to establish top partner productions


## Little Higgs model with T parity

- fermion partner instead of scalar partner
- top partner is important
- $\sigma($ boson $) / \sigma($ fermion $)=0.1$
- The difference comes from spin structure. scalar top production is "mostly" p wave. Evidence of top partner pair production immediately means non-SUSY BSM


## After jet reconstructions

 reconstruction with jet cone $\mathrm{R}=0.4$ In AcerDET(<ATLFAST)

## Top reconstruction

Signal, without b tagging

largest 2 jet wiswriant mass in the heminsphare


- hemisphare algorithm by Moortgat (CMS)
- take highest PT jet as seed of an axis. (A)


## $\downarrow$

- take 2 nd jet with max $P \times \Delta R$ from the 1st jet as the seed of the 2nd axis (B)
- assign jet and leptor activities to the "closer axis". (C)

$$
\downarrow \hat{\imath}
$$

- recalculate "axis=sum of particle in the hemisphere" , repeat. (D,E)
Note:the jet energy resolution is better in ATLAS $(50 \% / \sqrt{ }$ E
is better than CMS $(100 \%) / \sqrt{ } E$
Matsumoto Nojiri Nomura 2006


## Signal and BG




- tops are seen in both of the hemisphere
- probability of top reconstruction is small for the ttbar background (because of $\mathrm{E}_{\text {Tmiss }} \mathrm{cut}$ )
- top reconstruction helps to reduces QCD, W, Z background. (preliminary ATLAS simulations by Kiyamura at Kobe


## ATLAS simulation



- top reconstruction is actually useful to reduce W, and $Z$ backgrounds (Kiyamura et al Kobe very preliminary)


## signal distribution \& top background



- signal $\mathrm{E}_{\text {Tmiss }}$ distribution has a peak near $\mathrm{Meff}_{\text {ef }}$ /2
- BG peaks at $\mathrm{E}_{\text {tmiss }} \ll \mathrm{M}_{\text {eff }}$
we see DM!
- good margin for discovery due to the bump structure.


## top quark in SUSY non relativistic top quark

- top quark in SUSY

$$
\tilde{g} \rightarrow(t \tilde{t} \text { or } b \tilde{b}) \rightarrow t b \tilde{\chi}_{1}^{ \pm}
$$

- reconstruct tb end points tell us stop and sbottom

- Background to $t \Rightarrow b W \Rightarrow b j j$ is estimated from events in the sideband

$$
\begin{aligned}
& \mathrm{m}_{\mathrm{jj}}<\mathrm{Mw}-15 \mathrm{GeV} \\
& \mathrm{~m}_{\mathrm{jj}}>\mathrm{Mw}+15 \mathrm{GeV} .
\end{aligned}
$$

rescale the jet energies so that they are in W mass range. look for the consistent samples with $\mathrm{t} \Rightarrow \mathrm{bW} \Rightarrow \mathrm{bjj}$

(Hisano, Kawagoe, Nojiri 2003)

## KT vs Jet cone in stop reconstruction

- Measure the efficiency of clustering algorithm in New Physics signal simulation (it must contain truth. )
- high reconstruction efficiency of tb end point ( height) ~efficiency
- small $R$ is better because we expect many jets in the final state. $\mathrm{R}=0.4$.
- Low efficiency in simple jet cone(in ATLFAST at that time). overlapping? or infrared unsafety?
- KT works OK for small R. without underlying events. "Splash in effect"

$R$ end point height

|  |  |  |
| ---: | :---: | :---: |
| HW cone 0.4 | $434.5 \pm 5.8$ | $354.8 \pm 23.3$ |
| 0.5 | $460.2 \pm 4.9$ | $349.2 \pm 22.8$ |
| 0.6 | $440.9 \pm 7.1$ | $305.3 \pm 33.7$ |
| $K_{T} 0.4$ | $434.9 \pm 4.3$ | $406.5 \pm 22.1$ |
| 0.5 | $460.0 \pm 5.5$ | $379.6 \pm 33.5$ |
| 0.6 | $468.4 \pm 5.8$ | $314.3 \pm 20.7$ |

Hisano Kawagoe Nojiri (2003)

## Comparison between KT and cone

- To deal with jet-jet system, the definition of "jet" matters. How can we use measured end points to the mass spectrum?
- Cone--Simple cone alg. takes the hardest PT cluster, add the activities nearby the cluster. (Some adjustment for overlapping jets) infra unstable except SISCone.
- KT merges soft collinear activity to nearby hard ones. (infra safe) use $\min \left(k t k^{2} k t_{j}{ }^{2}\right) \times R_{k j}{ }^{2} / R^{2}$
- cf. Camblidge KT Use $\mathrm{R}_{\mathrm{ij}}$ to merge particles ( motivated by angular ordering.....)



## merged jets?


$\ldots p_{T} \sqrt{y} \ldots \ldots . \equiv d_{k l} / \cdot\left(p_{T^{n}}^{\mathrm{jet}}\right)^{2}$

- Butterworth et al looked into jet with mass $\sim 80 \mathrm{GeV}$ from $W$ with $\mathrm{P}_{\mathrm{T}} \sim 200 \mathrm{GeV}$ in KT .
- the subjet structure is defined as the last marged jet. This may be used to separate QCD jets and marged jet (like relativistic W in SUSY )


## New physics with LHC

- I do not care much about "inverse problem".
- Want to add more stable quantity that can be used at LHC, ..... toward something which has same taste of "what is expected at ILC".
- We can learn lots from lepton mode at LHC, but we need more improvement on understanding jet signals.

