# Validity of the soft approximation <sup>and</sup> the impact of higher order corrections in Higgs production in gluon fusion

#### Marco Bonvini

**DESY Hamburg** 

#### CERN, TH/LPCC Institute on SM at the LHC, October 4, 2012



Work in collaboration with:

Stefano Forte, Giovanni Ridolfi + Richard Ball, Simone Marzani

I will present a formalism to systematically answer the questions:

- 1. Is soft-gluon resummation needed?
- 2. Is soft-gluon resummation advisable? (Are soft terms dominant?)

[MB, Forte, Ridolfi 2012 (and 2010)]

I will present a formalism to systematically answer the questions:

- 1. Is soft-gluon resummation needed?
- 2. Is soft-gluon resummation advisable? (Are soft terms dominant?)
- 3. Is the pointlike (large- $m_t$ ) approximation adequate? (Higgs only)

[MB, Forte, Ridolfi 2012 (and 2010)]

I will present a formalism to systematically answer the questions:

- 1. Is soft-gluon resummation needed?
- 2. Is soft-gluon resummation advisable? (Are soft terms dominant?)
- 3. Is the pointlike (large- $m_t$ ) approximation adequate? (Higgs only)

[MB, Forte, Ridolfi 2012 (and 2010)]

I will also present:

 a prediction of N<sup>3</sup>LO Higgs cross-section (gg channel only) [Ball, MB, Forte, Marzani, Ridolfi (work in progress)]

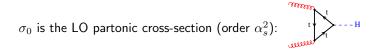
# Higgs production cross-section: notations

$$\sigma(\tau) = \tau \, \sigma_0 \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}\left(\frac{\tau}{z}\right) C(z, \alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$

I concentrate on the gluon fusion channel:  $\mathscr{L}=\mathscr{L}_{gg}$ 

Partonic cross section

$$\hat{\sigma}(z) = z \,\sigma_0 \underbrace{\left[ \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 C^{(2)}(z) + \ldots \right]}_{\text{coefficient function } C(z, \alpha_s)}$$



# Higgs production cross-section: notations

$$\sigma(\tau) = \tau \, \sigma_0 \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}\left(\frac{\tau}{z}\right) C(z, \alpha_s), \qquad \tau = \frac{m_H^2}{s}, \qquad z = \frac{m_H^2}{\hat{s}}$$

I concentrate on the gluon fusion channel:  $\mathscr{L}=\mathscr{L}_{gg}$ 

Partonic cross section

$$\hat{\sigma}(z) = z \,\sigma_0 \underbrace{\left[ \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 C^{(2)}(z) + \dots \right]}_{\text{coefficient function } C(z, \alpha_s)}$$

$$\sigma_0$$
 is the LO partonic cross-section (order  $\alpha_s^2$ ):

The questions 1., 2., 3. refer to the hadronic cross-section  $\sigma(\tau)$ , but the approximations apply to the partonic coefficient function  $C(z, \alpha_s)$ .

It is not straightforward to answer!

Marco Bonvini

$$\frac{\sigma(\tau)}{\tau \, \sigma_0} = \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}\left(\frac{\tau}{z}\right) C(z, \alpha_s)$$

The soft region  $z \to 1$  is always included in the integration region. When is the  $z \sim 1$  region actually relevant/dominant?

$$\frac{\sigma(\tau)}{\tau \, \sigma_0} = \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}\left(\frac{\tau}{z}\right) C(z, \alpha_s)$$

The soft region  $z \to 1$  is always included in the integration region. When is the  $z \sim 1$  region actually relevant/dominant?

Mellin transform:

$$\sigma(N) = \mathscr{L}(N) C(N, \alpha_s)$$

$$\frac{\sigma(\tau)}{\tau \,\sigma_0} = \int \frac{dN}{2\pi i} \, \tau^{-N} \mathscr{L}(N) \, C(N, \alpha_s)$$

$$\frac{\sigma(\tau)}{\tau \, \sigma_0} = \int_{\tau}^1 \frac{dz}{z} \, \mathscr{L}\left(\frac{\tau}{z}\right) C(z, \alpha_s)$$

The soft region  $z \to 1$  is always included in the integration region. When is the  $z \sim 1$  region actually relevant/dominant?

Mellin transform:

$$\sigma(N) = \mathscr{L}(N) C(N, \alpha_s)$$

$$\frac{\sigma(\tau)}{\tau \sigma_0} = \int \frac{dN}{2\pi i} \tau^{-N} \mathscr{L}(N) C(N, \alpha_s)$$
$$= \int \frac{dN}{2\pi i} \exp\left[N \log \frac{1}{\tau} + \log \mathscr{L}(N) + \log C(N, \alpha_s)\right]$$

$$\frac{\sigma(\tau)}{\tau \,\sigma_0} = \int_{\tau}^1 \frac{dz}{z} \,\mathscr{L}\left(\frac{\tau}{z}\right) C(z,\alpha_s)$$

The soft region  $z \to 1$  is always included in the integration region. When is the  $z \sim 1$  region actually relevant/dominant?

Mellin transform:

$$\sigma(N) = \mathscr{L}(N) C(N, \alpha_s)$$

$$\frac{\sigma(\tau)}{\tau \sigma_0} = \int \frac{dN}{2\pi i} \tau^{-N} \mathscr{L}(N) C(N, \alpha_s)$$
$$= \int \frac{dN}{2\pi i} \exp\left[N \log \frac{1}{\tau} + \log \mathscr{L}(N) + \log C(N, \alpha_s)\right]$$

The integral is dominated by the values of N in the proximity of the saddle point  $N=N_0{\rm :}$ 

$$\log \frac{1}{\tau} = -\left. \frac{d}{dN} \log \mathscr{L}(N) \right|_{N_0} - \left. \frac{d}{dN} \log C(N, \alpha_s) \right|_{N_0}$$

[MB, Forte, Ridolfi 2010]

Marco Bonvini

# Then we rephrase the questions:

- 1. Is soft-gluon resummation needed?
- 2. Is soft-gluon resummation advisable?
- 3. Is the pointlike (large- $m_t$ ) approximation adequate?

#### ₩

- 1. Is soft-gluon resummation needed at the saddle point  $N_0$ ?
- 2. Is soft-gluon resummation advisable at the saddle point  $N_0$ ?
- 3. Is the pointlike approximation adequate at the saddle point  $N_0$ ?

#### Then we rephrase the questions:

- 1. Is soft-gluon resummation needed?
- 2. Is soft-gluon resummation advisable?
- 3. Is the pointlike (large- $m_t$ ) approximation adequate?

#### ₩

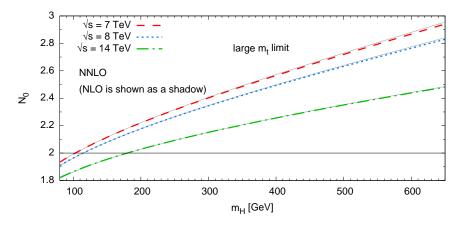
- 1. Is soft-gluon resummation needed at the saddle point  $N_0$ ?
- 2. Is soft-gluon resummation advisable at the saddle point  $N_0$ ?
- 3. Is the pointlike approximation adequate at the saddle point  $N_0$ ?

$$\sigma(N) = \mathscr{L}(N) C(N, \alpha_s)$$

Hadron level kinematics  $m_H, \sqrt{s}$  (or  $\tau$ )  $\Rightarrow$  saddle point  $N_0$ 

#### Position of the saddle point

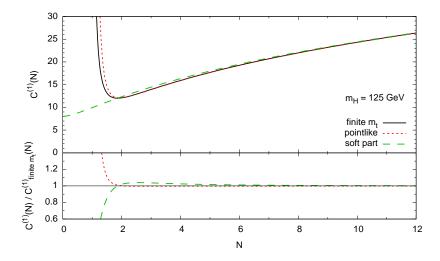
A unique real saddle  $N_0$  always exists.



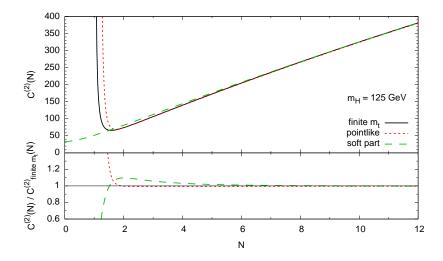
Note that these curves **do not depend on the perturbative order**. This is due to the fact that they are mainly determined by the PDFs.

#### Partonic comparison

$$C(z, \alpha_s) = \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 C^{(2)}(z) + \alpha_s^3 C^{(3)}(z) + \dots$$



$$C(z, \alpha_s) = \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 \frac{C^{(2)}(z)}{C} + \alpha_s^3 C^{(3)}(z) + \dots$$



# Partonic comparison

$$C(z, \alpha_s) = \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 C^{(2)}(z) + \alpha_s^3 C^{(3)}(z) + \dots$$

# ?

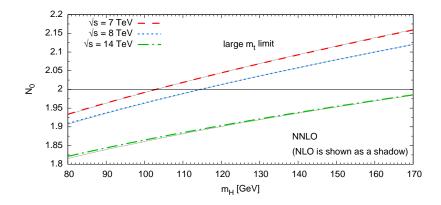
#### Partonic comparison

$$C(z, \alpha_s) = \delta(1-z) + \alpha_s C^{(1)}(z) + \alpha_s^2 C^{(2)}(z) + \alpha_s^3 C^{(3)}(z) + \dots$$

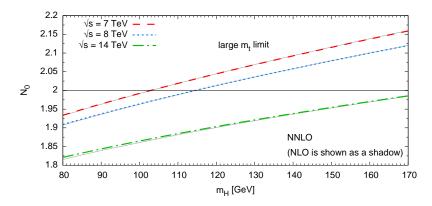
# ?

Based on all-order considerations, we can argue that soft terms dominate for  $N\gtrsim 2$ 

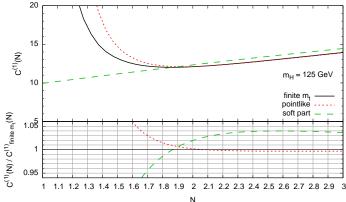
(and, a fortiori, the pointlike approx has to be good at least in the same region)



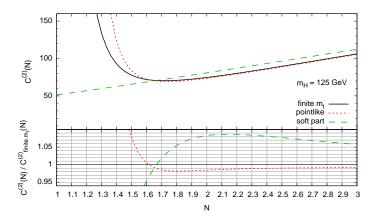
- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0 \simeq 2.03$  for LHC  $8~{\rm TeV}$
- $N_0\simeq 1.92~{\rm for}~{\rm LHC}~14~{\rm TeV}$



- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0 \simeq 2.03$  for LHC 8 TeV
- $N_0 \simeq 1.92$  for LHC 14 TeV



- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0 \simeq 2.03$  for LHC 8 TeV
- $N_0 \simeq 1.92$  for LHC  $14~{\rm TeV}$



- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0\simeq 2.03$  for LHC  $8~{\rm TeV}$
- $N_0 \simeq 1.92$  for LHC  $14~{\rm TeV}$

LHC	pointlike	soft
	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
8  TeV	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
$14 { m ~TeV}$	good ( $\sim 1-2\%$ )	not so good $(5-10\%)$

Assuming  $m_H = 125 \text{ GeV}$ 

- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0 \simeq 2.03$  for LHC 8 TeV
- $N_0 \simeq 1.92$  for LHC  $14~{\rm TeV}$

LHC	pointlike	soft
	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
8  TeV	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
$14 { m ~TeV}$	good ( $\sim 1-2\%$ )	not so good $(5-10\%)$

Is resummation needed?

Assuming  $m_H = 125 \text{ GeV}$ 

- $N_0 \simeq 2.06$  for LHC 7 TeV
- $N_0\simeq 2.03$  for LHC  $8~{\rm TeV}$
- $N_0 \simeq 1.92$  for LHC  $14~{\rm TeV}$

LHC	pointlike	soft
	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
8  TeV	very good ( $\sim 1\%$ )	quite good $(4-9\%)$
$14 \mathrm{TeV}$	good ( $\sim 1-2\%$ )	not so good $(5-10\%)$

Is resummation *needed*?

The loss-of-perturbativity condition is  $\alpha_s \log^2 N \gtrsim 1$  which gives  $N \gtrsim 10$ 

The saddle point is never so large!  $\Rightarrow$  Resummation never needed.

Soft-gluon resummation can just be used as a (useful!) tool to predict higher order terms.

Marco Bonvini

- We have a solid method to establish the dominant partonic region (saddle point)
- The saddle point is independent on the perturbative order and on other details  $\rightarrow$  property of the considered subprocess
- Based on all-order considerations, we can argue that the partonic region in which soft terms dominate is the region  $N\gtrsim 2$
- Then, when  $N_0 \gtrsim 2$ , we can predict higher orders using soft terms coming from soft-gluon resummation

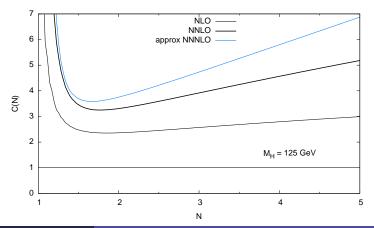
# N<sup>3</sup>LO prediction for Higgs production

Actually, we use terms coming from both large- and small-z resummations [Ball, MB, Forte, Marzani, Ridolfi (work in progress)]

# N<sup>3</sup>LO prediction for Higgs production

Actually, we use terms coming from both large- and small-z resummations [Ball, MB, Forte, Marzani, Ridolfi (work in progress)]

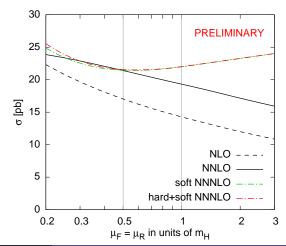
$$C(N, \alpha_s) = 1 + \alpha_s C^{(1)}(N) + \alpha_s^2 C^{(2)}(N) + \alpha_s^3 C^{(3)}(N) + \dots$$



Marco Bonvini

# N<sup>3</sup>LO prediction for Higgs production

Actually, we use terms coming from both large- and small-z resummations [Ball, MB, Forte, Marzani, Ridolfi (work in progress)]



Higgs hadronic cross-section

Validity of the soft approximation

- The saddle point method is fast, efficient and it works!
- It is a tool to easily check the validity of the soft (and other) approximations

• Approximate prediction for Higgs cross-section at N<sup>3</sup>LO

# Backup slides

In z-space: 
$$\mathcal{D}_k(z) = \left[\frac{\log^k(1-z)}{1-z}\right]_+$$

In z-space: 
$$\mathcal{D}_k(z) = \left[\frac{\log^k(1-z)}{1-z}\right]_+$$

 $2\mathcal{D}_1(N) =$ 

$$= \psi_0^2(N) - \psi_1(N) + 2\gamma\psi_0(N) + \zeta_2 + \gamma^2$$

$$\left[\frac{\log(1-z)}{1-z}\right]_+$$

In z-space: 
$$\mathcal{D}_k(z) = \left[\frac{\log^k(1-z)}{1-z}\right]_+$$

 $2\mathcal{D}_1(N) =$ 

$$= \psi_0^2(N) - \psi_1(N) + 2\gamma\psi_0(N) + \zeta_2 + \gamma^2$$

$$\simeq \psi_0^2(N) + 2\gamma\psi_0(N) + \zeta_2 + \gamma^2$$

$$\begin{bmatrix} \frac{\log(1-z)}{1-z} \end{bmatrix}_+$$
$$\begin{bmatrix} \frac{\log\frac{1-z}{\sqrt{z}}}{1-z} \end{bmatrix}_+' = \begin{bmatrix} \frac{\log(1-z)}{1-z} \end{bmatrix}_+ - \frac{\log\sqrt{z}}{1-z}$$

$$\ln z \text{-space: } \mathcal{D}_{k}(z) = \left\lfloor \frac{\log^{*}(1-z)}{1-z} \right\rfloor_{+}$$

$$= \psi_{0}^{2}(N) - \psi_{1}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2} \qquad \left[ \frac{\log(1-z)}{1-z} \right]_{+}$$

$$\simeq \psi_{0}^{2}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2} \qquad \left[ \frac{\log\frac{1-z}{\sqrt{z}}}{1-z} \right]_{+}' = \left[ \frac{\log(1-z)}{1-z} \right]_{+} - \frac{\log\sqrt{z}}{1-z}$$

$$\simeq \log^{2}\frac{1}{N-\frac{1}{2}} + 2\gamma\log\frac{1}{N-\frac{1}{2}} + \zeta_{2} + \gamma^{2} \qquad \frac{1}{\sqrt{z}} \left[ \frac{\log\log\frac{1}{z}}{\log\frac{1}{z}} \right]_{+} + (\zeta_{2} + \gamma^{2})\delta(1-z)$$

 $2\mathcal{D}_1(N) =$ 

 $=\psi_0^2(N)-\psi_1$ 

 $\simeq \log^2 \frac{1}{N - \frac{1}{2}}$ 

$$2\mathcal{D}_{1}(N) = \\ = \psi_{0}^{2}(N) - \psi_{1}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2} \qquad \begin{bmatrix} \frac{1}{N} \\ \frac{1}{N}$$

In z-space: 
$$\mathcal{D}_k(z) = \left[ rac{\log^k(1-z)}{1-z} 
ight]_+$$

$$\begin{bmatrix} \underline{\log(1-z)}\\ 1-z \end{bmatrix}_+ \\ \begin{bmatrix} \underline{\log\frac{1-z}}{\sqrt{z}}\\ 1-z \end{bmatrix}_+ = \begin{bmatrix} \underline{\log(1-z)}\\ 1-z \end{bmatrix}_+ - \frac{\log\sqrt{z}}{1-z} \\ \frac{1}{\sqrt{z}} \begin{bmatrix} \underline{\log\log\frac{1}{z}}\\ \frac{1}{2} \end{bmatrix}_+ + (\zeta_2 + \gamma^2)\delta(1-z) \\ \begin{bmatrix} \underline{\log\log\frac{1}{z}}\\ \frac{1}{2} \end{bmatrix}_+ + (\zeta_2 + \gamma^2)\delta(1-z) \end{bmatrix}$$

$$2\mathcal{D}_{1}(N) =$$

$$= \psi_{0}^{2}(N) - \psi_{1}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2}$$

$$\simeq \psi_{0}^{2}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2}$$

$$\simeq \log^{2} \frac{1}{N - \frac{1}{2}} + 2\gamma \log \frac{1}{N - \frac{1}{2}} + \zeta_{2} + \gamma^{2}$$

$$\simeq \log^{2} \frac{1}{N} + 2\gamma \log \frac{1}{N} + \zeta_{2} + \gamma^{2}$$

$$\simeq \psi_{0}^{2}(N + 1) + 2\gamma\psi_{0}(N + 1) + \zeta_{2} + \gamma^{2}$$

In z-space: 
$$\mathcal{D}_k(z) = \left[ rac{\log^k(1-z)}{1-z} 
ight]_+$$

$$\begin{bmatrix} \frac{\log(1-z)}{1-z} \\ 1-z \end{bmatrix}_{+}^{\prime} = \begin{bmatrix} \frac{\log(1-z)}{1-z} \\ 1-z \end{bmatrix}_{+}^{\prime} - \frac{\log\sqrt{z}}{1-z} \\ \frac{1}{\sqrt{z}} \begin{bmatrix} \frac{\log\log\frac{1}{z}}{\log\frac{1}{z}} \\ \frac{\log\log\frac{1}{z}}{\log\frac{1}{z}} \end{bmatrix}_{+}^{\prime} + (\zeta_{2} + \gamma^{2})\delta(1-z) \\ \begin{bmatrix} \frac{\log\log\frac{1}{z}}{\log\frac{1}{z}} \\ \frac{\log\frac{1-z}{\sqrt{z}}}{\log\frac{1}{z}} \end{bmatrix}_{+}^{\prime} + (\zeta_{2} + \gamma^{2})\delta(1-z) \\ z \begin{bmatrix} \frac{\log\frac{1-z}{\sqrt{z}}}{1-z} \end{bmatrix}_{+}^{\prime}$$

$$\ln z \text{-space: } \mathcal{D}_{k}(z) = \left[\frac{\log\left(1-z\right)}{1-z}\right]_{+}$$

$$= \psi_{0}^{2}(N) - \psi_{1}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2} \qquad \left[\frac{\log(1-z)}{1-z}\right]_{+}$$

$$\simeq \psi_{0}^{2}(N) + 2\gamma\psi_{0}(N) + \zeta_{2} + \gamma^{2} \qquad \left[\frac{\log\frac{1-z}{\sqrt{z}}}{1-z}\right]_{+}^{\prime} = \left[\frac{\log(1-z)}{1-z}\right]_{+} - \frac{\log\sqrt{z}}{1-z}$$

$$\simeq \log^{2}\frac{1}{N-\frac{1}{2}} + 2\gamma\log\frac{1}{N-\frac{1}{2}} + \zeta_{2} + \gamma^{2} \qquad \frac{1}{\sqrt{z}}\left[\frac{\log\log\frac{1}{z}}{\log\frac{1}{z}}\right]_{+} + (\zeta_{2} + \gamma^{2})\delta(1-z)$$

$$\simeq \log^{2}\frac{1}{N} + 2\gamma\log\frac{1}{N} + \zeta_{2} + \gamma^{2} \qquad \left[\frac{\log\log\frac{1}{z}}{\log\frac{1}{z}}\right]_{+} + (\zeta_{2} + \gamma^{2})\delta(1-z)$$

$$\simeq \psi_{0}^{2}(N+1) + 2\gamma\psi_{0}(N+1) + \zeta_{2} + \gamma^{2} \qquad z\left[\frac{\log\frac{1-z}{\sqrt{z}}}{1-z}\right]_{+}^{\prime}$$

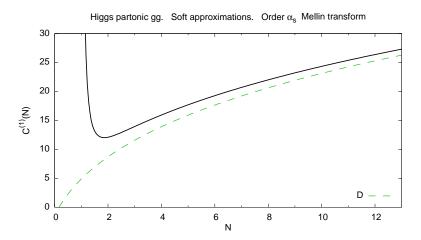
The last option allows to include all the terms  $\alpha_s^n \log^{2n-1}(1-z)$  in the soft terms (equivalent to the collinear improvement of [Krämer, Laenen, Spira 1997], [Catani, de Florian, Grazzini 2001], [Catani, de Florian, Grazzini, Nason 2003])

Marco Bonvini

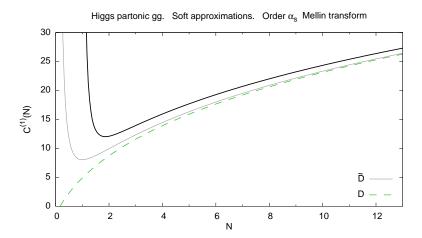
=

- ()  $\left[\log^{k}(1-\alpha)\right]$ 

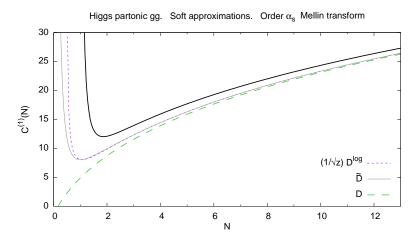
$$\mathcal{D}_k = \left[\frac{\log^k(1-z)}{1-z}\right]_+$$



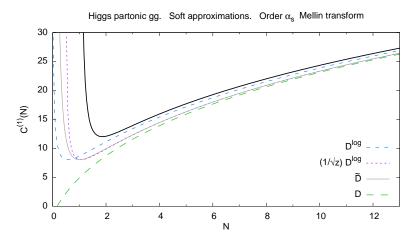
$$\mathcal{D}_{k} = \left[\frac{\log^{k}(1-z)}{1-z}\right]_{+} \qquad \tilde{\mathcal{D}}_{k} = \left[\frac{\log^{k}\frac{1-z}{\sqrt{z}}}{1-z}\right]_{+}^{\prime}$$



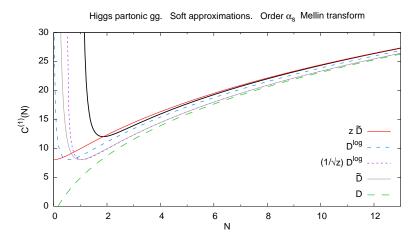
$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$



$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$

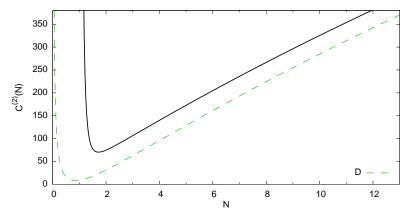


$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$



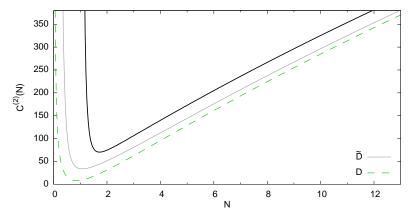
$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]'_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$

Higgs partonic gg. Soft approximations. Order  $\alpha_s^2$  Mellin transform



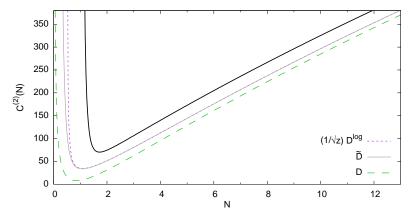
$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]'_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$

Higgs partonic gg. Soft approximations. Order  $\alpha_s^2$  Mellin transform



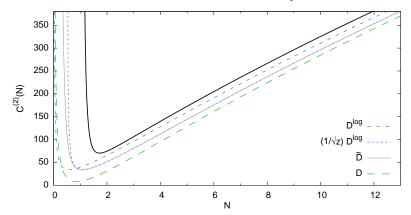
$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$

Higgs partonic gg. Soft approximations. Order  $\alpha_s^2$  Mellin transform



$$\mathcal{D}_{k} = \left[\frac{\log^{k}(1-z)}{1-z}\right]_{+} \qquad \tilde{\mathcal{D}}_{k} = \left[\frac{\log^{k}\frac{1-z}{\sqrt{z}}}{1-z}\right]_{+}' \qquad \mathcal{D}_{k}^{\log} = \left[\frac{\log^{k}\log\frac{1}{z}}{\log\frac{1}{z}}\right]_{+} + c_{k}\delta(1-z)$$

Higgs partonic gg. Soft approximations. Order  ${\alpha_s}^2$  Mellin transform



$$\mathcal{D}_k = \left[\frac{\log^k (1-z)}{1-z}\right]_+ \qquad \tilde{\mathcal{D}}_k = \left[\frac{\log^k \frac{1-z}{\sqrt{z}}}{1-z}\right]_+ \qquad \mathcal{D}_k^{\log} = \left[\frac{\log^k \log \frac{1}{z}}{\log \frac{1}{z}}\right]_+ + c_k \delta(1-z)$$

Higgs partonic gg. Soft approximations. Order  ${\alpha_s}^2$  Mellin transform

