

Light scalar at LHC: the Higgs or the dilaton?

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(with W. Goldberger and B. Grinstein)

hep-ph/0708:1463
(today)

Outline

- Why is the Higgs so similar to the dilaton?
- Conformal theories at the TeV scale
- Dilaton couplings:
 - Higgs-like: fermions and W/Z
 - Higgs-unlike: massless gauge bosons, self couplings
- The dilaton at colliders: LEP, LHC, ILC.
- Things to do ...

Why is the Higgs so similar to the dilaton?

Because the dilaton couples to $T_{\mu}^{\mu} = \sum_i m_i \bar{\psi} \psi + \dots$

Because a light Higgs is also the dilaton

- SM interactions are approximately conformal down to the QCD scale
- Higgs mass term - explicit breaking
- Higgs VEV - spontaneous breaking

However, in general conformal invariance can be broken at a higher scale than the EW symmetry

$$\Lambda_{CFT} \sim 4\pi f$$

The breaking of conformal invariance triggers EWSB

$$\Lambda_{EW} \sim 4\pi v \leq \Lambda_{CFT}$$

The scales v and f are not the same,
except for the Higgs

Conformal theories at the TeV scale

Classic example: walking technicolor
(some doubted that the dilaton would be light)

Things changed with AdS/CFT and RS model,
where there are plenty of examples of CFT's
that are spontaneously broken

There is no doubt a small parameter controlling
the dilaton mass exists in such theories

Dilaton couplings

Given the Lagrangian $\mathcal{L} = \sum_i g_i(\mu) \mathcal{O}_i(x)$,

the divergence of the scale current is:

$$\partial_\mu S^\mu = \sum_i g_i(\mu) (d_i - 4) \mathcal{O}_i(x) + \sum_i \beta_i(g) \frac{\partial}{\partial g_i} \mathcal{L}$$

Including the dilaton field, $\chi(x)$, makes the Lagrangian formally scale invariant

$$g_i(\mu) \rightarrow g_i \left(\mu \frac{\chi}{f} \right) \left(\frac{\chi}{f} \right)^{4-d_i}$$

The electroweak sector:

$$\mathcal{L}_{EW} = \mathcal{L}_{\chi EW} + \mathcal{L}_{\psi} + \mathcal{L}_Y$$

EW chiral Lagrangian kinetic terms for fermions Yukawa couplings

After replacing $\chi(x) \rightarrow f + \bar{\chi}(x)$

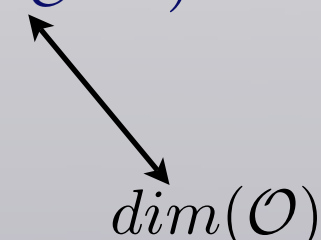
$$\mathcal{L} = \left(\frac{2\bar{\chi}}{f} + \frac{\bar{\chi}^2}{f^2} \right) \left[m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right] + \frac{\bar{\chi}}{f} \sum_{\psi} m_{\psi} \bar{\psi} \psi$$

(The usual Higgs couplings rescaled by $\frac{v}{f}$. Note only partial restoration of unitarity if $f > v$.)

Dilaton cubic self coupling

Suppose CS is explicitly broken: $\mathcal{L}_{CFT} + \lambda_{\mathcal{O}} \mathcal{O}(x)$

Usual spurion
analysis gives

$$V(\chi) = \chi^4 \sum_{n=0}^{\infty} c_n(\Delta_{\mathcal{O}}) \left(\frac{\chi}{f} \right)^{n(\Delta_{\mathcal{O}} - 4)}$$


There are two limits in which there is a small parameter
(a) $\lambda_{\mathcal{O}}$ small in units of f , (b) $|\Delta_{\mathcal{O}} - 4| \ll 1$

$$V(\bar{\chi}) = \frac{1}{2} m^2 \bar{\chi}^2 + \frac{\lambda}{3!} \frac{m^2}{f} \bar{\chi}^3 + \dots$$

$$\lambda = \begin{cases} (\Delta_{\mathcal{O}} + 1) + \mathcal{O}(\lambda_{\mathcal{O}}) & \text{case (a)} \\ 5 + \mathcal{O}(|\Delta_{\mathcal{O}} - 4|) & \text{case (b)} \end{cases}$$

$$\lambda = \begin{cases} (\Delta_{\mathcal{O}} + 1) + \mathcal{O}(\lambda_{\mathcal{O}}) & \text{when } \lambda_{\mathcal{O}} \ll 1 \\ 5 + \mathcal{O}(|\Delta_{\mathcal{O}} - 4|) & \text{when } |\Delta_{\mathcal{O}} - 4| \ll 1 \end{cases}$$

The Higgs case, $\Delta_{\mathcal{O}} = 2$, checks out $\lambda = 3$

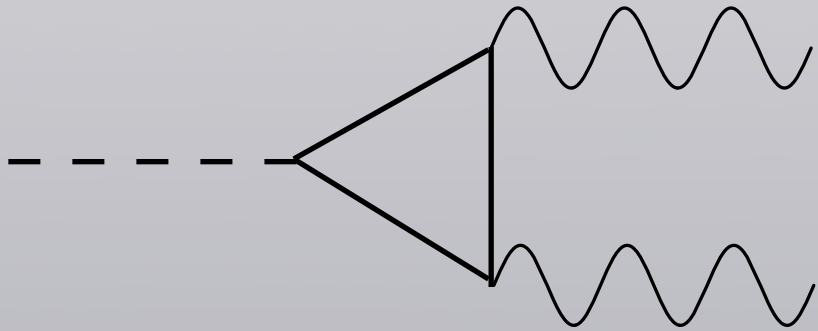
Irrelevant perturbations should not break conformal symmetry which implies an upper bound on the cubic

$$\lambda \leq 5$$

saturated for nearly marginal operators

Couplings to massless gauge bosons

At zero momentum the Higgs/dilaton couplings are related to the conformal anomaly



$$\mathcal{L}_{hGG} = \frac{\alpha_s}{8\pi} \sum_{\text{heavy}} b_0^i \frac{h}{v} (G_{\mu\nu}^a)^2$$

If SM is embedded in a conformal theory

$$\sum_{\text{light}} b_0 + \sum_{\text{heavy}} b_0 = 0$$

$$\mathcal{L}_{\chi gg} = -\frac{\alpha_s}{8\pi} b_0^{\text{light}} \frac{\bar{\chi}}{f} (G_{\mu\nu}^a)^2$$

- Large enhancement of the dilaton-gluon coupling possible, an order of magnitude compared to SM Higgs
- The coupling to the photons may be suppressed
- Not a clean dilaton signature since Higgs couplings can be altered by heavy particles as well
- An exact result for the couplings obtained using conformal compensator

$$\mathcal{L}_{\chi gg} = -\frac{\beta(g)}{2g} \frac{\bar{\chi}}{f} (G_{\mu\nu}^a)^2$$

The dilaton at colliders

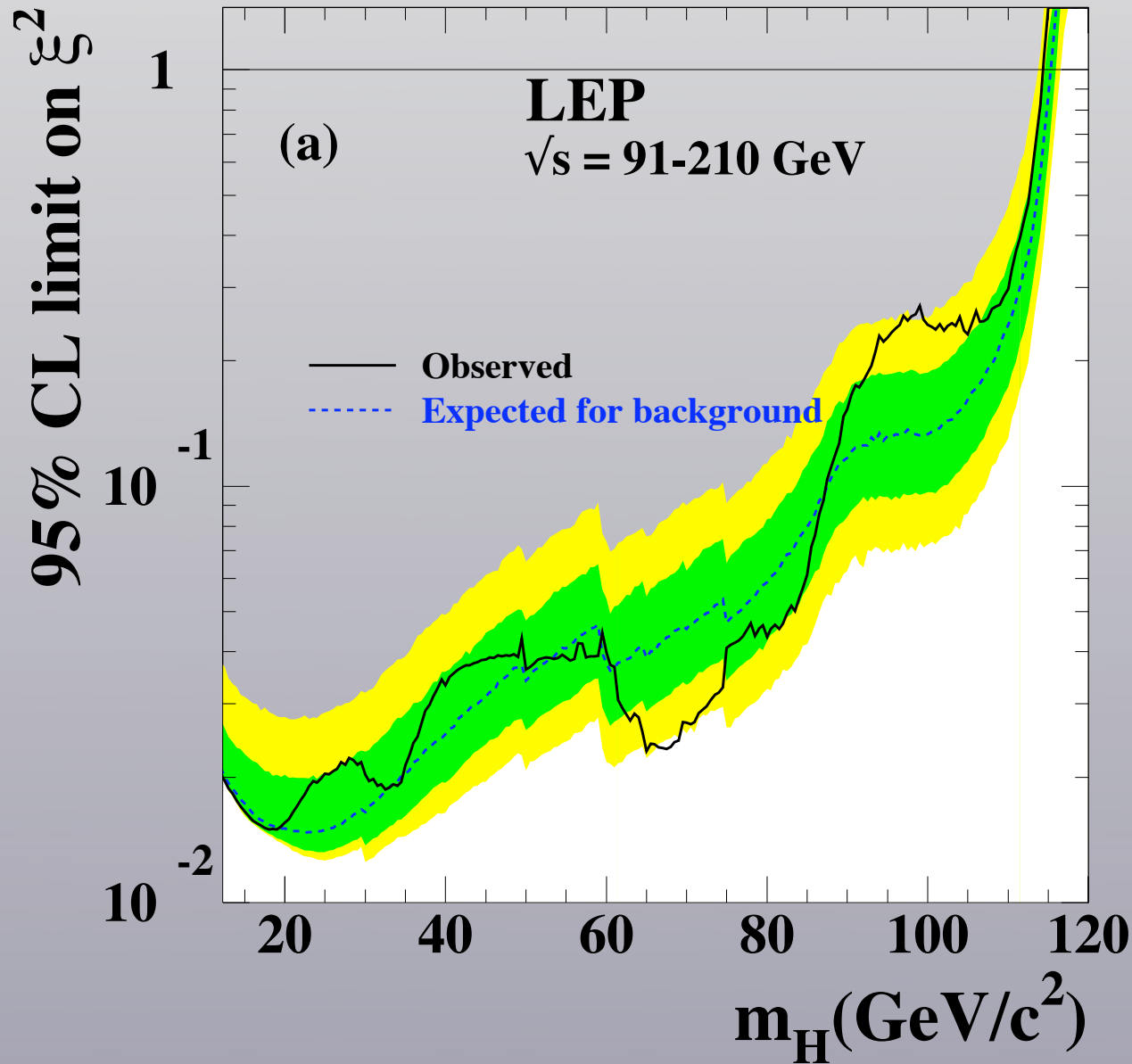
Branching ratios to fermions and WW, ZZ same as Higgs

The crucial parameters are f and m

complete Lagrangian also has three couplings: λ, c_G, c_{EM}

- LEP: bounds if $v^2/f^2 > 0.1 - 0.01$
- LHC: discovery that could be easier or harder than the Higgs case depending on the ratio v/f and the strength of the $\chi gg, \chi\gamma\gamma$ couplings, crude measurement of v/f
- ILC: precise measurements of f via couplings of gauge bosons, branching ratios, total width, a chance to measure the cubic

LEP bounds

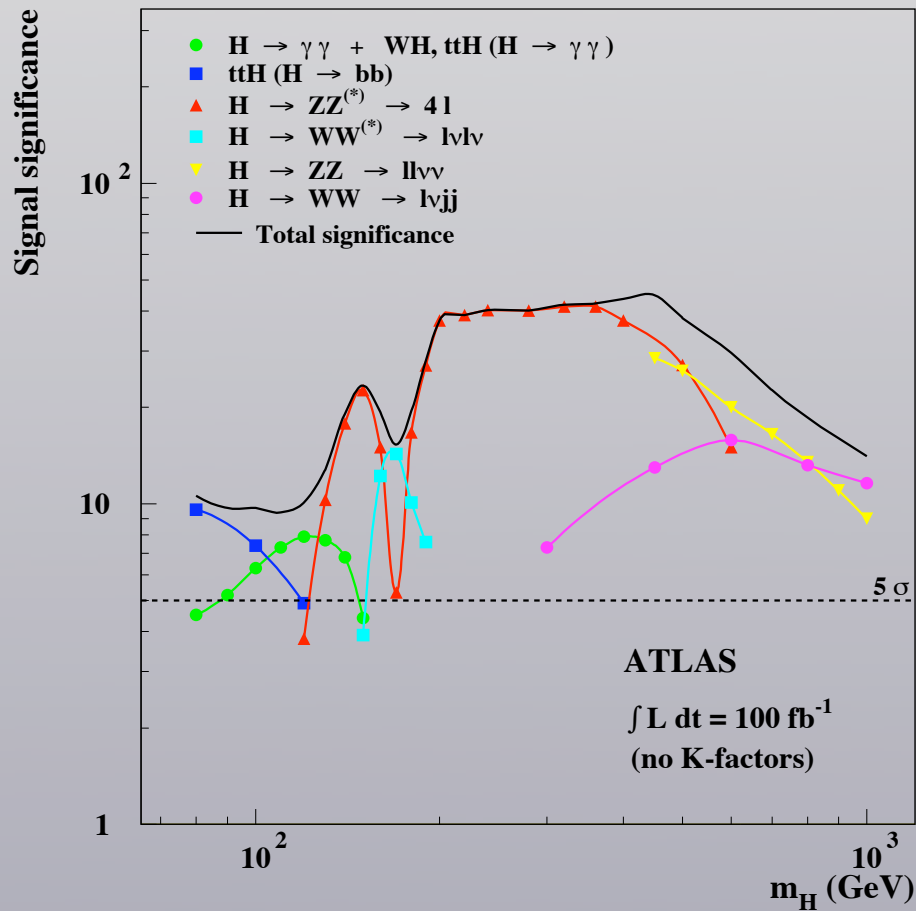


$$\xi = \frac{v}{f}$$

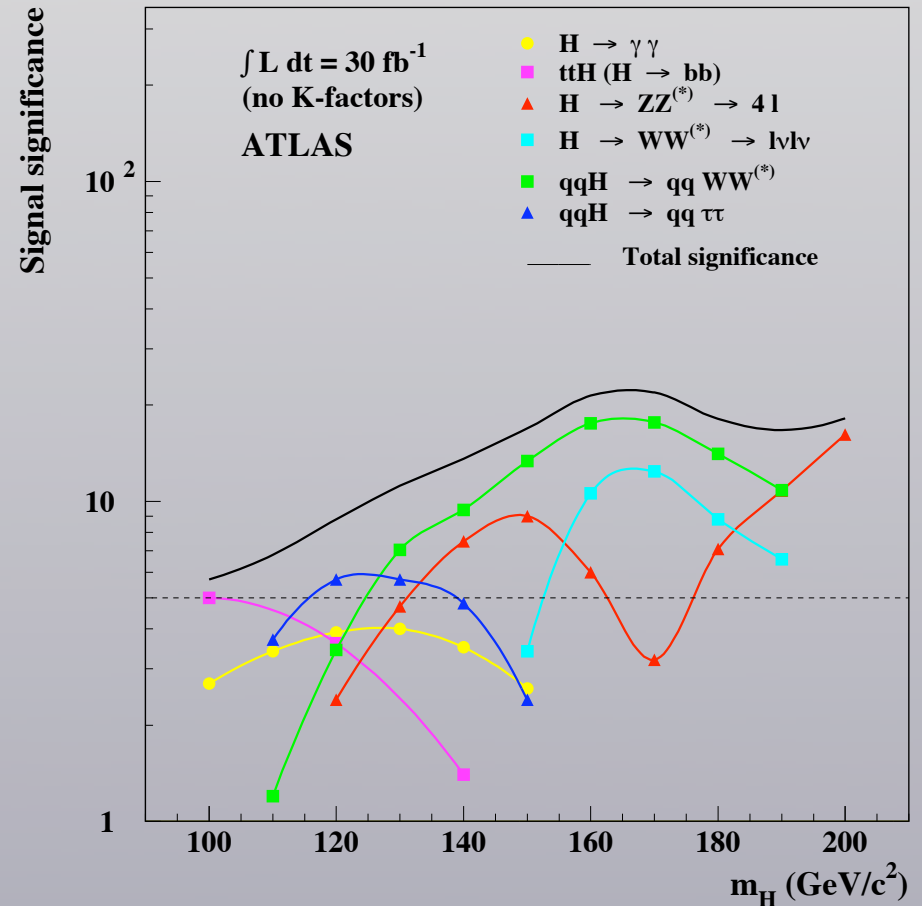
[LEP Higgs
working group]

LHC Higgs reach

(multiply by $c_G^2 \frac{v^2}{f^2}$)



(no VBF)

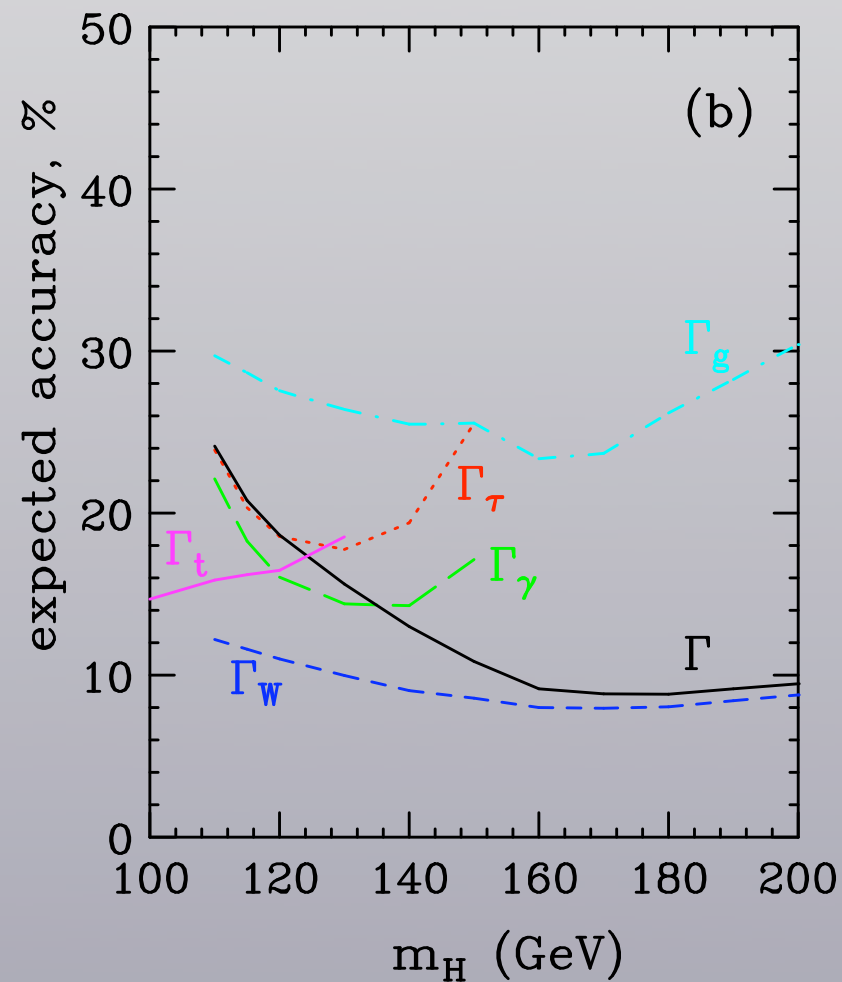
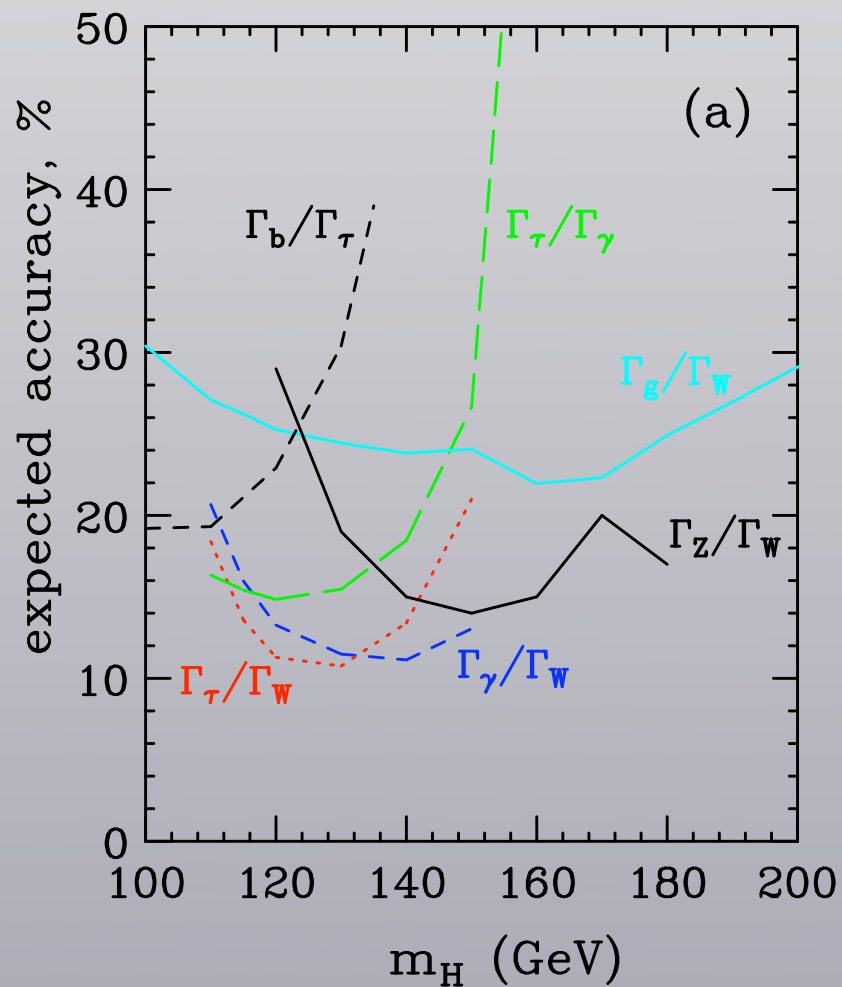


(qq-> qqh included)

LHC Higgs properties

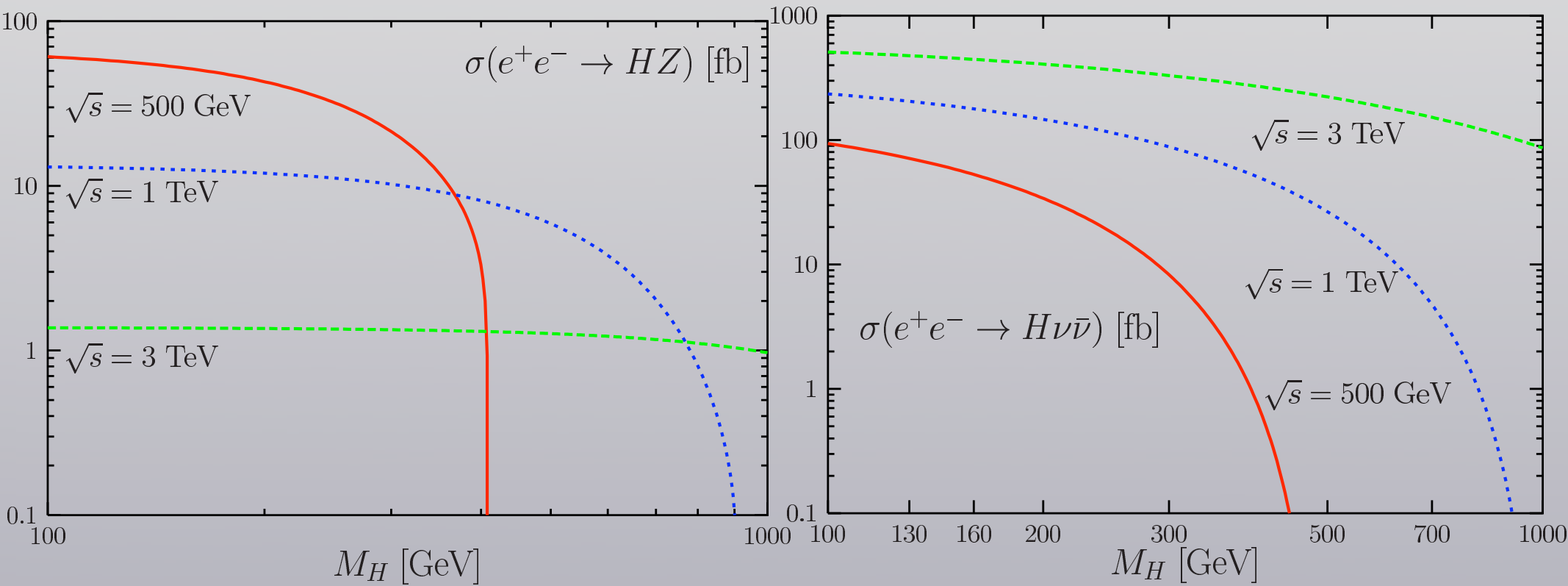
width ratios

(partial) widths



[M. Duhrssen et al.,
PRD 70, 113009]

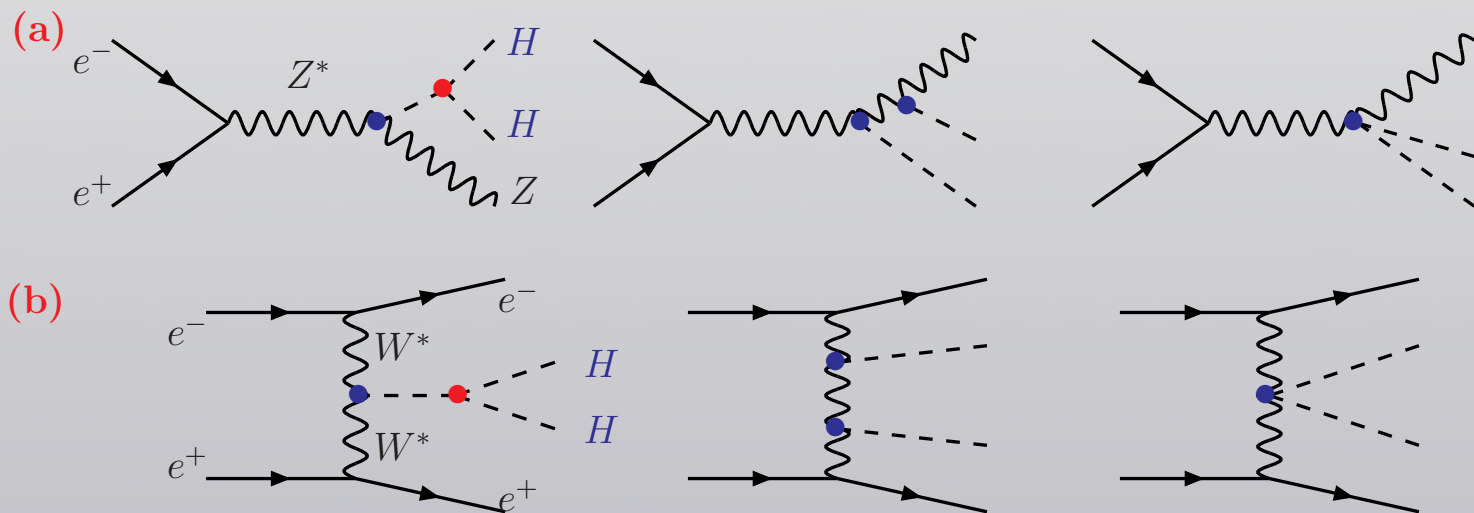
ILC Higgs production



(For the dilaton
rescale by v^2 / f^2)

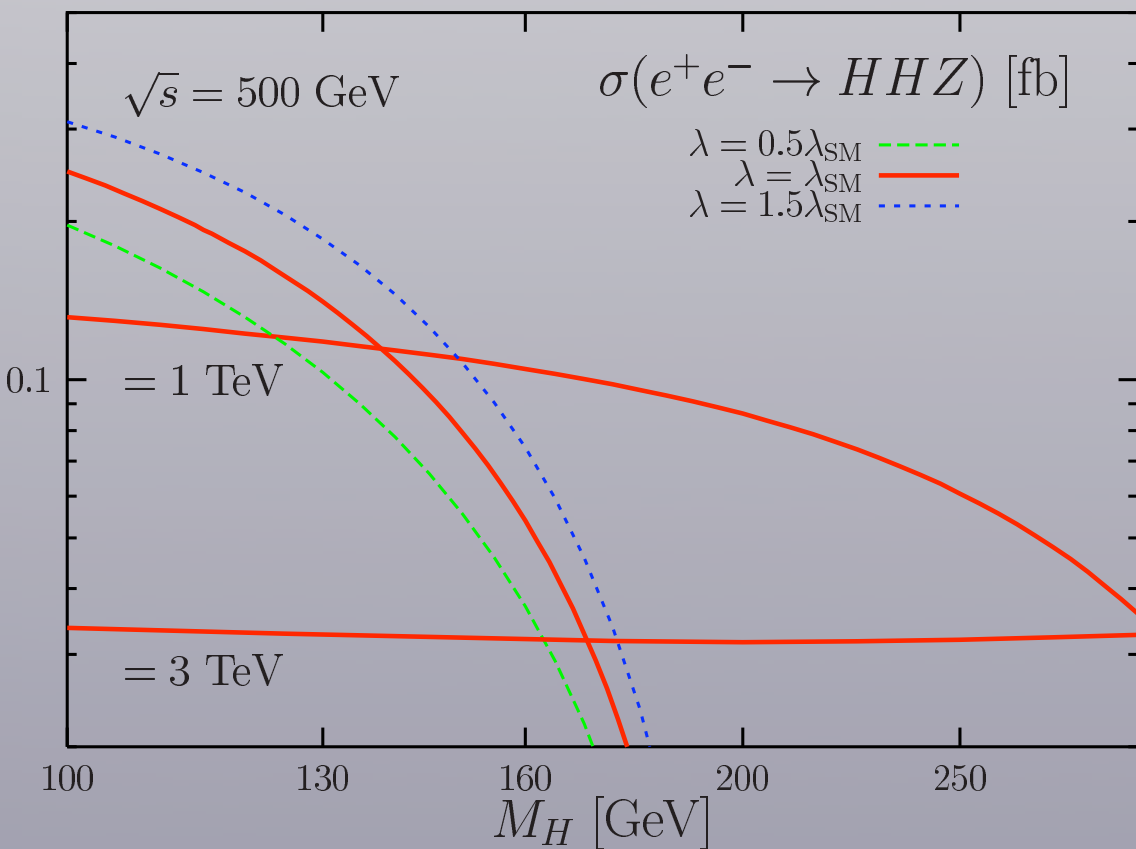
[A. Djouadi
hep-ph/0503172]

ILC 2 Higgs production



[A. Djouadi
hep-ph/0503172]

relative accuracy
(500@500)



Quantity	$M_H = 120$ GeV	$M_H = 140$ GeV
ΔM_H	± 0.00033	± 0.0005
Γ_H	± 0.061	± 0.045
ΔCP	± 0.038	—
λ_{HHH}	± 0.22	± 0.30
g_{HWW}	± 0.012	± 0.020
g_{HZZ}	± 0.012	± 0.013
g_{Htt}	± 0.030	± 0.061
g_{Hbb}	± 0.022	± 0.022
g_{Hcc}	± 0.037	± 0.102
$g_{H\tau\tau}$	± 0.033	± 0.048

Things to do ...

- Accurate estimates and search strategies at the LHC, what is the best way to determine the decay constant?
- A bound on the cubic coupling, what happens if there are several sources of symmetry breaking?
- What can we learn at the ILC?
- Dilaton mass in nearly conformal gauge theories
- Partial unitarization vs masses of heavy states
- Model-dependent questions