

Cosmological Matter-Antimatter Asymmetry and Neutrino Oscillations

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- A Conjecture + An Ansatz
- Seesaw + Leptogenesis
- ν -Mixing + Baryogenesis

Space Part 06, Beijing, 19-21 April 2006



Motivation: the new minimal SM

Davoudiasl, Kitano, Li, Murayama, [hep-ph/0405097](#)

NMSM = MSM + New Physics

(Minimal number of new degrees of freedom)

Experimental/Observational Evidence for NP:

- Dark Matter
- Dark Energy
- Cosmic Inflation
- Cosmic Baryon Asymmetry
- Atmospheric & Solar Neutrino Oscillations

Can one stone kill two big birds?

The minimal seesaw model:

Frampton, Glashow, Yanagida, [hep-ph/0208157](#);

2 Right-handed neutrinos added:

$$\nu_R = \begin{pmatrix} N_1 \\ N_2 \end{pmatrix}$$

$$-\mathcal{L}_{Y(\text{SM})} = \bar{l}_L Y_l e_R H + \bar{l}_L Y_\nu \nu_R H^c + \frac{1}{2} \bar{\nu}_R^c M_R \nu_R + \text{h.c.}$$

- Principle of minimal particle content
- $SU(2) \times U(1)$ gauge symmetry preserved

Seesaw + Leptogenesis

Neutrino Oscillations
Baryon Asymmetry

Is there special ν -mass hierarchy?

3 Right-handed neutrinos \longrightarrow more freedom

Conjecture: Universal Geometric Neutrino Mass Hierarchy

Z.Z.X., [hep-ph/0406047](#) (for both light and heavy ν);
Kaus, Meshkov, [hep-ph/0410024](#) (for light ν);
Tsujimoto, [hep-ph/0501023](#) (for heavy ν).

Light

$$\frac{m_1}{m_2} = \frac{m_2}{m_3} = \frac{M_1}{M_2} = \frac{M_2}{M_3}$$

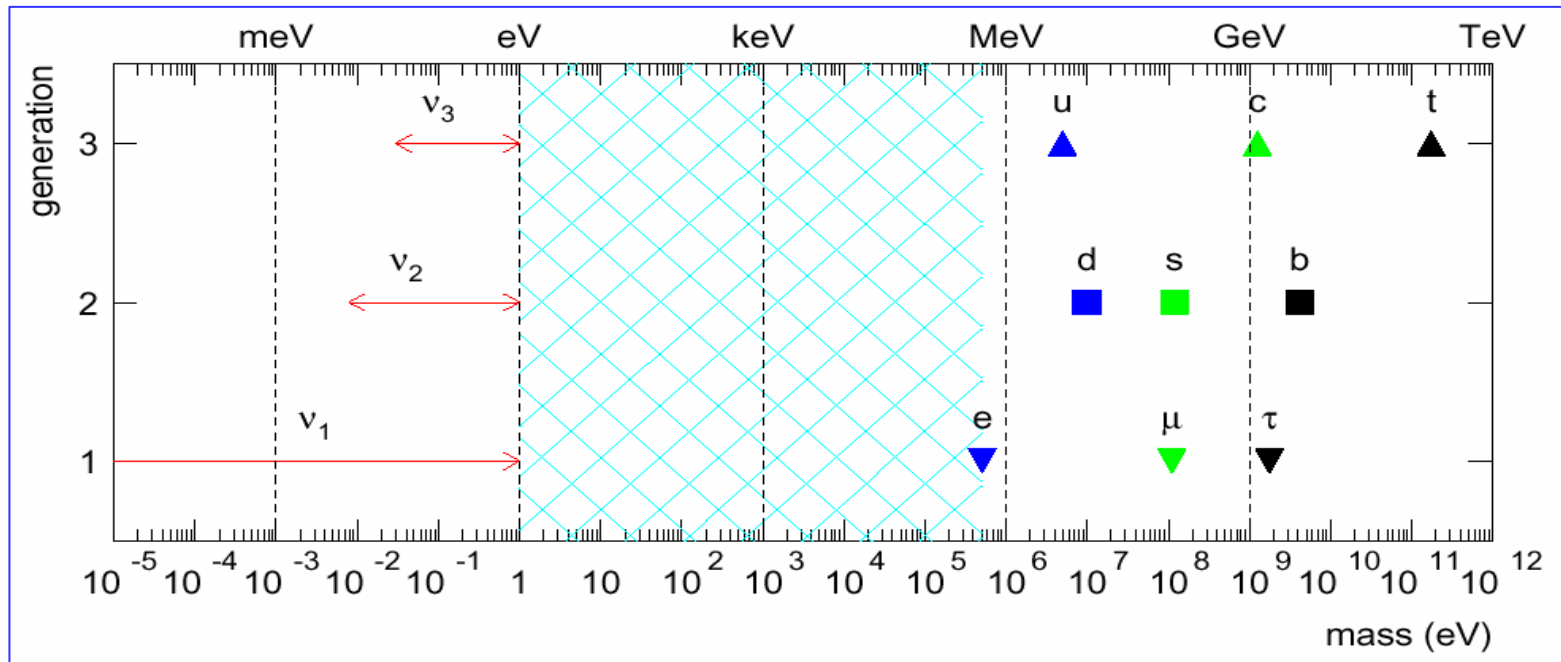
Heavy

Quarks

$$\frac{m_u}{m_c} \sim \frac{m_c}{m_t} \quad ; \quad \frac{m_d}{m_s} \sim \frac{m_s}{m_b}$$



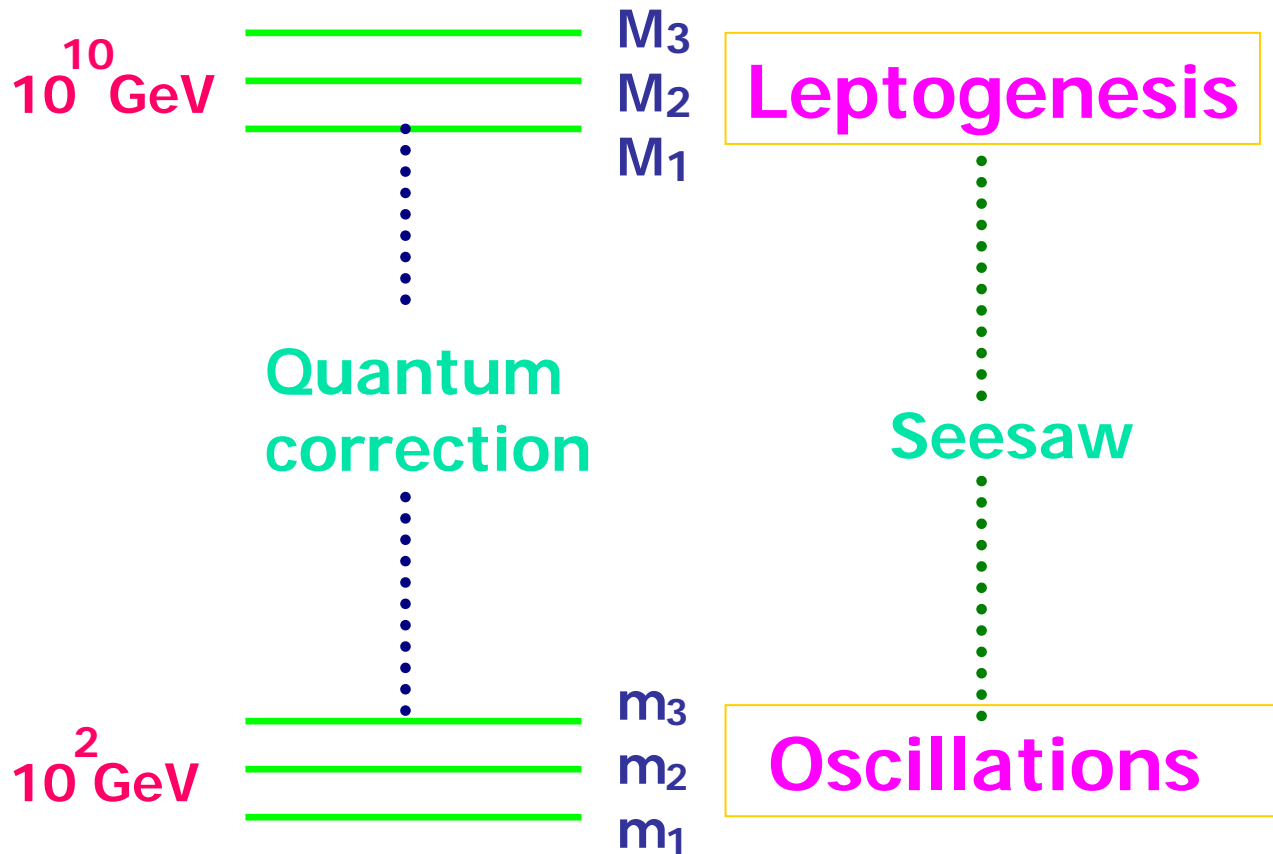
A plot of fermion mass spectrum



(a normal neutrino mass hierarchy is assumed)

Gaps: eV — MeV; and TeV — RH ν -mass scale

A phenomenological picture:



Geometric ν -mass hierarchy:

$$m_1 : m_2 : m_3 = r^2 : r : 1$$

at electroweak scale

★ **Atm** $\Delta m_{31}^2 = m_3^2 - m_1^2 \sim 2.3 \times 10^{-3} \text{eV}^2$
★ **Sun** $\Delta m_{21}^2 = m_2^2 - m_1^2 \sim 6.9 \times 10^{-5} \text{eV}^2$

$$m_1 = \frac{r^2}{\sqrt{1-r^4}} \sqrt{\Delta m_{31}^2}$$
$$m_2 = \frac{r}{\sqrt{1-r^4}} \sqrt{\Delta m_{31}^2}$$
$$m_3 = \frac{1}{\sqrt{1-r^4}} \sqrt{\Delta m_{31}^2}$$

$$r = \sqrt{\frac{\Delta m_{21}^2}{\Delta m_{31}^2 - \Delta m_{21}^2}}$$

$$0.122 \lesssim r \lesssim 0.270$$

3Q

Renormalization-group equations

$$m_i(M_1) \approx m_i(M_Z) \mathcal{I}_\alpha$$

Minimal SM case

$$\mathcal{I}_\alpha = \exp \left[\int_{\ln M_Z}^{\ln M_1} \frac{1}{16\pi^2} (\lambda_H - 3g_2^2 + 6f_t^2) d\tau \right]$$

→ $m_1(M_1)/m_2(M_1) \approx m_2(M_1)/m_3(M_1) \approx r$

Further Conjecture:

$$M_1/M_2 = M_2/M_3 = r$$

$$D_1/D_2 = D_2/D_3 = r$$

D_i : Eigenvalues of Y_ν

M_i : Eigenvalues of M_R

at the seesaw scale

$$m_i \propto \frac{D_i^2}{M_i} \langle \phi \rangle^2$$

Seesaw-invariant Fritzsch texture

Seesaw relation:

$$M_\nu \approx Y_\nu M_R^{-1} Y_\nu^T \langle \phi \rangle^2$$

Conjecture: Y_ν and M_R Fritzsch Texture

$$Y_\nu = D_3 F_D$$

$$M_R = M_3 F_R$$

$$F_\lambda = \begin{pmatrix} 0 & ce^{i\varphi_\lambda} & 0 \\ ce^{i\varphi_\lambda} & 0 & be^{i\varphi_\lambda} \\ 0 & be^{i\varphi_\lambda} & a \end{pmatrix}$$

$$a = 1 - r + r^2$$

$$b = (1 - r) \sqrt{r(1 + r^2)/a}$$

$$c = r \sqrt{r/a}$$

Phase condition

$$\phi_D - \phi_R = \varphi_D - \varphi_R$$

$$M_\nu \approx Y_\nu M_R^{-1} Y_\nu^T \langle \phi \rangle^2 = m_3 \mathcal{I}_\alpha F_\nu$$

Seesaw invariance!

$$D_3 = \sqrt{m_3 M_3 \mathcal{I}_\alpha} / \langle \phi \rangle$$

$$\mathcal{I}_\alpha \sim 1.4 \text{ for } M_1 \sim 10^{10} \text{ GeV}$$

MNS lepton flavor mixing matrix

**Charged leptons:
at the seesaw scale**

$$Y_l = \begin{pmatrix} \mathbf{0} & \tilde{c}e^{i\varphi_l} & \mathbf{0} \\ \tilde{c}e^{i\varphi_l} & \mathbf{0} & \tilde{b}e^{i\phi_l} \\ \mathbf{0} & \tilde{b}e^{i\phi_l} & \tilde{a} \end{pmatrix}$$

$$\tilde{a} \approx m_\tau / \langle \phi \rangle, \tilde{b} \approx \sqrt{m_\mu m_\tau} / \langle \phi \rangle \text{ and } \tilde{c} \approx \sqrt{m_e m_\mu} / \langle \phi \rangle$$

The **MNS** matrix V arises from mismatch between the diagonalizations of Y_l and Y_ν . Due to the normal hierarchy of m_i , **RGE** effects on V is negligibly small from seesaw to electroweak scales.

At low scales, the **Fritzsch-like** texture of lepton mass matrices is found to be compatible with the present neutrino data at the 3σ level.

(**Z.Z.X. 02**; **Z.Z.X., S. Zhou 04**)

$$0.215 \lesssim r \lesssim 0.270$$

Isomeric textures and FTY seesaw

If M_I and M_ν take a parallel and Fritzsche-like form, there are 6 different combinations. Their consequences on lepton flavor mixing are exactly the same --- an isomeric feature! (This feature is also true when the leptogenesis is concerned.)

A seesaw ansatz with Fritzsche form of M_I and M_D , but $M R = \mathbf{1}$ (3 degenerate heavy

neutrinos) by Fukugita, Tanimoto, Yanagida (93, 03):
 At low scales: better fit to current neutrino oscillations;
 At seesaw scale: no leptogenesis due to N -degeneracy.

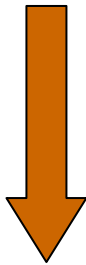
(A)	$M = \begin{pmatrix} 0 & Ce^{i\varphi} & 0 \\ Ce^{i\varphi} & 0 & Be^{i\phi} \\ 0 & Be^{i\phi} & A \end{pmatrix}$
(B)	$M = \begin{pmatrix} 0 & 0 & Ce^{i\varphi} \\ 0 & A & Be^{i\phi} \\ Ce^{i\varphi} & Be^{i\phi} & 0 \end{pmatrix}$
(C)	$M = \begin{pmatrix} 0 & Ce^{i\varphi} & Be^{i\phi} \\ Ce^{i\varphi} & 0 & 0 \\ Be^{i\phi} & 0 & A \end{pmatrix}$
(D)	$M = \begin{pmatrix} 0 & Be^{i\phi} & Ce^{i\varphi} \\ Be^{i\phi} & A & 0 \\ Ce^{i\varphi} & 0 & 0 \end{pmatrix}$
(E)	$M = \begin{pmatrix} A & 0 & Be^{i\phi} \\ 0 & 0 & Ce^{i\varphi} \\ Be^{i\phi} & Ce^{i\varphi} & 0 \end{pmatrix}$
(F)	$M = \begin{pmatrix} A & Be^{i\phi} & 0 \\ Be^{i\phi} & 0 & Ce^{i\varphi} \\ 0 & Ce^{i\varphi} & 0 \end{pmatrix}$

Leptogenesis at the seesaw scale

Lepton-number-violating decays:

$$N_i \rightarrow l + \phi^+$$

$$N_i \rightarrow l^c + \phi$$

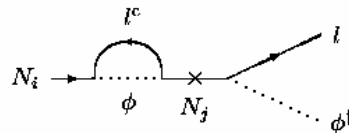


CP violation

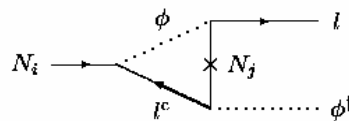
Fukugita,
Yanagida **86**



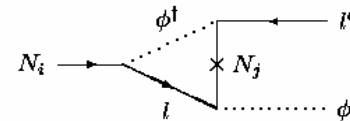
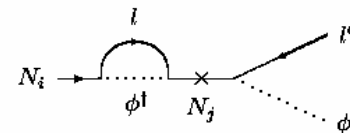
(a)



(b)



(c)



Baryogenesis via Leptogenesis

- ★ Out of thermal equilibrium and N_1 decay

$$T \sim M_1 \gg T_{EW}, \Gamma_1 < H(T \sim M_1)$$

- ★ Net lepton asymmetry

$$Y_L \equiv \frac{n_l - n_{\bar{l}}}{s} \approx \varepsilon_1 \frac{d}{g_*}$$

- ★ Net baryon asymmetry via sphaleron

$$Y_B \equiv \frac{n_B - n_{\bar{B}}}{s} \approx -0.35 Y_L$$

Observation:

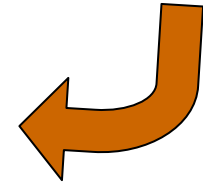
$$Y_B \equiv \frac{1}{7} \frac{n_B}{n_\gamma} \equiv \frac{10^{-10}}{7} \eta_{10}$$

$$\eta_{10} = \begin{cases} 5.6 \pm 0.5 & (\text{BBN, } T \sim 1\text{MeV}) \\ 5.1 \pm 1.6 & (\text{SNIa, } T \sim 1\text{meV}) \\ 6.0 \pm 0.6 & (\text{CMB, } T \sim 0.1\text{eV}) \\ 6.1^{+0.3}_{-0.2} & (\text{WMAP, }) \end{cases}$$

CP asymmetry in our ansatz

$$\varepsilon_1 \approx -\frac{3M_1}{16\pi} \left[\frac{\text{Im}[(Y_\nu^\dagger Y_\nu)_{12}]^2}{M_2(Y_\nu^\dagger Y_\nu)_{11}} + \frac{\text{Im}[(Y_\nu^\dagger Y_\nu)_{13}]^2}{M_3(Y_\nu^\dagger Y_\nu)_{11}} \right] \quad \begin{matrix} M_1^2 \ll M_2^2 \ll M_3^2 \\ r^2 \leq 0.1 \end{matrix}$$

$$\varepsilon_1 = -\frac{3M_1 m_3 \mathcal{I}_\alpha (1-r)^2 (1+r^2) \sin^3 \omega \cos \omega}{2\pi \langle \phi \rangle^2 \left[r^2 (1+r)^2 + 4(1-r)(1+r^2) \sin^2 \omega \right]}$$



where $\omega \equiv (\phi_D - \phi_R)/2$

Dilution factor $\left\{ \begin{array}{l} d \approx 0.02 \times (0.01 \text{ eV}/\tilde{m}_1)^{1.1} \quad \text{(A)} \\ \frac{1}{d} \approx \frac{3.3 \times 10^{-3} \text{ eV}}{\tilde{m}_1} + \left(\frac{\tilde{m}_1}{5.5 \times 10^{-4} \text{ eV}} \right)^{1.16} \quad \text{(B)} \end{array} \right.$

in which $\tilde{m}_1 \equiv \frac{(Y_\nu^\dagger Y_\nu)_{11}}{M_1} \langle \phi \rangle^2 = m_1 \mathcal{I}_\alpha \left[1 + \frac{4(1-r)(1+r^2) \sin^2 \omega}{r^2 (1+r)^2} \right]$

The allowed ranges of M_1 and ω

Input

$$\Delta m_{31}^2 = 2.5 \times 10^{-3} \text{eV}^2$$

$$m_H = 144 \text{ GeV}$$

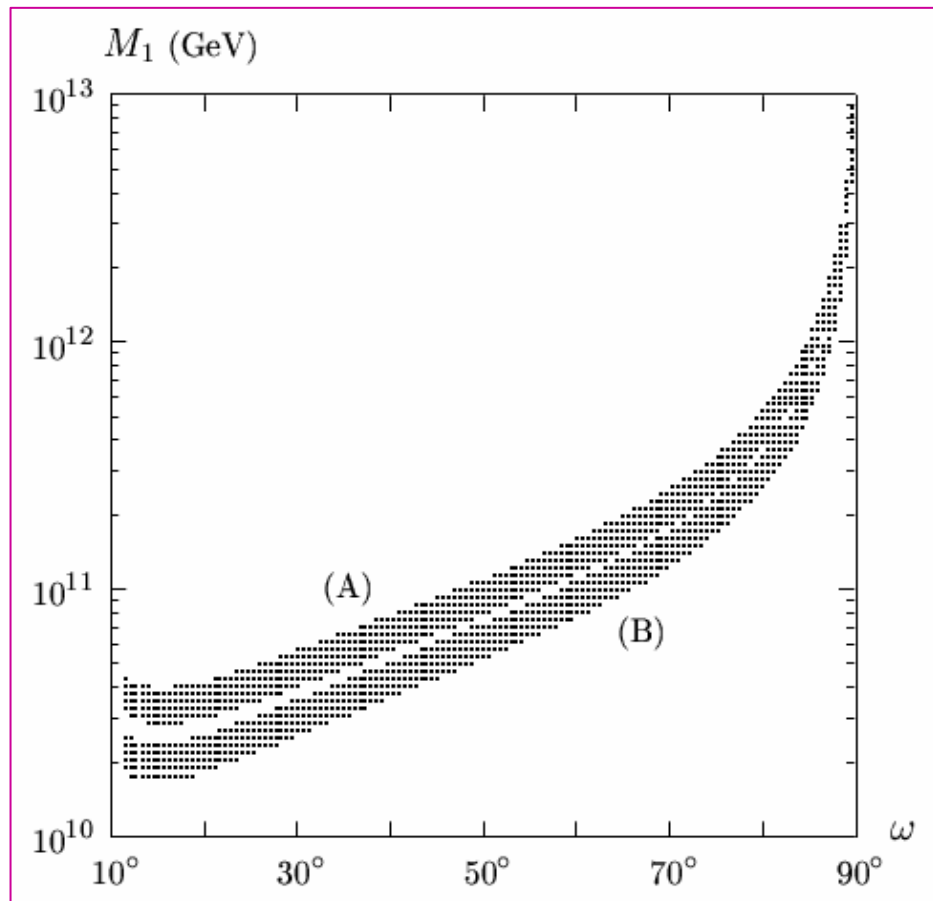
$$r = 0.25$$

$$g_* = 106.75$$

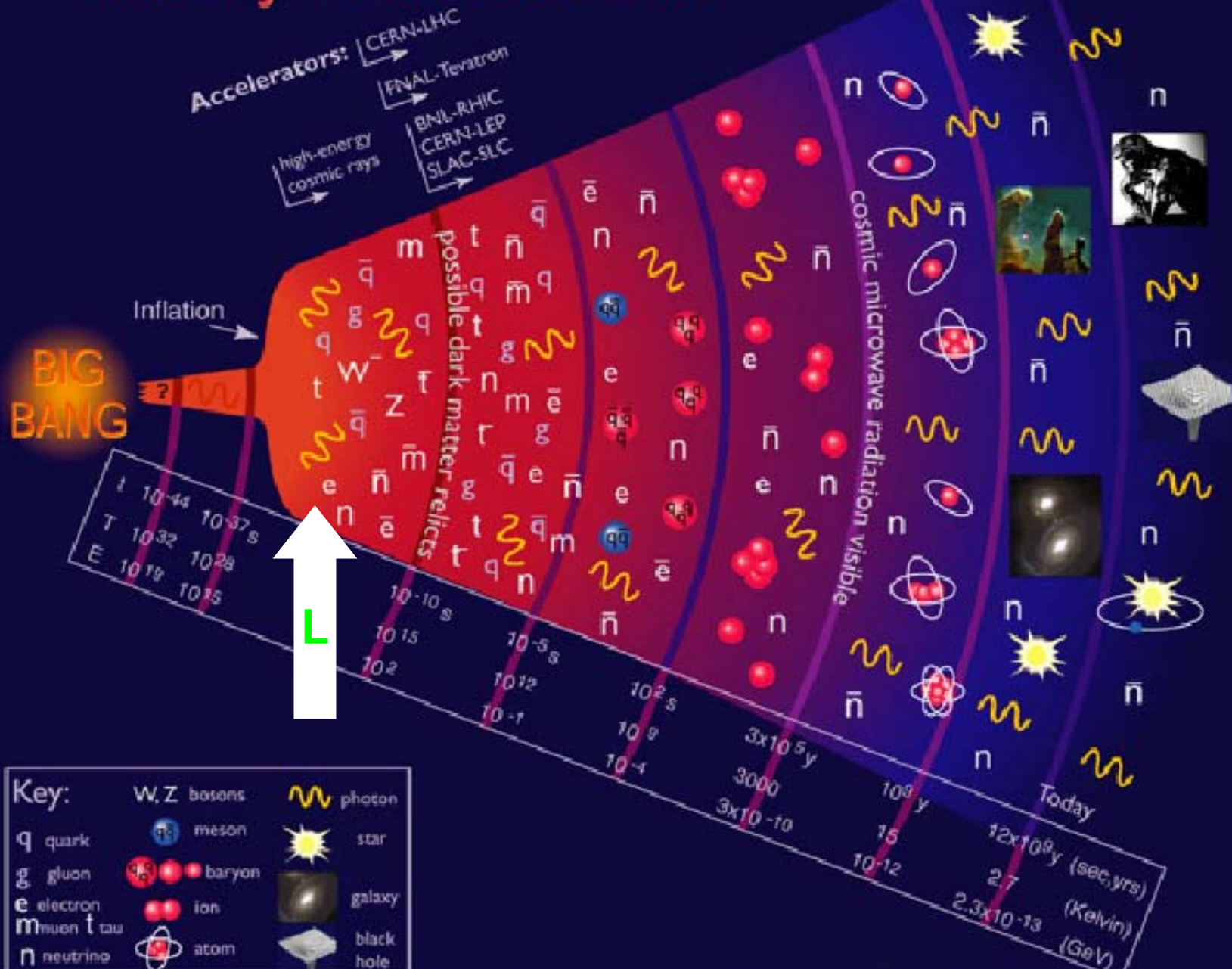
M_1, ω
random



$$7 \times 10^{-11} \lesssim Y_B \lesssim 10^{-10}$$



History of the Universe



Some further comments:

- The cosmological baryon asymmetry and the present ν -oscillation data can simultaneously be interpreted in this phenomenological ansatz

- Y_B can be extended to the MSSM for $\tan \beta \leq 50$

- The CP-violating asymmetry ε_1 has no direct connection with the low-energy CP violation in neutrino oscillations

- Replace $\frac{m_1}{m_2} = \frac{m_2}{m_3} = \frac{M_1}{M_2} = \frac{M_2}{M_3}$ by $\frac{m_1}{m_2} \sim \frac{m_2}{m_3} \sim \frac{M_1}{M_2} \sim \frac{M_2}{M_3}$

The result is stable, but more free parameters.

Concluding question

Leibniz: why is there something rather than nothing?



For neutrino masses, the considerations have always been qualitative, and, despite some interesting attempts, there has never been a convincing quantitative model of the neutrino masses.

Thanks