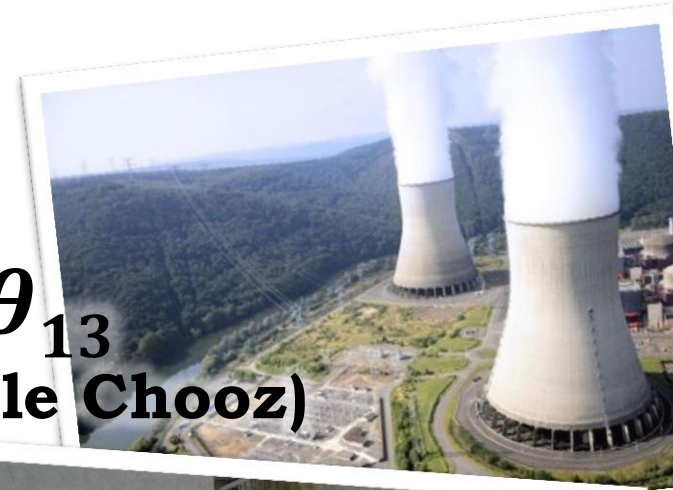




THE UNIVERSITY OF  
**WARWICK**

# **(Direct) Measurement of $\theta_{13}$** **(T2K, MINOS, Reno, Daya Bay, Double Chooz)**

Phillip Litchfield (Kyoto/Warwick)

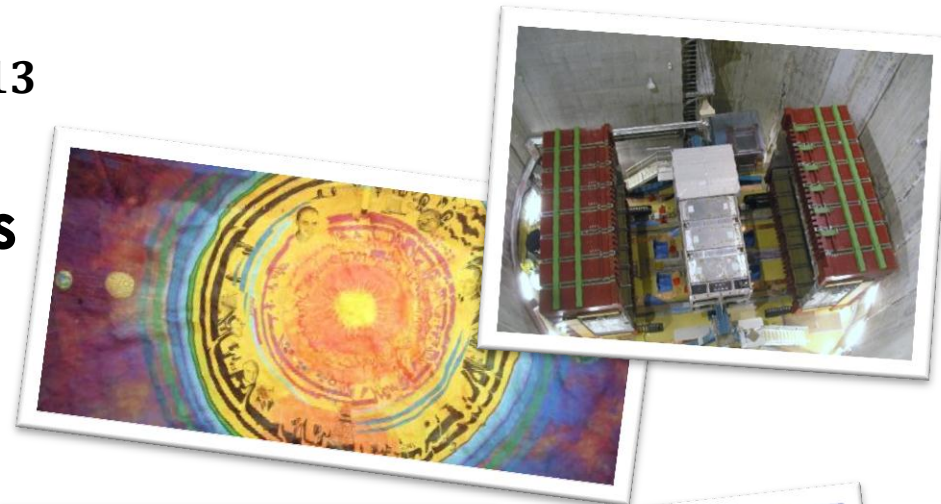


FPCP'12,  
USTC

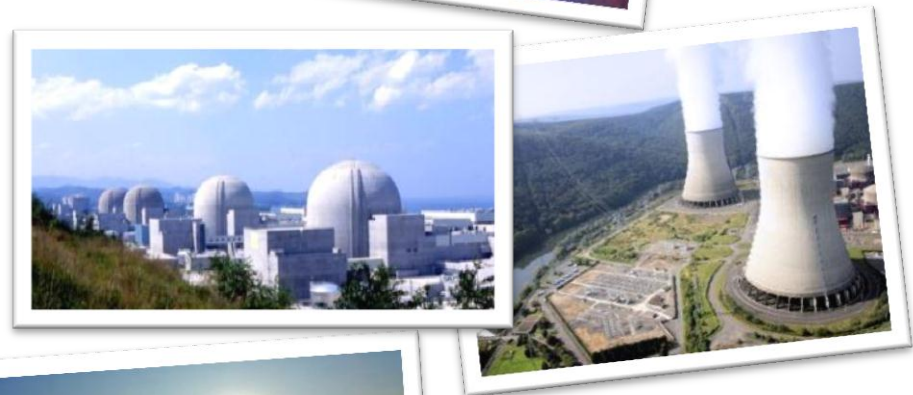


# Outline

- Role and significance of  $\theta_{13}$
- Long-baseline experiments



- Reactor experiments



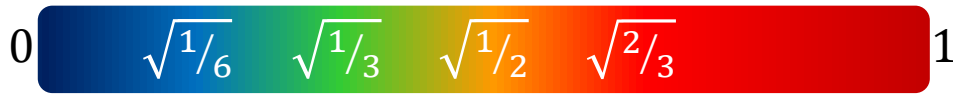
- Next talk  
[Jianglai LIU]



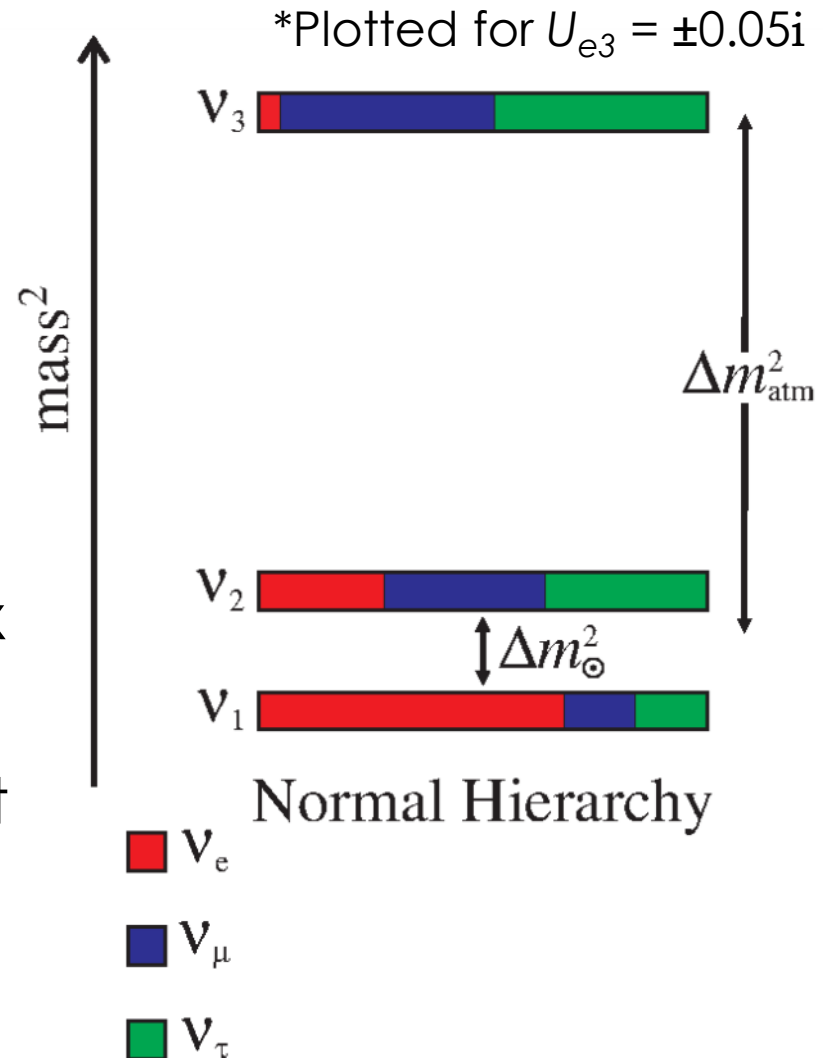
# Neutrino mixing

Neutrino mixing (PMNS) matrix is:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$



- Very different from CKM matrix
- 8 large elements
- $U_{e3}$  is significant as the smallest element, and the last to be measured.





# Angle parameterisation

---

The mixing matrix is commonly parameterised as the product of two rotations and a unitary transformation. Writing  $s_{ij} = \sin\theta_{ij}$ , and  $c_{ij} = \cos\theta_{ij}$ :

$$\begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix}$$

The choice of parameterisation is convenient as the **solar** and **atmospheric** disappearance amplitudes are well approximated as functions of  $\theta_{12}$  and  $\theta_{23}$ , respectively. This approximation only works because the third angle  $\theta_{13}$  is small.

# The role of $\theta_{13}$

---

In the standard parameterisation, it turns out that

$$U_{e3} = \sin \theta_{13} e^{-i\delta}, \quad \text{and therefore} \quad \sin \theta_{13} = |U_{e3}|.$$

The value of  $\sin \theta_{13}$  is particularly significant because a zero element in the mixing matrix would eliminate the possibility of (KM-mechanism) leptonic CP violation.

**The future program of neutrino physics is strongly dependent on a non-zero measurement ( $\rightarrow$  size) of  $\theta_{13}$ .**

To study, need channels involving  $\langle \nu_e | \nu_3 \rangle$ . The most accessible are  $\bar{\nu}_e \rightarrow \bar{\nu}_e$  and  $\nu_\mu \rightarrow \nu_e$  at first 'atmospheric' maximum ( $L/m \sim 500 \times E/\text{MeV}$ )

# Outline

- Role and significance of  $\theta_{13}$
- Long-baseline experiments



- Reactor experiments



- Next talk  
[Jianglai LIU]



# $\nu_\mu \rightarrow \nu_e$ appearance probability

---

The  $\nu_e$  appearance probability can be written approximately as a sum of terms quadratic in the small parameters  $\alpha = \Delta m_{21}^2 / \Delta m_{31}^2 \approx 1/32$ , and  $\sin 2\theta_{13}$ :

$$P(\nu_\mu \rightarrow \nu_e) \approx T_{\theta\theta} \sin^2 2\theta_{13} \frac{\sin^2([1-A]\Delta)}{[1-A]^2} + T_{\alpha\alpha} \alpha^2 \frac{\sin^2(A\Delta)}{A^2} \\ + T_{\alpha\theta} \alpha \sin 2\theta_{13} \frac{\sin([1-A]\Delta)}{(1-A)} \frac{\sin(A\Delta)}{A} \cos(\delta + \Delta)$$

where

$$T_{\theta\theta} = \sin^2 \theta_{23}, \quad T_{\alpha\alpha} = \cos^2 \theta_{23} \sin^2 2\theta_{12}, \\ T_{\alpha\theta} = \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{23}$$

and  $\Delta = \frac{\Delta m_{31}^2 L}{4E} \sim \frac{\pi}{2}$  at 1<sup>st</sup> osc. maximum.

$A (= \pm 2\sqrt{2}G_F n_e E / \Delta m_{31}^2)$  is the matter density parameter; MINOS  $|A| \sim 0.3$ , T2K  $|A| \sim 0.07$

# $\nu_e$ appearance measurements

---

Two experiments have recently published results from the  $\nu_\mu \rightarrow \nu_e$  channel: **MINOS** and **T2K**.

**MINOS** is an older experiment, not optimised for  $\nu_e$  appearance, but has been running for years.

**T2K** is newer & will be much more sensitive, but so far has only a small fraction of its planned integrated luminosity.

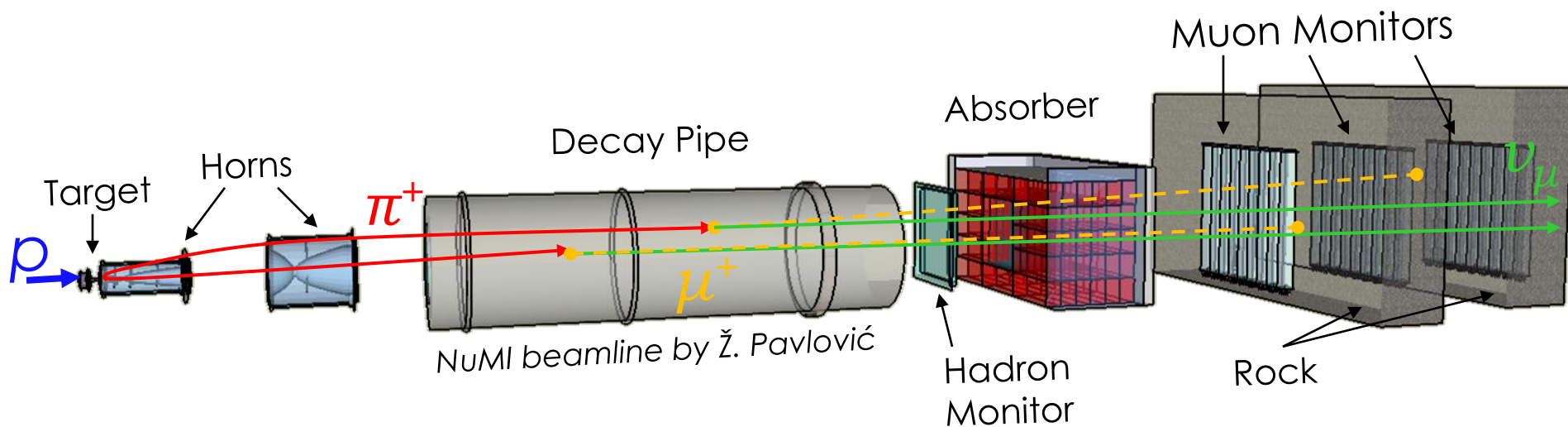
Both experiments follow the same basic principle:

- A beam  $\nu_\mu$  of from pion decay.
- A near detector to measure interaction rates
- A far detector to look for oscillations



# Muon neutrino beams.

- $\nu_\mu$  from pion decay.
- Pions produced in proton interactions on a target, and focussed by magnetic horns (NuMI/MINOS: 2, T2K: 3)
- Horn current and geometry most important in determining (on-axis) spectrum
- Wrong sign (anti- $\nu$ ) and  $\nu_e$  backgrounds are  $\sim$  few %

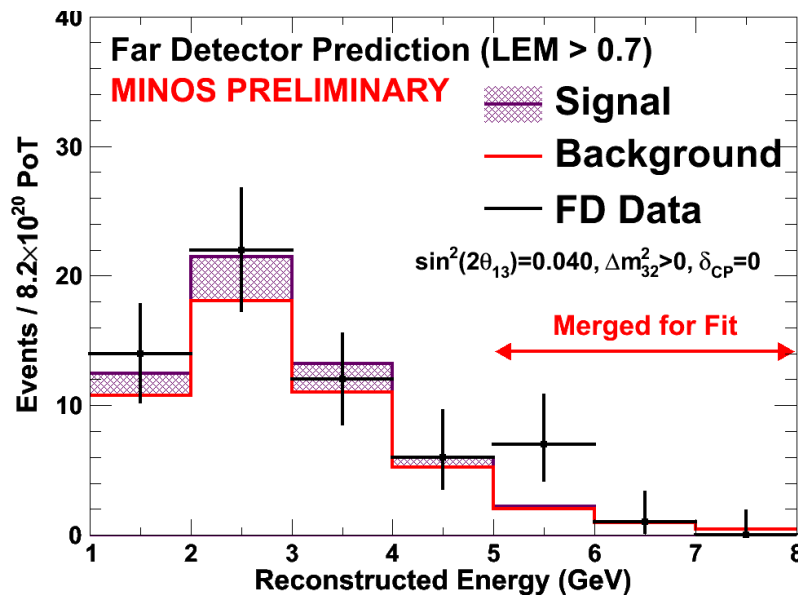
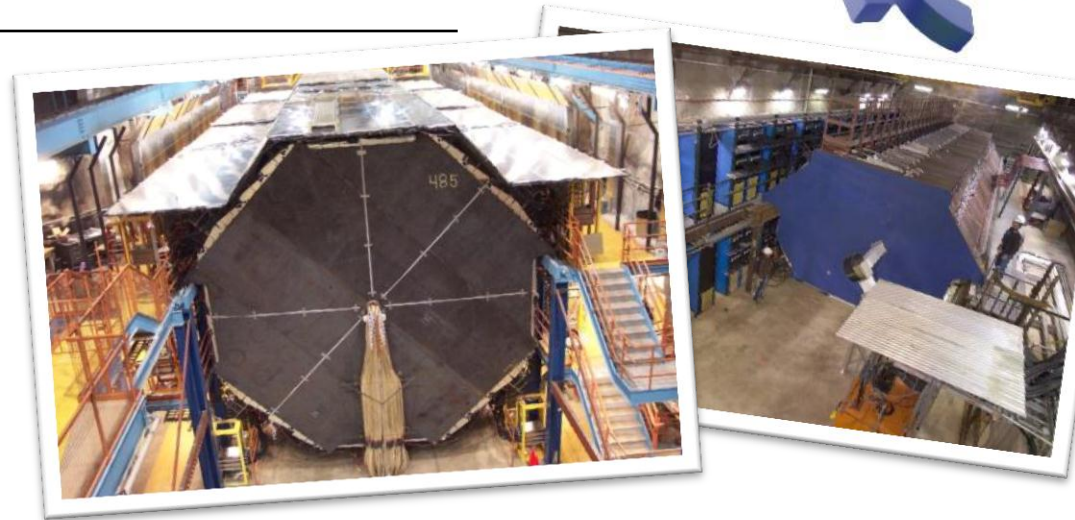


# MINOS analysis



Both MINOS detectors are steel/scintillator tracking calorimeters.

At peak energy ( $\sim 3\text{GeV}$ )  $\nu_e$  signal events appear as small, tight showers.



Analysis approach uses the fact both detectors are functionally identical.

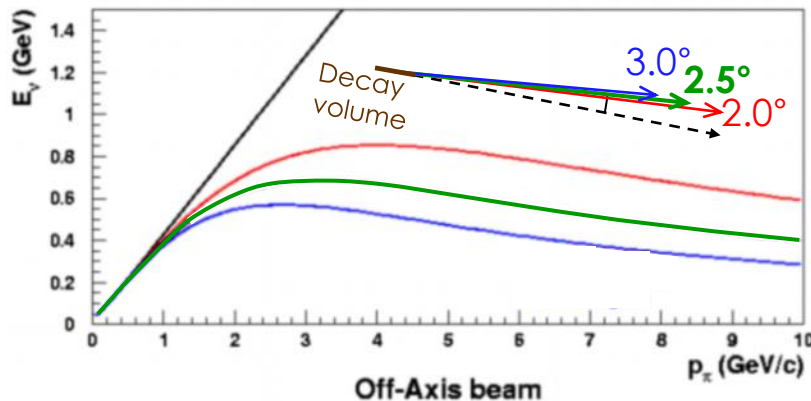
NC interactions of higher energy neutrinos are the major background: well-controlled, but limits sensitivity because of its size.

# T2K: The Off-axis 'trick'



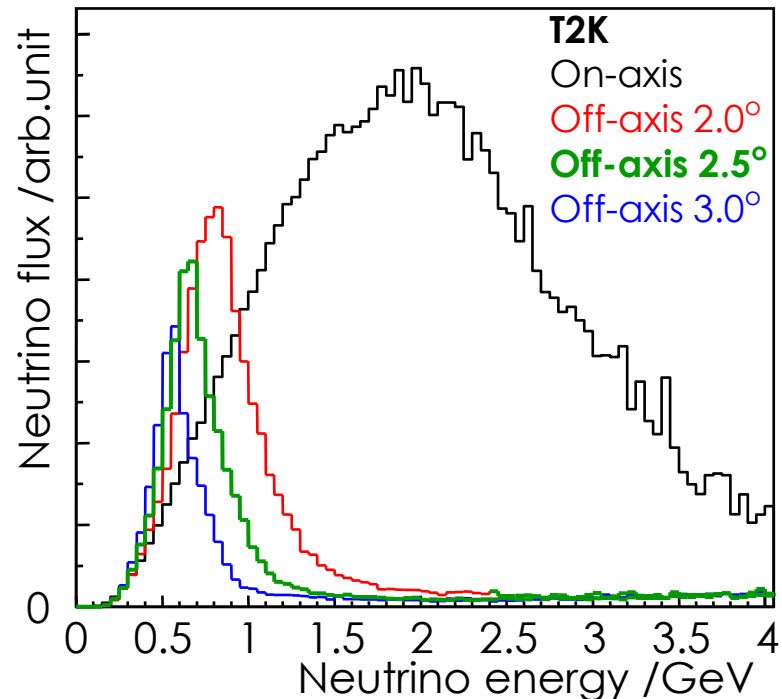
**T2K is the first experiment to have its detectors off-axis**

Relativistic kinematics  $\rightarrow$  at a small angle to the beam axis, neutrino energy is insensitive to parent pion energy.



Gives slightly narrower flux peak, and **drastically reduces high energy tail.**

- Ideal for  $\nu_e$  appearance (much reduced NC BG)



# T2K analysis

---

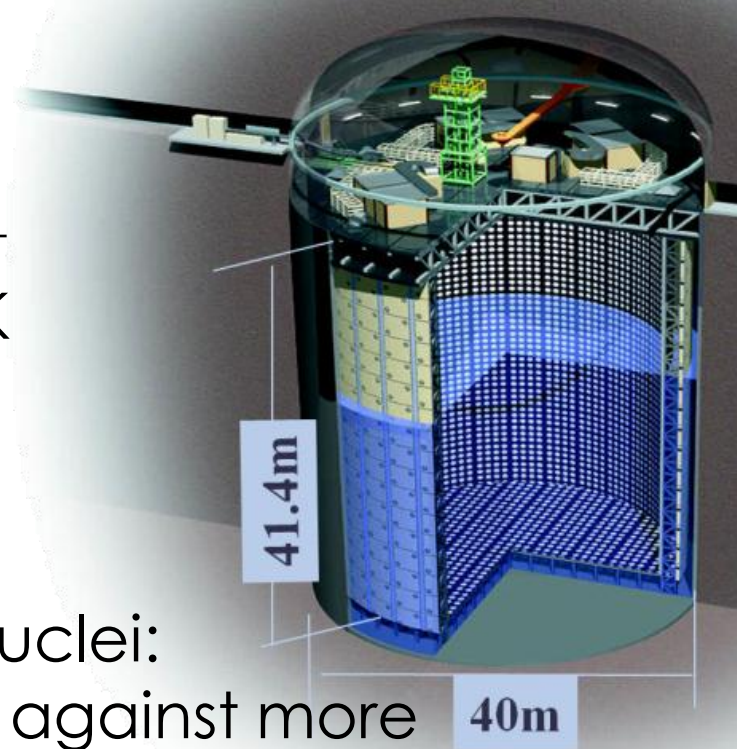
T2K uses the 22.5 kt (fiducial) Super-K water Čerenkov detector to look for  $\nu_e$  appearance.

Signal is  $\nu, n \rightarrow p, e^-$  (CCQE) on  $^{16}\text{O}$  nuclei:  
Results in a fuzzy Čerenkov ring. PID against more common  $\mu^-$  events is based on the fuzziness of this ring.

Other important backgrounds from

- **Beam  $\nu_e$  contamination** (removed by energy cut)
- **NC- $\pi^0$  events** (removed by fitting for a 2<sup>nd</sup>  $\gamma$  ring)

T2K uses exclusive channels, so need good understanding of cross-sections  $\rightarrow$  drives design of the Near Detector.





# T2K-ND280 near detectors

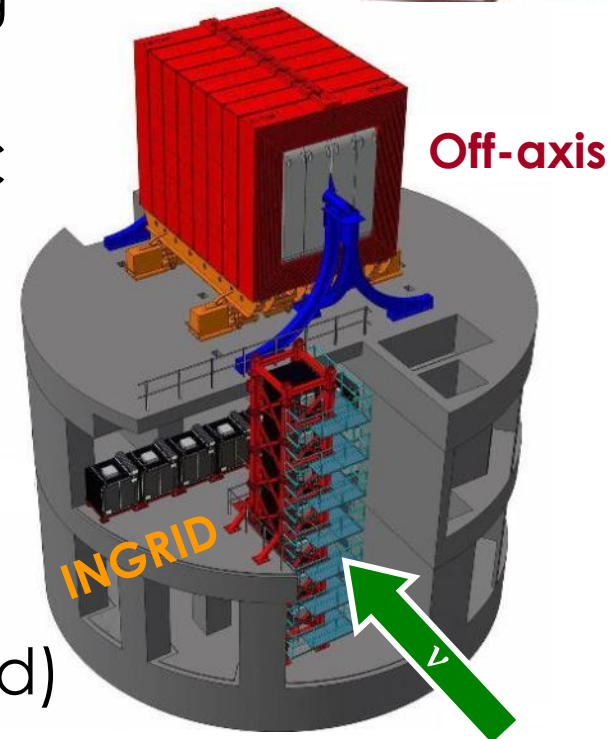
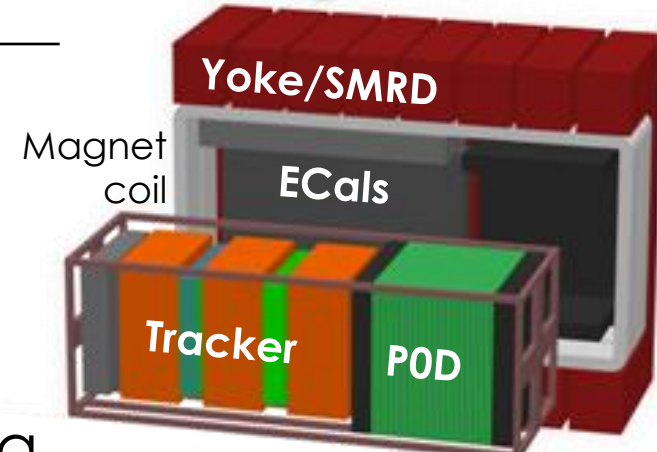


**Off-axis** detector uses fine-resolution scintillator detectors and TPCs for momentum measurement and PID.

- Much better estimate of interaction rates than using MC dead-reckoning
- Extensive program of cross-section measurements to improve future MC

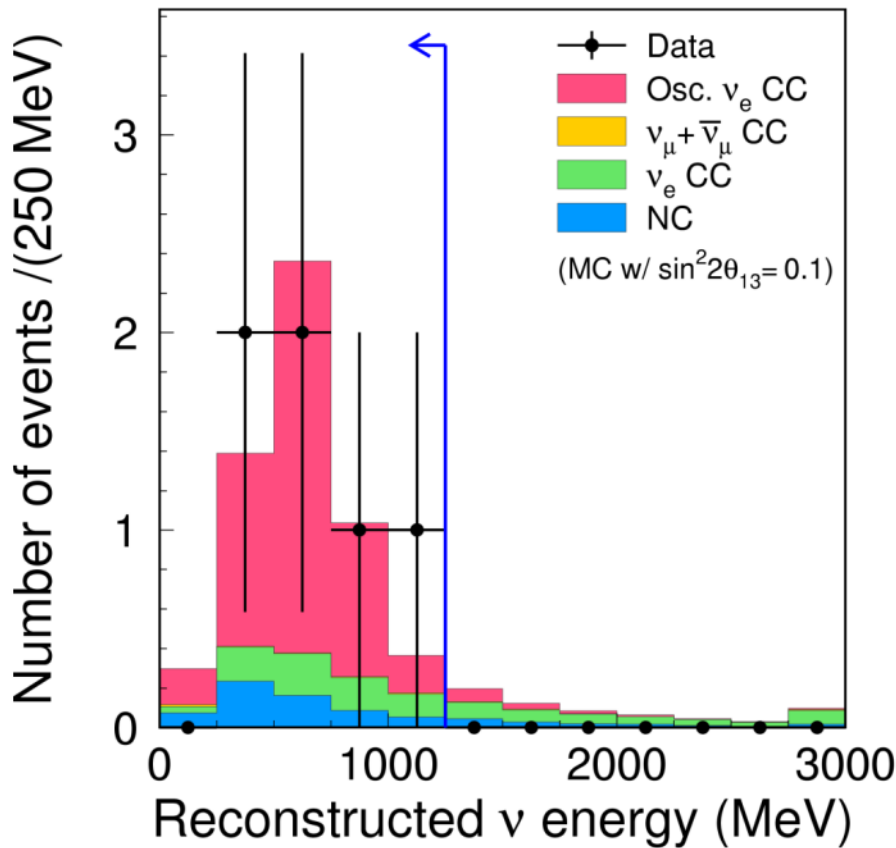
**INGRID:** A cross arrangement of iron-scintillator detectors, centred on the beam axis, to check beam direction.

- Stable to better than required (1 mrad)



## Current results dominated by statistical error.

(Data set up to 11/03/11 earthquake)



Backgrounds are much smaller than MINOS,  
↳ sensitive measurement with only six FD events.

Current results use ND280 for rate normalisation.

Future results will need the improved constraints from detailed ND820 analyses

# T2K & MINOS: side-by-side



## T2K

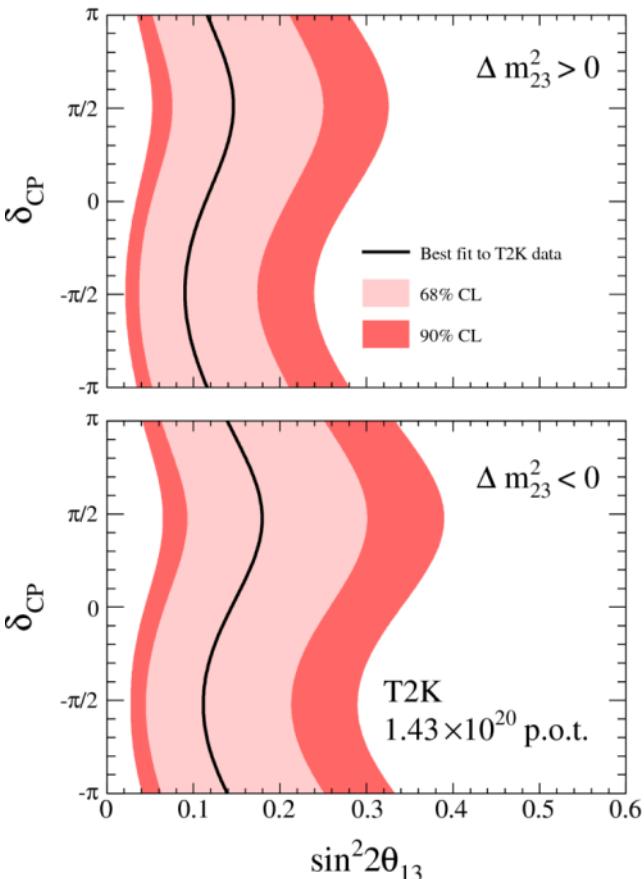
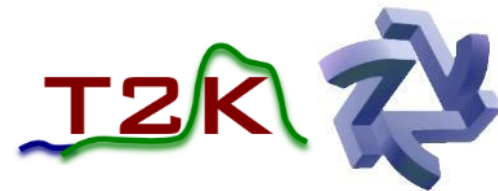
- Small data sample, but very low background.
- Systematics controlled by high-resolution Near Detector measurements.
- Precise flux inputs from NA61 experiment.
- Exclusive analysis in sub-GeV region.
- New experiment, much more (stats, analysis) to come!

## MINOS

- Large data sample, but high backgrounds.
- Many systematics cancel out because of similarity of Near & Far detectors.
- Flexible beam helps separate flux from cross-section errors.
- Inclusive analysis in few-GeV region.
- Mature—data taking finished last month.

**Although experiments are similar in principle, in practice their analyses & systematics are quite different.**

# Results from T2K & MINOS



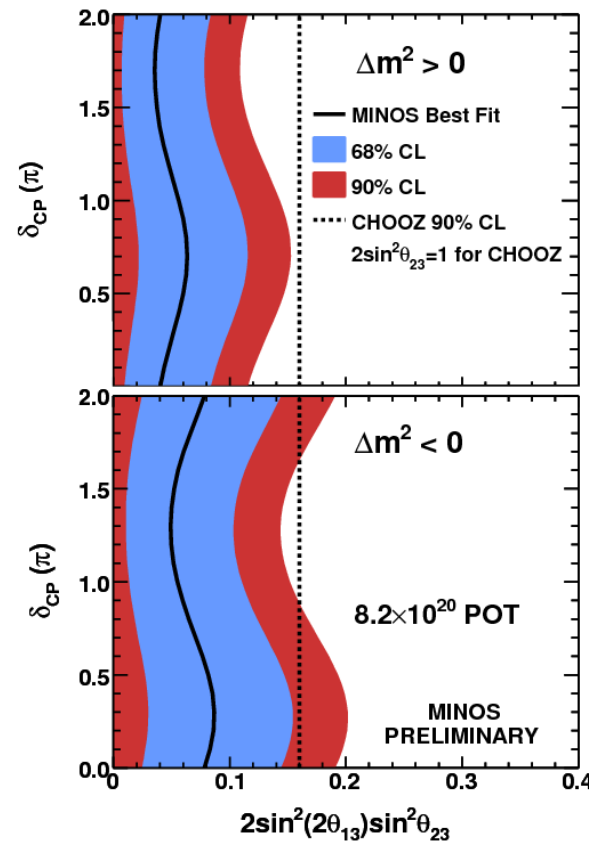
Results vs  $\delta$

► **MINOS**

◄ **T2K**

Note:  
y-axes are offset by  $\pi$

Note:  
T2K holds  $\theta_{12}, \theta_{23},$  &  $\Delta m^2$   
fixed; MINOS varies  
within global errors.  
Difference is small.



T2K:  $\sin^2(2\theta_{13}) = 0$  disfavoured at 99.3%

MINOS:  $\sin^2(2\theta_{13}) = 0$  disfavoured at 89%

arXiv:1106.2822

arXiv:1108.0015



# Outline

- Role and significance of  $\theta_{13}$
- Long-baseline experiments



- Reactor experiments



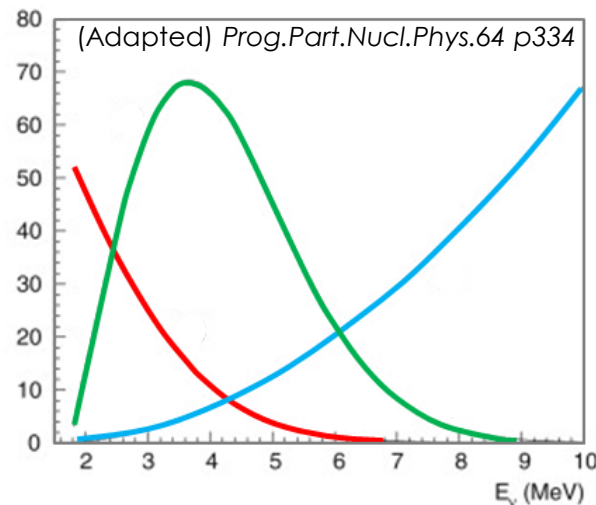
- Next talk  
[Jianglai LIU]



# $\bar{\nu}_e \rightarrow \bar{\nu}_e$ survival

Disappearance channel. Electron is light enough that we can use nuclear reactors.

- Cheap, intense source.
- **Interaction rate** peaks at 3~4 MeV  
(**falling flux** × **rising cross-section**)



(Some) previous reactor experiments:

- **Savannah River**: first  $\bar{\nu}_e$  detection
- **CHOOZ**: previous best measurement of  $\sin 2\theta_{13}$
- **KamLAND**: same channel on longer baseline measures  $\Delta m_{21}^2$  and  $\sin 2\theta_{12}$

New generation of reactor experiments provide much better sensitivity through control of systematics.

# $\bar{\nu}_e$ survival probability

---

For reactor antineutrinos in the earth's crust we can ignore matter effects, and the probability is very simple:

$$P(\nu_e \rightarrow \nu_e) \approx 1 - \sin^2 2\theta_{13} \sin^2 \Delta - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \alpha \Delta$$

On a  $\sim 1$  km baseline the 3<sup>rd</sup> solar-scale term is small, (around 0.001) so the survival probability allows **direct access** to the parameter  $\sin^2 2\theta_{13}$ .

In contrast with the appearance channel,  $\bar{\nu}_e$  disappearance at reactors gives a **theoretically clean determination of  $\sin^2 2\theta_{13}$**  but, on it's own, **cannot determine the mass hierarchy or look for CP violation.**

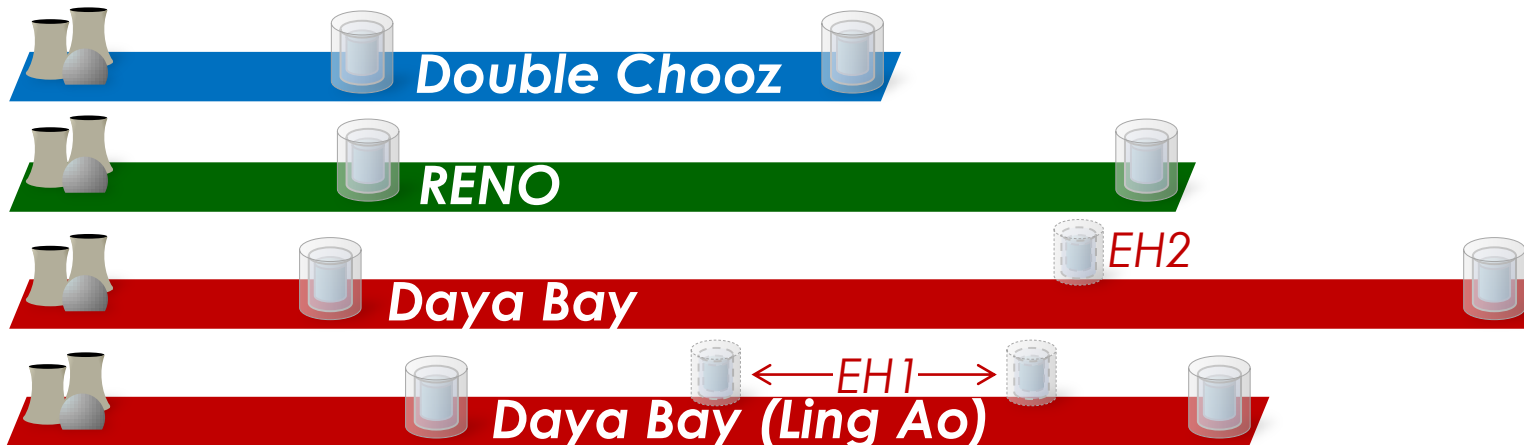
# New reactor experiments



Three new reactor experiments turned on in 2011:

**Double Chooz** (France), **RENO** (S.Korea), **Daya Bay** (China)  
Basic design is very similar among all three, differences are mostly in detector mass & reactor flux.

- But backgrounds, calibration systems, photosensor coverage (etc.) will affect ultimate performance.
- All experiments use a **near detector** to reduce systematics.



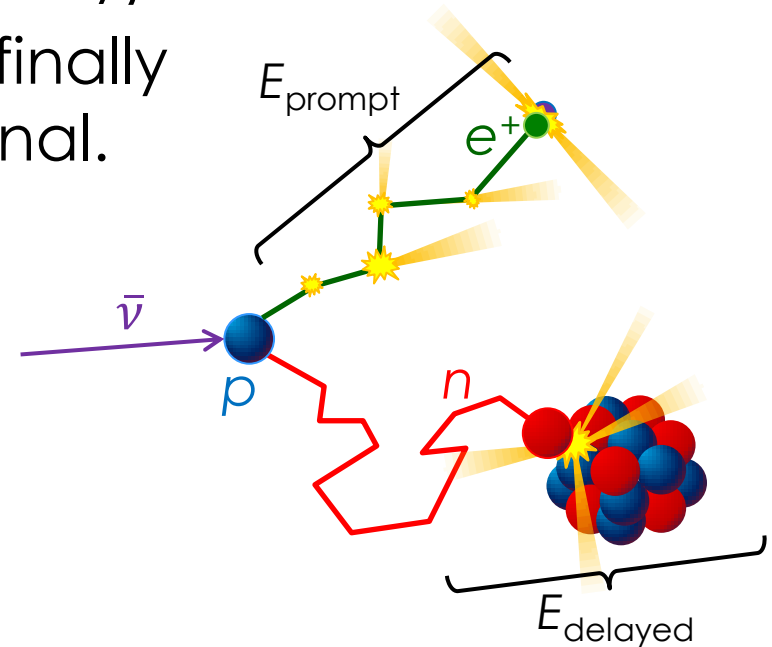


# Principles of $\bar{\nu}_e$ detection



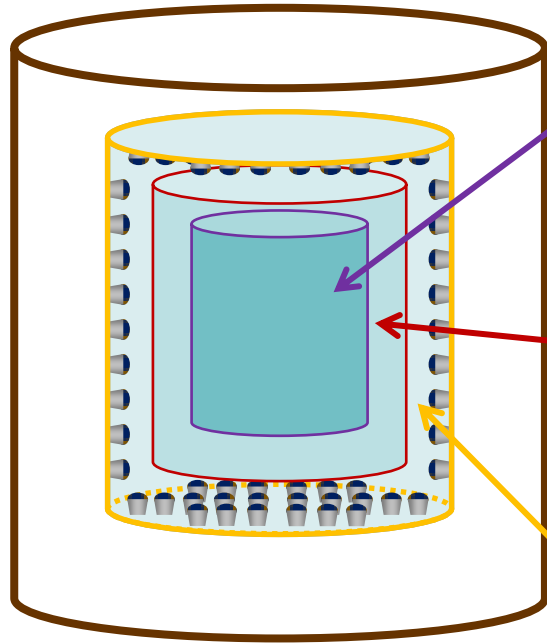
Anti-neutrino interacts on **protons** in a liquid scintillator target (inverse beta decay)

- **Positron** excites scintillator and finally annihilates, giving a prompt signal.
- **Neutron** is detected using Gadolinium dopant (Gd has high neutron capture cross section); this gives a large delayed coincidence signal.



The delayed coincidence signal heavily suppresses backgrounds, so the same principle is used by all three of the new generation reactor experiments.

# A $\bar{\nu}_e$ detector



## Veto tank

Pure water & PMTs  
to veto cosmics and  
external radioactivity

**Target region**, Liquid Scintillator (LS) doped with Gd. Around 2~3m (diameter & height)

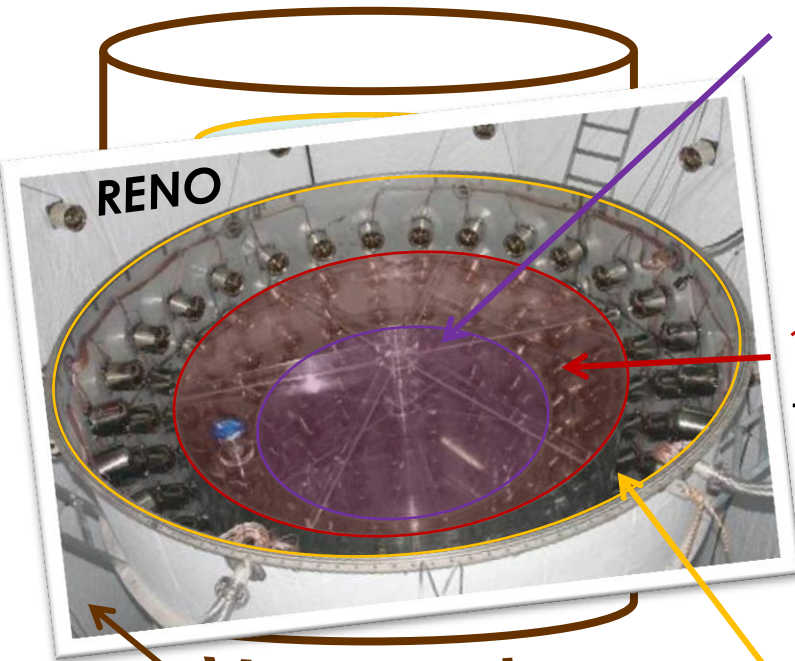
**$\gamma$ -catcher**, LS without Gd doping, typically about 50cm thick

- Improves energy resolution.

**Buffer region**, Mineral oil without scintillator, up to ~1m thick

- Shields active regions

# A $\bar{\nu}_e$ detector



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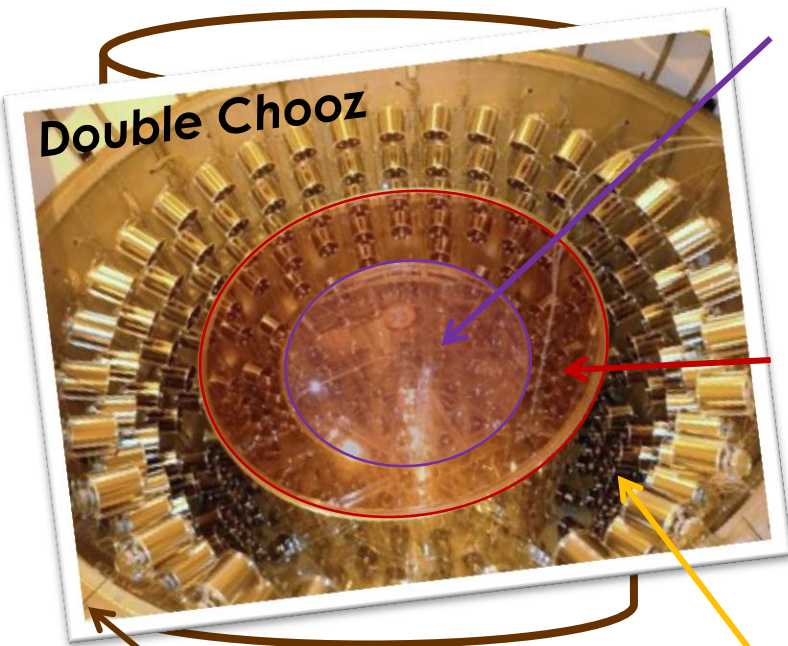
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**Veto tank**

Pure water & PMTs to veto cosmics and external radioactivity

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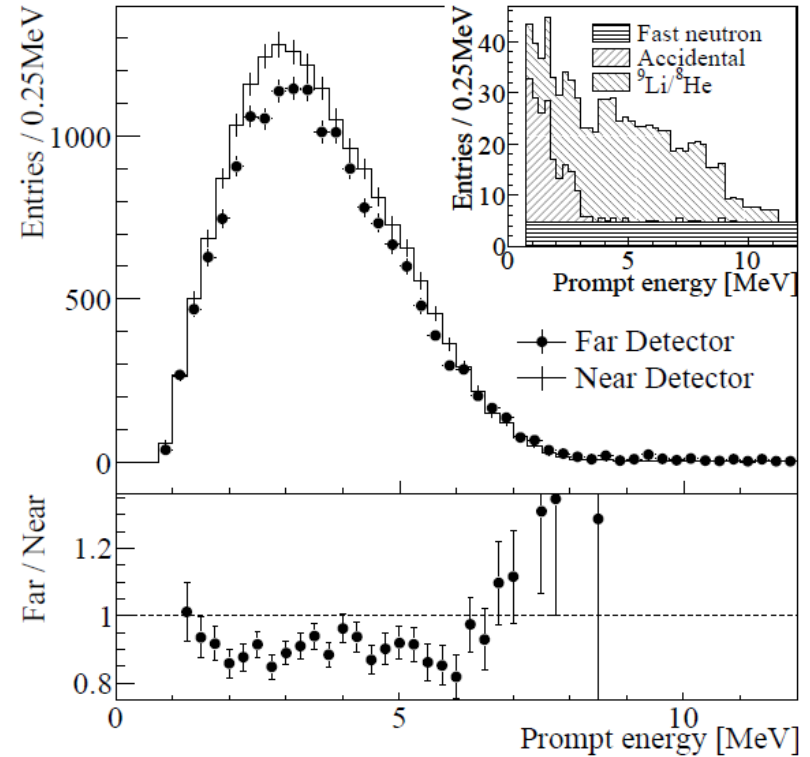
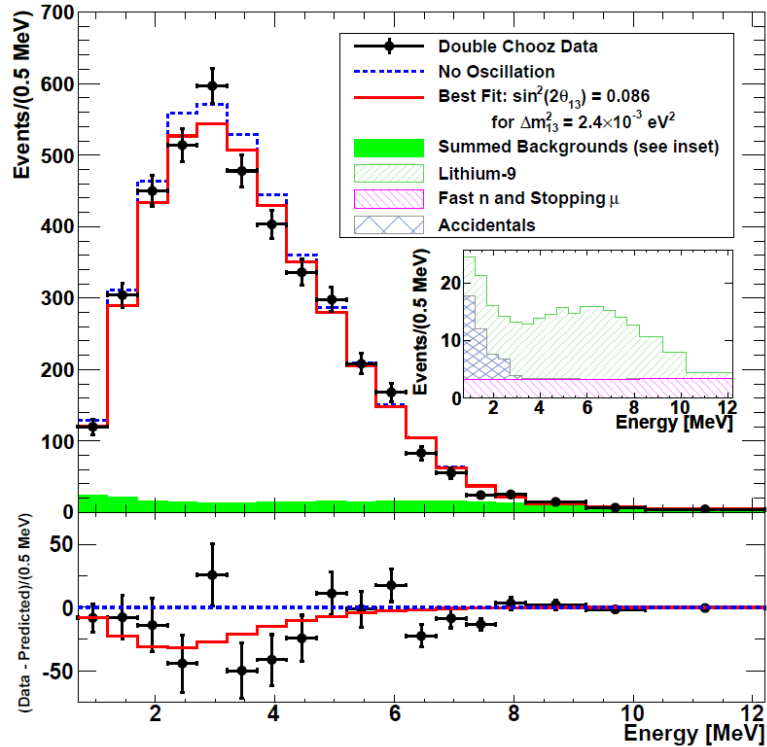
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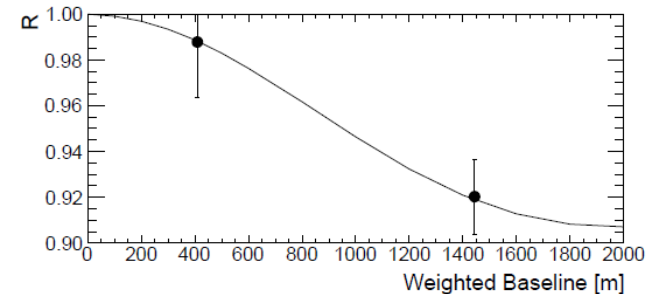
**Veto tank**

Pure water & PMTs to veto cosmics and external radioactivity

# Results from DC and RENO

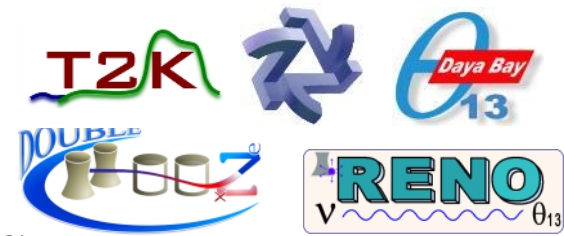


- ▲ Double Chooz result from 1050m only period (2011)
  - No oscillation disfavoured at 94.6%
- ▶ RENO results with two detectors.
  - No oscillation disfavoured at  $4.9\sigma$



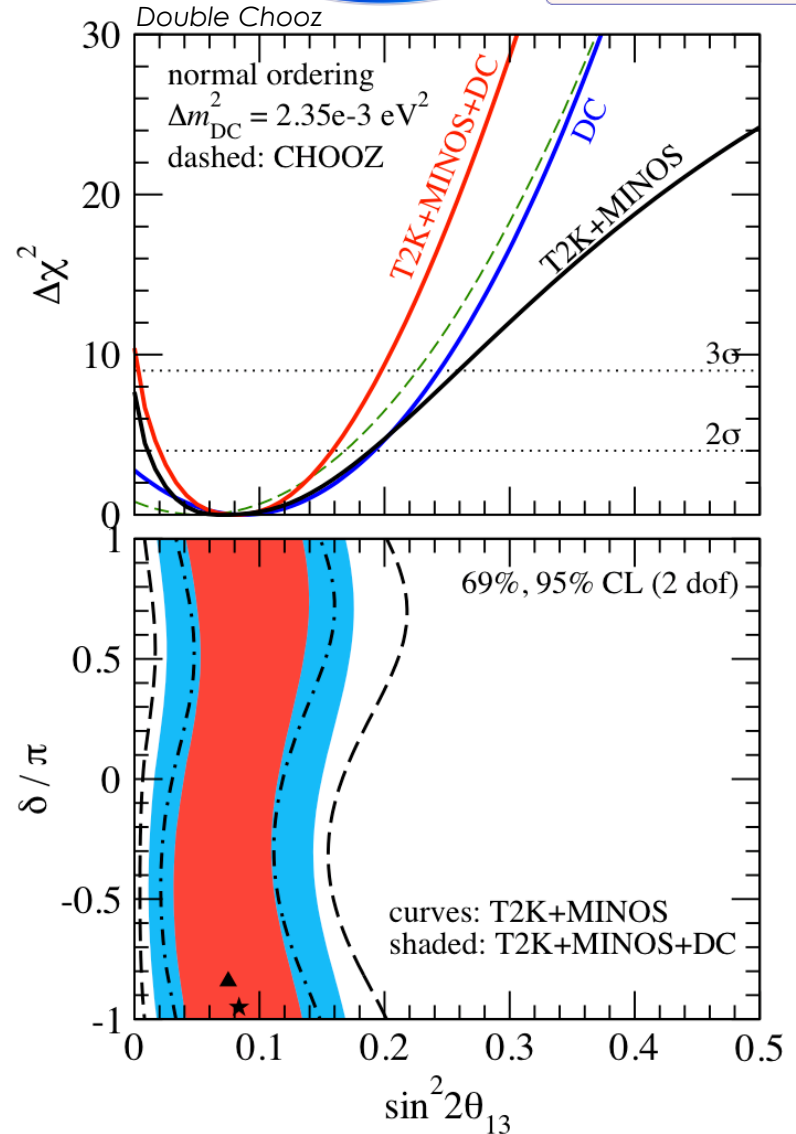
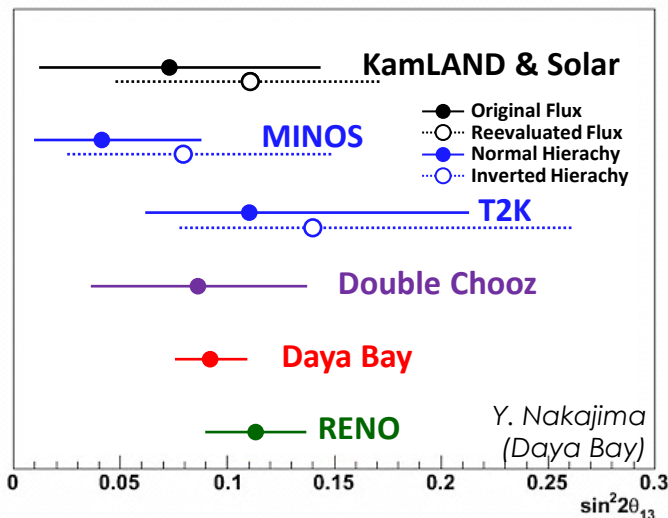


# Our new knowledge



In under a year  $\theta_{13}$  has gone from unknown to well measured.

- At the top of expected range.
- Focus is rapidly shifting to CP phase and mass-hierarchy.
  - Need precision on  $|U_{e3}|^2$ , but also  $|U_{\mu3}|^2$  and  $\Delta m_{atm}^2$



# Thank you for listening!

---



Long baseline results

**T2K** Phys.Rev.Lett.**107** (2011) 041801  
arXiv:1106.2822 [hep-ex]

**MINOS** Phys.Rev.Lett.**107** (2011) 181802  
arXiv:1108.0015 [hep-ex]

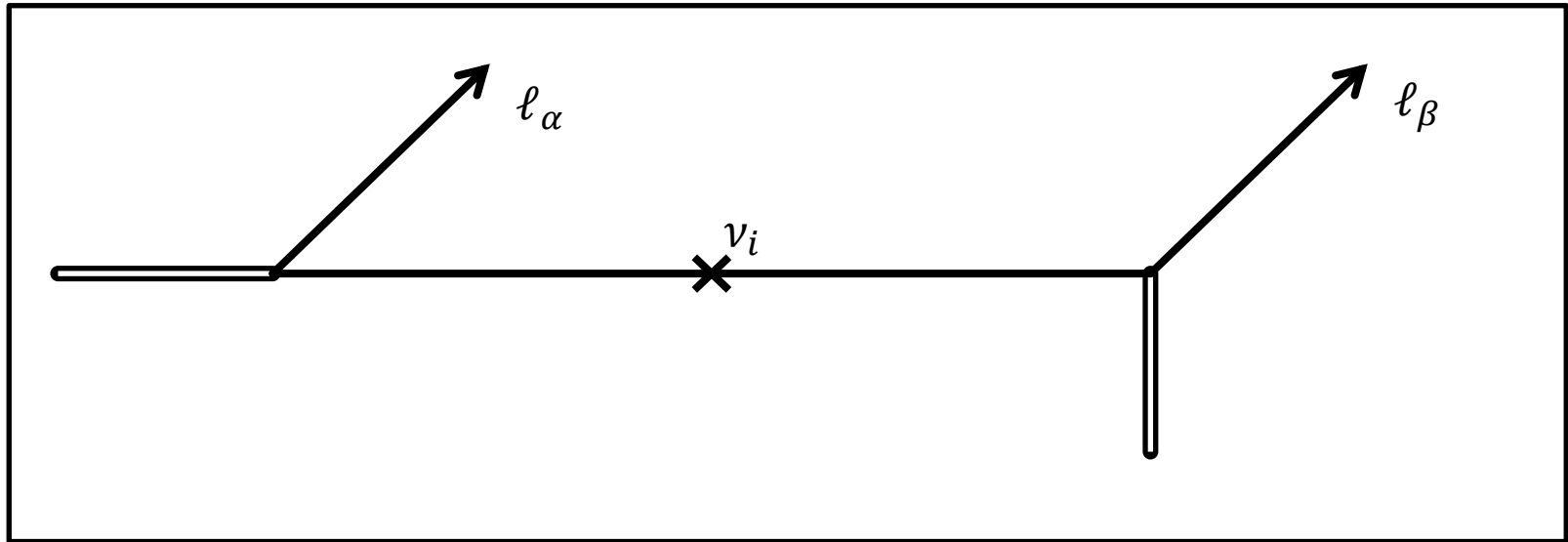
Reactor results

**Double Chooz** Phys.Rev.Lett.**108** (2012) 131801  
arXiv:1112.6353 [hep-ex]

**Daya Bay** Phys.Rev.Lett.**108** (2012) 171803  
arXiv:1203.1669 [hep-ex]

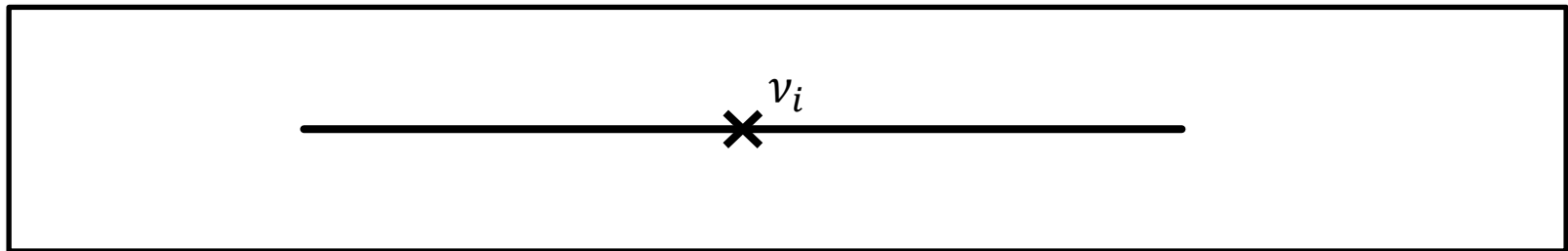
**RENO** Phys.Rev.Lett.**108** (2012) 191802  
arXiv:1204.0626 [hep-ex]

# Neutrino 'box' diagram



My preferred version:

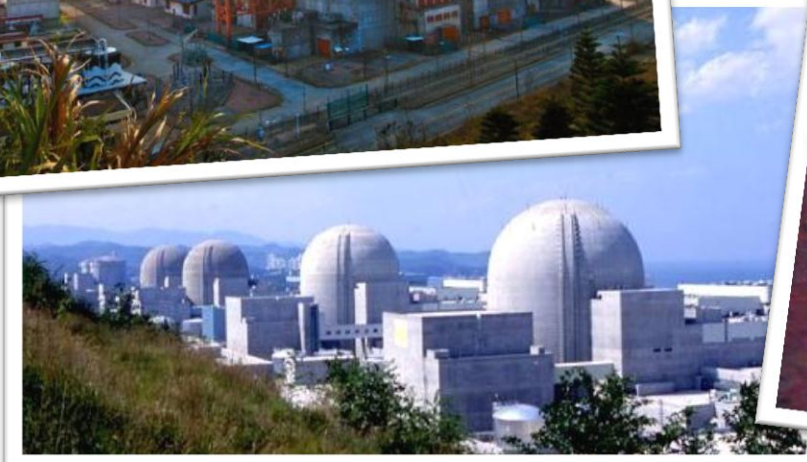
Here is the box



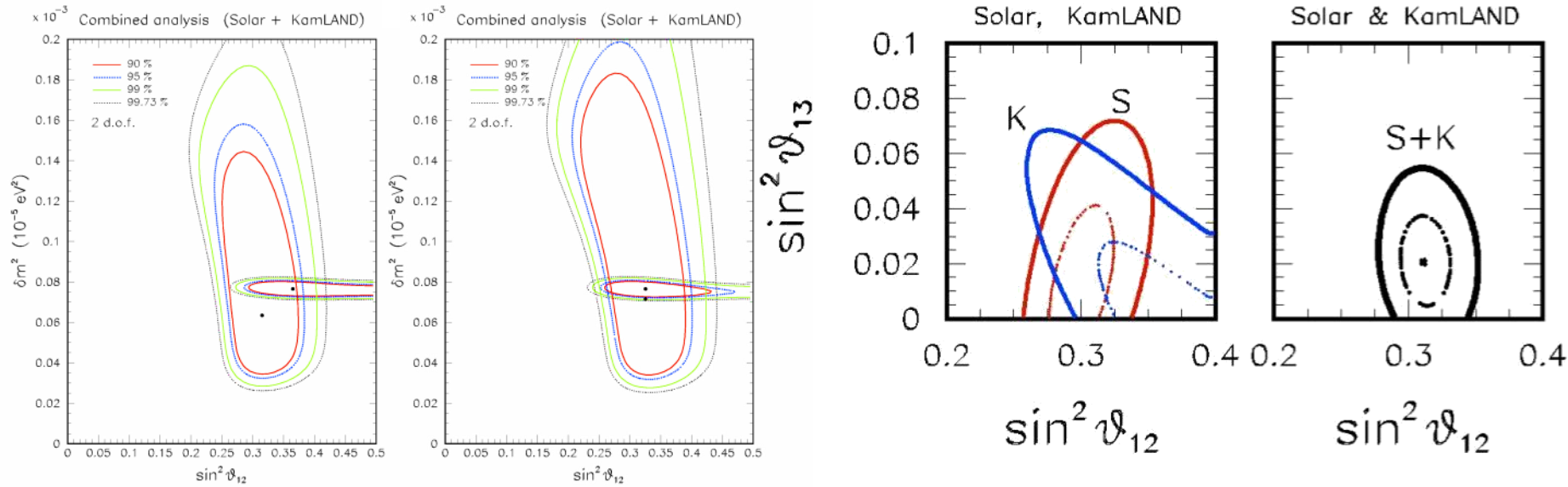


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# Extra slides



# Indirect measurements of $\theta_{13}$

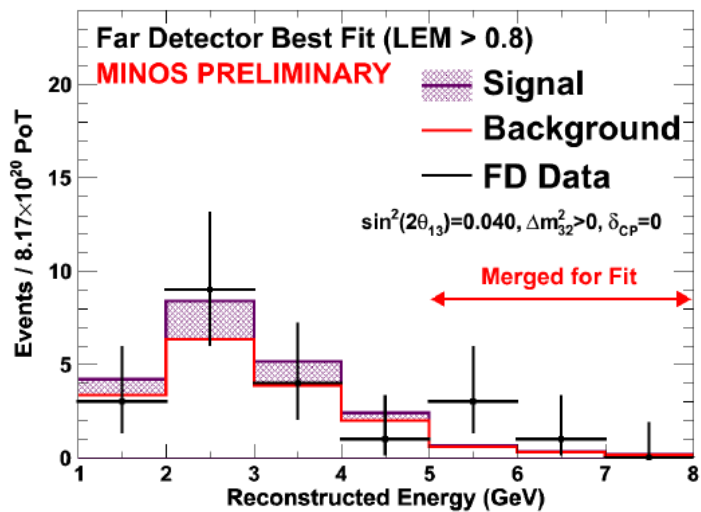
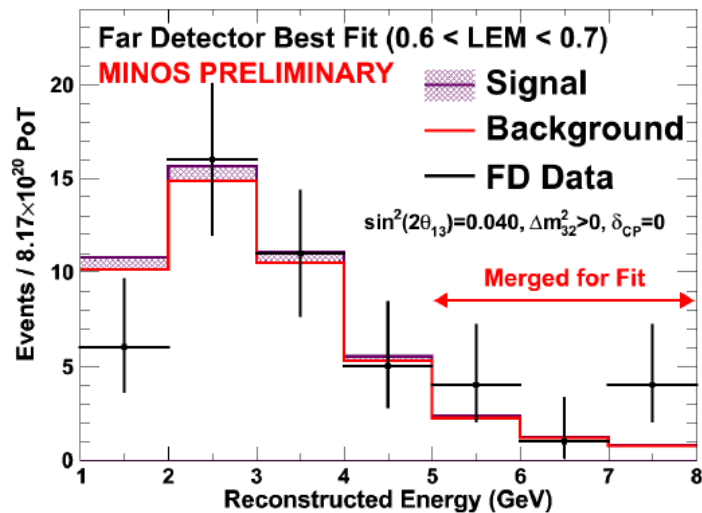
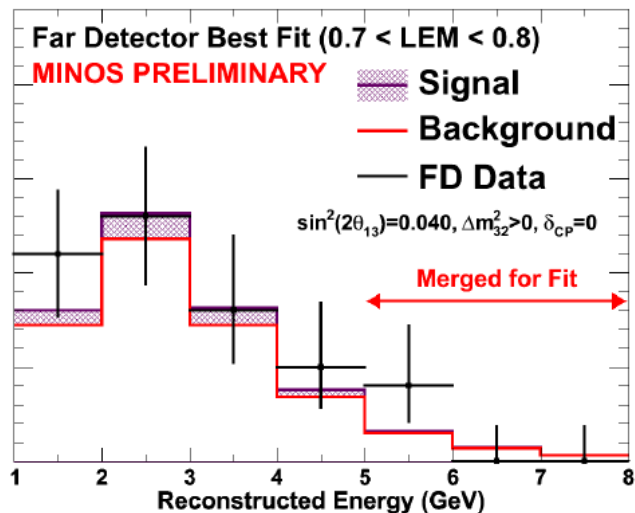


Before recent measurements there was a  $\sim 1.5\sigma$  hint from combining Solar and KamLAND data.

Turned out to agree very well with new measurements, but weak compared to new data.



# Minos results in detail



**For  $\text{LEM} > 0.7$**

Expected background events:

$$49.5 \pm 2.8 \text{ (syst)} \pm 7.0 \text{ (stat)}$$

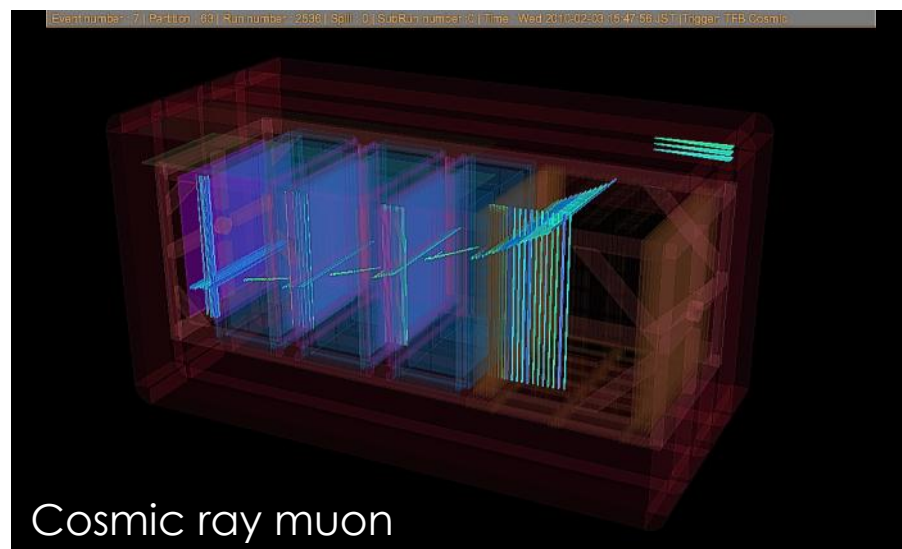
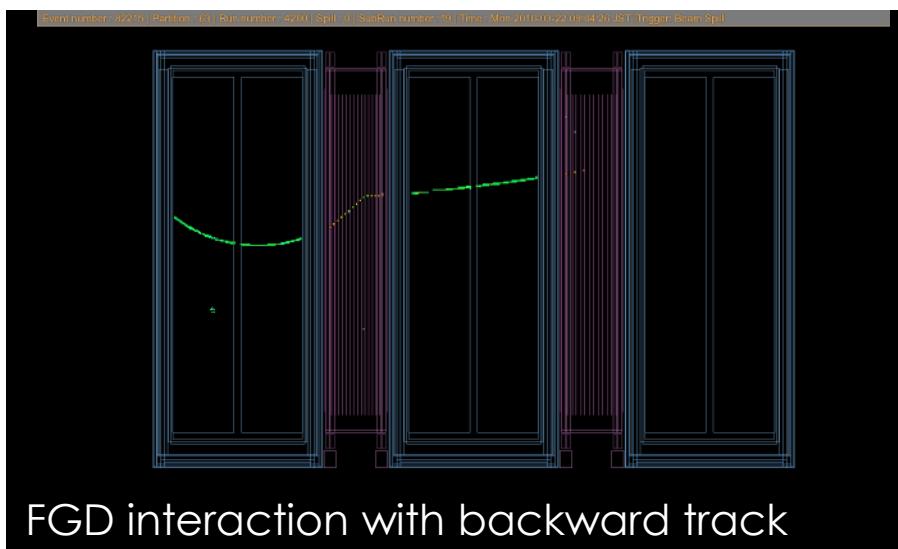
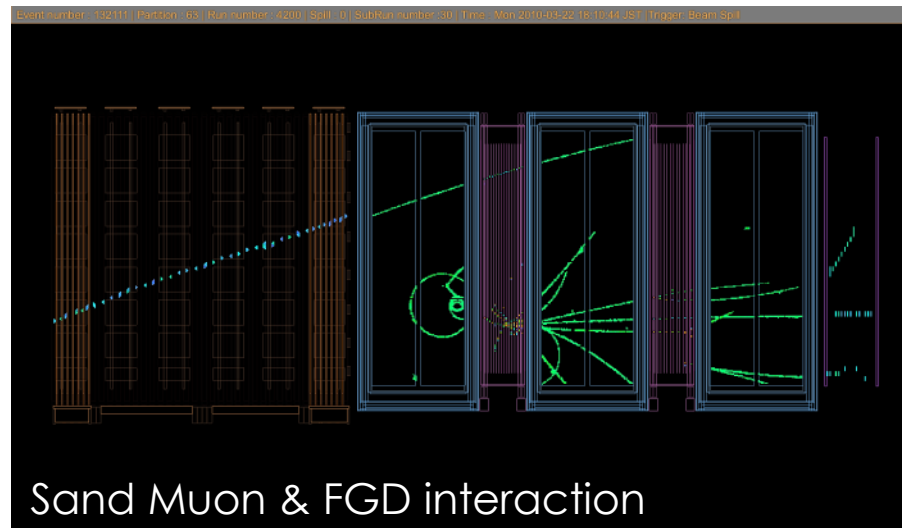
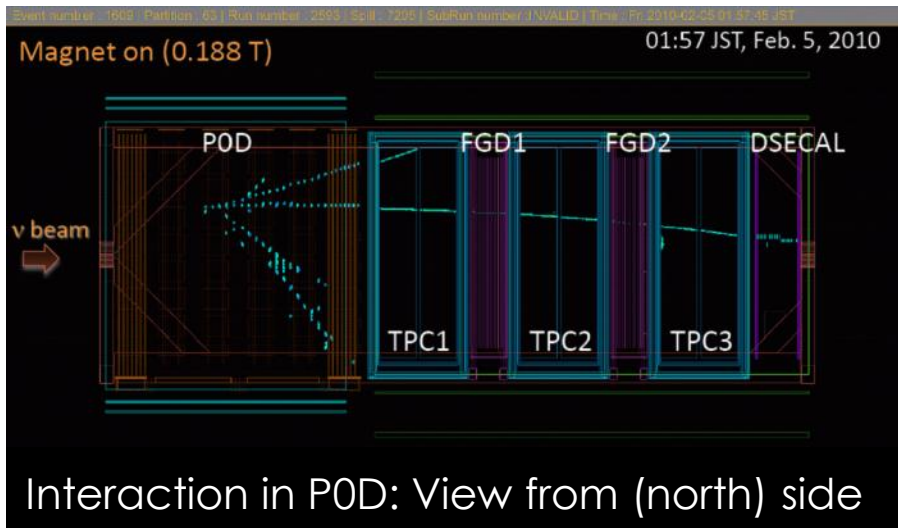
Observed events in FD data:

**62**

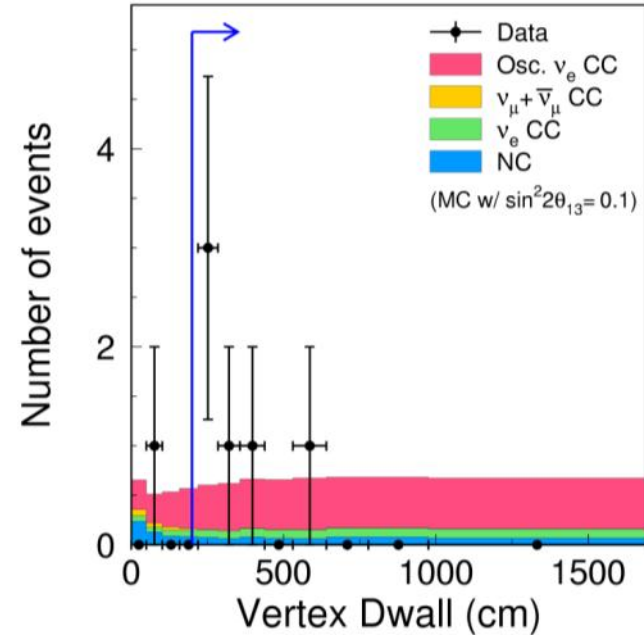
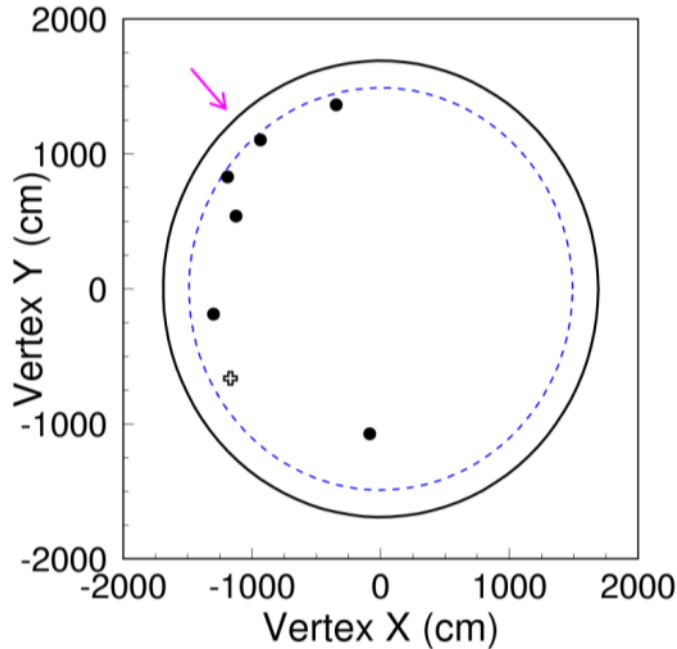
**$1.7\sigma$**  excess above background

# Example ND280 events

[ ND280 Data ]



# T2K Event distributions



Distribution is unusual:  $d_{\text{upstream-wall}}$  prob. is 0.14%,  
 $d_{\text{nearest-wall}}$  ('Dwall') prob. is 3.7%

But, addition of non-fiducial region improves the probabilities, opposite to a 'leak-in' hypothesis.

**T2K currently judges this to be a chance artefact.**

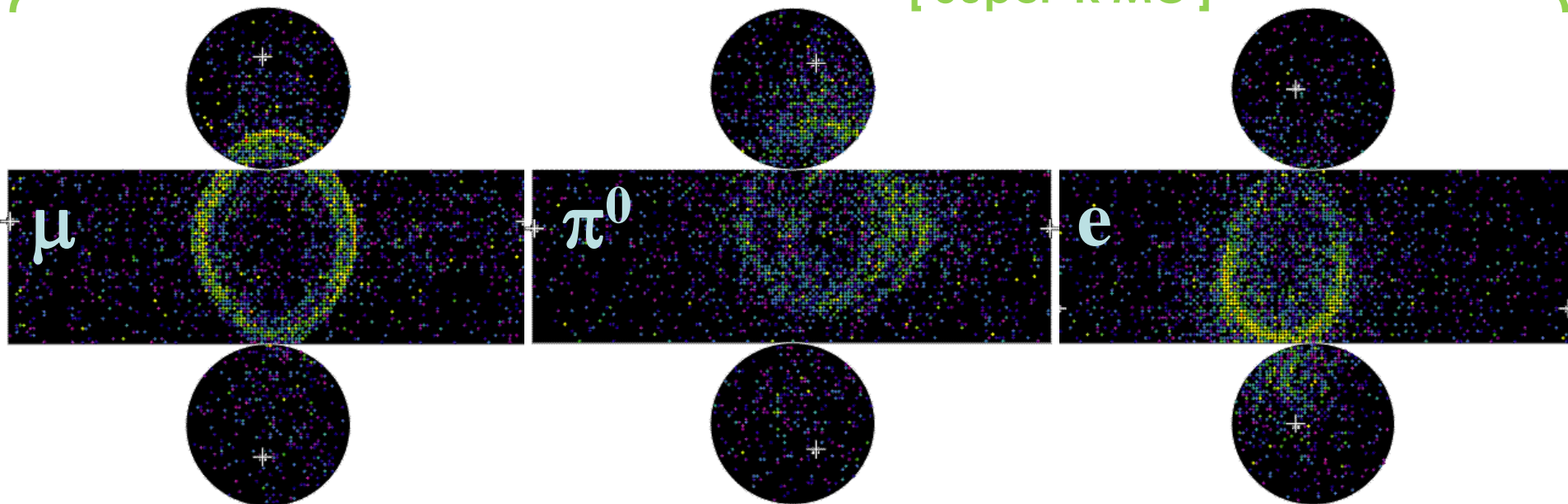
# T2K Pi-zero background

Isolated neutral pions from  $\nu_\mu$ -NC events:

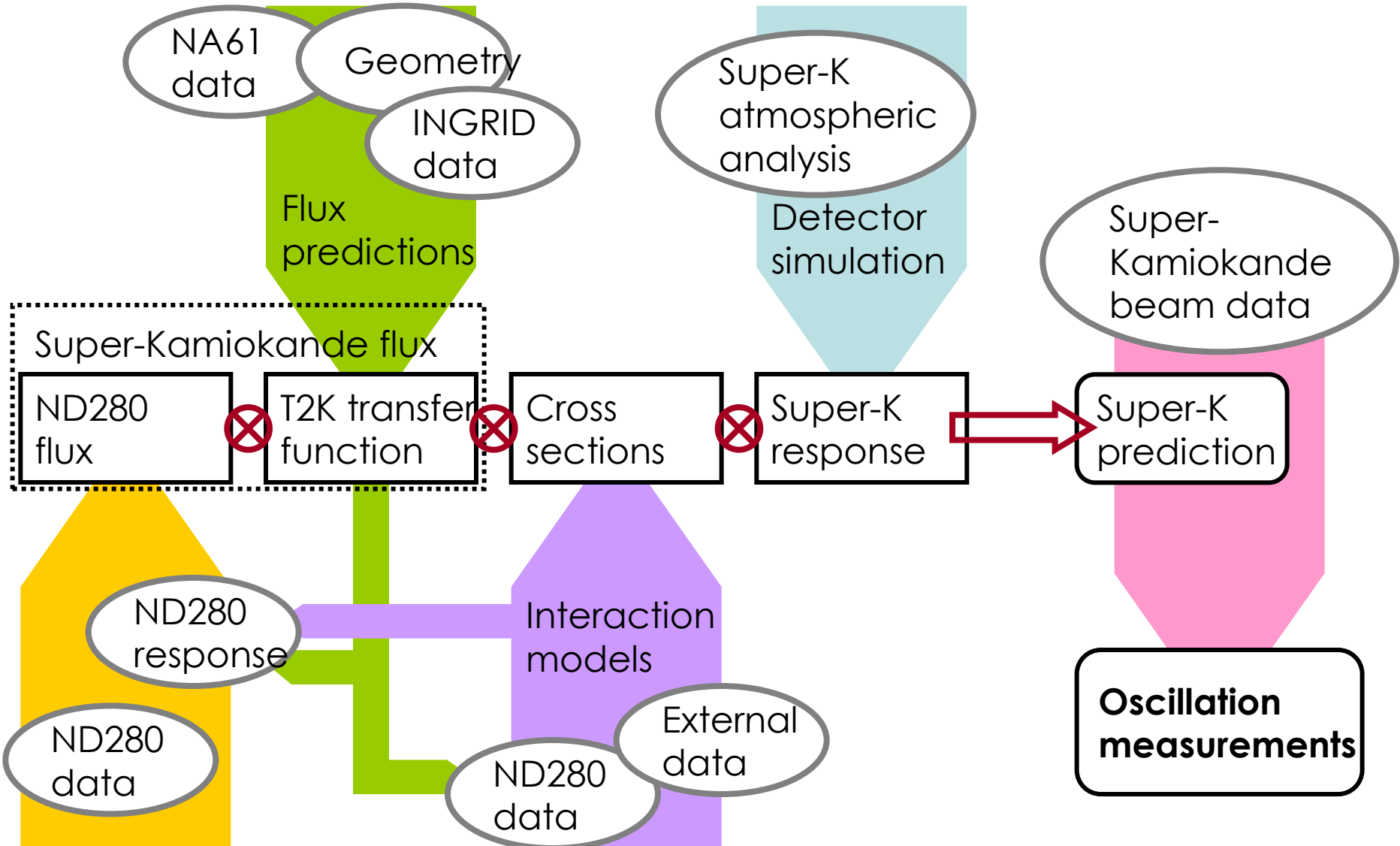
Neutral pion  $\rightarrow$  photon pair  $\rightarrow$  2 EM showers

- If the EM showers have same direction they mimic a single EM shower (electron signal)

[ Super-K MC ]



# T2K analysis strategy





# Averaging

LBL experiments hard to include in an average.

(how to interpret..?)

Reactor experiments are cleaner, and likelihood curves are closer to Gaussian.

A simple Gaussian combination (by me!) looks like this:

