

Measuring DPI in ATLAS in Wjj

Ellie Dobson

MPI@TAU

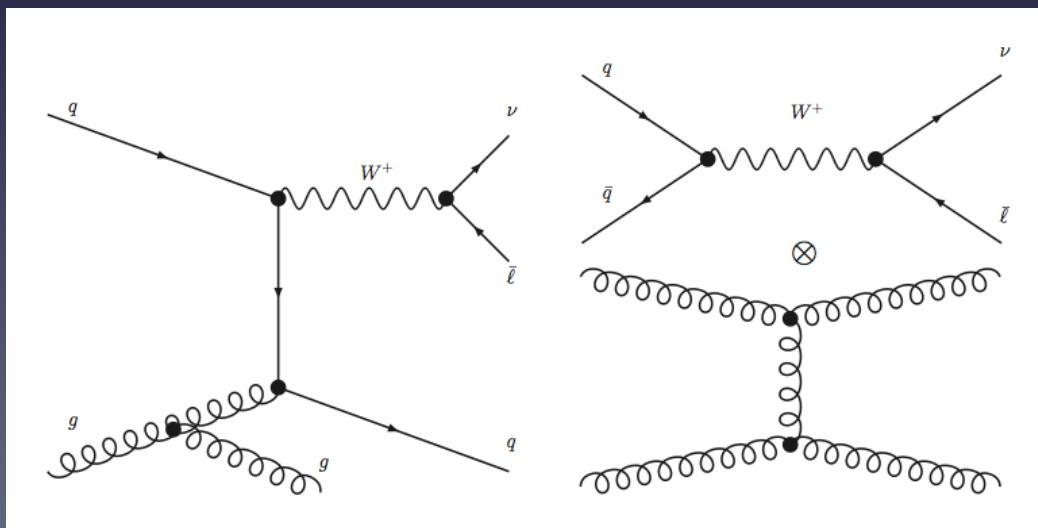
16/08/12

Motivation and method

Motivation: To quantify the probability of **hard** secondary scatter

- ? Hard DPI (double parton interactions) forms an irreducible BG to new physics searches and is not a well understood process
- ? Is DPI rate process independent?
- ? (How) does DPI rate depend on the collision energy?

Method: Exploit kinematic difference in DPI events to measure fraction of W+DPI contamination in W+2jet events, and use to extract σ_{eff}



These two will both pass W +2jet selection
- but in what fraction?

Samples and selection

Sample	Details
Pythia inclusive	v6, AMBT tune 1
Sherpa inclusive	v1.3.1, default UE, CKKW matching scale=30GeV
Alpgen+Herwig+Jimmy inclusive	MLM matching, Jimmy v4.31, AUET tune, Herwig v6.510
Sherpa MPI off	As above + MI_HANDLER=NONE
Alpgen+Herwig+Jimmy MPI off	As above + remove events where both jets' closest outgoing parton with $P_T > 3.5$ GeV is not primary
Data (W sample)	All 2010 data run
Data (jet sample)	All 2010 data run

W selection

Single lepton trigger

1 lepton (e, μ) $P_T > 20$ GeV, $\eta < 2.5$

MET > 25 GeV, $M_T > 40$ GeV

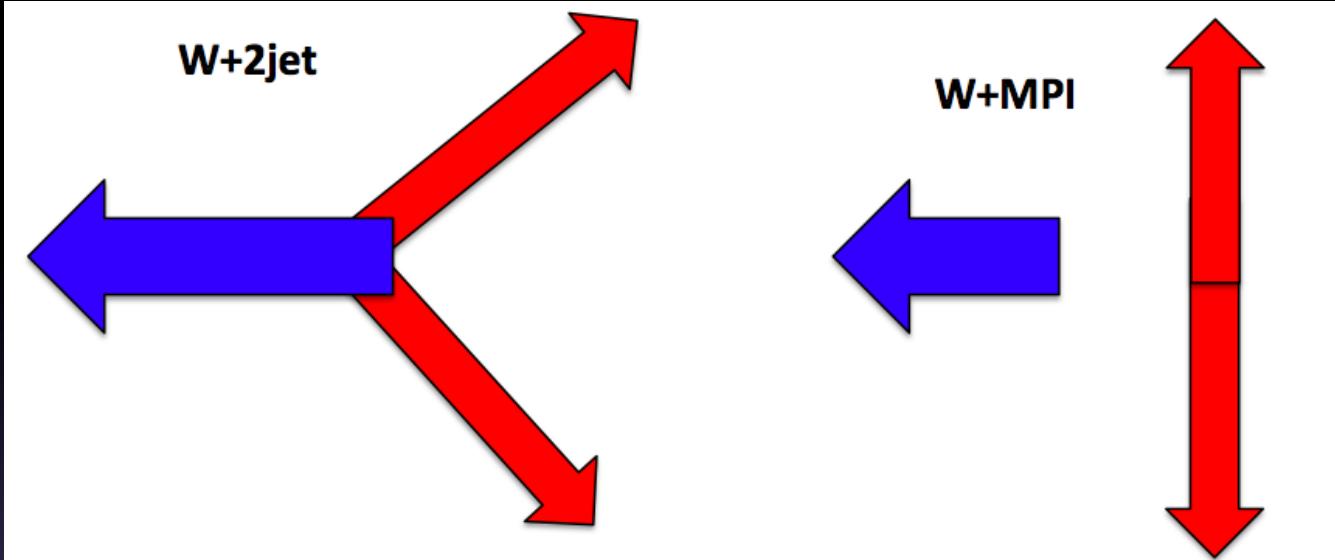
2 kt6 jets, $P_T > 20$ GeV, $y < 2.8$

Jet selection

Minimum bias trigger

Exactly 2 kt6 jets, $P_T > 20$ GeV, $y < 2.8$

Wjj topology I



W+MPI events differ from the W+2jet events in several ways:

- Total energy deposited in the event
- Transverse W momenta
- Recoil of jet system from W system
- Transverse jet momenta
- Angle between jets
- **Collinearity of jets**



Experimentally problematic for a DPI analysis

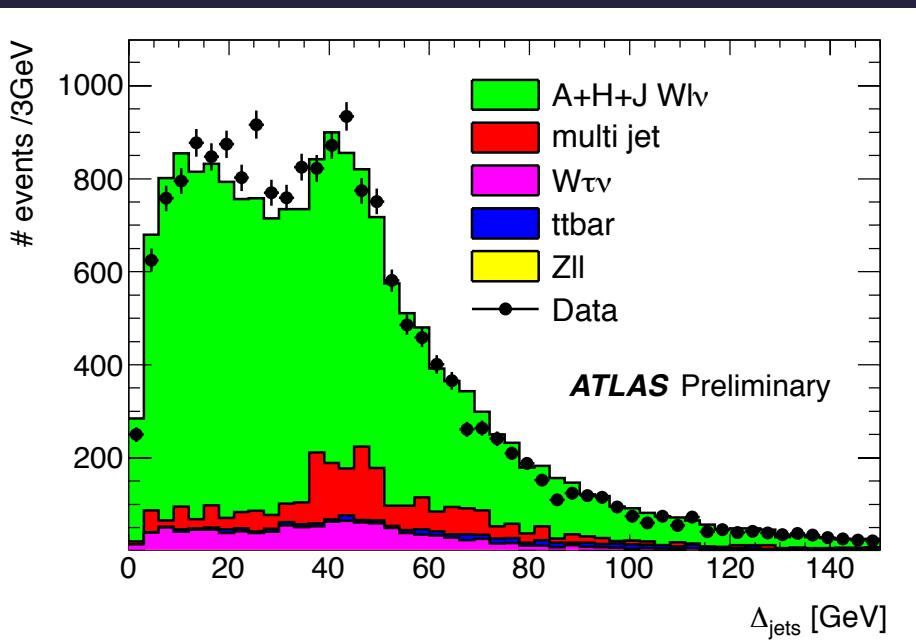
Wjj topology II

$$\Delta_{jets} = \left| \vec{P}_T^{J1} + \vec{P}_T^{J2} \right|$$

Component of jets back to back



Component of jets recoiling from W



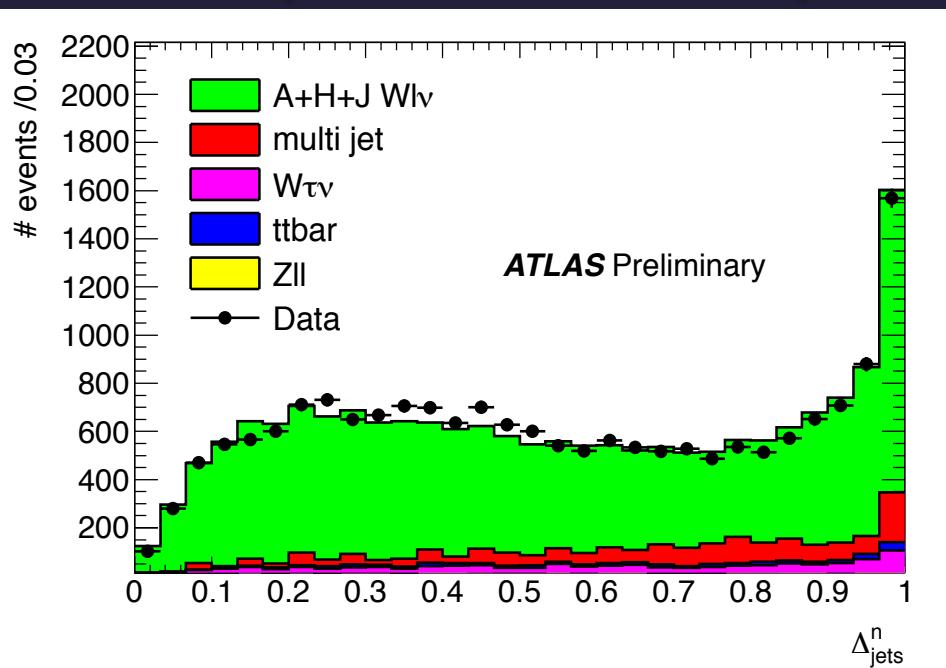
Wjj topology III

$$\Delta_{jets}^n = \frac{|\vec{P}_T^{J1} + \vec{P}_T^{J2}|}{|\vec{P}_T^{J1}| + |\vec{P}_T^{J2}|}$$

Component of jets back to back



Component of jets recoiling from W



Extracting DPI rate f_{DP}^R

$$f_{DP}^R = \frac{N_{W_0+2j_{MPI}}}{N_{W+2j}}$$



Numbers of events seen in the detector **passing selection cuts**

Overall distribution = $(1-f_{DP}^R) \bullet \text{Template A} + f_{DP}^R \bullet \text{Template B}$

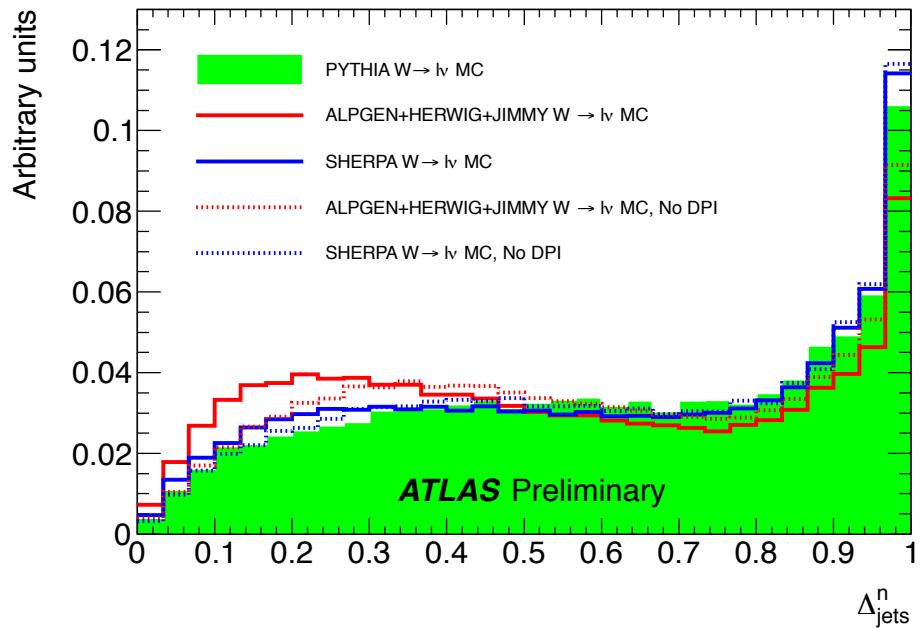
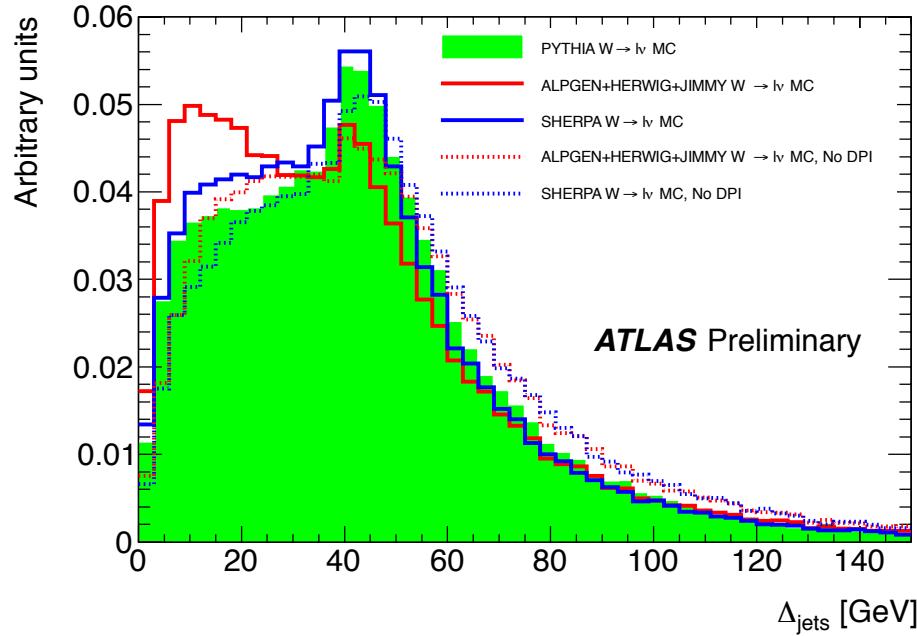


W+2jet (direct)



W+2jet (MPI)

Template A: W+2jets



Sherpa mismodelling now understood to be due to placement of CKKW matching cut at 30GeV

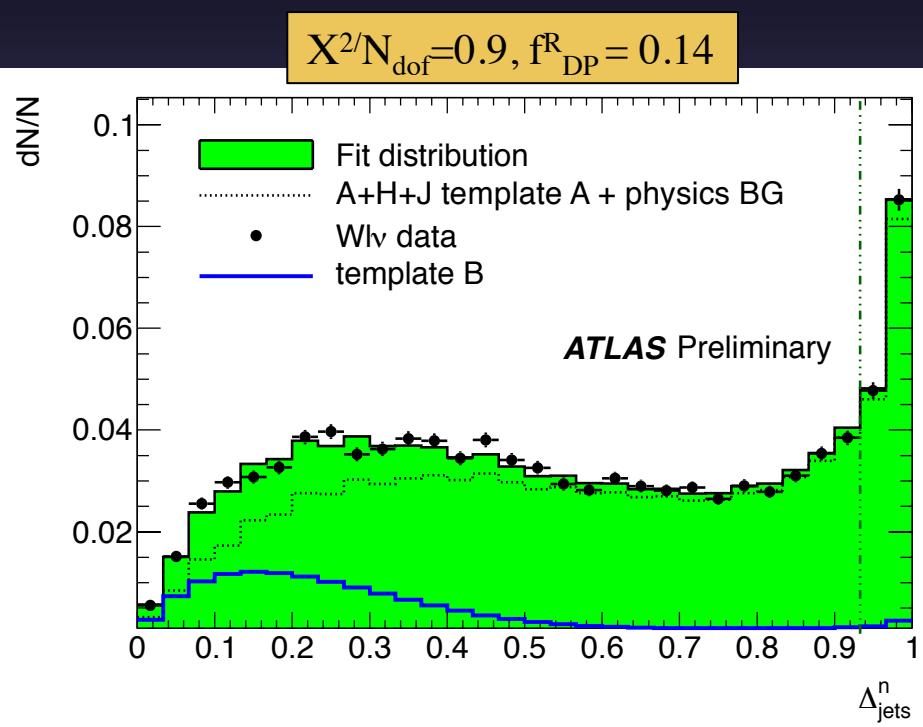
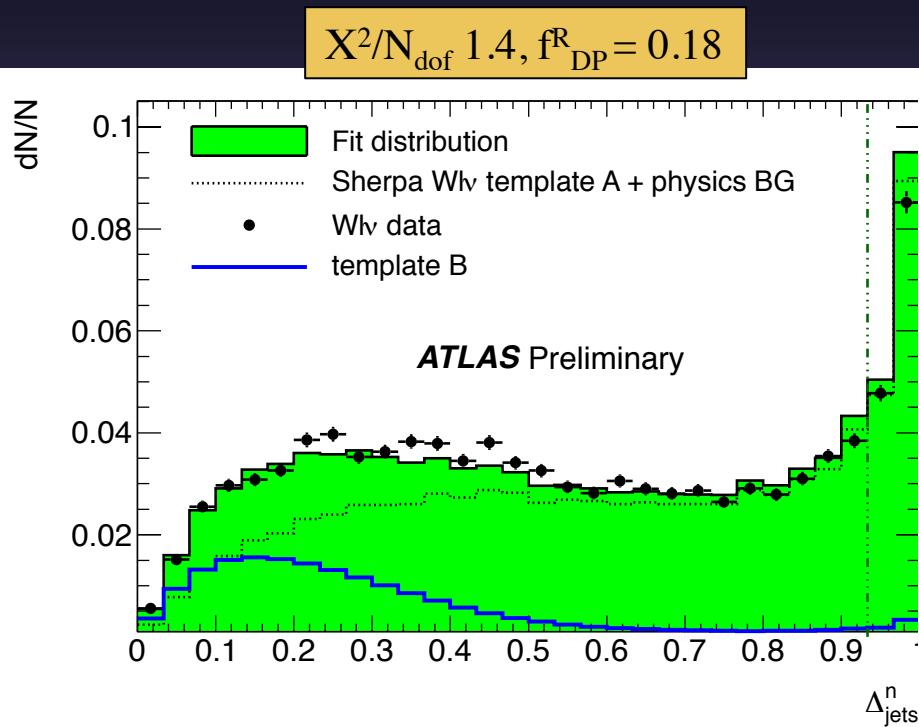
Template B: W+DPI

Dijet selection in data

Extraction of f^R_{DP}

Overall distribution = $(1-f^R_{DP}) \cdot \text{Template A} + f^R_{DP} \cdot \text{Template B}$

↑
 χ^2 minimisation

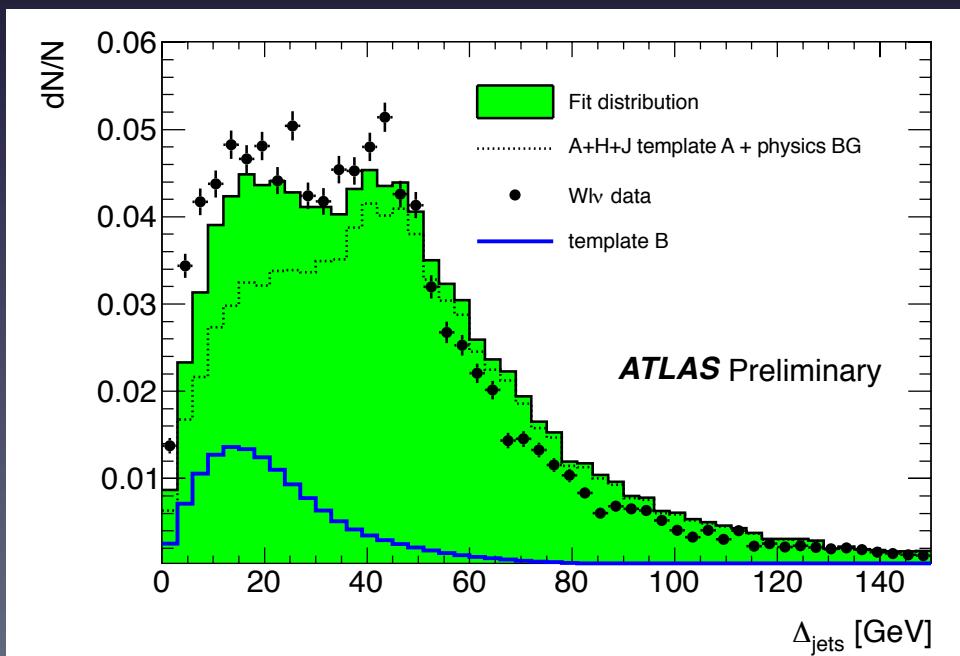


Comparison to Δ_{jets}

From previous fit

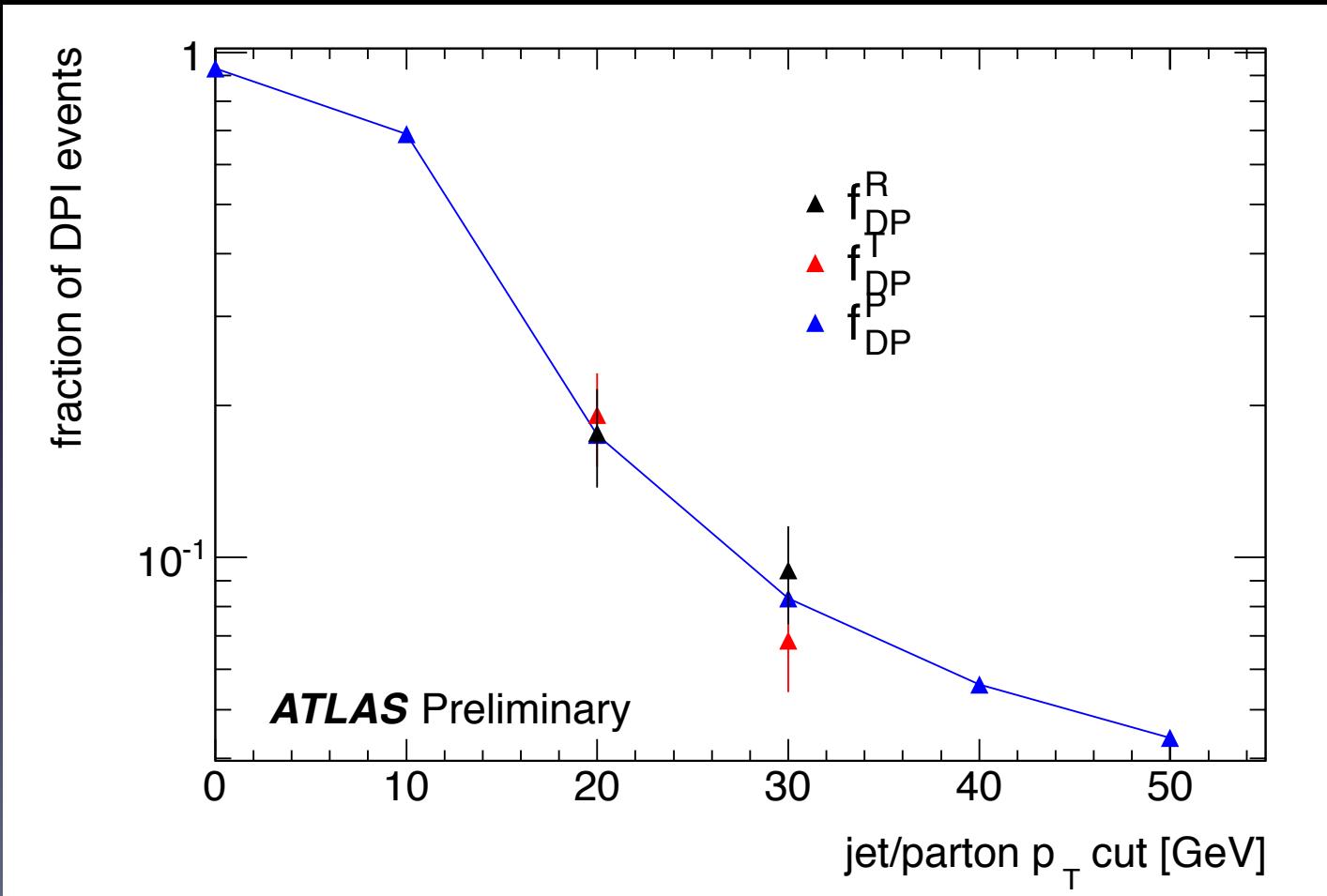


$$\text{Overall distribution} = (1-f^R_{DP}) \cdot \text{Template A} + f^R_{DP} \cdot \text{Template B}$$



Variation of f_{DP}^R with phase space

Both predicted and extracted DPI rate decrease as P_T cut is raised



f_{DP}^{R} results

Source of uncertainty	Method of evaluation	Fractional uncertainty / %
Generator modelling	ALPGEN+HERWIG+JIMMY vs SHERPA	12
Transition to parton level	Monte Carlo studies	10
Jet reconstruction	Jet energy scale shift	10
Pileup	Varying vertex number requirement	8
Trigger bias	Comparison of data streams	5
Background modelling	Varying multi jet background normalisation	1
Total systematic	Quadratic sum of the above	21
Total statistical	$\chi^2 + 1$	7

Table 2: Summary of the uncertainties on the extraction of f_{DP}^{R} .

$$f_{\text{DP}}^{\text{R}} = 0.16 \pm 0.01 \text{ (stat.)} \pm 0.03 \text{ (sys.)}.$$

Introducing σ_{eff}

$$\sigma_{\text{eff}} := \frac{\sigma_A \cdot \sigma_B}{\sigma_{AB}}$$

Can think of σ_{eff} as the effective area of the proton
-> Larger σ_{eff} : smaller DPI rate

Collab	E_{COM}	PT cuts /GeV	$\sigma_{\text{eff}}/\text{mb}$
AFS	63GeV	4	~ 5
UA2	630GeV	15	> 8.3
CDF (jjjj)	1.8TeV	25	$12.1^{+10.5}_{-5.4}$
CDF (γjjj)	1.8TeV	15 on γ , 5-7 on j	$14.5^{+1.7}_{-2.3}$
Do (γjjj)	1.96TeV	60->80 on γ	$16.4^{+0.3}_{-2.3}$
ATLAS (Wjj) ?	7TeV	20 on $ v $, 20 on j	??

σ_{eff} is postulated to be process independent – its measurement in W+2jets allows calculation of DPI background to any physics analysis

Converting to σ_{eff}

Taking input definitions

$$f_{\text{DP}}^{\text{R}} = \frac{N_{W_0+2j_{\text{DPI}}}}{N_{W+2j}}, \quad \sigma_{\text{eff}} = \frac{\sigma_{W_0} \cdot \sigma_{2j}}{\sigma_{W_0+2j_{\text{DPI}}}},$$

writing i.t.o cross sections

$$\sigma_{\text{eff}} = \frac{1}{f_{\text{DP}}^{\text{R}}} \cdot \frac{N_{W_0} N_{2j}}{N_{W+2j}} \cdot \frac{A_{W_0+2j_{\text{DPI}}}}{A_{W_0} A_{2j}} \cdot \frac{\epsilon_{W_0+2j_{\text{DPI}}}}{\epsilon_{W_0} \epsilon_{2j}} \cdot \frac{\mathcal{L}_{W_0+2j_{\text{DPI}}}}{\mathcal{L}_{W_0} \mathcal{L}_{2j}}.$$

and using input assumptions of analysis*

$$A_{W_0+2j_{\text{DPI}}} = A_{W_0} \cdot A_{2j_{\text{DPI}}},$$

$$A_{2j_{\text{DPI}}} = A_{2j_D} \cdot$$

Yields**

$$\sigma_{\text{eff}} = \frac{1}{f_{\text{DP}}^{\text{R}}} \cdot \frac{N_{W_0} N_{2j_D}}{N_{W+2j}} \cdot \frac{1}{\epsilon_{2j_D}} \cdot \frac{1}{\mathcal{L}_{2j_D}}.$$

* need small correction for overlap removal

** include additional ₁₄ systematic for trigger bias

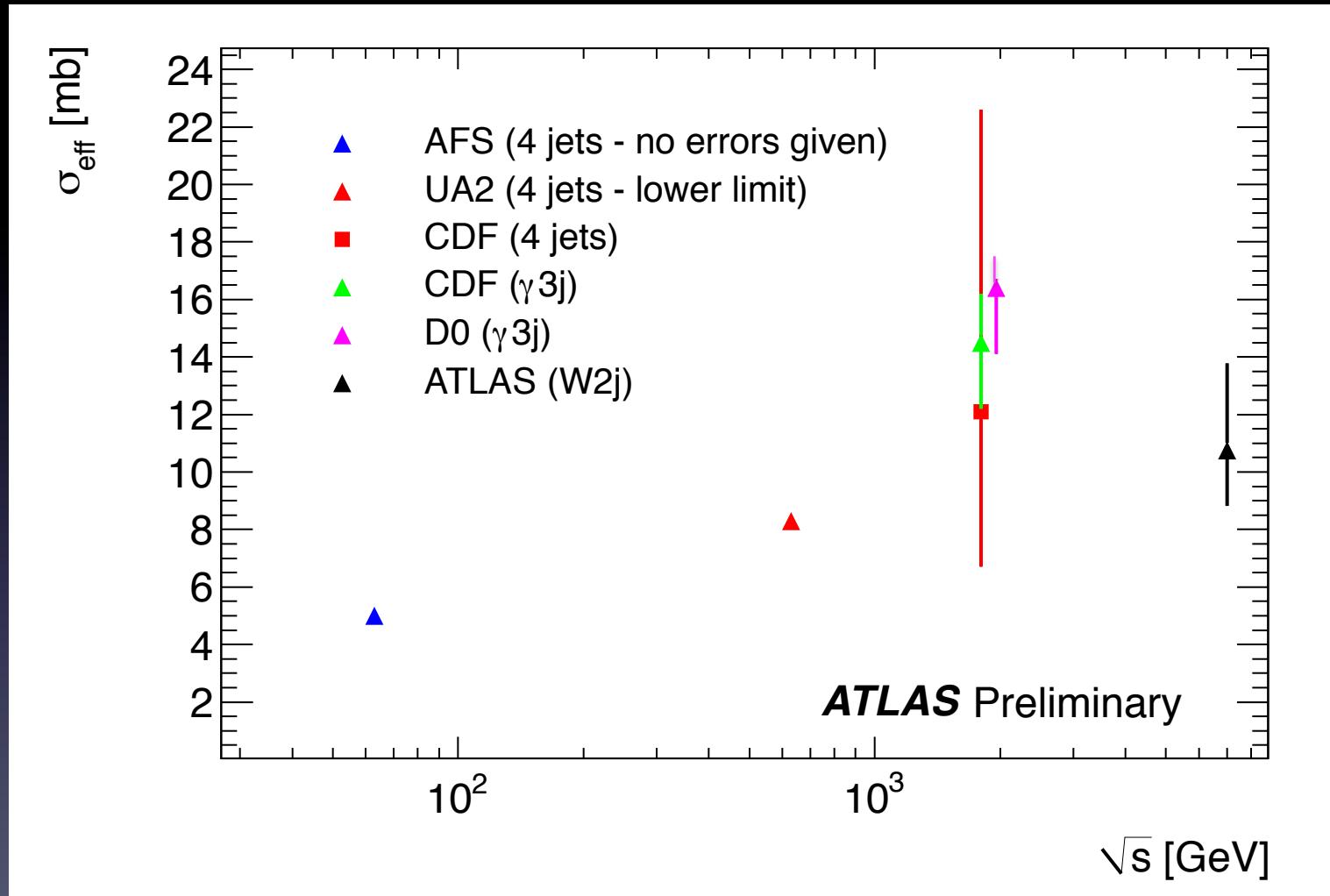
σ_{eff} results

Quantity	Systematic source	Method of evaluation	Fractional uncertainty /%
$N_{W0}/N_{W2} \cdot N_{jj}$	Acceptance cancellation	Section 6.1	< 3
N_{W0}/N_{W2}	Background modelling	Reference [53]	5
\mathcal{L}_{jj}	Luminosity	Beam parameters [52]	3.4
f_{DP}	Total	As in Table 2	21

Table 3: Summary of the systematic uncertainties on σ_{eff} .

$$\sigma_{\text{eff}}(7 \text{ TeV}) = 11 \pm 1 \text{ (stat.)} {}^{+3}_{-2} \text{ (sys.) mb.}$$

Putting the result into context....



Results consistent with other measurements

Conclusions

The relative DPI rate is extracted for W+2jet events in the ATLAS detector:

$$f_{\text{DP}}^{\text{R}} = 0.16 \pm 0.01 \text{ (stat.)} \pm 0.03 \text{ (sys.)}.$$

From this, the effective cross section is measured in 7 TeV pp collisions

$$\sigma_{\text{eff}}(7 \text{ TeV}) = 11 \pm 1 \text{ (stat.)} {}^{+3}_{-2} \text{ (sys.) mb.}$$

which is consistent with results obtained in different channels at the Tevatron.

Coming up in the paper...

- New Sherpa samples with lower matching cut (increased σ_{eff})
- Better understanding of how f_{DP} translates from parton to reconstruction level
- Truth-level distributions

backup

f_{DP}^R =parton level f_{DP}^P ?

Fit

ALPGEN+HERWIG+JIMMY=

$$(1-f_{DP}^T) A \text{ (Sherpa)} + f_{DP}^T B$$

- vary f_{DP}^P in ALPGEN+HERWIG+JIMMY

