



Serpentine acceleration in zero-chromatic FFAG

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Outline

- Motivations and introduction
- Serpentine acceleration scheme in scaling FFAG
- Experiments for demonstration of serpentine acceleration with electron scaling FFAG
- Application of using a serpentine acceleration scheme
- Summary and Conclusion

Requirements for accelerators

- Accelerators are used in many fields these days...for instance,
- Elementary Physics



Muon accelerator & muon collider are good candidates for the next generation of elementary physics experiments.



Rapid-acceleration scheme is required to accelerate short-lived particles within lifetimes.

- Accelerator Driven System (ADS)

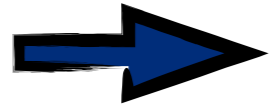


ADS : A system which keeps nuclear fission chain reaction induced by spallation neutrons obtained by irradiation of a target with high energy proton beams generated by **accelerators**.

High power proton beam is needed to generate enough neutrons

Necessary conditions for accelerators

Requirements for accelerators



- Achievable of **Rapid acceleration** ($\sim 2\mu\text{sec}$)
- Generating of **High power beam** ($\sim 10\text{ MW}$)

Accelerators are necessary satisfied with the following conditions:

Requirements	Necessary conditions	Reason
Rapid acceleration	Static magnetic field	Variable magnetic field \rightarrow Short-lived particles cannot be accelerated within the lifetime.
High-intensity beam	Fixed rf frequency acceleration	Continuous wave (CW) acceleration can be achieved
High-energy beam	Strong focusing	Large dynamic aperture is secured for secondary particle acceleration scheme.

Candidate of accelerators

● Candidate so far

	Problems
Linac	Lots of cavities are needed
	Large space to put many rf cavities is required



Expensive

● Circular accelerators

	Strong focusing	Static magnetic field	Fixed rf frequency acceleration
Cyclotron	NO	YES	NO
Synchrotron	YES	NO	MAYBE ($\beta \sim 1$: Possible)

● New candidate of accelerators

Fixed Field Alternating Gradient (FFAG) Accelerators

FFAG accelerators

Fixed Field Alternating Gradient (FFAG) Accelerators

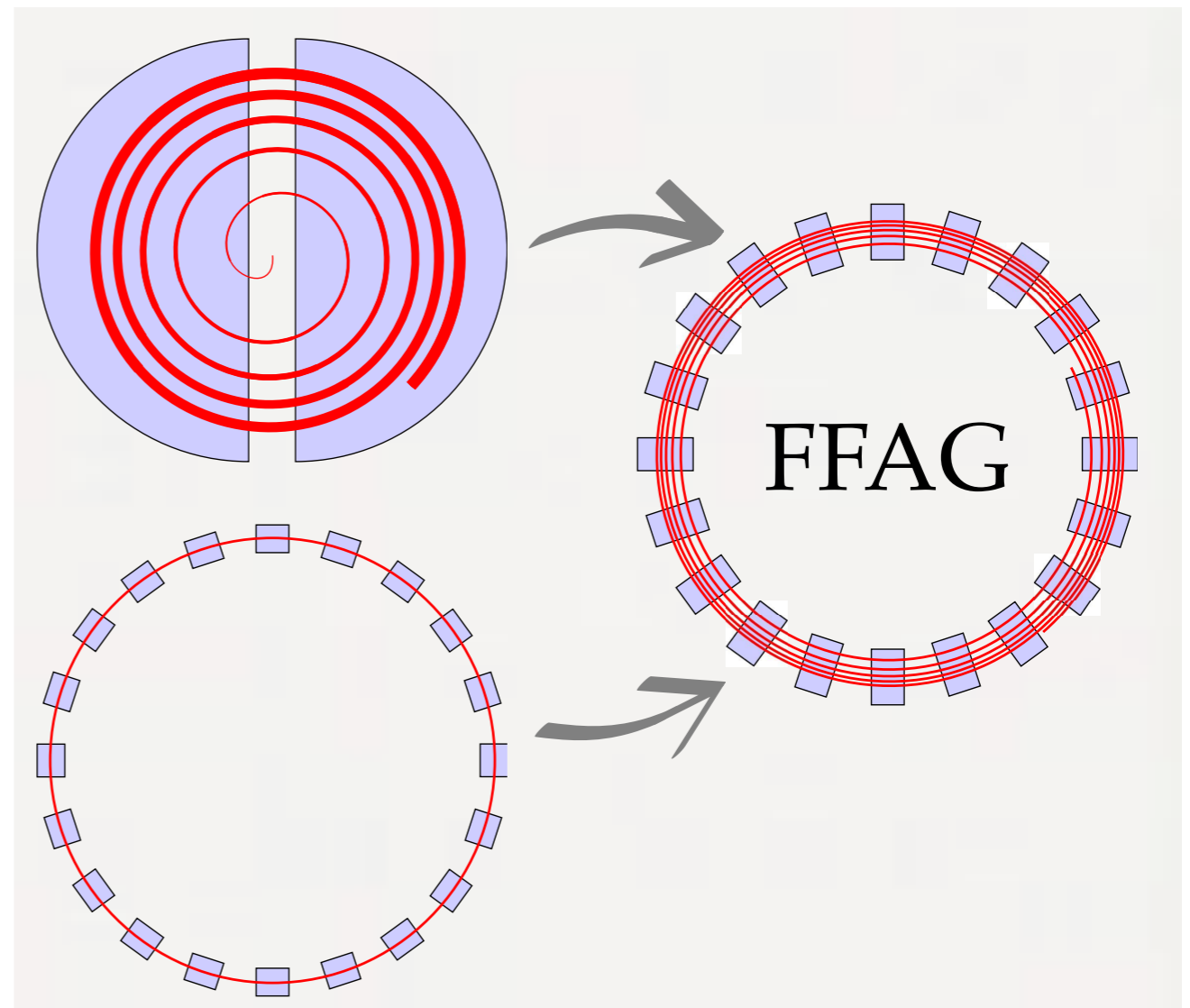
**Fixed field in time
(Like cyclotron)**

+

**Strong focusing
(Like synchrotron)**

+

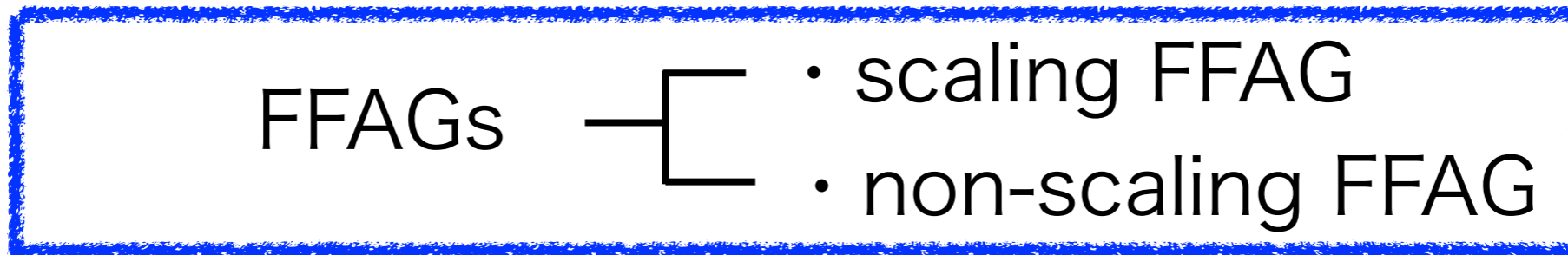
Fixed rf frequency acceleration



FFAG satisfies three necessary conditions to realize high-intensity and rapid-acceleration scheme.

Features of FFAGs

Category of FFAGs



		Features
scaling FFAG	non-linear magnetic field ($B \propto r^k$)	Constant betatron tune
non-scaling FFAG	Linear magnetic field	Betatron tune is not constant

Fixed rf frequency acceleration scheme in FFAGs

Accelerators	Acceleration scheme	Energy
scaling FFAG	Stationary bucket acceleration	Relativistic energy ($\beta \sim 1$) beam is suitable.
non-scaling FFAG	Serpentine acceleration	

Purpose of the study

● Beam acceleration in non-relativistic energy region

	Energy	Requirements
Low energy muon acceleration	non-relativistic energy region	Rapid acceleration
Proton driver for ADS		High-intensity proton beam



Fixed rf frequency acceleration scheme is required
Non-relativistic ($\beta < 1$) energy region



Serpentine acceleration is applied in **scaling FFAG**

The purpose of this study is to examine a **serpentine acceleration** scheme in **scaling FFAG** both **theoretically and practically**, allowing fixed rf frequency acceleration in **non-relativistic energy** region.

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 - Longitudinal hamiltonian for fixed rf frequency in scaling FFAG
 - Longitudinal phase space
 - Minimum rf voltage to make a serpentine channel
 - Total energy gain and phase acceptance in serpentine acceleration
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Longitudinal hamiltonian for fixed rf frequency in scaling FFAG

● Phase equation

Phase difference ($\Delta\Phi$) per turn at the cavity.

➔ $\Delta\phi = 2\pi(f_{rf} \cdot T - h)$ T : revolution period of no-synchronous particle
 h : harmonic number f_{rf} : rf frequency

The relation between the equivalent radius (R) and momentum (P) is exactly given by

➔
$$R = R_0 \left(\frac{P}{P_0} \right)^{\frac{1}{k+1}}$$

➔
$$\begin{aligned} \frac{T}{T_s} &= 1 + \frac{\Delta\phi}{2\pi h} = \left(\frac{R}{R_s} \right) / \frac{P/E}{P_s/E_s} \\ &= \left(\frac{P}{P_s} \right)^\alpha \frac{P_s}{E_s} \left(\frac{E^2}{E^2 - m^2} \right)^{\frac{1}{2}} \\ &= \left(\frac{E^2 - m^2}{P_s^2} \right)^{\frac{\alpha}{2}} \frac{P_s}{E_s} \left(\frac{E^2}{E^2 - m^2} \right)^{\frac{1}{2}} \\ &= \left(E_s^2 - m^2 \right)^{\frac{1-\alpha}{2}} \frac{E}{E_s} \left(E^2 - m^2 \right)^{\frac{\alpha-1}{2}} \end{aligned}$$

α : momentum compaction factor

$$\alpha = \frac{1}{k+1}$$

➔ Solve the equation for $\Delta\Phi$

Longitudinal hamiltonian for fixed rf frequency in scaling FFAG

● Phase equation

$$\Delta\phi = 2\pi h \left[\frac{(E_s^2 - m^2)^{\frac{1-\alpha}{2}}}{E_s} E (E^2 - m^2)^{\frac{\alpha-1}{2}} - 1 \right]$$

● Energy equation

Energy gain (ΔE) per turn at the cavity

$$\Delta E = eV_{rf} \sin \phi$$

We assume that there are many rf cavity in a ring

➔ Phase and energy differences are very small.

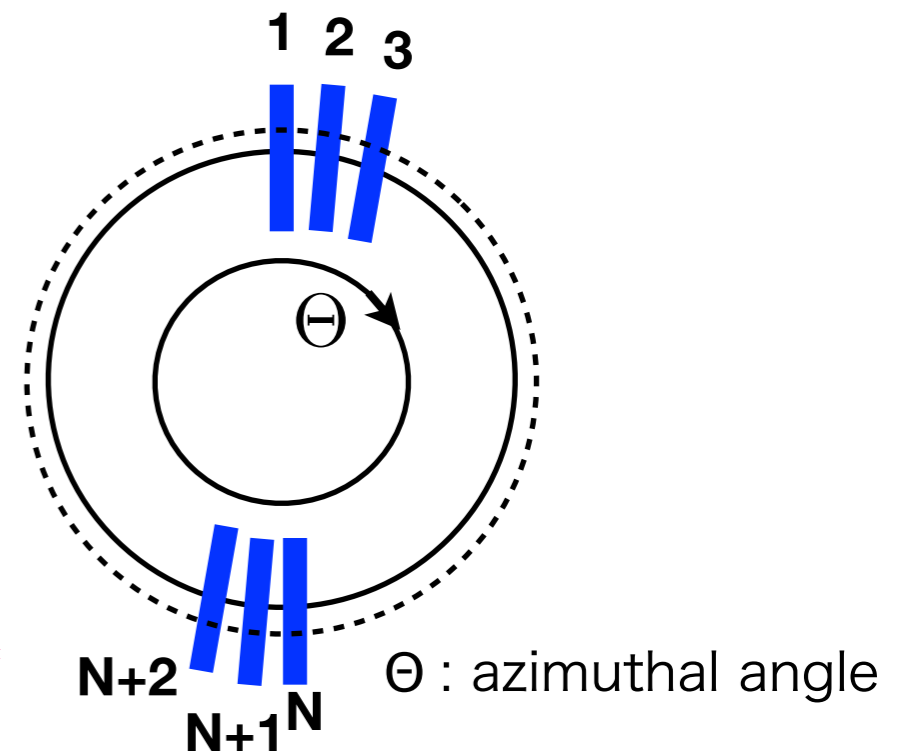
$$\frac{\Delta\phi}{2\pi} \leftrightarrow \frac{d\phi}{d\Theta} \quad \frac{\Delta E}{2\pi} \leftrightarrow \frac{dE}{d\Theta}$$

Equation of motions



$$\frac{d\phi}{d\Theta} = h \left[\frac{(E_s^2 - m^2)^{\frac{1-\alpha}{2}}}{E_s} E (E^2 - m^2)^{\frac{\alpha-1}{2}} - 1 \right]$$

$$\frac{dE}{d\Theta} = \frac{eV_{rf}}{2\pi} \sin \phi$$



Longitudinal hamiltonian for fixed rf frequency in scaling FFAG

Equation of motions



$$\frac{d\phi}{d\Theta} = h \left[\frac{(E_s^2 - m^2)^{\frac{1-\alpha}{2}}}{E_s} E (E^2 - m^2)^{\frac{\alpha-1}{2}} - 1 \right]$$

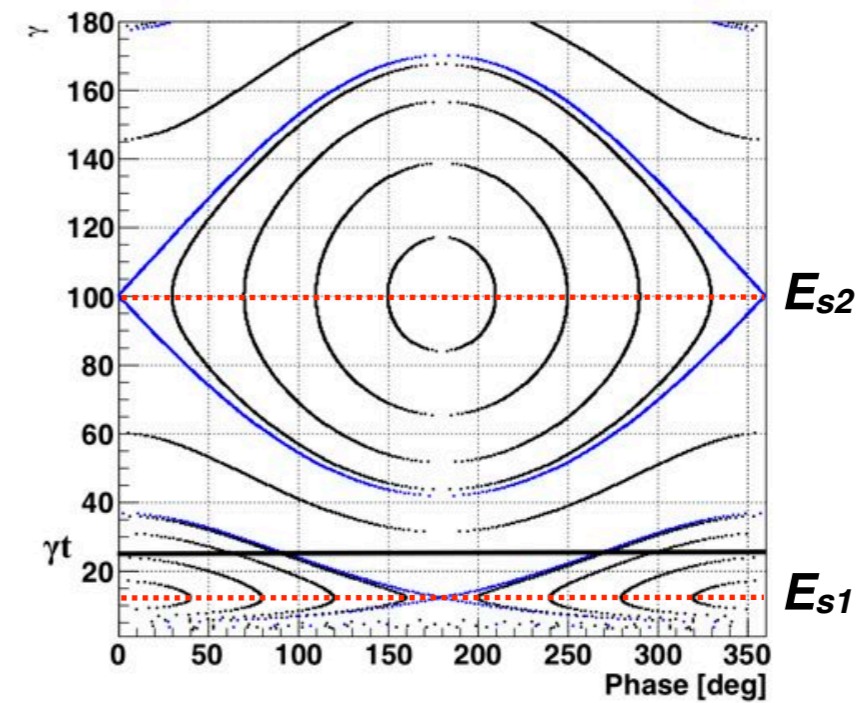
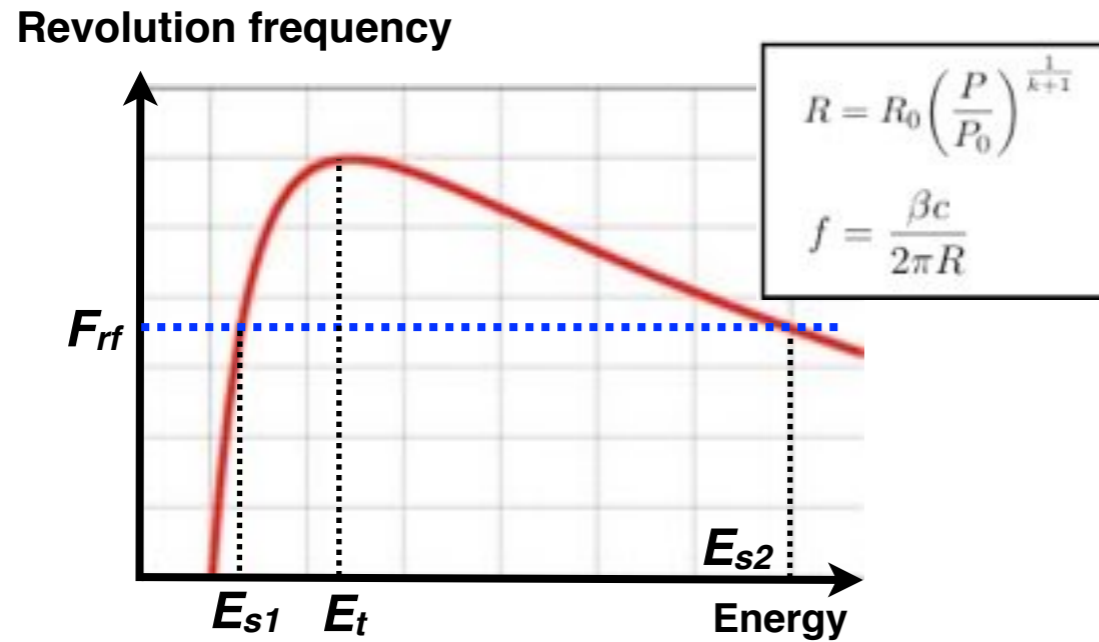
$$\frac{dE}{d\Theta} = \frac{eV_{rf}}{2\pi} \sin \phi$$

We choose the energy variable E canonically conjugate to the coordinate variable ϕ .

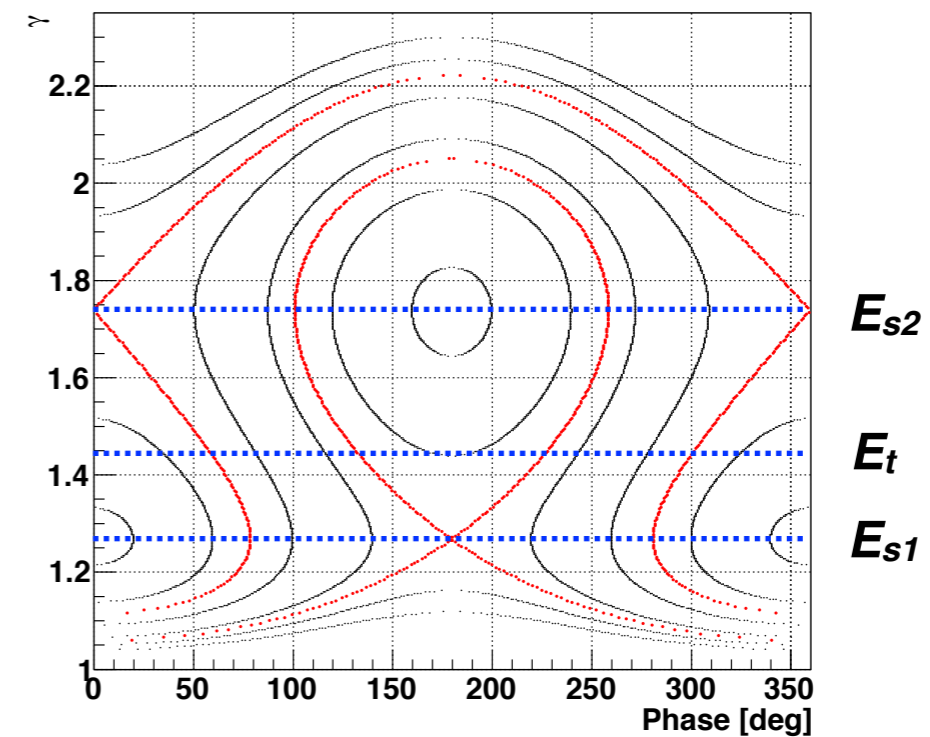
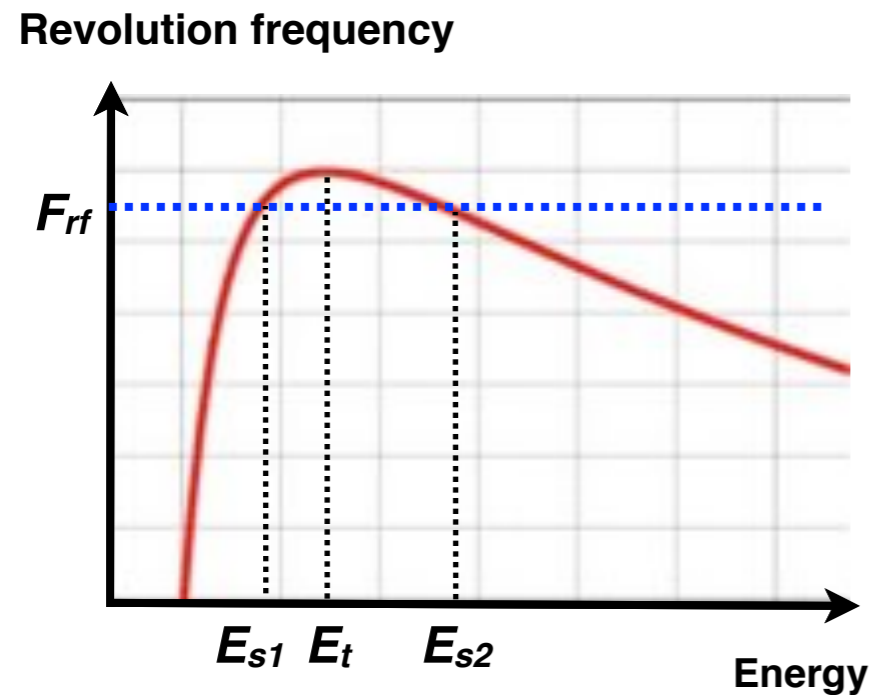
Longitudinal hamiltonian

$$H(E, \phi; \Theta) = h \left[\frac{1}{\alpha + 1} \frac{(E^2 - m^2)^{\frac{\alpha+1}{2}}}{E_s (E_s^2 - m^2)^{\frac{\alpha-1}{2}}} - E \right] + \frac{eV_0}{2\pi} \cos \phi.$$

Longitudinal phase space



Rf frequency is fixed near the transition energy



Serpentine acceleration can be achieved in non-relativistic energy region with appropriate selection of transition energy.

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Minimum rf voltage

- Limit condition of making serpentine channel

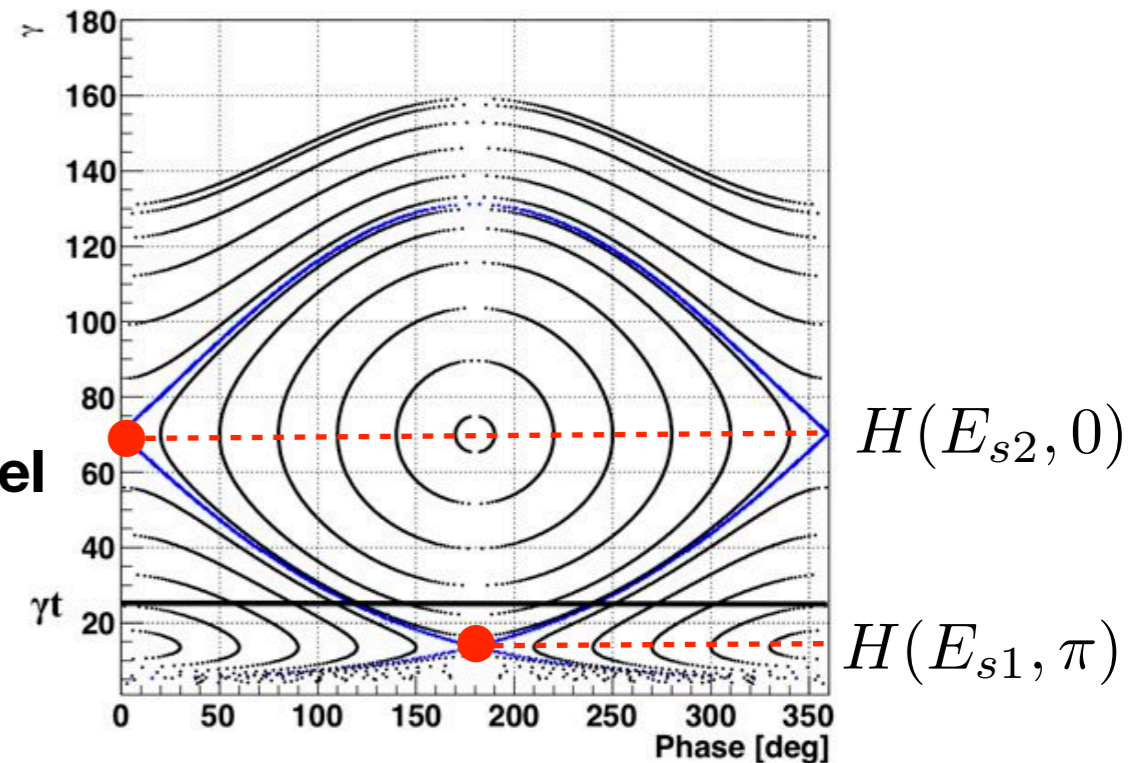
Both stationary buckets are just connected.

Separatrix satisfies the equation :

$$H(E_{s1}, \pi) = H(E_{s2}, 0)$$

Lowest limit of rf voltage to make a serpentine channel

$$V_{min} = \pi h \left[\frac{1}{\alpha + 1} \left(\frac{(E_{s1}^2 - m^2)}{E_{s1}} - \frac{(E_{s2}^2 - m^2)}{E_{s2}} \right) + (E_{s2} - E_{s1}) \right]$$

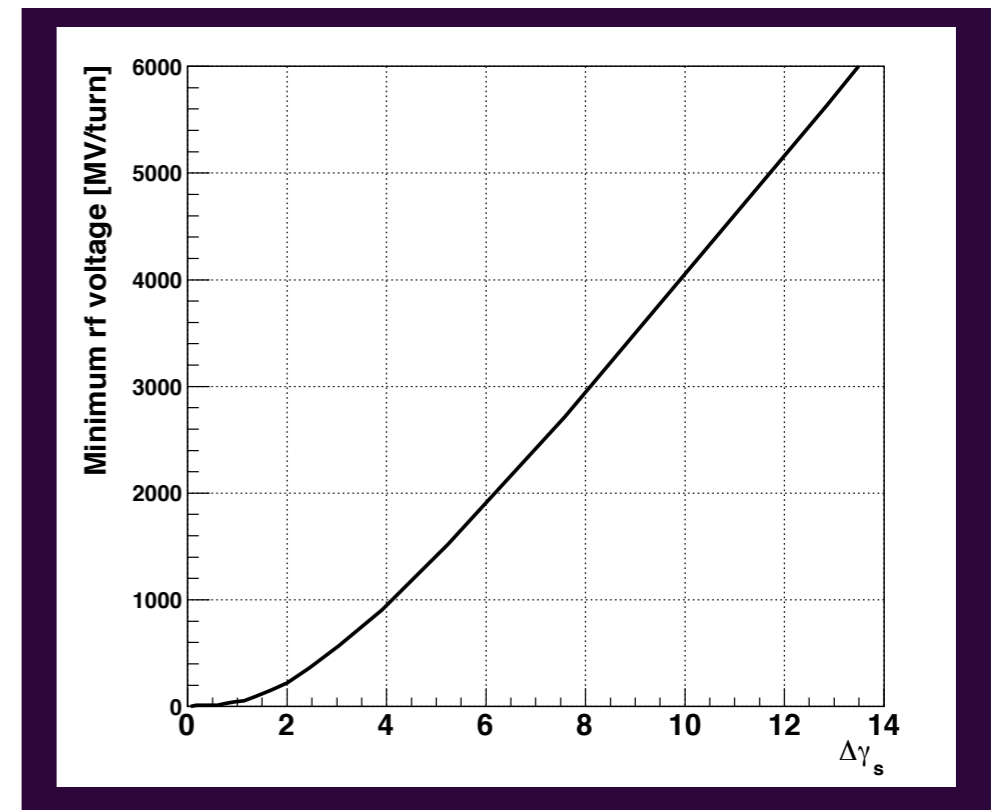


Relation between minimum rf voltage and energy difference of stationary energies $\Delta\gamma_s$

(k-value & h are fixed)



Minimum rf voltage becomes big with difference of stationary energies ($\Delta\gamma_s$).



Maximum energy gain of serpentine acceleration

Maximum energy gain

Maximum energy range is determined by minimum and maximum energy on serpentine channel.

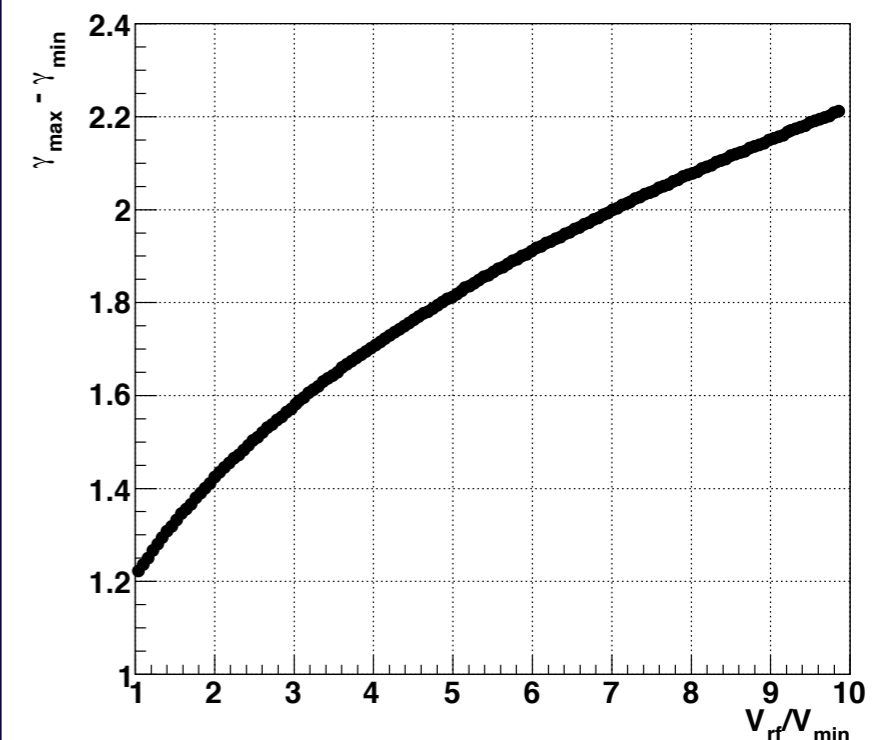
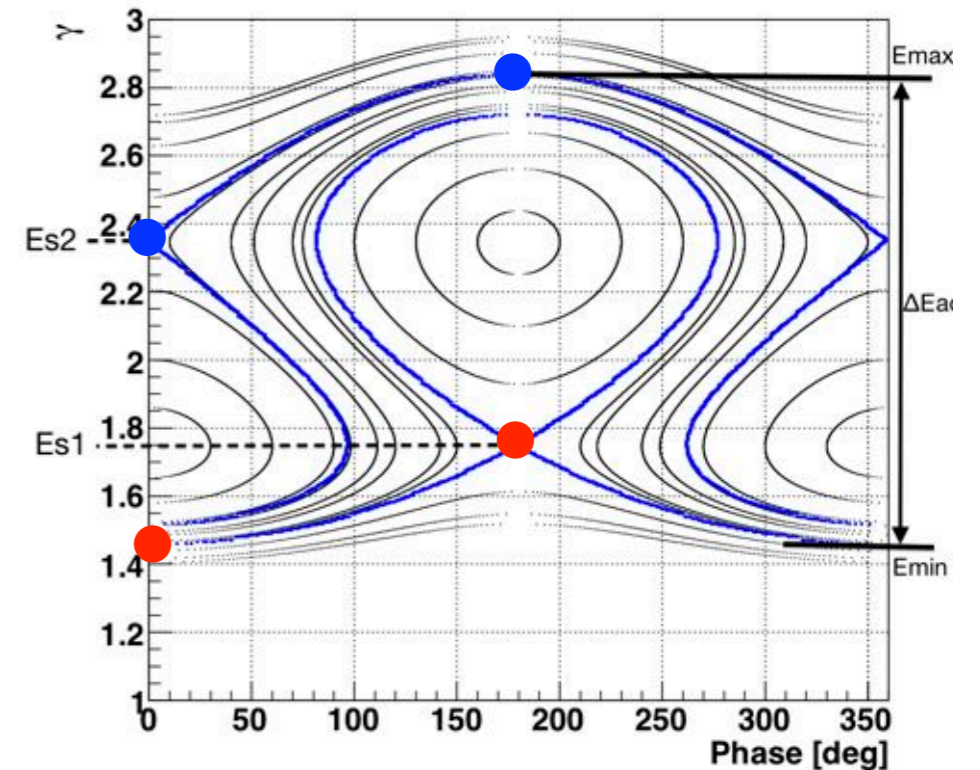
Separatrices for Minimum energy E_{min} and Maximum energy E_{max} satisfy these equations:

$$\begin{cases} H(E_{max}, \pi) = H(E_{s2}, 0), \\ H(E_{min}, 0) = H(E_{s1}, \pi). \end{cases}$$

Relation between maximum energy gain ΔE and rf voltage V_{rf}
(Stationary energies, harmonic number and k-value are fixed.)

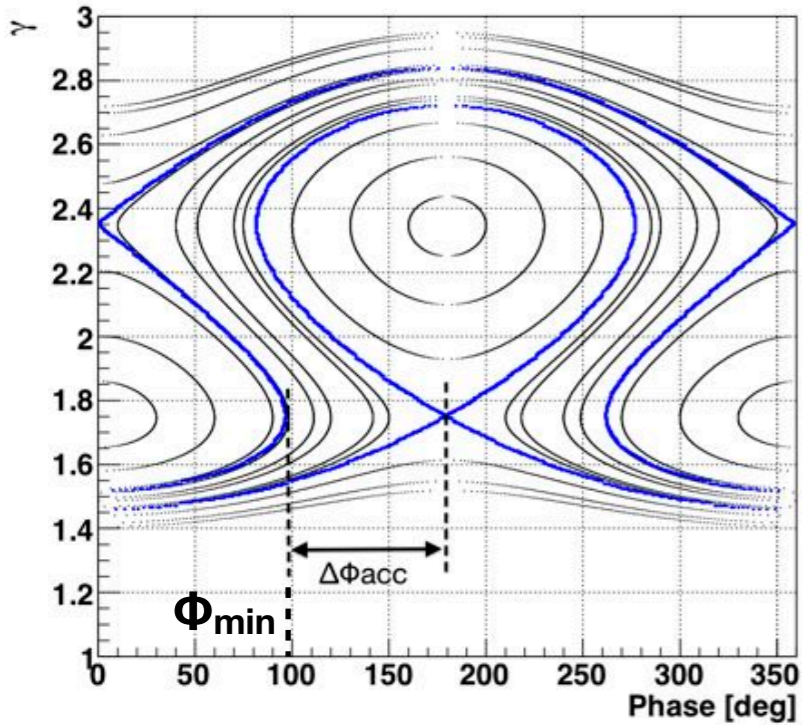


Maximum energy gain becomes large with rf voltage.



Phase acceptance

- Phase acceptance of serpentine channel at stationary energy below transition



Separatrix for minimum phase (Φ_{min}) of phase acceptance satisfies the equation:

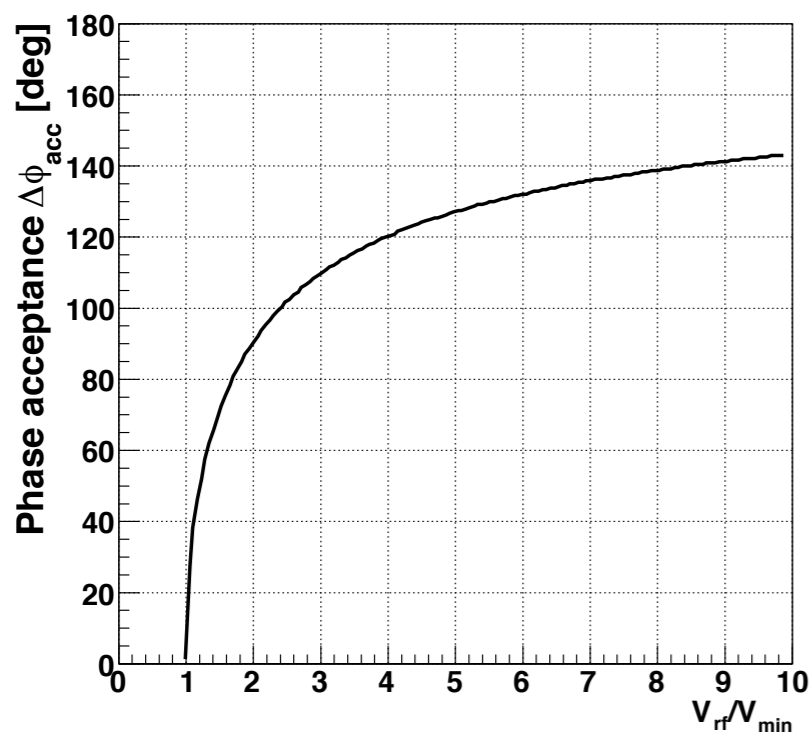
$$H(E_{s1}, \phi_{min}) = H(E_{s2}, 0)$$

Phase acceptance $\Delta\Phi_{acc}$

$$\begin{aligned} \Delta\phi_{acc} &= \pi - \phi_{min} \\ &= \pi - \arccos \left[\frac{2\pi h}{eV_{rf}} \left(\frac{1}{\alpha + 1} \left(\frac{(E_{s2}^2 - m^2)}{E_{s2}} - \frac{(E_{s1}^2 - m^2)}{E_{s1}} \right) + (E_{s1} - E_{s2}) \right) + 1 \right] \end{aligned}$$



Relation between phase acceptance $\Delta\Phi$ and rf voltage V_{rf}
(Stationary energies, harmonic number and k-value are fixed.)

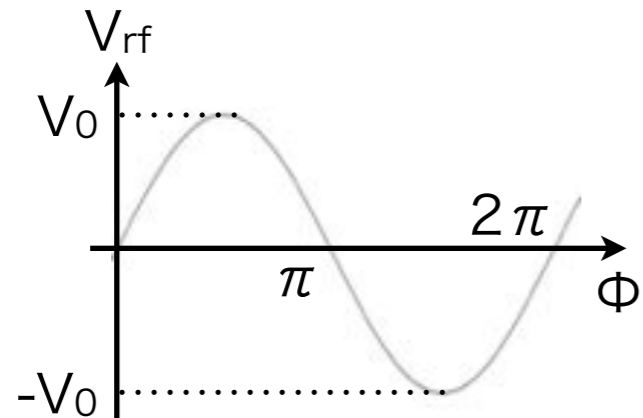


Phase acceptance becomes large with rf voltage.

Square rf voltage wave form in serpentine acceleration

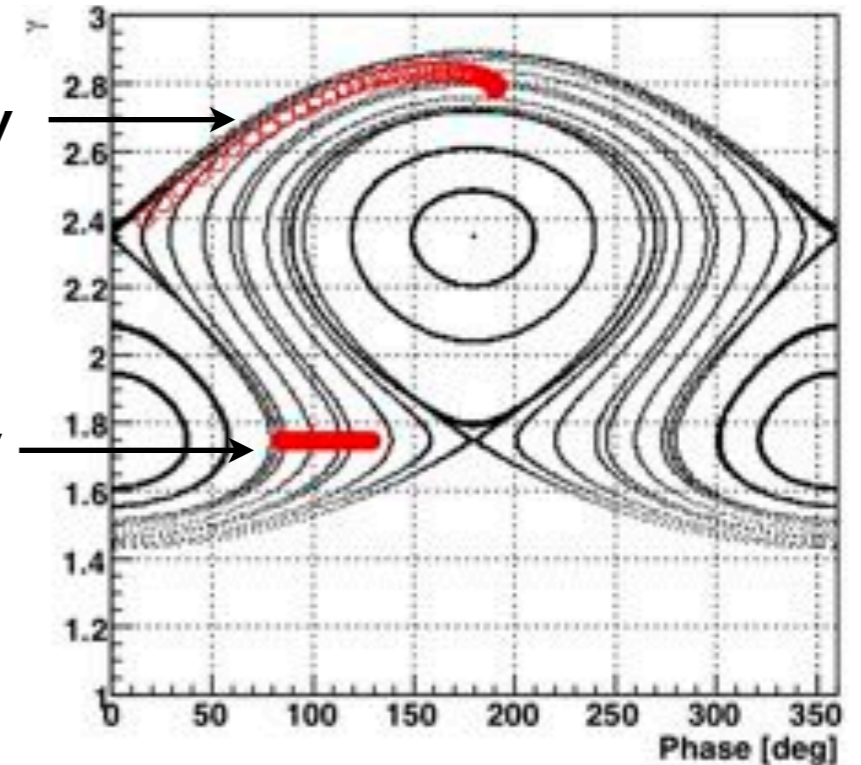
- Energy spread at extraction with sinusoidal rf voltage wave

$$V_{rf} = eV_0 \sin \phi$$



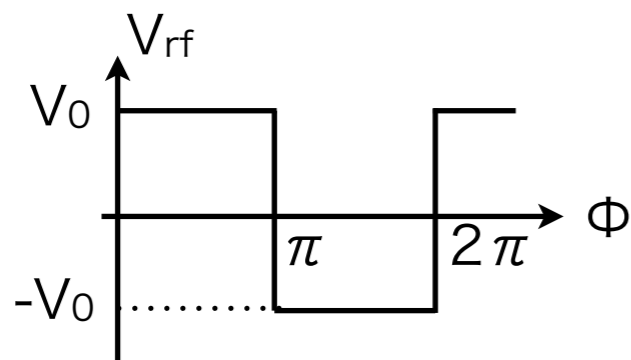
Extraction energy

Injected energy



Square rf voltage wave form is applied to obtain a mono-energetic extraction beam.

$$V_{rf} = \begin{cases} V_0 & : 0 < \phi \leq \pi, \\ -V_0 & : \pi < \phi \leq 2\pi. \end{cases}$$

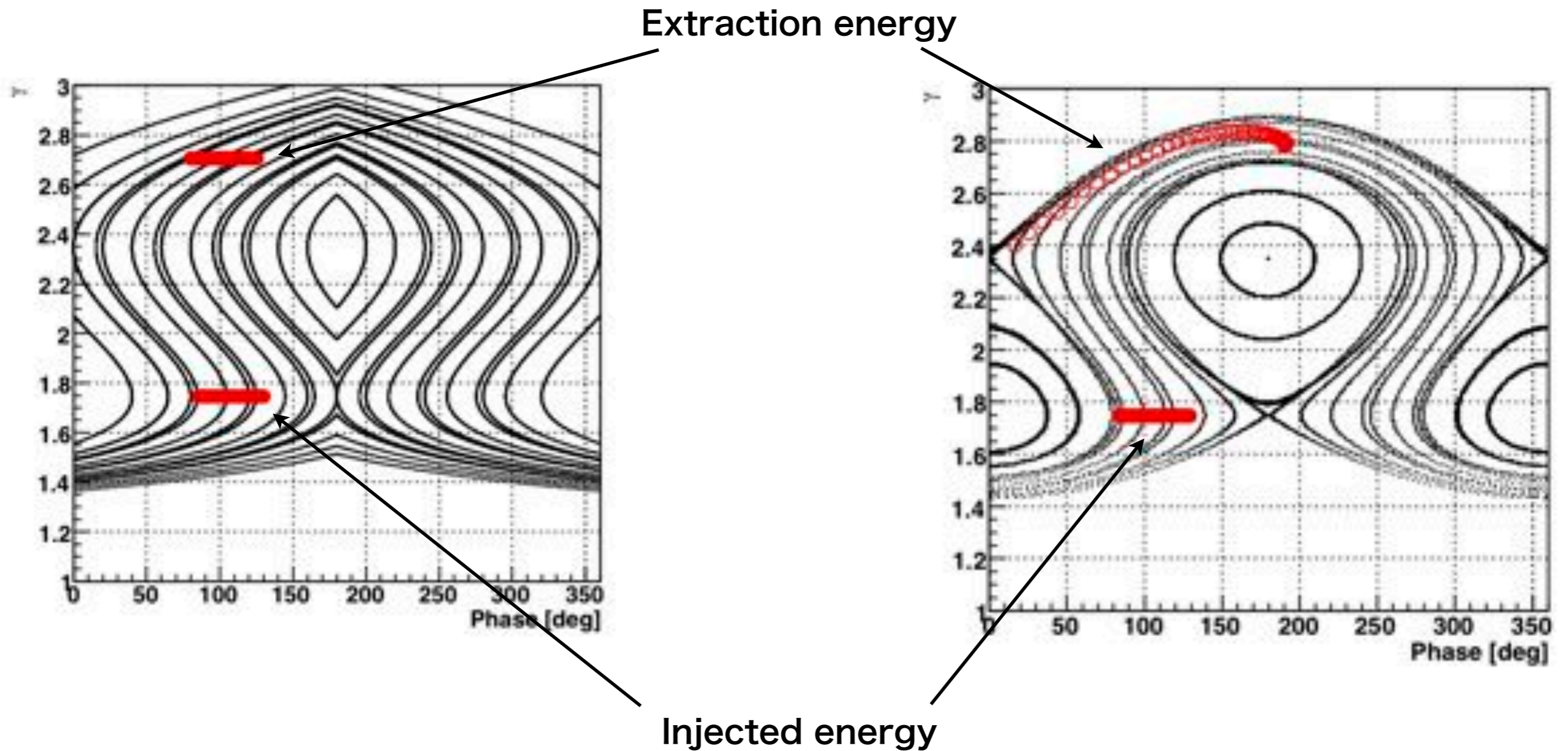


Hamiltonian

$$H(E, \phi; \Theta) = \begin{cases} h \left[\frac{1}{1 + \alpha} \frac{(E^2 - m^2)^{\frac{1+\alpha}{2}}}{E_s (E_s^2 - m^2)^{\frac{\alpha-1}{2}}} - E \right] - \frac{V_0}{2\pi} \phi & : 0 < \phi \leq \pi, \\ h \left[\frac{1}{1 + \alpha} \frac{(E^2 - m^2)^{\frac{1+\alpha}{2}}}{E_s (E_s^2 - m^2)^{\frac{\alpha-1}{2}}} - E \right] + \frac{V_0}{2\pi} \phi - V_0 & : \pi < \phi \leq 2\pi, \end{cases}$$

Square rf voltage wave form in serpentine acceleration

- Energy spread at extraction with sinusoidal and square rf voltage wave



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1. Demonstration of serpentine acceleration in scaling FFAG

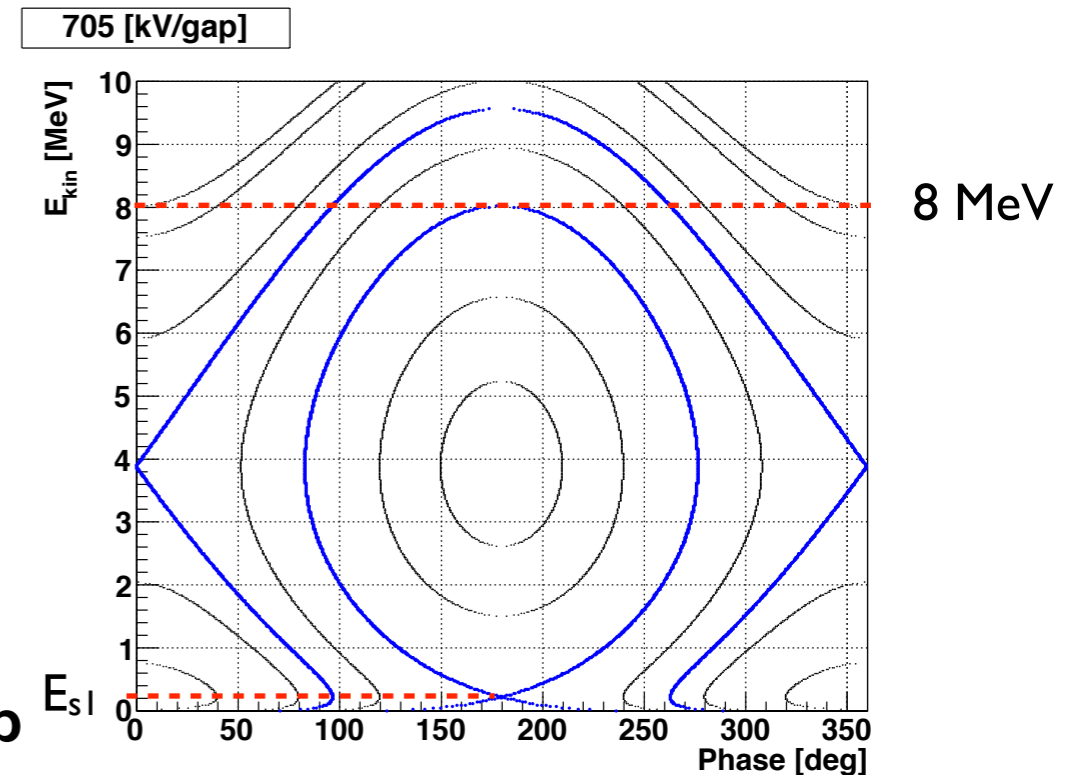
2. Phase acceptance study with rf voltage in serpentine acceleration

- Application of serpentine acceleration
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Demonstration of serpentine acceleration

● Injected beam parameters

Injected energy	160 keV
Stationary energy (E_{s1})	205 keV
Injected beam phase	< 20% of rf period
Rf voltage	705 kV/gap



● Measurement conditions

- Beam current in the ring is measured by Faraday cup
Faraday cup is fixed at 8 MeV closed orbit
- Central phase of injected beam is changed with respect to the rf phase at the cavity.



Beam cannot be injected in the stationary bucket above transition.

Particles injected in the stationary bucket below transition cannot be detected by the Faraday cup.

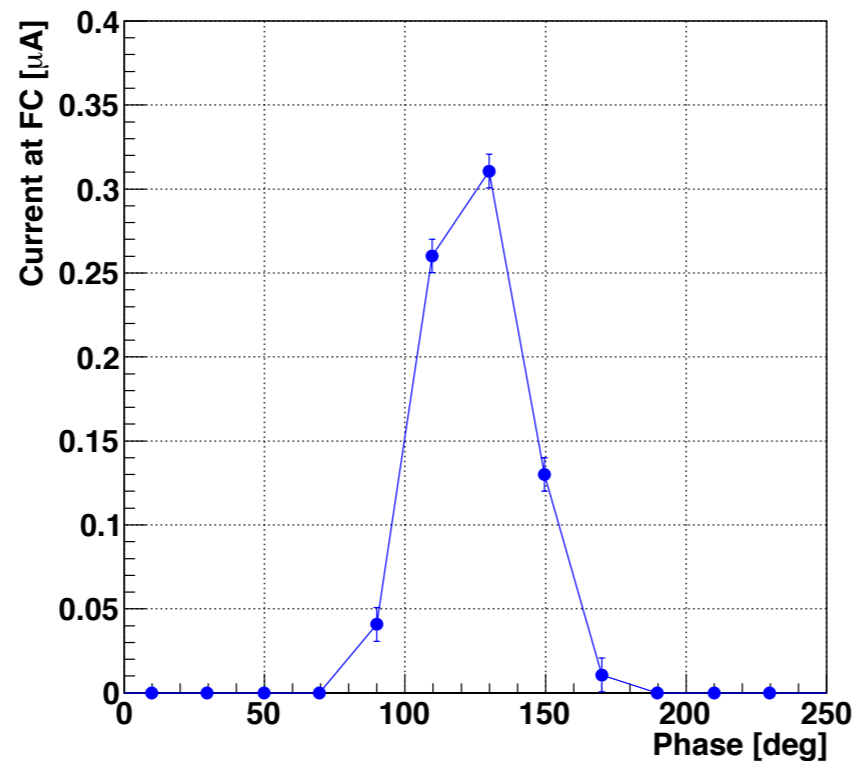


Only particles accelerated in serpentine channel can be measured by the Faraday cup.

Serpentine加速法の実証実験



Results



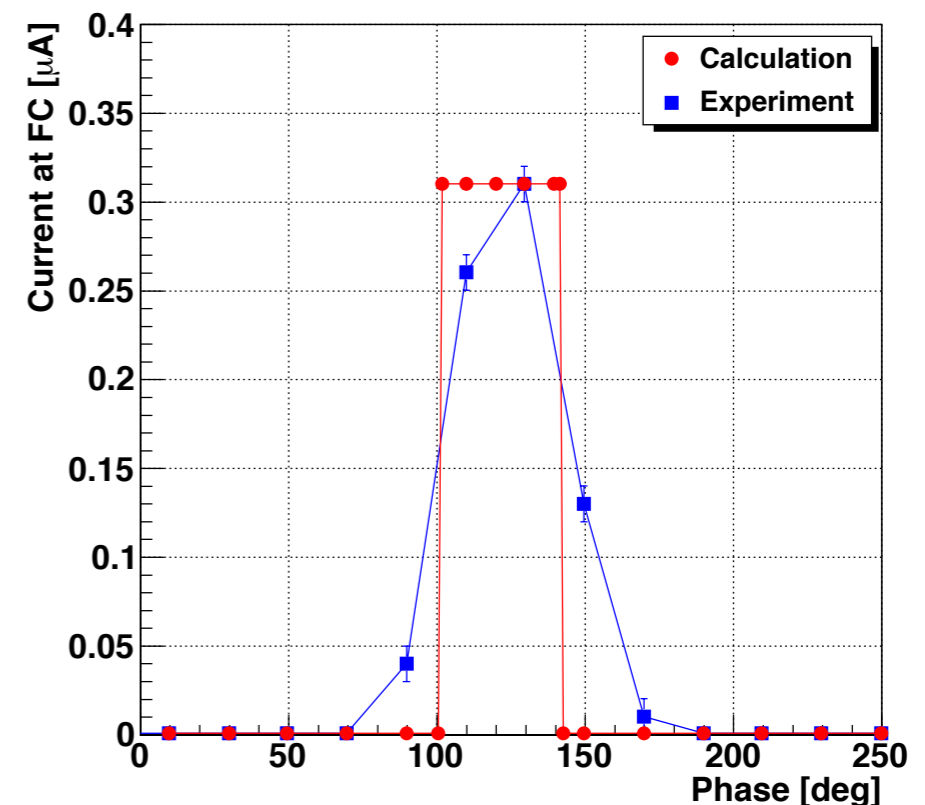
Simulation

Single particle longitudinal tracking has been done.

When the particle reaches 8 MeV,
current is set to the maximum measured value.



There is a discrepancy due to the phase range of
injected beam.

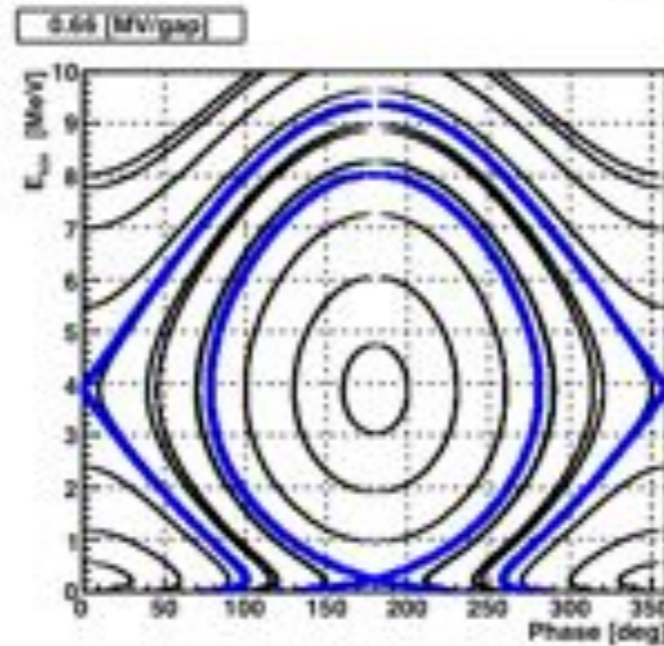
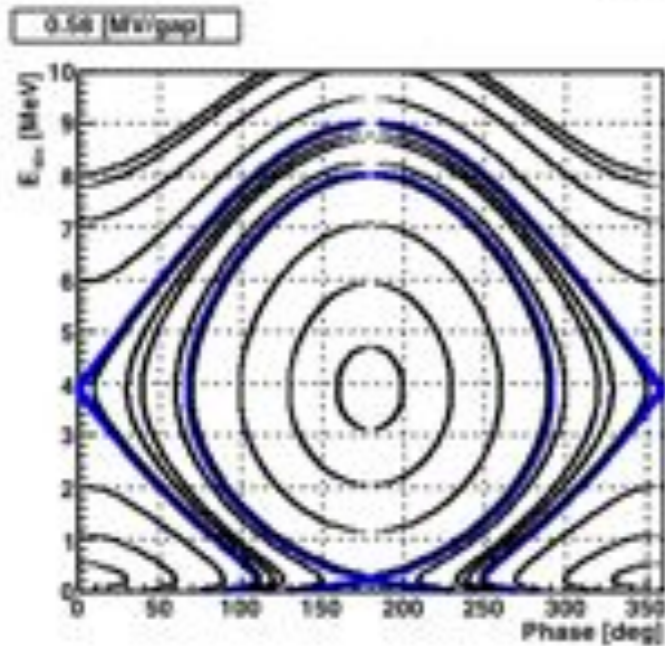
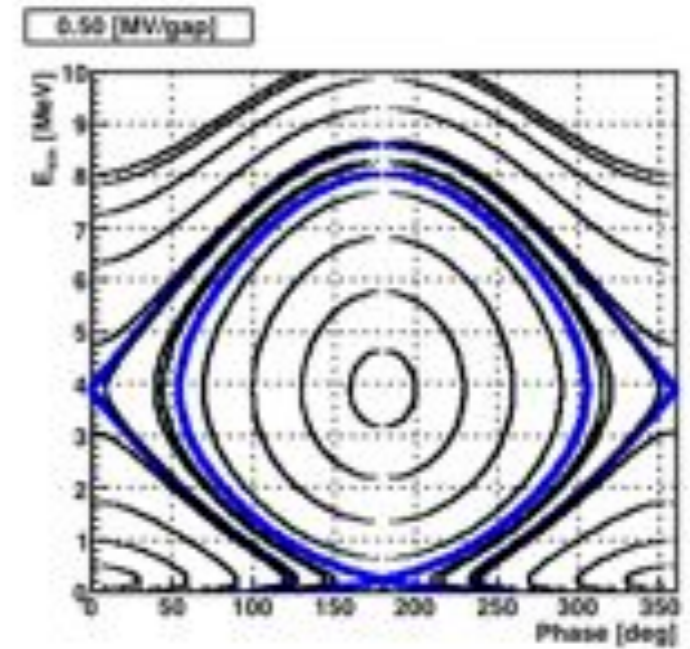
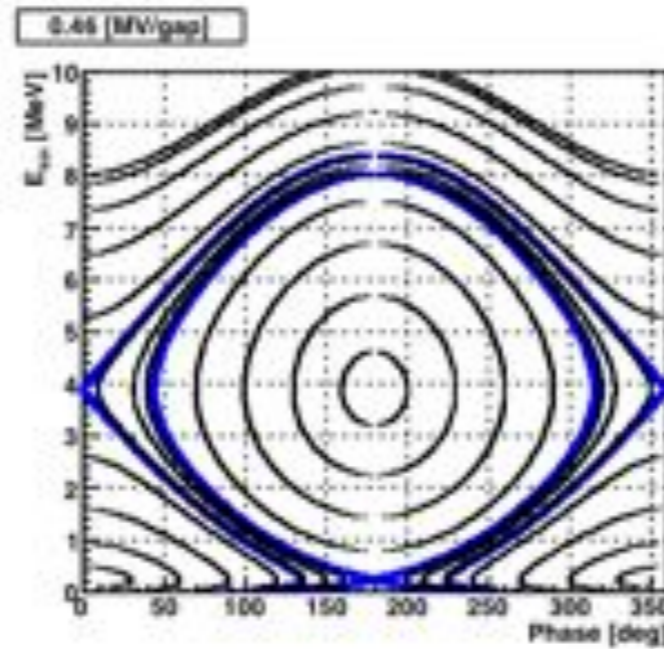
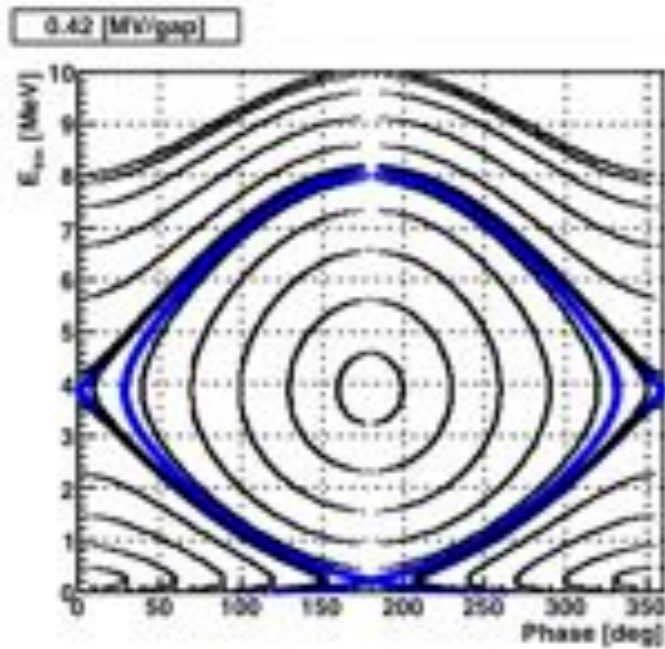


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Phase acceptance measurement

- Phase acceptance becomes large with rf voltage



Phase acceptance measurement

● Injected beam parameters

Injected energy	200 keV
Stationary energy (E_{s1})	205 keV
Injected beam phase	< 20% of rf period
Rf voltage	420 to 660 kV/gap

● Measurement conditions

- Beam current in the ring is measured by Faraday cup

Faraday cup is fixed at **8 MeV** closed orbit

- Central phase of injected beam is changed with respect to the rf phase at the cavity.



Beam cannot be injected in the stationary bucket above transition.

Particles injected in the stationary bucket below transition cannot be detected by the Faraday cup.

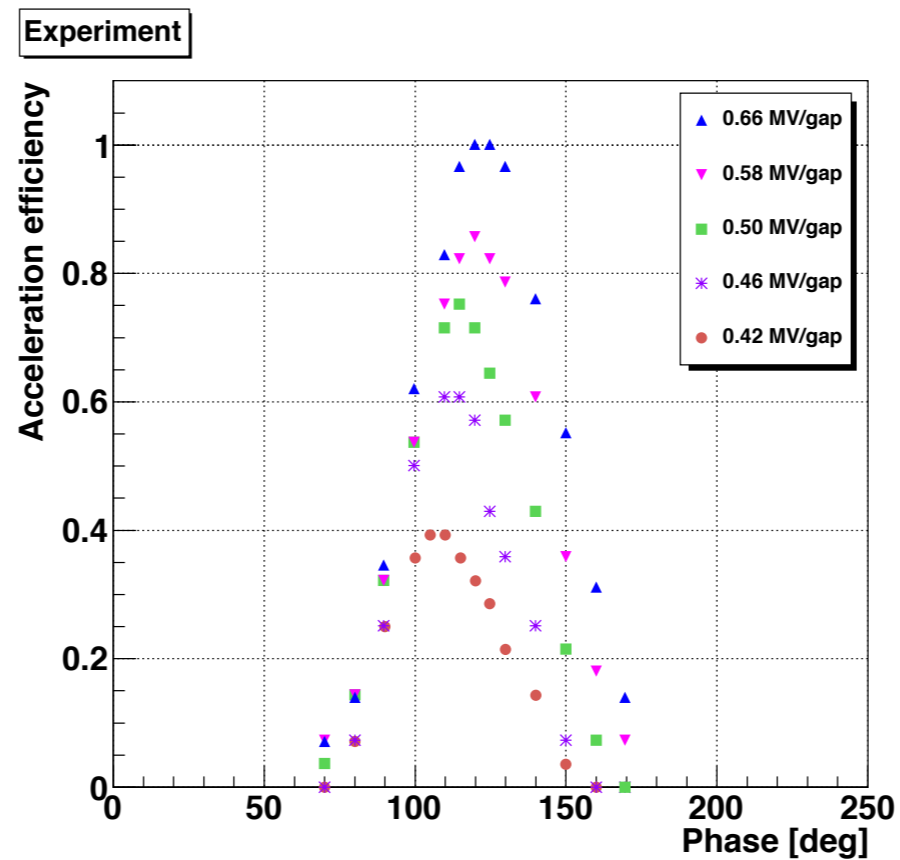


Only particles accelerated in serpentine channel can be measured by the Faraday cup.

Phase acceptance measurement



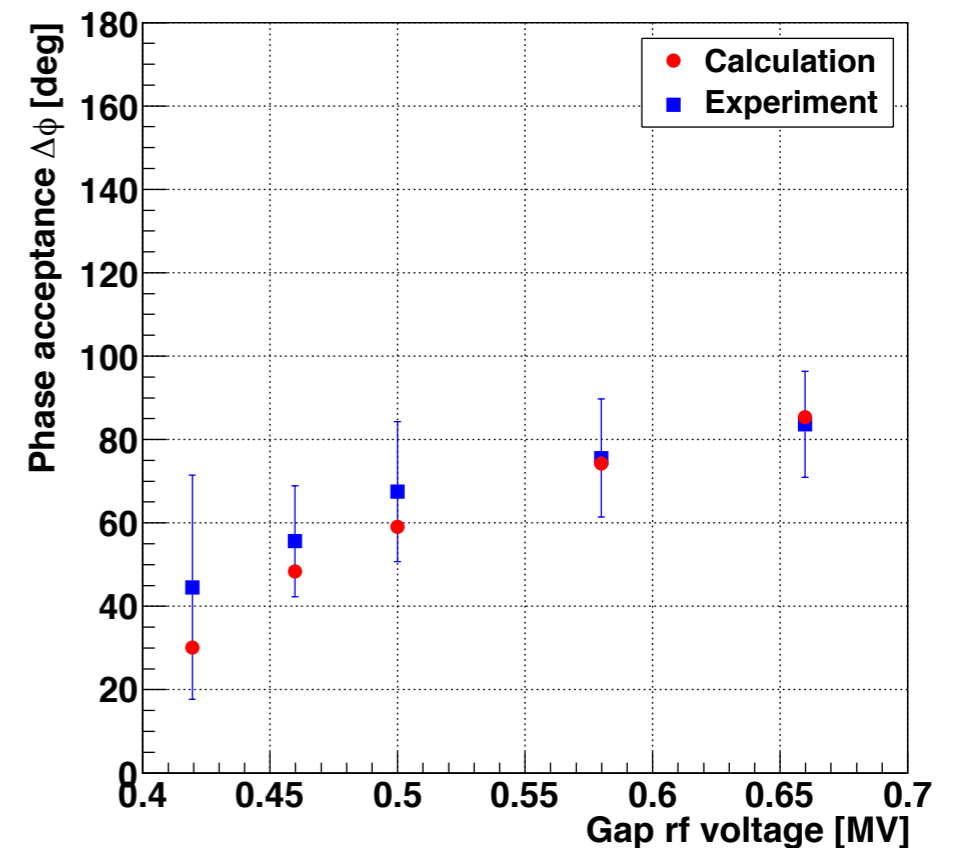
Results



Simulations

Single particle longitudinal tracking has been done.

Tendency of measurement results agree with the simulations.



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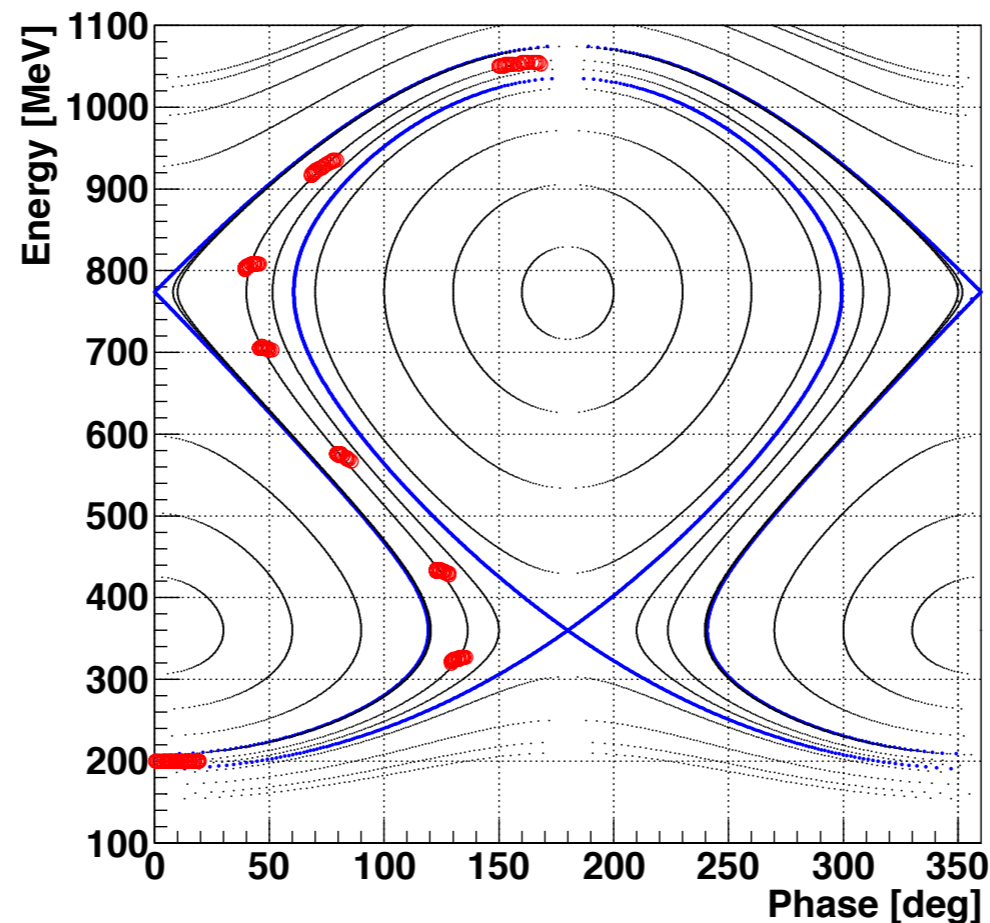
Application of serpentine acceleration

- Summary and Conclusion

Application of serpentine acceleration

● Proton driver for ADS

k -value	1.45
Equivalent mean radius at 200 MeV [m]	3
Equivalent mean radius at 1 GeV [m]	5.9
Stationary kinetic energy below transition [MeV]	360
rf voltage [MV/turn]	15 ($h=1$)
rf frequency [MHz]	9.6($h=1$)



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Summary and Conclusion

Summary and Conclusion

- **A new type of fixed rf acceleration in scaling FFAG has been developed analytically.**
- **Proof of principle has been conducted experimentally with electron machine.**
- **Application of serpentine acceleration has been proposed with longitudinal design of proton driver for ADS**
- **This success also opens new possibilities for neutrino factory.**

