



Study of the K^+K^- FSI in proton – proton and electron - positron collisions



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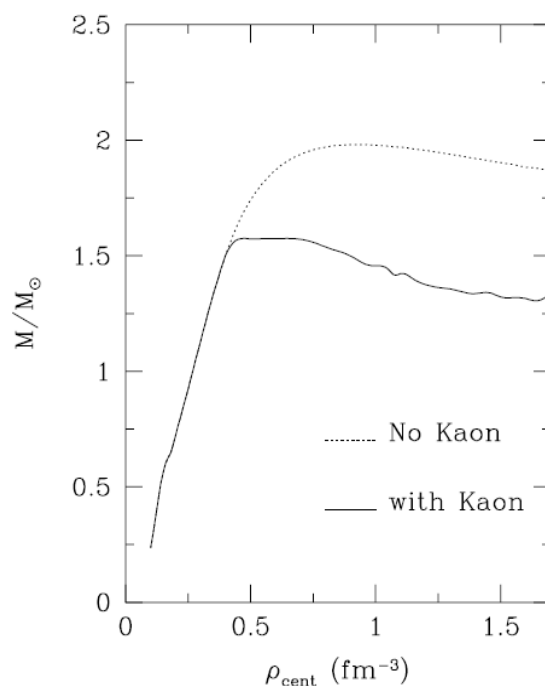
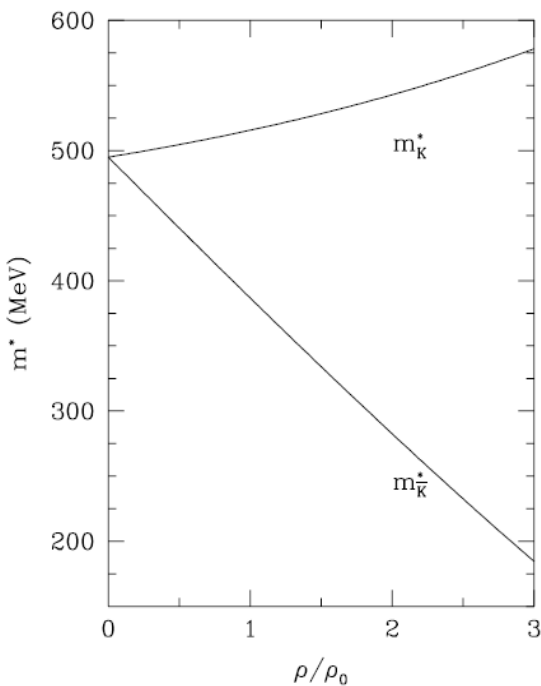


- ❖ **Motivation**
- ❖ **Proton-proton collisions at K^+K^- threshold: COSY**
- ❖ **Near future: KLOE-2 @ DAΦNE**

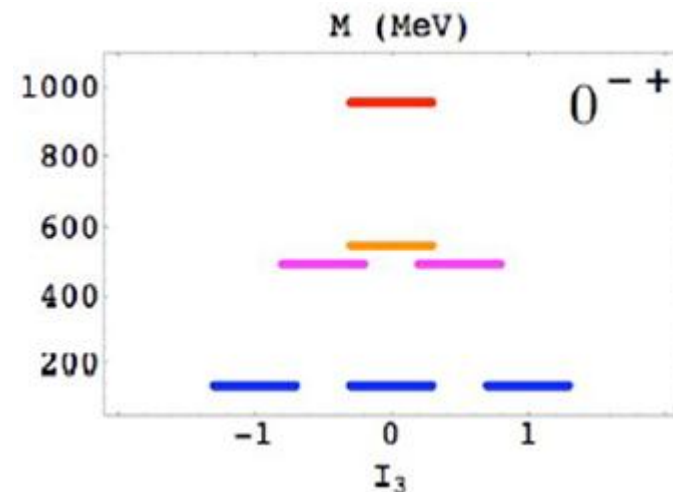
3-9 February 2013 Bjelasnica Mountain, Sarajevo

Motivation

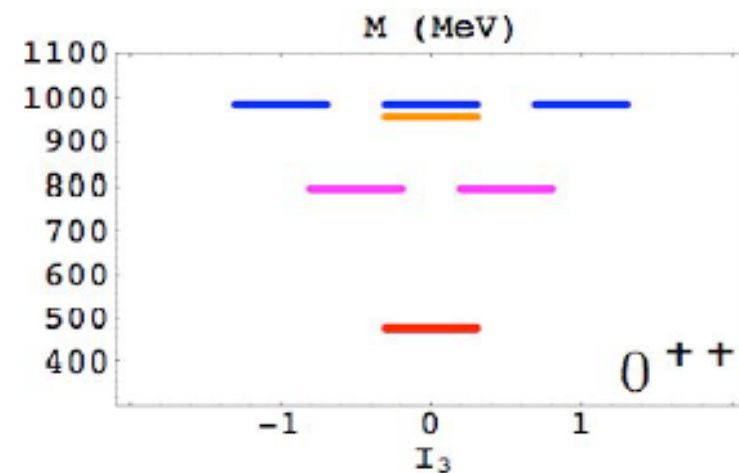
- ❖ a_0 and f_0 mesons as a K^+K^- molecules
- ❖ Physics of neutron stars:
kaon condensates



Pseudoscalar mesons



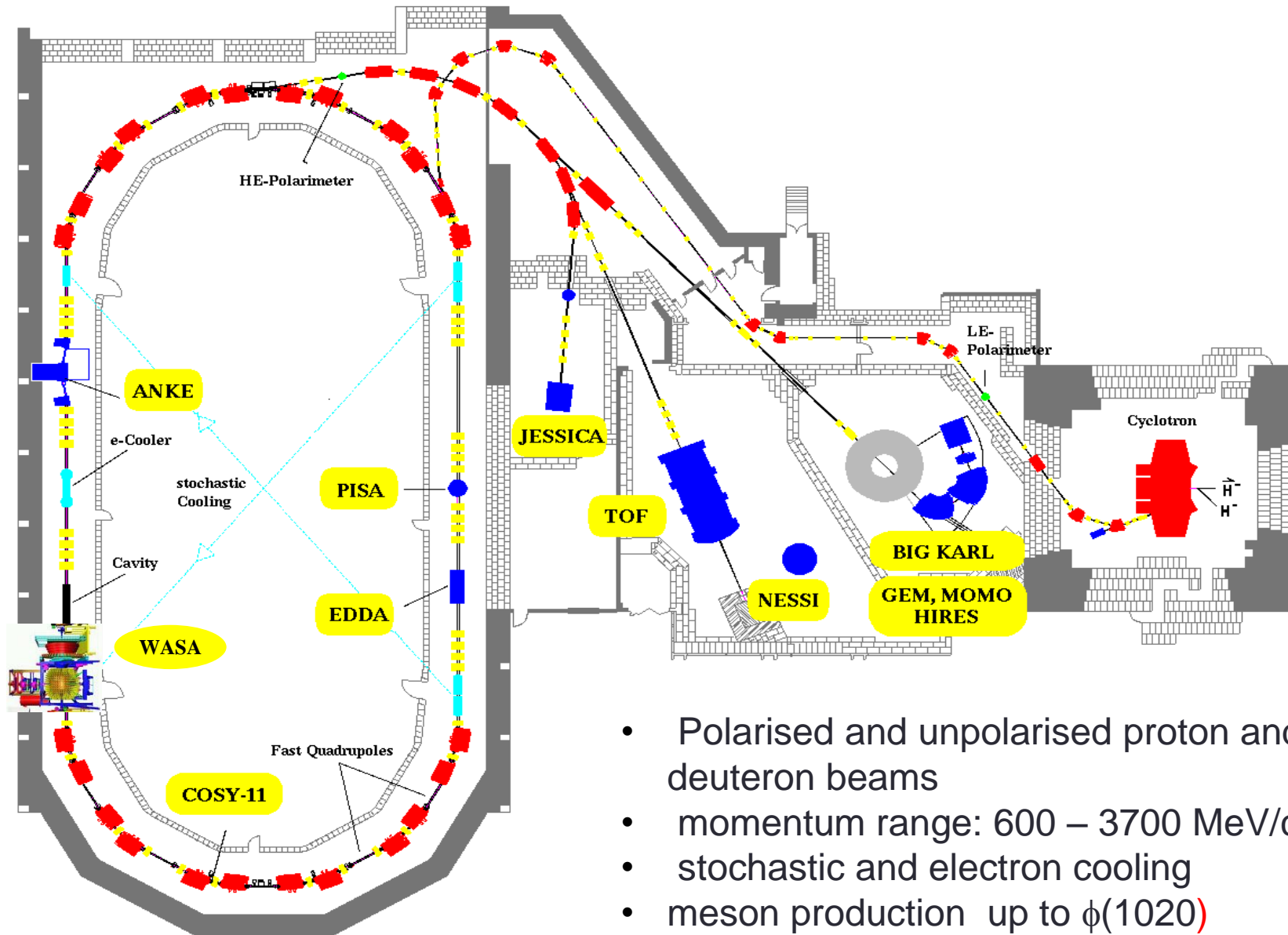
Scalar multiplet:
 $\sigma(500)$, $\kappa(700)$, $f_0(980)$, $a_0(980)$



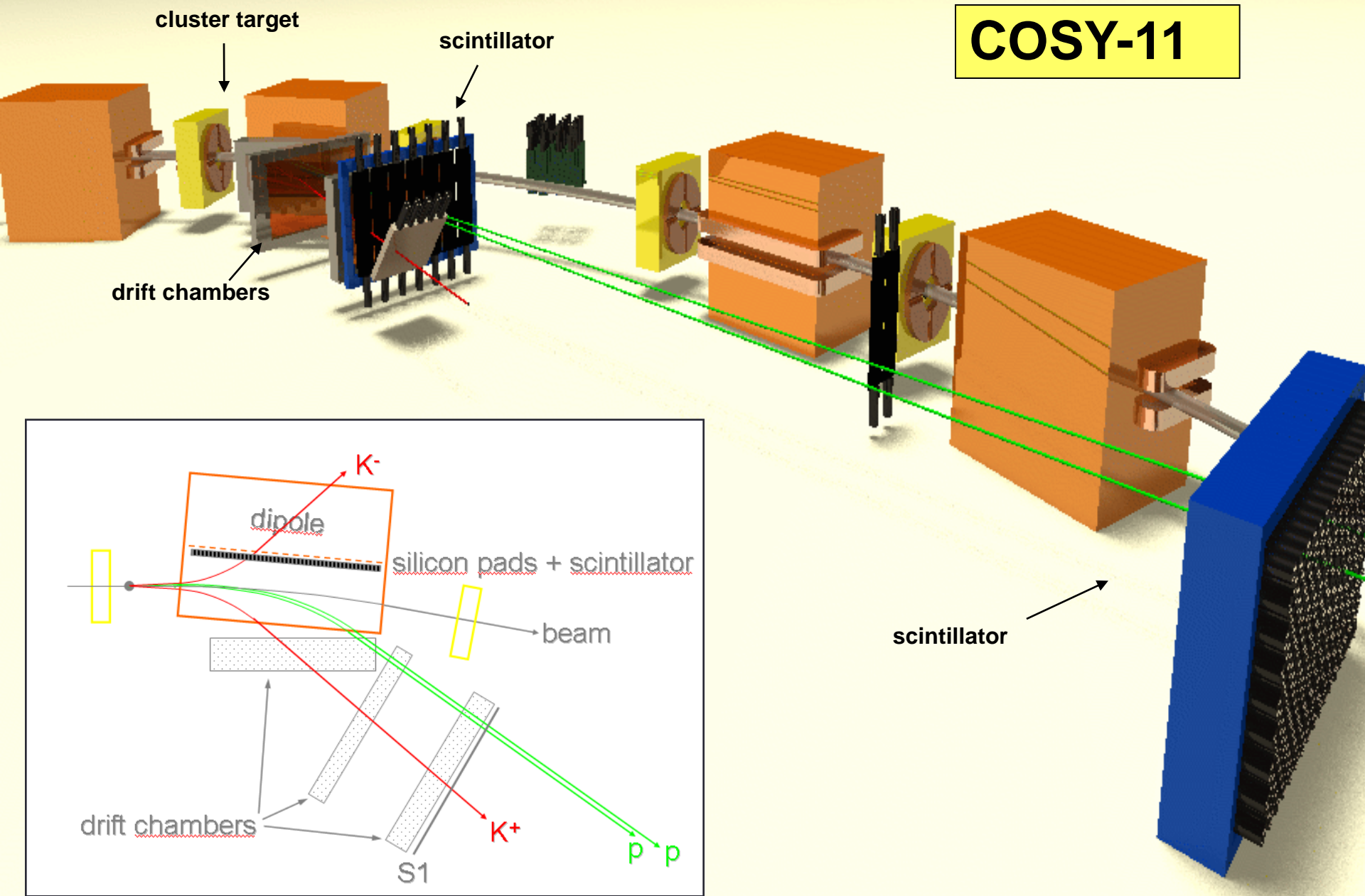


proton-proton collisions at K^+K^- threshold: COSY

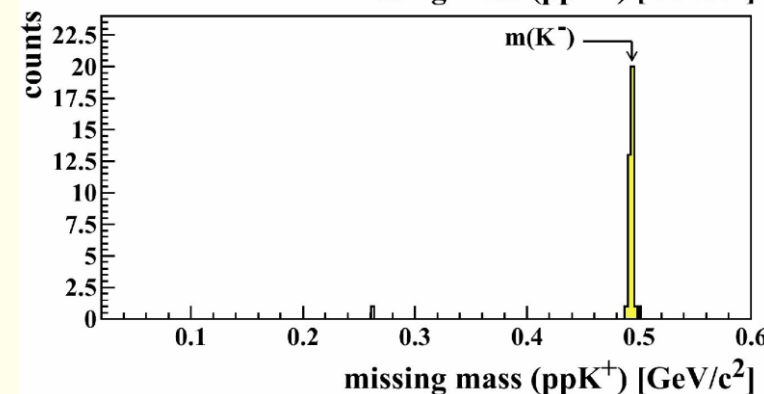
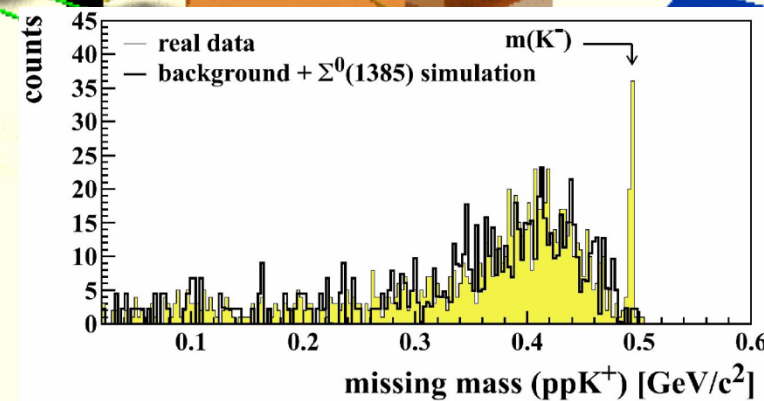
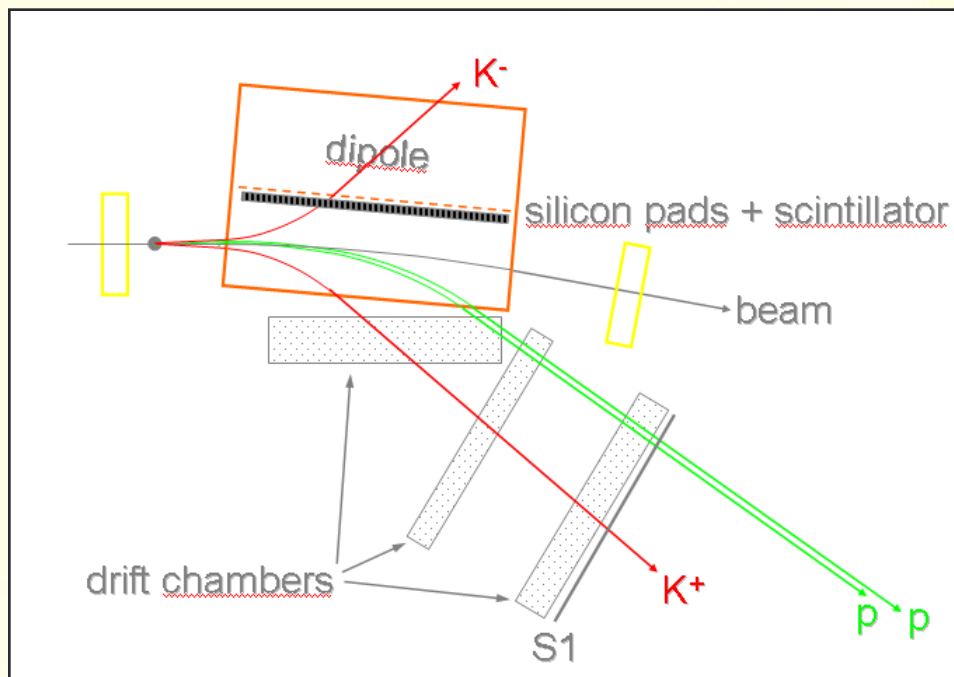
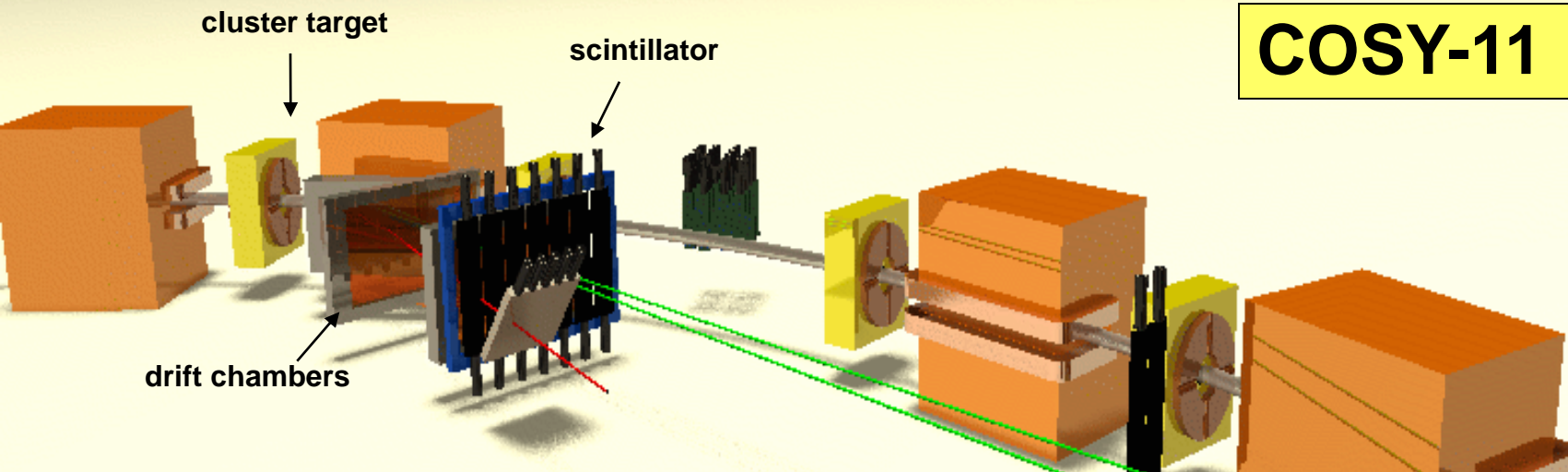
COoler SYNchrotron COSY



COSY-11



COSY-11



Excitation function

- ❖ FSI indication in both total and differential cross sections at K^+K^- threshold

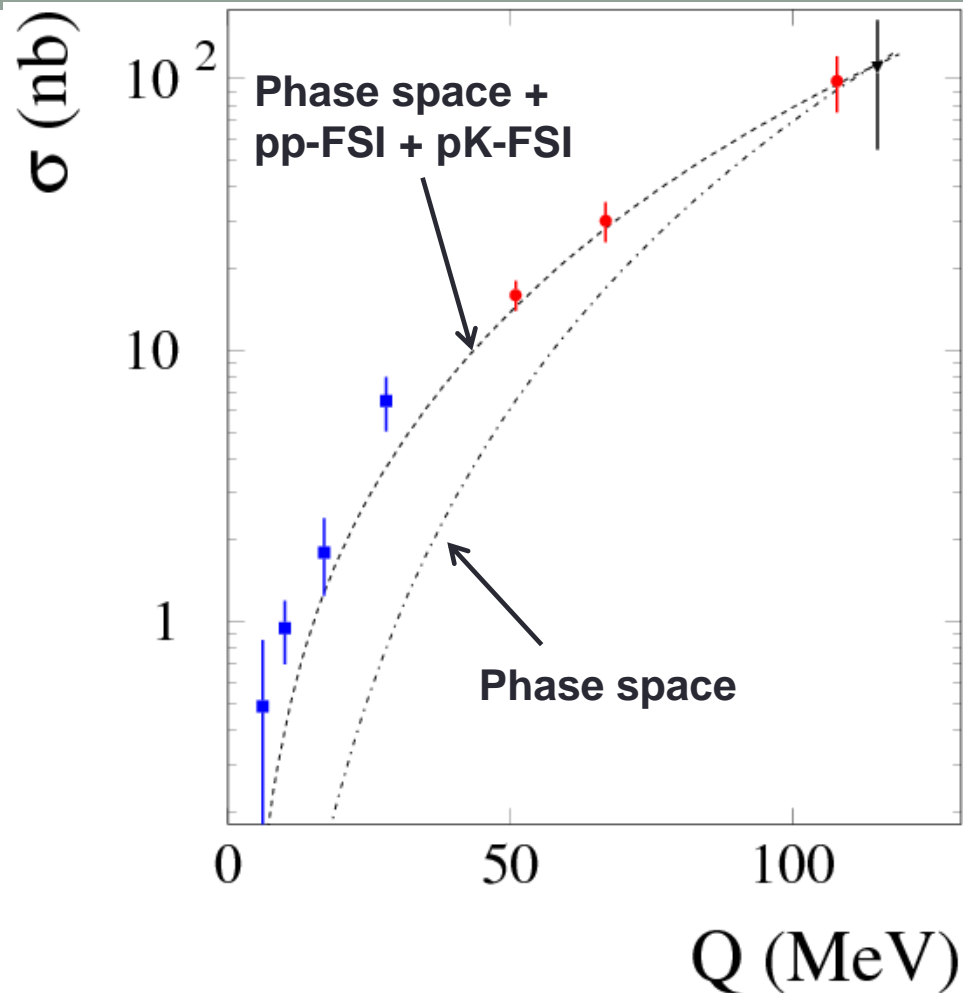
$$\left| M_{pp \rightarrow ppK^+K^-} \right|^2 \approx |M_0|^2 |F_{FSI}|^2$$

$$F_{FSI} = F_{pp}(q) \times F_{p_1K^-}(k_1) \times F_{p_2K^-}(k_2)$$

$$F_{pp}(q) = \frac{e^{i\delta_{pp}(^1S_0)} \times \sin \delta_{pp}(^1S_0)}{C \times q}$$

$$F_{pK^-}(k) = \frac{1}{1 - ika}$$

$$a = (0 + i1.5) [\text{fm}]$$



DISTO: F. Balestra et al., Phys. Rev. C 63, 024004 (2001)

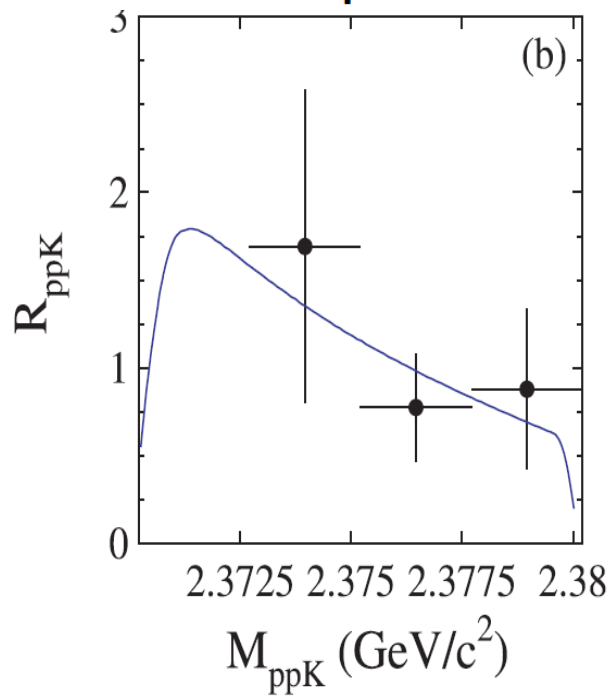
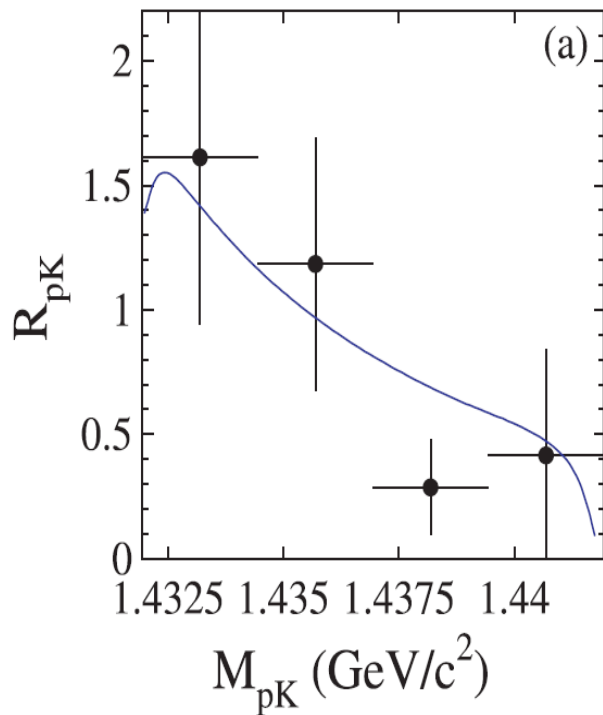
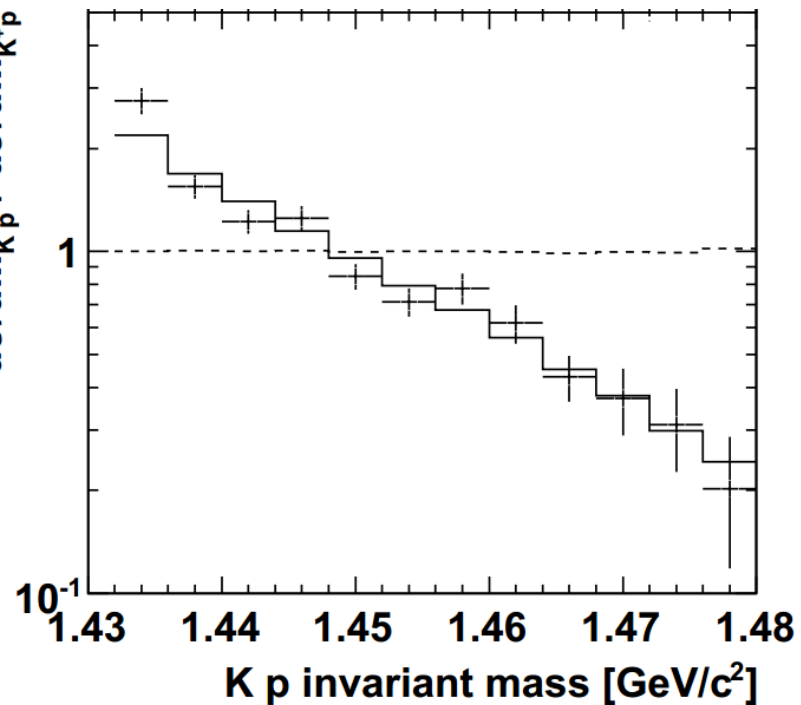
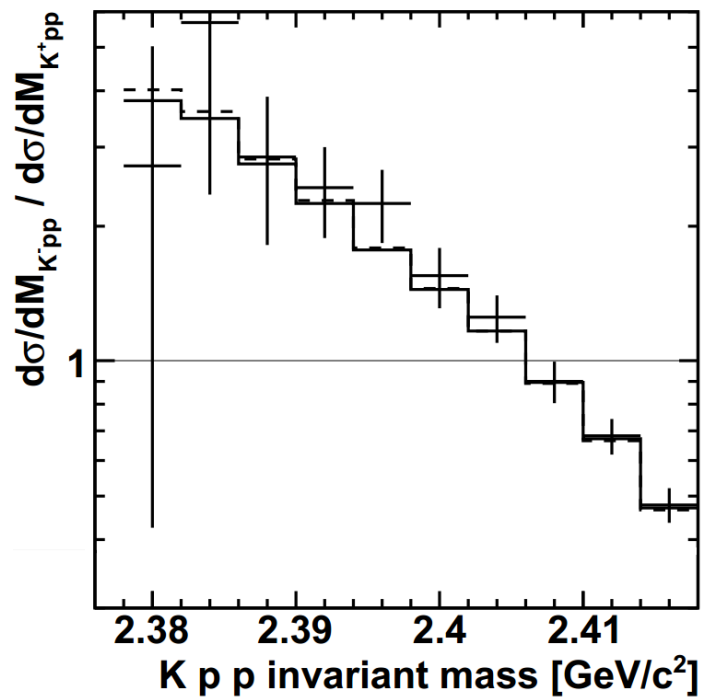
ANKE: Y. Maeda et al. Phys., Rev. C 77, 01524 (2008)

ANKE: Q. J. Ye et al., Phys. Rev. C 85, 035211 (2012)

COSY-11: C. Quentmeier et al., Phys. Lett. B 515 (2001) 276-282

COSY-11: P. Winter et al., Phys. Lett. B 635 (2006) 23-29

COSY-11: M. Wolke, PhD thesis

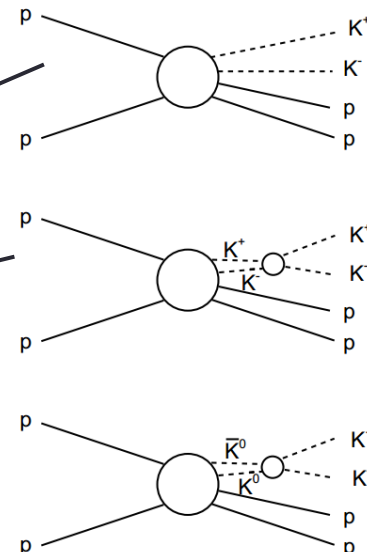


Coupled channel effects

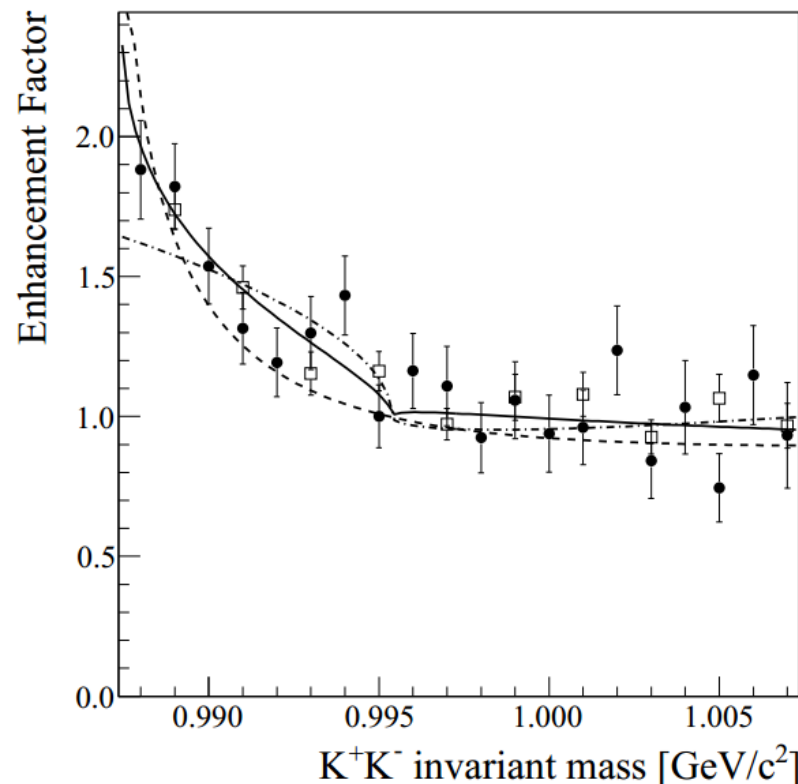
$$\mathcal{F} = \left| \frac{B_1/(B_1 + B_0)}{(1 - i\frac{1}{2}q[A_1 - A_0])(1 - ikA_1)} + \frac{B_0/(B_1 + B_0)}{(1 - i\frac{1}{2}q[A_0 - A_1])(1 - ikA_0)} \right|$$

ANKE: A. Dzyuba et al., Phys. Lett. B668, 315 (2008).

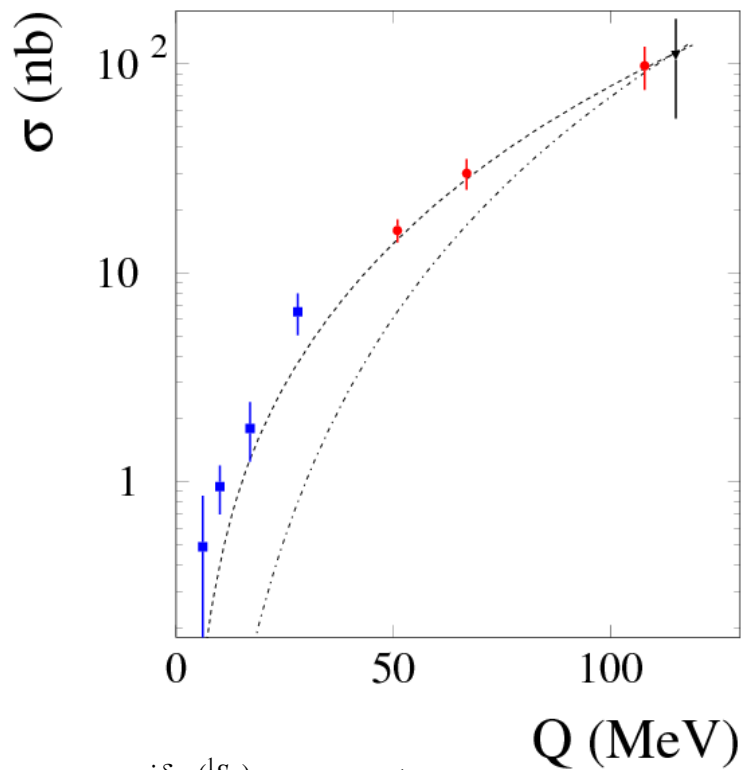
$A_0 = (-0.45 + i1.63)$ fm ; $A_1 = (0.1 + i0.7)$ fm
(M. Ablikim et al., Phys. Lett. B 607 (2005) 243;)



- ❖ With the ANKE statistics the expected cusp effects are not distinguishable from the elastic scattering of K^+ and K^-
- ❖ Izospin $I=0$ state is favourable
- ❖ No indication of the $f_0(980)/a_0(980)$ influence
- ❖ More statistics at lower excess energy needed



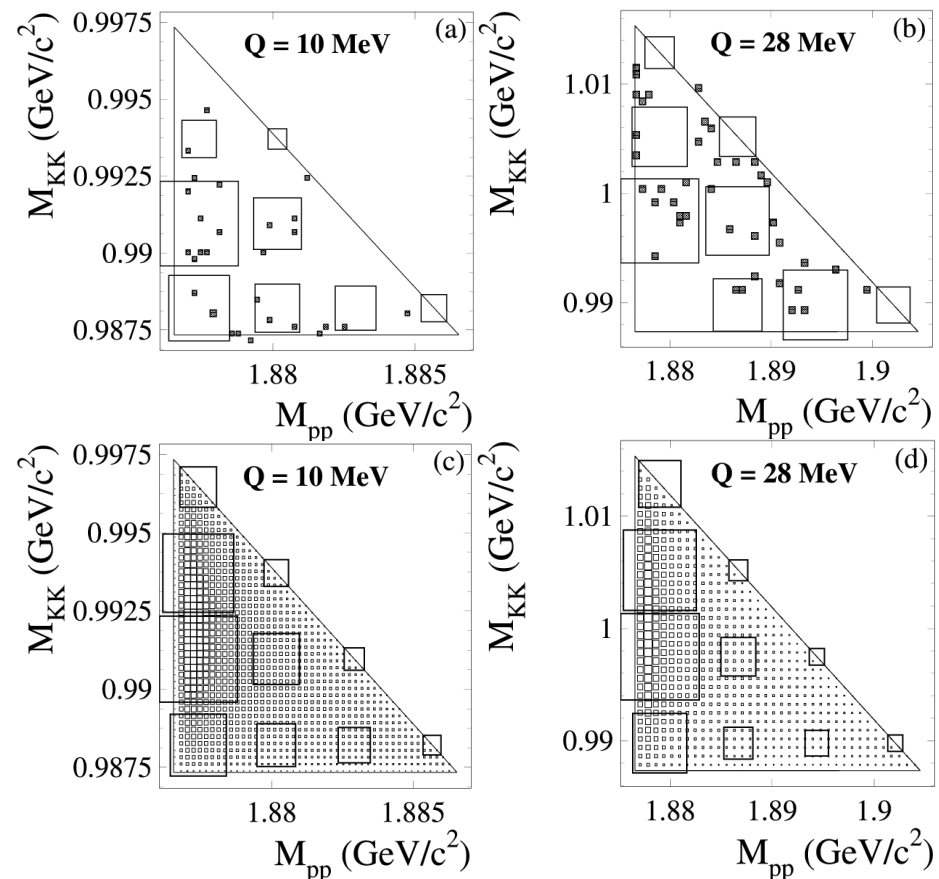
Analysis of the K^+K^- -FSI at COSY-11



$$F_{pp}(q) = \frac{e^{i\delta_{pp}(^1S_0)} \sin \delta_{pp}(^1S_0)}{Cq}$$

$$a_{pK^-} = (-0.65 + i0.78) [\text{fm}]$$

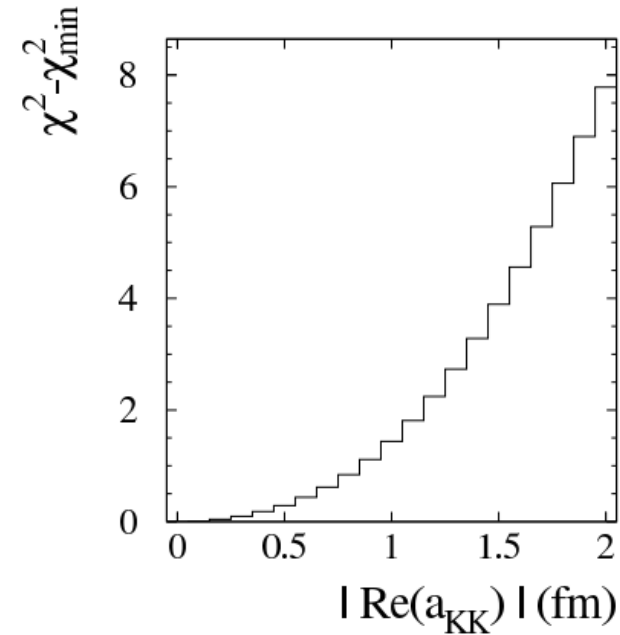
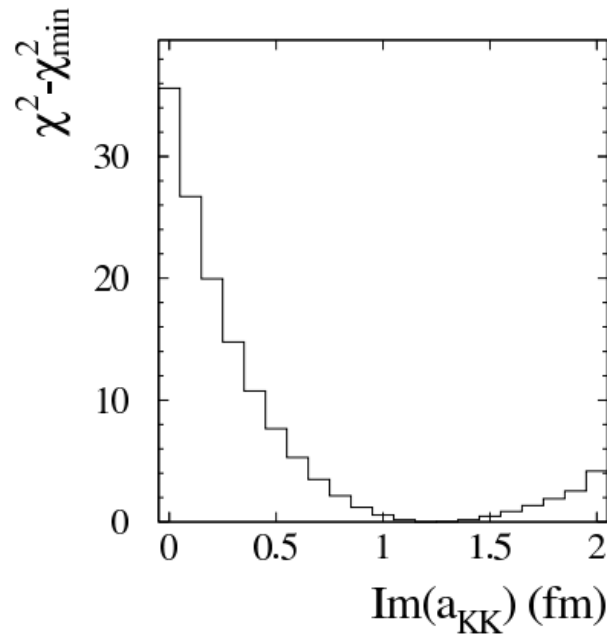
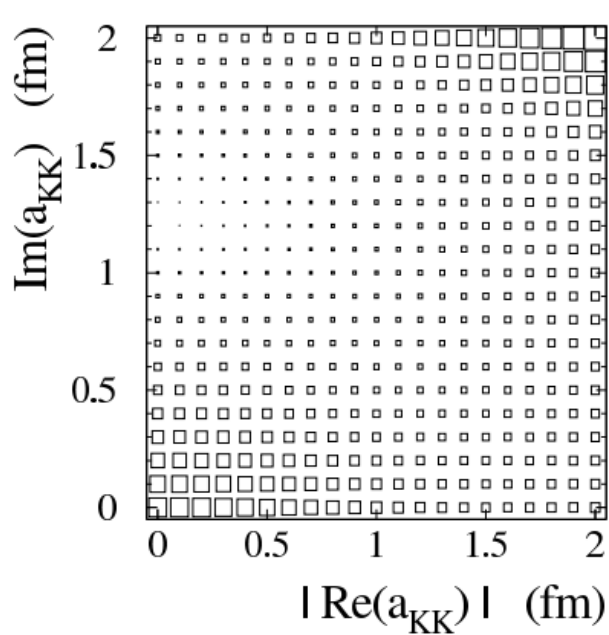
(Y. Yan, arXiv:0905.4818 [nucl-th])



M. Silarski, et al., Phys. Rev. C 80, 045202 (2009)

- ❖ Analysis of the Goldhaber plots measured at $Q = 10$ MeV (27 events) and $Q = 28$ MeV (30 events) + near threshold excitation function

Analysis of the K^+K^- -FSI at COSY-11



$$\chi^2(a_{K^+K^-}, \alpha) = \sum_{i=1}^8 \frac{(\sigma_i^{exp} - \alpha \sigma_i^m)^2}{(\Delta \sigma_i^{exp})^2} + 2 \cdot \sum_{j=1}^2 \sum_{k=1}^{10} [\beta_j N_{jk}^s - N_{jk}^e + N_{jk}^e \ln(\frac{N_{jk}^e}{\beta_j N_{jk}^s})]$$

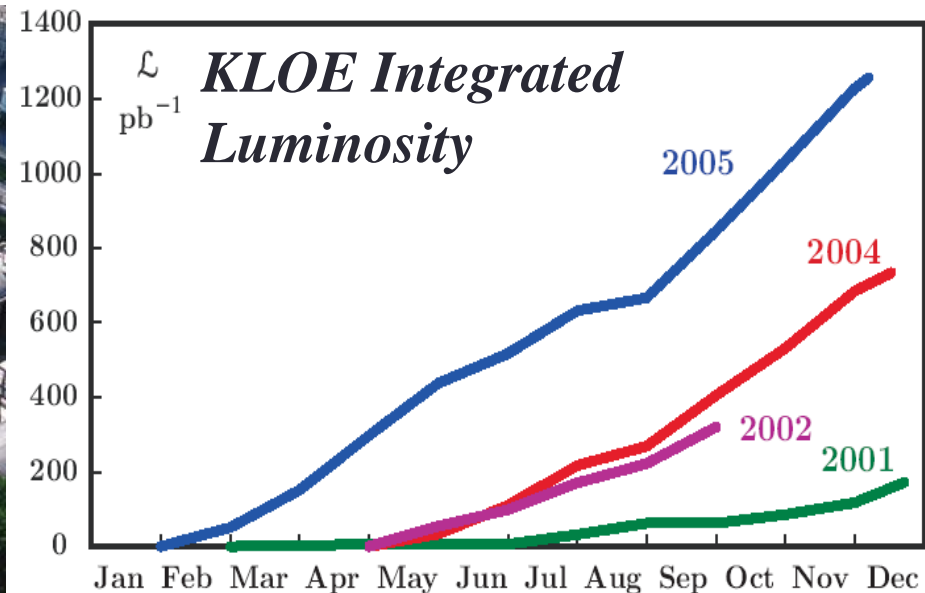
$$|\text{Re}(a_{K^+K^-})| = 0.0 \begin{matrix} +1.1_{stat} \\ -0.0_{stat} \end{matrix} \text{ fm}$$

$$\text{Im}(a_{K^+K^-}) = 1.1 \begin{matrix} +0.6_{stat} & +0.9_{sys} \\ -0.5_{stat} & -0.6_{sys} \end{matrix} \text{ fm}$$



**Near future:
KLOE-2 @ DAΦNE**

DAΦNE Luminosity history



KLOE run:

- ❑ Daily performance: 7-8 pb⁻¹
- ❑ Best month $\int L dt \sim 200 \text{ pb}^{-1}$
- ❑ Total KLOE $\int L dt \sim 2400 \text{ pb}^{-1}$ at ϕ mass peak
+ 250 pb⁻¹ off peak (@ 1 GeV)

BR's for selected Φ decays

K^+K^-	49.1%
$K_S K_L$	34.1%
$\rho\pi + \pi^+\pi^-\pi^0$	15.5%

KLOE (K Long Experiment)

Large cylindrical drift chamber

- ❑ Uniform tracking and vertexing in all volume
- ❑ Helium based gas mixture (90% He - 10% IsoC₄H₁₀)
- ❑ Stereo wire geometry

$$\sigma_p/p = 0.4 \%$$

$$\sigma_{xy} = 150 \mu\text{m}; \sigma_z = 2 \text{ mm}$$

$$\sigma_{\text{vtx}} \sim 3 \text{ mm}$$

$$\sigma(M_{\pi\pi}) \sim 1 \text{ MeV}$$

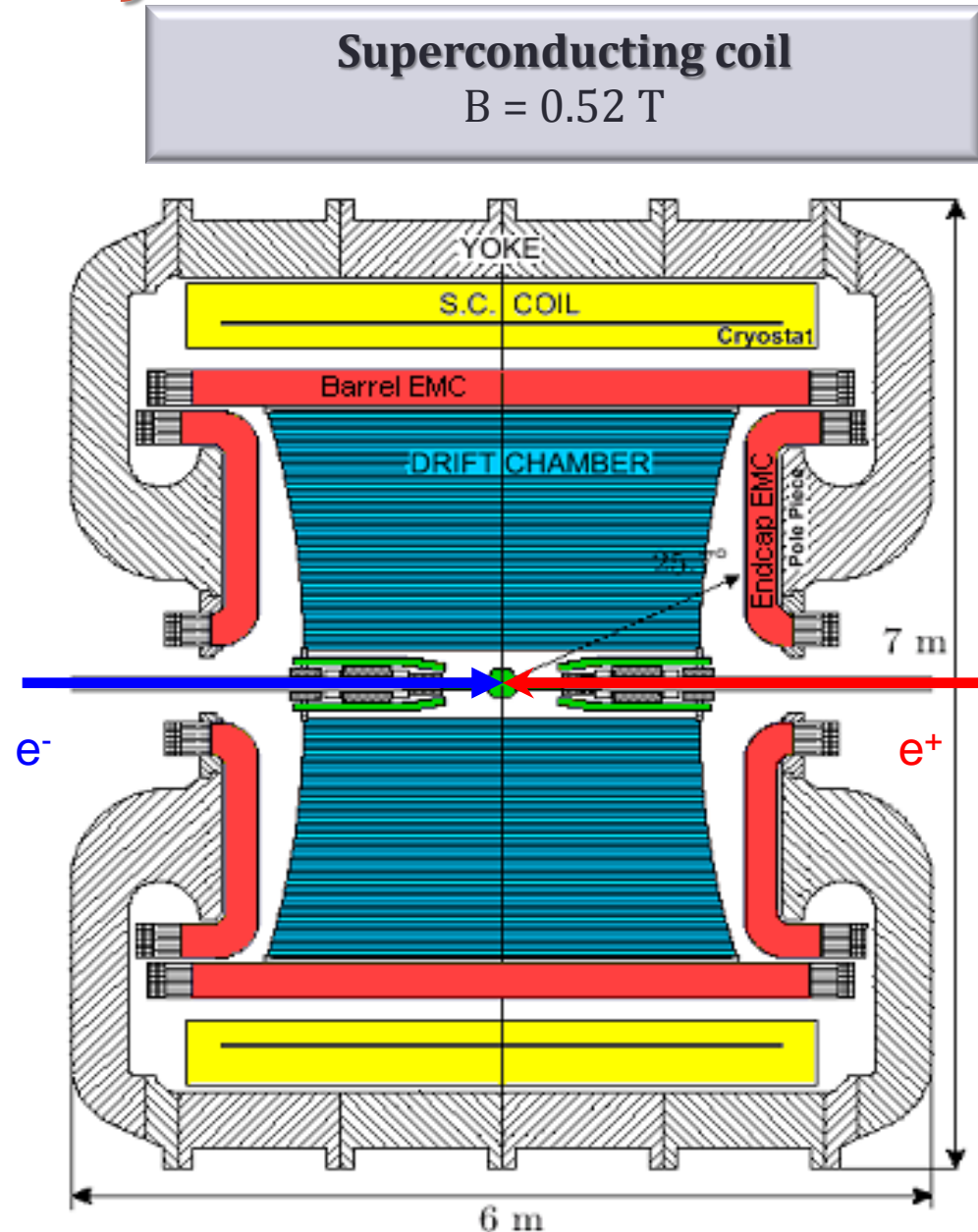
Lead/scintillating-fiber calorimeter

- ❑ Hermetical coverage
- ❑ High efficiency for low energy photons

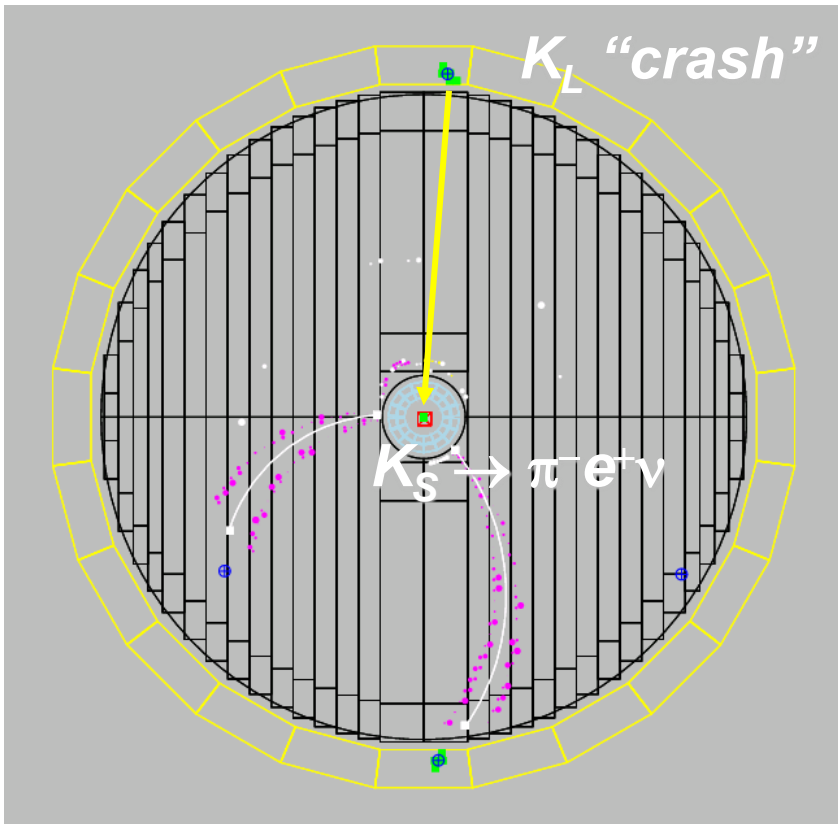
$$\sigma_E/E = 5.7\% / \sqrt{E(\text{GeV})}$$

$$\sigma_t = 57 \text{ ps} / \sqrt{E(\text{GeV})} \oplus 140 \text{ ps}$$

$$\sigma_{\text{vtx}}(\gamma\gamma) \sim 1.5 \text{ cm}$$



A Φ -factory offers the possibility to select pure kaon beams:

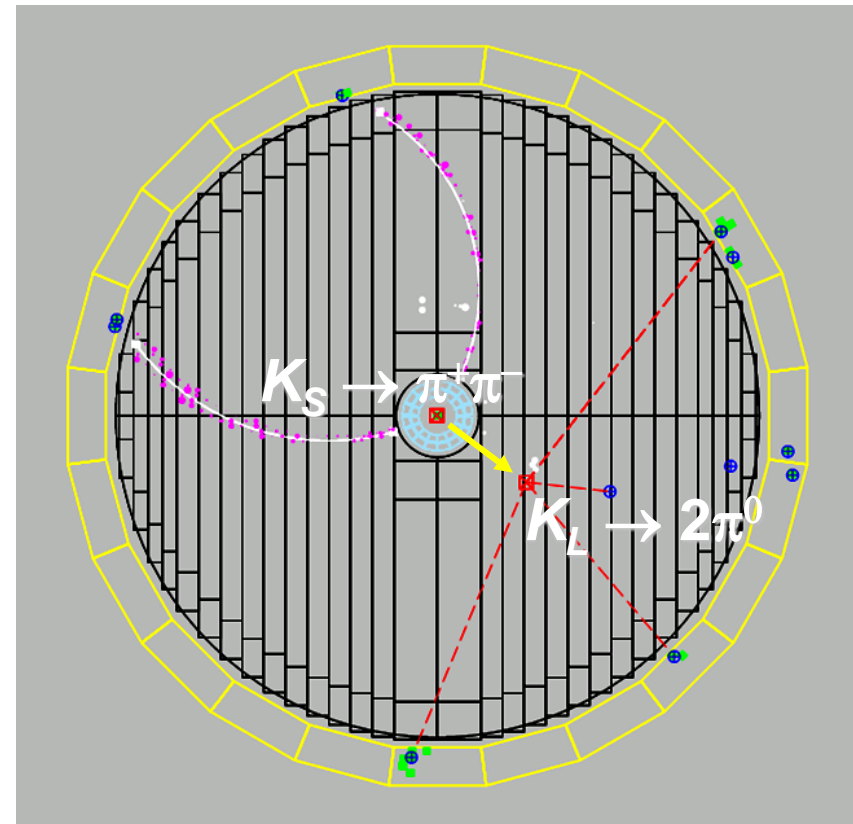


K_S tagged by K_L interaction in EmC

Efficiency $\sim 30\%$

K_S angular resolution: $\sim 1^\circ$ (0.3° in φ)

K_S momentum resolution: ~ 2 MeV



K_L tagged by $K_S \rightarrow \pi^+\pi^-$ vertex at IP

Efficiency $\sim 70\%$

K_L angular resolution: $\sim 1^\circ$

K_L momentum resolution: ~ 2 MeV

KK -FSI @ KLOE-2

$$e^+e^- \rightarrow K^+K^-\gamma$$

- ❖ **Advantage with respect to $pp \rightarrow ppK^+K^-$: only two interacting particles** (estimated scattering length independent from the FSI model)
- ❖ The cross section including both ISR & FSR $\sim 7\text{nb}$ (calculated with PHOKHARA at 1 GeV) $\Rightarrow 200 \text{ pb}^{-1}$ of integrated luminosity $\sim 1.4 \cdot 10^6$ events ($10^2 - 10^3$ higher statistics than COSY-11)
- ❖ To fully describe the K^+K^- -FSI we need also to measure:

$$e^+e^- \rightarrow \pi^0\pi^0\gamma \text{ [EPJC49(2007)473]}$$

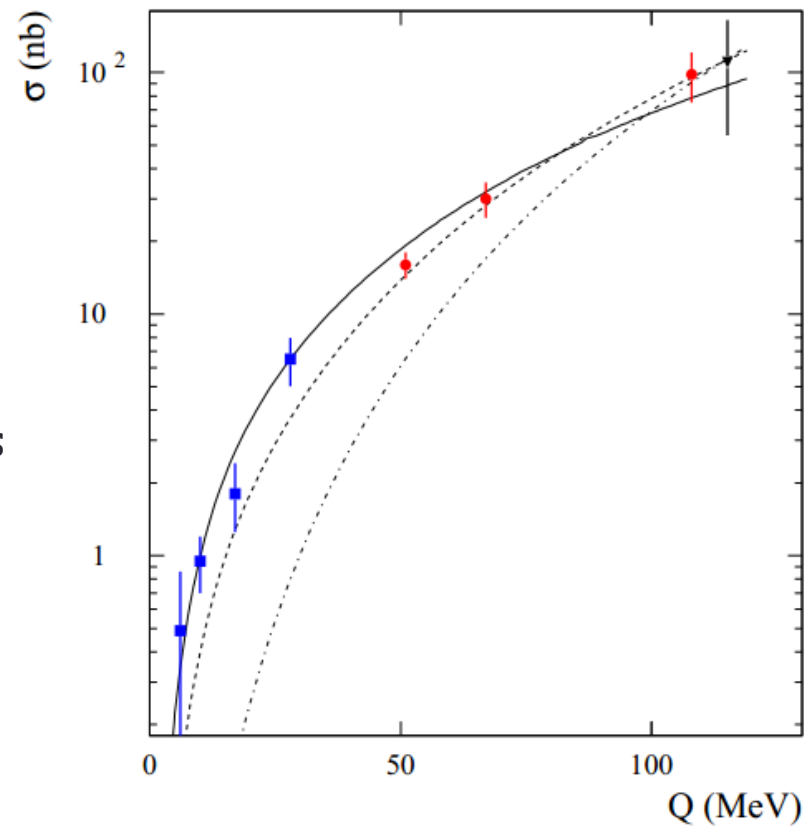
$$e^+e^- \rightarrow \pi^+\pi^-\gamma \text{ [PLB606(2005)12, PLB670(2009)285, PLB700(2011)102]}$$

$$e^+e^- \rightarrow \pi^0\eta\gamma \text{ [PLB681(2009)5]}$$

$$e^+e^- \rightarrow K_S K_S\gamma \text{ [PLB679(2009)10]}$$

Conclusions & outlook

- ❑ The excitation function for the $pp \rightarrow ppK^+K^-$ reaction reveal an enhancement which may be assigned to the influence of the pK^- and K^+K^- interaction
- ❑ The ANKE factorization ansatz underestimates experimental data very close to threshold
- ❑ The coupled channel effects and production of $f_0(980)/a_0(980)$ are up to now not distinguishable even with high statistic measurements
- ❑ We have estimated the K^+K^- scattering length based on the near threshold data independently from a_{pK^-} obtained by the ANKE group
- ❑ Rough estimates show that with KLOE-2 we could study the KK final state interaction with high precision



SPARES

$$\beta_j = \frac{L_j \alpha \sigma_j^m}{N_j^{gen}}$$

$$\sigma^m = \int \frac{\pi^2 |M|^2}{8s\sqrt{-B}} dM_{pp}^2 dM_{K^+K^-}^2 dM_{pK^-}^2 dM_{ppK^-}^2 dM_{ppK^+}^2$$

$$\left| M_{pp \rightarrow ppK^+K^-} \right|^2 \approx |M_0|^2 |F_{FSI}|^2$$

$$F_{FSI} = F_{pp}(q) \times F_{p_1K^-}(k_1) \times \\ \times F_{p_2K^-}(k_2) \times F_{K^+K^-}(k_3)$$

$$F_{K^+K^-}(k_3) = \frac{1}{1 - ik_3 a_{K^+K^-}}$$

Generalization of the Dalitz Plot

- Probability of reaction yielding a state with the i -th particle in momentum range dp_i (in CM):

$$d^{12}R = d^3 p_1 d^3 p_2 d^3 p_3 d^3 p_4 \frac{1}{16E_1 E_2 E_3 E_4} \delta^3 \left(\sum_j \vec{p}_j \right) \delta \left(\sum_j E_j - \sqrt{s} \right) f^2$$

- Assuming that f depends only on invariant masses of the particles one obtains (Nyborg et al. Phys. Rev. 140 922 (1965)):

$$d^5 R = f^2 \frac{\pi^2}{8s\sqrt{-B}} dM_{12}^2 dM_{14}^2 dM_{34}^2 dM_{124}^2 dM_{134}^2$$



$$\left| M_{pp \rightarrow ppK^+K^-} \right|^2 \approx |M_0|^2 |F_{FSI}|^2$$

$$F_{FSI} = F_{pp}(q) \times F_{p_1K^-}(k_1) \times F_{p_2K^-}(k_2)$$

$$F_{pp}(q) = \frac{e^{-i\delta_{pp}(^1S_0)} \times \sin \delta_{pp}(^1S_0)}{C \times q}$$

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