



eQCD, Sarajevo

Feb, 2013

GLUEBALLS IN THE BETHE-SALPETER FORMALISM





LATTICE RESULTS

Y. Chen et al., Phys. Rev. D73, 014516 (2006).



EXPERIMENT

 $f_0(1500) \ IJ^{PC} = 00^{++}$





EXPERIMENT



J. Meyers and E.S. Swanson, PRD [arxiv:1211.4648]

R. Alkofer and L. von Smekal, Phys. Rept. 353, 281 (2001).

















required elements:









gluon propagator:

$$D_{\mu\nu}(k) \equiv -i\left(g_{\mu\nu} - \frac{k_{\mu}k_{\nu}}{k^2}\right) \cdot G(k) - i\xi \frac{k^{\mu}k^{\nu}}{k^4 + i\epsilon}$$



I.L. Bogolubsky, E.M. Ilgenfritz, M. Muller-Preussker and A. Sternbeck, Phys. Lett. **B676**, 69 (2009)

vertices:

various models taken from

A.C. Aguilar and J. Papavassiliou, Phys. Rev. D83, 014013 (2011)L. von Smekal, A. Hauck, and R. Alkofer, Ann. Phys. 267, 1 (1998)M.R. Pennington and D.J. Wilson, Phys. Rev. D84, 119901 (2011)

attempt to take into account Slavnov-Taylor identities, multiplicative renormalisability, etc

results: ghost



I.L. Bogolubsky, E.M. Ilgenfritz, M. Muller-Preussker and A. Sternbeck, Phys. Lett. **B676**, 69 (2009)

results: quark



$f_{\pi} = 240 \text{ MeV}$ $\langle \bar{\psi}\psi \rangle (1 \text{ GeV}) = (-251 \text{ MeV})^3$

$$S(k) = i \frac{Z(k)}{\not k - M(k)}$$

P.O. Bowman et al., Phys. Rev. D71, 054507 (2005).

GLUEBALL EQUATIONS

$$\begin{split} \chi_{\mu\nu}(k_{+},k_{-}) &= ig^{2}N \int \frac{d^{4}q}{(2\pi)^{4}} \chi_{\alpha\beta}(q_{+},q_{-}) \mathcal{C}^{\alpha\beta}_{..\mu\nu} G(q_{+})G(q_{-}) \\ &+ ig^{2}N \int \frac{d^{4}q}{(2\pi)^{4}} \chi_{\alpha\beta}(q_{+},q_{-}) \mathcal{T}^{\alpha\beta}_{..\mu\nu}(q,k;P) G(q_{+})G(q_{-})G(Q) \\ &+ ig^{2}N \int \frac{d^{4}q}{(2\pi)^{4}} \chi(q_{+},q_{-}) \mathcal{G}_{\mu\nu}(q,k;P) H(q_{+})H(q_{-})H(Q) \\ &- \frac{g^{2}}{2} \int \frac{d^{4}q}{(2\pi)^{4}} \operatorname{tr} [\gamma_{\mu}S(q_{+})\chi(q_{+},q_{-})S(-q_{-})\gamma_{\nu}S(Q)] h^{n}(k_{+})h^{n}(k_{-})\bar{A}(Q,q_{+})\bar{A}(Q,q_{-}) \\ &+ \operatorname{crossed} \end{split}$$
(1)
$$\chi(k_{+},k_{-}) &= ig^{2}N \int \frac{d^{4}q}{(2\pi)^{4}} \chi(q_{+},q_{-}) \mathcal{H}(q,k;P) H(q_{+})H(q_{-})G(Q) \\ &+ ig^{2}N \int \frac{d^{4}q}{(2\pi)^{4}} \chi_{\alpha\beta}(q_{+},q_{-}) \mathcal{B}^{\alpha\beta}(q,k;P) G(q_{+})G(q_{-})H(Q) \end{aligned}$$
(2)
$$\chi(k_{+},k_{-}) &= g^{2}C_{F} \int \frac{d^{4}q}{(2\pi)^{4}} \gamma_{\alpha}S(-Q)\gamma_{\beta}G(q_{+})G(q_{-})\chi_{\alpha\beta}(q_{+},q_{-})h^{n}(q_{+})h^{n}(q_{-})A(Q,k_{+})\bar{A}(Q,k_{-}) \\ &+ ig^{2}C_{F} \int \frac{d^{4}q}{(2\pi)^{4}} \gamma_{\mu}S(q_{+})\chi(q_{+},q_{-})S(-q_{-})\gamma_{\nu}G(Q)P_{\mu\nu}(Q) h^{2n}(Q) \bar{A}(k_{+},q_{+}) \bar{A}(k_{-},q(3)) \end{split}$$

$$\mathcal{C}^{\alpha\beta\mu\nu} = 2g^{\mu\nu}g^{\alpha\beta} - g^{\alpha\nu}g^{\beta\mu} - g^{\alpha\mu}g^{\beta\nu}$$

$$\mathcal{T}_{\alpha\beta\mu\nu} = P^{\gamma\gamma'}(Q) V_{\alpha\mu\gamma}(q_+, k_+) V_{\beta\nu\gamma'}(q_-, k_-)$$

$$P^{\gamma\gamma'}(Q) = g^{\gamma\gamma'} - \frac{Q^{\gamma}Q^{\gamma}}{Q^2}$$

 $V_{\alpha\mu\gamma}(q_{+},k_{+}) = A_{+}(q_{+},k_{+},q_{+}-k_{+})g^{\alpha\mu}(q_{+}+k_{+})^{\gamma} + A_{+}(k_{+},q_{+}-k_{+},q_{+})g^{\mu\gamma}(q_{+}-2k_{+})^{\alpha} + A_{+}(q_{+}-k_{+},q_{+},k_{+})g^{\gamma\alpha}(k_{+}-2q_{+})^{\mu}.$

$$\mathcal{G}^{\mu\nu} = \left(q_{+}^{\mu}\frac{h(k_{+})}{h(q_{+})} + Q^{\mu}\left(\frac{h(k_{+})}{h(Q)} - 1\right)\right)\left(Q^{\nu}\frac{h(k_{-})}{h(Q)} - q_{-}^{\nu}\left(\frac{h(k_{-})}{h(q_{-})} - 1\right)\right)$$

$$\mathcal{H} = P_{\mu\nu}(Q)k_{+}^{\mu}k_{-}^{\nu} \cdot \left(\frac{h(Q)}{h(k_{+})} + \frac{h(Q)}{h(q_{+})} - 1\right) \cdot \left(\frac{h(Q)}{h(k_{-})} + \frac{h(Q)}{h(q_{-})} - 1\right)$$

$$\mathcal{B}^{\alpha\beta} = -k_{+}^{\alpha}k_{-}^{\beta} \cdot \left(\frac{h(q_{+})}{h(Q)} + \frac{h(q_{+})}{h(k_{+})} - 1\right) \cdot \left(\frac{h(q_{-})}{h(Q)} + \frac{h(q_{-})}{h(k_{-})} - 1\right)$$

impose transversality and parity

$$\chi_{R}^{\mu\nu}(k_{+},k_{-}) = \epsilon(\mu,\nu) \cdot \Pi_{R} \cdot \chi_{R}^{\mu\nu}(\tilde{k}_{+},\tilde{k}_{-})$$

 $\chi^{\mu\nu}(k_+,k_-) = A_0 g^{\mu\nu} + A_1 k_+^{\mu} k_+^{\nu} + A_2 k_-^{\mu} k_-^{\nu} + A_3 k_+^{\mu} k_-^{\nu} + A_4 k_-^{\mu} k_+^{\nu} + A_5 \epsilon^{\mu\nu\alpha\beta} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{+\alpha} k_{-\beta} k_{+\alpha} k_{+\alpha} k_{+\alpha} k_{-\beta} k_{+\alpha} k$

$$\chi^{\mu\nu}(k_{+},k_{-}) = \epsilon^{\mu\nu\alpha\beta}k_{\alpha}P_{\beta}F(k,P) \quad [0^{-+}]$$

$$\chi^{\mu\nu}(k_{+},k_{-}) = A(k,P)A^{\mu\nu} + B(k,P)B^{\mu\nu} \quad [0^{++}]$$





Wilson loop



the Bethe-Salpeter equation simplifies in the heavy quark limit

$$V \to -iG(-\vec{p}^2)$$



[screened Coulomb with log corrections]

- kernel structure
- instantons
- reggeons
- flux tubes
- N-ality confinement
- pomeron



It is methodology is in its infancy -- it will likely require sophisticated modelling

 the importance of topological vacuum effects and multigluon dynamics may be revealed.

PANDA@GSI, JLab, BESIII, superB, LHCb provide hope for progress

+ ÆRIC MEC HEHT GEWYRCAN

