

Double parton distributions and their evolution

Krzysztof Golec-Biernat

(in collaboration with Emilia Lewandowska)

Institute of Nuclear Physics, Krakow, Poland

Low X Meeting, Paphos, 27 June - 1 July 2012

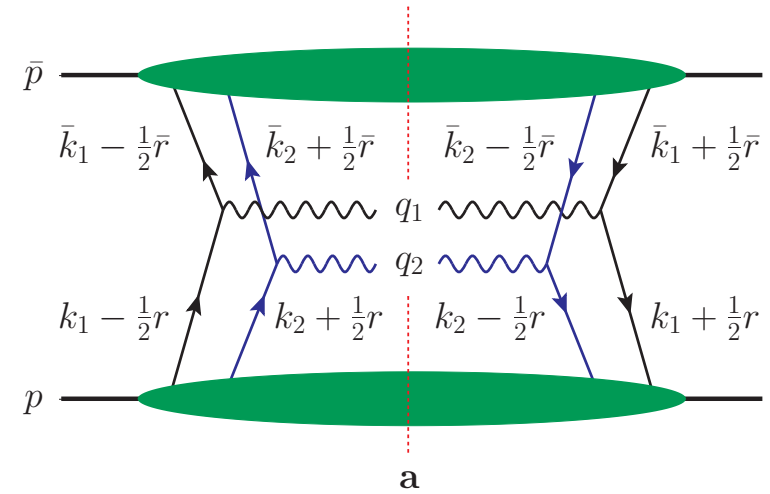
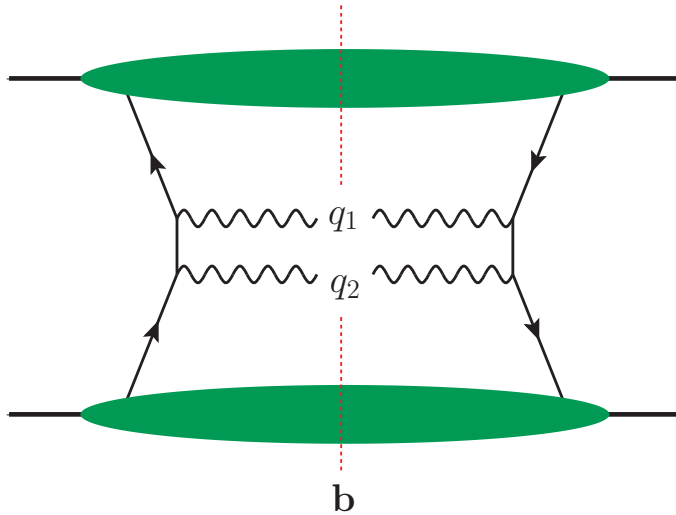
Motivation

- Growing interest in multiparton scattering in hadronic collisions
- Two parton scattering: Snigirev, Gaunt, Stirling, Ceccopieri, Ryskin, Bartels, Diehl, A. Schaefer, Dokshitzer, Frankfurt, Strikman, Blok, Kulesza, Szczurek,.....
- Double parton distributions in the cross section, e.g.,

$$\sigma_{pp \rightarrow WWX} = D(x_1, x_2) \otimes \hat{\sigma}_{qq \rightarrow W} \otimes \hat{\sigma}_{qq \rightarrow W} \otimes D(x_1, x_2)$$

Single PDF versus double PDF

- Two vector bosons from two and four partons scattering.



- One and two hard scatterings with single and double PDF:

$$D_f(x, \mu),$$

$$D_{f_1 f_2}(x_1, x_2, \mu_1, \mu_2)$$

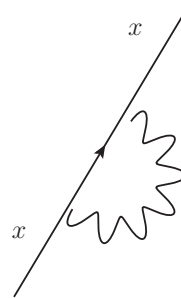
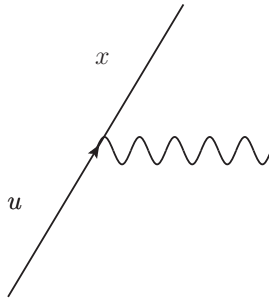
- Obvious condition: $0 < x_1 + x_2 \leq 1$.

- Ansatz $D_{f_1 f_2}(x_1, x_2, \mu_1, \mu_2) = D_{f_1}(x_1, \mu_1) \cdot D_{f_2}(x_2, \mu_2)$ is not general.

Evolution equations for single PDF

- Markov process with time $t = \ln \mu$ and $f = q_i, \bar{q}_i, g$

$$D_f(x, t + \delta t) = \sum_{f'} \int_x^1 du \underbrace{\delta t P_{ff'}^r(x, u) D_{f'}(u, t)}_{\text{real emission}} + \underbrace{(1 - \delta t P_f^v(x)) D_f(x, t)}_{\text{virtual emission}}$$



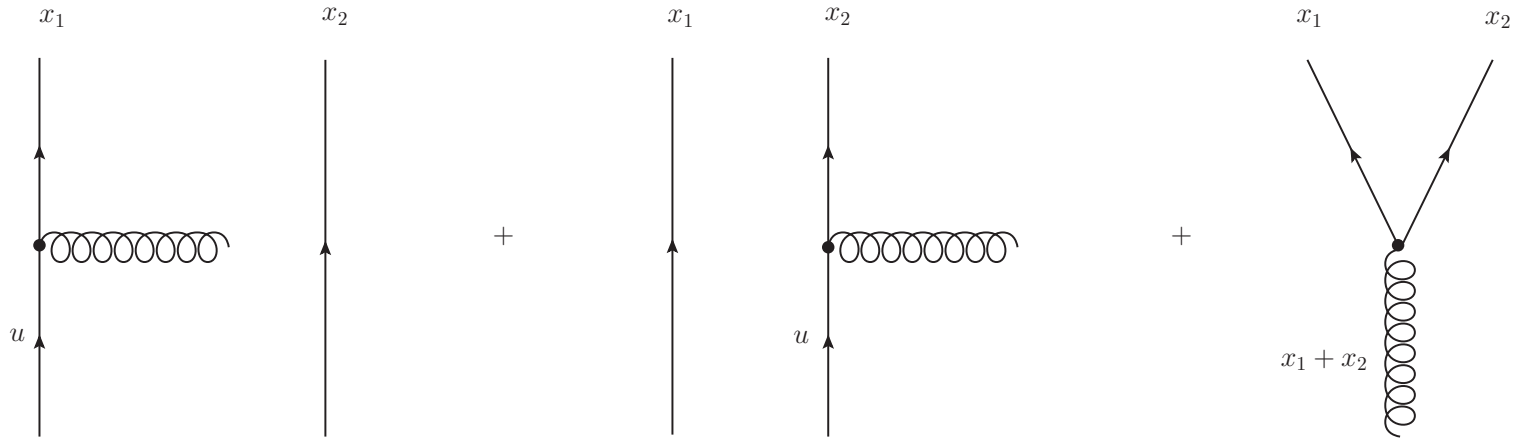
- In the limit $\delta t \rightarrow 0$, a set of linear differential equations (DGLAP)

$$\partial_t D_f(x, t) = \sum_{f'} \int_x^1 du P_{ff'}^r(x, u) D_{f'}(u, t) - P_f^v(x) D_f(x, t)$$

- Virtual $P_f^v(x)$ from momentum sum rule: $\int_0^1 dx \sum_f x D_f(x, t) = 1.$

Evolution equations for dPDF (LLA)

● Pictorially: $D_{f_1 f_2}(x_1, x_2, \mu, \mu), \quad t = \ln \mu$



● Analytically:

$$\begin{aligned}
 \partial_t D_{f_1 f_2}(x_1, x_2, t) &= \sum_{f'} \int_0^{1-x_2} du P_{f_1 f'}(x_1, u) D_{f' f_2}(u, x_2, t) \\
 &+ \sum_{f'} \int_0^{1-x_1} du P_{f_2 f'}(x_2, u) D_{f_1 f'}(x_1, u, t) \\
 &+ \sum_{f'} P_{f' \rightarrow f_1 f_2} \left(\frac{x_1}{x_1 + x_2} \right) \underbrace{D_{f'}(x_1 + x_2, t)}_{sPDF}
 \end{aligned}$$

Evolution equations for dPDF

- For different scales: $t_1 = \ln \mu_1 < t_2 = \ln \mu_2$

$$D_{f_1 f_2}(x_1, x_2, t_0, t_0) \rightarrow D_{f_1 f_2}(x_1, x_2, t_1, t_1) \rightarrow D_{f_1 f_2}(x_1 = \textit{fixed}, x_2, t_1, t_2)$$

- sPDF evolution for x_2 from $t_1 \rightarrow t_2$ with fixed x_1 :

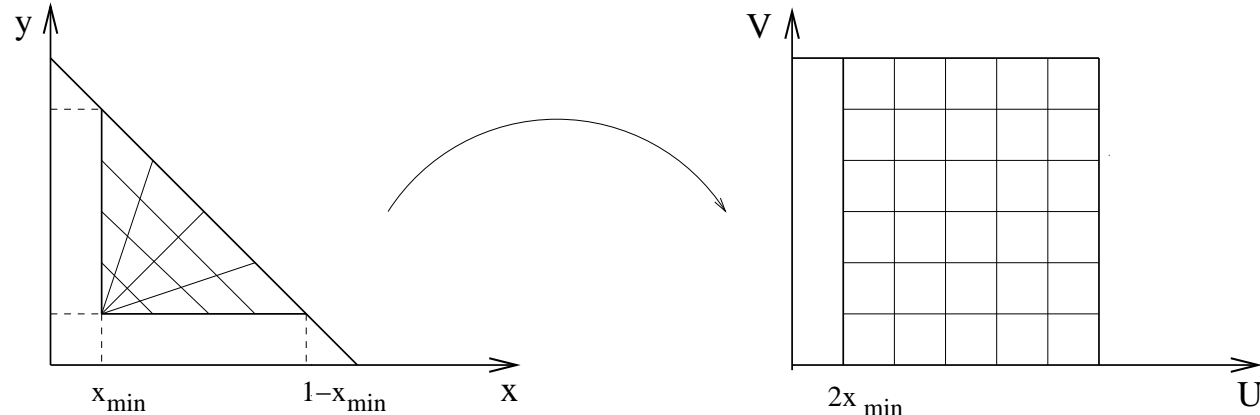
$$\partial_{t_2} D_{f_1 f_2}(x_1, x_2, t_1, t_2) = \sum_{f'} \int_0^{1-x_1} du P_{f_2 f'}(x_2, u) D_{f_1 f_2}(x_1, u, t_1, t_2)$$

- No parton exchange symmetry:

$$D_{f_1 f_2}(x_1, x_2, t_1, t_2) \neq D_{f_2 f_1}(x_2, x_1, t_1, t_2)$$

Finding solutions

- Chebyshev polynomial approximation method on a grid of points in the (x_1, x_2) plane.

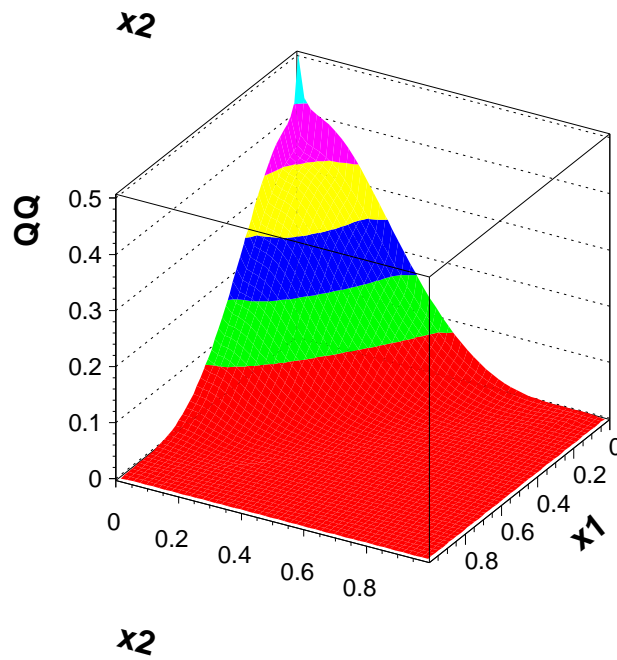
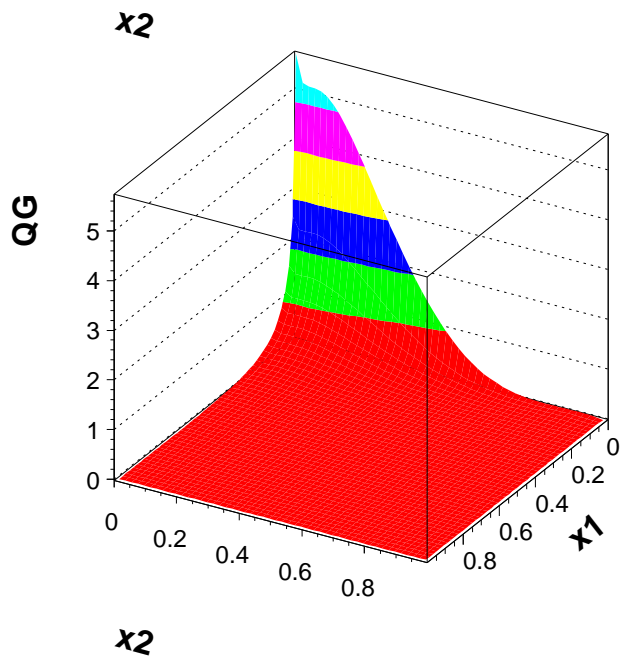
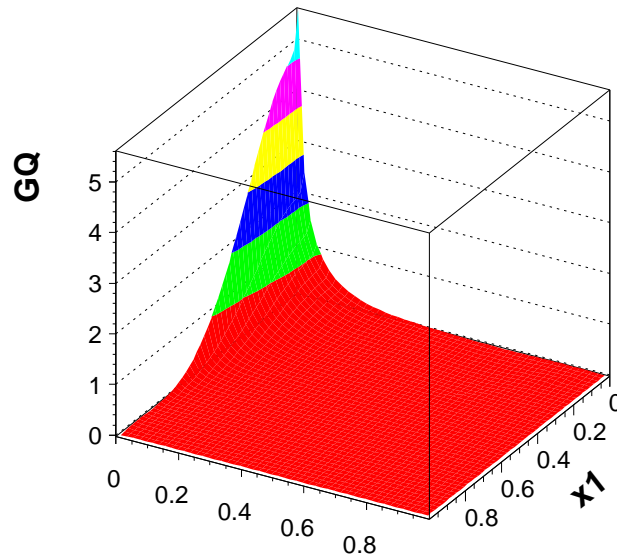
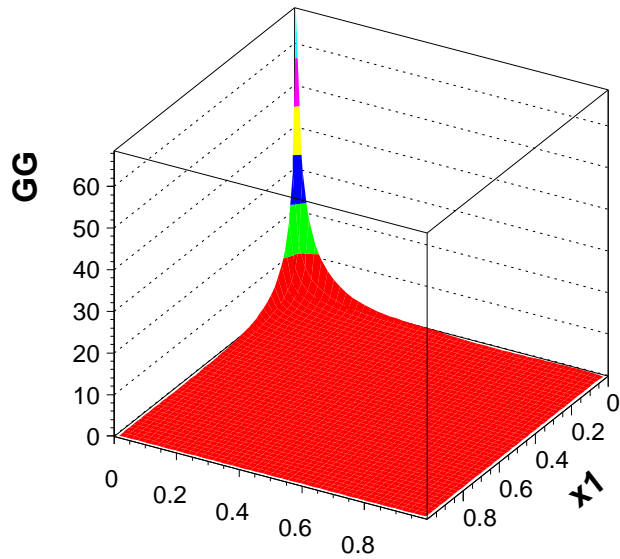


- Set of first order differential equations for dPDF at grid points

$$\frac{d}{dt} D_{nm}(t) = \sum_{k,l} A_{nm|kl} D_{kl}(t) + \sum_k B_{nmk} D_k(t)$$

- Problems: stability, agreement with carefully chosen tests,...

An example



Momentum sum rule

- Conditional probability with known second parton:

$$\sum_{f_1} \int_0^{1-x_2} dx_1 x_1 \frac{D_{f_1 f_2}(x_1, x_2)}{D_{f_2}(x_2)} = (1 - x_2)$$

- Relation conserved by evolution equations once imposed at initial t_0

$$\sum_{f_1} \int_0^{1-x_2} dx_1 x_1 D_{f_1 f_2}(x_1, x_2, t) = (1 - x_2) D_{f_2}(x_2, t)$$

- Sum rule for dPDF after using sPDF momentum sum rule:

$$\sum_{f_1, f_2} \int_0^1 dx_2 \int_0^{1-x_2} dx_1 \frac{x_1 x_2}{1 - x_2} D_{f_1 f_2}(x_1, x_2, t) = 1$$

How to fulfill momentum sum rule?

- Initial conditions (Korotkin, Snigirev, Gaunt, Stirling):

$$D_{f_1 f_2}(x_1, x_2) = D_{f_1}(x_1) D_{f_2}(x_2) \frac{(1 - x_1 - x_2)^n}{(1 - x_1)^{n_1} (1 - x_2)^{n_2}}$$

- Our initial conditions (no parton exchange symmetry):

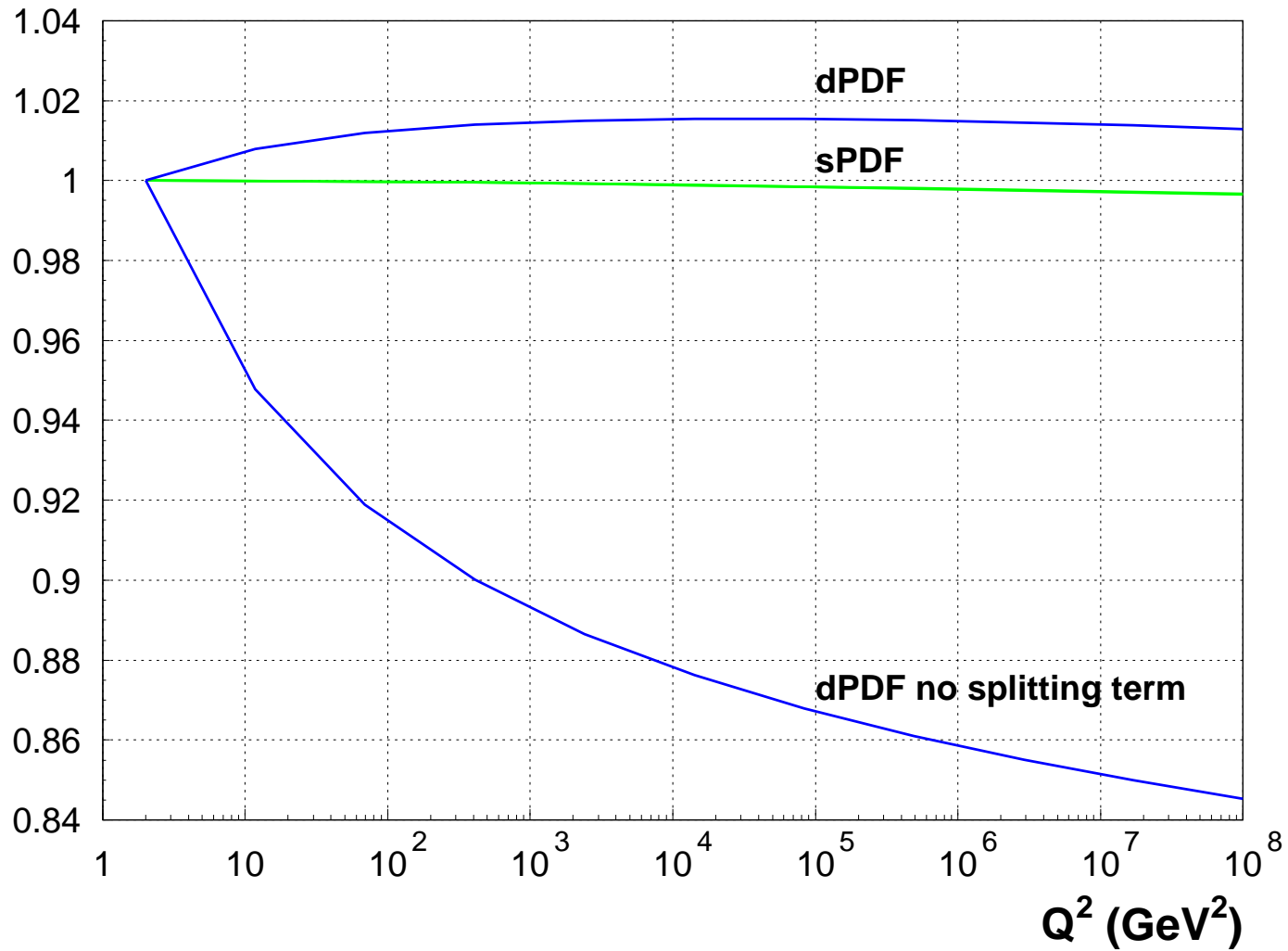
$$D_{f_1 f_2}(x_1, x_2) = \frac{1}{1 - x_2} D_{f_1} \left(\frac{x_1}{1 - x_2} \right) D_{f_2}(x_2)$$

- Satisfies also valence number quark rule

$$\int_0^{1-x_2} dx_1 \{ D_{f_1 f_2}(x_1, x_2) - D_{\bar{f}_1 f_2}(x_1, x_2) \} = N_{f_1} D_{f_2}(x_2)$$

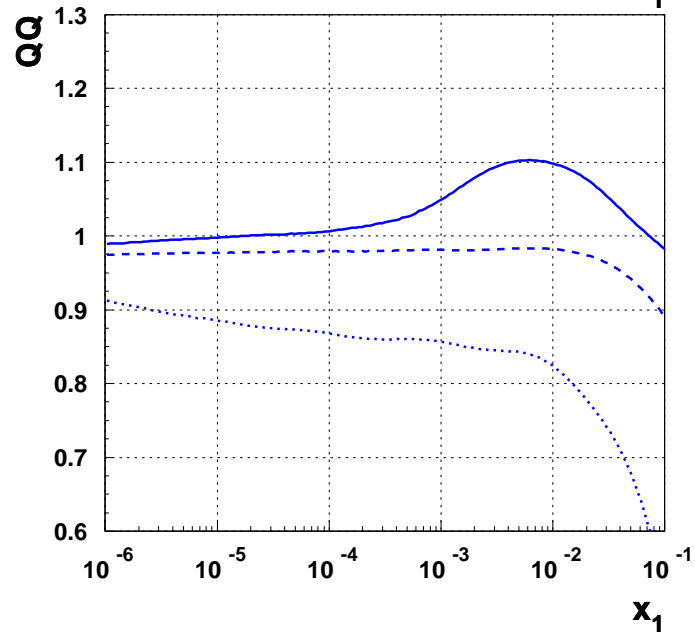
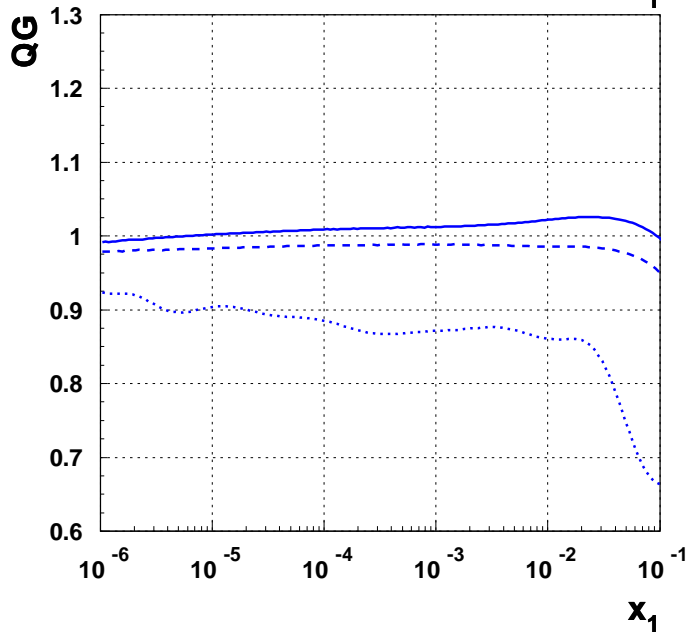
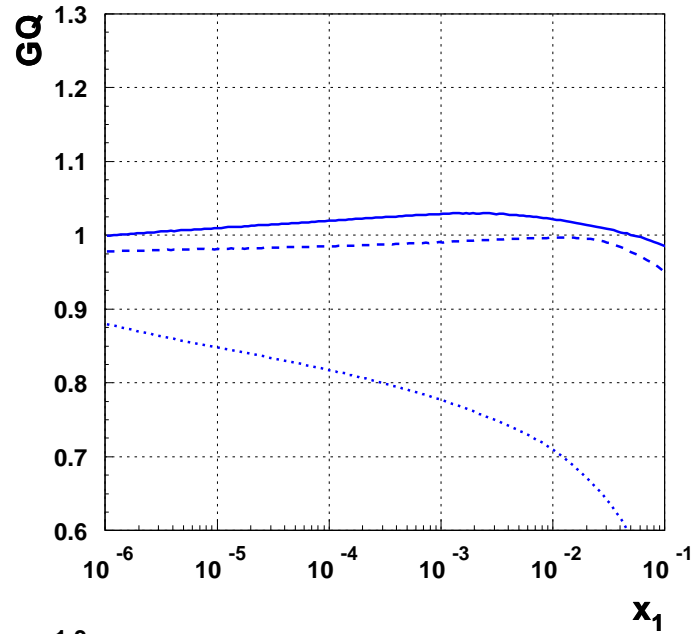
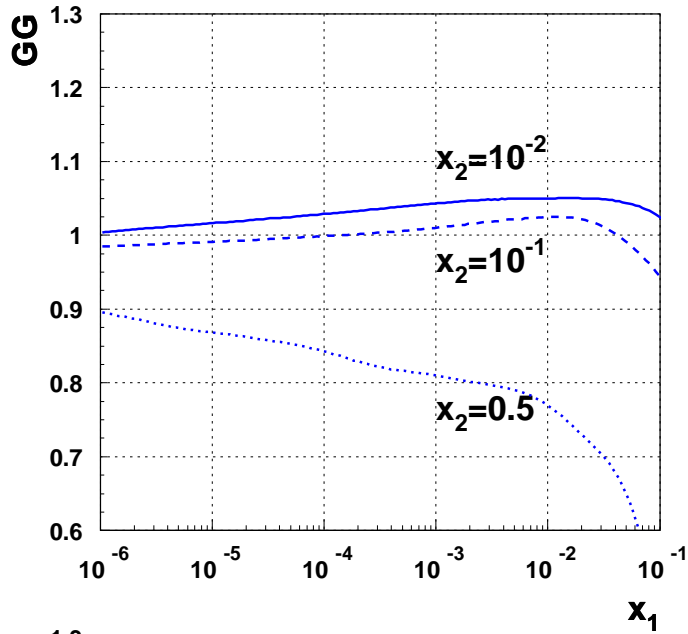
with $N_u = 2$ and $N_d = 1$ for the proton.

Momentum sum rule in practice



$D_{f_1 f_2}(x_1, x_2)$ versus $D_{f_1}(x_1)D_{f_2}(x_2)$

ratio for $Q^2=10^2$



Summary

- Sum rules play important role in dPDF evolution LLA.
- For small x_1, x_2 , a naive product of sPDF not so bad, but ...
- ...remember about higher order corrections, parton recombination effects, transverse momentum dependend dPDF,...
- We are still at the beginnig of the understanding of multiparton scattering.