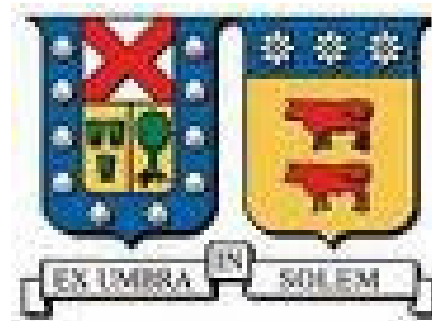


Nuclear modification factor for gluon jets

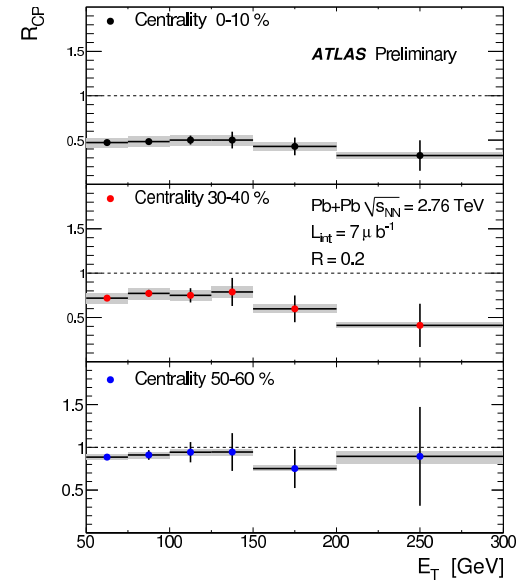
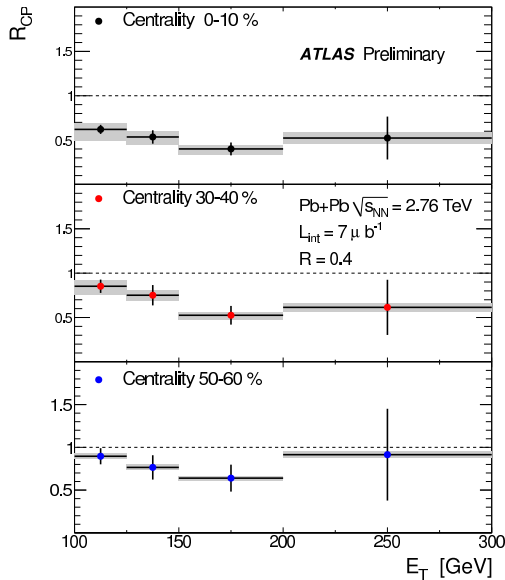
Eugene Levin, UTFSM / TAU



June 27- July 1, 2012 ,Paphos, Cyprus

E. L.: “Nuclear modification factor for gluon jets,” *arXiv:1203.4467 [hep-ph]*)

ATLAS data.



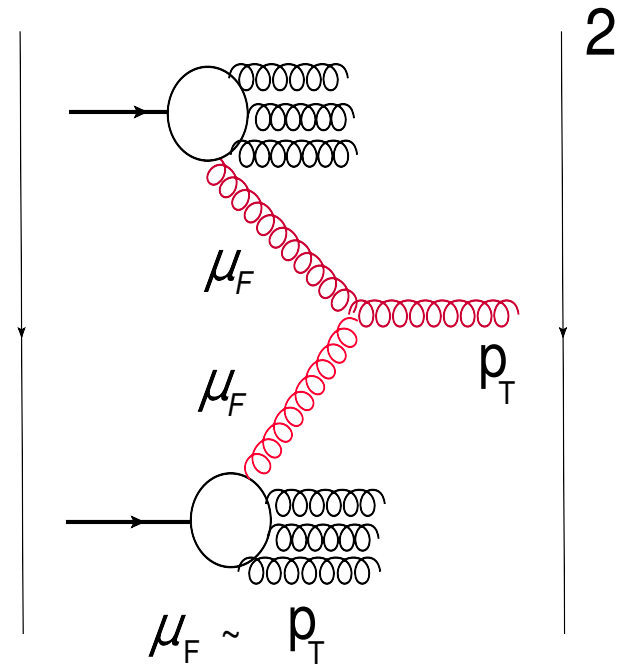
Two theoretical questions:

1. Could be NMF < 1 at low x and large $p_T \gg Q_s$ without violation of the factorization theorem ?
2. Can be the value of NMF of the order of the value measured by ATLAS?

QCD factorization.

- $$\frac{d^2\sigma(Y', p_T)}{dY' d^2p_T} = \sigma_{hard} \otimes x_1 G_{Pr}(x_1, \mu_F^2) x_2 G_{Tr}(x_2, \mu_F^2)$$
- $$xG_A(x, \mu_F^2) = AxG_P(x, \mu_F^2)$$
- $$\text{NMF} = \frac{d^2\sigma(AA)(Y', p_T)}{dY' d^2p_T} / A^2 \frac{d^2\sigma(pp)(Y', p_T)}{dY' d^2p_T}$$

$$= (xG_A(x, \mu_F^2))^2 / (AxG_p(x, \mu_F^2))^2 = 1$$



Violation of QCD factorization ?!

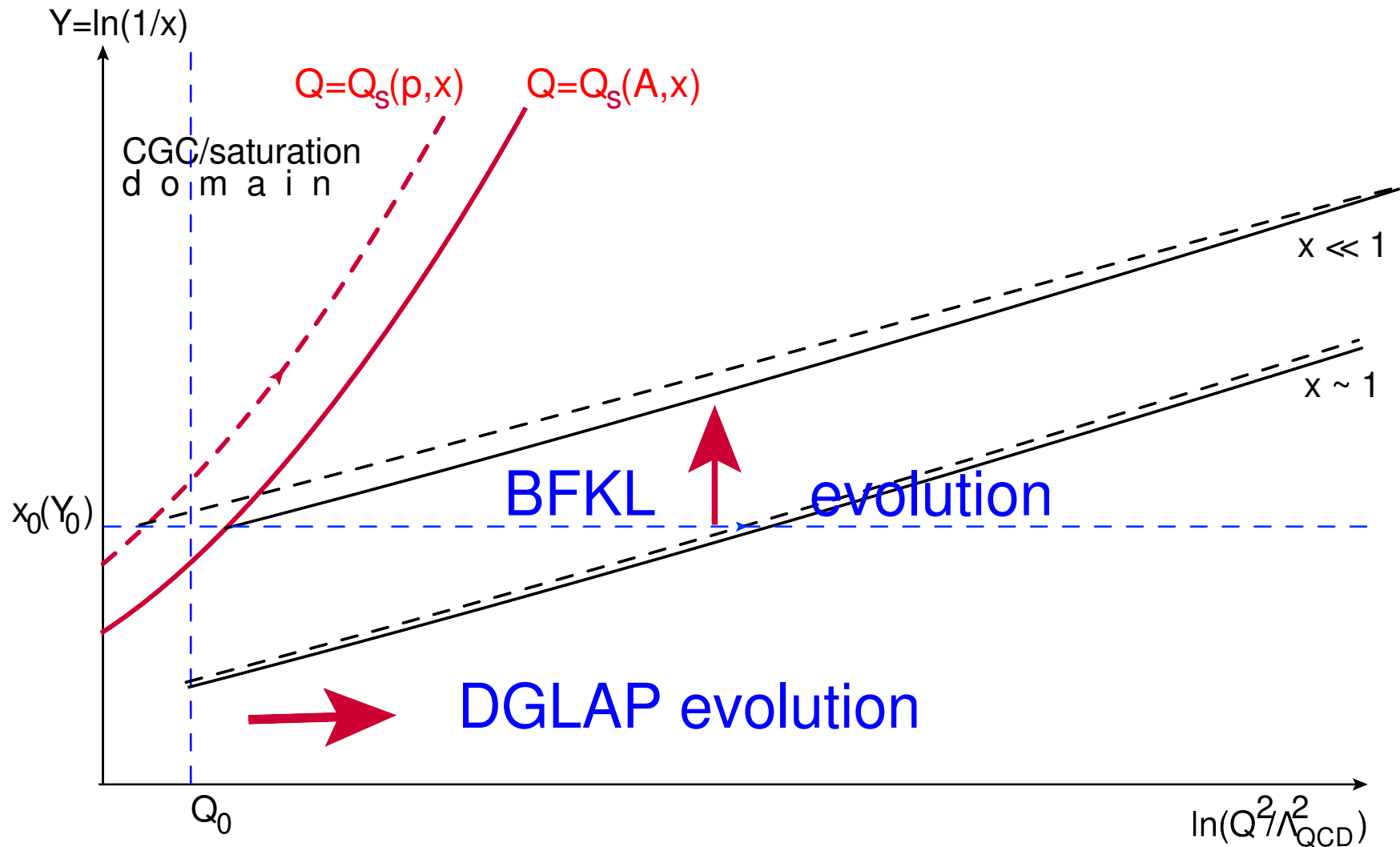
No violation of QCD factorization, but

- $xG_A(x, p_T^2/Q_s^2(A; x_0)) = AxG_P(x, p_T^2/Q_s^2(A; x_0))$
 $= AxG_P(x, p_T^2/(A^{1/3}Q_s^2(P; x_0)))$ for $p_T \gg Q_s(A; x)$

- $R_{AA} = \frac{x_1 G_P(x_1, p_T^2/Q_s^2(A; x_{10})) x_2 G_P(x_2, p_T^2/Q_s^2(A; x_{20}))}{x_1 G_P(x_1, p_T^2/Q_s^2(P; x_{10})) x_2 G_P(x_2, p_T^2/Q_s^2(P; x_{20}))} < 1$

Linear equations at large p_T

Evolution in $\ln Q^2$ (DGLAP evolution):



Evolution in $\ln(1/x)$ (BFKL-type evolution):

For low x both BFKL and DGLAP evolution equations have the form:

$$\bullet \quad \frac{\partial xG(Y; p_T^2)}{\partial Y} = \int d^2k_T K(p_T, k_T) xG(Y, k_T^2)$$

Solutions have the form

$$xG\left(Y, \xi \equiv \ln\left(p_T^2/\Lambda_{QCD}^2\right)\right) = \int d\xi' G(Y - Y_0, \xi - \xi') xG_{in}(Y = Y_0, \xi')$$

- $G(Y - Y_0, \xi - \xi') \implies$ Green's function,
- $G(Y = Y_0, \xi - \xi') = \delta(\xi - \xi')$.
- $xG_{in}(Y = Y_0, \xi')$ is the initial condition for xG

General property of the solutions:

- $xG\left(Y, \xi \equiv \ln\left(p_T^2/\Lambda_{QCD}^2\right)\right) = \int d\xi' G(Y - Y', \xi - \xi') xG(Y', \xi')$

Indeed, the general solution has the form

- $xG(Y, \xi) = \int_{\epsilon-i\infty}^{\epsilon+i\infty} \frac{d\gamma}{2\pi i} e^{\omega(\gamma)Y + \gamma\xi} g(\gamma) = \int_{\epsilon-i\infty}^{\epsilon+i\infty} \frac{d\omega}{2\pi i} e^{\omega Y + \gamma(\omega)\xi} g(\omega)$

- $\omega = \omega(\gamma) \longrightarrow$ Mellin image of K ; $\gamma = \gamma(\omega) \longrightarrow$ solution of $\omega = \omega(\gamma)$

- $G(Y - Y', \xi) = \int_{\epsilon-i\infty}^{\epsilon+i\infty} \frac{d\gamma}{2\pi i} e^{\omega(\gamma)(Y - Y') + \gamma(\xi - \xi')}$

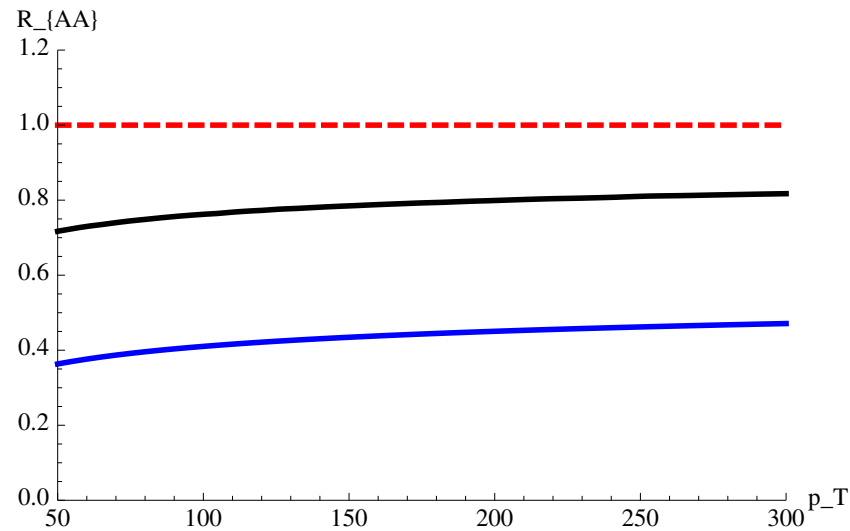
Assume: $x_0 G(Y_0, k_T) = S Q_s^2(Y_0) g(\tau = k_T^2 / Q_s^2(Y_0))$

$$\begin{aligned}
 & \bullet xG\left(Y, \xi \equiv \ln\left(p_T^2 / \Lambda_{QCD}^2\right)\right) \\
 &= \int \frac{dk_T^2}{k_T^2} G(Y - Y_0, \xi - \xi') xG\left(Y_0, k_T^2\right) \\
 &= S Q_s^2(Y_0) \int \frac{dk^2}{k_T^2} G\left(Y - Y_0, p_T^2 / k_T^2\right) g\left(k_T^2 / Q_s^2(Y_0)\right) \\
 &= S Q_s^2(Y_0) \int_{\epsilon - i\infty}^{\epsilon + i\infty} \frac{d\gamma}{2\pi i} e^{\omega(\gamma)(Y - Y_0)} \left(\frac{p_T^2}{k_T^2}\right)^\gamma g\left(\frac{k_T^2}{Q_s^2(Y_0)}\right) \\
 &= S Q_s^2(Y_0) \int_{\epsilon - i\infty}^{\epsilon + i\infty} \frac{d\gamma}{2\pi i} e^{\omega(\gamma)(Y - Y_0)} \left(\frac{p_t^2}{Q_s^2(Y_0)}\right)^\gamma \int \frac{d\tau}{\tau} \tau^{-\gamma} g(\tau) \\
 &= S Q_s^2(Y_0) \int_{\epsilon - i\infty}^{\epsilon + i\infty} \frac{d\gamma}{2\pi i} e^{\omega(\gamma)(Y - Y_0)} \left(\frac{p_t^2}{Q_s^2(Y_0)}\right)^\gamma f(\gamma)
 \end{aligned}$$

$$\text{Proof: } x_0 G(Y_0, k_T) = S Q_s^2(Y_0) g(\tau = k_T^2 / Q_s^2(Y_0))$$

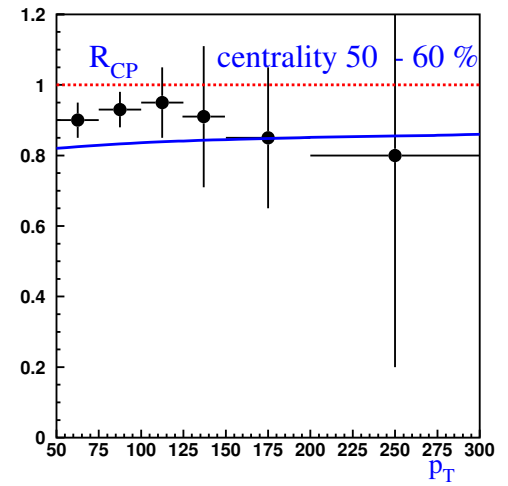
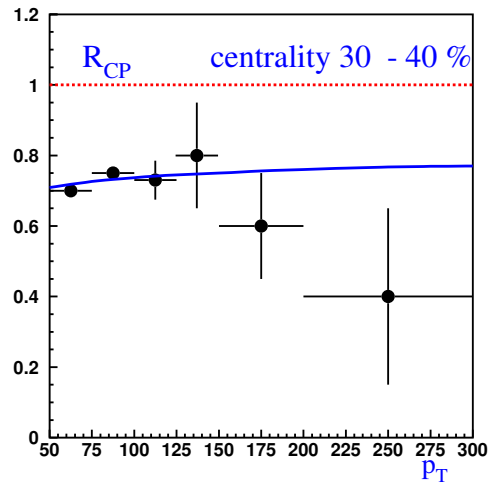
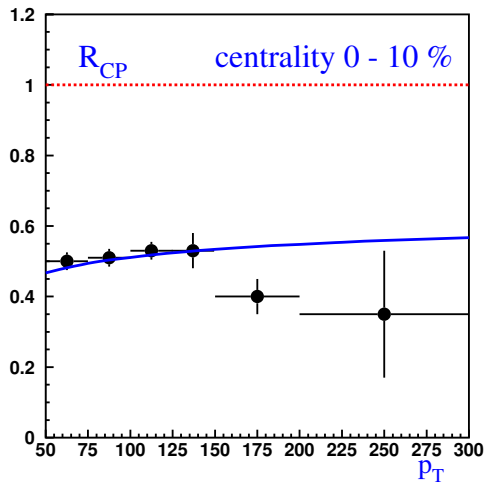
- $x G(x, p_T^2) = \int^{p_T^2} dk^2 \phi(Y, k^2)$
- $\phi_G^A(x_0; \vec{k}_T) = \frac{1}{\bar{\alpha}_S 4\pi} \int d^2b d^2r e^{i\vec{k}_T \cdot \vec{r}} \nabla_\perp^2 N_G^A(Y_0; r; b)$
- $N_G^A(Y_0; r, b) = 2 N(Y_0; r, b) - N^2(Y_0; r, b)$
 $\rightarrow 1 - \exp(-r^2 Q_s^2(Y_0; b) / 2)$
- for $N(Y_0; r, b) = 1 - \exp(-r^2 Q_s^2(Y_0; b) / 4)$ (McLerran-Venugopalan)
- $x_0 G(x_0, p_T^2) = \frac{\pi}{2} S Q_s^2(Y_0) \left(1 - (1 + 2\tau) \exp(-2\tau) \right)$
 $= S Q_s^2(Y_0) g(\tau)$

Estimates for NMF

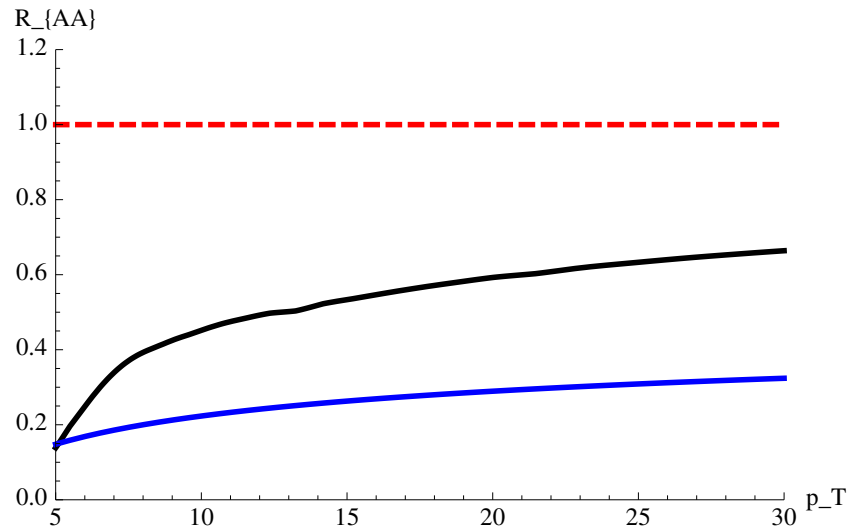


- $\gamma = \bar{\alpha}_S \left(\frac{1}{\omega} - 1 \right)$ or $\omega(\gamma) = \frac{\alpha_S}{\gamma + \bar{\alpha}_S}$
- $G(Y, \tau) = f(\gamma_{SP}) \exp \left(2 \sqrt{\bar{\alpha}_S Y \ln(\tau)} - \bar{\alpha}_S \ln(\tau) \right)$

with $\gamma_{SP} = \sqrt{\frac{\bar{\alpha}_S Y}{\ln(\tau)}}$



● $x = 10^{-3}$!!!



Conclusions

1. Could be $\text{NMF} < 1$ at low x and large $p_T \gg Q_s$ without violation of the factorization theorem ?

2. Can be the value of NMF of the order of the value measured by ATLAS?

YES