

Unpolarized azimuthal asymmetries in SIDIS: experimental overview

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Structure and Spectroscopy

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SIDIS x-section

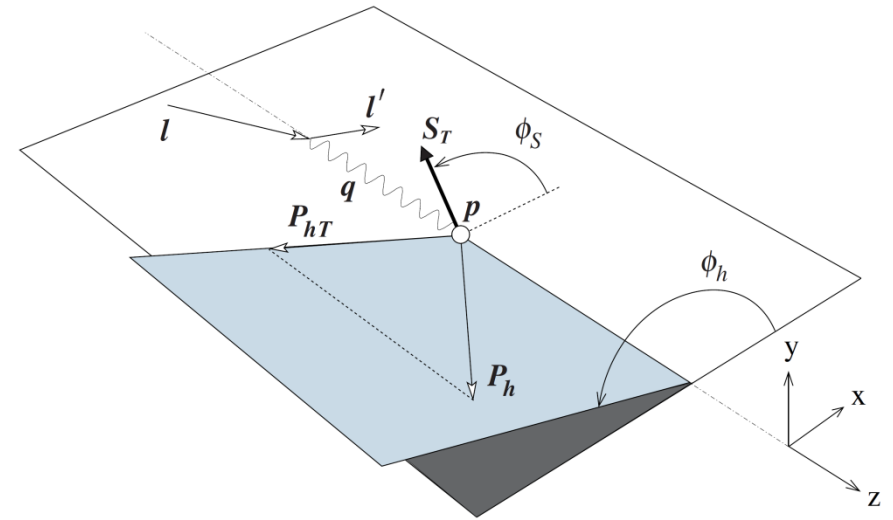
A. Kotzinian, Nucl. Phys. B441, 234 (1995). Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007)

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$A_{U(L),T}^{w(\varphi_h, \varphi_s)} = \frac{F_{U(L),T}^{w(\varphi_h, \varphi_s)}}{F_{UU,T} + \varepsilon F_{UU,L}}$$

$$\varepsilon = \frac{1 - y - \frac{1}{4} \gamma^2 y^2}{1 - y + \frac{1}{2} y^2 + \frac{1}{4} \gamma^2 y^2}, \quad \gamma = \frac{2Mx}{Q}$$

$$\left\{ \begin{array}{l} 1 + \cos \varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)} + \\ \lambda \sin \varphi_h \times \sqrt{2\varepsilon(1-\varepsilon)} A_{LU}^{\sin \varphi_h} + \\ S_L \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \varphi_h A_{UL}^{\sin \varphi_h} + \varepsilon \sin(2\varphi_h) A_{UL}^{\sin(2\varphi_h)} \right] + \\ S_L \lambda \left[\sqrt{1-\varepsilon^2} A_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \varphi_h A_{LL}^{\cos \varphi_h} \right] + \\ \left. \begin{array}{l} S_T \left[\begin{array}{l} \sin \varphi_s \times \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin \varphi_s} \right) + \\ \sin(\varphi_h - \varphi_s) \times \left(A_{UT}^{\sin(\varphi_h - \varphi_s)} \right) + \\ \sin(\varphi_h + \varphi_s) \times \left(\varepsilon A_{UT}^{\sin(\varphi_h + \varphi_s)} \right) + \\ \sin(2\varphi_h - \varphi_s) \times \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin(2\varphi_h - \varphi_s)} \right) + \\ \sin(3\varphi_h - \varphi_s) \times \left(\varepsilon A_{UT}^{\sin(3\varphi_h - \varphi_s)} \right) \end{array} \right] + \\ S_T \lambda \left[\begin{array}{l} \cos \varphi_s \times \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos \varphi_s} \right) + \\ \cos(\varphi_h - \varphi_s) \times \left(\sqrt{1-\varepsilon^2} A_{LT}^{\cos(\varphi_h - \varphi_s)} \right) + \\ \cos(2\varphi_h - \varphi_s) \times \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos(2\varphi_h - \varphi_s)} \right) \end{array} \right] \end{array} \right\}$$



15 amplitudes:

2-''UU'', 1-''LU'', 2-''UL'', 2-''LL'', 5-''UT'', 3-''LT''

SIDIS x-section

A. Kotzinian, Nucl. Phys. B441, 234 (1995). Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007)

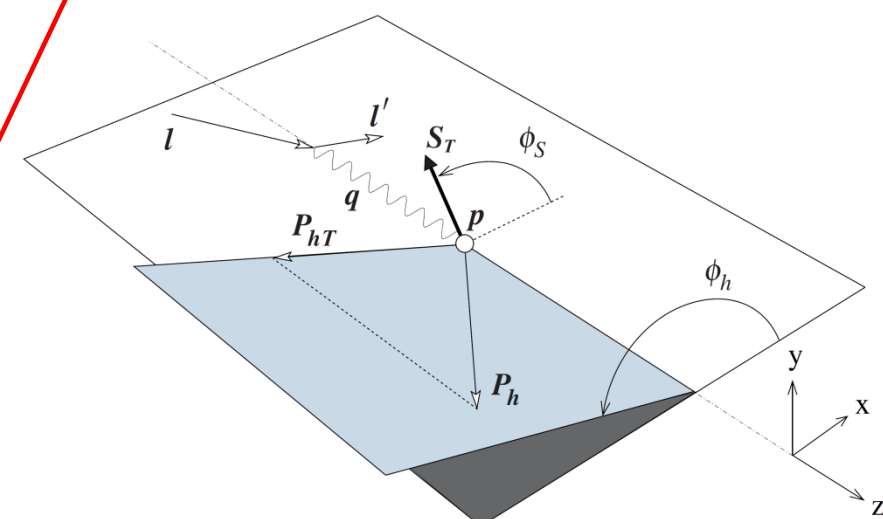
$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$A_{U(L),T}^{w(\phi_h, \phi_s)} = \frac{F_{U(L),T}^{w(\phi_h, \phi_s)}}{F_{UU,T} + \varepsilon F_{UU,L}}$$

$$\left\{ \begin{array}{l} 1 + \cos \phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \\ \lambda \sin \phi_h \times \sqrt{2\varepsilon(1-\varepsilon)} A_{LU}^{\sin \phi_h} + \\ S_L \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h A_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) A_{UL}^{\sin(2\phi_h)} \right] + \\ S_L \lambda \left[\sqrt{1-\varepsilon^2} A_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h A_{LL}^{\cos \phi_h} \right] + \\ \left[\begin{array}{l} S_T \left[\begin{array}{l} \sin \phi_s \times \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin \phi_s} \right) + \\ \sin(\phi_h - \phi_s) \times \left(A_{UT}^{\sin(\phi_h - \phi_s)} \right) + \\ \sin(\phi_h + \phi_s) \times \left(\varepsilon A_{UT}^{\sin(\phi_h + \phi_s)} \right) + \\ \sin(2\phi_h - \phi_s) \times \left(\sqrt{2\varepsilon(1+\varepsilon)} A_{UT}^{\sin(2\phi_h - \phi_s)} \right) + \\ \sin(3\phi_h - \phi_s) \times \left(\varepsilon A_{UT}^{\sin(3\phi_h - \phi_s)} \right) \end{array} \right] + \\ \left[\begin{array}{l} S_T \lambda \left[\begin{array}{l} \cos \phi_s \times \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos \phi_s} \right) + \\ \cos(\phi_h - \phi_s) \times \left(\sqrt{1-\varepsilon^2} A_{LT}^{\cos(\phi_h - \phi_s)} \right) + \\ \cos(2\phi_h - \phi_s) \times \left(\sqrt{2\varepsilon(1-\varepsilon)} A_{LT}^{\cos(2\phi_h - \phi_s)} \right) \end{array} \right] \end{array} \right] \end{array} \right\}$$

← This talk

← Presented by K. Kurek
Presented by G. Schnell

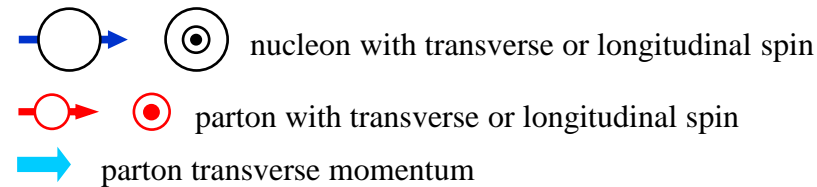
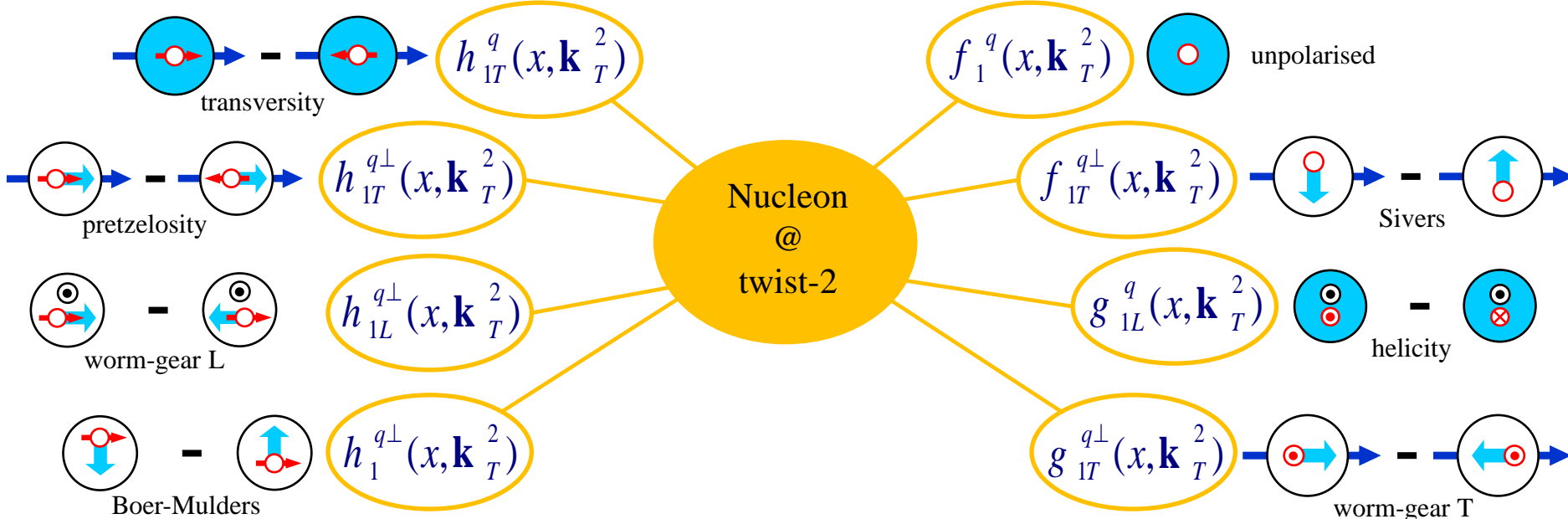


15 amplitudes:
2-''UU'', 1-''LU'', 2-''UL'', 2-''LL'', 5-''UT'', 3-''LT''

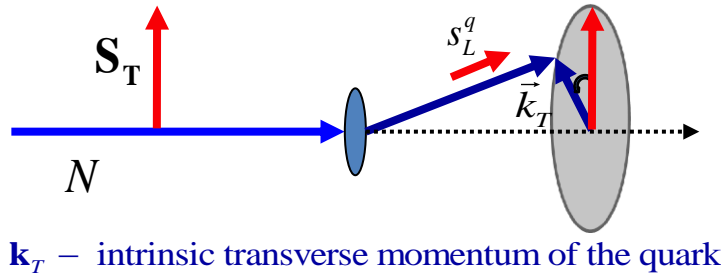
TMD parton distribution functions

Collins, Soper (81,82); Ji, Ma, Yuan (03, 04); Collins, Metz (04) – Definition of TMD DF's and FF's & QCD factorization of SIDIS.

LO QCD = Simple parton model + Factorized twist-2 PDF & FF

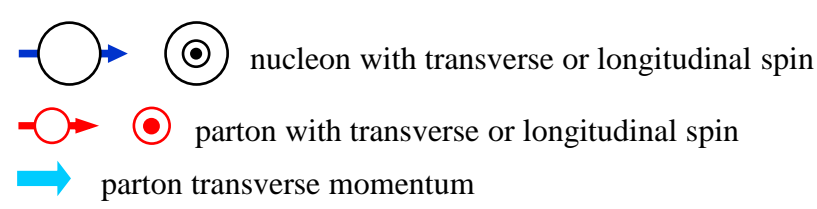
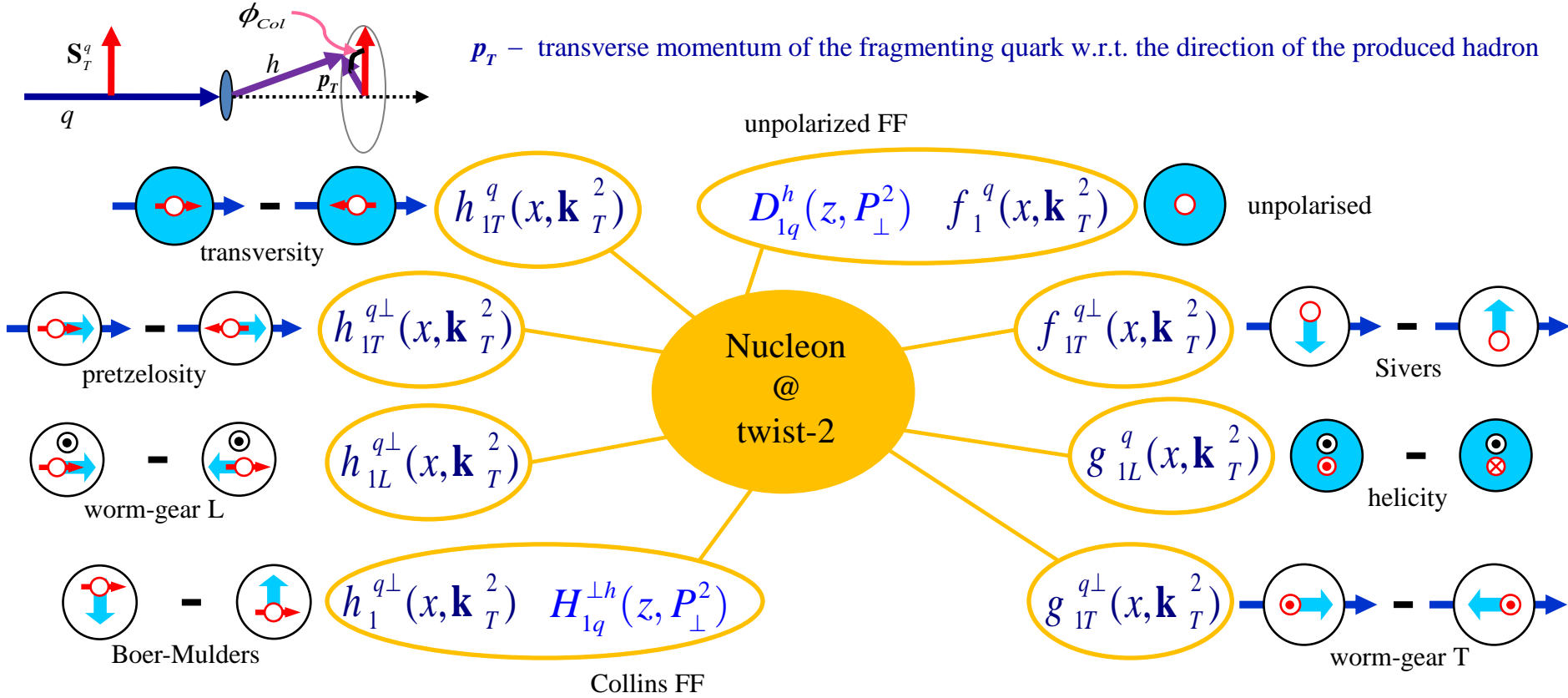


Proton goes out of the screen. Photon goes into the screen

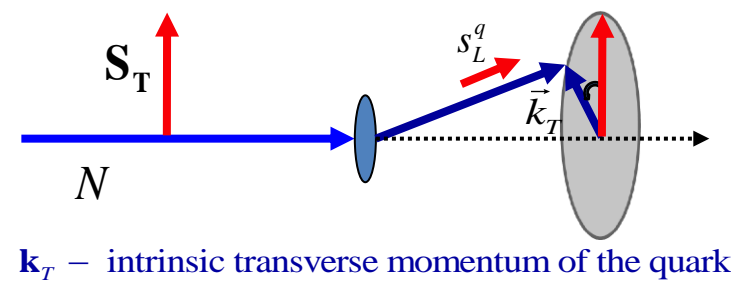


TMD parton distribution functions and FFs

Collins, Soper (81,82); Ji, Ma, Yuan (03, 04); Collins, Metz (04) – Definition of TMD DF's and FF's & QCD factorization of SIDIS.

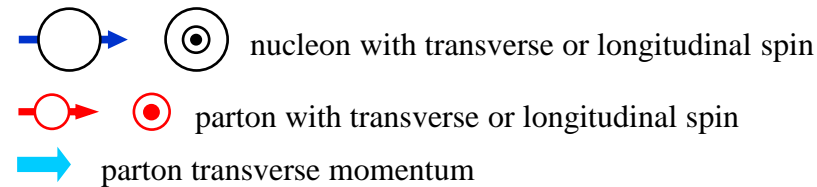
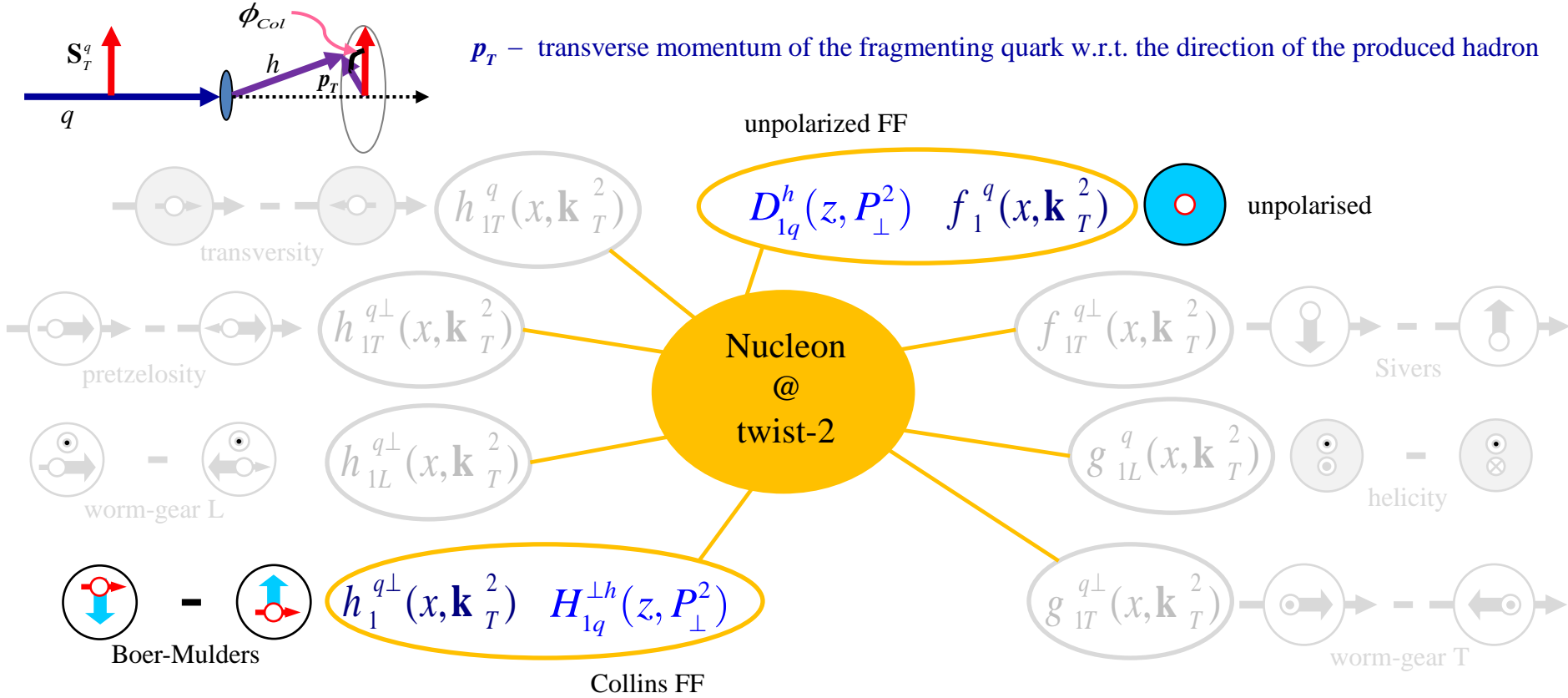


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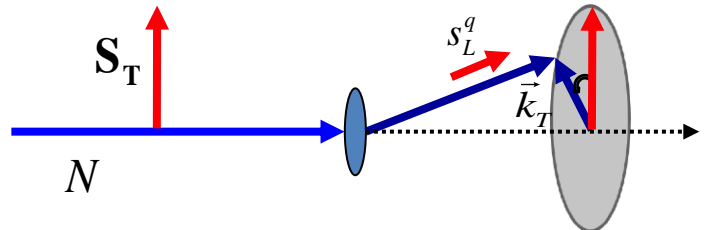


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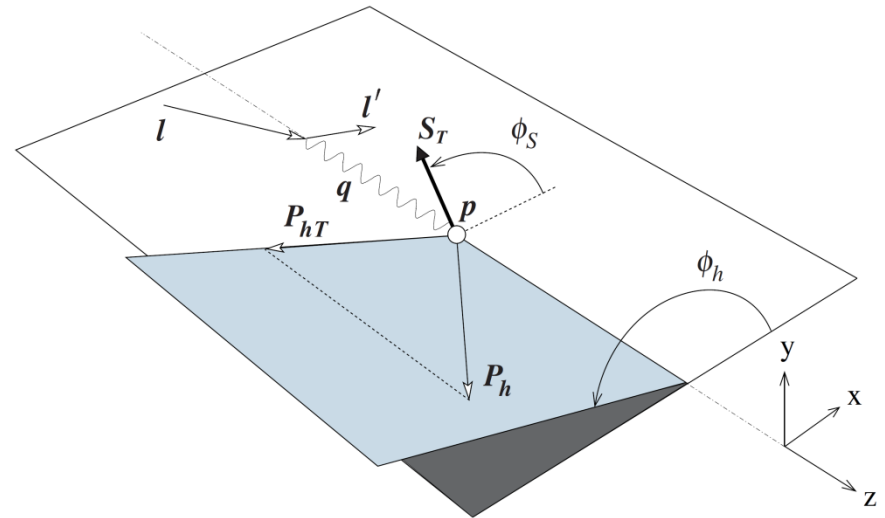


\mathbf{k}_T – intrinsic transverse momentum of the quark

SIDIS x-section (unpolarized part)

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \cos\varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos\varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)} + \dots$$



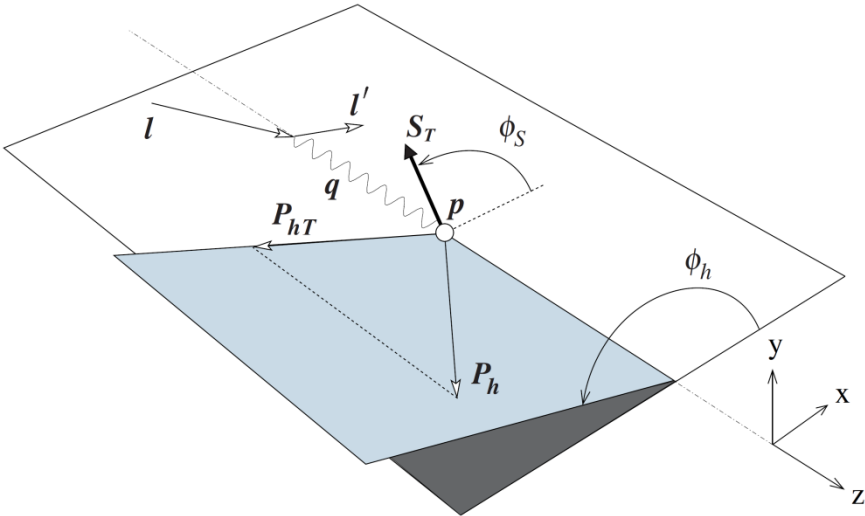
Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)}} + \dots$$



Cahn effect
R. N. Cahn, PLB 78 (1978)



Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)}}_{\text{Cahn effect}} + \dots$$

Cahn effect

R. N. Cahn, **PLB 78 (1978)**

The point that there are azimuthal dependences which arise from the transverse momenta of the partons was clearly stated in this papers:

T.P. Cheng and A. Zee, *Phys. Rev. D6* (1972) 885;

F. Ravndal, *Phys. Lett. 43B* (1973) 301.

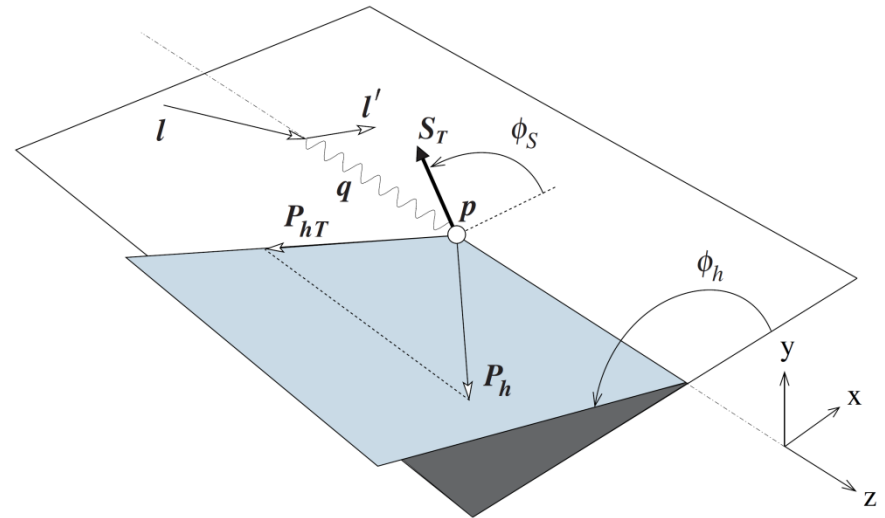
R.L. Kingsley, *Phys. Rev. D10* (1974) 1580;

A.M. Kotsinyan, *Teor. Mat. Fiz. 24* (1975) 206;

Engl. transl. *Theor. Math. Phys. 24* (1976) 776.



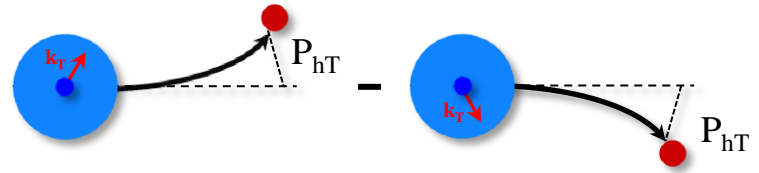
A. Kotzinian On behalf of:
T.P. Cheng, A. Zee, F. Ravndal,
R.L. Kingsley and himself



Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

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Kinematic effect: non-zero k_T induces an azimuthal modulation



Cahn effect
R. N. Cahn, PLB 78 (1978)

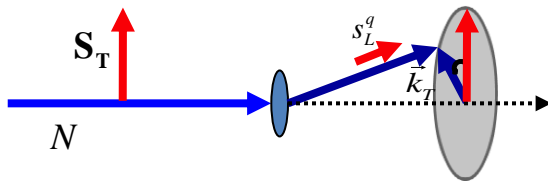
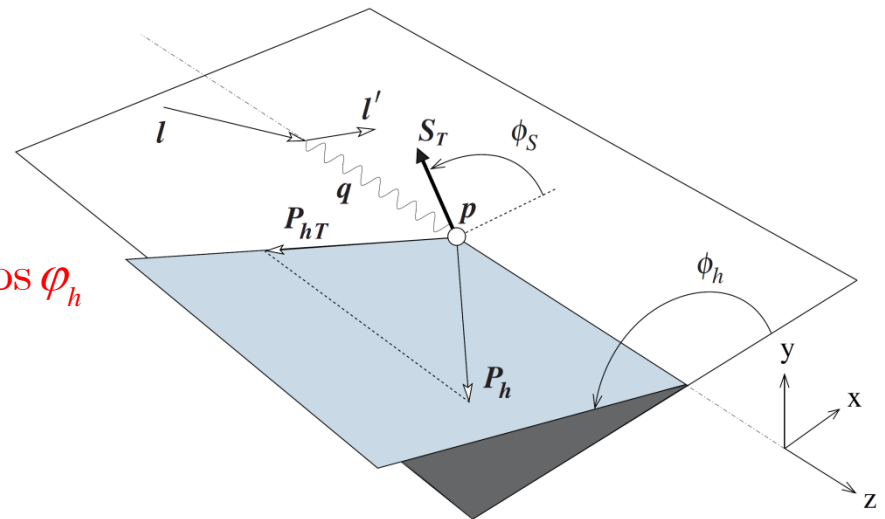
$$\hat{s} \approx xs \left[1 - 2\sqrt{1-y} \frac{k_T}{Q} \cdot \cos \varphi_q \right]$$

$$\hat{u} \approx -xs(1-y) \left[1 - \frac{2k_T}{Q\sqrt{1-y}} \cdot \cos \varphi_q \right]$$

$$\hat{t} = -Q^2 = -xys, \quad \text{where } s = (l + P)^2$$

$$d\sigma^{lp \rightarrow l'hX} \propto d\sigma^{lq \rightarrow lq} \propto \frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2}$$

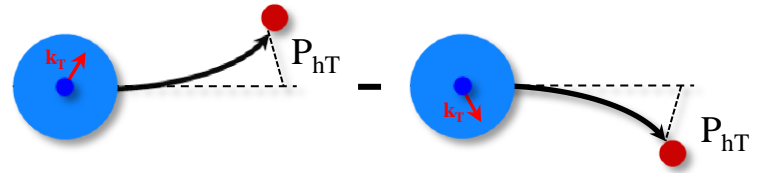
$$k_T \rightarrow \cos \varphi_q \rightarrow \cos \varphi_h$$



Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xy Q^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)}}_{\text{Cahn effect}} + \dots$$



Kinematic effect: non-zero k_T induces an azimuthal modulation



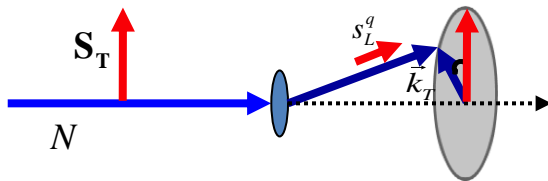
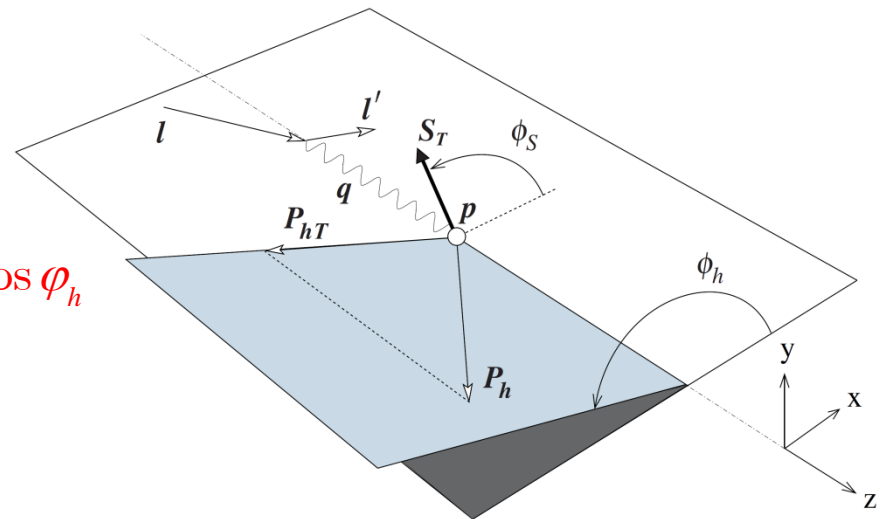
Cahn effect
R. N. Cahn, PLB 78 (1978)

$$\hat{s} \approx xs \left[1 - 2\sqrt{1-y} \frac{k_T}{Q} \cdot \cos \varphi_q \right] + O\left(\frac{k_T^2}{Q^2}\right)$$

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$$\hat{t} = -Q^2 = -xys, \quad \text{where } s = (l + P)^2$$

$$d\sigma^{lp \rightarrow l'hX} \propto d\sigma^{lq \rightarrow lq} \propto \frac{\hat{s}^2 + \hat{u}^2}{\hat{t}^2} \quad k_T \rightarrow \cos \varphi_q \rightarrow \cos \varphi_h$$

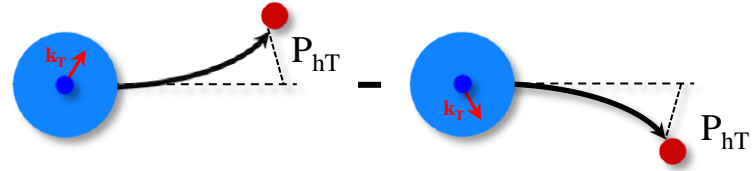


Cahn effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

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Kinematic effect: non-zero k_T induces an azimuthal modulation



Cahn effect
R. N. Cahn, PLB 78 (1978)

$$F_{UU}^{\cos\varphi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot \mathbf{p}_T}{M_h} \left(xh H_{1q}^{\perp h} + \frac{M_h}{M} f_1^q \frac{\tilde{D}_q^{\perp h}}{z} \right) - \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(xf^{\perp q} D_{1q}^h + \frac{M_h}{M} h_1^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

$$C[wfD] = x \sum_q e_q^2 \int d^2\mathbf{k}_T d^2\mathbf{p}_T \delta^{(2)}(\mathbf{k}_T - \mathbf{p}_T - \mathbf{P}_{hT}/z) w(\mathbf{k}_T, \mathbf{p}_T) f^q(x, \mathbf{k}_T^2) D_q^h(z, \mathbf{k}_T^2)$$

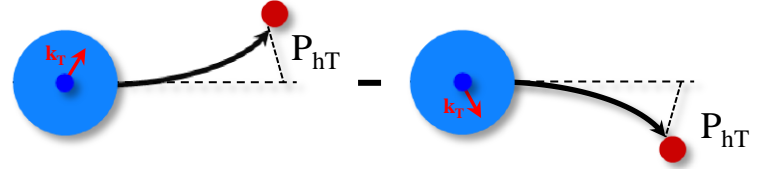
$$\hat{h} = \frac{\vec{P}_{hT}}{|\vec{P}_{hT}|}, \mathbf{p}_T - TM \text{ of the quark w.r.t. the direction of the produced hadron}$$

Cahn effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

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Kinematic effect: non-zero k_T induces an azimuthal modulation

Cahn effect

R. N. Cahn, PLB 78 (1978)

$$\tilde{x}\hat{h} + \frac{k_T^2}{M^2} h_1^{\perp q}$$

$$x f^{\perp q} + f_1^q$$

$$F_{UU}^{\cos \phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot \mathbf{p}_T}{M_h} \left(x h H_{1q}^{\perp h} + \frac{M_h}{M} f_1^q \frac{\tilde{D}_q^{\perp h}}{z} \right) - \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(x f^{\perp q} D_{1q}^h + \frac{M_h}{M} h_1^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

$$F_{UU}^{\cos \phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot \mathbf{p}_T}{M_h} \left(\left(x \tilde{x}\hat{h} + \frac{k_T^2}{M^2} h_1^{\perp q} \right) H_{1q}^{\perp h} + \frac{M_h}{M} f_1^q \frac{\tilde{D}_q^{\perp h}}{z} \right) - \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(\left(x f^{\perp q} + f_1^q \right) D_{1q}^h + \frac{M_h}{M} h_1^{\perp q} \frac{\tilde{H}_q^h}{z} \right) \right\}$$

P. J. Mulders and R. D. Tangerman,
Nucl. Phys. B461 (1996) 197–237
D. Boer, P. J. Mulders, and O. V. Teryaev,
Phys. Rev. D57 (1998) 3057–3064
Bacchetta et al. JHEP 0702:093,2007

$$C[wfD] = x \sum_q e_q^2 \int d^2 \mathbf{k}_T d^2 \mathbf{p}_T \delta^{(2)}(\mathbf{k}_T - \mathbf{p}_T - \mathbf{P}_{hT}/z) w(\mathbf{k}_T, \mathbf{p}_T) f^q(x, \mathbf{k}_T^2) D_q^h(z, \mathbf{k}_T^2)$$

$$\hat{h} = \vec{P}_{hT} / |\vec{P}_{hT}|, \mathbf{p}_T = TM \text{ of the quark w.r.t. the direction of the produced hadron}$$

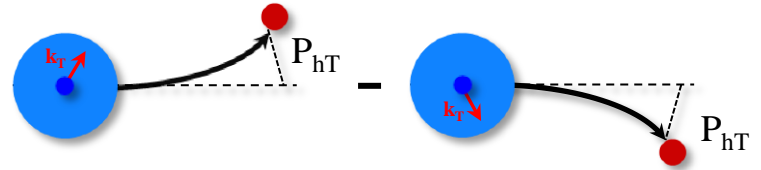
Bakur Parsamyan

Cahn effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

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Kinematic effect: non-zero k_T induces an azimuthal modulation



Cahn effect
R. N. Cahn, PLB 78 (1978)

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P. J. Mulders and R. D. Tangerman,
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Phys. Rev. D57 (1998) 3057–3064
Bacchetta et al. JHEP 0702:093,2007

Wandzura-Wilczek approximation
neglecting quark-gluon-quark correlators
(setting all functions with a tilde to zero)

$$F_{UU}^{\cos \varphi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot \mathbf{p}_T}{M_h} \left(\cancel{x\tilde{h} + \frac{k_T^2}{M^2} h_1^{\perp q}} \right) H_{1q}^{\perp h} + \frac{M_h}{M} \cancel{f_1^q} \frac{\cancel{\tilde{D}_q^{\perp h}}}{z} - \frac{\hat{h} \cdot \mathbf{k}_T}{M} \left(\cancel{x f^{\perp q} + f_1^q} \right) D_{1q}^h + \frac{M_h}{M} h_1^{\perp q} \frac{\cancel{\tilde{H}_q^h}}{z} \right\}$$

$$F_{UU}^{\cos \varphi_h} = \frac{2M}{Q} C \left\{ -\frac{(\hat{h} \cdot \mathbf{p}_T) k_T^2}{M_h M^2} h_1^{\perp q} H_{1q}^{\perp h} - \frac{\hat{h} \cdot \mathbf{k}_T}{M} f_1^q D_{1q}^h + \dots \right\}$$

sub-leading Cahn+Boer-Mulders effect

$$C[wfD] = x \sum_q e_q^2 \int d^2 \mathbf{k}_T d^2 \mathbf{p}_T \delta^{(2)}(\mathbf{k}_T - \mathbf{p}_T - \mathbf{P}_{hT}/z) w(\mathbf{k}_T, \mathbf{p}_T) f^q(x, \mathbf{k}_T^2) D_q^h(z, \mathbf{k}_T^2)$$

$\hat{h} = \vec{P}_{hT} / |\vec{P}_{hT}|$, $\mathbf{p}_T = TM$ of the quark w.r.t. the direction of the produced hadron

“Longitudinal” Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \varphi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \varphi_h} + \cos(2\varphi_h) \times \varepsilon A_{UU}^{\cos(2\varphi_h)} + \dots}_{\text{Cahn effect}} + \dots + \underline{S_L \lambda \sqrt{2\varepsilon(1-\varepsilon)} \cos \varphi_h \times A_{LL}^{\cos \varphi_h} + \dots}$$



Cahn effect
R. N. Cahn, PLB 78 (1978)

“Longitudinal” Cahn effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \phi_h}}_{\text{Cahn effect}} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \dots + \underbrace{S_L \lambda \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h \times A_{LL}^{\cos \phi_h}}_{\text{Longitudinal Cahn effect}} + \dots$$



Cahn effect
R. N. Cahn, PLB 78 (1978)

Longitudinal Cahn effect
Kotzinian et al. Phys. Rev. D 74, 074015 (2006)

$$x\tilde{e}_L$$



$$F_{LL}^{\cos \phi_h} = \frac{2M}{Q} C \left\{ \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M_h} \left(x\tilde{e}_L H_{1q}^{\perp h} - \frac{M_h}{M} g_{1L}^q \frac{\tilde{D}_q^{\perp h}}{z} \right) - \frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M} \left(xg_L^{\perp q} D_{1q}^h + \frac{M_h}{M} h_{1L}^{\perp q} \frac{\tilde{E}_q^h}{z} \right) \right\}$$

$$x\tilde{g}_L^{\perp q} + \frac{m}{M} h_{1L}^{\perp q} + g_{1L}^q$$



$$F_{LL}^{\cos \phi_h} = \frac{2M}{Q} C \left\{ \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M_h} \left(\cancel{x\tilde{e}_L} H_{1q}^{\perp h} + \frac{M_h}{M} \cancel{g_{1L}^q} \frac{\tilde{D}_q^{\perp h}}{z} \right) - \frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M} \left(\left(\cancel{x\tilde{g}_L^{\perp q} + \frac{m_q}{M} \cancel{h_{1L}^{\perp q}} + g_{1L}^q \right) D_{1q}^h + \frac{M_h}{M} \cancel{h_{1L}^{\perp q}} \frac{\tilde{E}_q^h}{z} \right) \right\}$$



$$F_{LL}^{\cos \phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M} g_{1L}^q D_{1q}^h + \dots \right\}$$

sub-leading effect

$$C[wfD] = x \sum_q e_q^2 \int d^2 \mathbf{k}_T d^2 \mathbf{p}_T \delta^{(2)}(\mathbf{k}_T - \mathbf{p}_T - \mathbf{P}_{hT}/z) w(\mathbf{k}_T, \mathbf{p}_T) f^q(x, \mathbf{k}_T^2) D_q^h(z, \mathbf{k}_T^2)$$

$\hat{\mathbf{h}} = \vec{P}_{hT} / |\vec{P}_{hT}|$, $\mathbf{p}_T = TM$ of the quark w.r.t. the direction of the produced hadron

$$\hat{s} \approx xs \left[1 - 2\sqrt{1-y} \frac{k_T}{Q} \cdot \cos \phi_q \right] + O\left(\frac{k_T^2}{Q^2}\right)$$

$$\hat{u} \approx -xs(1-y) \left[1 - \frac{2k_T}{Q\sqrt{1-y}} \cdot \cos \phi_q \right] + O\left(\frac{k_T^2}{Q^2}\right)$$

$$\hat{t} = -Q^2 = -xys, \quad \text{where } s = (l+P)^2$$

$$d\sigma^{lp \rightarrow l'hx} \propto d\sigma^{lq \rightarrow l'q} \propto \frac{\hat{s}^2 + \hat{u}^2 + \lambda \lambda_q (\hat{s}^2 - \hat{u}^2)}{\hat{t}^2}$$

Wandzura-Wilczek approximation

+ neglecting m_q scaled terms

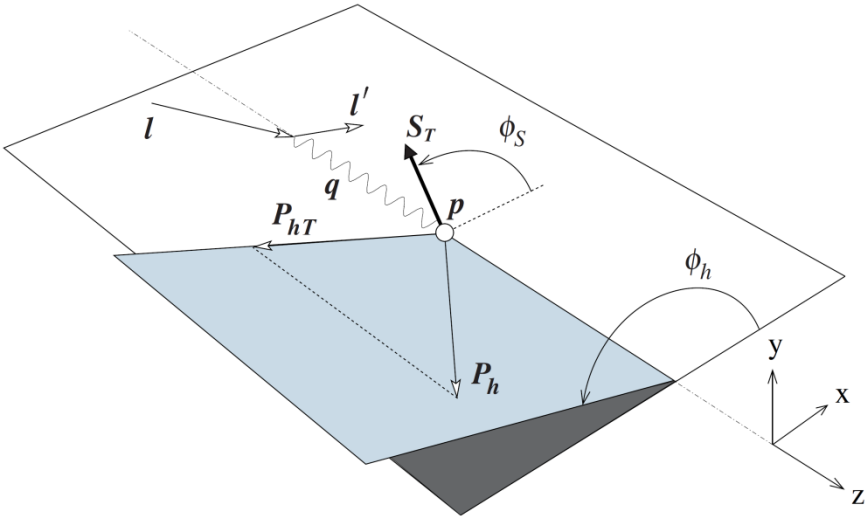
Boer-Mulders effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \underbrace{\cos \phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)}} + \dots$$



Boer-Mulders effect
D. Boer and P. J. Mulders, PRD 57 (1998)



Boer-Mulders effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

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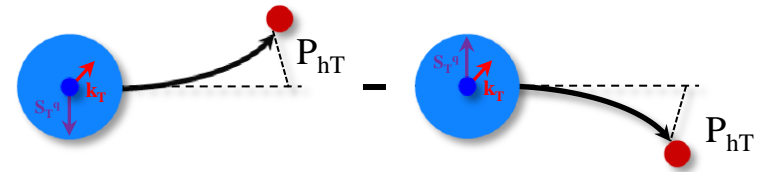
$$1 + \underbrace{\cos \phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos \phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \dots}$$



Boer-Mulders-Collins effect
D. Boer and P.J. Mulders, PRD 57 (1998)

Boer-Mulders PDF Collins FF

$$F_{UU}^{\cos 2\phi_h} = C \left\{ - \frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} \left(h_1^{\perp q} H_{1q}^{\perp h} \right) \right\}$$



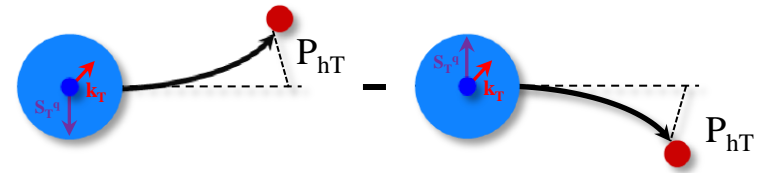
Arises due to the correlations between quark transverse spin and intrinsic transverse momentum
Is a leading order effect

Boer-Mulders effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

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Boer-Mulders-Collins effect
D. Boer and P.J. Mulders, PRD 57 (1998)

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$$F_{UU}^{\cos 2\varphi_h} = C \left\{ -\frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} h_1^{\perp q} H_{1q}^{\perp h} \right\}$$

$$F_{UU}^{\cos 2\varphi_h} = C \left\{ -\frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} h_1^{\perp q} H_{1q}^{\perp h} \right\} + \left(\frac{M}{Q} \right)^2 C \left\{ -\frac{2(\hat{h} \cdot k_T)^2 - k_T^2}{M^2} f_1^q D_{1q}^h + \dots \right\}$$



Cahn effect



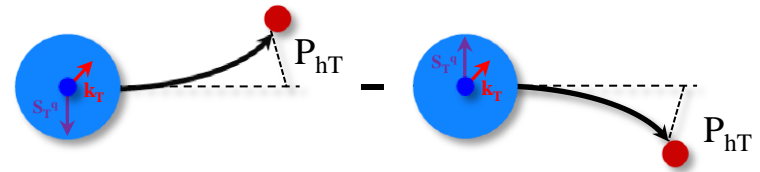
Boer-Mulders effect + twist-4 Cahn effect

Boer-Mulders effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\varphi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

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Arises due to the correlations between quark transverse spin and intrinsic transverse momentum



Boer-Mulders-Collins effect
D. Boer and P.J. Mulders, PRD 57 (1998)

Boer-Mulders PDF Collins FF

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$$F_{UU}^{\cos 2\varphi_h} = C \left\{ -\frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} h_1^{\perp q} H_{1q}^{\perp h} \right\} + \left(\frac{M}{Q} \right)^2 C \left\{ -\frac{2(\hat{h} \cdot k_T)^2 - k_T^2}{M^2} f_1^q D_{1q}^h + \dots \right\}$$

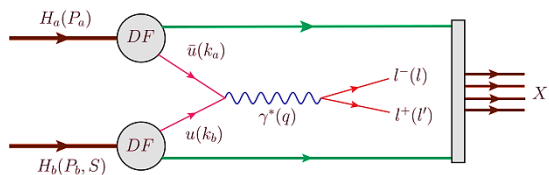


Cahn effect



Boer-Mulders effect + twist-4 Cahn effect

can be accessed also through DY



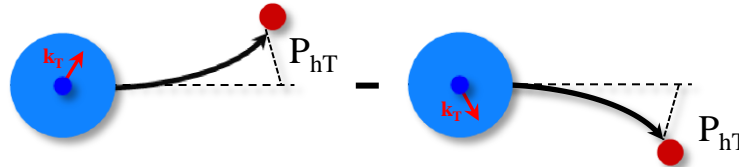
Cahn and Boer-Mulders effects

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \cos\phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos\phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \dots$$



Cahn effect
R. N. Cahn, PLB 78 (1978)



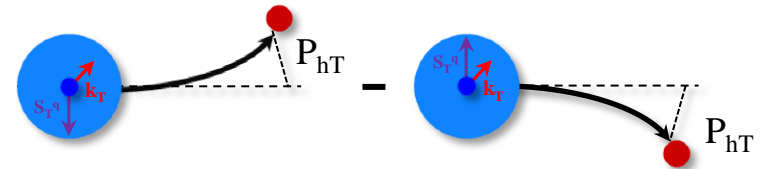
Cahn effect: kinematic effect arising due to the intrinsic transverse motion of quark

$$F_{UU}^{\cos\phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot k_T}{M} f_1^q D_{1q}^h - \frac{(\hat{h} \cdot p_T) k_T^2}{M_h M^2} h_1^{\perp q} H_{1q}^{\perp h} \right\}$$

sub-leading Cahn + Boer-Mulders effects



Boer-Mulders effect
D. Boer and P. J. Mulders, PRD 57 (1998)



Boer-Mulders effect: correlations between quark transverse spin and intrinsic transverse momentum

$$F_{UU}^{\cos 2\phi_h} = C \left\{ -\frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} h_1^{\perp q} H_{1q}^{\perp h} \right\} + \left(\frac{M}{Q} \right)^2 C \left\{ -\frac{2(\hat{h} \cdot k_T)^2 - k_T^2}{M^2} f_1^q D_{1q}^h + \dots \right\}$$

Boer-Mulders effect + twist-4 Cahn effect

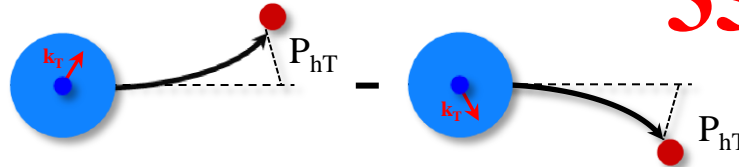
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Cahn effect
R. N. Cahn, PLB 78 (1978)



35 years !!!

Cahn effect: kinematic effect arising due to the intrinsic transverse motion of quark

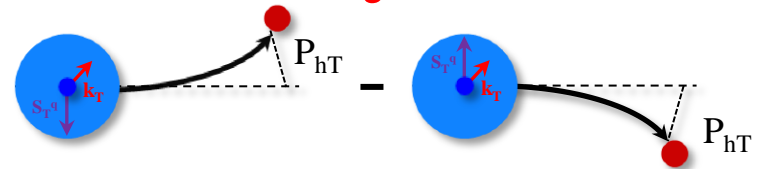
$$F_{UU}^{\cos\phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot k_T}{M} f_1^q D_{1q}^h - \frac{(\hat{h} \cdot p_T) k_T^2}{M_h M^2} h_1^{\perp q} H_{1q}^{\perp h} \right\}$$

sub-leading Cahn + Boer-Mulders effects

15 years !!!



Boer-Mulders effect
D. Boer and P. J. Mulders, PRD 57 (1998)



Boer-Mulders effect: correlations between quark transverse spin and intrinsic transverse momentum

$$F_{UU}^{\cos 2\phi_h} = C \left\{ -\frac{2(\hat{h} \cdot p_T)(\hat{h} \cdot k_T) - p_T \cdot k_T}{MM_h} h_1^{\perp q} H_{1q}^{\perp h} \right\} + \left(\frac{M}{Q} \right)^2 C \left\{ -\frac{2(\hat{h} \cdot k_T)^2 - k_T^2}{M^2} f_1^q D_{1q}^h + \dots \right\}$$

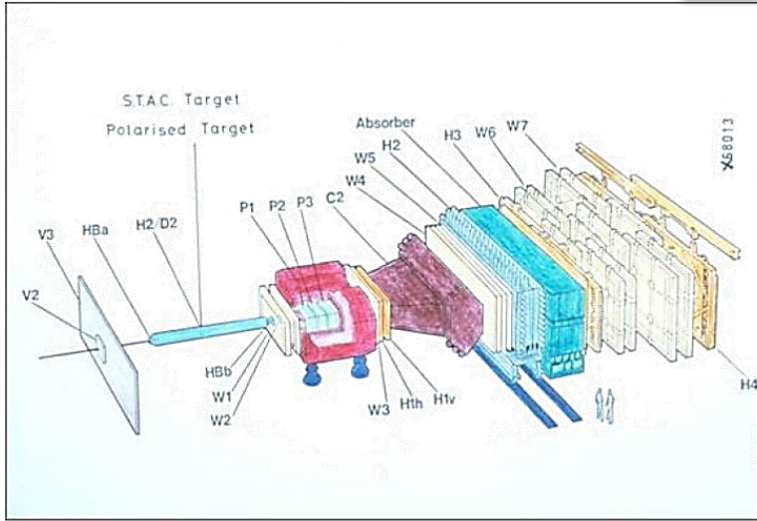
Boer-Mulders effect + twist-4 Cahn effect

Experimental data: part I

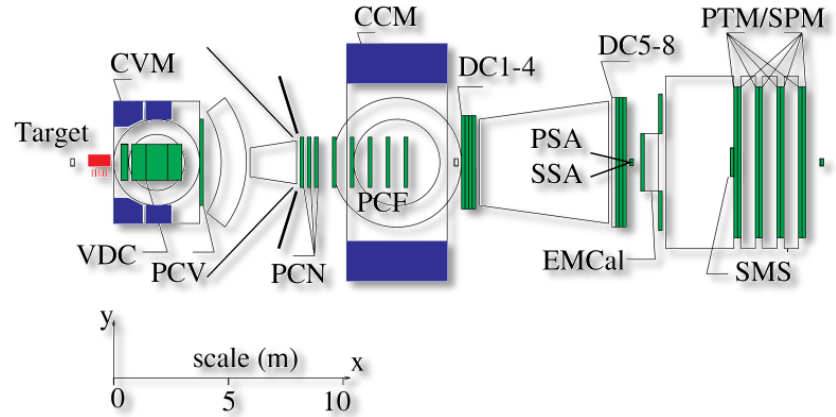
“First” results (year \leq 2009)

Experiments in last 35 years: part I

EMC CERN (μ - p , μ - d) @ 280 GeV




Fermilab E665 (μ - p , μ - d) @ 490 GeV

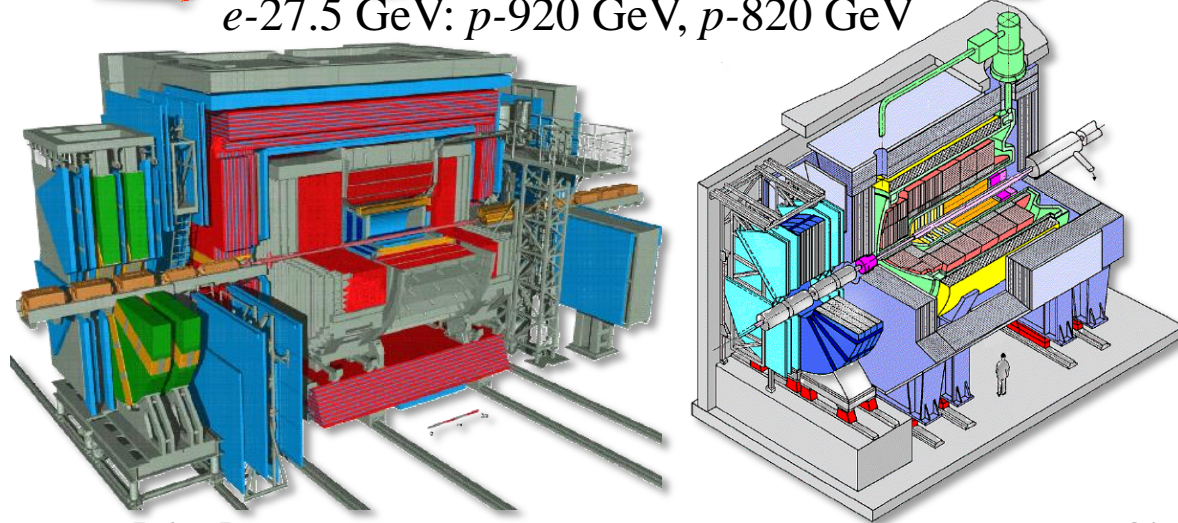
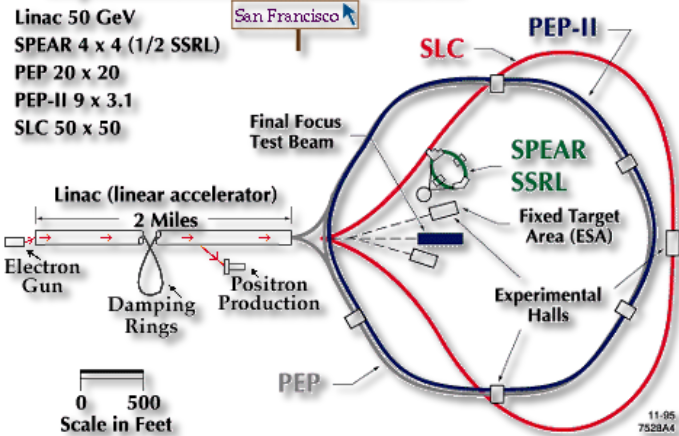


SLAC (e - p , e - d) @ 19.5 GeV

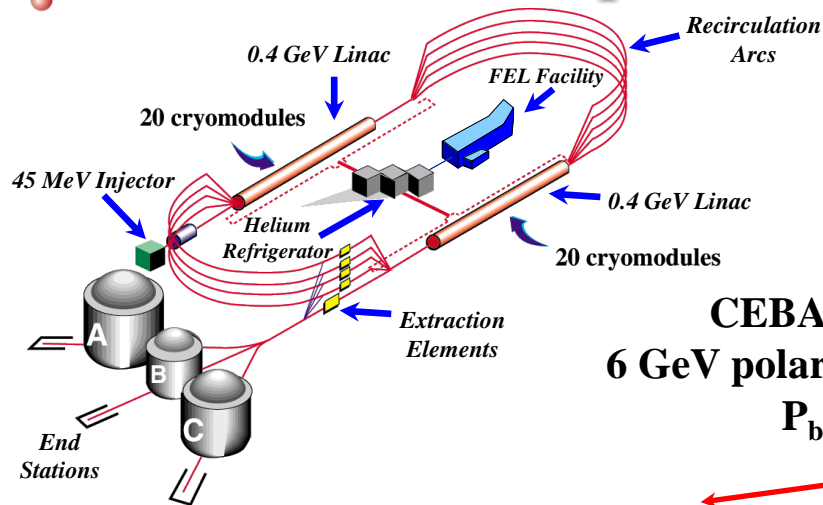


and  e - p collider HERA, DESY 
 e -27.5 GeV: p -920 GeV, p -820 GeV

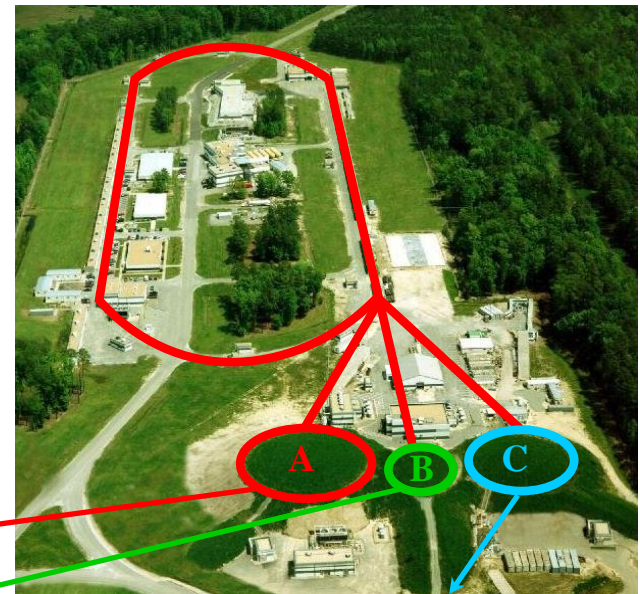
Experimental Areas at SLAC



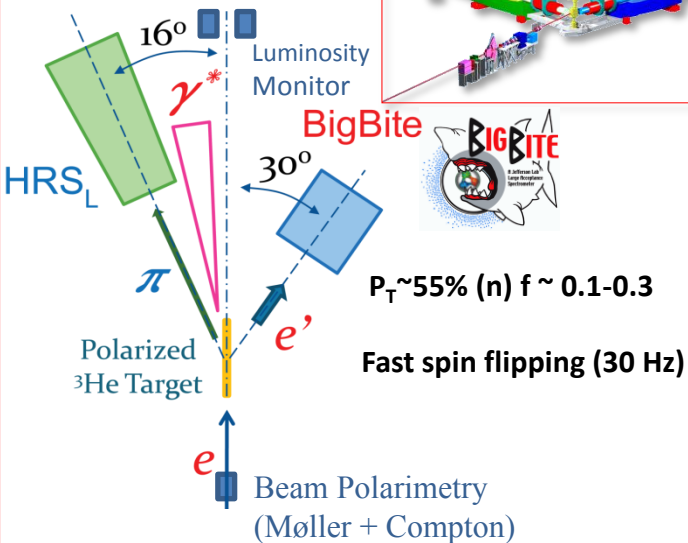
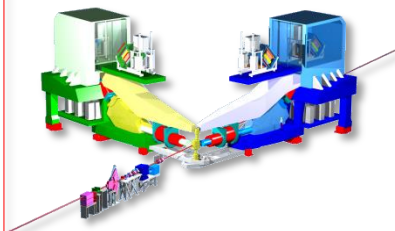
Jefferson Lab experimental halls



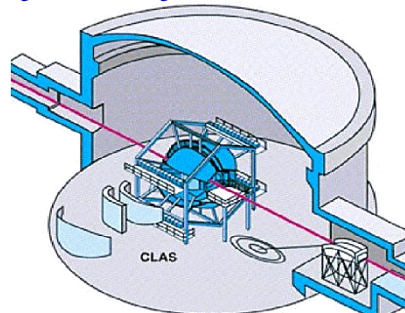
CEBAF accelerator
6 GeV polarized electron beam
 $P_{\text{beam}} \approx 85\%$



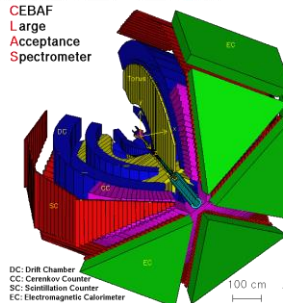
Hall A: two HRS'
 ${}^3\text{He}$ gas target (40 cm)



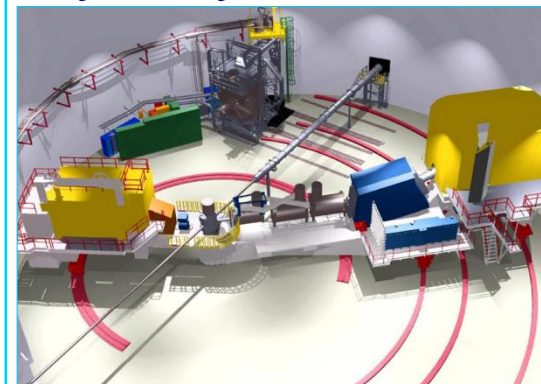
Hall B: CLAS
 NH_3 and ND_3 HD-Ice targets



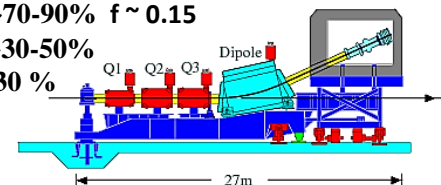
Polarizations
 Beam: $\sim 80\%$
 NH_3 proton 80%
 $\text{ND}_3 \sim 30\%$
 HD-Ice (H-75%, D-25%)
 $f \sim 0.15$



Hall C: HMS+SOS
 NH_3 and ND_3 LiD targets



Polarizations
 NH_3 : $\sim 70-90\%$ $f \sim 0.15$
 ND_3 : $\sim 30-50\%$
 LiD: $\sim 30\%$



(SLAC) *Phys. Rev. Lett.* **31**, 786 (1973)
(EMC) *Phys. Lett. B* 130 (1983) 118,
(EMC) *Z. Phys. C34* (1987) 277
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(E665) *Phys. Rev. D48* (1993) 5057
(ZEUS) *Eur. Phys. J. C11*, 251 (1999)
(ZEUS) *Phys. Lett. B* **481**, 199 (2000)
(H1) *Phys. Lett. B654*, 148 (2007)

Experiments in last 35 years: first results

EMC, E665, H1
and ZEUS

High beam energy, broad kinematic range
No hadron type and charge distinction
(averaged over any possible flavor dependence)
EMC, ZEUS – only hydrogen target
E665 – combined hydrogen and deuterium targets
Not enough statistics to look at differential x-sections in more than two kinematic variables

SLAC, JLab hall C

Relatively low beam energy, restricted kinematic range
x-sections measured only at a few kinematic points

CLAS Collaboration
(JLab hall B)

Relatively low beam energy
access to 4D multi-differential x-section

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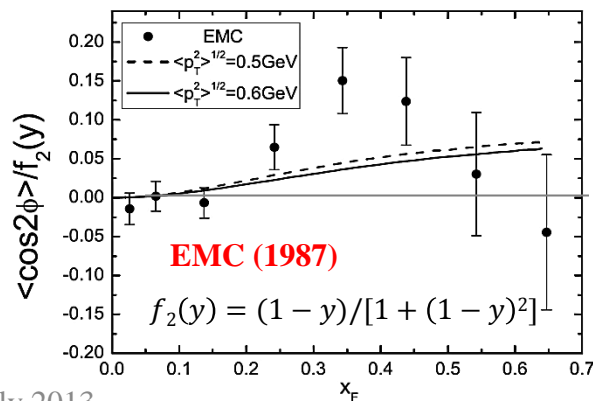
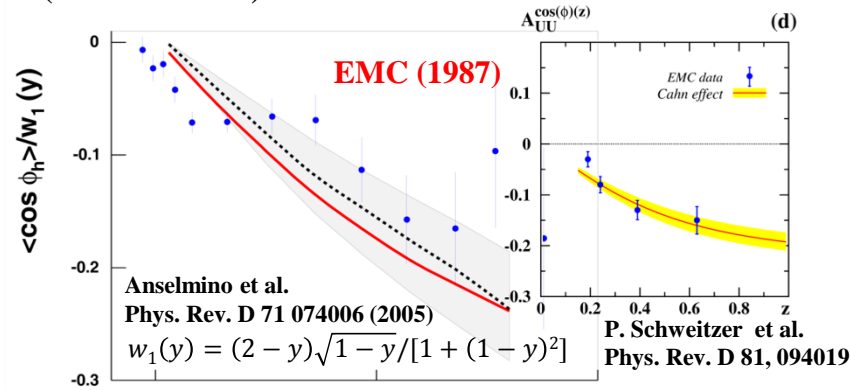
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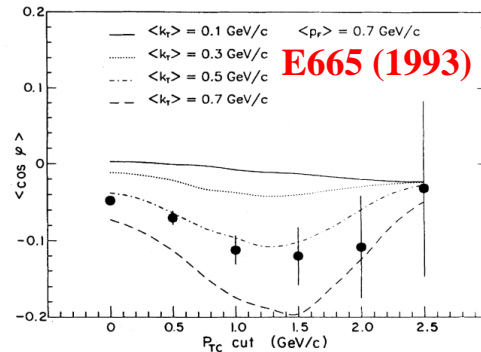
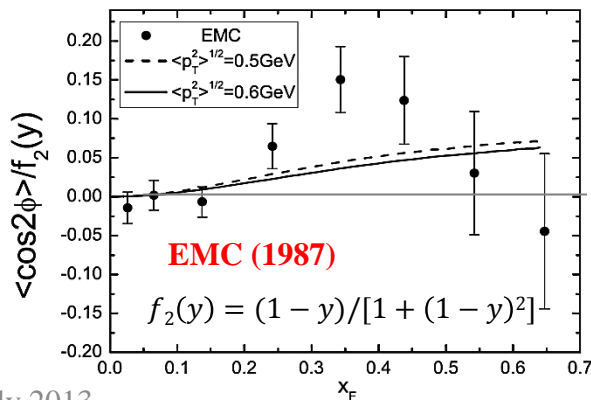
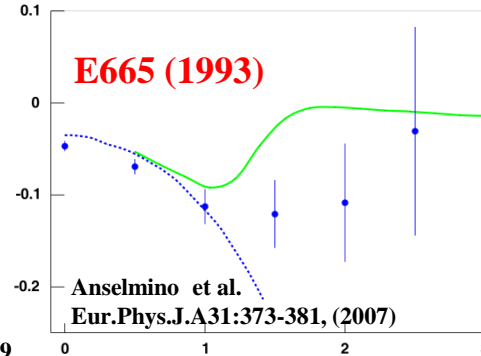
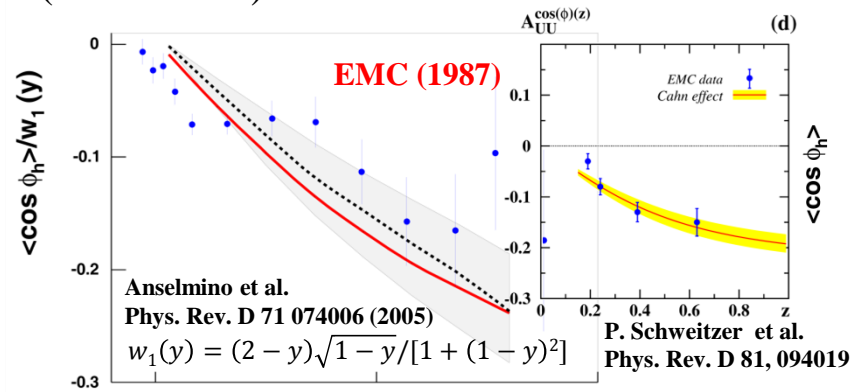
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J. Chay, S. D. Ellis, and J. W. Stirling,
 Phys. Rev. D **45**, 46 (1992), Phys. Lett. B **269**, 175 (1991).
 Bakur Parsamyan

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 EMC, ZEUS – only hydrogen target

E665 – combined hydrogen and deuterium targets

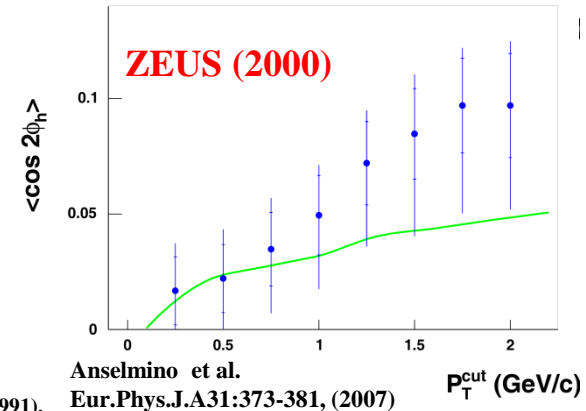
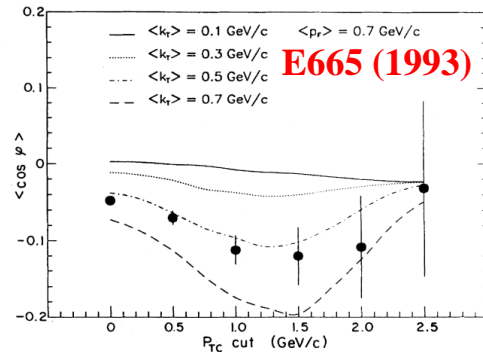
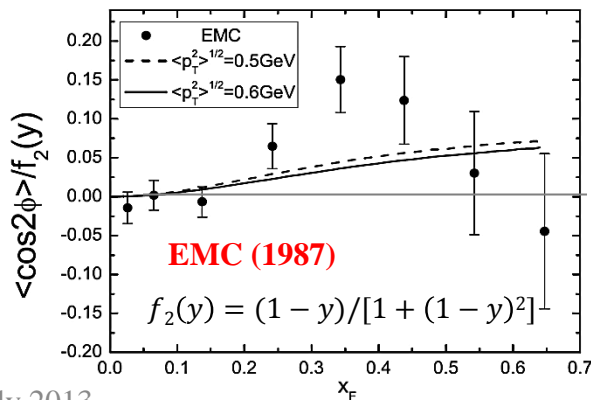
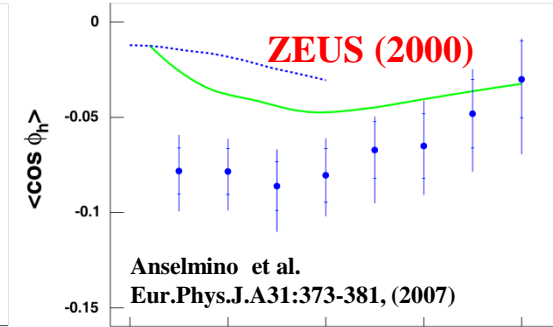
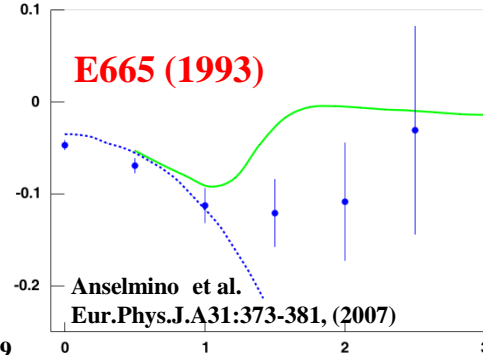
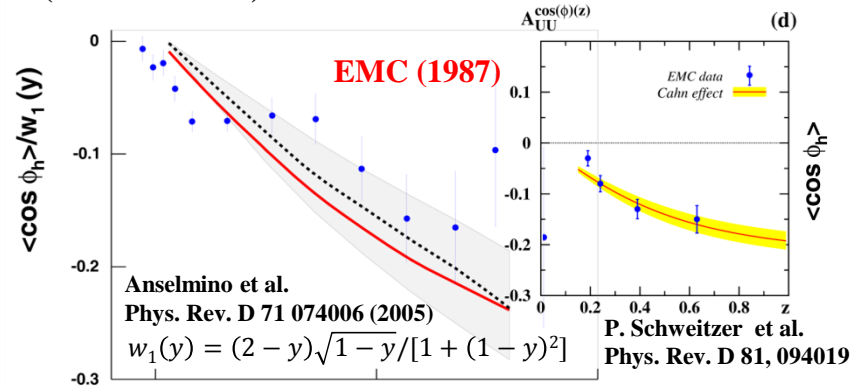
Not enough statistics to look at differential x-sections in more than two kinematic variables

SLAC, JLab hall C

Relatively low beam energy, restricted kinematic range
 x-sections measured only at a few kinematic points

CLAS Collaboration
 (JLab hall B)

Relatively low beam energy
 access to 4D multi-differential x-section



JLab 6: Hall C results

H. Mkrtychyan et al., Phys. Lett. B 665, 20 (2008).

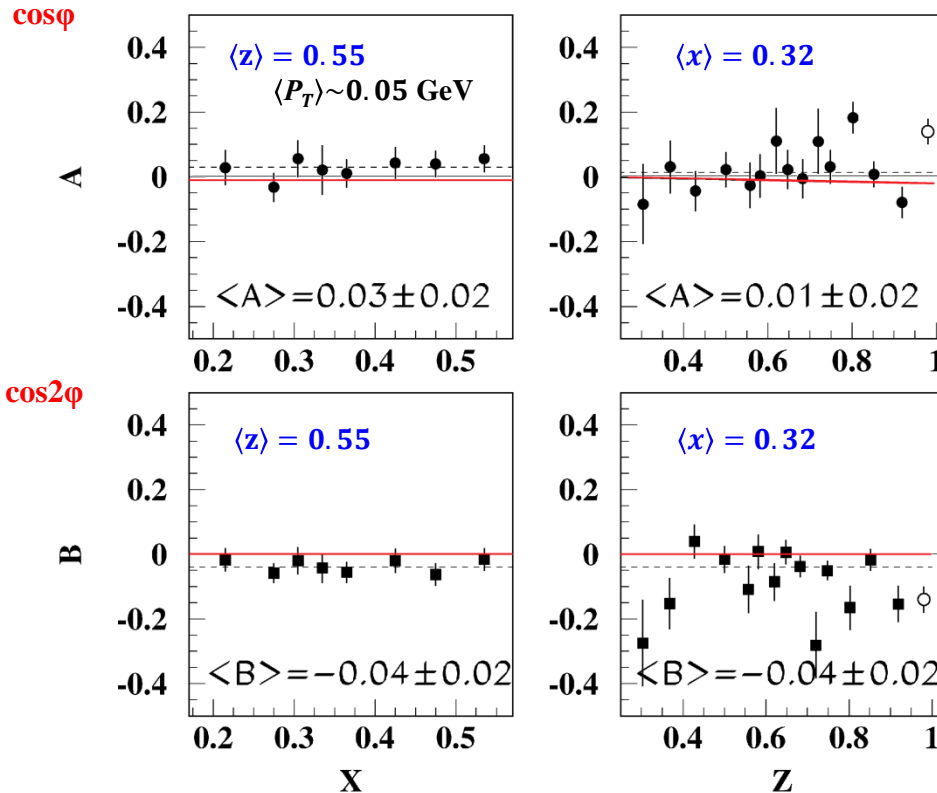
Kinematics:

$$2 < Q^2 < 4 \text{ (GeV/c)}^2; 0.2 < x < 0.5; 0.3 < z < 1$$

$$P_{hT}^2 < 0.2 \text{ (GeV/c)}^2$$

Semi-inclusive electro-production of charged pions (π^\pm) from both proton and deuteron targets, using 5.5 GeV electron beam

Small effects, consistent with theoretical predictions
Amplitudes are averaged over π^+ and π^- detected from proton and deuteron targets



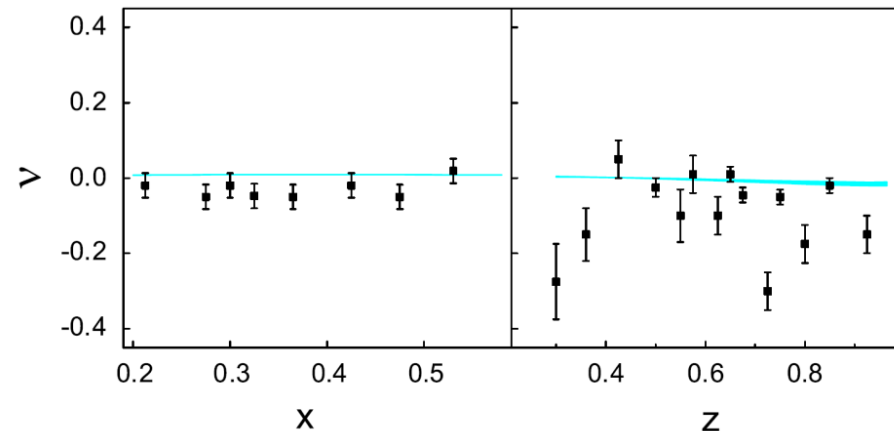
Systematic errors are estimated at ~ 0.03 on both A and B

Solid red lines are the theoretical predictions:

R.N. Cahn, Phys. Lett. B 78 (1978) 269;

R.N. Cahn, Phys. Rev. D 40 (1989) 3107

Zhang et al. Phys. Rev. D78:034035, 2008



Ingredients:

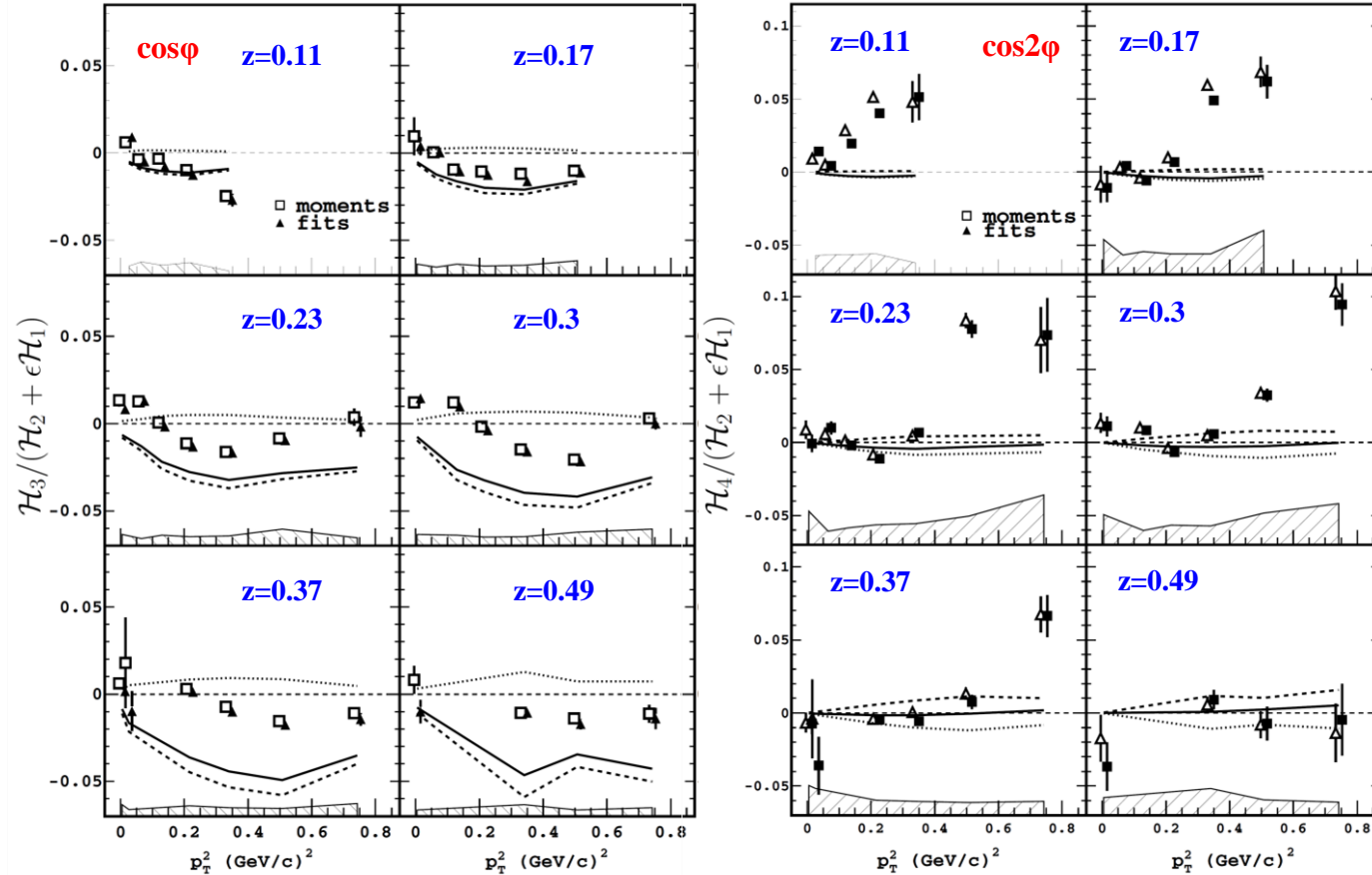
- Boer-Mulders functions parameterized from unpolarized p D Drell-Yan data by the FNAL E866/NuSea Collaboration
- combined with extracted Collins functions from *M. Anselmino et. al, Phys. Rev. D 75, 054032 (2007).*

CLAS (JLab hall B) results

M. Osipenko et al. (CLAS Collaboration)
 Phys.Rev.D80:032004,2009

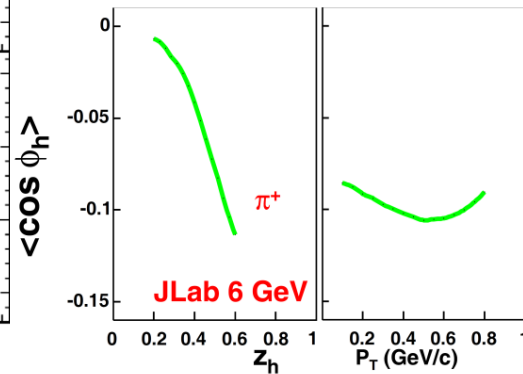
Positive pions

$1.4 < Q^2 < 7 \text{ (GeV/c)}^2$
 $0.15 < x < 1$
 $0.07 < z < 1$
 $0.005 < P_{hT}^2 < 1.5 \text{ (GeV/c)}^2$
 Beam energy 5.75 GeV



$\cos\phi$ amplitude (nonzero)
 is in strong disagreement with the
 theoretical predictions

$\cos 2\phi$ amplitude
 is compatible with zero except
 low z region where it is positive



Theoretical predictions: Cahn effect + Berger effect

R. N. Cahn, Phys. Rev. D40, 3107 (1989).

M. Anselmino et al., Phys. Rev. D71, 074006 (2005).

A. Brandenburg, V. V. Khoze, and D. Mueller, Phys. Lett. B347, 413 (1995).

Curves for Cahn contribution only
 Anselmino et al. Eur. Phys. J. A 31, 373-381 (2007)

Experimental data: part I

“First” results (year \leq 2009)

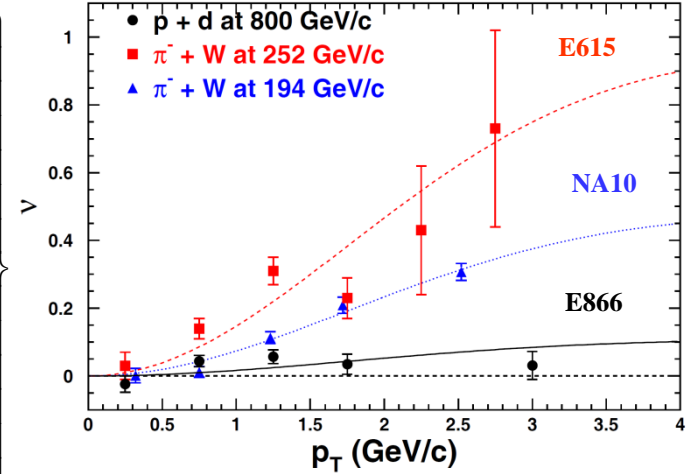
Boer-Mulders PDF from the Drell-Yan data

Boer-Mulders from Drell-Yan

S. Arnold, A. Metz and M. Schlegel, **Phys. Rev. D** **79**, 034005 (2009).

DY x-section (single-polarized)

$$\frac{d\sigma}{d\Omega} = \frac{\alpha_{em}^2}{Fq^2} \hat{\sigma}_U \left\{ \begin{array}{l} \left(1 + D_{[\sin 2\theta]} A_U^{\cos\varphi} \cos\varphi + D_{[\sin^2\theta]} A_U^{\cos 2\varphi} \cos 2\varphi \right) \\ + S_L \left(D_{[\sin 2\theta]} A_L^{\sin\varphi} \sin\varphi + D_{[\sin^2\theta]} A_L^{\sin 2\varphi} \sin 2\varphi \right) \\ \pm \left| \vec{S}_T \right| \left[\begin{array}{l} \left(D_{[1]} A_T^{\sin\varphi_S} + D_{[\cos^2\theta]} \tilde{A}_T^{\sin\varphi_S} \right) \sin\varphi_S \\ + D_{[\sin 2\theta]} \left(A_T^{\sin(\varphi+\varphi_S)} \sin(\varphi+\varphi_S) + A_T^{\sin(\varphi-\varphi_S)} \sin(\varphi-\varphi_S) \right) \\ + D_{[\sin^2\theta]} \left(A_T^{\sin(2\varphi+\varphi_S)} \sin(2\varphi+\varphi_S) + A_T^{\sin(2\varphi-\varphi_S)} \sin(2\varphi-\varphi_S) \right) \end{array} \right] \end{array} \right.$$



$$\hat{\sigma}_U = (F_U^1 + F_U^2) (1 + A_U^1 \cos^2 \theta); \quad D_{[f(\theta)]} = \frac{f(\theta)}{1 + A_U^1 \cos^2 \theta}; \quad \lambda = A_U^1, \mu = A_U^{\cos\varphi}, \nu = 2A_U^{\cos 2\varphi} \propto h_{1q}^\perp \otimes h_{1\bar{q}}^\perp$$

Experimental results:

NA10 Collaboration ($\pi^- N @ 194 \text{ GeV}/c$)

S. Falciano et al., **Z. Phys. C** **31**, 513 (1986).

M. Guanziroli et al., **Z. Phys. C** **37**, 545 (1988).

E615 Collaboration ($\pi^- N @ 252 \text{ GeV}/c$)

J. S. Conway et al., **Phys. Rev. D** **39**, 92 (1989).

E866/NuSea Collaboration at **FNAL** ($pd, pp @ 194 \text{ GeV}/c$)

pd - **FNAL-E866/NuSea**, L. Y. Zhu et al., **Phys. Rev. Lett.** **99**, 082301 (2007)

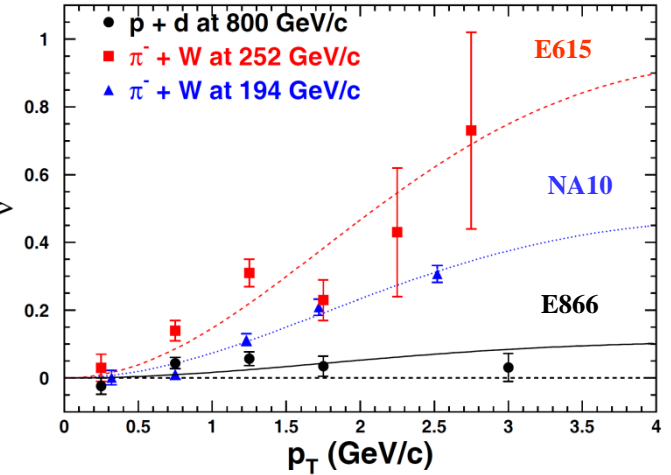
pp - **FNAL E866/NuSea**, L. Y. Zhu et al., **Phys. Rev. Lett.** **102**, 182001 (2009)

Boer-Mulders from Drell-Yan

S. Arnold, A. Metz and M. Schlegel, **Phys. Rev. D** **79**, 034005 (2009).

DY x-section (single-polarized)

$$\frac{d\sigma}{dY} = \frac{\alpha_{em}^2}{Fq^2} \hat{\sigma}_U \left\{ \begin{array}{l} \left(1 + D_{[\sin 2\theta]} A_U^{\cos\varphi} \cos\varphi + D_{[\sin^2\theta]} A_U^{\cos 2\varphi} \cos 2\varphi \right) \\ + S_L \left(D_{[\sin 2\theta]} A_L^{\sin\varphi} \sin\varphi + D_{[\sin^2\theta]} A_L^{\sin 2\varphi} \sin 2\varphi \right) \\ \pm \left| \vec{S}_T \right| \left[\left(D_{[1]} A_T^{\sin\varphi_S} + D_{[\cos^2\theta]} \tilde{A}_T^{\sin\varphi_S} \right) \sin\varphi_S \right. \\ \left. + D_{[\sin 2\theta]} \left(A_T^{\sin(\varphi+\varphi_S)} \sin(\varphi+\varphi_S) + A_T^{\sin(\varphi-\varphi_S)} \sin(\varphi-\varphi_S) \right) \right. \\ \left. + D_{[\sin^2\theta]} \left(A_T^{\sin(2\varphi+\varphi_S)} \sin(2\varphi+\varphi_S) + A_T^{\sin(2\varphi-\varphi_S)} \sin(2\varphi-\varphi_S) \right) \right] \end{array} \right\}$$



$$\hat{\sigma}_U = (F_U^1 + F_U^2)(1 + A_U^1 \cos^2\theta); \quad D_{[f(\theta)]} = \frac{f(\theta)}{1 + A_U^1 \cos^2\theta}; \quad \lambda = A_U^1, \mu = A_U^{\cos\varphi}, \nu = 2A_U^{\cos 2\varphi} \propto h_{1q}^\perp \otimes h_{1\bar{q}}^\perp$$

Experimental results:

NA10 Collaboration ($\pi^- N$ @ 194 GeV/c)

S. Falciano et al., **Z. Phys. C** **31**, 513 (1986).

M. Guanziroli et al., **Z. Phys. C** **37**, 545 (1988).

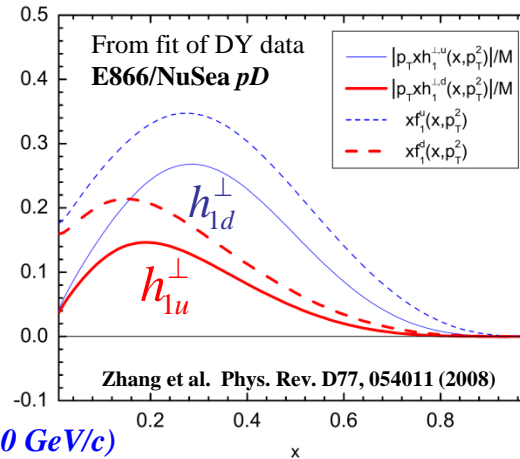
E615 Collaboration ($\pi^- N$ @ 252 GeV/c)

J. S. Conway et al., **Phys. Rev. D** **39**, 92 (1989).

E866/NuSea Collaboration at FNAL (pd, pp @ 800 GeV/c)

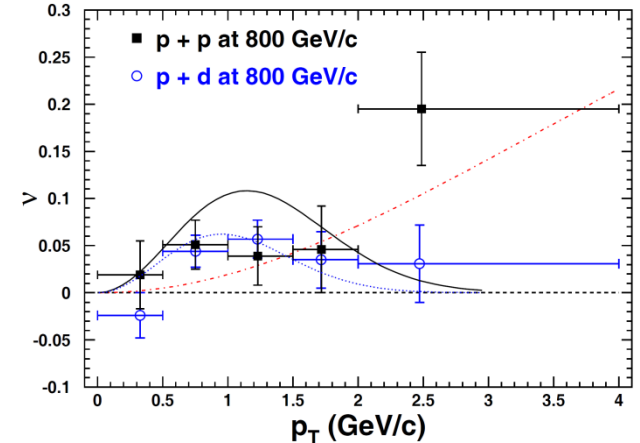
pd - FNAL-E866/NuSea, L. Y. Zhu et al., **Phys. Rev. Lett.** **99**, 082301 (2007)

pp - FNAL E866/NuSea, L. Y. Zhu et al., **Phys. Rev. Lett.** **102**, 182001 (2009)



Theoretical curves are from:

Zhang et al. **Phys. Rev. D** **78**:034035, 2008

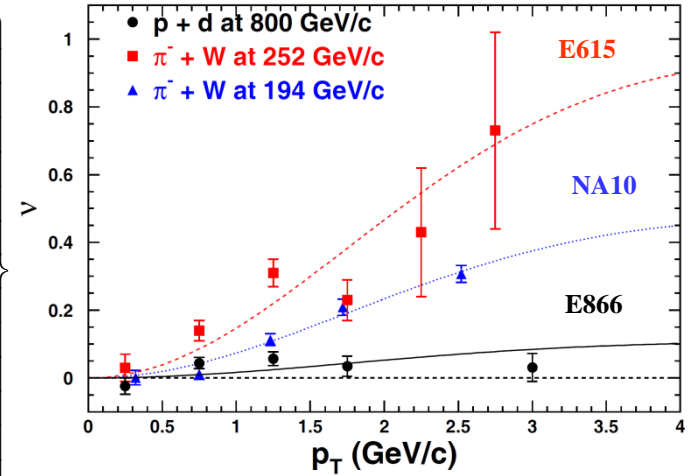


Boer-Mulders from Drell-Yan

S. Arnold, A. Metz and M. Schlegel, *Phys. Rev. D* **79**, 034005 (2009).

DY x-section (single-polarized)

$$\frac{d\sigma}{d\Omega} = \frac{\alpha_{em}^2}{Fq^2} \hat{\sigma}_U \left\{ \begin{array}{l} \left(1 + D_{[\sin 2\theta]} A_U^{\cos\varphi} \cos\varphi + D_{[\sin^2\theta]} A_U^{\cos 2\varphi} \cos 2\varphi \right) \\ + S_L \left(D_{[\sin 2\theta]} A_L^{\sin\varphi} \sin\varphi + D_{[\sin^2\theta]} A_L^{\sin 2\varphi} \sin 2\varphi \right) \\ \pm \left| \vec{S}_T \right| \left[\begin{array}{l} \left(D_{[1]} A_T^{\sin\varphi_S} + D_{[\cos^2\theta]} \tilde{A}_T^{\sin\varphi_S} \right) \sin\varphi_S \\ + D_{[\sin 2\theta]} \left(A_T^{\sin(\varphi+\varphi_S)} \sin(\varphi+\varphi_S) + A_T^{\sin(\varphi-\varphi_S)} \sin(\varphi-\varphi_S) \right) \\ + D_{[\sin^2\theta]} \left(A_T^{\sin(2\varphi+\varphi_S)} \sin(2\varphi+\varphi_S) + A_T^{\sin(2\varphi-\varphi_S)} \sin(2\varphi-\varphi_S) \right) \end{array} \right] \end{array} \right\}$$



$$\hat{\sigma}_U = (F_U^1 + F_U^2)(1 + A_U^1 \cos^2\theta); \quad D_{[f(\theta)]} = \frac{f(\theta)}{1 + A_U^1 \cos^2\theta}; \quad \lambda = A_U^1, \mu = A_U^{\cos\varphi}, \nu = 2A_U^{\cos 2\varphi} \propto h_{1q}^\perp \otimes h_{1\bar{q}}^\perp$$

D. Boer, *Phys. Rev. D* **60**, 014012 (1999)

“crude” model for the BM function of the pion from the fit of the NA10 data

$$h_1^{\perp a}(x, \mathbf{p}_T^2) = \frac{\alpha_T}{\pi} c_H^a \frac{M_C M_H}{\mathbf{p}_T^2 + M_C^2} e^{-\alpha_T \mathbf{p}_T^2} f_1(x)$$

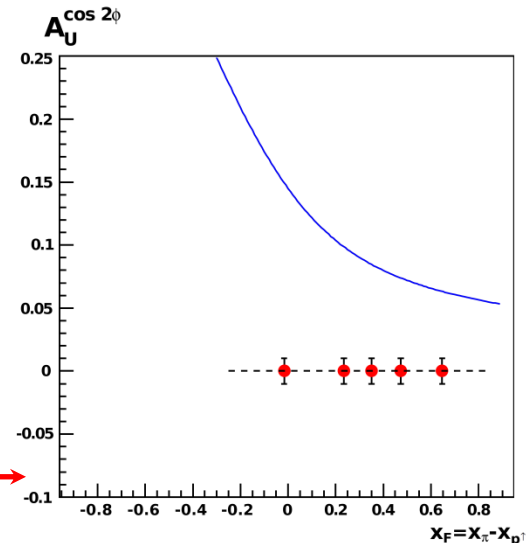
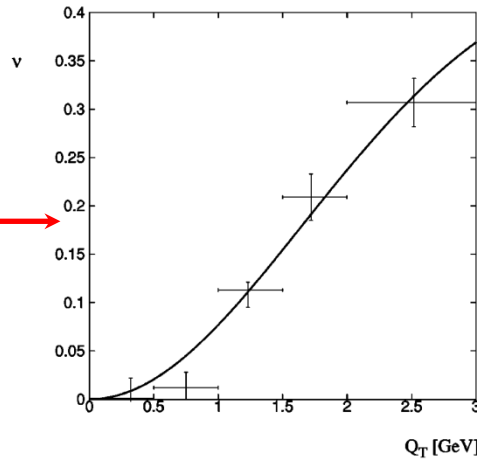
$$M_C = 2.3 \text{ GeV}, c_H^a = 1 \text{ and } \alpha_T = 1 \text{ GeV}^{-2},$$

+

The parameterization of the BM for the proton from

Zhang et al. *Phys. Rev. D* **77**, 054011 (2008)

For the $A_U^{\cos 2\varphi}$ at COMPASS

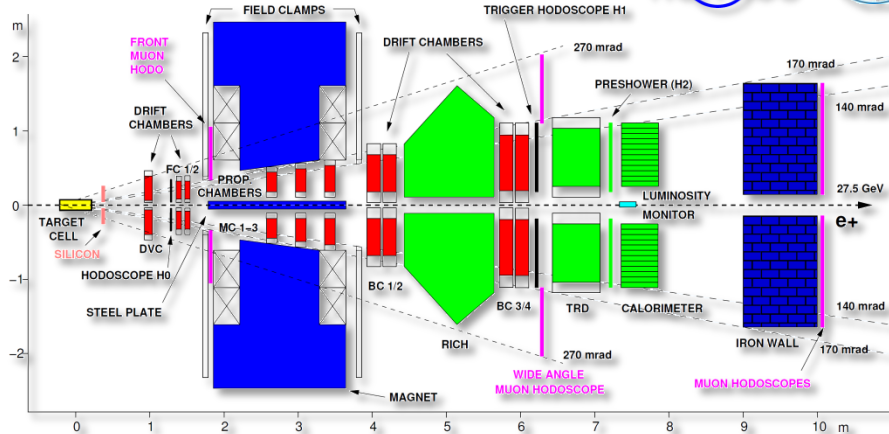


Experimental data: part II

Recent results

Experiments in last 35 years: part II

HERA MEasurement of Spin



Location: DESY, HERA

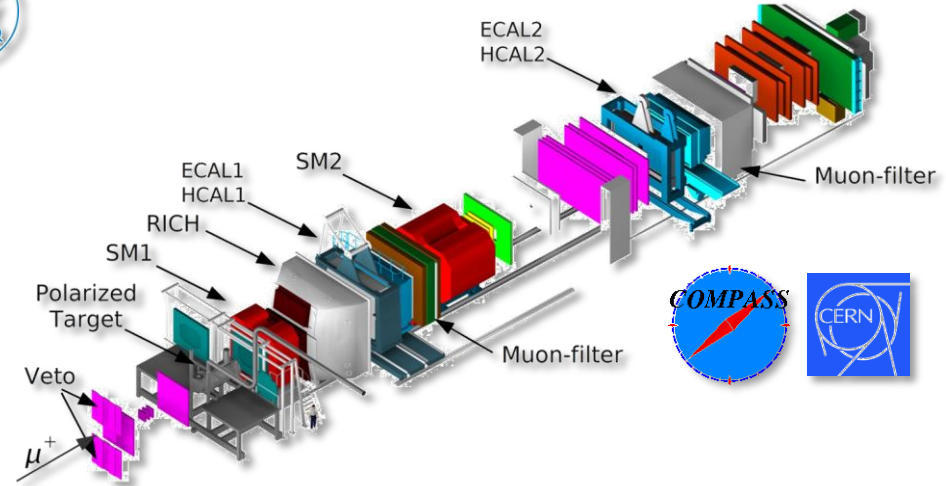
Beam: e^+/e^- , polarized (both helicity states) (<60%), 27.5 GeV

Target: Gaseous target (H/D)

- H/D Polarization (L & T) ~ 70-85%, $f \sim 1$
- Direct access to hydrogen or deuterium
- Fast spin reversal (<1s)
- Same acceptance for different polarization states
- single cell configuration
- Hydrogen - measurements only with transverse polarization
- Deuterium - both transverse and longitudinal polarization measurements



COMmon MUon PRoton Apparatus for Structure and Spectroscopy



Location: CERN SPS North Area. (2-stage spectrometer LAS-SAS)

Beam: μ^+ , longitudinally polarized (~80%), 160 GeV

Target: Solid state target (${}^6\text{LiD}$ or NH_3)

- ${}^6\text{LiD}$ Polarization (L & T) ~ 50%, $f \sim 0.38$
- NH_3 Polarization (L & T) ~ 80%, $f \sim 0.14$
- 2-cell target configuration for ${}^6\text{LiD}$ and 3-cell for NH_3
- Neighboring cells are polarized in opposite directions
- Data is collected simultaneously for the two target spin orientations
- Spin reversal after each ~4-5 days
- Such a construction allows to reduce systematic effects due to the acceptance



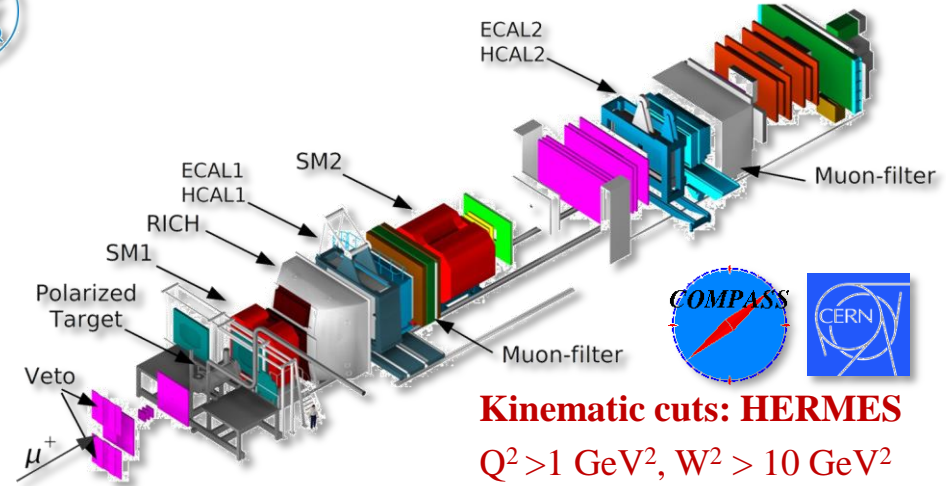
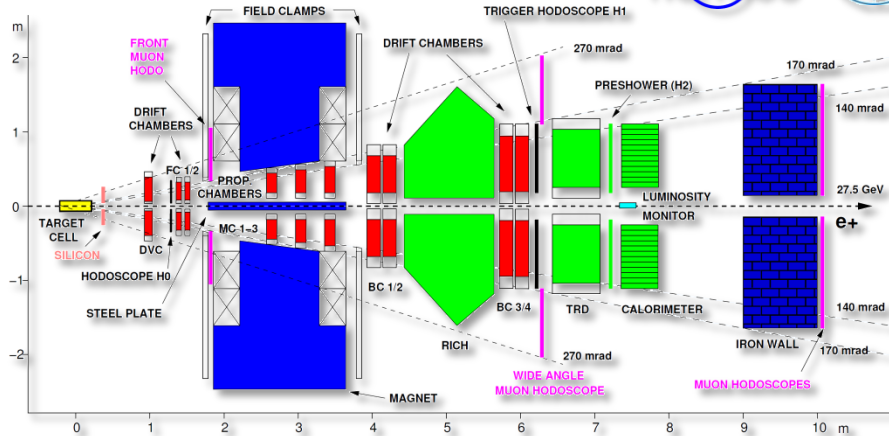
Both collaborations have put a substantial effort in the study of possible acceptance effects

Experiments in last 35 years: part II

HERA MEasurement of Spin



COMmon MUon PRoton Apparatus for STRUCTure and SPectroscopy



Kinematic cuts: HERMES

$Q^2 > 1 \text{ GeV}^2$, $W^2 > 10 \text{ GeV}^2$
 $0.023 < x < 0.6$, $0.2 < y < 0.85$
 $z > 0.2$ and $x_F > 0.2$

Pions $1 \text{ GeV} < P_h < 15 \text{ GeV}$
 Kaons $2 \text{ GeV} < P_h < 15 \text{ GeV}$

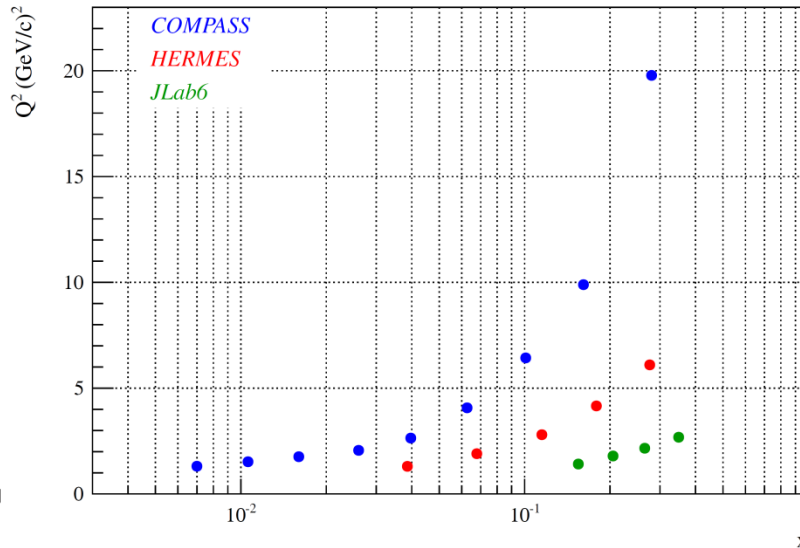
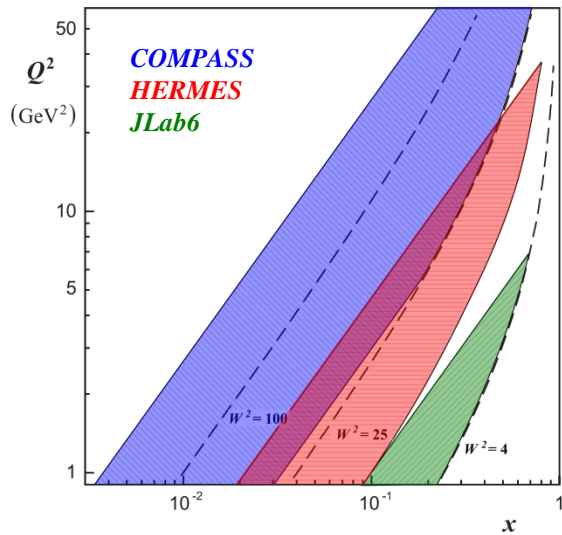
Kinematic cuts: COMPASS

$Q^2 > 1 \text{ GeV}^2$, $W^2 > 25 \text{ GeV}^2$
 $\theta_{\gamma}^{\text{lab}} < 0.06$
 $0.003 < x < 0.13$, $0.2 < y < 0.9$
 $0.2 < z < 0.85$

$0.1 < P_{hT} < 1 \text{ GeV}/c$

Kinematic cuts: JLab

$Q^2 > 1 \text{ GeV}^2$, $W^2 > 4 \text{ GeV}^2$
 $0.14 < x < 0.48$, $0.4 < y < 0.7$
 $0.4 < z < 0.7$



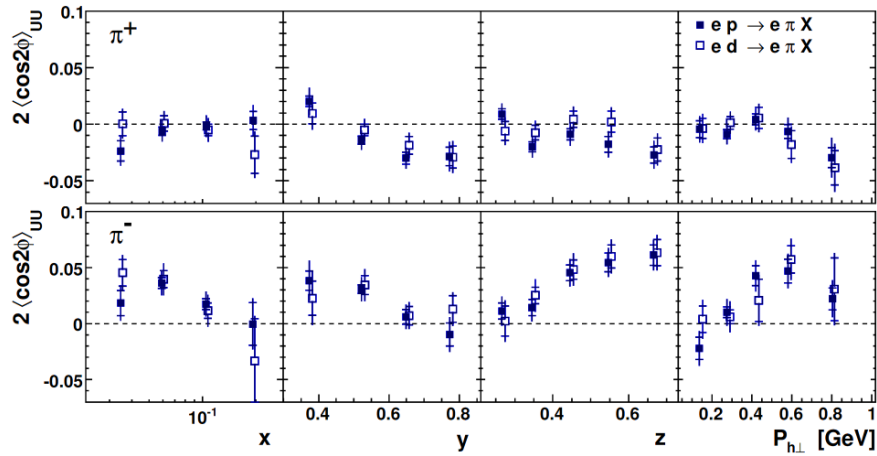
Experimental data: part II

HERMES results

(Proton, Deuteron, π^+ / π^- , K^+ / K^- , h^+ / h^-)

$A_{UU}^{\cos 2\phi}$ -amplitude on p & d: pions

HERMES Collaboration, Phys. Rev. D 87, 012010 (2013)



← Zero or negative for π^+

← Positive amplitudes for π^-

increase with P_{hT}

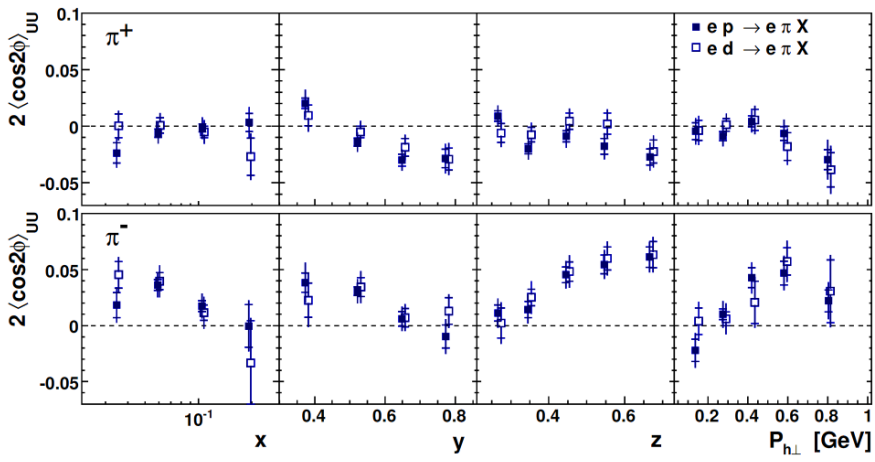
Data is available at:

Durham HEP database, <http://durpdg.dur.ac.uk>

INSPIRE, <http://inspirebeta.net/record/1111237/>

$A_{UU}^{\cos 2\phi}$ -amplitude on p & d: pions

HERMES Collaboration, Phys. Rev. D 87, 012010 (2013)



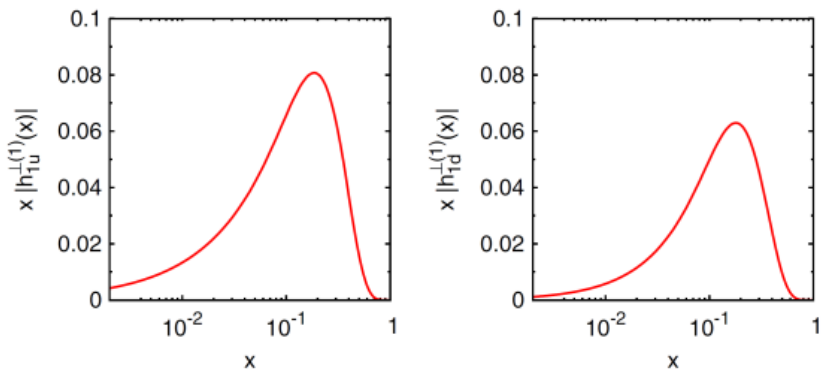
← Zero or negative for π^+
 ← Positive amplitudes for π^-

increase with P_{hT}

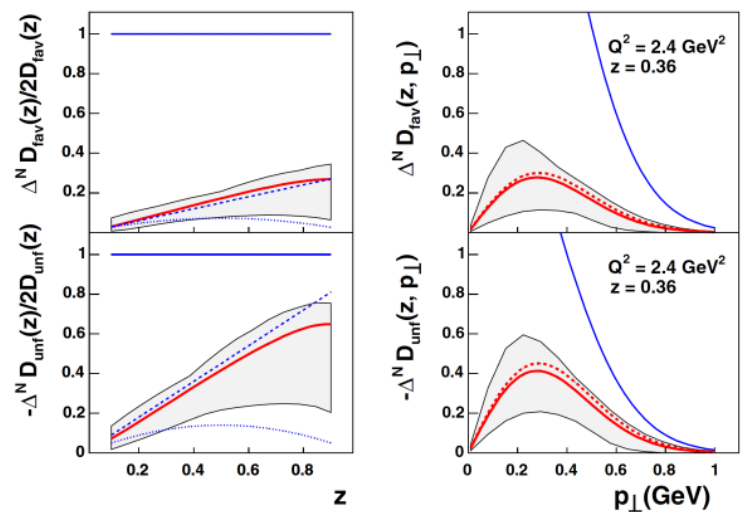
$$A_{UU}^{\cos 2\phi_h} \propto -h_1^{\perp q} \otimes H_{1q}^{\perp h} + \left(\frac{M}{Q}\right)^2 f_1^q \otimes D_{1q}^h + \dots$$

Opposite sign for favored - disfavored Collins FF

Similarity of H & D \leftrightarrow BM PDF has same sign for u and d quarks

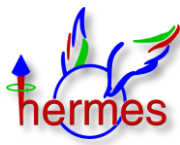


Barone, Melis and Prokudin Phys. Rev. D 81, 114026 (2010)

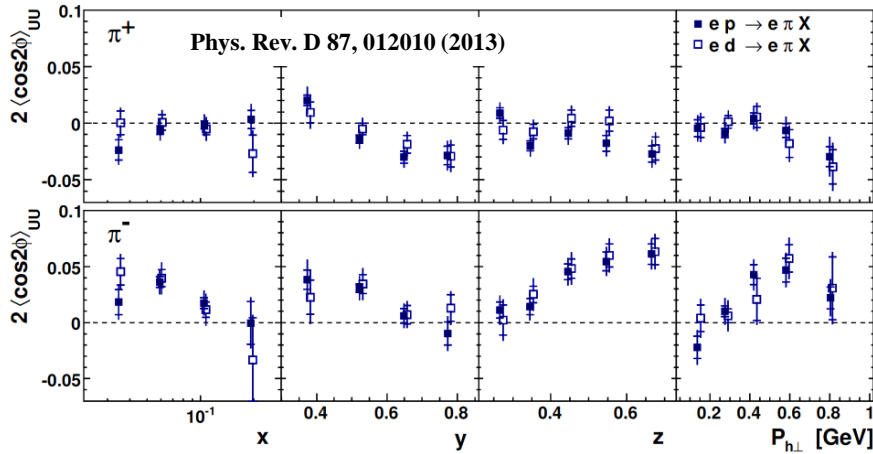


Anselmino et al. Phys. Rev. D 75, 054032 (2007)

$A_{UU}^{\cos 2\phi}$ -amplitude on p & d: pions



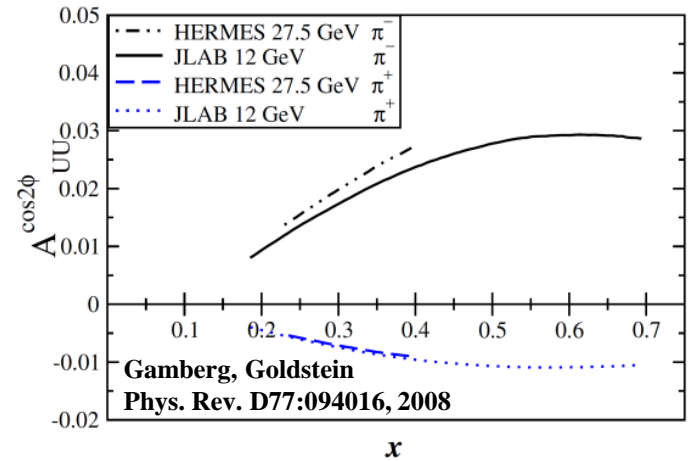
HERMES Collaboration, Phys. Rev. D 87, 012010 (2013)



← Zero or negative for π^+

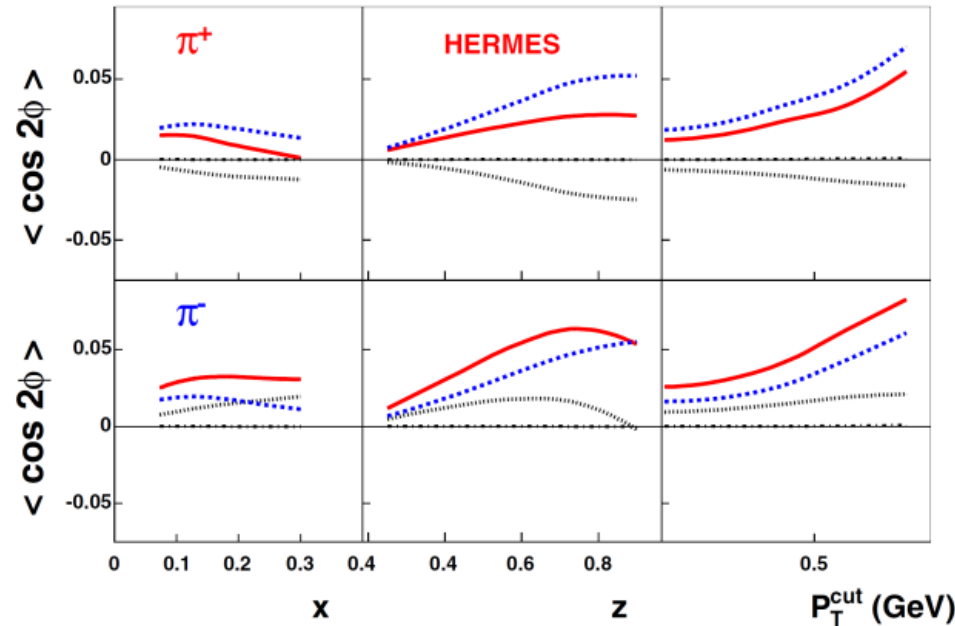
← Positive amplitudes for π^-

increase with P_{hT}

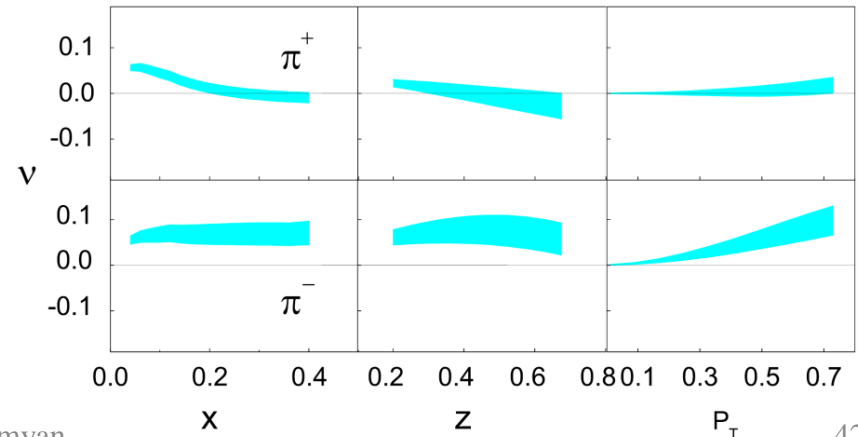


Gamberg, Goldstein
Phys. Rev. D77:094016, 2008

Barone, Prokudin, and Ma, Phys. Rev. D 78, 045022 (2008).



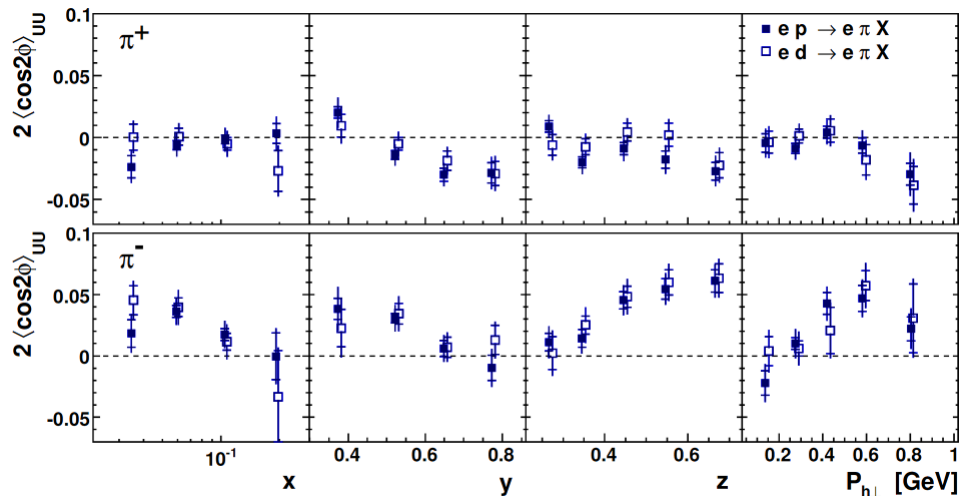
Zhang et al. Phys. Rev. D78:034035, 2008



$A_{UU}^{\cos 2\phi}$ -amplitude on p & d: pions, kaons

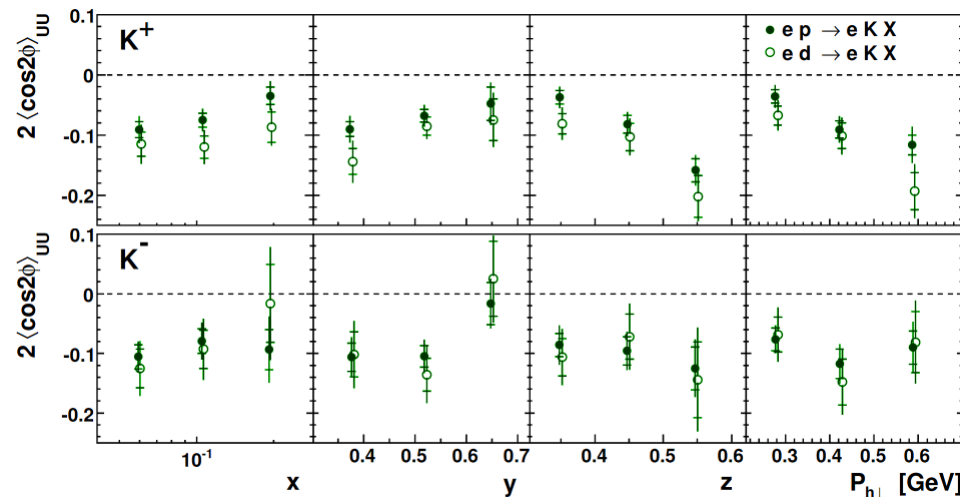


HERMES Collaboration, Phys. Rev. D 87, 012010 (2013)



- Large amplitude \leftrightarrow non-zero BM PDF (spin effect)
- Opposite sign for $\pi^+/\pi^- \leftrightarrow$ opposite sign for favored - disfavored Collins FF
- Similarity of H & D $\leftrightarrow h_1^{\perp,u} \approx h_1^{\perp,d}$

$$A_{UU}^{\cos 2\phi_h} \propto -h_1^{\perp q} \otimes H_{1q}^{\perp h} + \left(\frac{M}{Q}\right)^2 f_1^q \otimes D_{1q}^h + \dots$$



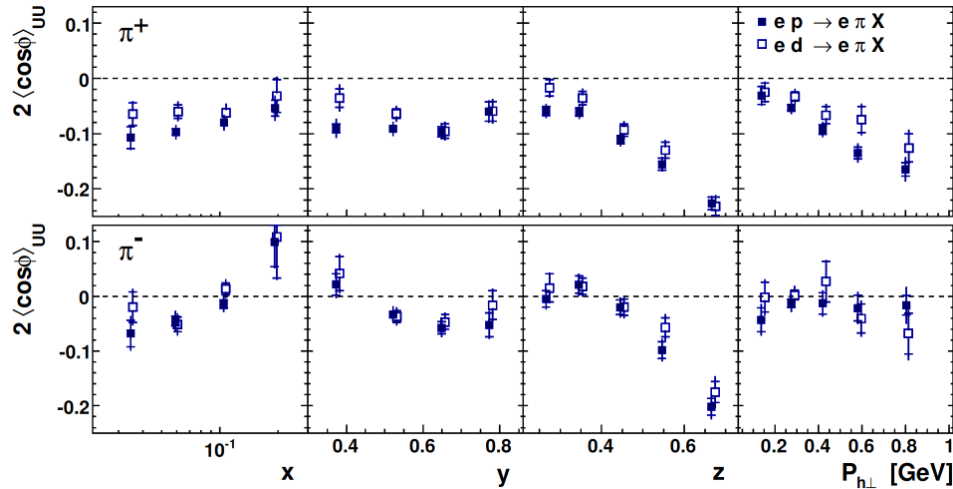
- K^+/K^- amplitudes are larger than π^+/π^-
- Different trends, but same sign as for π^+/π^-
- K^+ - u -dominance (same sign with π^+)
- K^- - fully sea object
- Collins FF for kaons – unknown
- As well as strange quark contribution

Kinematic range Pions (all hadrons)			
x	y	z	$P_{h\perp}$
0.023–0.27	0.3–0.85	0.2–0.75	0.05–1.0
Kinematic range Kaons			
x	y	z	$P_{h\perp}$
0.042–0.27	0.3–0.7	0.2–0.6	0.2–0.7

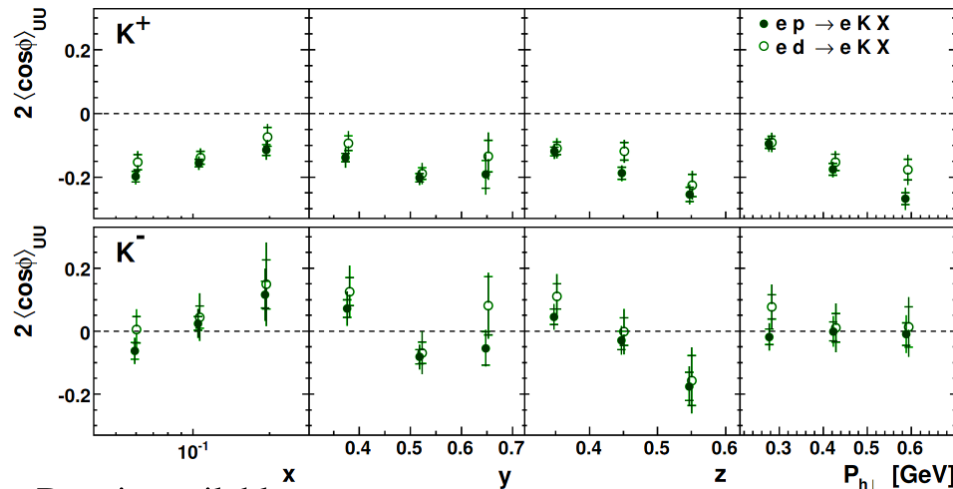
Data is available at:

Durham HEP database, <http://durpdg.dur.ac.uk>
INSPIRE, <http://inspirebeta.net/record/1111237/>

$A_{UU}^{\cos\phi}$ -amplitude on p & d: pions, kaons

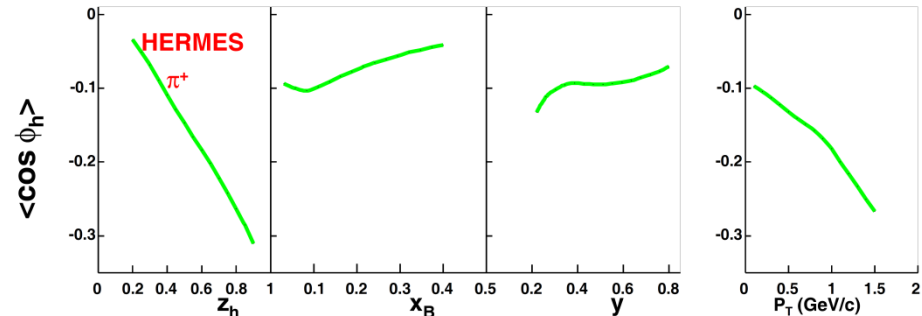


$$A_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} \left\{ -f_1^q \otimes D_{1q}^h - h_1^{\perp q} \otimes H_{1q}^{\perp h} \right\}$$



HERMES Collaboration, Phys. Rev. D 87, 012010 (2013)

- Negative amplitudes
- π^+/π^- difference \leftrightarrow due to the BM effect (Cahn expected to be flavor blind)
- Predictions for Cahn only are much larger BM contribution...



Curves for Cahn contribution only
Anselmino et al. Eur. Phys. J. A 31, 373-381 (2007)

- K^+ amplitudes are larger than π^+ contribution from BM which is large for K^+
- K^- amplitudes are compatible with zero while K^- BM is large... non trivial

Data is available at:

Durham HEP database, <http://durpdg.dur.ac.uk>
INSPIRE, <http://inspirebeta.net/record/1111237/>

23 July 2013

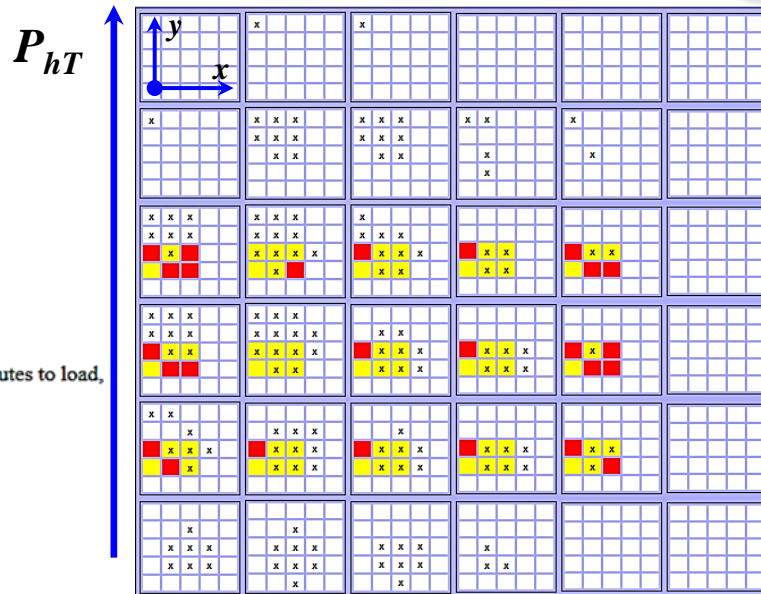
Bakur Parsamyan

Kinematic range Pions (all hadrons)			
x	y	z	P_{hT}
0.023–0.27	0.3–0.85	0.2–0.75	0.05–1.0
Kinematic range Kaons			
x	y	z	P_{hT}
0.042–0.27	0.3–0.7	0.2–0.6	0.2–0.7

Multi-dimensional analysis



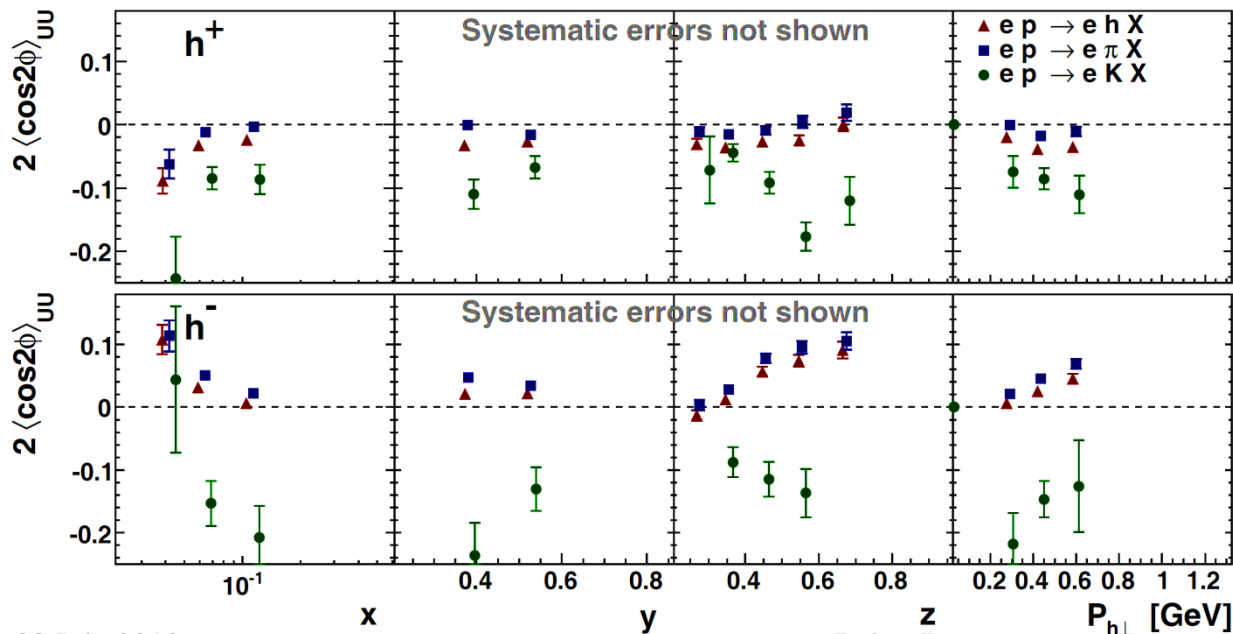
<http://www-hermes.desy.de/cosnphi/>



- x:
- y:
- z:
- pt:

Please enable pop-ups
Results may take several minutes to load,
please do not refresh

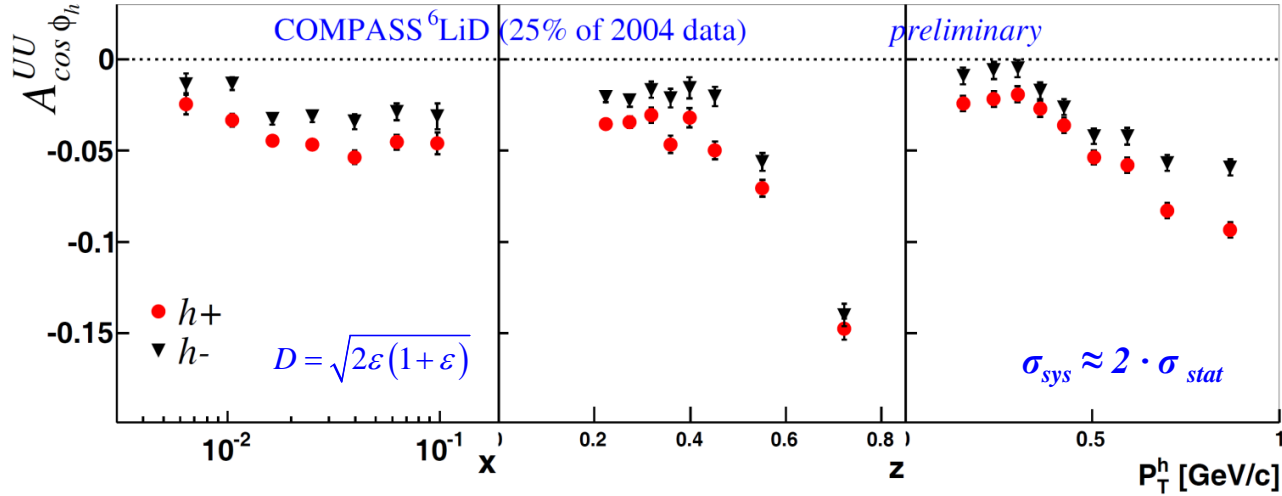
Binning							
900 kinematic bins x 12 ϕ -bins							
Variable	Bin limits						#
x	0.023	0.042	0.078	0.145	0.27	0.6	5
y	0.2	0.3	0.45	0.6	0.7	0.85	5
z	0.2	0.3	0.4	0.5	0.6	0.75	1
P_{hT}	0.05	0.2	0.35	0.5	0.7	1	1.3



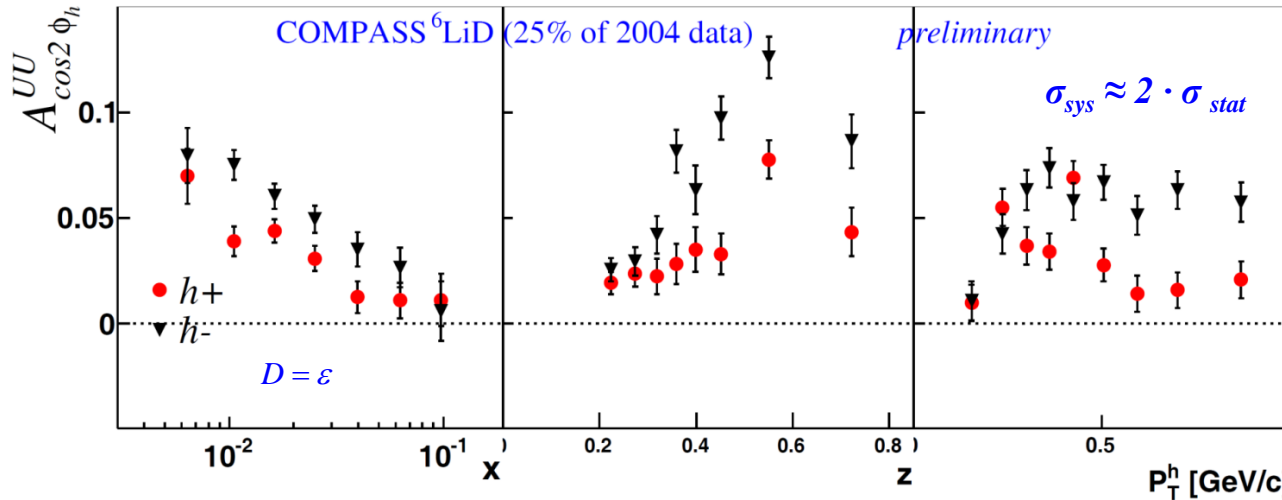
Experimental data: part II

COMPASS results (Deuteron, h^+/h^-)

$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos 2\phi}$ amplitudes h^+/h^-



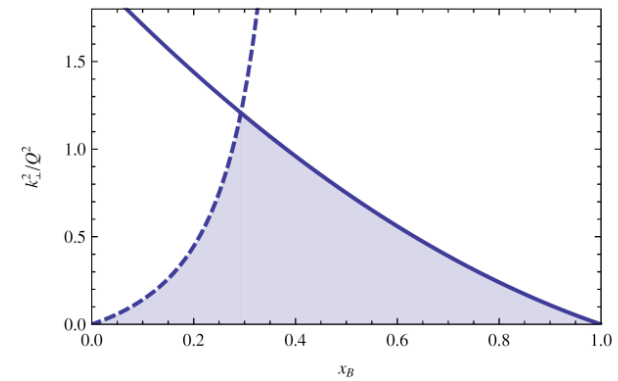
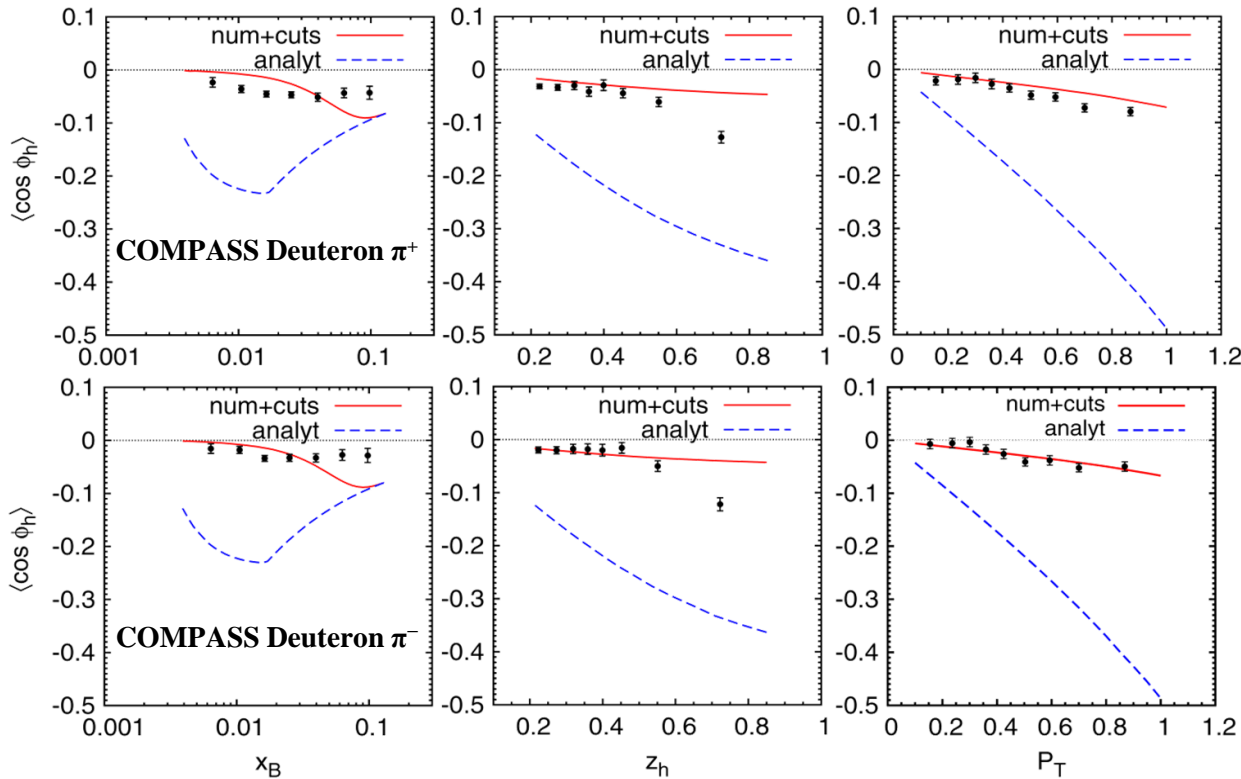
- Negative amplitudes h^+/h^-
- Clear differences between h^+/h^-
 - Larger amplitude for h^+



- Positive amplitudes h^+/h^-
- Clear differences between h^+/h^-
 - Larger amplitude for h^-

$A_{UU}^{\cos\phi}$ -amplitude: comparison with theory

M. Boglione, S. Melis, and A. Prokudin, Phys. Rev. D 84, 034033 (2011)



$$f_{q/p}(x, k_{\perp}) = f_{q/p}(x) \frac{1}{1 - e^{-(k_{\perp}^{\max})^2 / \langle k_{\perp}^2 \rangle}} \frac{e^{-k_{\perp}^2 / \langle k_{\perp}^2 \rangle}}{\pi \langle k_{\perp}^2 \rangle}$$



Description improves a lot

- 1) the energy of the parton to be less than the energy of the parent hadron

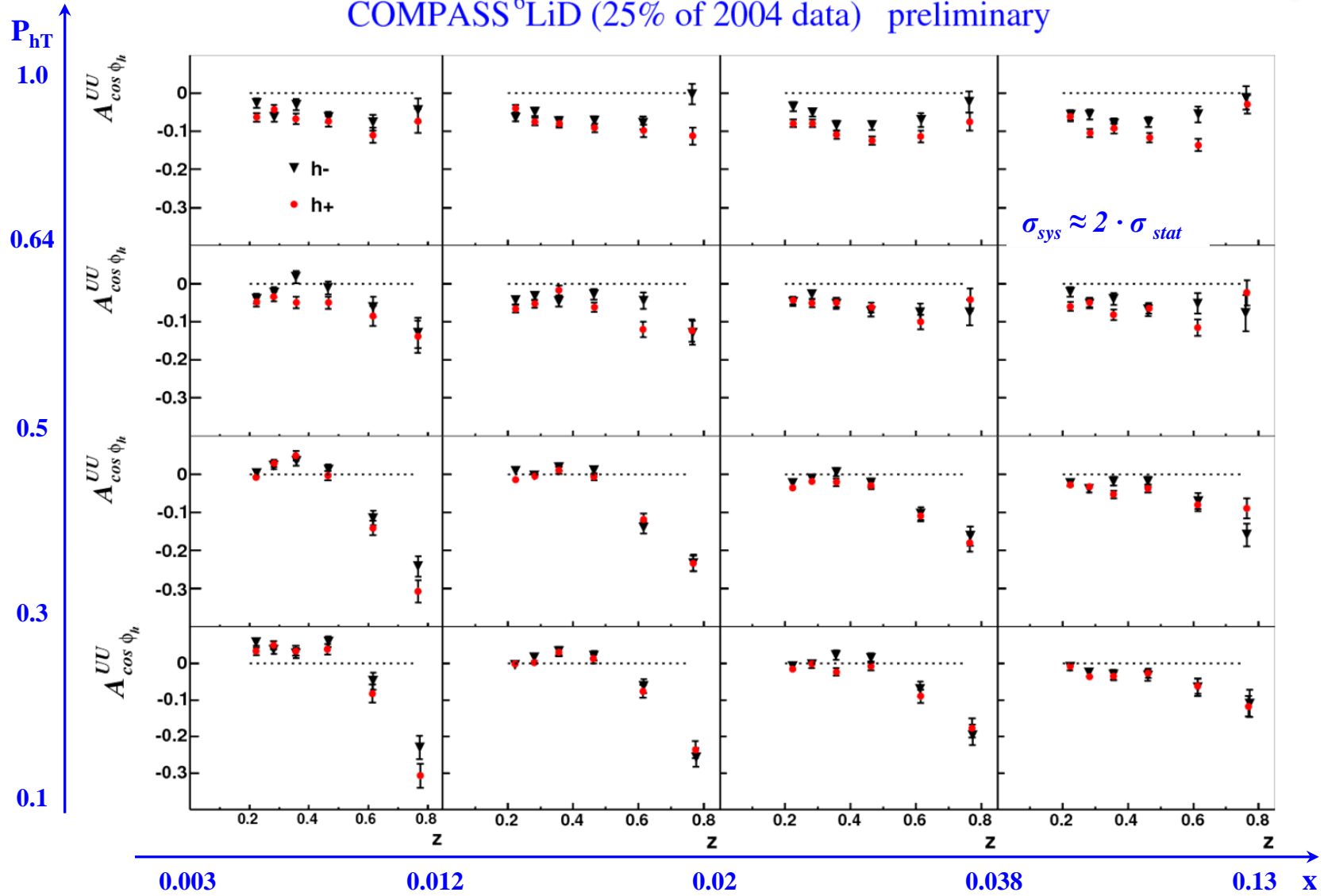
$$k_{\perp}^2 \leq (2 - x_B)(1 - x_B)Q^2, \quad 0 < x_B < 1.$$

- 2) the parton to move in the forward direction with respect to the parent hadron

$$k_{\perp}^2 \leq \frac{x_B(1 - x_B)}{(1 - 2x_B)^2} Q^2, \quad x_B < 0.5.$$

$A_{UU}^{\cos\phi_h}$ - asymmetry (z - dependence)

COMPASS⁶LiD (25% of 2004 data) preliminary

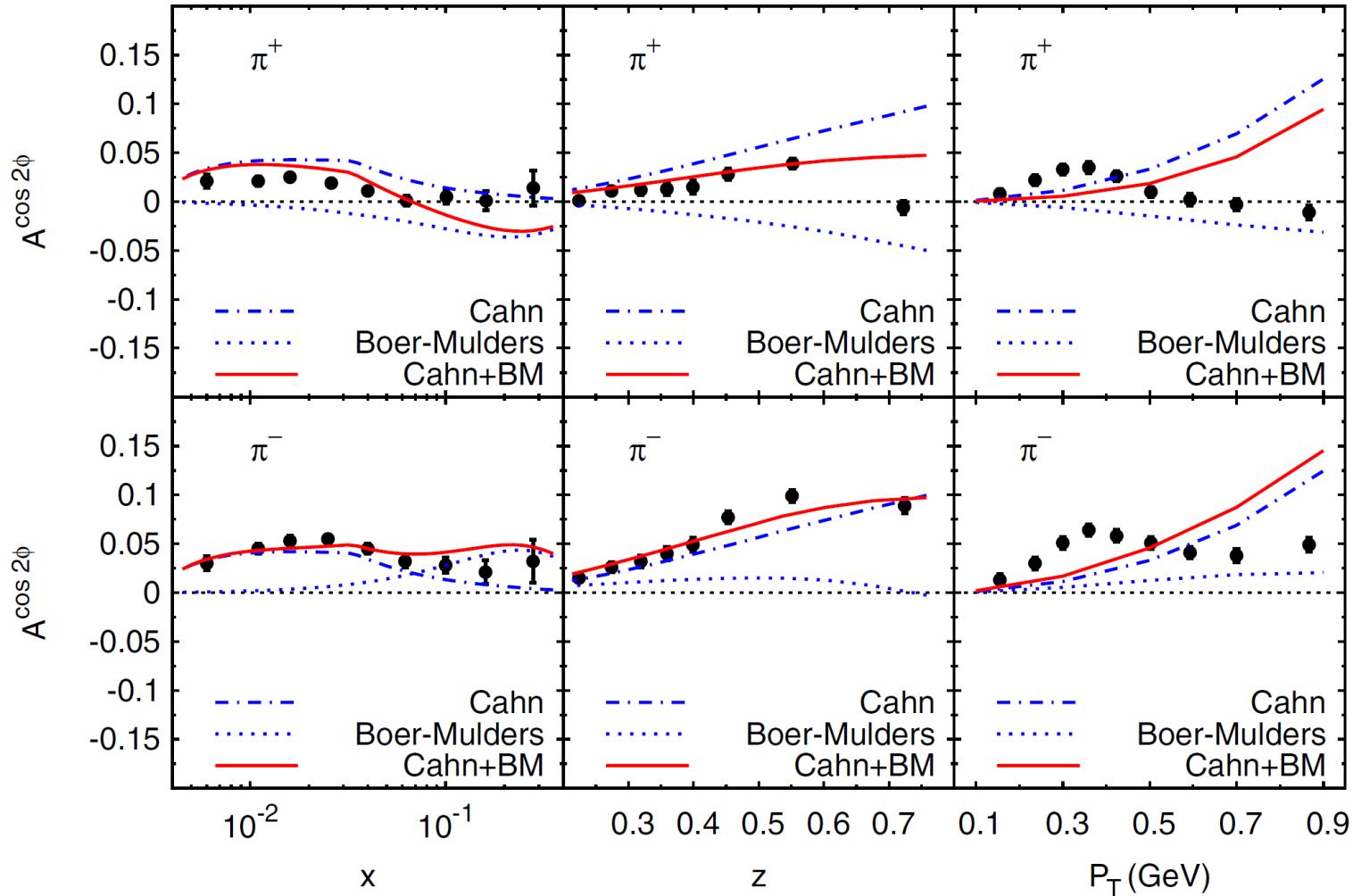


z strong dependence more evident at small x and small P_{hT}

$A_{UU}^{\cos 2\phi}$ -amplitude

V. Barone, S. Melis and A. Prokudin, Phys. Rev. D 81, 114026 (2010)

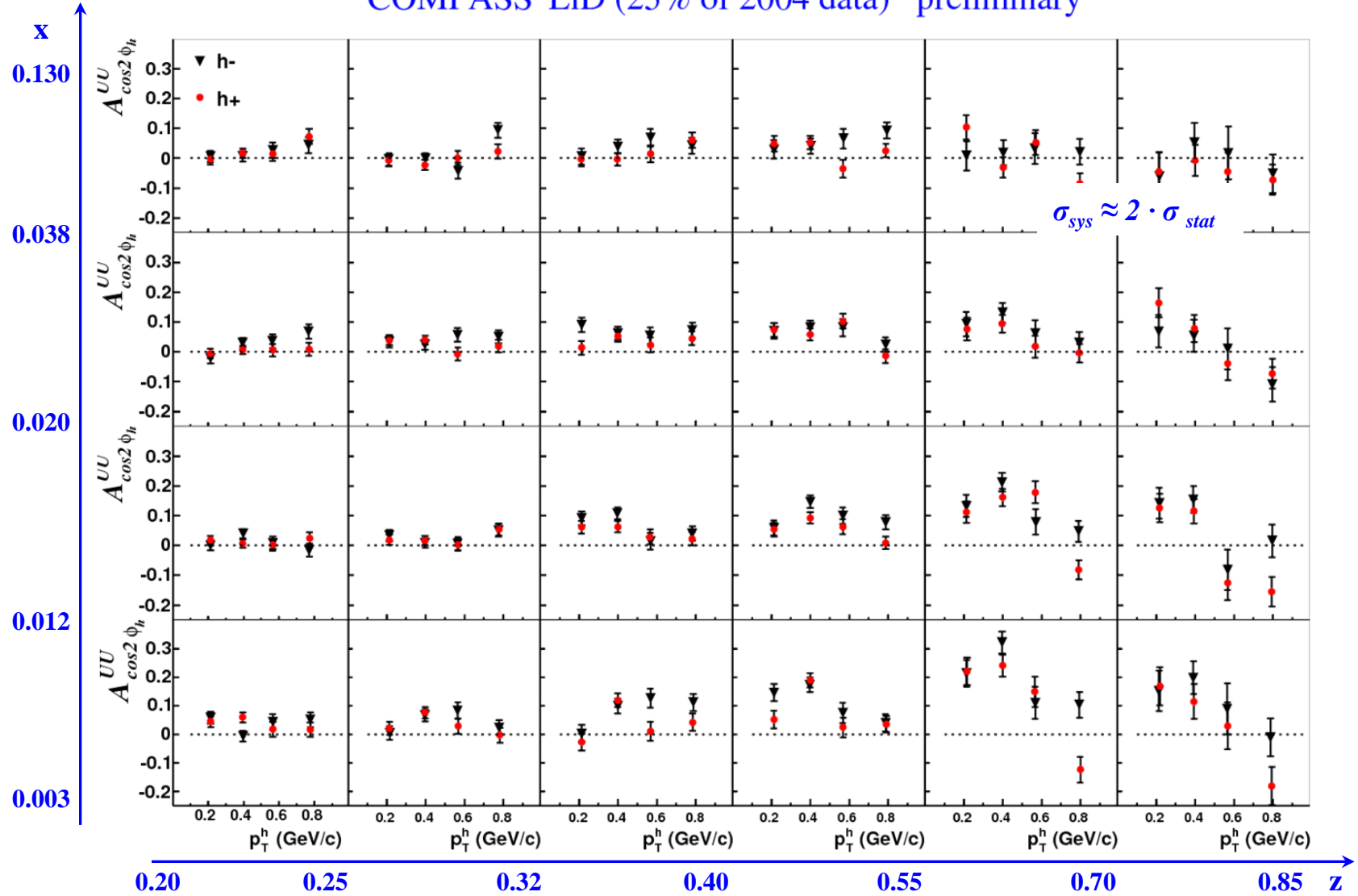
COMPASS Deuteron



P_{hT} dependence difficult to reproduce

$A_{UU}^{\cos 2\phi_h}$ - asymmetry (P_{hT} - dependence)

COMPASS⁶LiD (25% of 2004 data) preliminary



P_{hT} trend not described by the models arises in large z and low x region

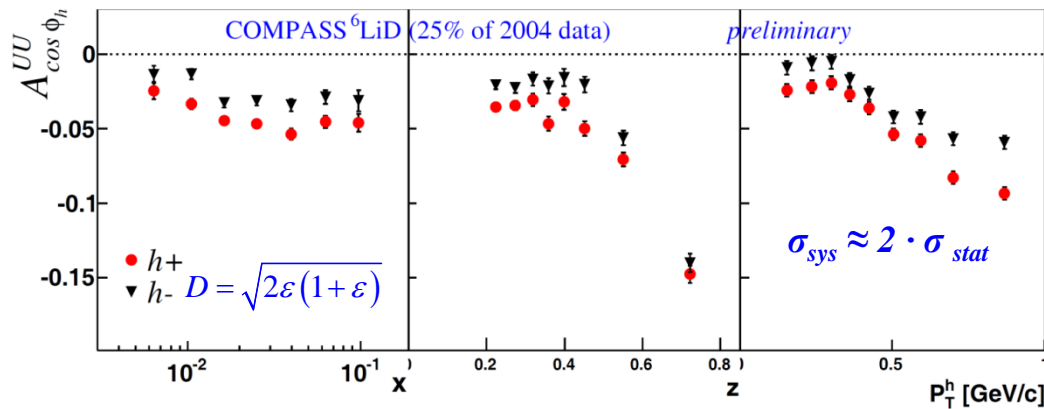
Experimental data: part II

Comparison of COMPASS and HERMES results

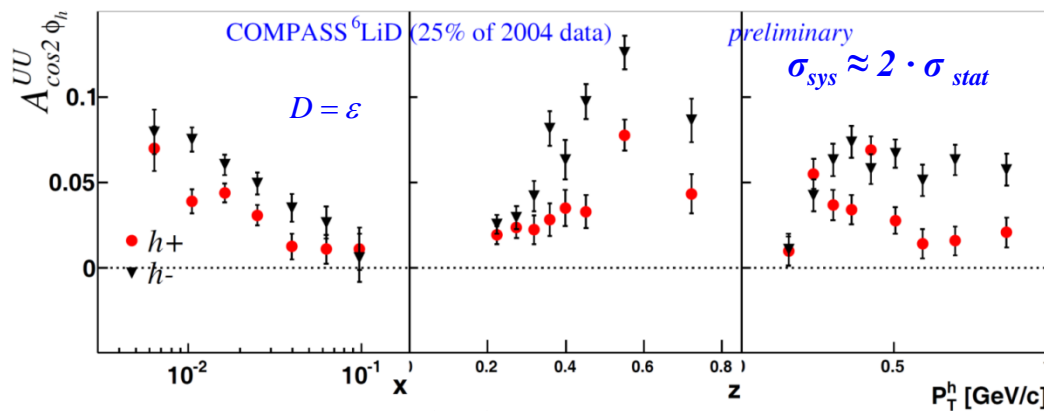
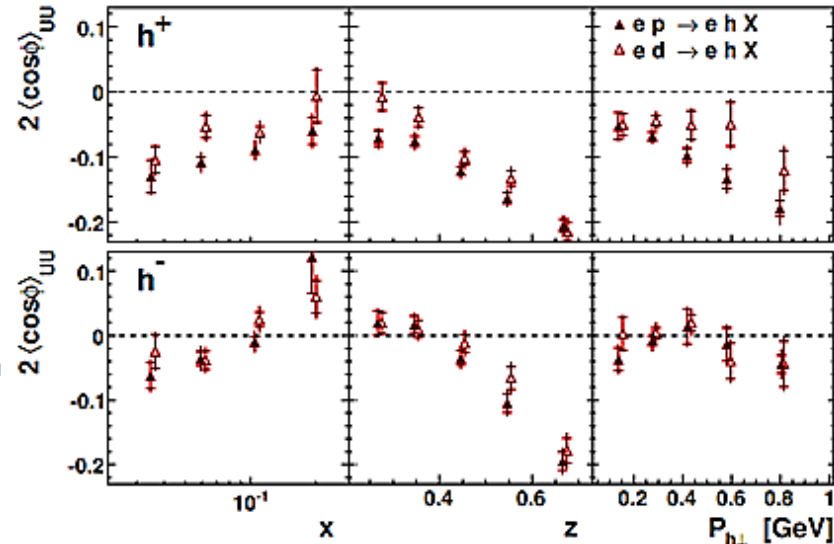
$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos 2\phi}$ amplitudes h^+/h^-



Different kinematic regions!

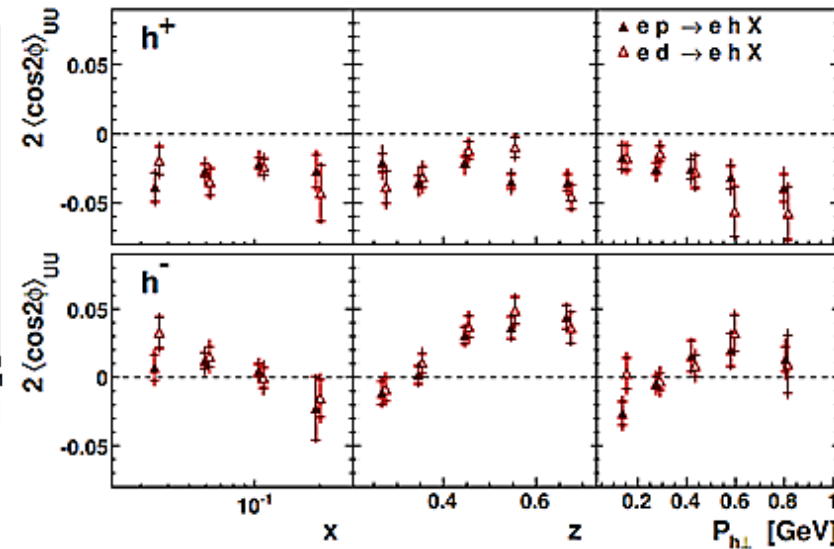


➤ Similar trends for h^+/h^-



➤ Similar trends for h^+/h^-

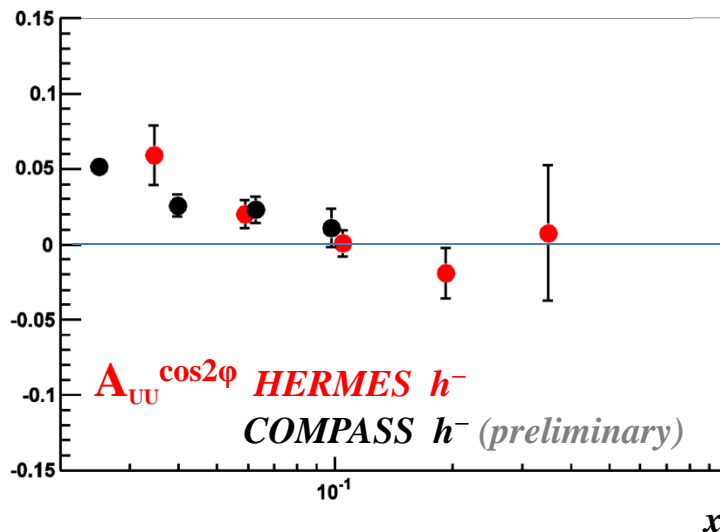
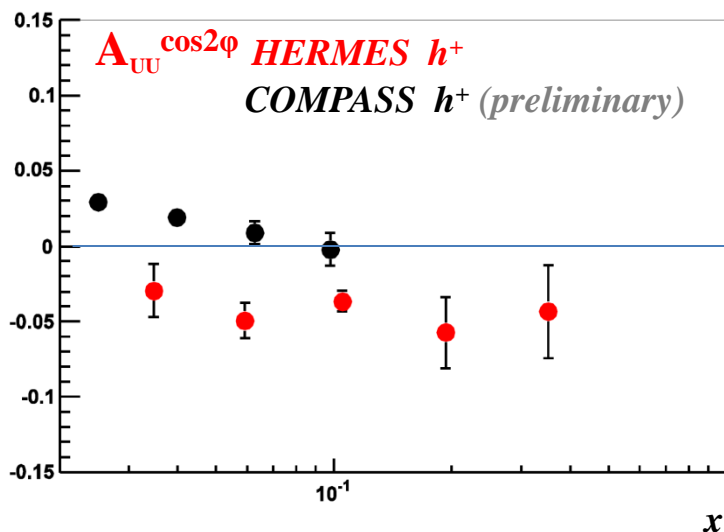
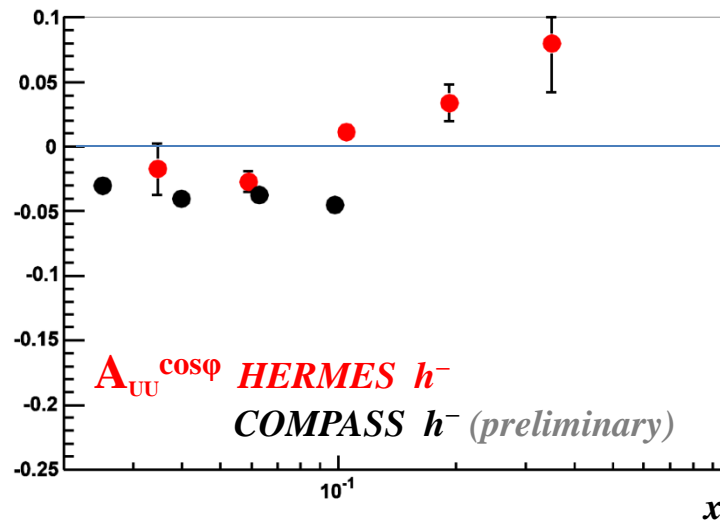
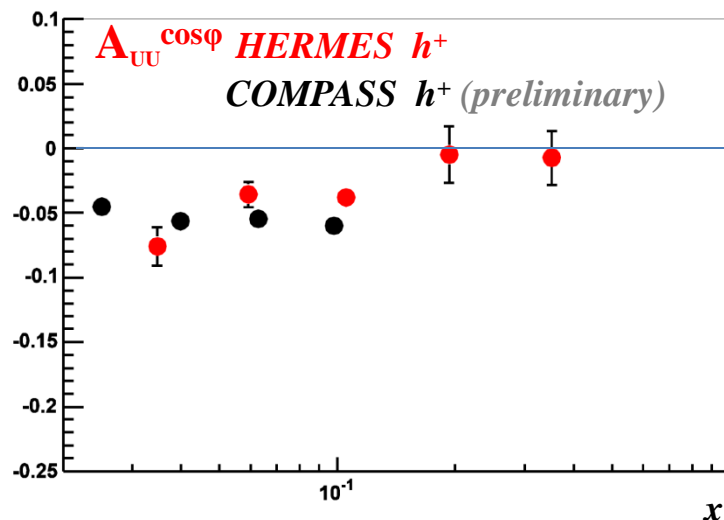
➤ No sign change for h^+/h^- at COMPASS



$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos2\phi}$ amplitudes h^+/h^-



Shrinking to the same range

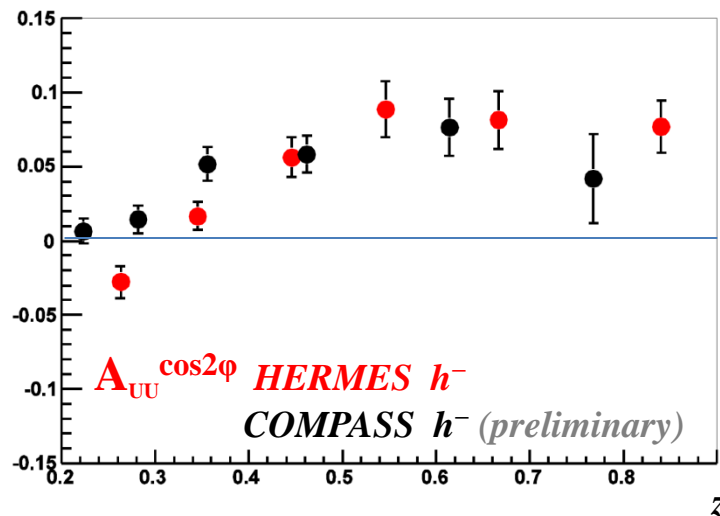
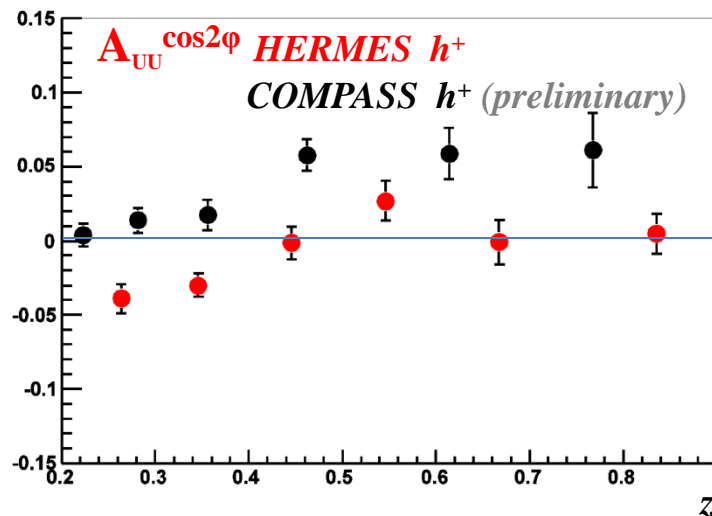
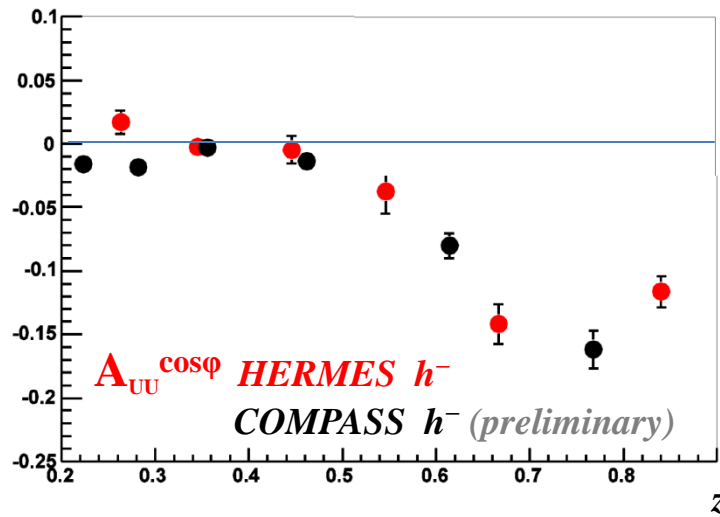
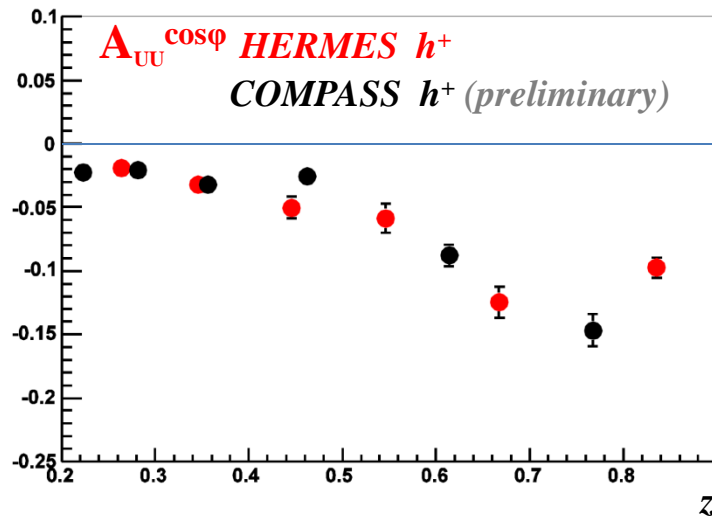


Selected range for COMPASS-HERMES: $0.2 < z < 1$; $0.05 < P_{hT} < 1$

$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos2\phi}$ amplitudes h^+/h^-



Shrinking to the same range



COMPASS $0.02 < x < 0.13$ ($\langle Q^2 \rangle \approx 4$); HERMES $0.023 < x < 0.145$; ($\langle Q^2 \rangle \approx 2$)
COMPASS-HERMES $0.05 < P_{hT} < 1$

Longitudinal Cahn effect

$A_{LL}^{\cos(\phi_h)}$ “Longitudinal” Cahn effect

Bacchetta, Diehl, Goeke, Metz, Mulders and Schlegel JHEP 0702:093 (2007).

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \cos\phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos\phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \dots + S_L \lambda \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h \times A_{LL}^{\cos\phi_h} + \dots$$

Longitudinal Cahn effect

Kotzinian et al. Phys. Rev. D 74, 074015 (2006)

$$F_{LL}^{\cos\phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot \mathbf{k}_T}{M} g_{1L}^q D_{1q}^h \right\}$$

$A_{LL}^{\cos(\phi_h)}$ “Longitudinal” Cahn effect

$$\frac{d\sigma}{dx dy dz dP_{hT}^2 d\phi_h d\psi} = \left[\frac{\alpha}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \right] \times (F_{UU,T} + \varepsilon F_{UU,L}) \times$$

$$1 + \cos\phi_h \times \sqrt{2\varepsilon(1+\varepsilon)} A_{UU}^{\cos\phi_h} + \cos(2\phi_h) \times \varepsilon A_{UU}^{\cos(2\phi_h)} + \dots + S_L \lambda \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h \times A_{LL}^{\cos\phi_h} + \dots$$

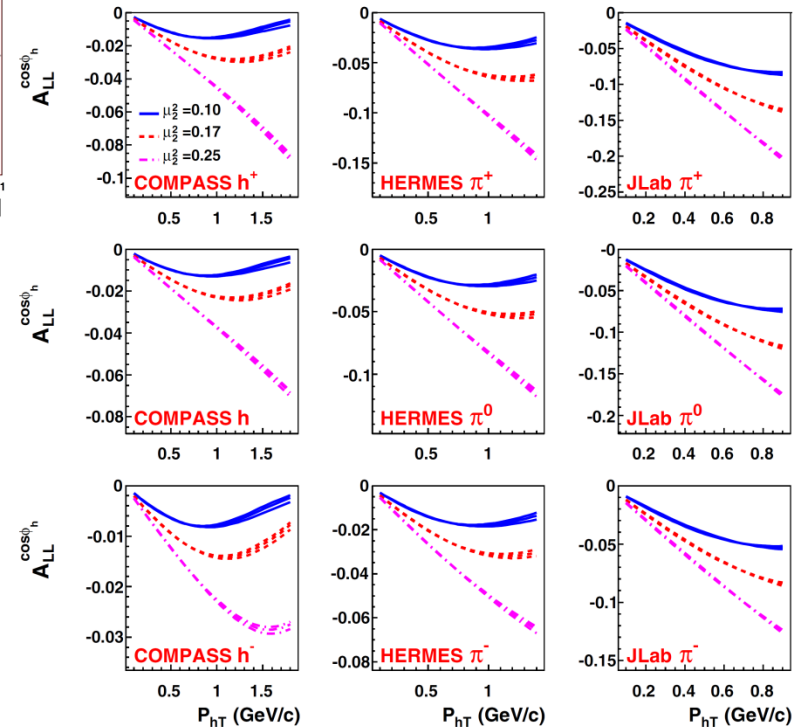
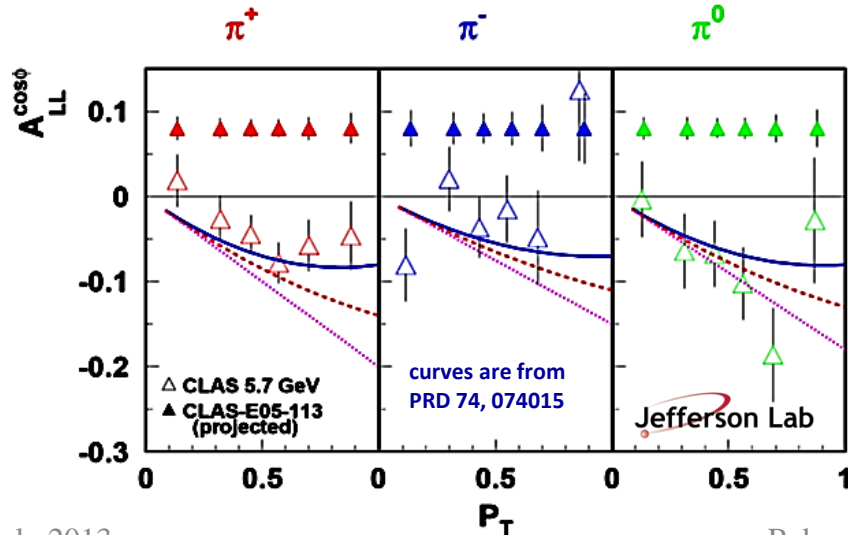
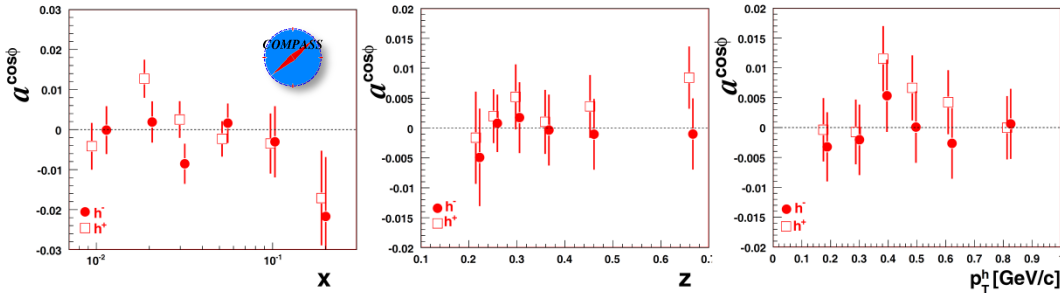
Longitudinal Cahn effect

Kotzinian et al. Phys. Rev. D 74, 074015 (2006)

$$F_{LL}^{\cos\phi_h} = \frac{2M}{Q} C \left\{ -\frac{\hat{h} \cdot k_T}{M} g_{1L}^q D_{1q}^h \right\}$$

$$g_{1L}^q(x, k_T) = g_1^q(x) \frac{1}{\pi\mu_2^2} \exp\left(-\frac{k_T^2}{\mu_2^2}\right)$$

COMPASS Collaboration
Eur.Phys.J.C70:39-49,2010



Summary

Experiment

- In past 35 years a lot of data accumulated from different experiments
- Both $A_{UU}^{\cos\varphi}$ and $A_{UU}^{\cos 2\varphi}$ are not zero!
- Interesting behavior for pion and kaon asymmetries from HERMES
- Trends measured at HERMES are in general confirmed by COMPASS
- Multidimensional approach
 - **The ball is on the “theoretical” side of the court!**
- Promising future measurements JLab12, COMPASS (LH), EIC..

Theory

- In past years, in parallel to the experimental efforts, a lot of theoretical and phenomenological studies
- Many attempts for extractions and fits
- Intriguing predictions for future measurements

Thank you!

While “googleing” for Robert Cahn’s photo...

2001 UEFA Champions League Final
Bayern Munich  1–1  Valencia
penalties 5–4

last penalty
4-4 goes to 5-4

Transverse component

Man of the Match:
Oliver Kahn (Bayern Munich)
During his career:
139 penalties 34 (24,46 %) saved.

It was not a SIDIS reaction and it had nothing to do with the “Cahn-effect”, but an educated guess about the size of the transverse component played a crucial role!

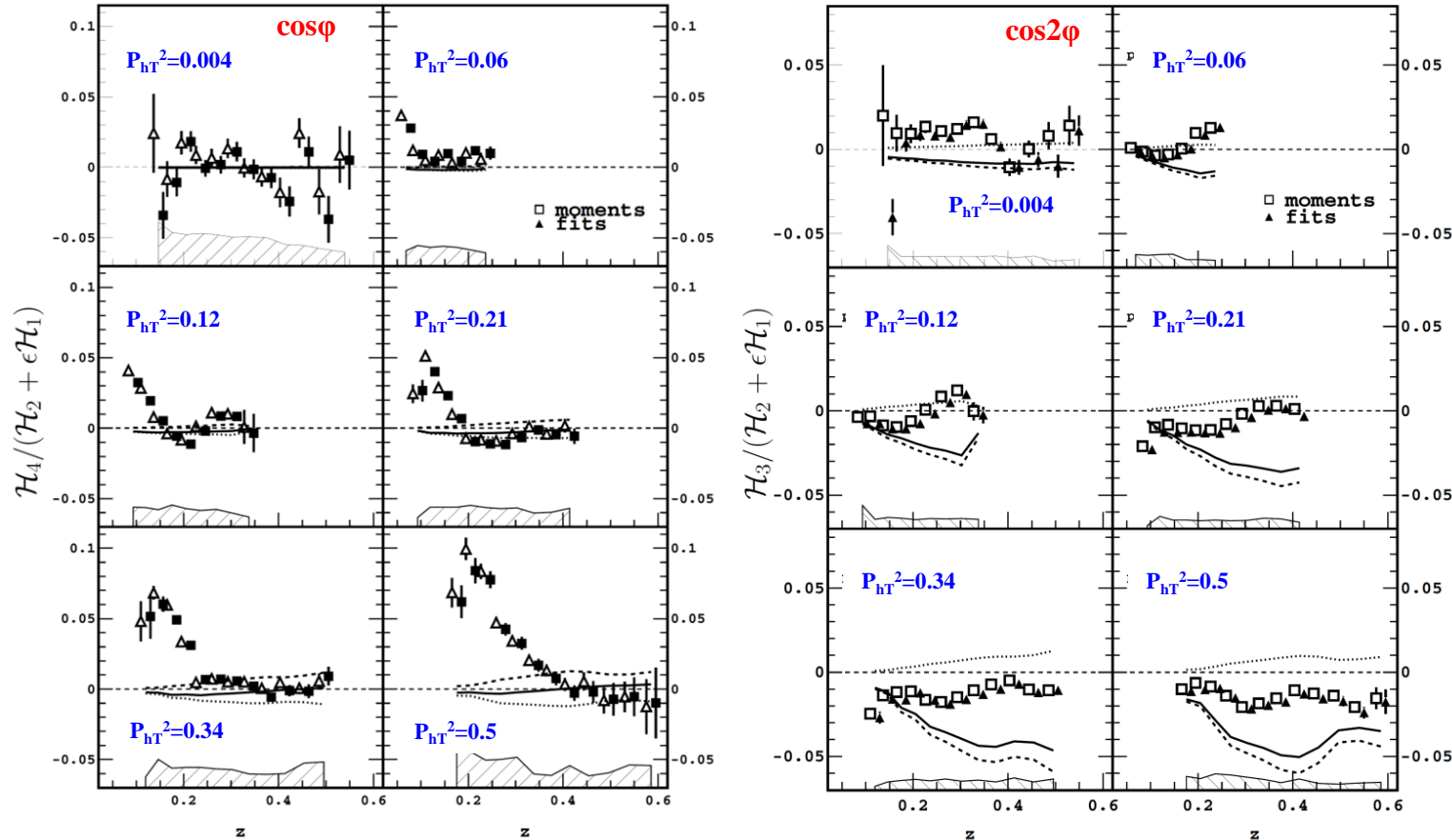
Spare slides

CLAS (Jlab hall B) results z-dependence

M. Osipenko et al. (CLAS Collaboration)

Phys.Rev.D80:032004,2009

Positive pions



Theoretical predictions: Cahn effect + Berger effect

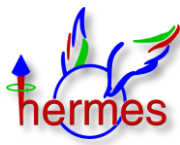
R. N. Cahn, Phys. Rev. D40, 3107 (1989).

M. Anselmino et al., Phys. Rev. D71, 074006 (2005).

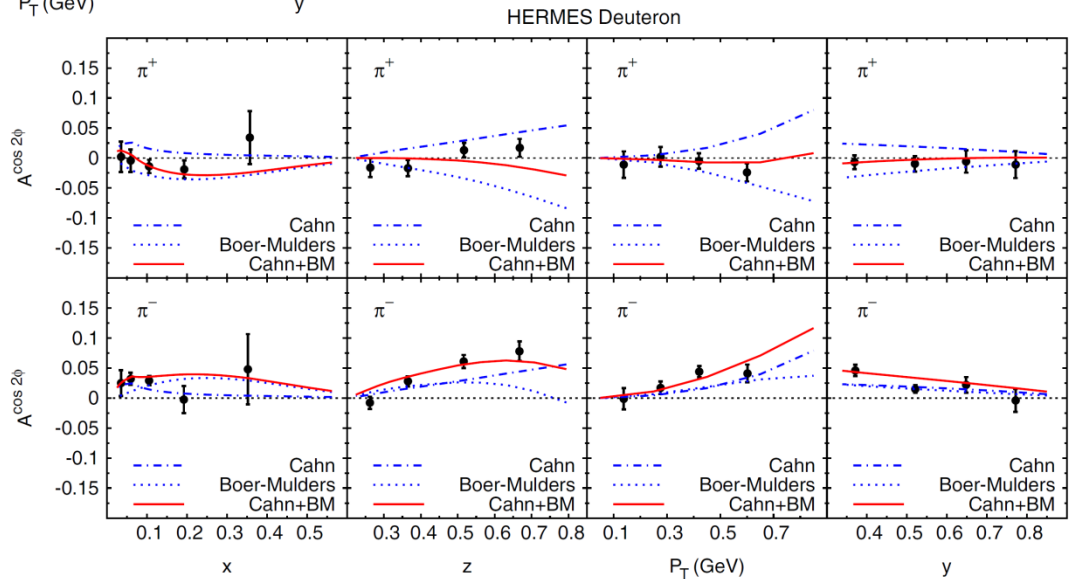
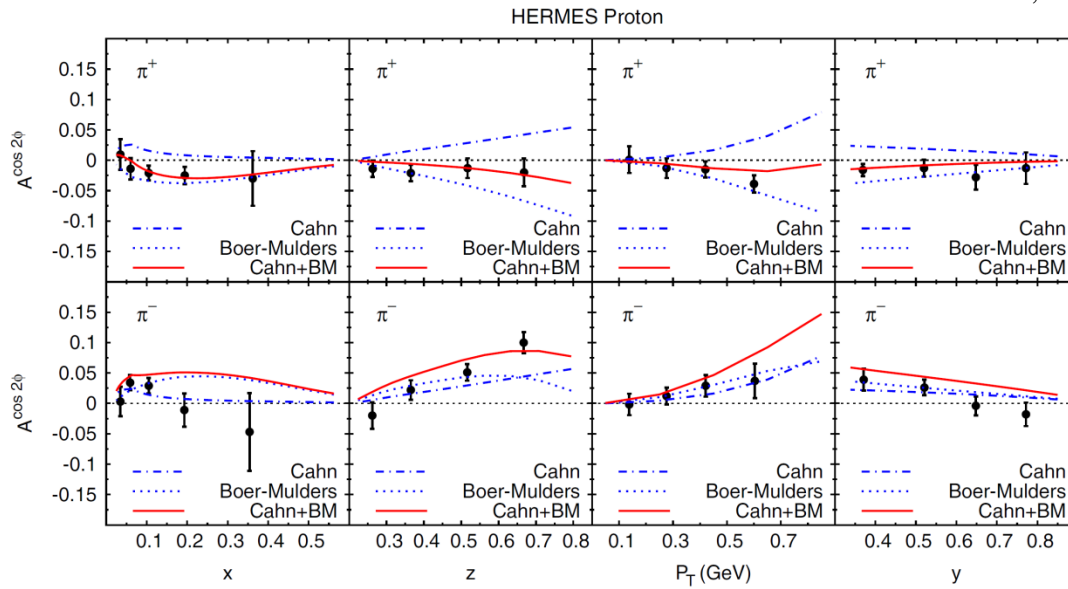
A. Brandenburg, V. V. Khoze, and D. Mueller, Phys. Lett. B347, 413 (1995).

The Berger effect is the exclusive production of a single pion from a free, struck quark that radiates a gluon, produces a $q\bar{q}$ pair, and recombines with the \bar{q} .

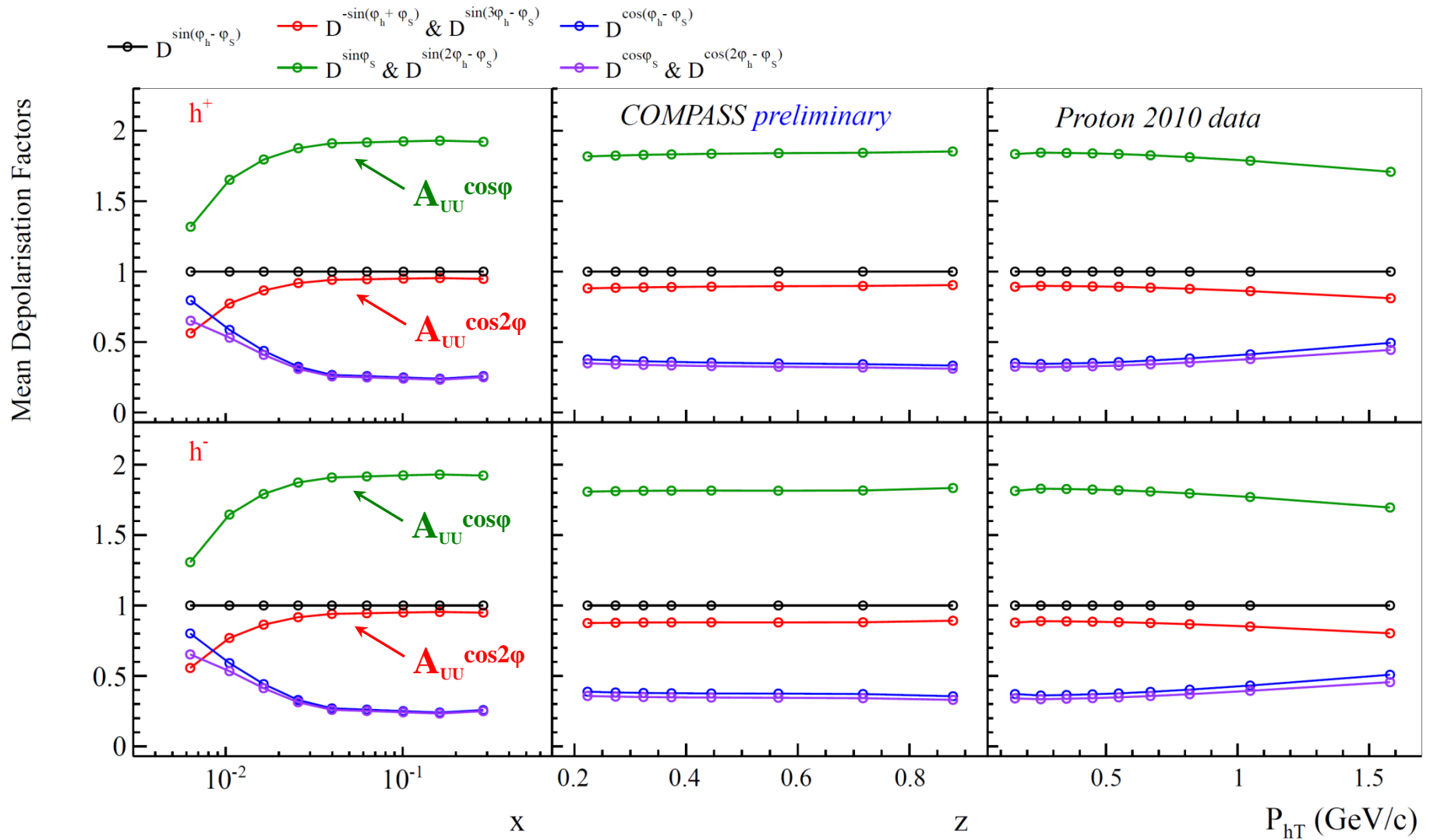
$A_{UU}^{\cos 2\phi}$ -amplitude



Barone, Melis and Prokudin Phys. Rev. D 81, 114026 (2010)



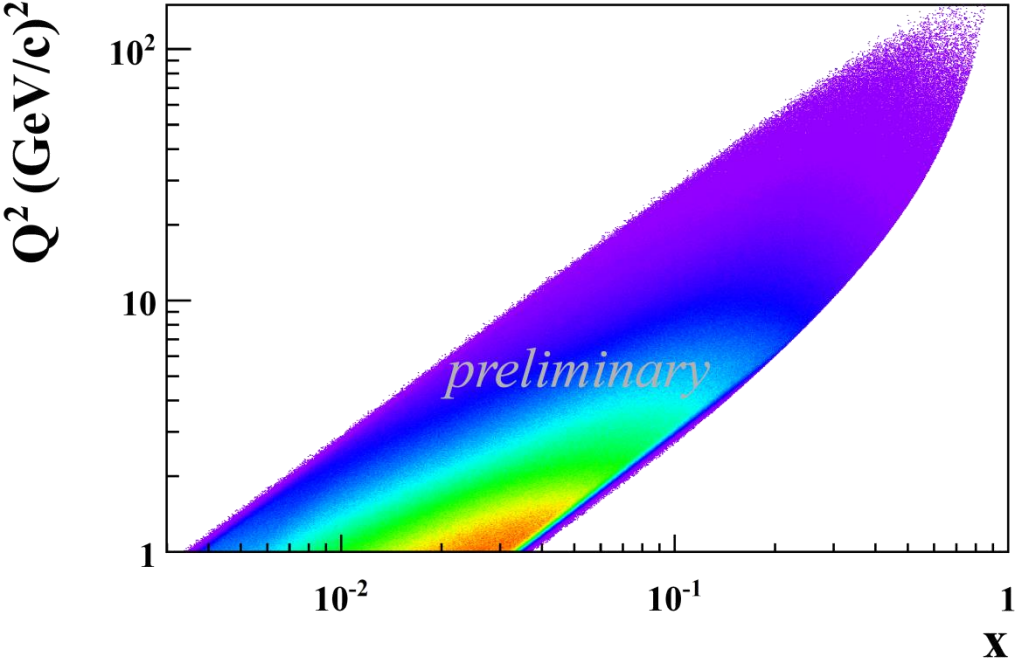
Mean Depolarization Factors



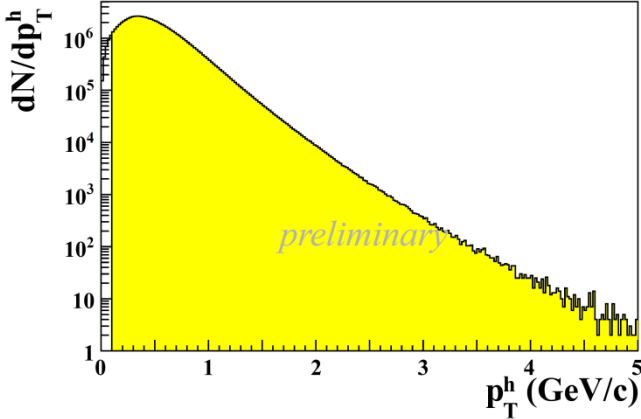
Data selection

- DIS cuts :
 - $Q^2 > 1 \text{ GeV}^2$
 - $0.1 < y < 0.9$
 - $W > 5 \text{ GeV}$
- Hadron cuts :
 - $z > 0.2$
 - $P_{hT} > 0.1 \text{ GeV}/c$

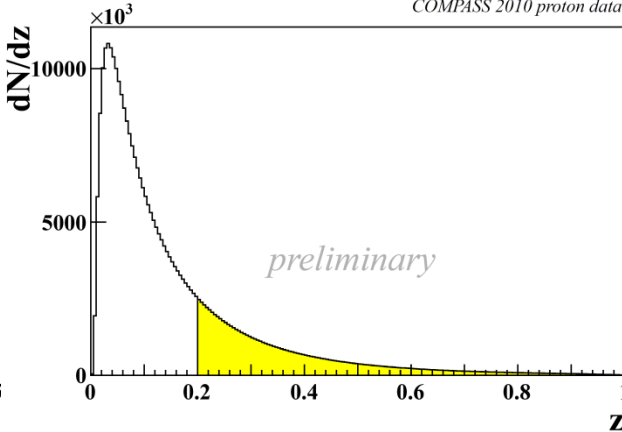
COMPASS 2010 proton data



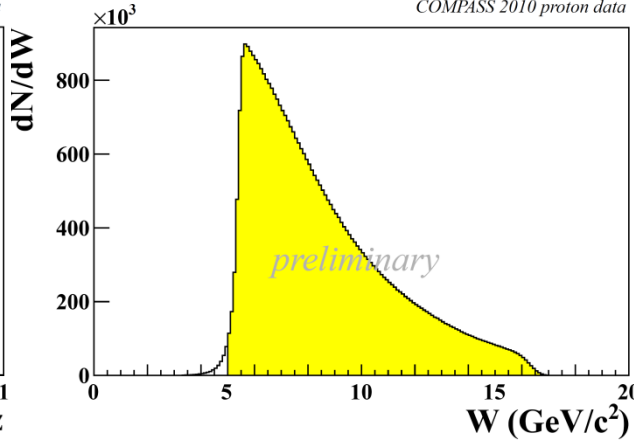
COMPASS 2010 proton data



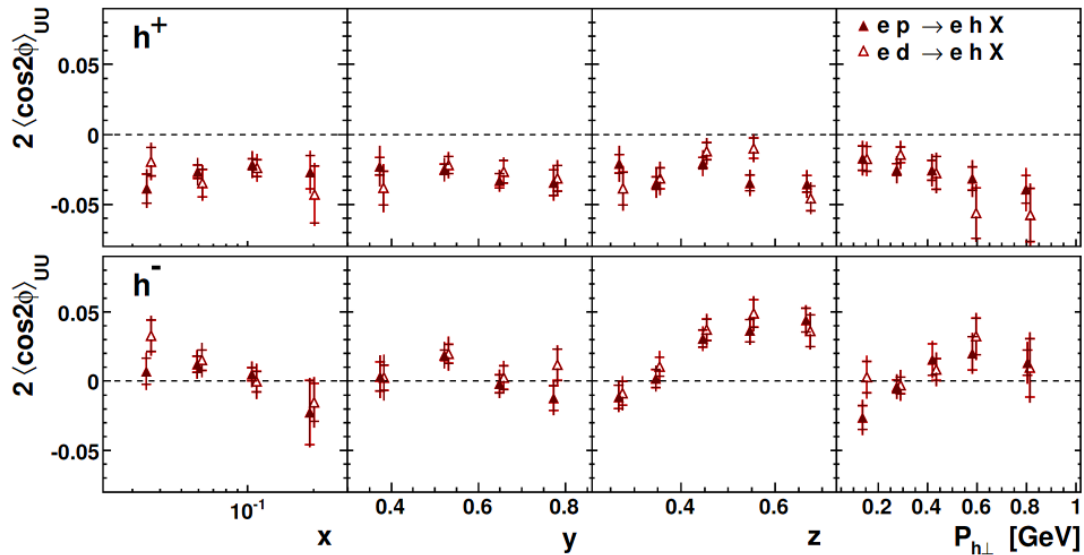
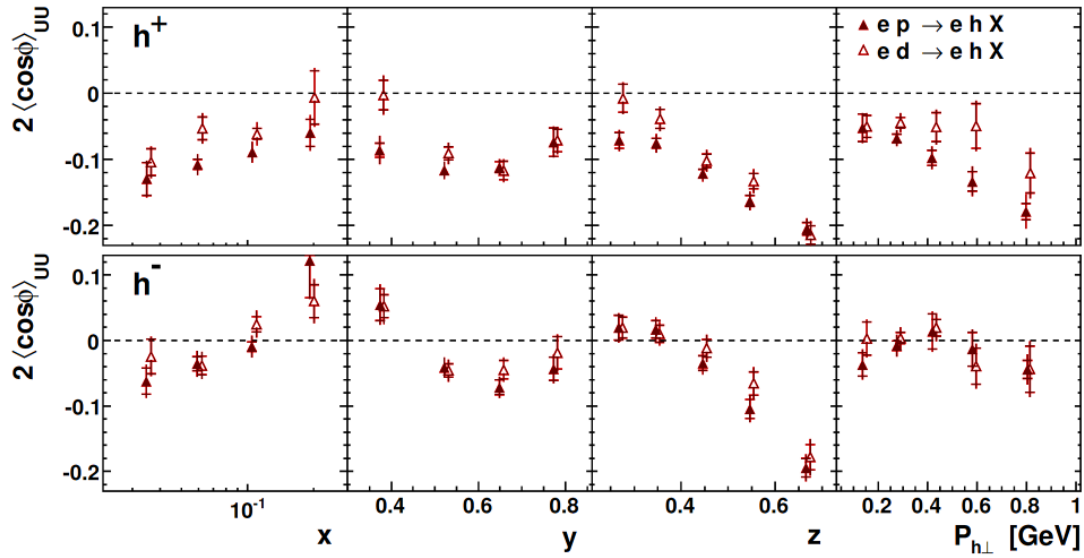
COMPASS 2010 proton data



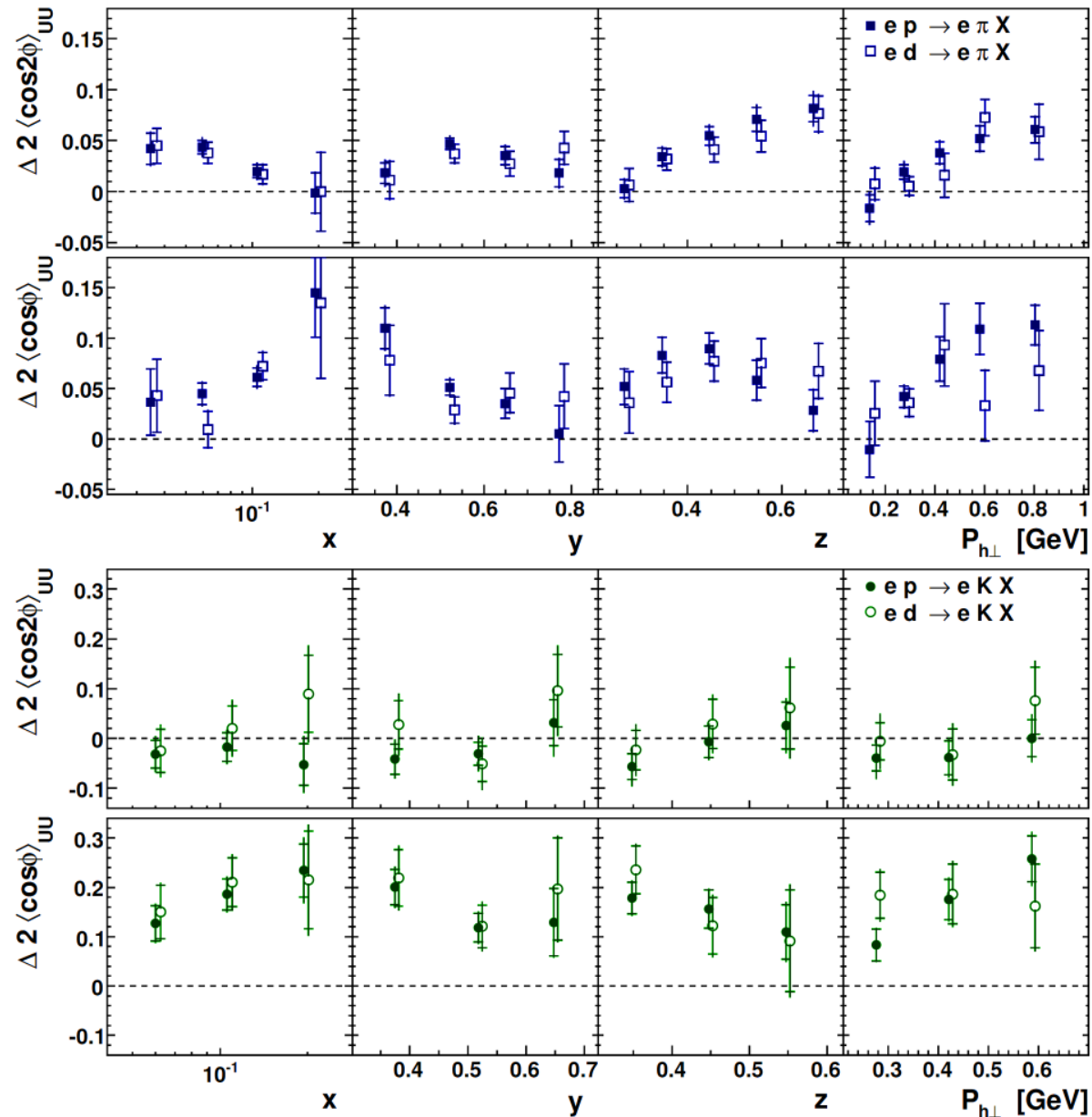
COMPASS 2010 proton data



$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos 2\phi}$ -amplitude on p & d: all hadrons

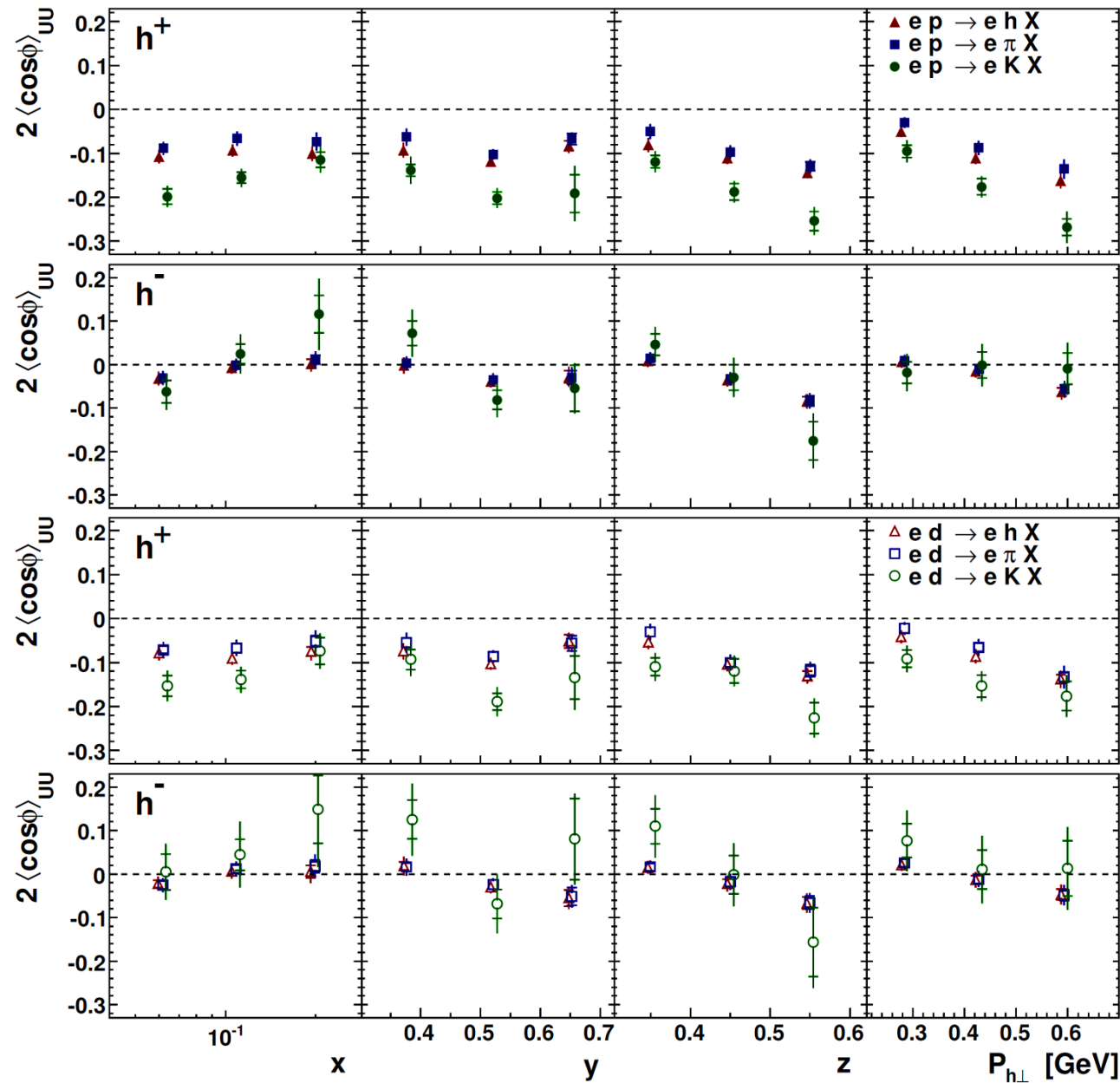


$A_{UU}^{\cos\phi}$ and $A_{UU}^{\cos 2\phi}$ h $^{+/-}$ difference on p & d



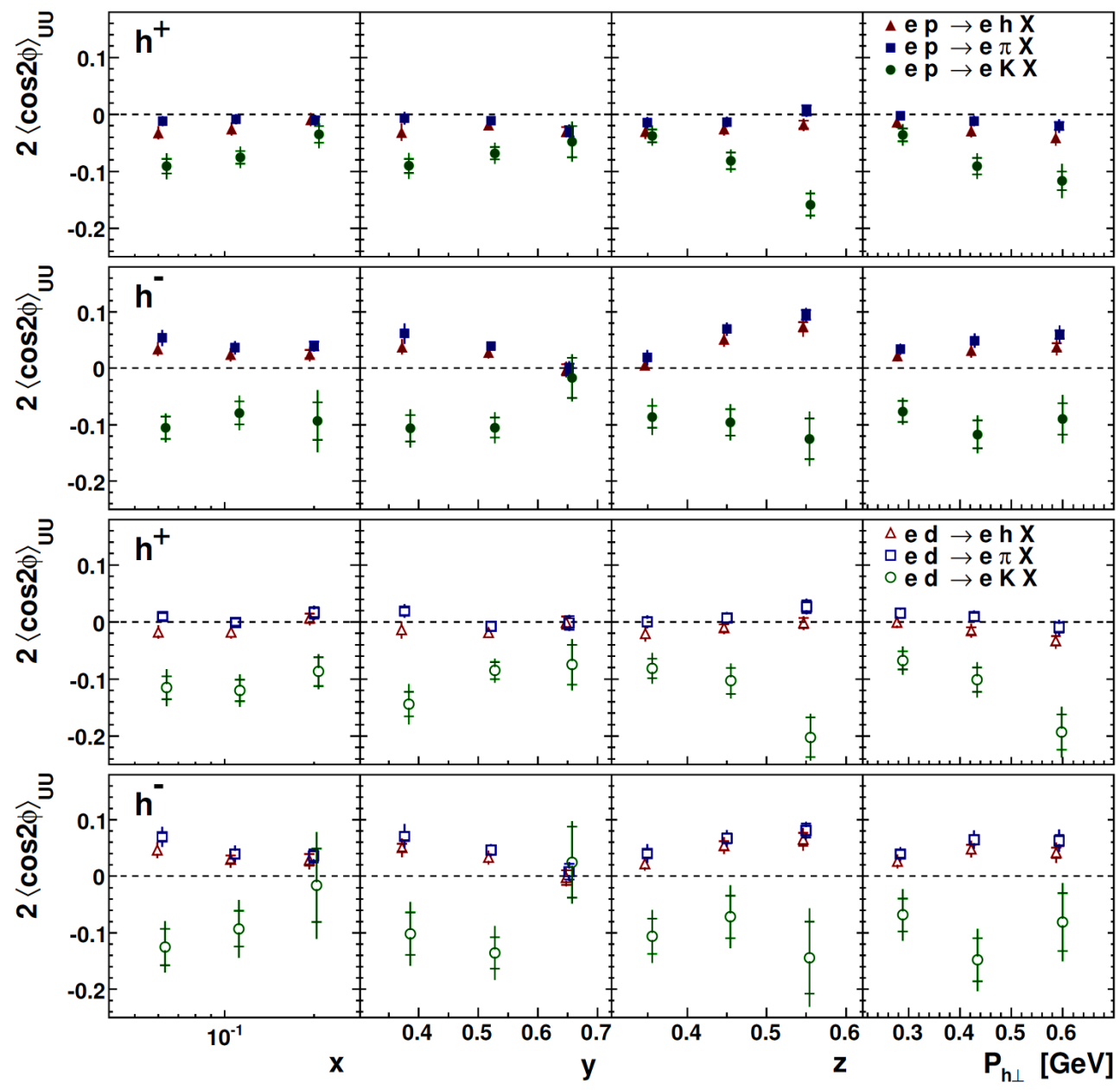
Kinematic range Pions (all hadrons)			
x	y	z	P_{hT}
0.023–0.27	0.3–0.85	0.2–0.75	0.05–1.0
Kinematic range Kaons			
x	y	z	P_{hT}
0.042–0.27	0.3–0.7	0.2–0.6	0.2–0.7

$A_{UU}^{\cos\phi}$ -amplitude on p & d (all hadrons, pions, kaons)



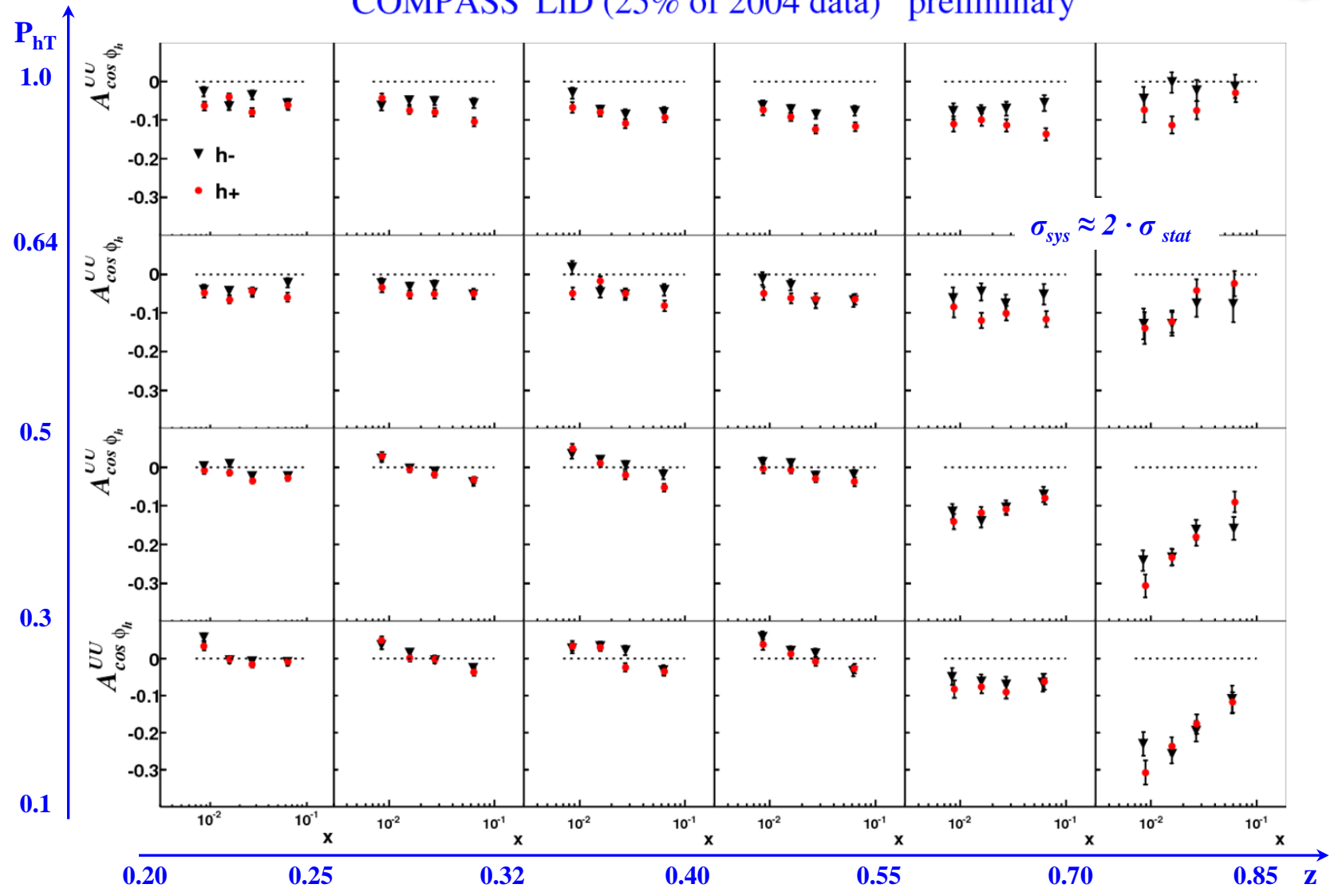


$A_{UU}^{\cos 2\phi}$ -amplitude on p & d (all hadrons, pions, kaons)



$A_{UU}^{\cos\phi_h}$ - asymmetry (x - dependence)

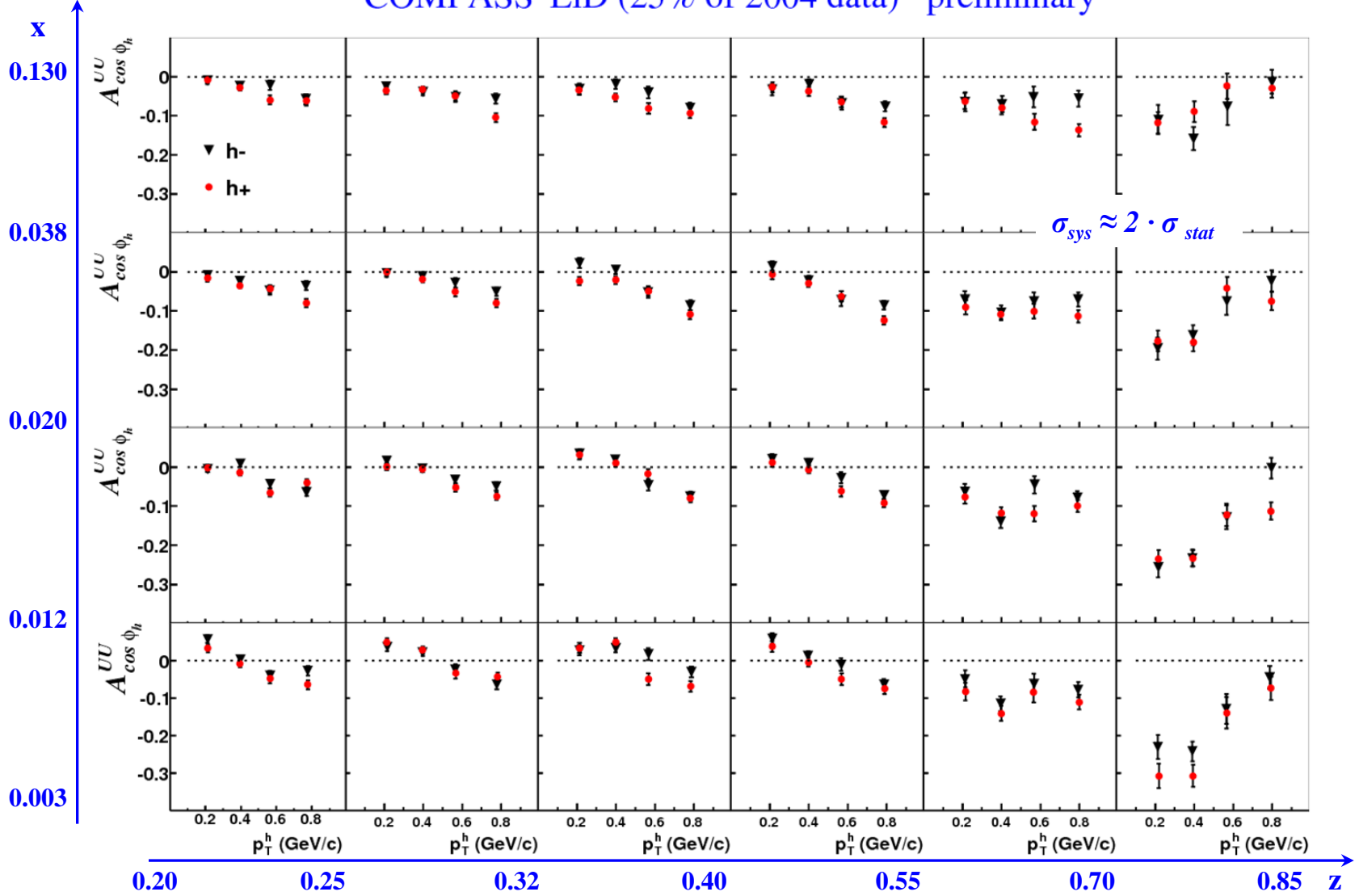
COMPASS⁶LiD (25% of 2004 data) preliminary



largest difference between positive and negative hadrons at large P_{hT} x-trend changes going from small to large z values

$A_{UU}^{\cos\phi_h}$ - asymmetry (P_{hT} - dependence)

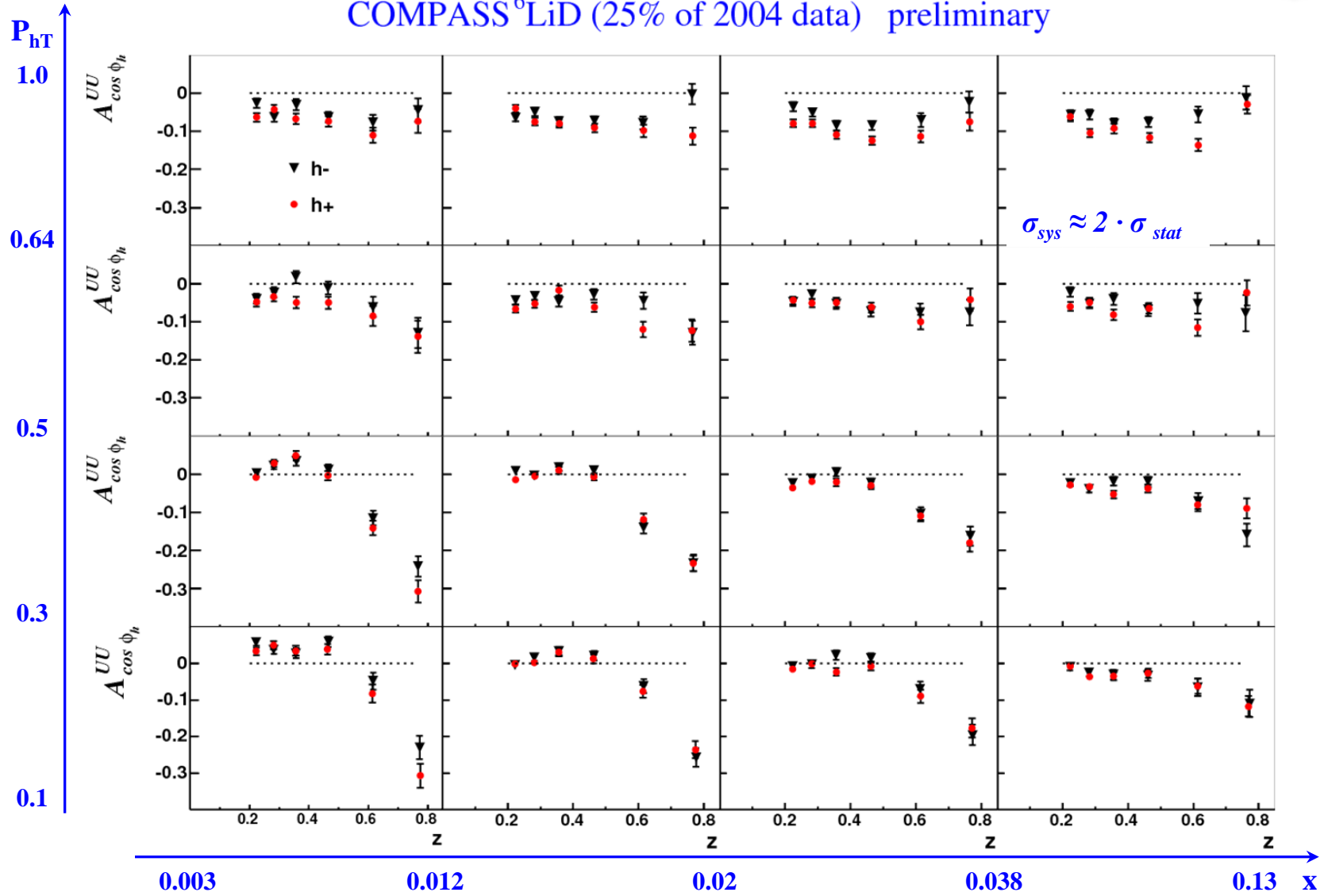
COMPASS ^6LiD (25% of 2004 data) preliminary



P_{hT} trend changes going from small to large z values and it is roughly the same for all x intervals

$A_{UU}^{\cos\phi_h}$ - asymmetry (z - dependence)

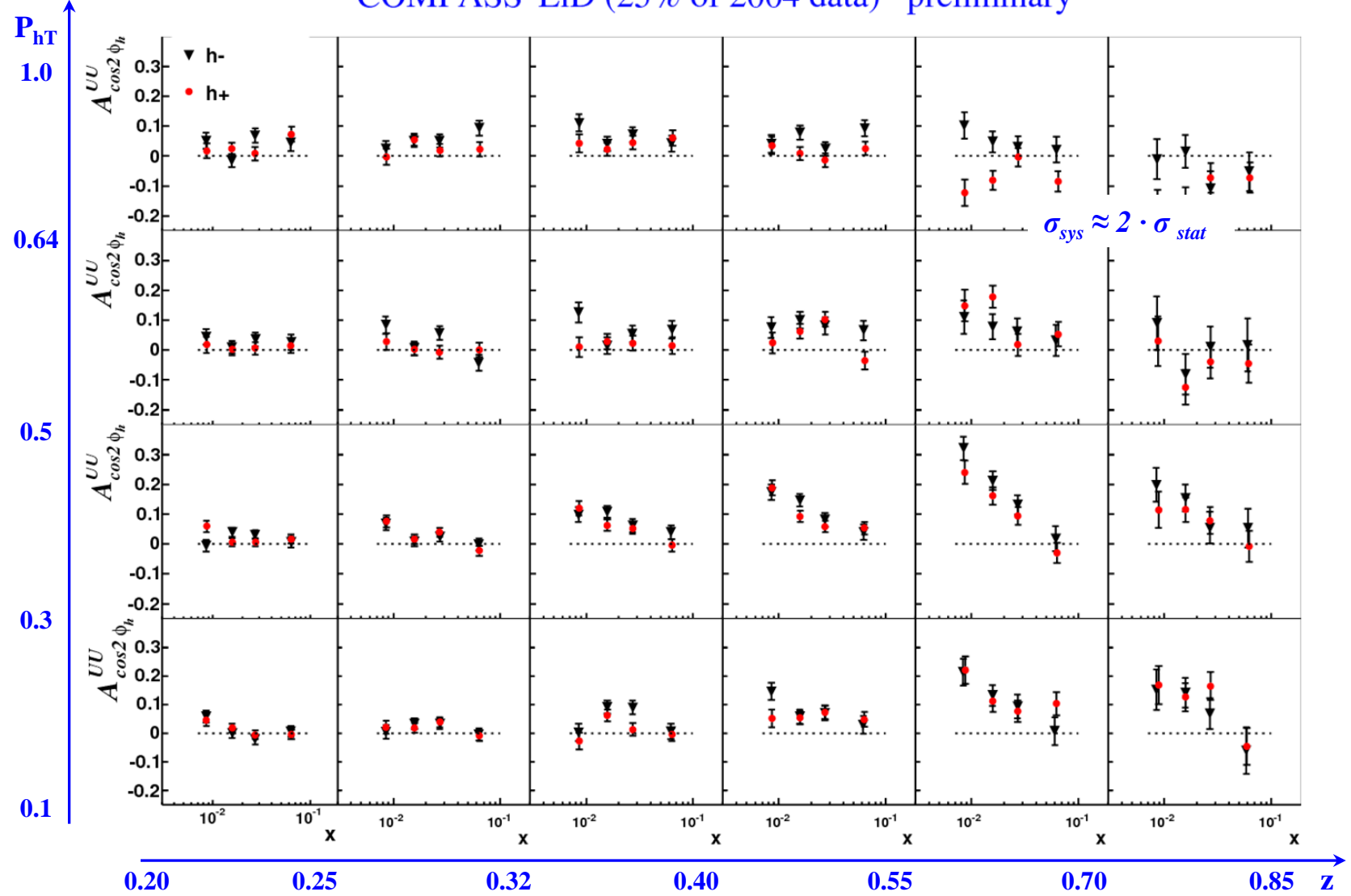
COMPASS⁶LiD (25% of 2004 data) preliminary



z strong dependence more evident at small x and small P_{hT}

$A_{UU}^{\cos 2\phi_h}$ - asymmetry (x - dependence)

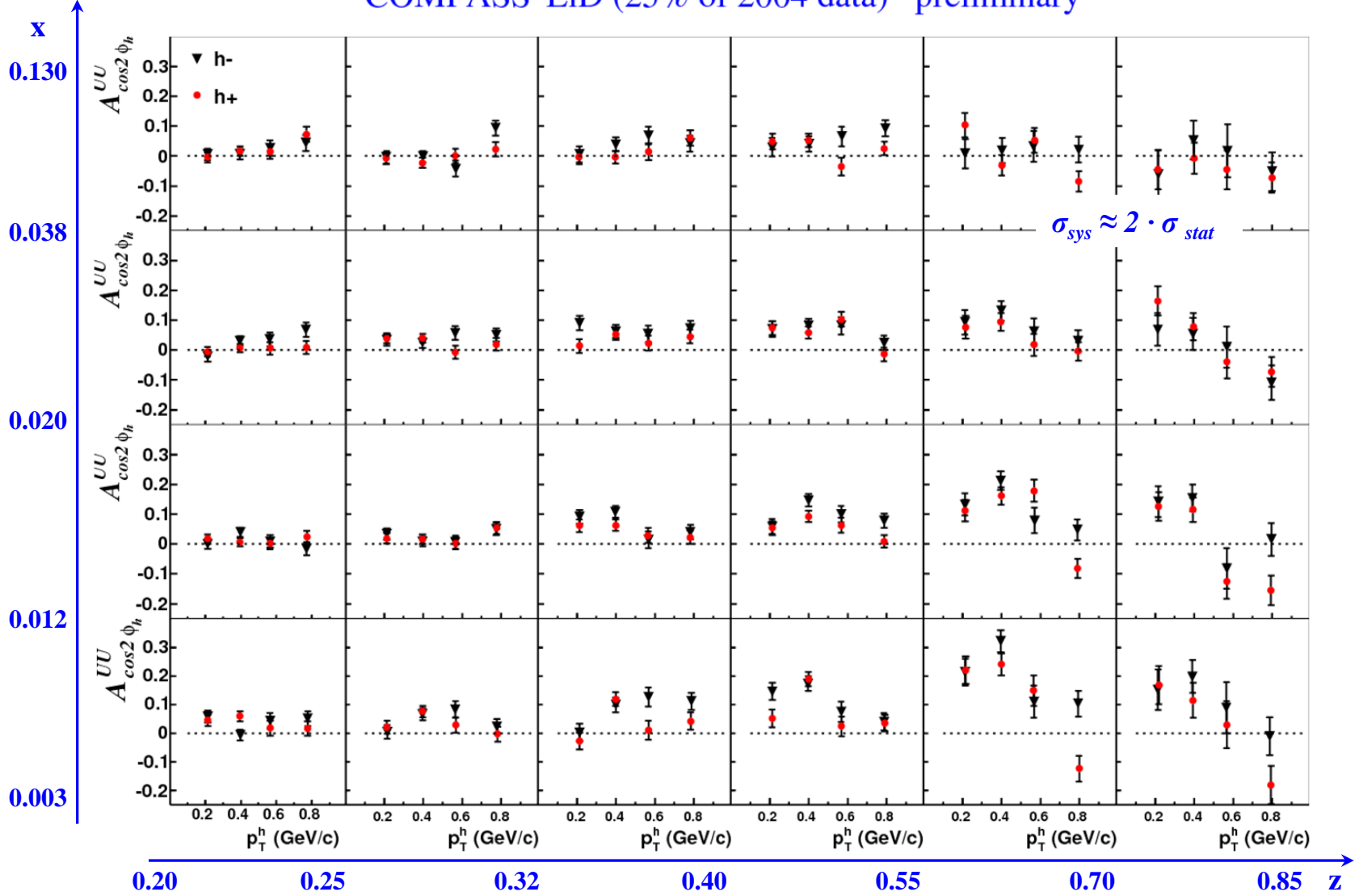
COMPASS⁶LiD (25% of 2004 data) preliminary



x trend changes from small to large z values

$A_{UU}^{\cos 2\phi_h}$ - asymmetry (P_{hT} - dependence)

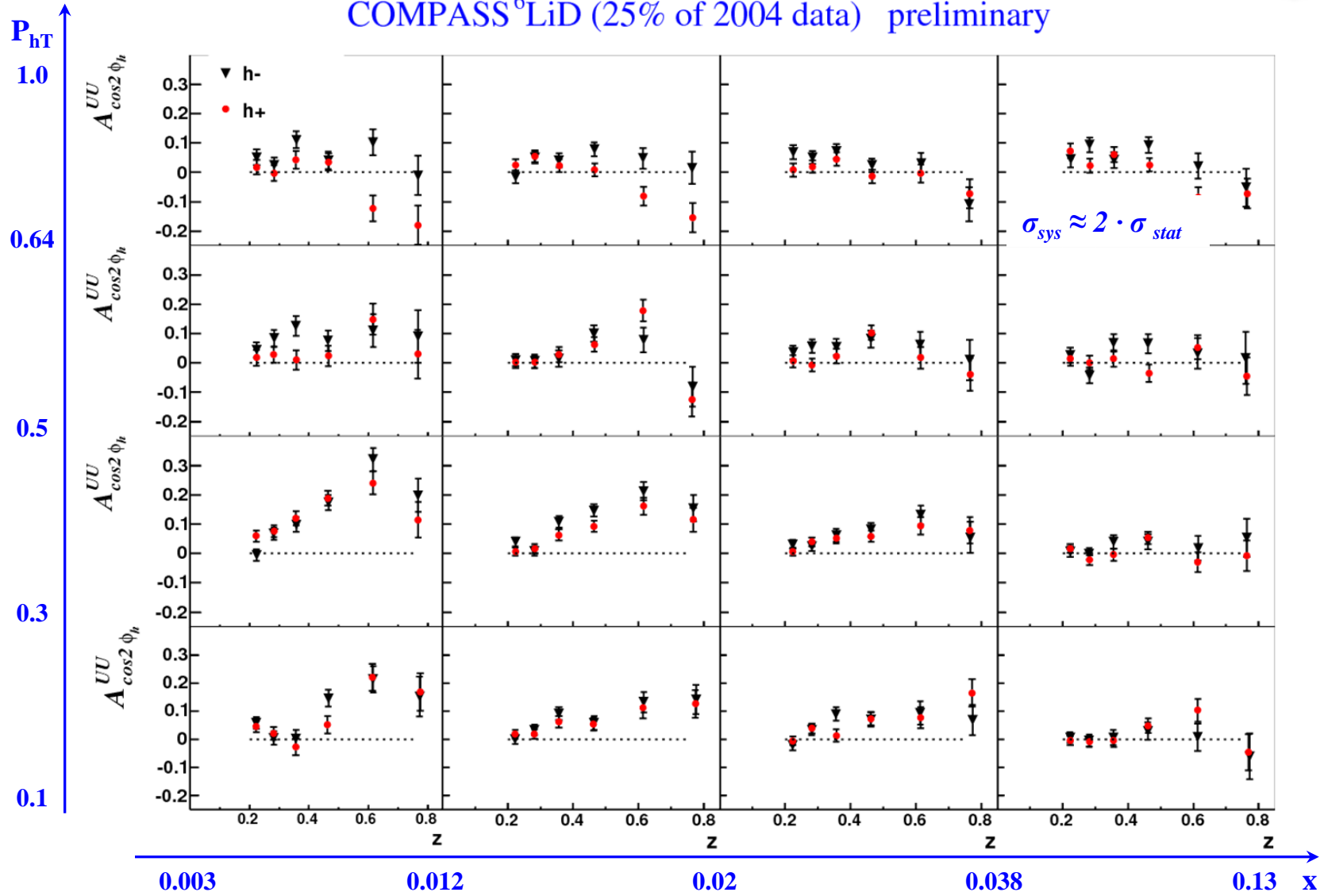
COMPASS ^6LiD (25% of 2004 data) preliminary



the P_{hT} trend difficult to reproduce by models is there for large z and low x

$A_{UU}^{\cos 2\phi_h}$ - asymmetry (z - dependence)

COMPASS⁶LiD (25% of 2004 data) preliminary

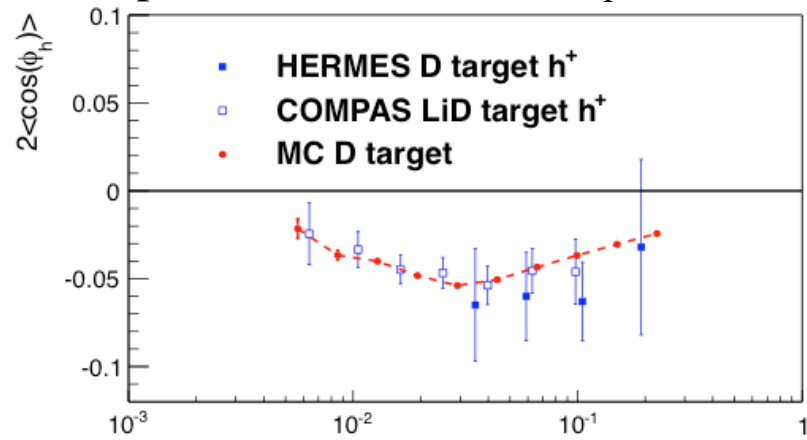


strongest effect at low x and low P_{hT}

Cahn from HERMES and COMPASS (MC studies)

M. Aghasyan arXiv:1307.3500v1 [hep-ex]

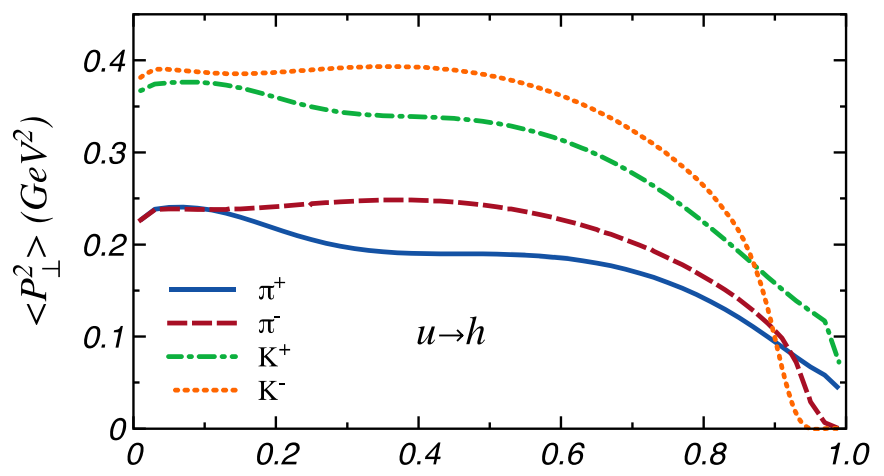
Procedure: fit HERMES Cahn \rightarrow extract k_T -widths \rightarrow use it in MC with correct phase space description and describe the P_T distribution and Cahn effect for COMPASS and HERMES



$f_{1,q}(x)$ Fixed MSTW

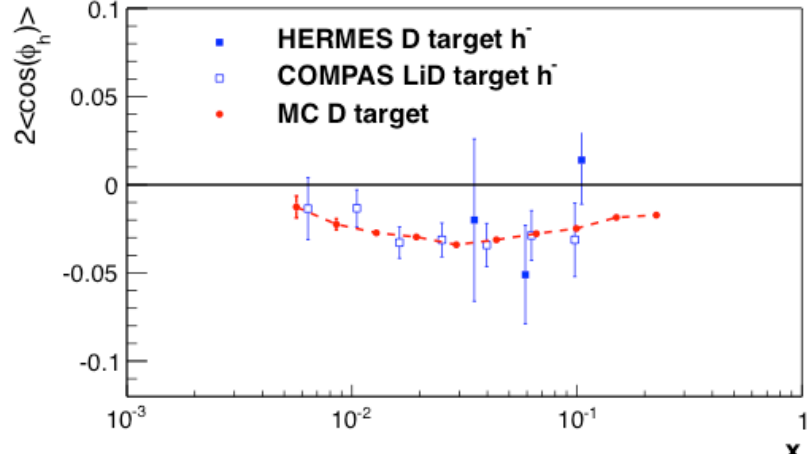
$D_1^{q \rightarrow h}(z)$ Fixed DSS

Fixed - Matevosyan:PRD85,014021 (2012)



The averaged transverse momentum of π and K mesons emitted by a u quark

HERMES k_T - COMPASS Cahn effect



Using parameters from HERMES proton data, the Cahn asymmetry for positive and negative hadrons from D target is presented from COMPASS, HERMES and MC.