

UHE COSMIC NEUTRINOS: failures and hopes

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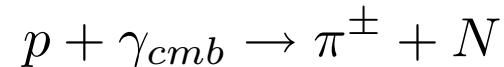
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UHE NEUTRINOS: PRODUCTION and SOURCES

- **Astrophysical neutrinos** (produced by CRs)
pp: $p + N_{\text{tar}} \rightarrow \pi^\pm + \text{all}$, **p γ :** $p + \gamma_{\text{tar}} \rightarrow \pi^\pm + \text{all}$
- **Top-Down neutrinos** (direct pion production :)
TDs, annihilation of DM, decay of SHDM, oscillation of mirror neutrinos.
- **Hidden astrophysical sources:**
Cocooned black hole: **VB, Ginzburg 1981**,
Stecker AGN model: **Stecker et al 1991**,
Collapsing galactic nuclei: **VB, Dokuchaev 2001**,
Hidden jets: **Razzaque, Smirnov 2010**
- **Hidden Top-Down sources:**
Annihilation of DM in the Earth and Sun,
Mirror matter sources (oscillation of neutrinos)
- **Bright phase** (Pop III stars at $z \sim 10 - 20$, $z_{\text{reion}} = 11.0 \pm 0.14$ WMAP).
V.B., Ozernoy 1981, V.B., Blasi 2012.

UHECR and COSMOGENIC NEUTRINOS

UHE protons propagating through CMB produce cosmogenic neutrinos:



Proof of this process is detection of primary protons in UHECR.

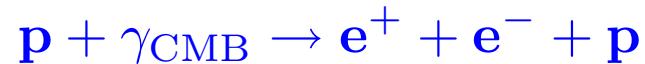
Spectral signatures of proton propagation through CMB are:

pair-production dip and GZK cutoff

COSMOGENIC NEUTRINOS IN THE DIP MODEL FOR UHECR

V.B. and Grigorieva 1988; V.B., Gazizov, Grigorieva 2005 - 2006.

The **dip** is a feature in the spectrum of UHE protons propagating through CMB:

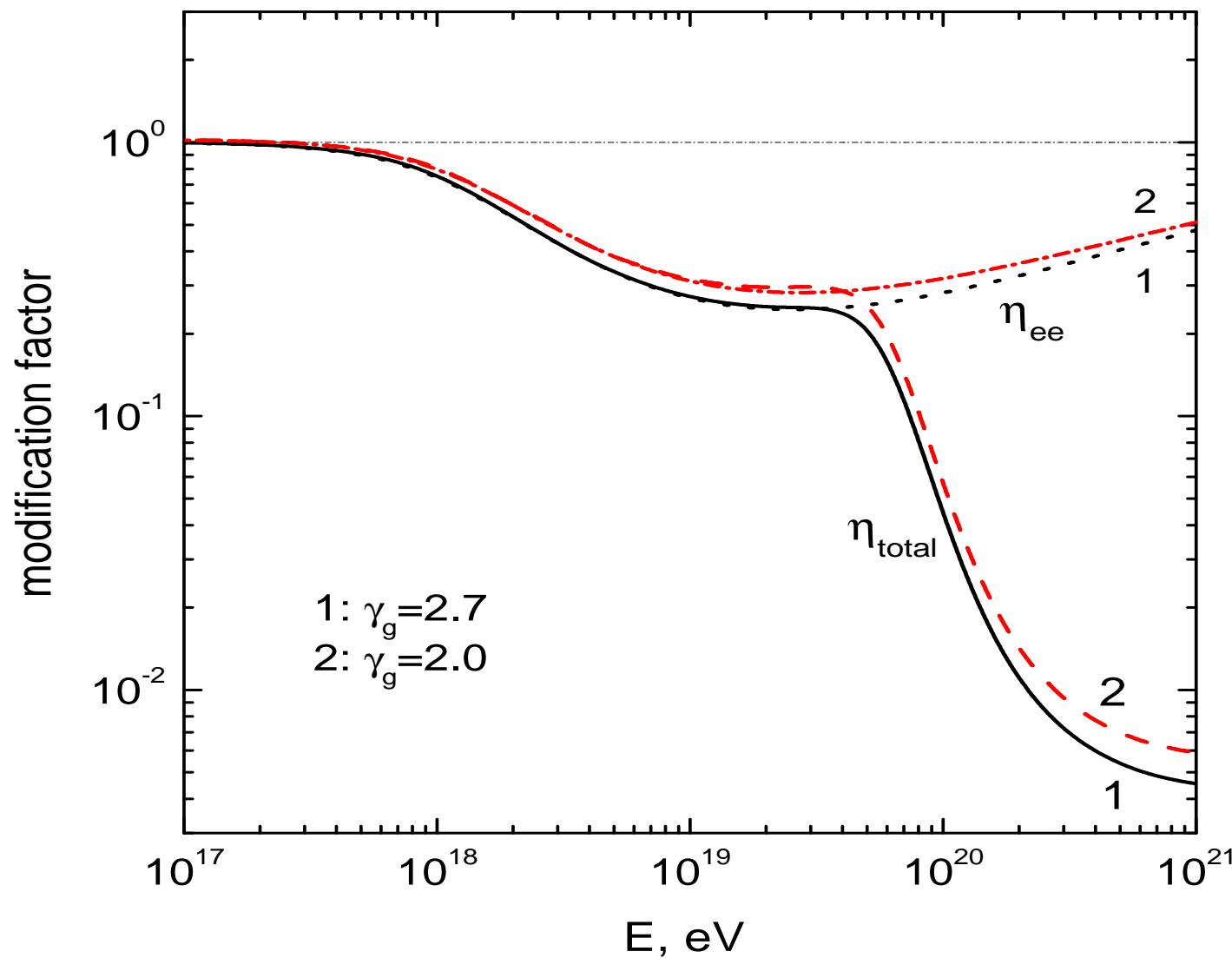


Dip is a faint spectral feature, seen better in terms of **modification factor**

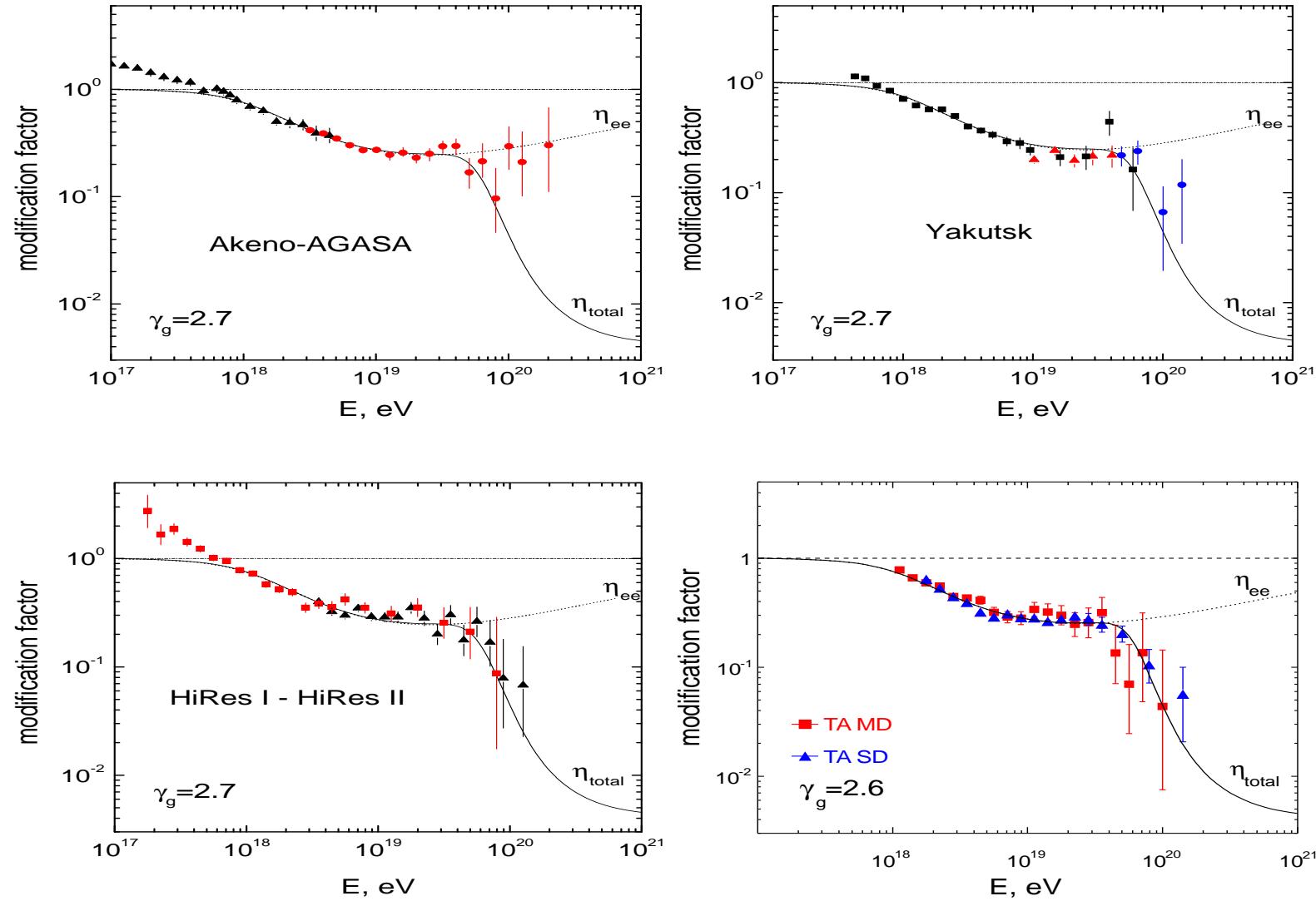
$$\eta(\mathbf{E}) = \frac{\mathbf{J}_p(\mathbf{E})}{\mathbf{J}_p^{\text{unm}}(\mathbf{E})},$$

where $J_p^{\text{unm}}(E) = KE^{-\gamma_g}$ includes only adiabatic energy losses (redshift),
and $J_p(E)$ all energy losses.

DIP AND GZK CUTOFF IN TERMS OF MODIFICATION FACTOR

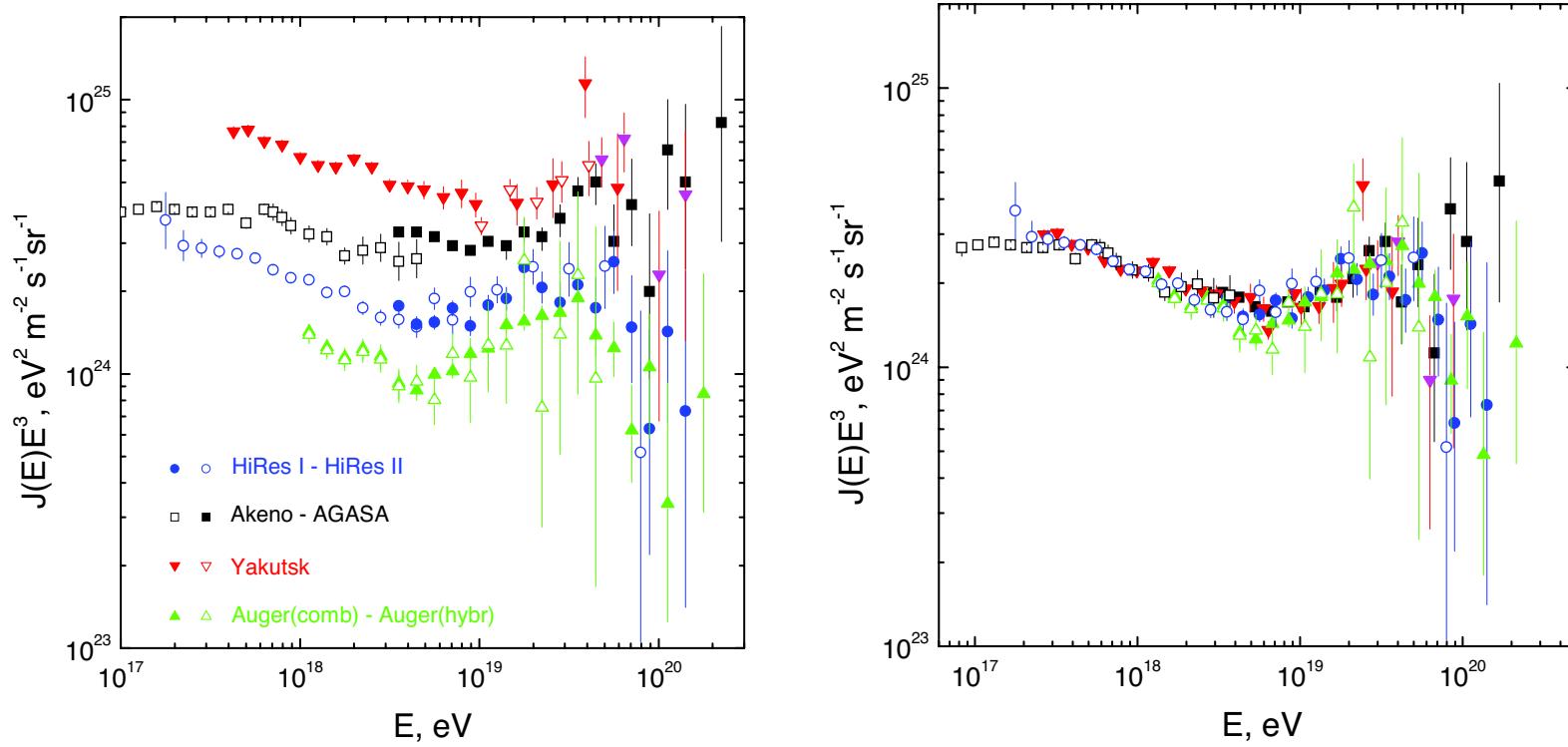


COMPARISON OF PAIR-PRODUCTION DIP WITH OBSERVATIONS



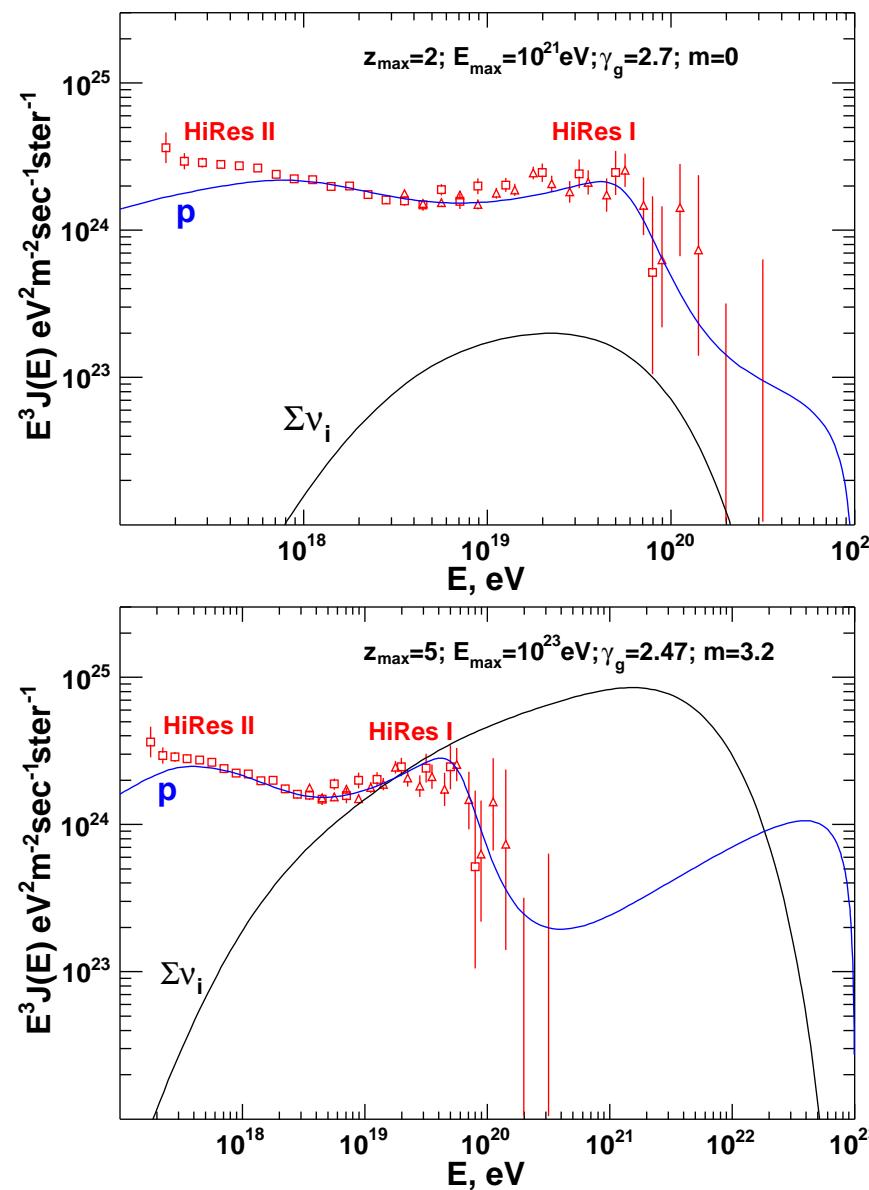
RECALIBRATION of AUGER, HIRES, AGASA, YAKUTSK FLUXES

V.B., A. Gazizov, S. Grigorieva

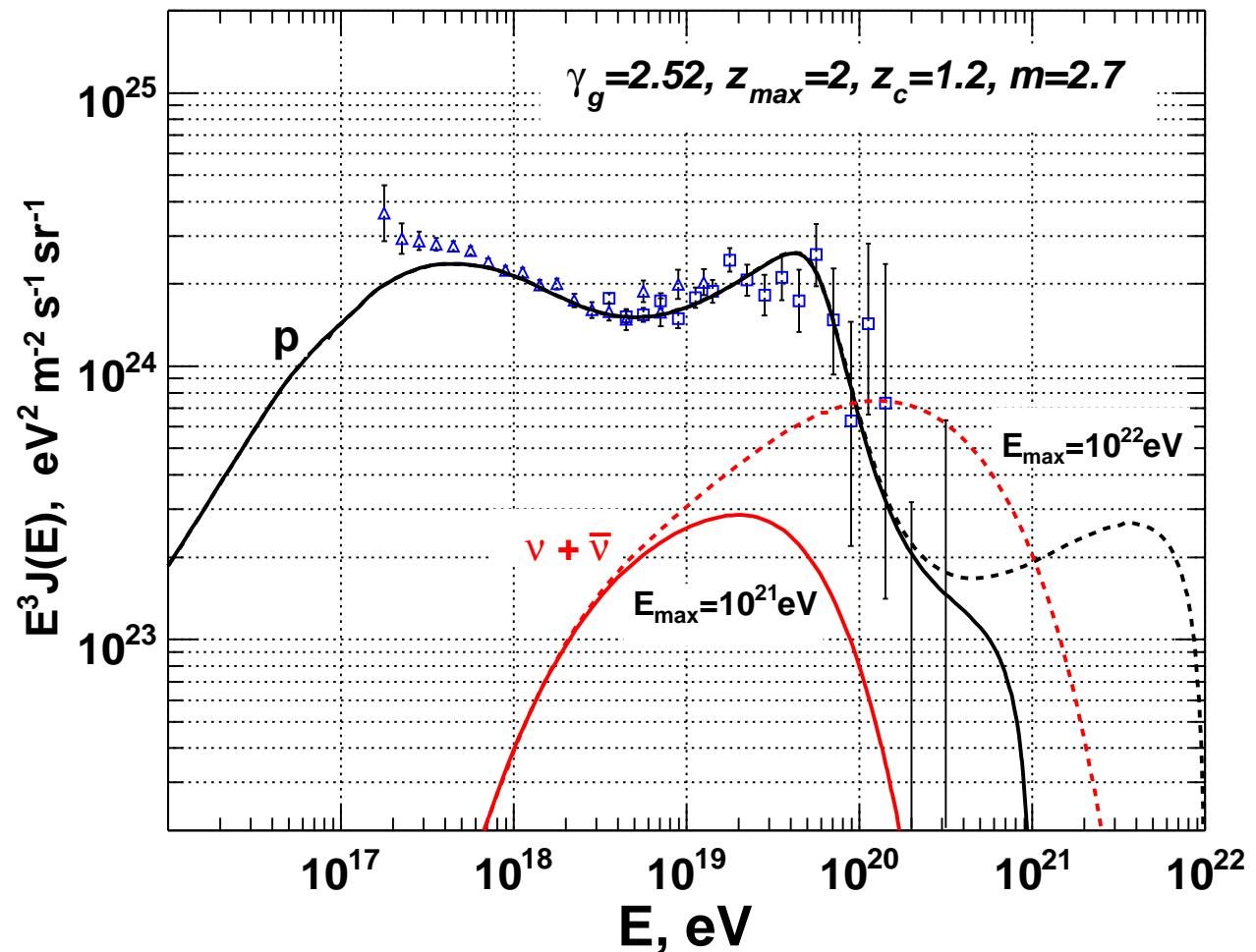


Recalibration factors: $\lambda = 1.2$ (Auger), $\lambda = 1.0$ (HiRes), $\lambda = 0.75$ (AGASA), $\lambda = 0.625$ (Yakutsk).

COSMOGENIC NEUTRINO FLUXES IN THE DIP MODEL

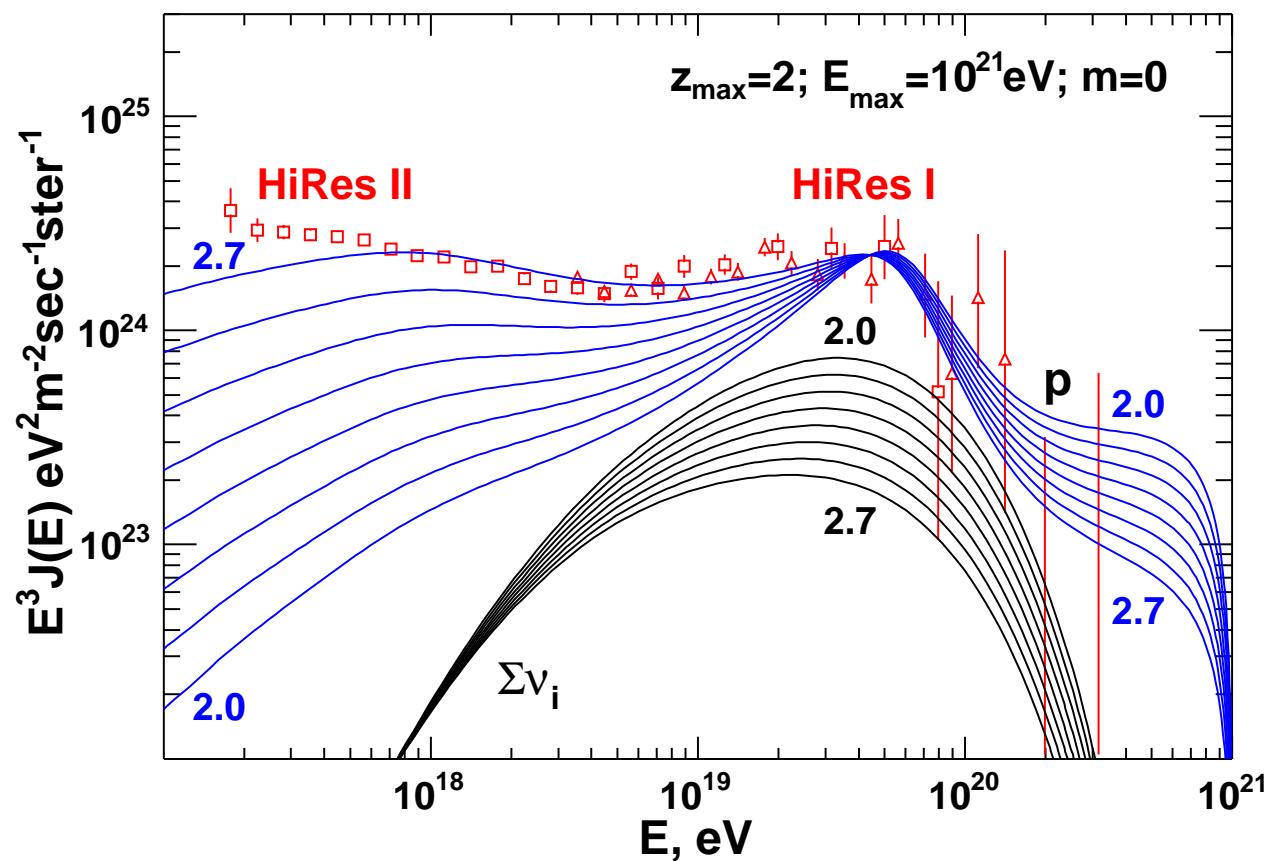


COSMOGENIC NEUTRINO FLUXES FROM AGN



LOWER LIMIT ON NEUTRINO FLUXES IN THE PROTON MODELS

V.B. and A. Gazizov 2009

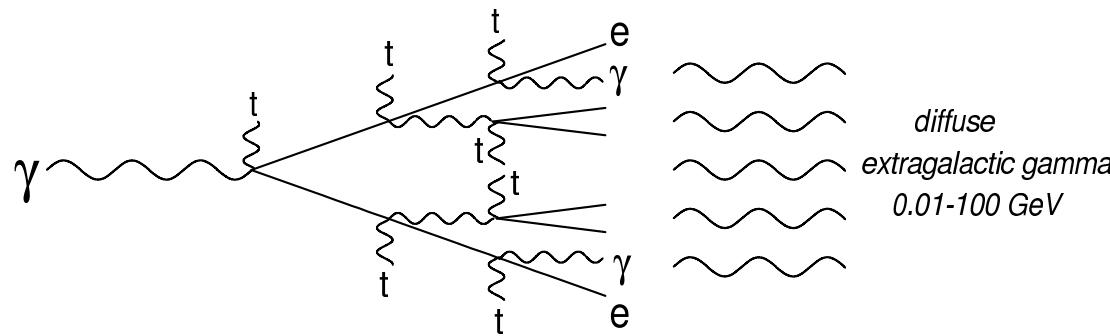


CASCADE UPPER LIMIT

V.B. and A.Smirnov 1975

e – m cascade on target photons :

$$\left\{ \begin{array}{l} \gamma + \gamma_{\text{tar}} \rightarrow e^+ + e^- \\ e + \gamma_{\text{tar}} \rightarrow e' + \gamma' \end{array} \right.$$



Spectrum of cascade photons

$$J_\gamma^{\text{cas}}(E) = \begin{cases} K(E/\varepsilon_X)^{-3/2} & \text{for } E \leq \varepsilon_X, \\ K(E/\varepsilon_X)^{-2} & \text{for } \varepsilon_X \leq E \leq \varepsilon_a, \end{cases} \quad (1)$$

with a steepening at $E > \varepsilon_a$, and $\varepsilon_X = 1/3 (\varepsilon_a/m_e)^2 \varepsilon_{\text{cmb}}$.

EGRET: agreement with spectrum (1) and $\omega_\gamma^{\text{obs}} \sim 3 \times 10^{-6} \text{ eV/cm}^3$.

UPPER LIMIT ON NEUTRINO FLUX

$$\omega_{\text{cas}} > \frac{4\pi}{c} \int_E^\infty E J_\nu(E) dE > \frac{4\pi}{c} E \int_E^\infty J_\nu(E) dE \equiv \frac{4\pi}{c} E J_\nu(> E)$$

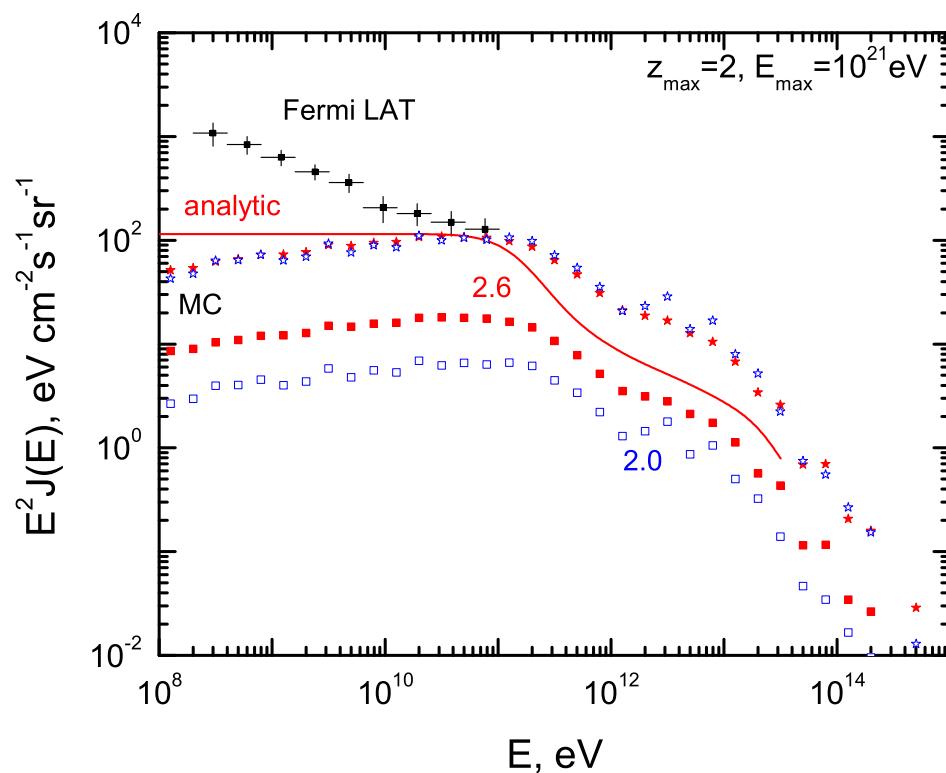
$$E^2 I_\nu(E) < \frac{c}{4\pi} \omega_{\text{cas}}.$$

E^{-2} - generation spectrum:

$$E^2 J_\nu(E) < \frac{c}{4\pi} \frac{\omega_{\text{cas}}}{\ln E_{\text{max}}/E_{\text{min}}}.$$

CASCADE UPPER LIMIT FROM FERMI LAT

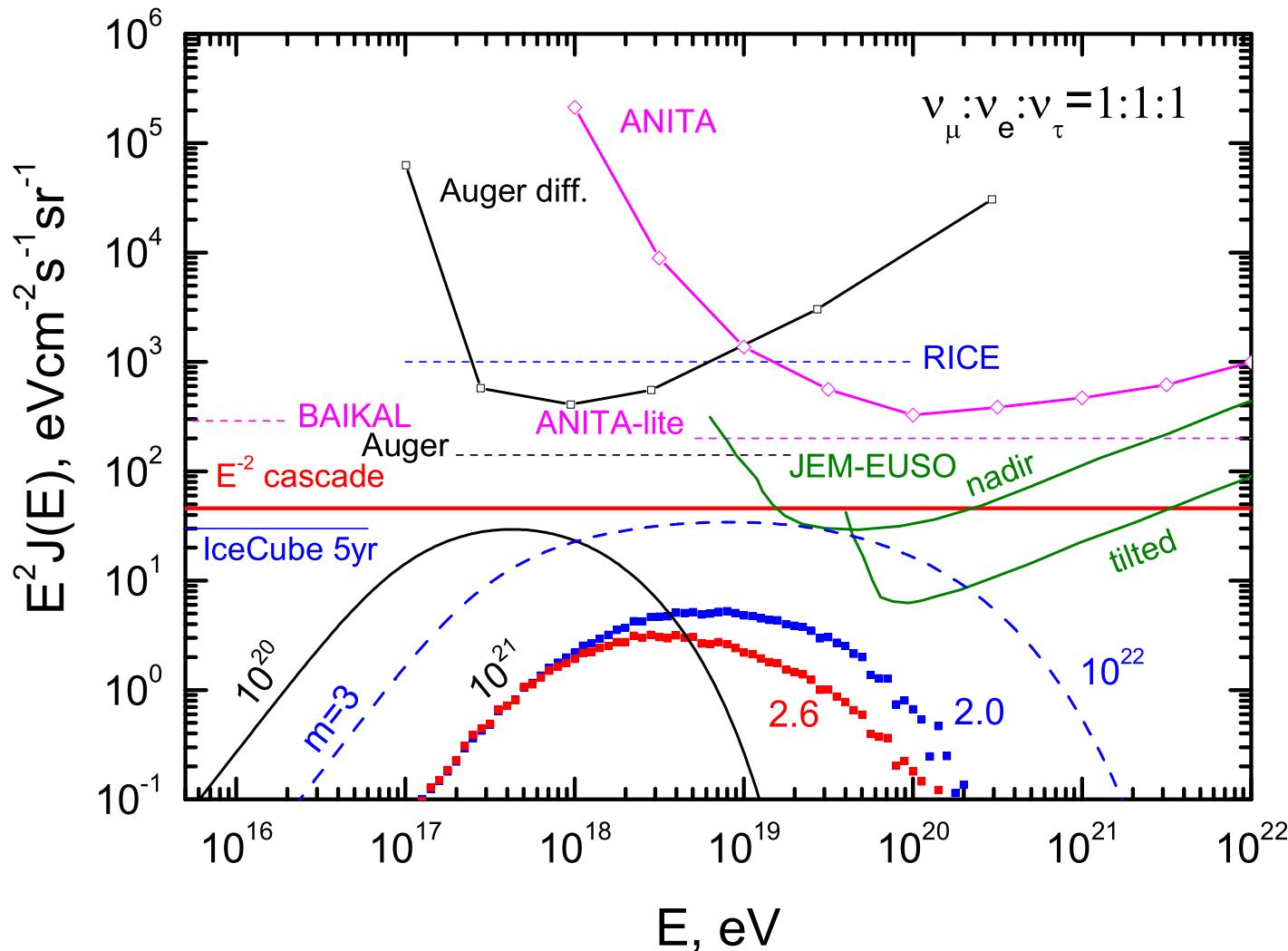
V.B., Gazizov, Kachelriess, Ostapchenko Phys. Lett. B 695 (2011) 13.
Ahlers, Anchordoqui, Gonzales-Garcia, Halzen, Sarkar Astrop. Phys. 34 (2010) 106.
G. Gelmini, O. Kalashev, D. Semikoz arXiv:1107.1672



$$\omega_{\text{cas}}^{\max} = 5.8 \times 10^{-7} \text{ eVcm}^{-3}$$

OBSERVATIONAL AND THEORETICAL UPPER LIMITS

V.B., Gazizov, Kachelriess, Ostapchenko 2010.



Reionization of Universe: Bright Phase

Burst of first massive star formation

- Cooling of universe and **recombination** at $T_{\text{dec}} \sim 3600 \text{ K}$, $z_{\text{dec}} \sim 1100$.
- **DARK AGES:** Evolution of DM structures with neutral hydrogen.
- Cooling of baryonic gas by H_2 formation $z \leq 20 - 30$
- Formation of the first **Pop III** stars hosted by $M \sim 10^6 M_\odot$ DM halos.
Properties: metal poor, massive ($M \geq 100 M_\odot$), hot ($\varepsilon \sim 30 \text{ eV}$), short-lived, strong wind, finishing evolution by SN explosion with W_{SN} up to 10^{53} erg .
- **Reionization** by Pop III radiation and by Pop III SN, observed by WMAP. For model of **Instantaneous reionization** $z_{\text{reion}} = 11.0 \pm 1.4$ (WMAP).

Pop III scenario fills in several blanks in astrophysics:

- Produces **metals** needed for evolution of normal stars.
- Produces **reionization** observed by WMAP.
- Produces **magnetic fields** in Universe.

Pioneering proposal: Bisnovatiy-Kogan, Ruzmaikin, Syunyaev 1973.

Magn. field production **without pre-existing magn. field**:

Collisionless shock without magn. field (e.g. Spitkovsky 2008).

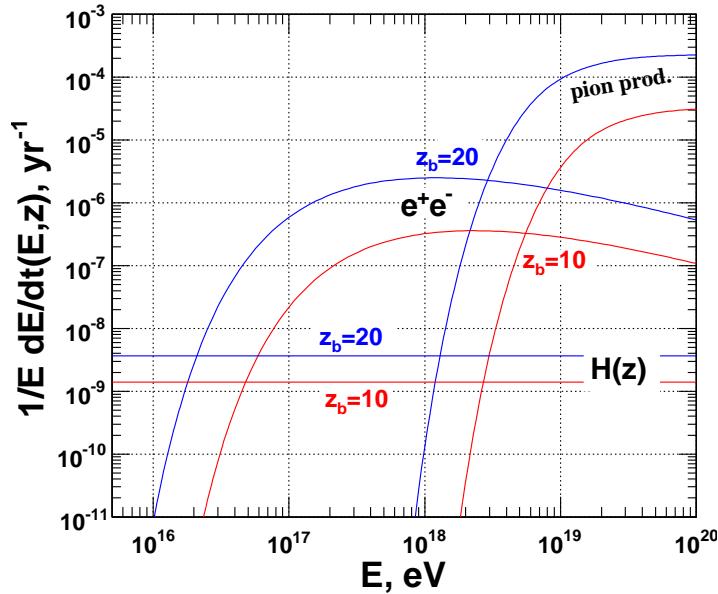
Weibel instability in shocks and winds (Medvedev et al 1999, 2006).

Resistive mechanism in shocks and winds (Miniati and Bell 2010).

UHE NEUTRINOS from BRIGHT PHASE

V.B., Ozernoy 1981; Gao et al 2011, V.B., Blasi 2011.

At $z_b \sim 10 - 20$ CRs (mostly protons) are accelerated in Pop III SN.
 Cosmogenic neutrinos are produced in $p\gamma_{\text{cmb}}$.



Neutrino flux has a maximum at:

$$E_\nu \sim 0.05 \frac{E_{\text{GZK}}}{(1+z_b)^2} \approx 7.5 \times 10^{15} \text{ eV},$$

$$E_\nu^{\max} \sim 0.05 \frac{E_p^{\max}}{(1+z_b)} \approx 2.5 \times 10^{17} \text{ eV}.$$

Energy losses of protons at z_b .

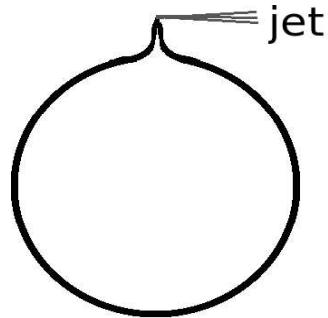
$$E_\nu^2 J_\nu(E) = 0.1 \frac{c}{4\pi} \frac{\omega_p(z_b)}{(1+z_b)^4} \frac{1}{\ln E_{\max}/E_{\min}}$$

The value of **basic parameter** $\frac{\omega_p(z_b)}{(1+z_b)^4} = 9.5 \times 10^{-7} \text{ eV/cm}^3$ corresponds to future IceCube sensitivity $E^2 J_\nu(E) = 3 \times 10^{-9} \text{ GeV/cm}^2 \text{ s sr}$ and respects the Fermi upper limit $\omega_{\text{cas}}^{\max} = 5.8 \times 10^{-8} \text{ eV/cm}^3$.

TOPOLOGICAL DEFECTS

Symmetry breaking in early universe results in phase transitions (D.A. Kirzhnitz 1972), accompanied by TDs. Their common feature is production of HE particles.

Ordinary and superconducting strings



Produced at $U(1)$ symmetry breaking.

Particles are massless inside the string and massive outside.

Loop oscillates with periodically produced cusp ($v = c$) and with large Lorentz factors, e.g. above $\Gamma \sim 10^{10}$, at nearby points.

Particles escaping from cusp segment have energy $E \sim \Gamma m_X$,

Particles are emitted as jets with $\theta \sim 1/\Gamma_c$.

In a wide class of particle physics strings are superconducting (Witten 1985)

UHE neutrino jets from superconducting strings

V.B., K.Olum, E.Sabancilar and A.Vilenkin 2009

Electric current is generated in magnetic fields (B , f_B), dominated by clusters.

Clusters of galaxies dominate.

Lorentz factor of cusp $\Gamma_c \sim 1 \times 10^{12}$.

Diffuse neutrino flux :

$$E^2 J_\nu(E) = 2 \times 10^{-8} \mathbf{i}_c B_{-6} f_{-3} \text{ GeV cm}^{-2} \text{s}^{-1}$$

UHE NEUTRINOS FROM ORDINARY STRINGS.

1. Ordinary strings with EW Higgs condensate. Vachaspati 2010

Interaction of EW Higgs ϕ with the string field Φ (κ is coupling constant):

$$S_{\text{int}} = \kappa \int d^4x (\Phi^+ \Phi - \eta^2) \phi^+ \phi$$

After GUT symmetry breaking ($\langle \Phi \rangle = \eta$):

$$S_{\text{int}} = -\kappa \eta \int d^2\sigma \sqrt{-\gamma} \phi,$$

where $d^2\sigma$ is string world-sheet space, γ_{ab} is the world-sheet metric.

The higgses are emitted through the cusp.

1. UHE neutrinos emitted from ordinary strings via dilatons and moduli.

VB, Sabancilar, Vilenkin, 2011 following Damour and Vilenkin 1997.

$$S_{\text{int}} = (\sqrt{4\pi}\alpha/M_{\text{Pl}}) \int d^4x \phi T_\nu^\nu,$$

$$T_\nu^\nu(x) = -2\eta^2 \int d^2\sigma \sqrt{-\gamma} \delta^4(x^\alpha - x^\alpha(\sigma))$$

is the trace of energy-momentum tensor of string.

Dilatons and moduli are produced as radiation quanta from the cusp.

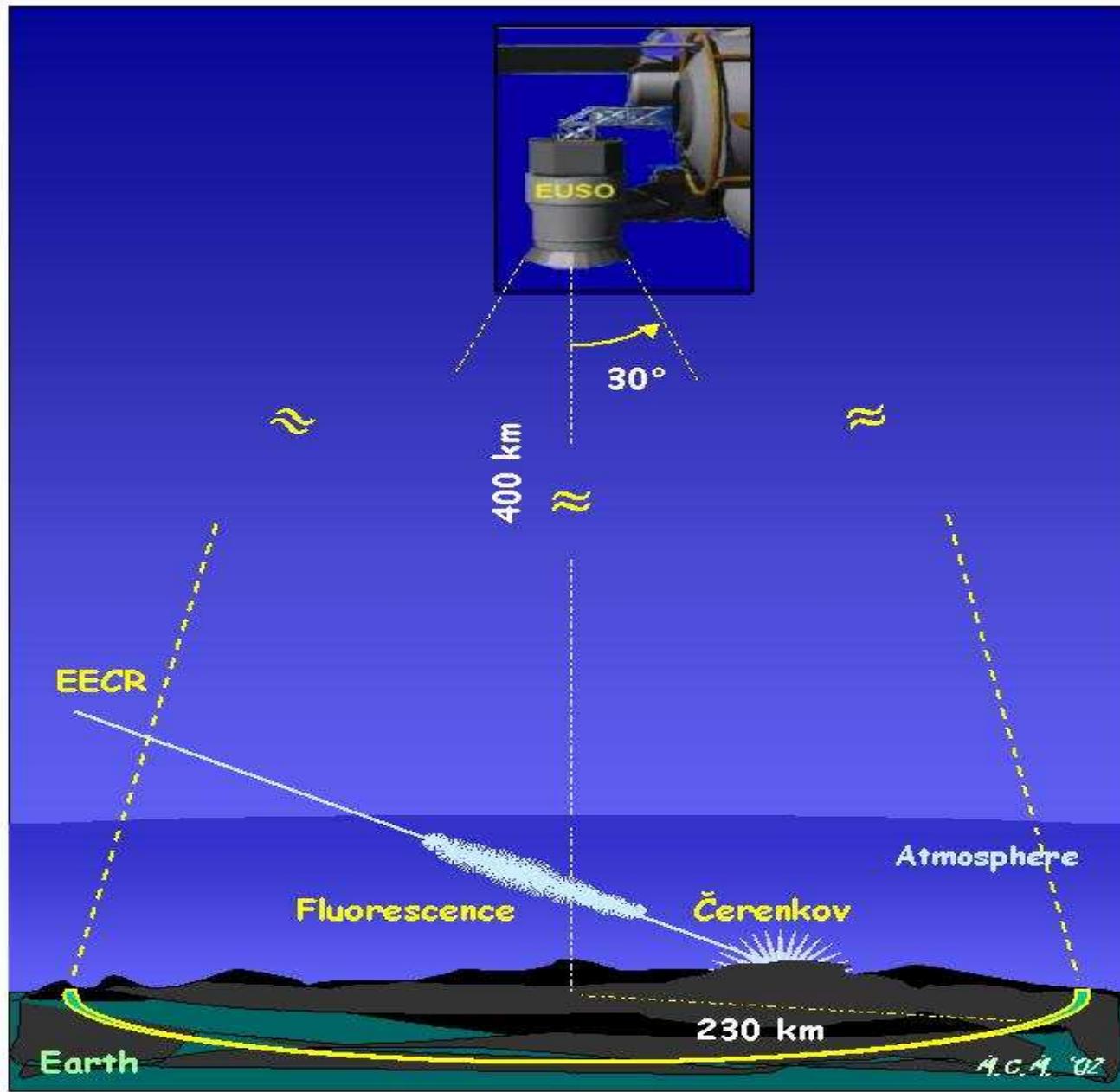
In terms of the Fourier momenta k : $dN(k) = \alpha^2 G \eta^4 \ell^{2/3} k^{-7/3} dk$.

CONCLUSIONS

- Fluxes of **diffuse cosmogenic neutrinos** strongly depend on the mass composition of UHECR. According to **Auger data** mass composition becomes **heavy** at highest energy, which dramatically decreases the predicted neutrino flux.
- **HiRes data** are compatible with **pure proton composition** and with much larger fluxes of cosmogenic neutrinos. However, these fluxes are strongly bounded by **the cascade upper limit** based on latest Fermi-LAT data. With this upper limit cosmogenic neutrinos can be detected only marginally by **JEM-EUSO** in the extreme models.
- **IceCube detector** also can detect only marginally cosmogenic fluxes in case of extreme models. However, IceCube can detect cosmogenic neutrinos from the **bright phase**. These fluxes are limited weaker by the cascade upper bound.
- **IceCube is the first detector which crossed the cascade upper bound and entered the physically allowed domain of cosmogenic neutrino fluxes.**

- In the light of new stronger limit on diffuse flux of cosmogenic neutrinos **search for the sources** becomes the priority goal of neutrino astronomy and JEM-EUSO and IceCube detectors. This task is viable even if protons constitute a small part of primary radiation. Discovery of HE neutrino radiation from **SNR** (in case of IceCube) will give the final proof of **GCR SM**, from **AGN** and **GRB** - proof of UHECR sources.
- There is an impressive progress in theoretical study of **TDs** as UHE neutrino sources: The ordinary cosmic strings, which are the simplest TDs, can produce the large fluxes with extremely high E_{\max} . This prediction directly follows from fundamental properties of the strings: existence of **cusp**, gravitational interaction of intermediate particles (higgses, dilatons, moduli) with the string field Φ and basic string parameter $\eta^2 = \mu$, satisfying $G\mu \gtrsim 10^{-20}$, while observational limits are $G\mu \lesssim 10^{-6}$.

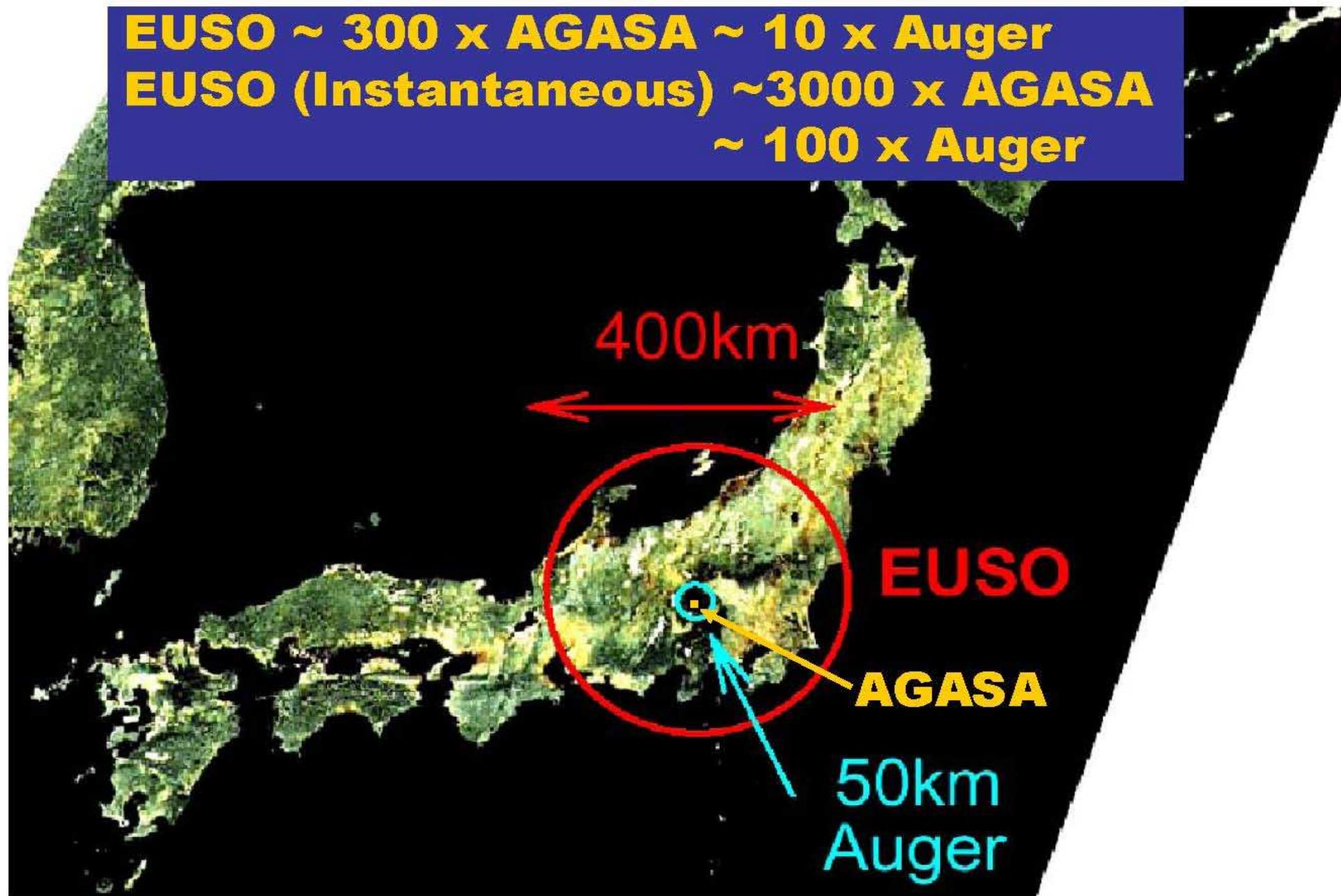
PRINCIPLES OF EUSO OBSERVATIONS



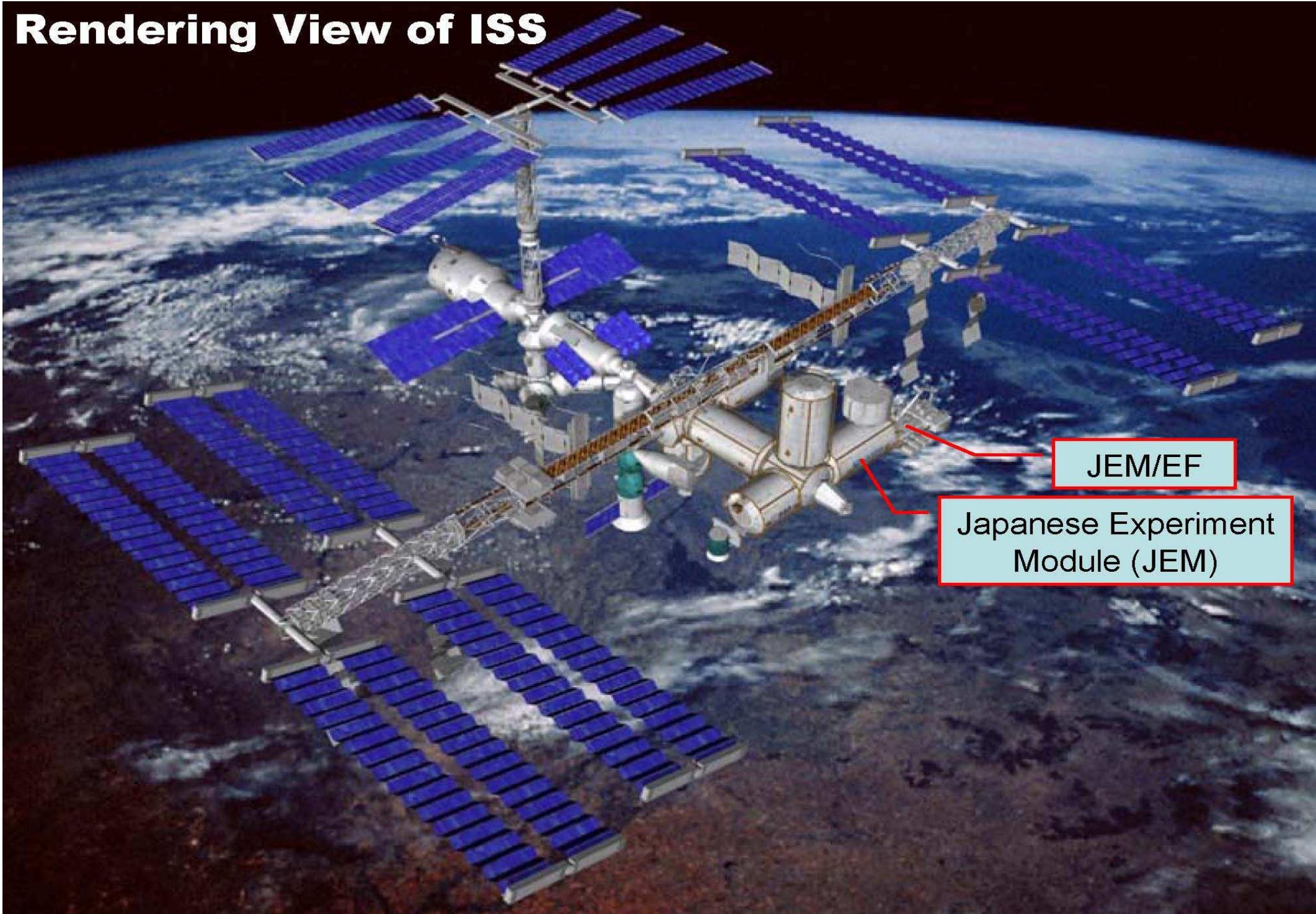
Field of View of EUSO

EUSO ~ 300 x AGASA ~ 10 x Auger

**EUSO (Instantaneous) ~3000 x AGASA
~ 100 x Auger**



Rendering View of ISS



H-IIA Launch Vehicle



**Nov. 29, 2003
Accident happened
for the H-IIA Launch Vehicle No 6**

**Feb. 26, 2005
The H-IIA Launch Vehicle No. 7 with
MTSAT-1R was launched
successfully.**

**Jan. 24, 2006
The H-IIA Launch Vehicle No. 8 with
the Advanced Land Observing
Satellite "Daichi" (ALOS) was
launched successfully.**

**Feb. 18, 2006
The H-IIA Launch Vehicle No. 9
with the Multi-functional Transport
Satellite 2 (MTSAT-2) was launched
successfully.**