#### Scattering in Fixed Point Gravity

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#### Introduction

Unitarity in Single Graviton Exchange Outlook Multiple Graviton Exchange Conclusion

#### Ingredients

# Aim: Examination of scalar scattering with Graviton exchange

- Starting point: Linearized Einstein-Hilbert-Lagrangian in d = 4 + n dimensions
  - metric  $g_{\mu\nu}$  is the carrier of the gravitational force
  - n extra dimensions are large and compact  $\rightarrow$  d-dimensional Planck Mass M = O(TeV)
  - extra dimensions generate Kaluza-Klein tower of Gravitons

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#### Fixed Point Gravity

But: KK-integral is not UV-safe  $\rightarrow$  use results from Fixed Point Gravity (cf. Weinberg '79) UV governed by interacting fixed point  $g^*$  of coupling g

$$\mathbf{I^{(n)}(x)} \propto \int_{0}^{\infty} dm \, m^{n-1} P(x,m,\Lambda) \quad , P_{\mathrm{IR}}(x,m,\Lambda) = rac{1}{x-m^2}$$

- KK-integral becomes UV-safe (no cutoff!)
- use different parametrization for IR-UV transition motivated by FPG [3]
- important parameter: crossover scale  $\Lambda$ , related to  $g^*$

#### Feynman Diagrams for Higgs (Single Graviton)



## Higgs Scattering

$$M = I^{(n)}(s) \left( \frac{n s^2}{n+2} + 2t^2 + 2ts \right) + I^{(n)}(t) \left( \frac{n t^2}{n+2} + 2s^2 + 2ts \right)$$
$$+ I^{(n)}(-s-t) \left( s^2 + t^2 - \frac{2(s+t)^2}{n+2} \right)$$

remember: 
$$I^{(n)}(x) \propto \int_{0}^{\infty} dm \, m^{n-1} P(x, m, \Lambda)$$

 application of unitarity constraints to examine breakdown energy s resp. bounds on parameters

#### Unitarity Constraints I

partial wave amplitude for J = 0

$$a_0(s) = rac{1}{16 \, \pi \, s} \int\limits_{-s}^{0} M(s,t) \, dt$$

unitarity of the S-matrix implies here:

$$|\operatorname{\mathsf{Re}}(a_0)| < 1 \quad \lor \quad |a_0| < 2$$

s-channel unitarity leads to upper bound on  $\Lambda$ :

$$\Lambda_{crit} = sup \left\{ \Lambda \ | \ \left| a_0^{(s)}(s,\Lambda) \right| < 2 \right\}$$

#### Unitarity Constraints II

t/u channel unitarity is violated at  $s_{crit}$ :

$$\left| \mathsf{Re}\left[ a_0^{(\mathsf{t},\mathsf{u})}(s_{\mathsf{crit}}) 
ight] 
ight| = 1$$

- $s_{crit}$  depends on  $\Lambda$  and n
- t/u channel unitarity always violated for large s

#### Results: Critical s



### Eikonal Approximation

#### Multiple Graviton Exchange

- for spin 2 exchanges, the eikonal approximation is the nlo contribution in the kinematic regime of low angle scattering t leq s
- eikonal approximation is the summation of an infinite class of ladder diagrams
- restriction: only Graviton exchanges, only t-channel diagrams, only leading kinematic contribution in s/t

### Eikonal Amplitude

$$A_{\text{Eik}} = -4 \pi i s \int_{0}^{\infty} db J_{0}(q b) (\exp[i \chi(b)] - 1)$$
$$\chi(b) = \frac{1}{2 s} \int_{0}^{\infty} \frac{dq}{2 \pi} J_{0}(b q) A_{\text{Born}}(q^{2})$$

- q: absolute exchanged momentum  $t = -q^2$
- b: impact parameter, [length], Fourier conjugate variable to q

#### Aims

- calculate  $\chi$  and  $A_{\rm Eik}$  and capture the essential behaviour in arbitrary extra dimensions
- compare results to other scenarios for Eikonal Graviton Exchange, namely Dimensional Regularisation (Giudice, Rattazzi, Wells), Finite Brane Width Scenario (Sjödahl, Gustafson), Effective Field Theory (Litim, Old - work in progress)
- examine unitarity behaviour



- KK-integration can be made UV-safe using results from Fixed Point Gravity
- unitarity breakdown as studied in Higgs scattering via single-graviton exchange typically occurs beyond the fundamental Planck scale: √s<sub>crit</sub> > M
- crossover scale  $\Lambda$  has upper bound:  $\Lambda_{max} = \mathcal{O}(M)$
- Eikonal calculation will give further insight into the FPG phenomenology of scalar scattering

#### References

- M. Niedermaier, The Asymptotic Safety Scenario in Quantum Gravity - An Introduction, arXiv:gr-qc/0610018v2, 19 Jul 2007
- [2] J. Brinkmann, Unitarity of Higgs Scattering in Extra Dimensional Models within Fixed Point Gravity, Diploma Thesis, Dec 2009
- [3] E. Gerwick, D. Litim and T. Plehn, Asymptotic safety and Kaluza-Klein gravitons at the LHC, arXiv:1101.5548 [hep-ph].

#### Parametrizations of the KK Integral

renormalization approach:

$$\int_{0}^{\infty} dm \, m^{n-1} \frac{1}{x - m^2} \, Z_{\rm i}^{-1}(\mu)$$

anomalous dimension approach:

$$\int_{0}^{\infty} dm \, m^{n-1} \frac{\Lambda^{2\,\Delta_i}}{(x-m^2)^{1+\Delta_i}}$$

minus approach:

$$\int_{0}^{\infty} dm \, m^{n-1} \frac{\Lambda^{2\,\Delta_{i}}}{(x-m^{2})(|x|+m^{2})^{\Delta_{i}}}$$

### Approximations for Z and $\eta$

quench:

$$Z(\mu) = egin{cases} 1 & \mu < \Lambda \ rac{\mu^{2+n}}{\Lambda^{2+n}} & \mu \geq \Lambda \end{cases}, \quad \eta(\mu) = egin{cases} 0 & \mu < \Lambda \ -2-n & \mu \geq \Lambda \end{cases}$$

linear:

$$Z(\mu) = 1 + rac{\mu^{2+n}}{\Lambda^{2+n}}, \quad \eta(\mu) = -(2+n)(1-Z^{-1}(\mu))$$

quadratic:

$$Z(\mu) = \sqrt{1 + \left(\frac{\mu^{2+n}}{2\Lambda^{2+n}}\right)^2} + \frac{\mu^{2+n}}{2\Lambda^{2+n}}, \quad \eta(\mu) = -(2+n)\frac{Z^2(\mu) - 1}{Z^2(\mu) + 1}$$

# Λ Matching

different  $\Lambda_{param}$  are matched through IR condition:

$$\lim_{x\to 0}\int_0^\infty dm\,m^{n-1}P_i(x,m,\Lambda_i)=\lim_{x\to 0}\int_0^\infty dm\,m^{n-1}P_k(x,m,\Lambda_k)$$