

***EFFECT OF JETS ON***  
 ***$v_4/v_2^2$  RATIO AND***  
***CONSTITUENT QUARK SCALING***  
***IN RELATIVISTIC HEAVY-ION***  
***COLLISIONS***

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# OUTLINE

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**I. Motivation**

**II. HYDJET++ model (hydro + jets)**

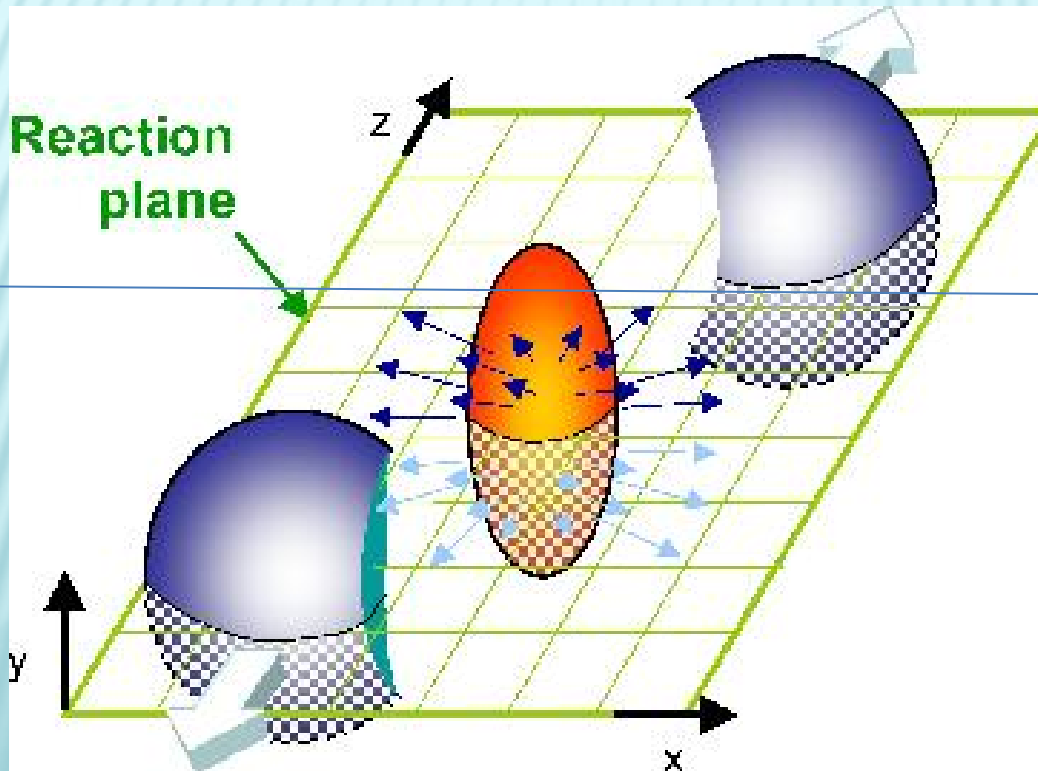
**III. Model results for the ratio  
 $v_4/(v_2)^2$  at RHIC and LHC**

**IV. NCQ-scaling at RHIC and LHC**

# I. $v_4/(v_2)^2$ ratio

## Anisotropic flow

$$\frac{dN}{d\varphi} = \frac{1}{2\pi} \left( 1 + \sum_{n=1}^{\infty} 2v_n(p_t) \cos[n(\varphi - \psi_r)] \right)$$



# Predictions

*N. Borghini, J.-Y. Ollitrault, PLB 642 (2006) 227*

- Within the *approximation* that the particle momentum  $\mathbf{p}$  and the fluid velocity  $\mathbf{v}$  are parallel (valid for *large momentum*  $p_{\perp}$  and *low freeze-out temperature*  $T$ )

$$dN/d\varphi = \exp(2\varepsilon p_{\perp} \cos(2\varphi)/T)$$

- Expanding to order  $\varepsilon$ , the  $\cos(2\varphi)$  term is

$$v_2 = \varepsilon p_{\perp}/T$$

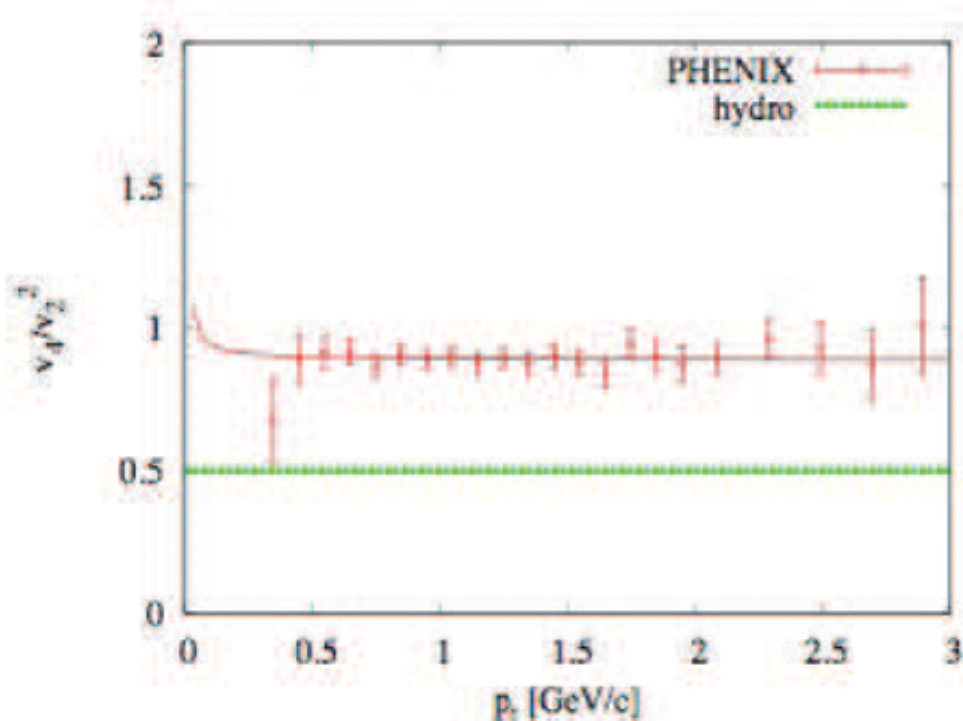
- Expanding to order  $\varepsilon^2$ , the  $\cos(4\varphi)$  term is

$$v_4 = \frac{1}{2} (v_2)^2$$

Hydrodynamics has a universal prediction for  $v_4/(v_2)^2$  !

Should be independent of equation of state, initial conditions, centrality, rapidity, particle type

## Comparison with data



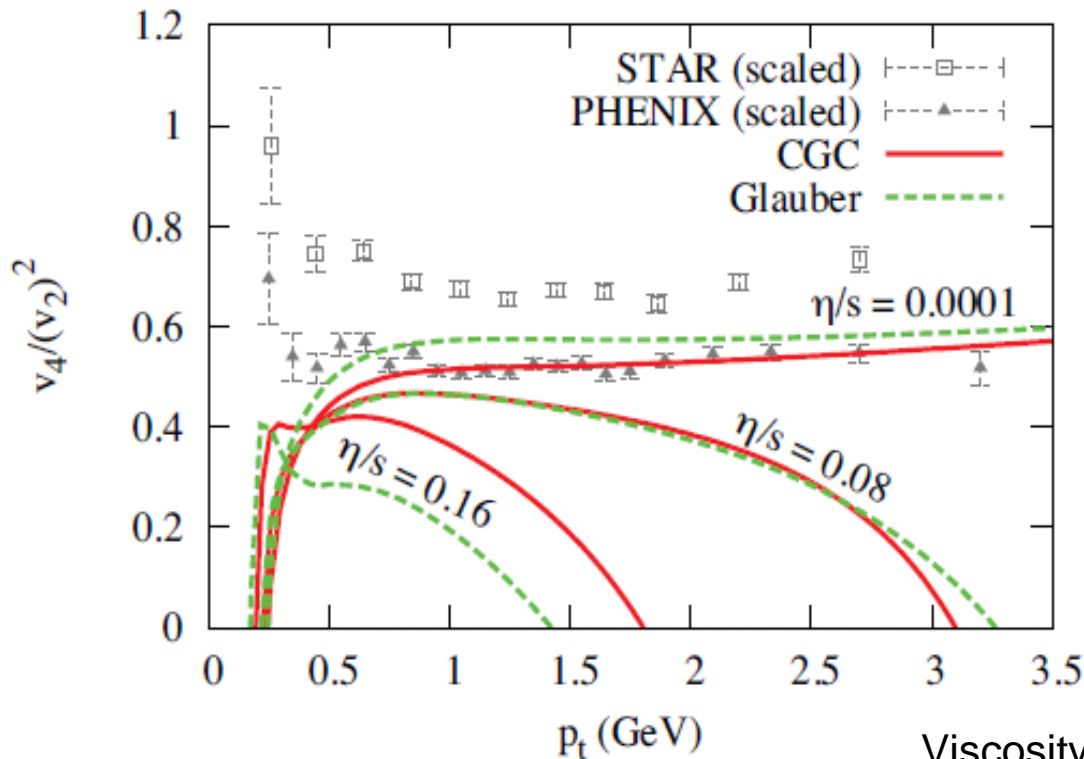
PHENIX data for charged pions

Au-Au collisions at 100+100 GeV

20-60% most central

The ratio is significantly larger than 0.5.  
Can this be explained by viscous corrections?

# Effects of initial profile and viscosity

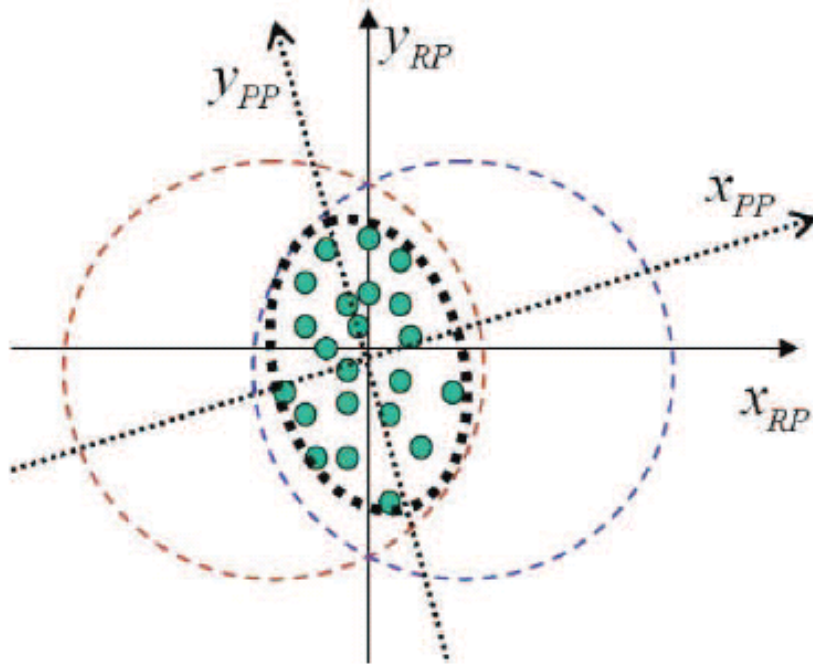


Initial profile has little effect although eccentricities differ.

results strongly depend on viscosity

Viscosity lowers  $v_4/(v_2)^2$  for a realistic  $T_f$

## Eccentricity fluctuations



Depending on where the participant nucleons are located within the nucleus at the time of the collision, the actual shape of the overlap area may vary: the **orientation and eccentricity** of the ellipse defined by participants fluctuates.

Assuming that  $v_2$  scales like the eccentricity, **eccentricity fluctuations** translate into  **$v_2$  fluctuations**

Eccentricity fluctuation can be computed in MC Glauber model or derived from experiment by comparing different methods for flow calculation.

# Why $\varepsilon$ fluctuations change $v_4/v_2^2$

Experimentally, no direct measure of  $v_2$  and  $v_4$

$v_2$  and  $v_4$  are measured via azimuthal correlations

$$v_2 \text{ from } \langle \cos(2\phi_1 - 2\phi_2) \rangle = \langle (v_2)^2 \rangle$$

$$v_4 \text{ from } \langle \cos(4\phi_1 - 2\phi_2 - 2\phi_3) \rangle = \langle v_4 (v_2)^2 \rangle$$

Experimentally measured

$$\frac{v_4}{v_2^2} = \frac{\langle v_4 (v_2)^2 \rangle}{\langle (v_2)^2 \rangle^2} = \frac{1}{2} \frac{\langle (v_2)^4 \rangle}{\langle (v_2)^2 \rangle^2} > \frac{1}{2}$$

fluctuations

hydro

Similar results obtained using Event Plane method



**II. HYDJET++ =  
FASTMS + HYDJET**

# HYDJET++ event generator

I.Lokhtin, L.Malinina, S.Petrushanko, A.Snigirev, I.Arsene, K.Tywoniuk,  
Comp. Phys. Commun.180 (2009) 779-799 (arXiv:0809.2708[hep-ph])

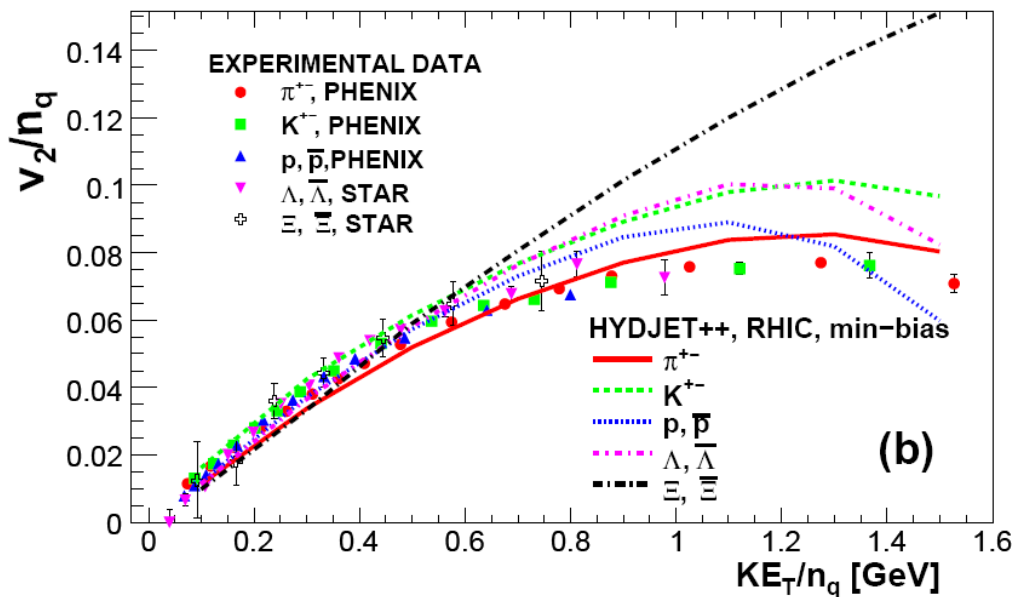
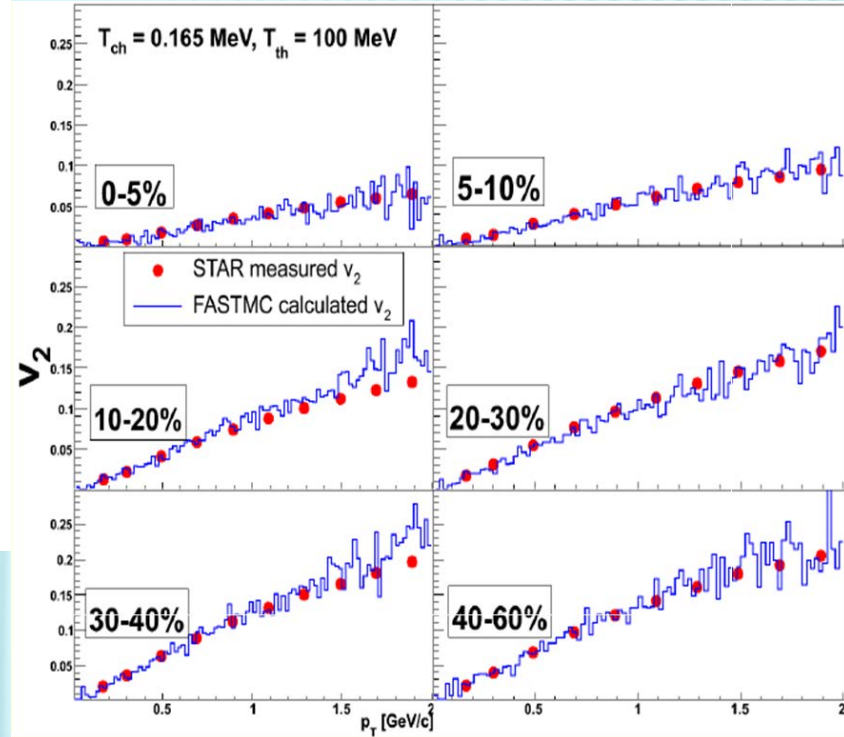
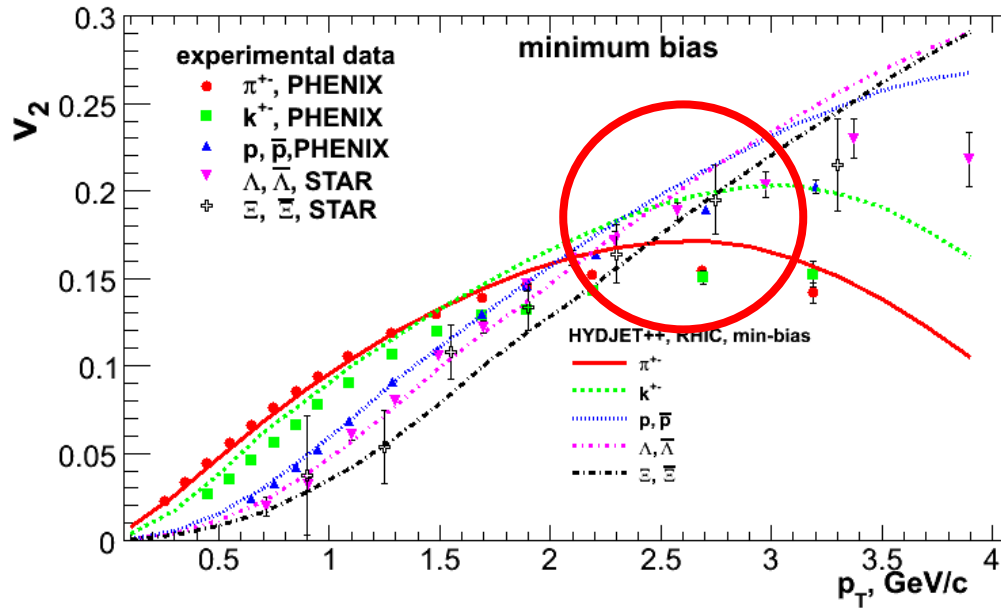
- The soft part of HYDJET++ event represents the "thermal" hadronic state.
  - ✓ multiplicities are determined assuming thermal equilibrium
  - ✓ hadrons are produced on the hypersurface represented by a parameterization of relativistic hydrodynamics with given freeze-out conditions
  - ✓ chemical and kinetic freeze-outs are separated
  - ✓ decays of hadronic resonances are taken into account (360 particles from SHARE data table) with "home-made" decayer

*the model reproduces soft hadroproduction features at RHIC  
(particle spectra, elliptic flow, HBT)*

- The hard, multi-partonic part of HYDJET++ event is identical to the hard part of Fortran written HYDJET (PYTHIA6.4xx + PYQUEN1.5).  
PYQUEN event generator is used for simulation of rescattering, radiative and collisional energy loss of hard partons in expanding quark-gluon plasma created in ultrarelativistic heavy ion AA collisions. HYDJET++ includes nuclear shadowing correction for parton distributions (important at LHC!)  
Impact-parameter dependent parameterization of *nuclear shadowing* (K.Tywoniuk, I.Arsene, L.Bravina, A.Kaidalov and E.Zabrodin, Phys. Lett. B 657 (2007) 170)

# RHIC DATA VS. HYDJET++ MODEL

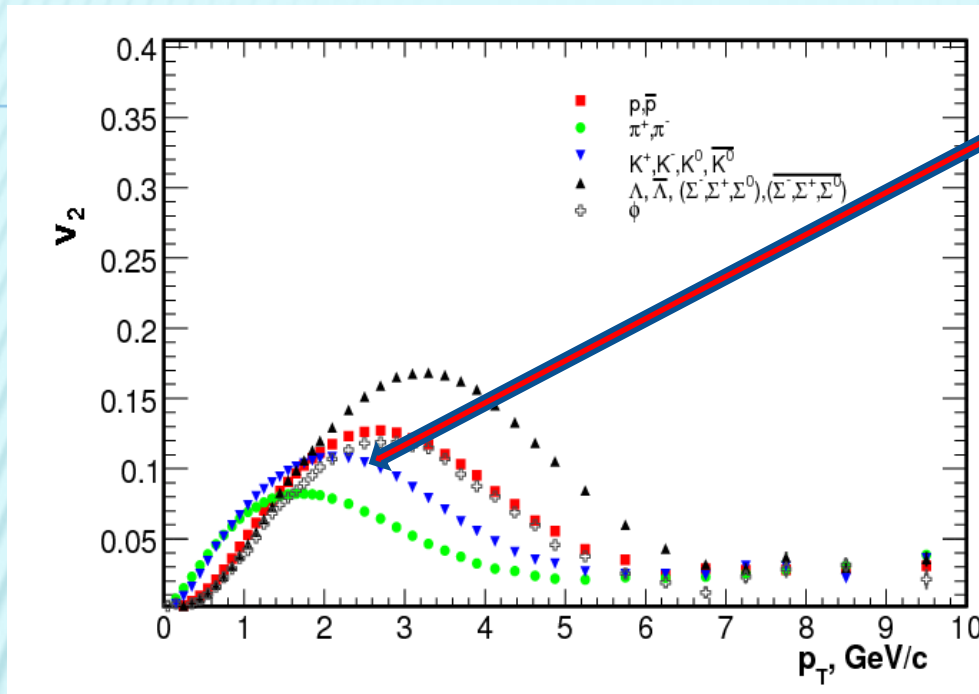
Au+Au @ 200 AGeV



Elliptic flow

G. Eyyubova et al., PRC 80 (2009) 064907;  
N.S. Amelin et al., PRC 77 (2008) 014903

# $V_2$ in HYDJET++ for different particles (centrality 30%)



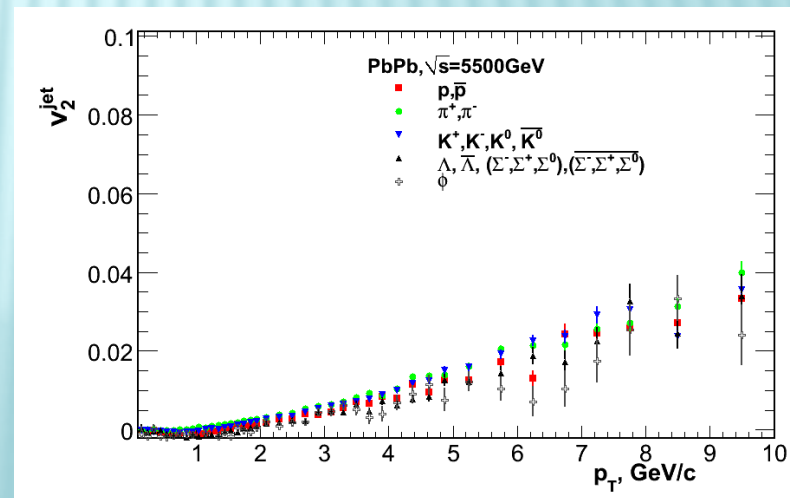
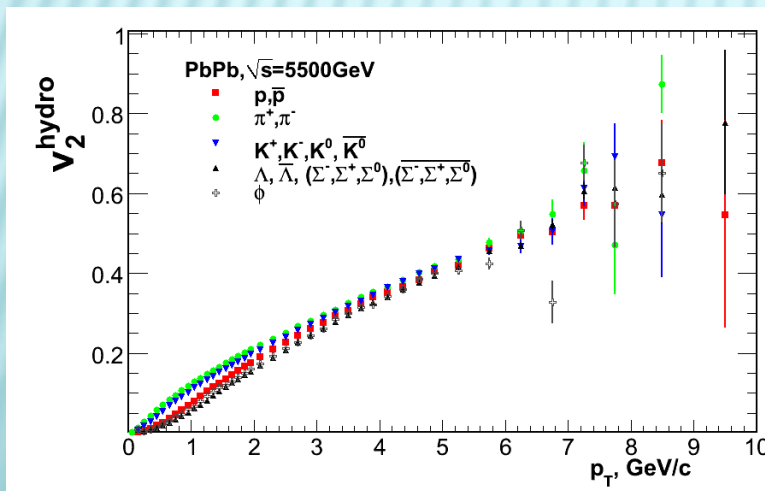
*Mass ordering in soft  $p_T$  regions then breaks.*

Why?

Hydrodynamics gives **mass ordering** of  $v_2$ .  
The model possesses **crossing of baryon and meson branches.**

*Hydrodynamics*

*Jet part + quenching*



Interplay between hydrodynamics and jets

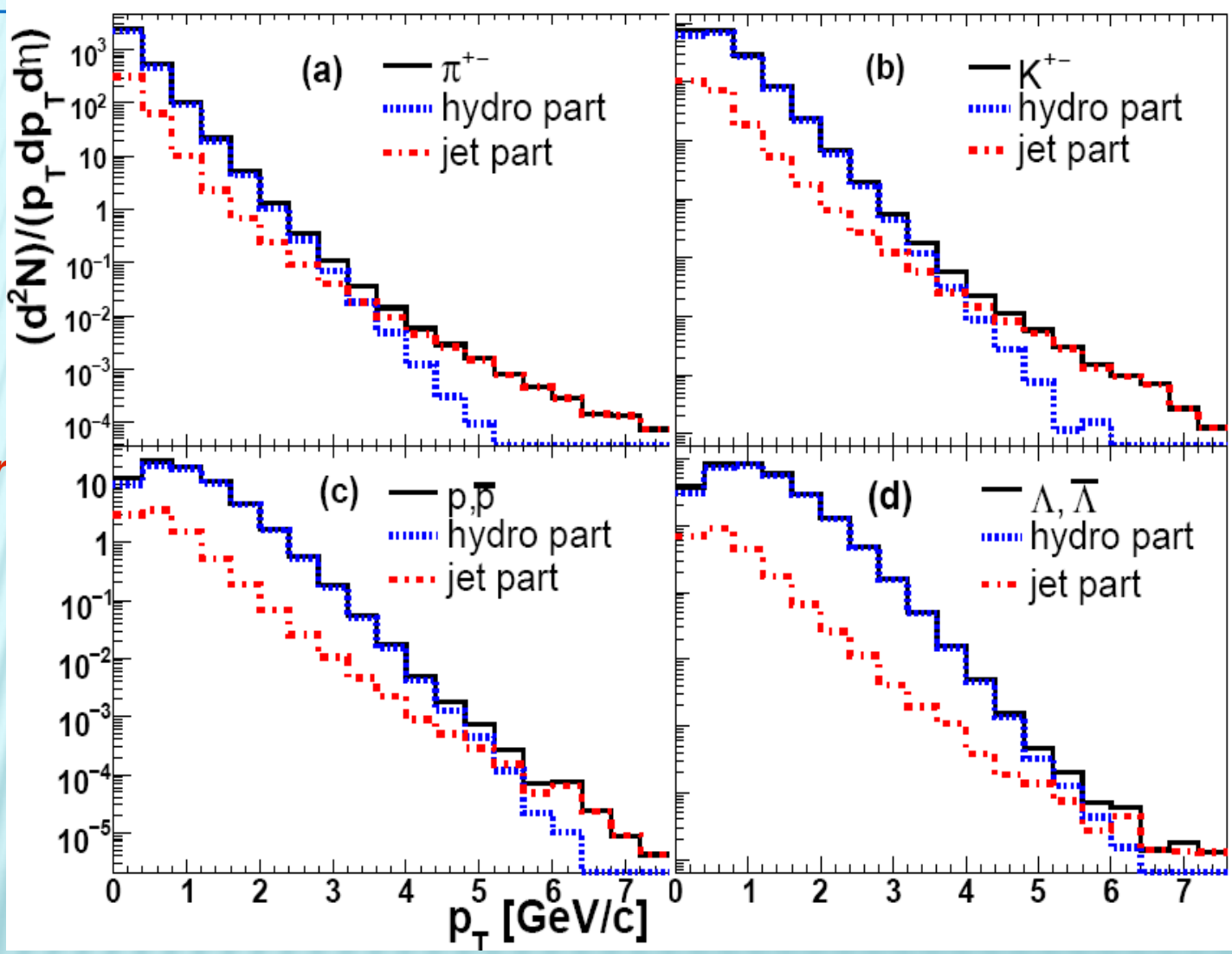
# The $p_T$ spectra of $\pi, K, p, \Lambda$ with HYDJET++ model, $\sqrt{s}=200\text{GeV}$

The slope for the hydro part depends strongly on mass:

- the heavier the particle -- the harder the spectrum



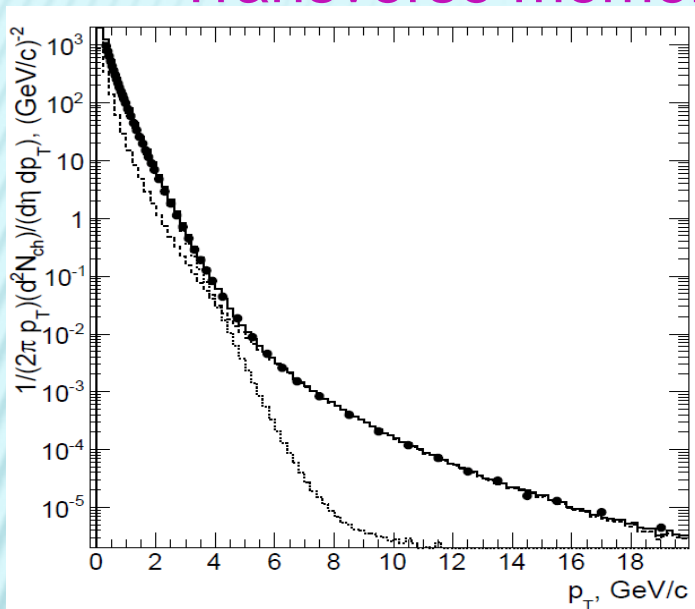
The hydro part dies out earlier for light particles than for heavy ones



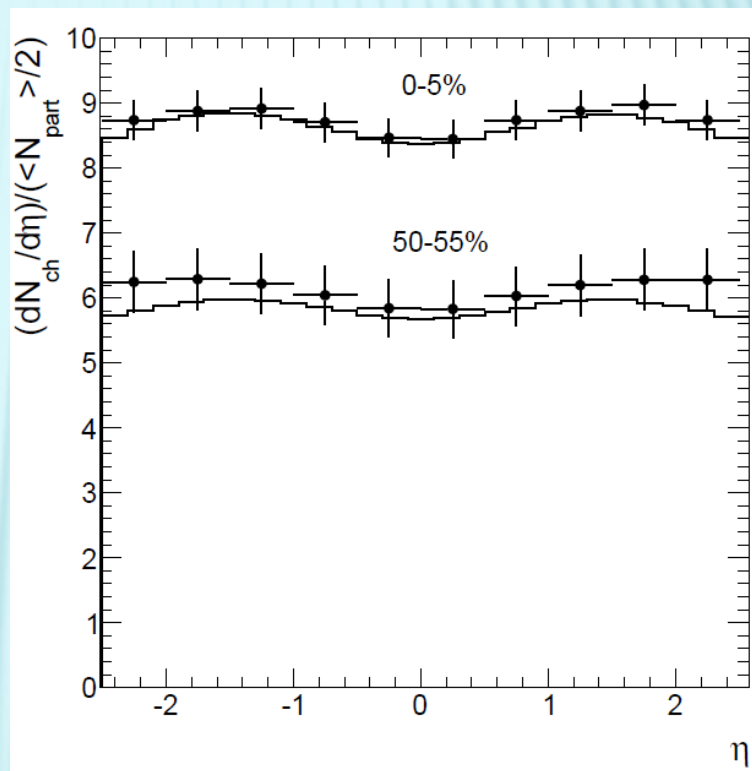
# LHC DATA VS. HYDJET++ MODEL

Transverse momentum

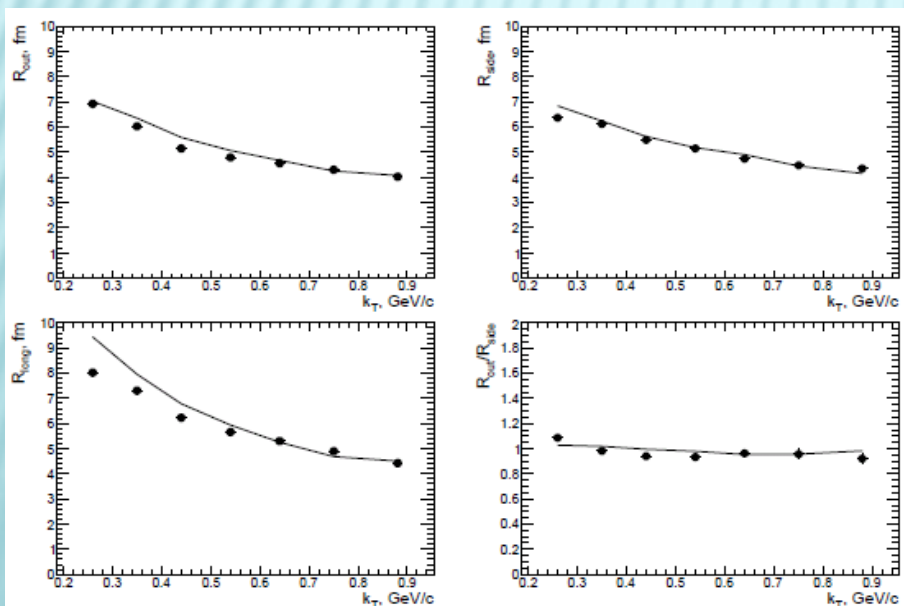
Pb+Pb @ 2.76 ATeV



Rapidity



I. Lokhtin et al., arXiv:1204.4820

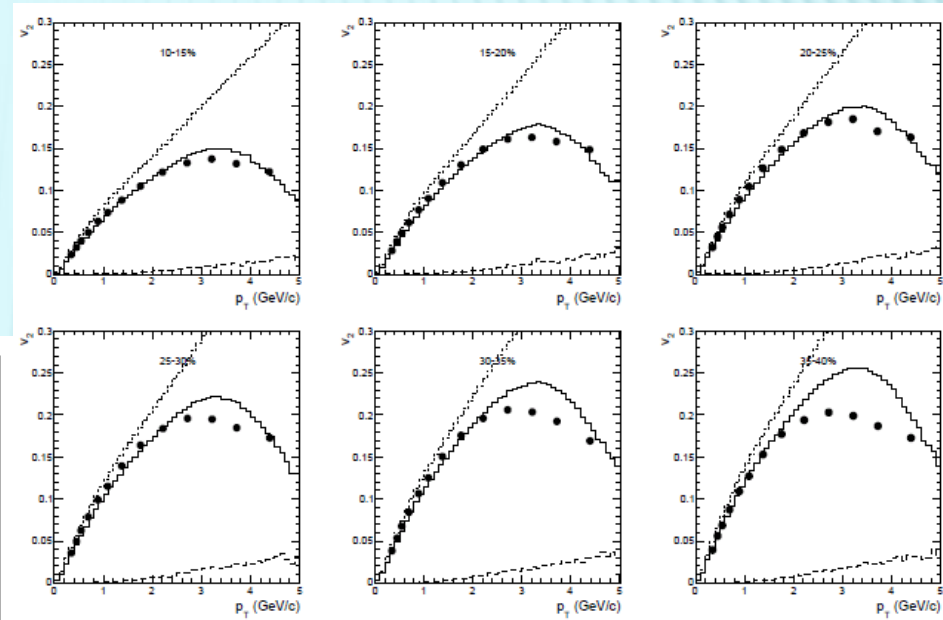
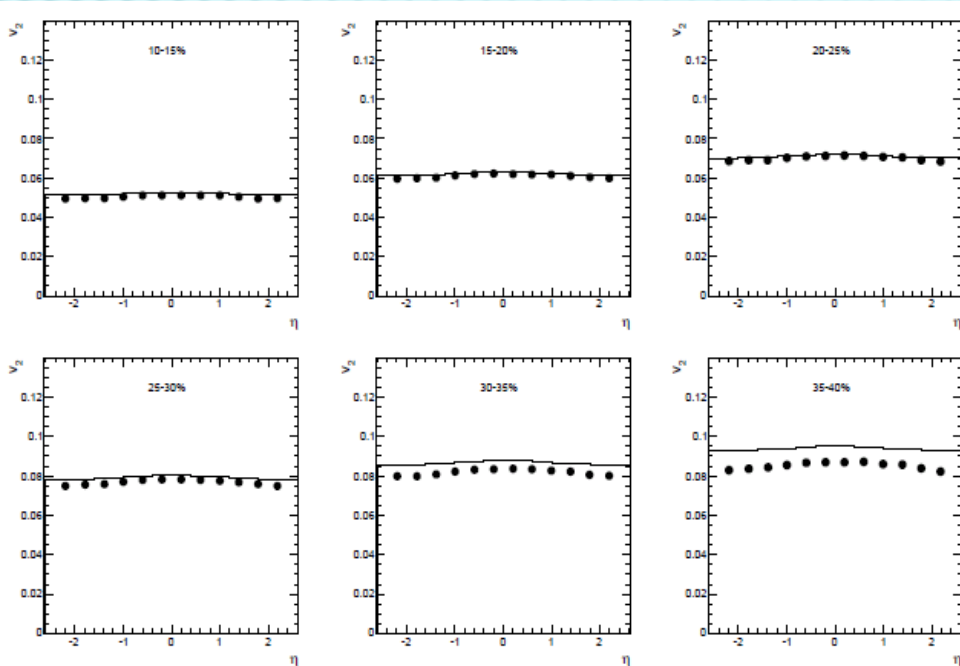


Correlation radii (femtoscscopy)

# LHC DATA VS. HYDJET++ MODEL

Pb+Pb @ 2.76 ATeV

## Elliptic flow



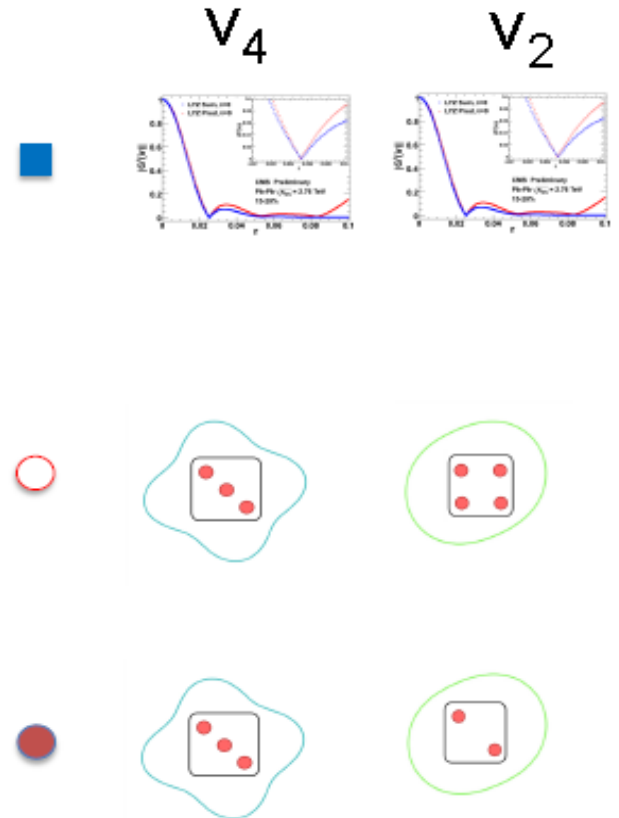
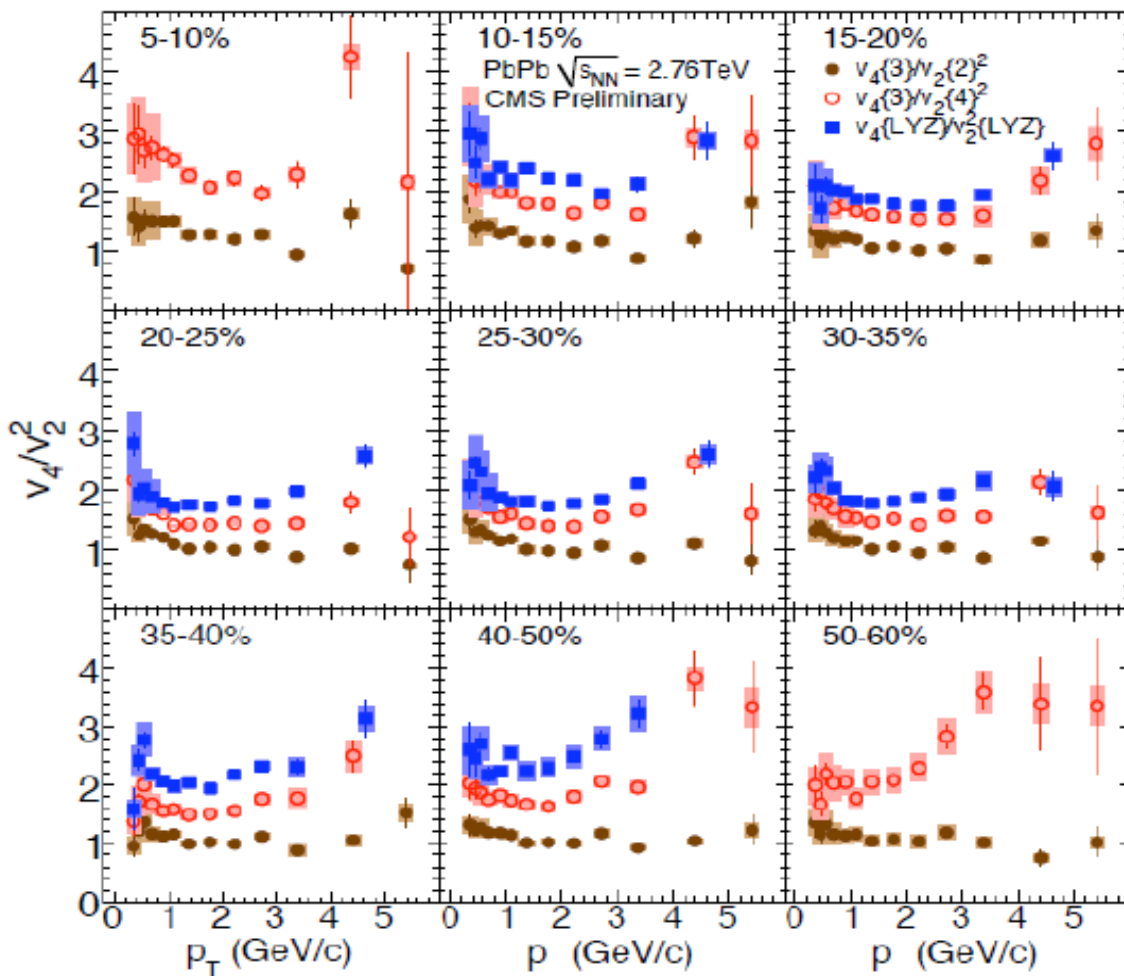
I. Lokhtin et al., arXiv:1204.4820

Model gives a fair description of various observables at both RHIC and LHC

### III. $V_4/(V_2*V_2)$ RATIO



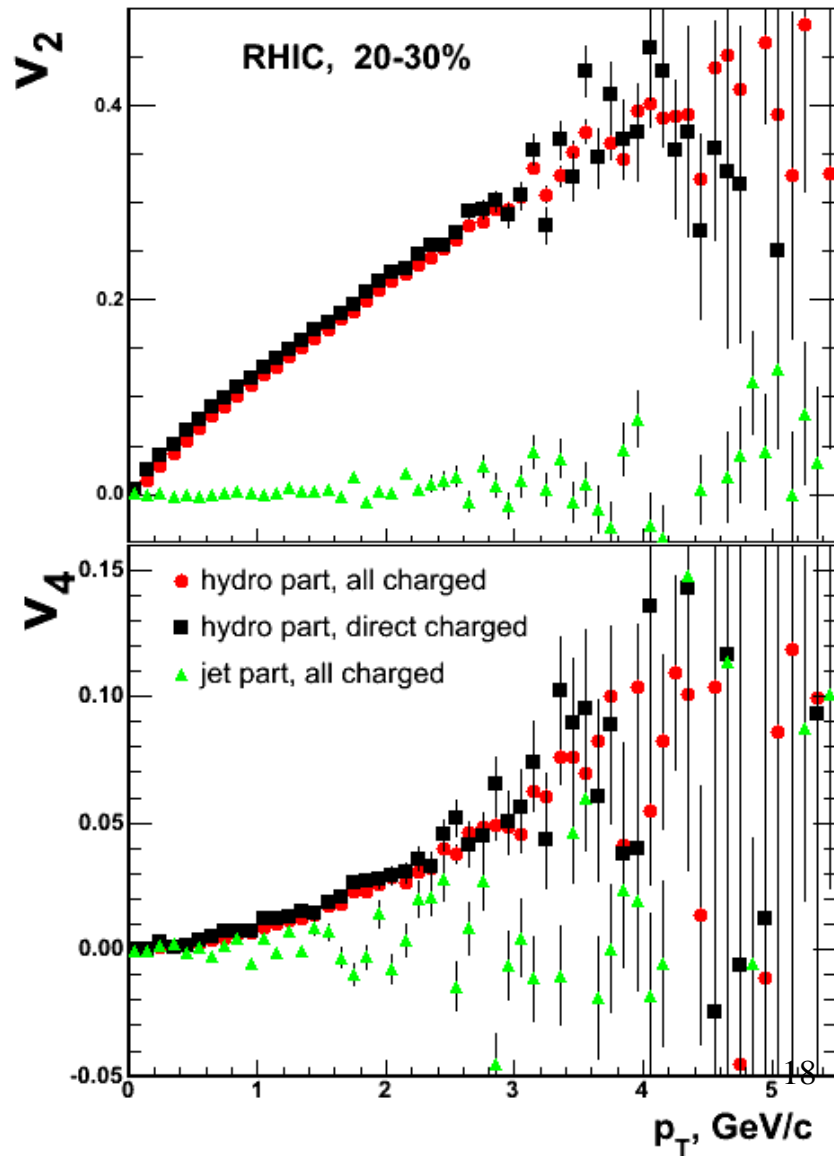
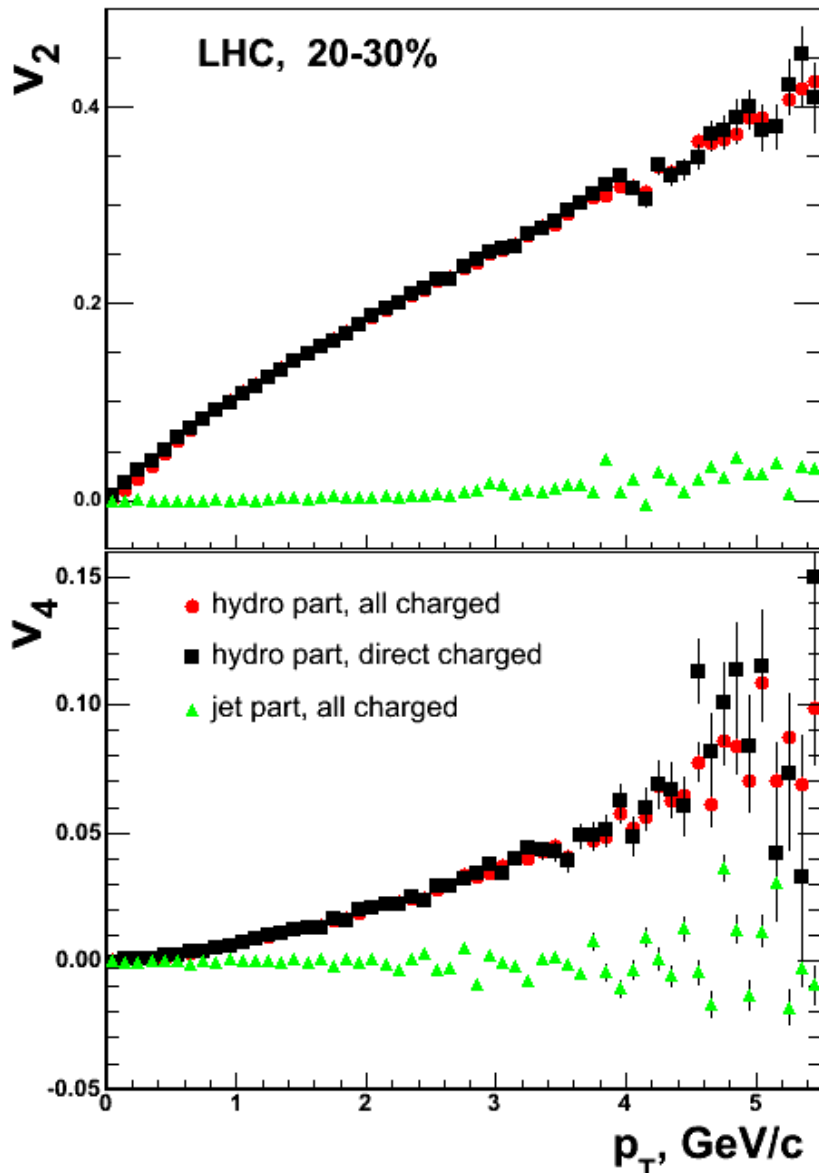
# $v_4 / v_2^2(p_T)$ at mid-rapidity $|\eta| < 0.8$



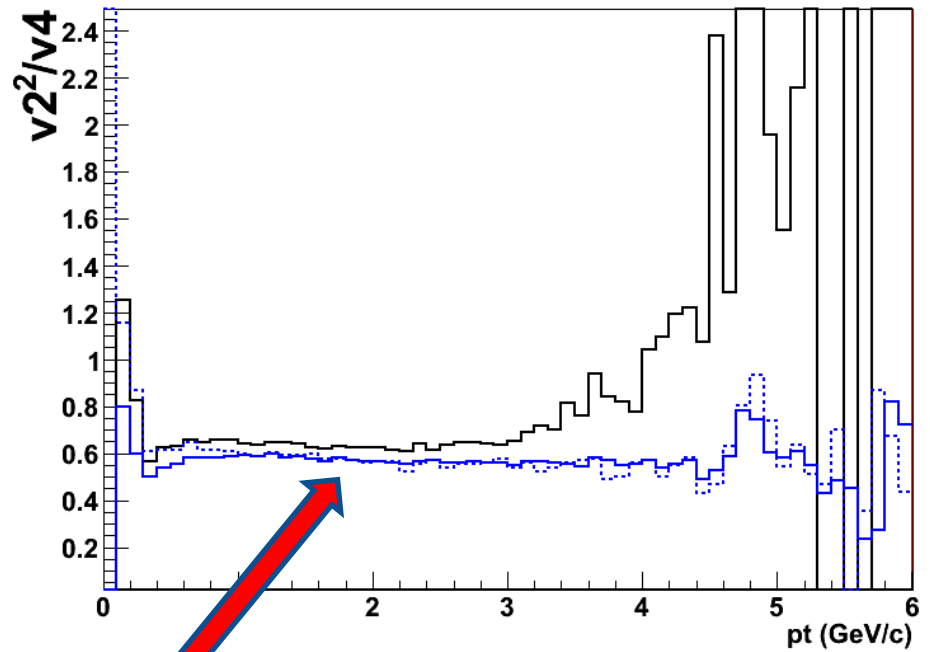
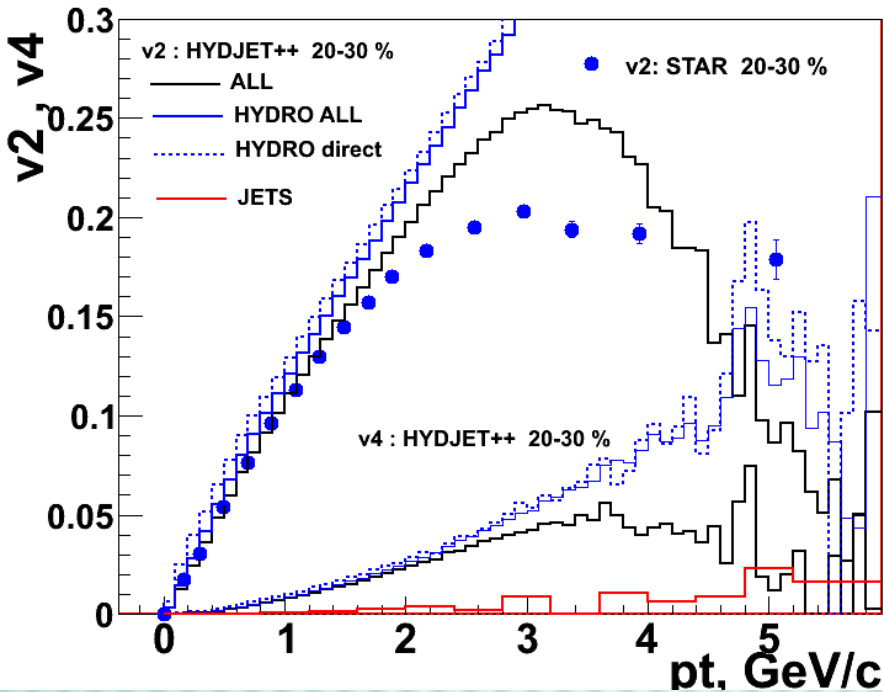
Significantly higher than RHIC: experimental method dependent

# HYDJET++

Effects to be studied: resonance decay and hard part influence



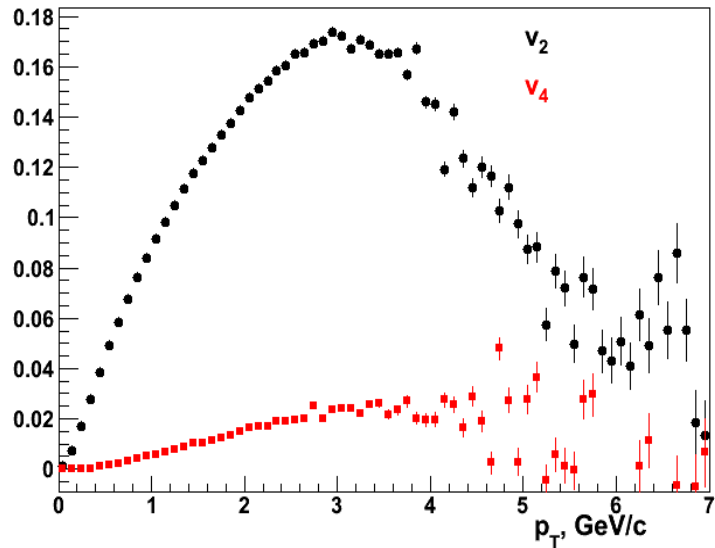
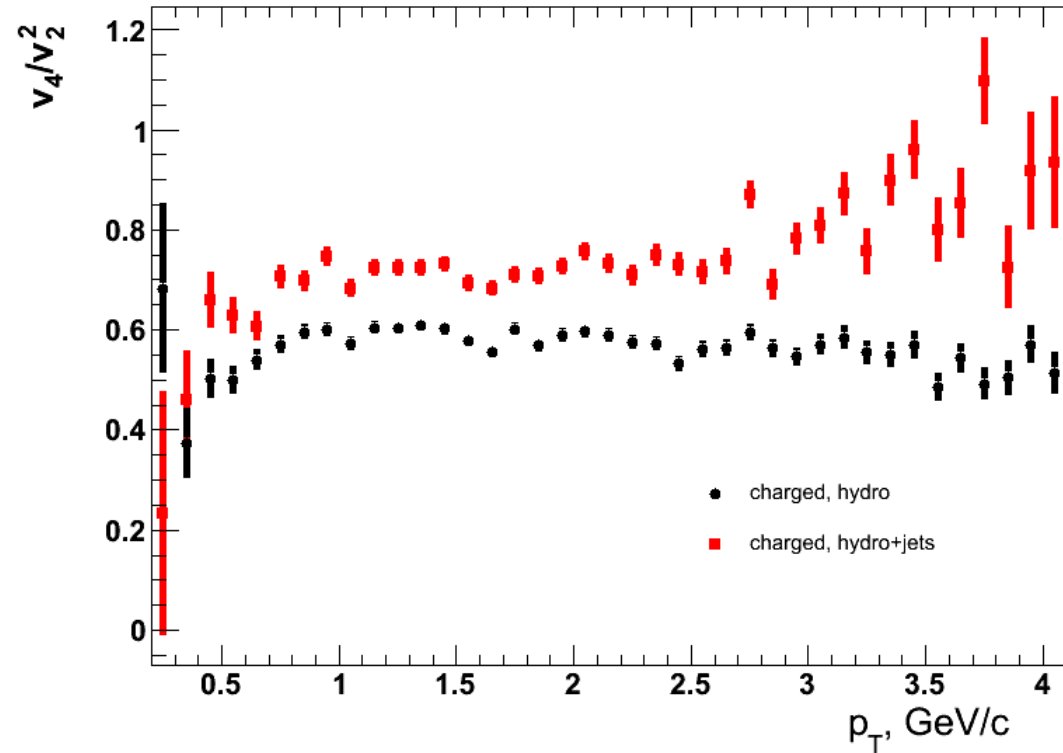
# HYDJET++ RESULTS FOR RHIC



Jets increase the ratio

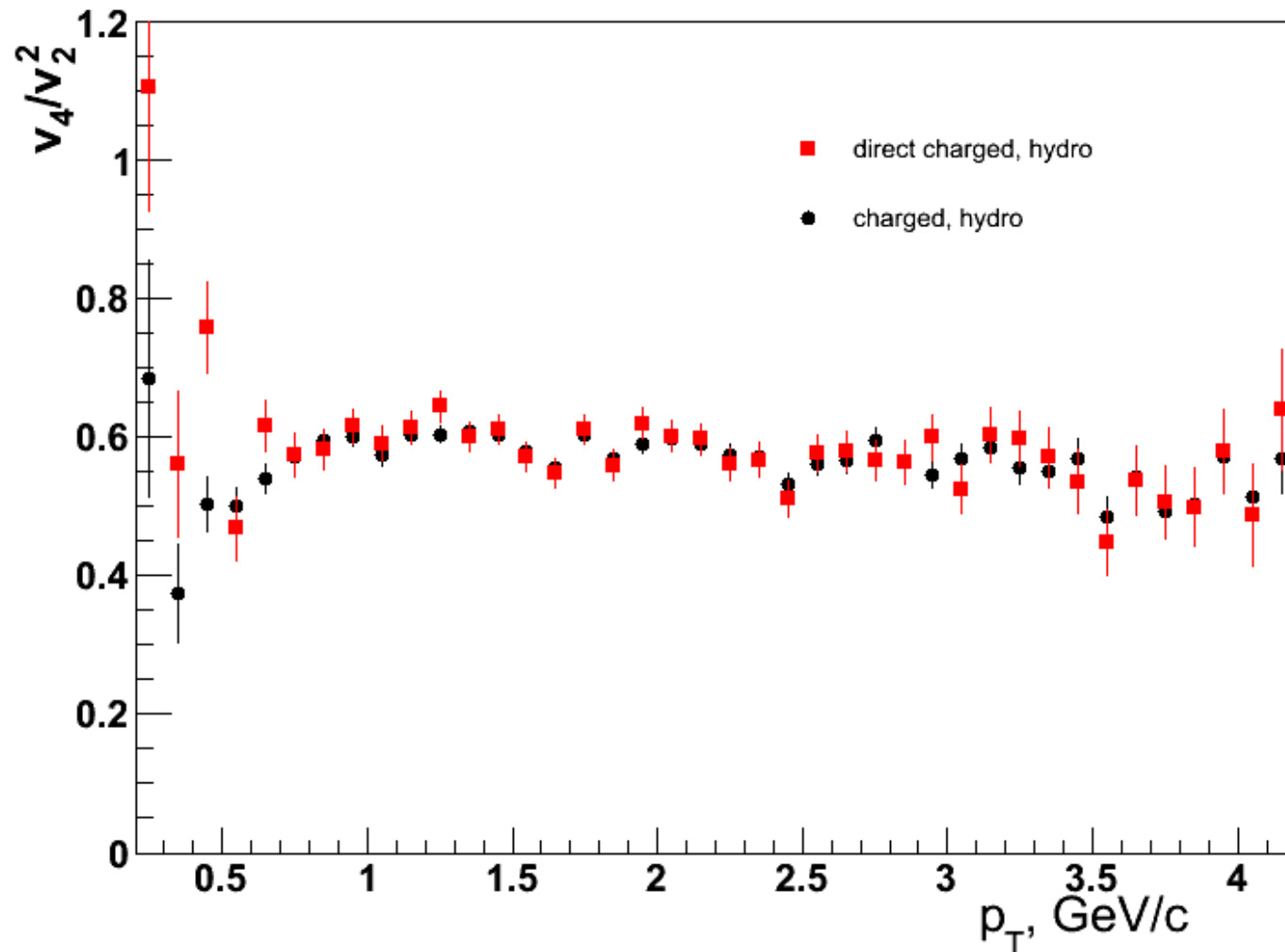
# HYDJET++ RESULTS FOR LHC

The same tendency is observed in Pb+Pb at LHC



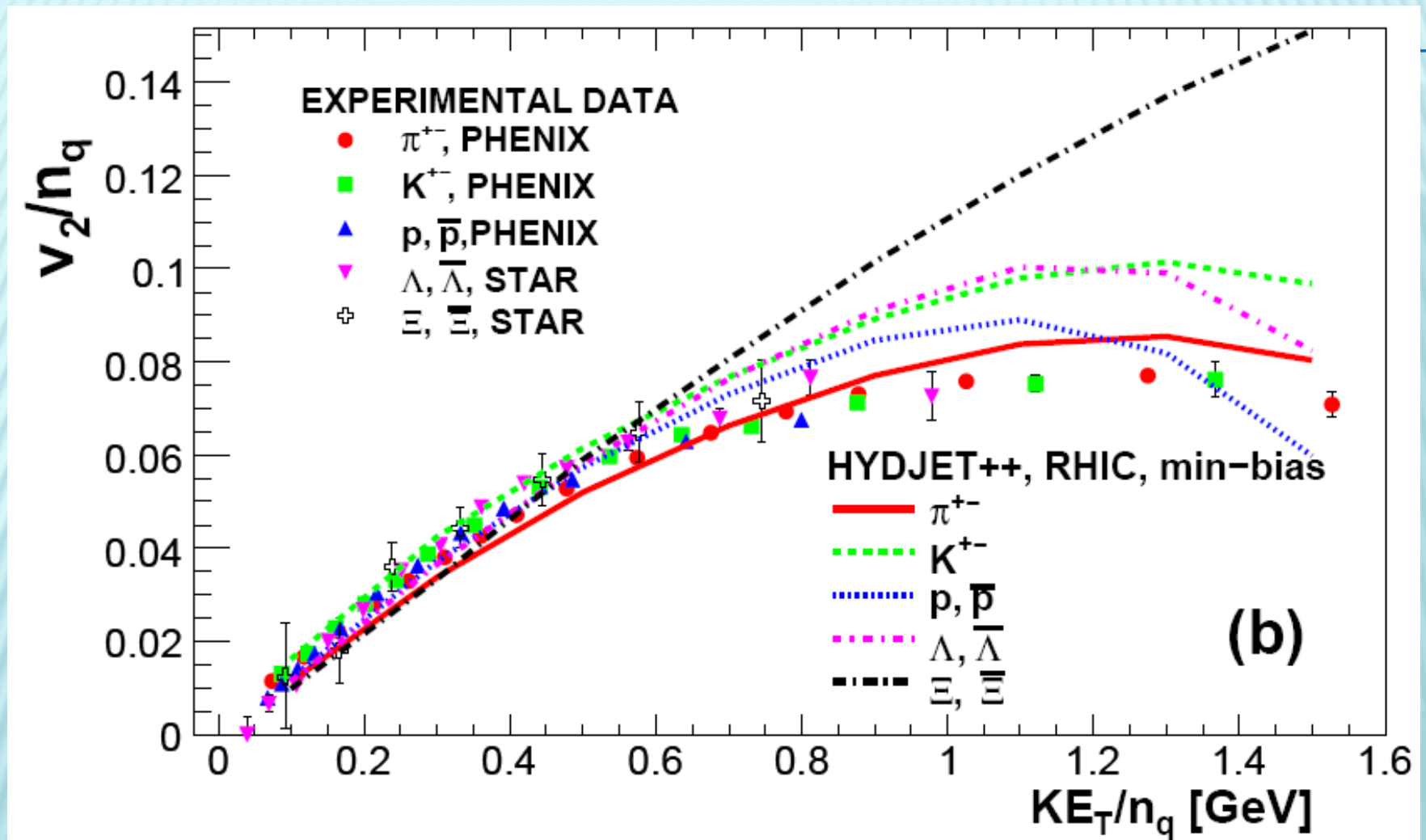
Still, the ratio is below 1

# DECAYS OF RESONANCES PLAY MINOR ROLE



# **IV. Number-of- constituent- quark (NCQ) scaling**

## COMPARISON WITH RHIC DATA

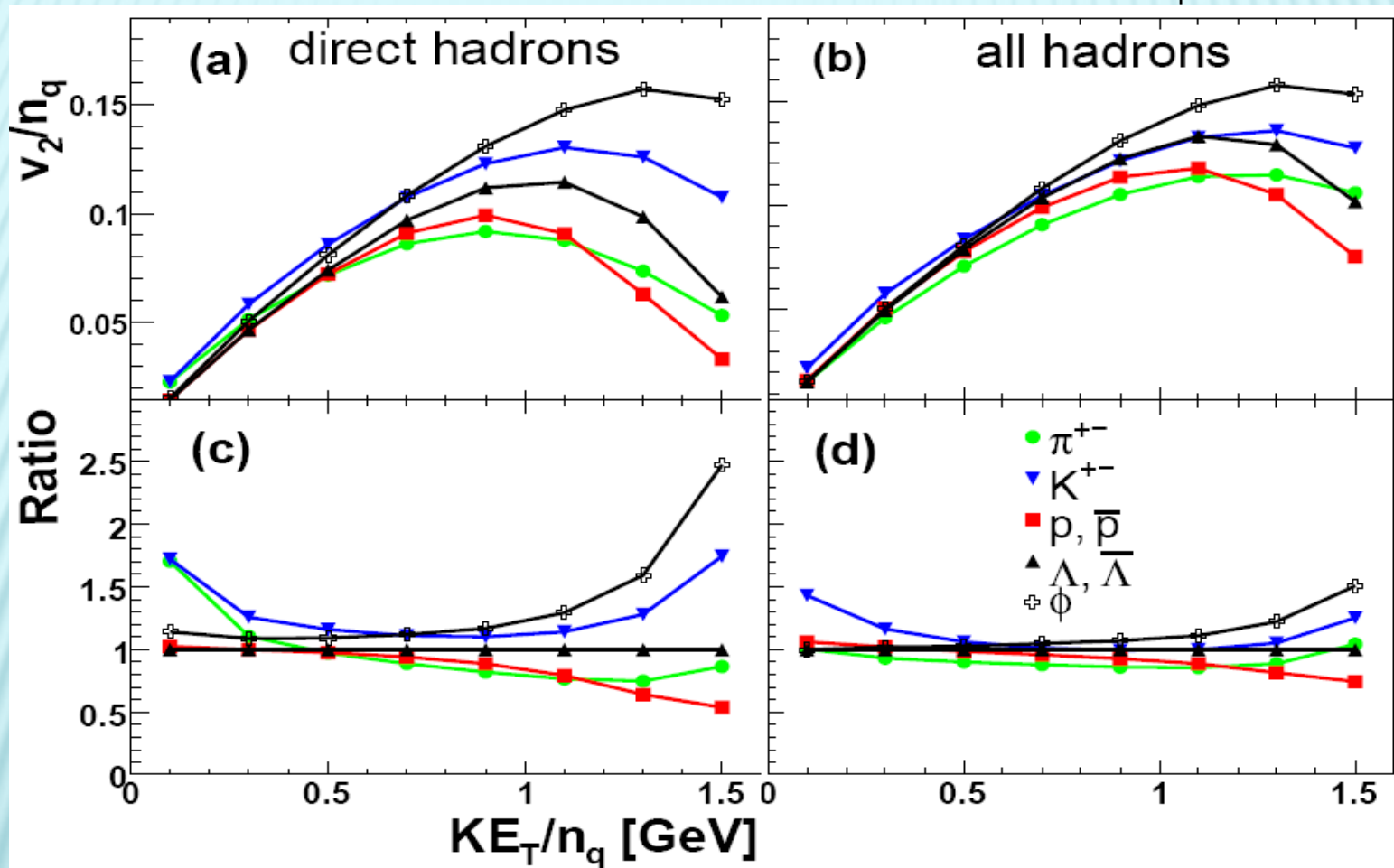


The agreement seems to be good at  $KE_T/n_q < 0.7$  GeV

# Number-of-constituent-quark scaling at RHIC

Direct particles: scaling is not good.

All particles:  $KE_T/n_q$  scaling



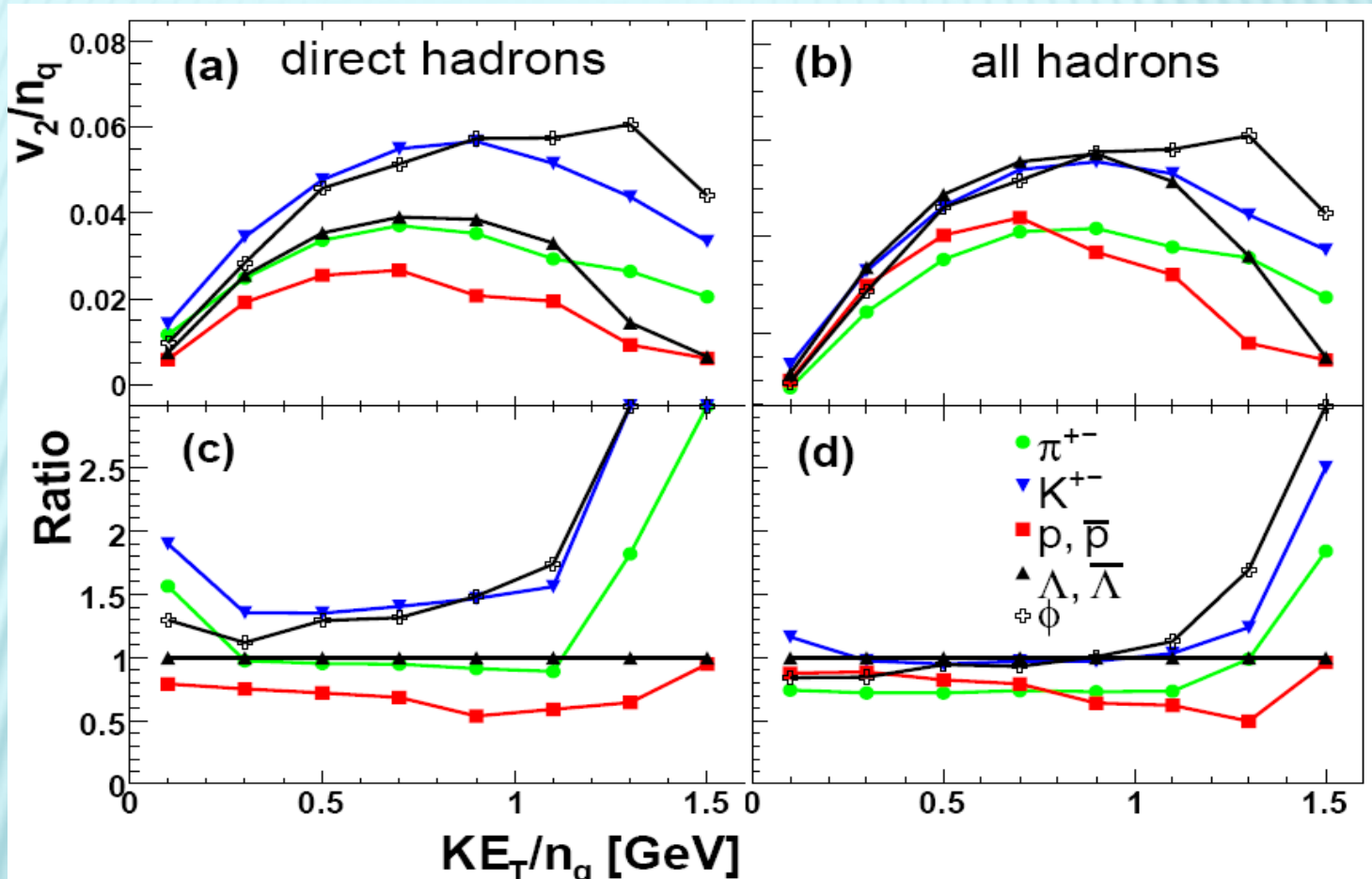
One of the explanations of  $KE_T/n_q$  scaling is partonic origin of the elliptic flow. *However, final state effects (such as resonance decays and jets) may also lead to appearance of the scaling*



# NCQ scaling at LHC

No scaling for  
direct particles

Appearance of the approximate  
scaling for all particles



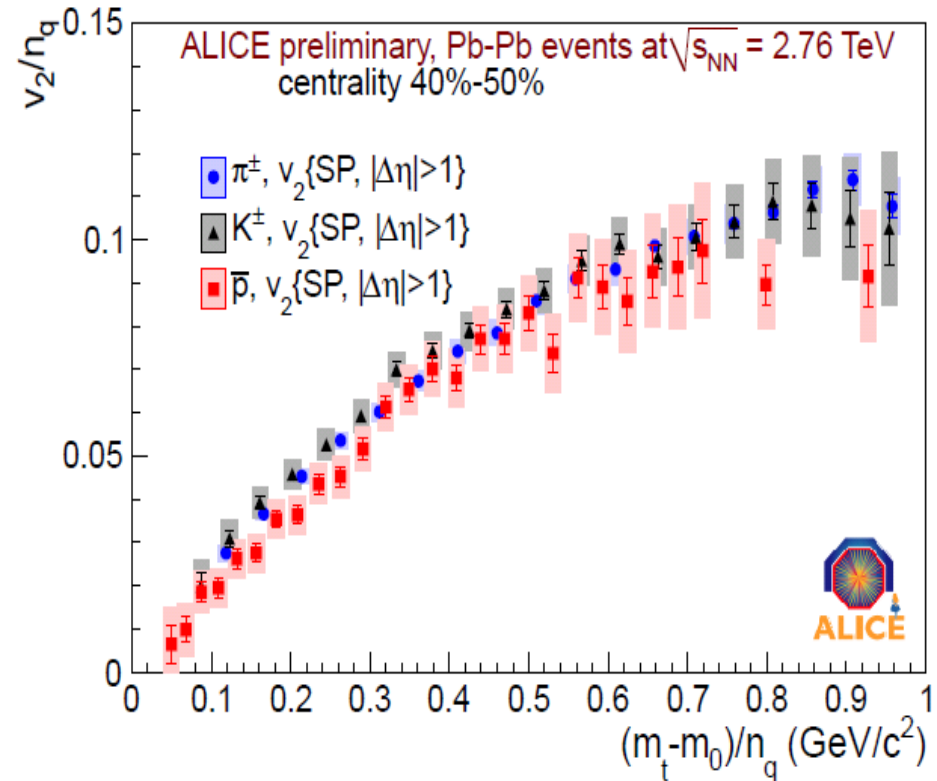
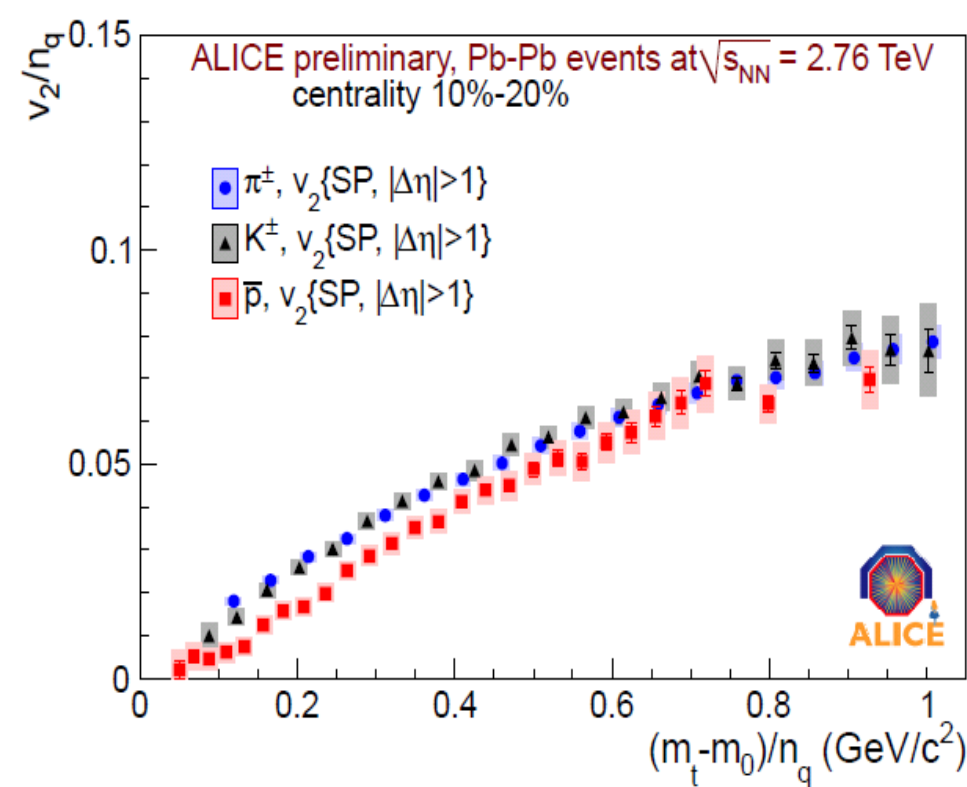
**LHC: NCQ scaling will be only approximate (prediction, 2009)**

# Experimental results (LHC)

ALICE collab., M. Krzewicki et al., JPG 38 (2011) 124047

## Semi-central collisions

## Semi-peripheral collisions



The NCQ scaling is indeed only approximate (2011)

# CONCLUSIONS

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The HYDJET++ model allows to investigate flow of hydro and jet parts separately, to look at reconstruction of pure hydro flow and its modification due to jet part.

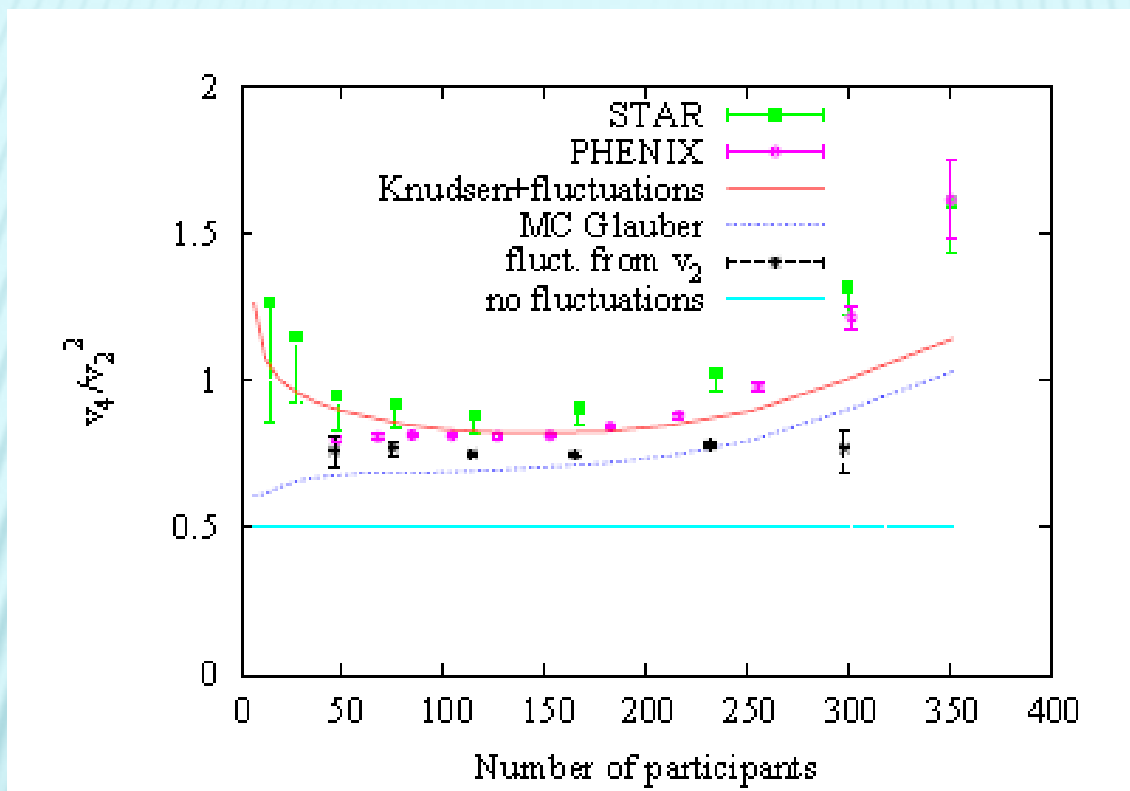
- *Jets result to increase by 25% - 30% of the ratio  $v_4/(v_2*v_2)$*
- *Eccentricity fluctuations can increase the ratio by factor 1.5*
- *Jets + eccentricity fluctuations are enough to explain RHIC data*
- *For LHC we can explain 75% of the signal. Other effects are needed*
- *The predicted violation of the NCQ scaling at LHC is observed*

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# Back-up Slides

# Effects of flow fluctuations and partial thermalization

M. Luzum, C. Gombeaud, J.-Y. Ollitrault, Phys.Rev.C81:054910,2010.

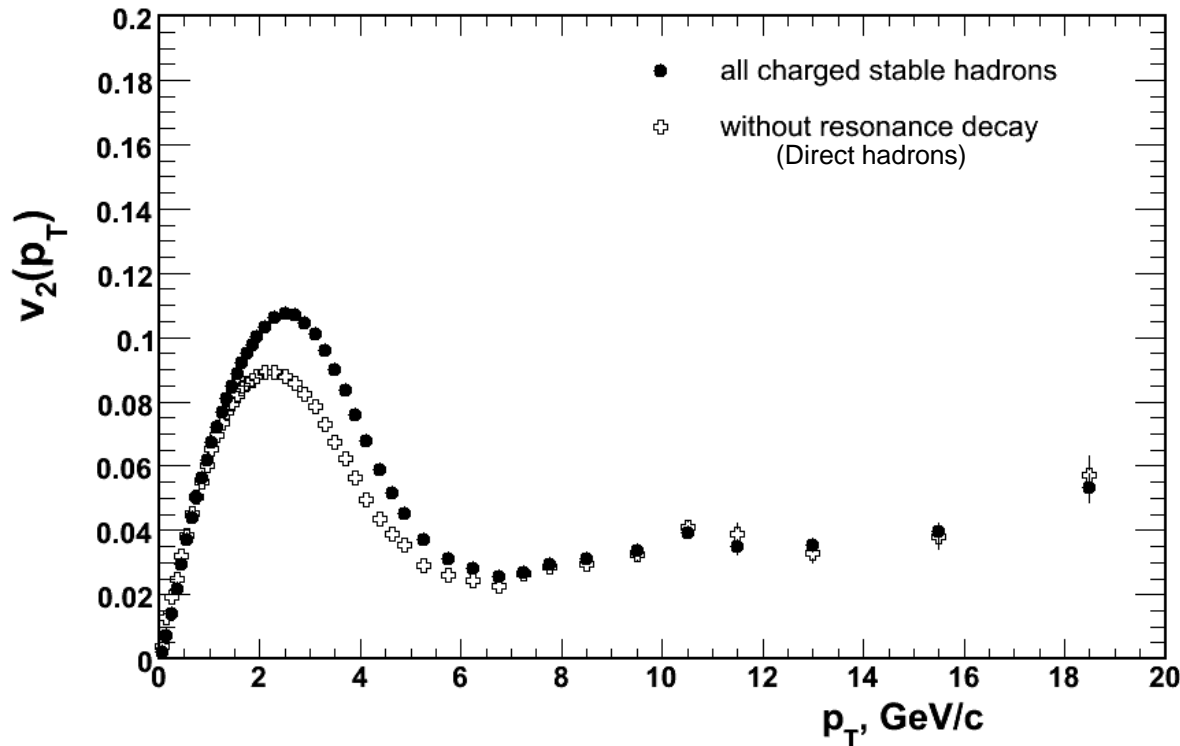


Stars: with fluctuations inferred from the difference between  $v_2\{2\}$  and  $v_2\{LYZ\}$ .

Dotted line: eccentricity fluctuations from a Monte-Carlo Glauber

# **III . INFLUENCE OF RESONANCE DECAYS**

## Influence of resonance decay on $v_2$ value



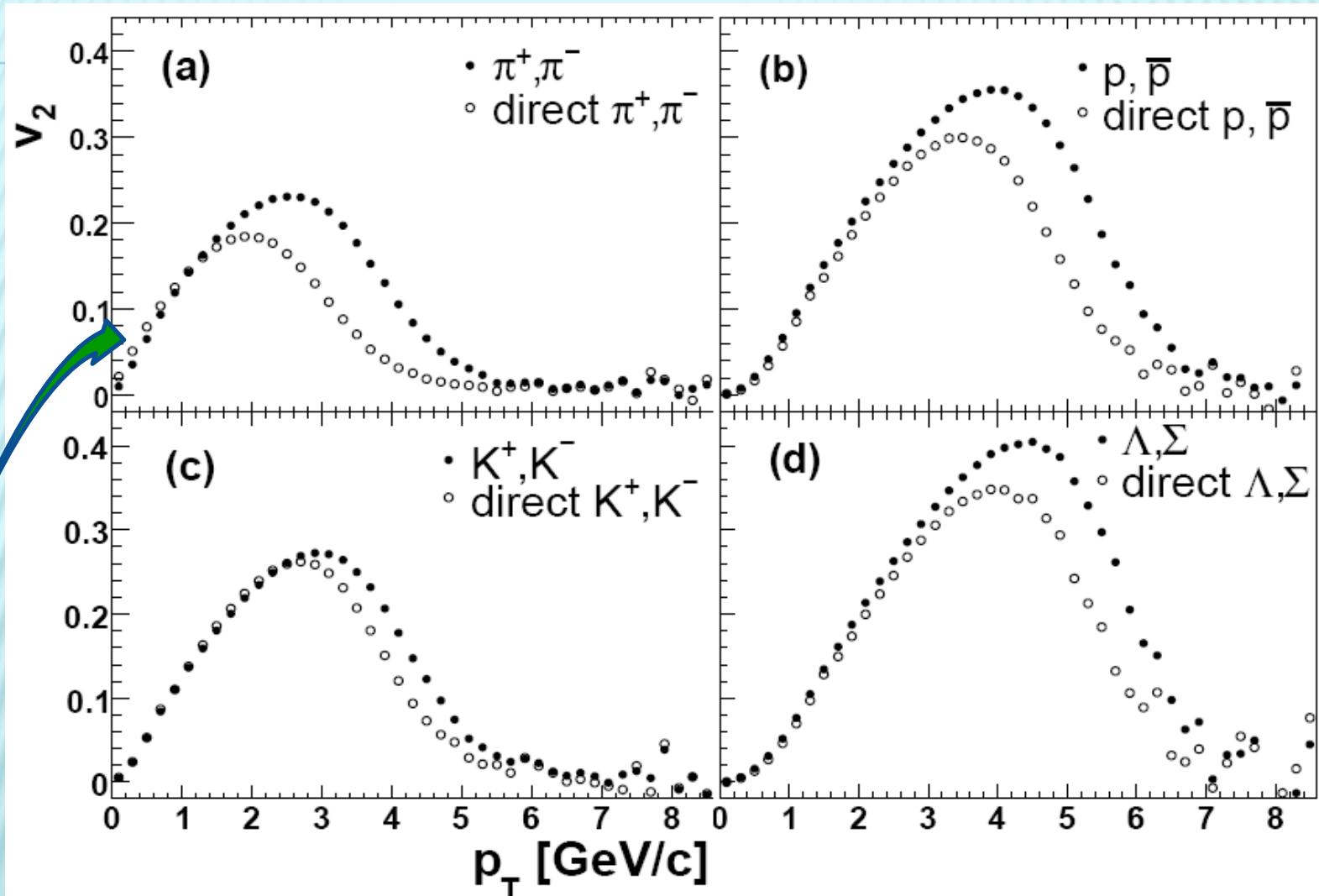
PbPb collisions,  $c=30\%$

The elliptic flow of directly produced particles is smaller than that for all particles.

TABLE I: Yields of the particles produced directly and with resonance decays,  $5.6 \cdot 10^6$  events,  $c=42\%$ , midrapidity

	$\pi^\pm$	$K + \bar{K}$	$p + \bar{p}$	$\Lambda + \bar{\Lambda} + \Sigma + \bar{\Sigma}$	$\phi$
all	860	185	63.8	42.3	6.55
direct	169	81.4	18.6	14.2	6.5
direct %	20 %	44 %	30 %	39 %	99 %

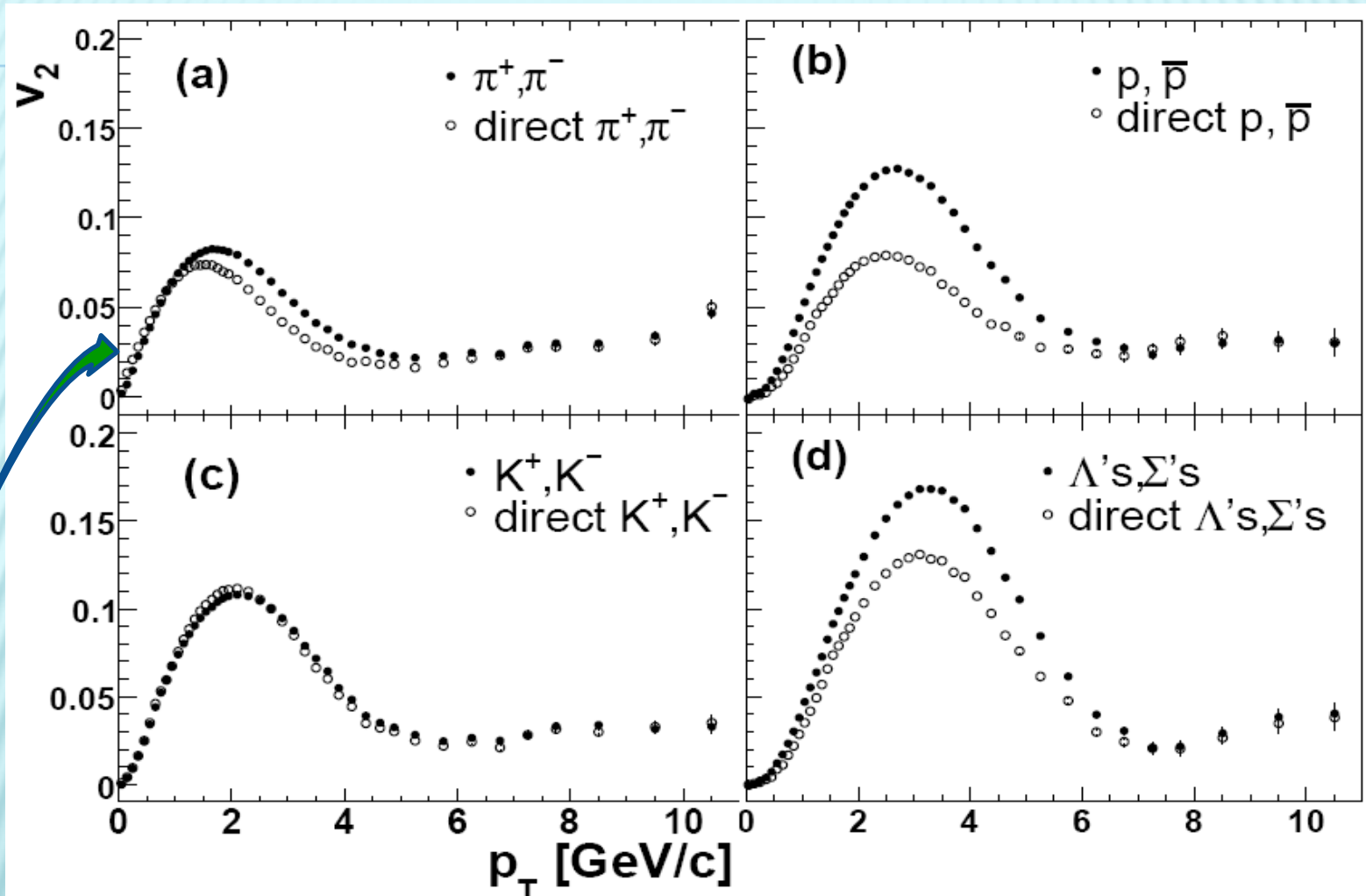
# Influence of resonance decays for different type of particles at RHIC



**Pions and kaons:** the resulting flow is weaker at low-pt and larger at high-pt  
**Baryons:** the resulting flow is stronger than the flow of direct particles



# Influence of resonance decays for different type of particles at LHC

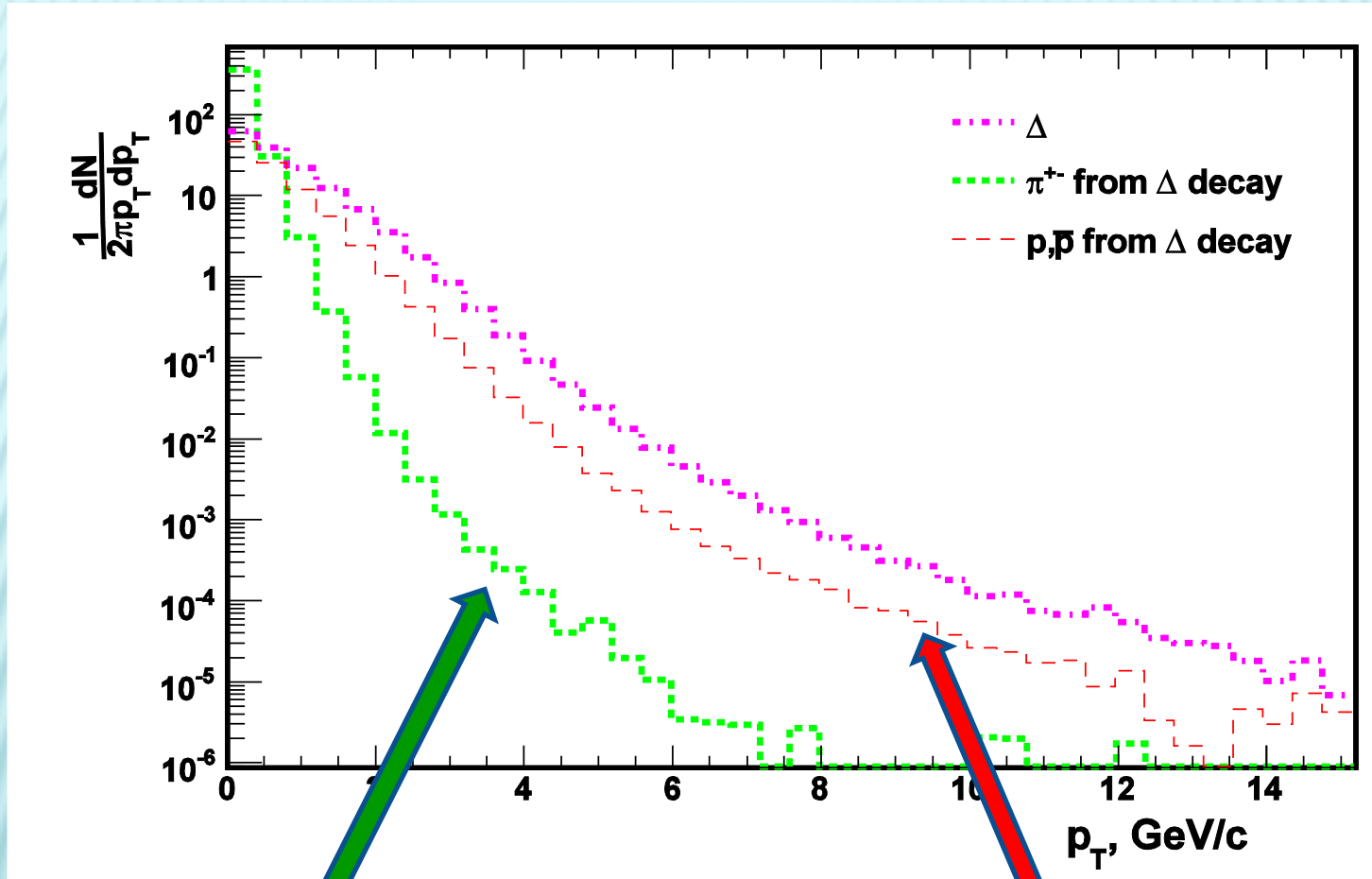


**Pions:** the resulting flow is weaker at low-pt and larger at high-pt

**Kaons:** both flows almost coincide

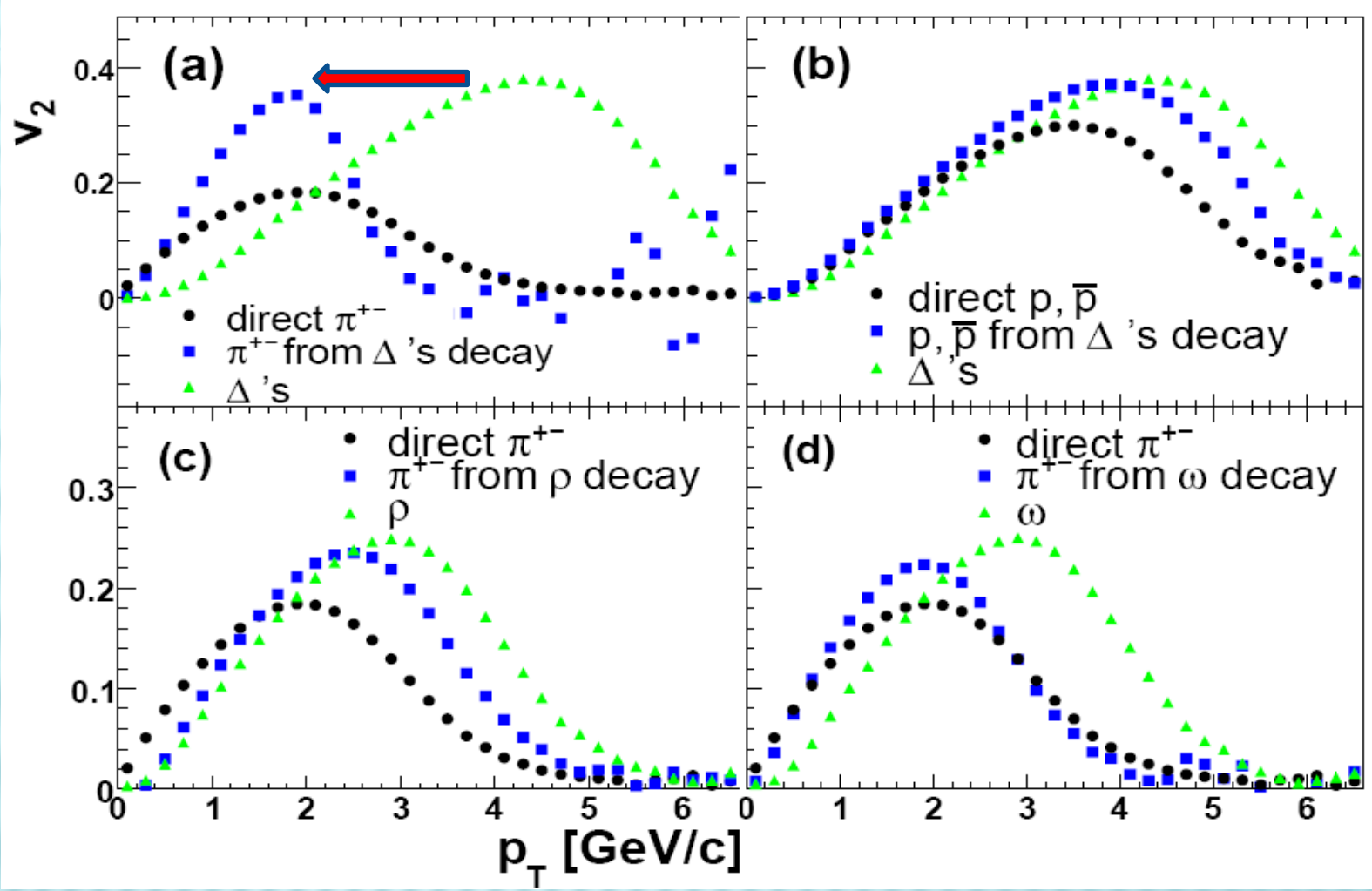
**Baryons:** the resulting flow is stronger than the flow of direct particles

# TRANSVERSE MOMENTUM OF SECONDARY PARTICLES



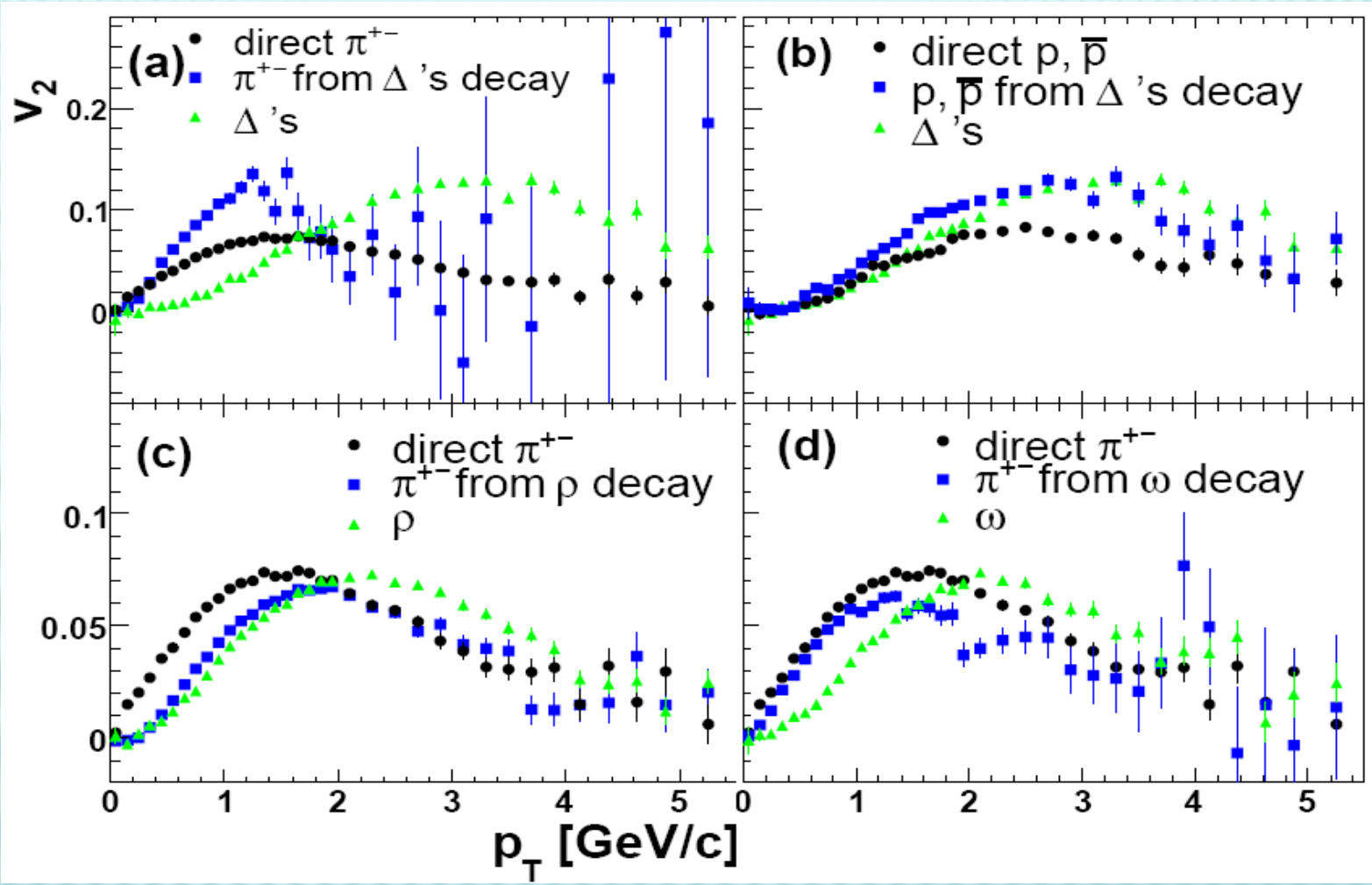
The secondary pion spectrum is much softer than proton spectrum

# ELLIPTIC FLOW OF DIRECT AND SECONDARY PARTICLES AT RHIC



The heavier resonances have larger  $v_2$  at high transverse momenta  
The decay kinematics keeps this high  $v_2$  for products of resonance decays

# ELLIPTIC FLOW OF DIRECT AND SECONDARY PARTICLES AT LHC



At low transverse momenta: pions from baryon resonances enhance the flow; pions from meson resonances reduce it

# **V . PARAMETERS OF THE MODEL**

## Model parameters.

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1. Thermodynamic parameters at chemical freeze-out:  $T_{ch}$  ,  $\{\mu_B, \mu_S, \mu_Q\}$
2. If thermal freeze-out is considered:  $T_{th}$  ,  $\mu\pi$ -normalisation constant
3. Volume parameters:  $T, \Delta T, R$
1.  $\rho_{\perp}^{\max}$  -maximal transverse flow rapidity for Bjorken-like parametrization
5.  $\eta_{\perp}^{\max}$  -maximal space-time longitudinal rapidity which determines the rapidity interval  $[-\eta_{\max}, \eta_{\max}]$  in the collision center-of-mass system.
6. Impact parameter range: minimal  $b_{\min}$  and maximal  $b_{\max}$  impact parameters
7. Flow anisotropy parameters  $\delta(b), \epsilon(b)$

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### PYTHIA+PYQUEN obligatory parameters

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9. Beam and target nuclear atomic weight  $A$
10.  $\sqrt{s_{NN}}$  -c.m.s. energy per nucleon pair (PYTHIA initialization at given energy)
11. **ptmin** – minimal pt of parton-parton scattering in PYTHIA event (ckin(3) in /pysubs/)
12. **nhsel** flag to include jet production in hydro-type event:

0 - jet production off (pure FASTMC event),  
1 - jet production on, jet quenching off (FASTMC+njet\*PYTHIA events),  
2 - jet production & jet quenching on (FASTMC+njet\*PYQUEN events),  
3 - jet production on, jet quenching off, FASTMC off (njet\*PYTHIA events),  
4 - jet production & jet quenching on, FASTMC off (njet\*PYQUEN events);

13. **ishad** flag to switch on/off nuclear shadowing

# Parameters of energy loss model in PYQUEN

(default, but can be changed from the default values by the user)

1. **T0** - initial temperature of quark-gluon plasma for central Pb+Pb collisions at mid-rapidity (initial temperature for other centralities and atomic numbers will be calculated automatically) at LHC: **T0=1 GeV**, at RHIC(200 AGeV) **T0=0.300 GeV**
2. **tau0** - proper time of quark-gluon plasma formation at LHC: **tau0=0.1 fm/c**, at RHIC(200 AGeV) **tau0=0.4 fm/c**
3. **nf** - number of active quark flavours in quark-gluon plasma (nf=0, 1, 2 or 3) at LHC: **nf=0**, at RHIC(200 AGeV) **nf=2**
4. **ienglu** - flag to fix type of medium-induced partonic energy loss (ienglu=0 - radiative and collisional loss, ienglu=1 - radiative loss only, ienglu=2 - collisional loss only, default value is ienglu=0);  
**ianglu** - flag to fix type of angular distribution of emitted gluons (ianglu=0 - small-angular, ianglu=1 - wide-angular, ianglu=2 - collinear, default value is ianglu=0).  
**ienglu=0**

# Methods for $v_2$ calculation

## (1) Event plane method

$$v_2^{obs} \{EP\} = \langle \cos 2(\varphi_i - \Psi_2) \rangle$$

$\Psi_2$  is the calculated reaction plane angle:  $\tan n \psi_n = \frac{\sum_i \omega_i \sin n \varphi_i}{\sum_i \omega_i \cos n \varphi_i}$ ,  $n \geq 1$ ,  $0 \leq \psi_n < 2\pi/n$

$$v_2 \{EP\} = \frac{v_2^{obs} \{EP\}}{R} = \frac{v_2^{obs} \{EP\}}{\langle \cos 2(\Psi_2 - \Psi_R) \rangle}$$

## (2) Two particle correlation method

$$v_2 \{2\} = \sqrt{\langle \cos 2(\varphi_i - \varphi_j) \rangle}$$

## (3) Lee-Yang zero method

$$G(ir) = \langle e^{irQ} \rangle, Q = \sum \cos(2\varphi)$$

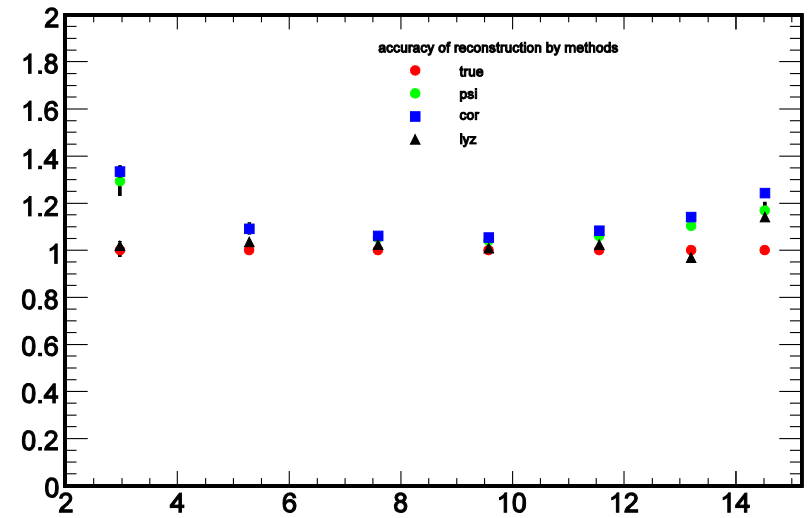
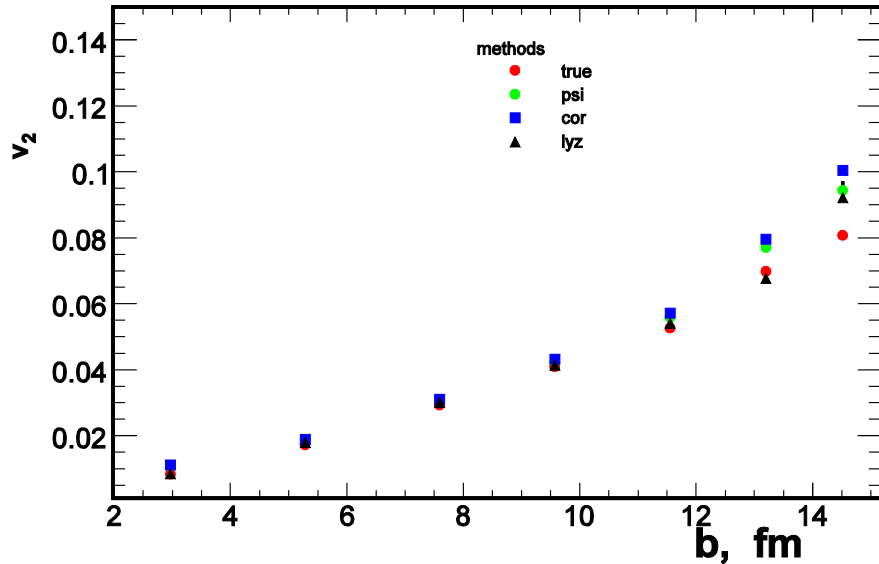
Integral  $v_2$  is connected with the first minimum  $r_0$  of the module of the  $G(ir)$ :

$$v_2 = \frac{j_0}{Nr_0}$$

Differential flow is calculated by the formula: 
$$\frac{v_2(p_T)}{Nv_2} = \text{Re} \left( \frac{\langle \cos(2\varphi) e^{ir_0 Q} \rangle}{\langle Q e^{ir_0 Q} \rangle} \right)$$



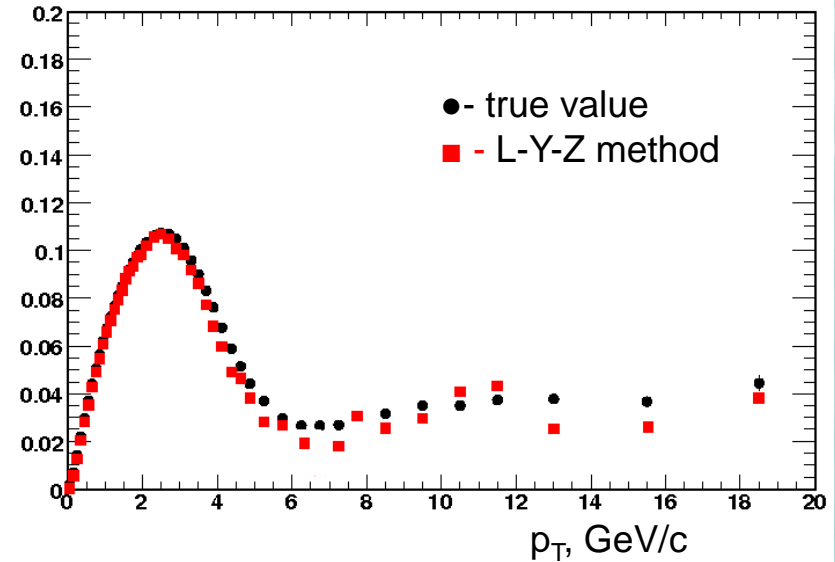
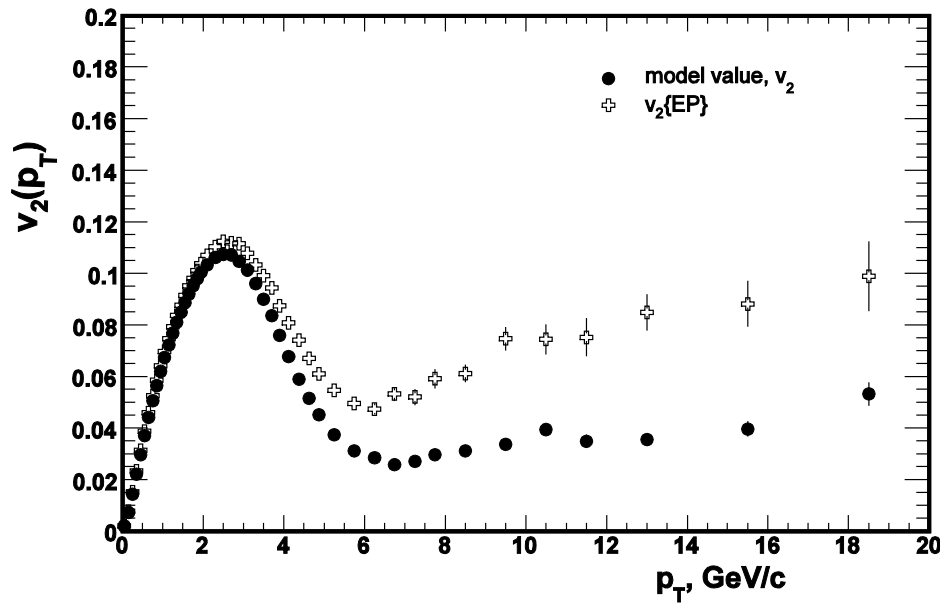
# RECONSTRUCTION OF INTEGRAL VALUE OF $V_2$ BY THE METHODS



The better reconstruction is achieved in midcentral collision for the methods, while Lee-Yang zero method tends to reconstruct true value at more central and more peripheral collision.

# Comparison of Event Plane and Lee-Yang zeroes methods ( $c=30\%$ )

## EventPlane method



## Lee-Yang zeroes Method

Event Plane method overestimates  $v_2$  at high  $p_t$  due to non-flow correlation (mostly because of jets).