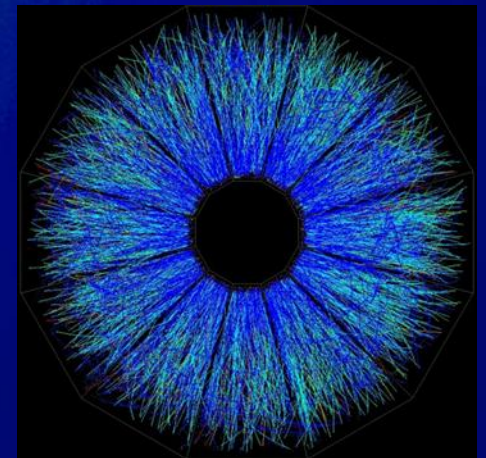


Open heavy flavor and quarkonia measurements in heavy ion collisions at RHIC



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FNSPE, Czech Technical University in Prague



International Conference on New Frontiers in Physics , Κρήτη, June 2012

Outline

- Motivation for **heavy flavor** physics
- Open heavy flavor
 - Charm mesons
 - Non-photonic electrons
- Quarkonia
 - J/ψ and Υ measurements
- Summary and Outlook

Relativistic Heavy Ion Collider

RHIC site in BNL on Long Island - taking data from 2000



RHIC has been exploring nuclear matter at extreme conditions over the last years

Lattice QCD predicts a phase transition from hadronic matter to a deconfined state, the **Quark-Gluon Plasma**

Colliding systems:

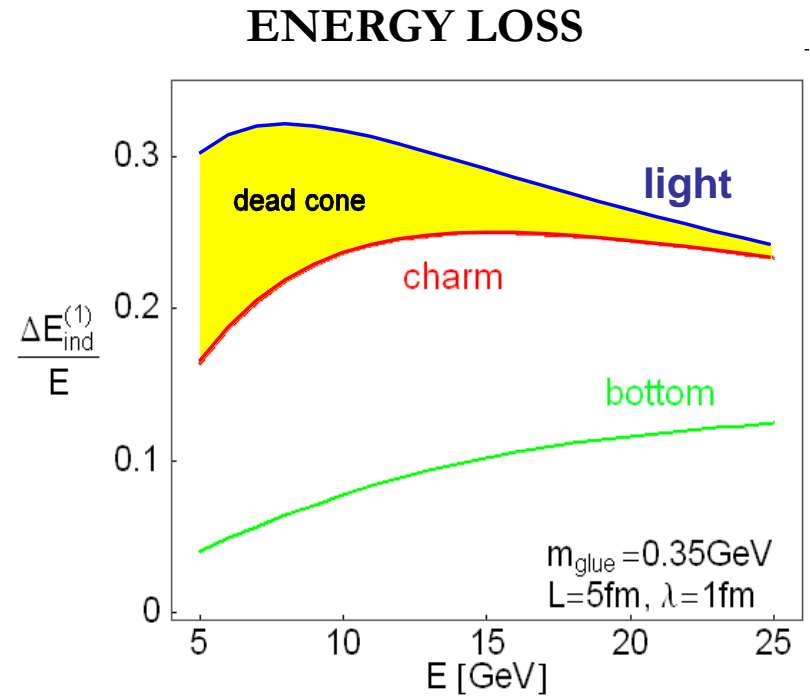
$p\uparrow+p\uparrow$, $d+Au$, $Cu+Cu$, $Au+Au$
 $Cu+Au$, $U+U$

Energies

$\sqrt{s_{NN}} = 20, 62, 130, 200 \text{ GeV (500 GeV)}$
 $+ 7.7, 11.5, 27, 39 \text{ GeV}$

Heavy quarks as a probe of QGP

- **p+p data:**
 - baseline of heavy ion measurements.
 - test of pQCD calculations.
- Due to their **large mass** heavy quarks are primarily **produced** by **gluon fusion** in early stage of collision.
 - production rates calculable by pQCD.
M. Gyulassy and Z. Lin, PRC 51, 2177 (1995)
- **heavy ion data:**
 - Studying **energy loss** of heavy quarks.
 - independent way to **extract properties** of the **medium**.



M.Djordjevic PRL 94 (2004)

Quarkonia states in A+A

Charmonia: J/ψ , Ψ' , χ_c

Bottomonia: $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$

Key Idea: Quarkonia melt in the QG plasma due to color screening of potential between heavy quarks

- Suppression of states is determined by T_c and their binding energy
- Lattice QCD: Evaluation of spectral functions $\Rightarrow T_{\text{melting}}$

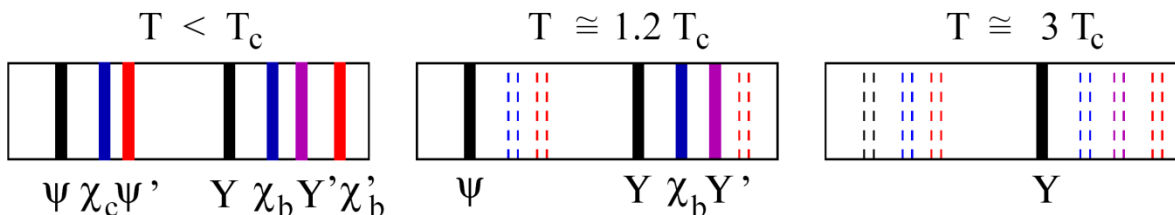
Sequential disappearance of states:

\Rightarrow Color screening \Rightarrow Deconfinement

\Rightarrow QCD thermometer \Rightarrow Properties of QGP

When do states really melt?

$$T_{\text{diss}}(\Psi') \approx T_{\text{diss}}(\chi_c) < T_{\text{diss}}(\Upsilon(3S)) < T_{\text{diss}}(J/\psi) \approx T_{\text{diss}}(\Upsilon(2S)) < T_{\text{diss}}(\Upsilon(1S))$$



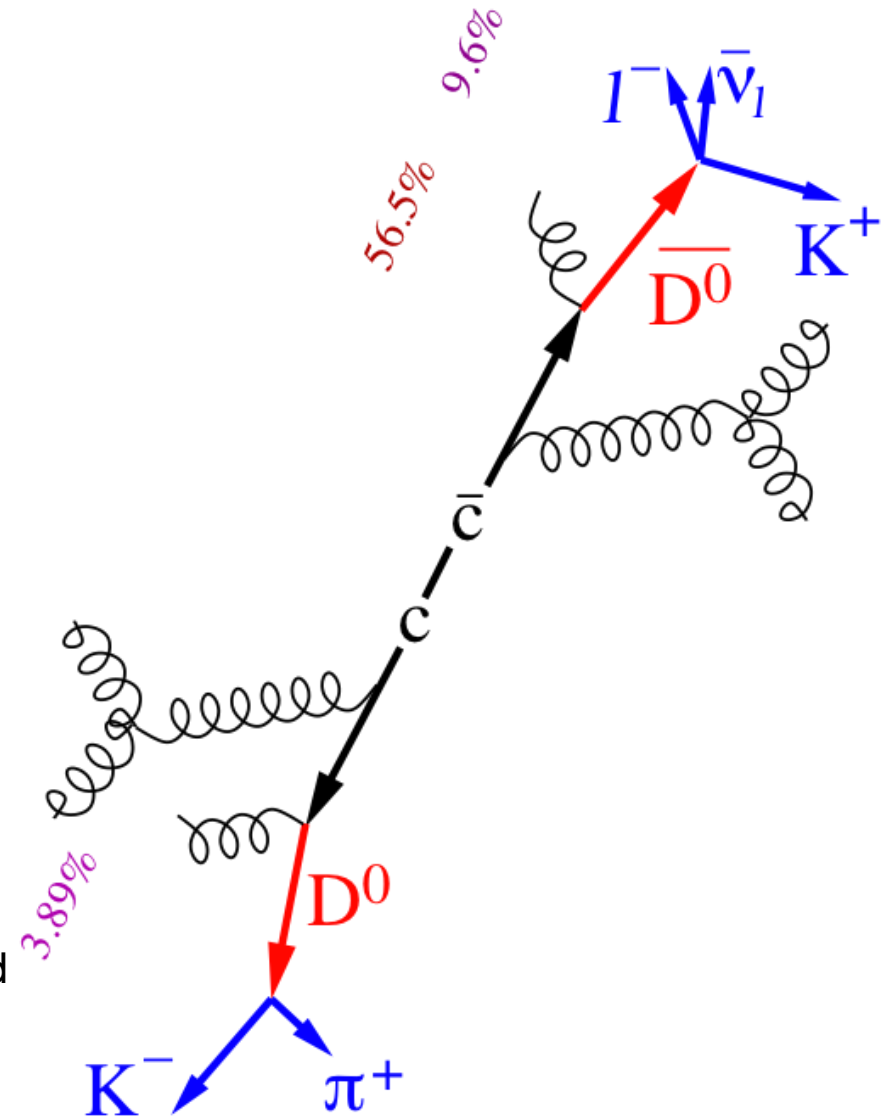
Open heavy flavor

- **Direct reconstruction**

- direct access to heavy quark kinematics
- hard to trigger (high energy trigger only for correlation measurements)
- smaller Branching Ratio (B.R.)
- large combinatorial background (need handle on decay vertex)

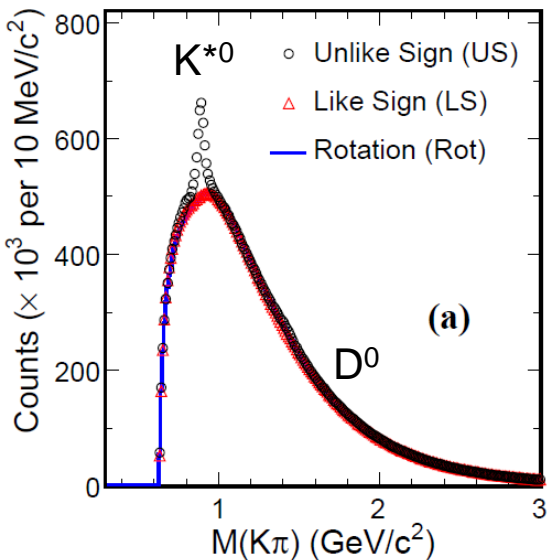
- **Indirect measurements through decay leptons**

- can be triggered easily (high p_T)
- higher B.R.
- indirect access to the heavy quark kinematics
- mixing contribution from all charm and bottom hadron decays

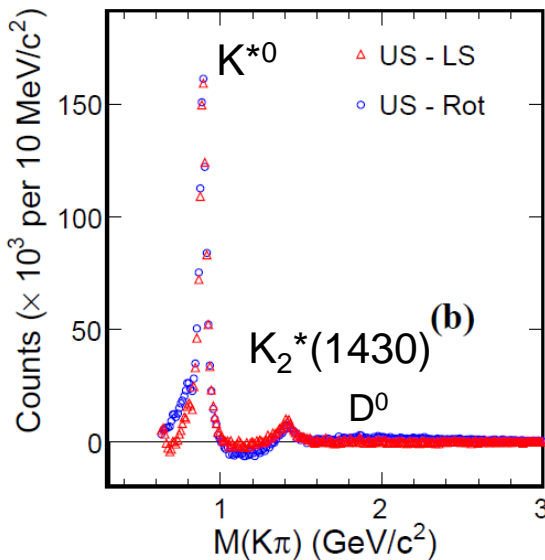


D⁰ and D^{*} signal in p+p and Au+Au 200 GeV

STAR arXiv: 1204.4244

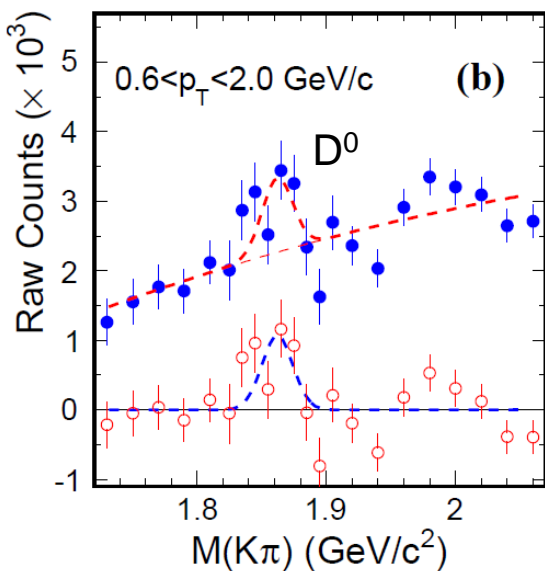
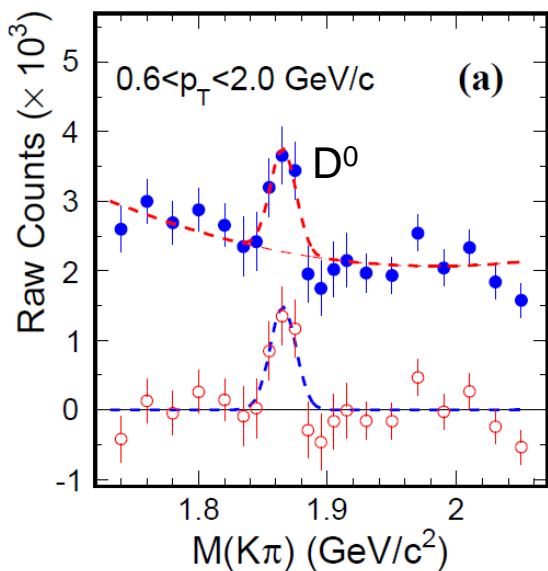
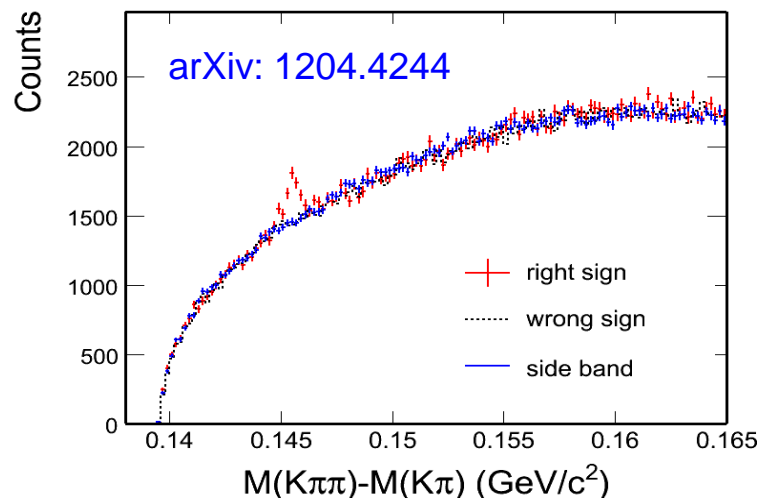


D⁰ mass window zoom



D⁰ mass window zoom

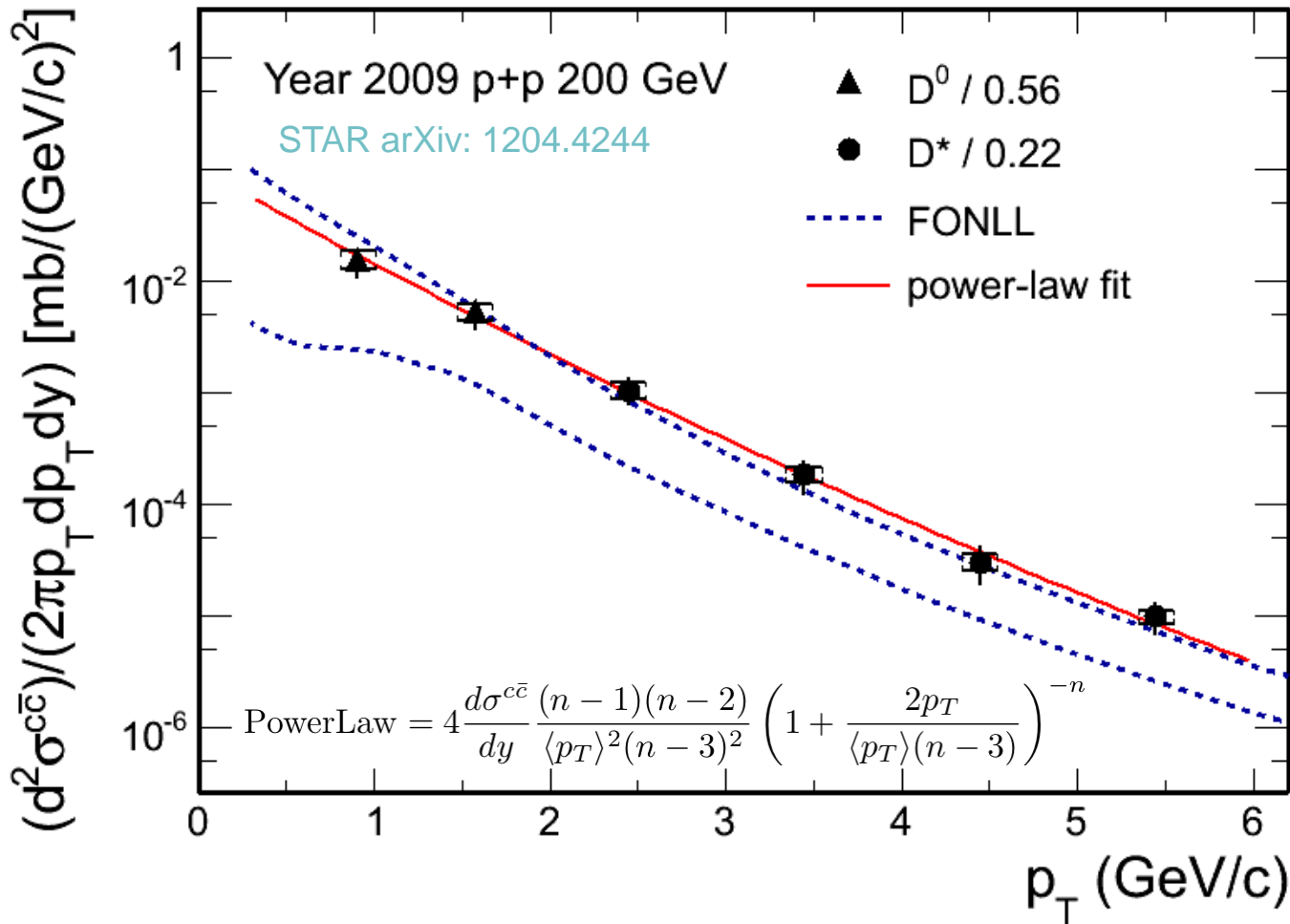
$$D^{*\pm} \rightarrow D^0(\overline{D}^0) + \pi^\pm \rightarrow K^\mp \pi^\pm \pi^\pm$$



Different methods reproduce combinatorial background.

Consistent results from two background methods.

D⁰ and D* p_T spectra in p+p 200 GeV



D⁰ scaled by $N_{D^0}/N_{cc} = 0.56^{[1]}$

D* scaled by $N_{D^*}/N_{cc} = 0.22^{[1]}$

Consistent with FONLL^[2]
upper limit.

$X_{sec} = dN/dy|_{y=0}^{cc} * F * \sigma_{pp}$

$F = 4.7 \pm 0.7$ scale to full
rapidity.

$\sigma_{pp}(\text{NSD}) = 30 \text{ mb}$

The charm cross section
at mid-rapidity is:

$170 \pm 45(\text{stat.})_{-51}^{+37}(\text{sys.}) \mu\text{b}$

The charm total cross
section is extracted as:

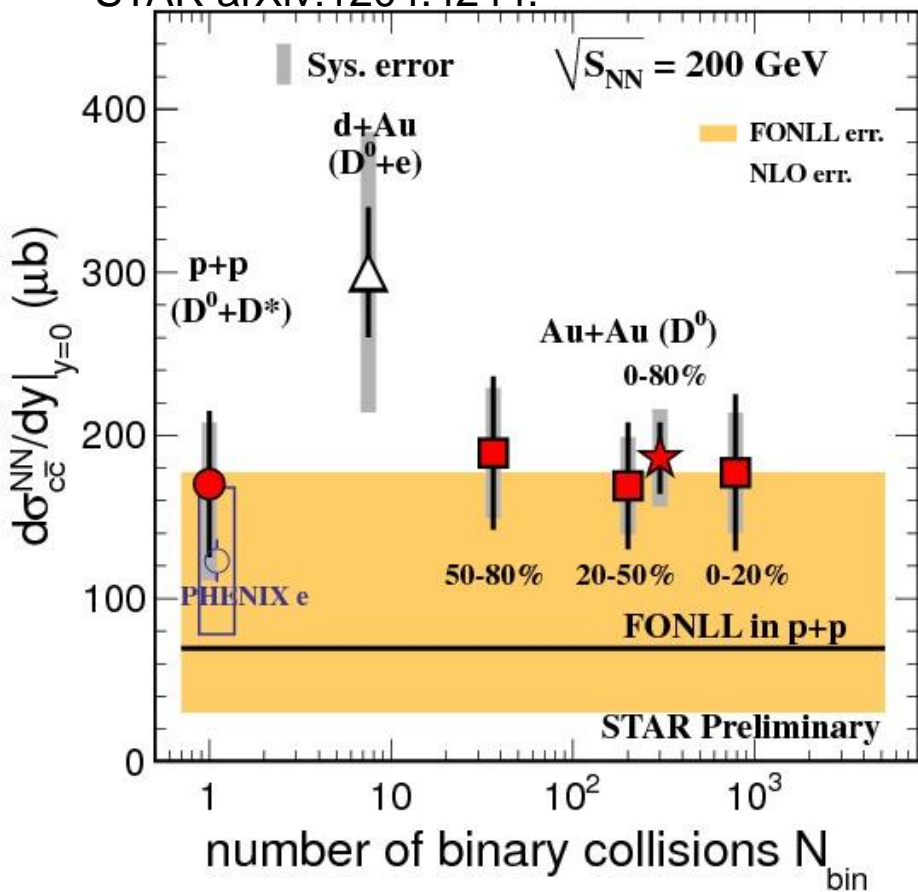
$797 \pm 210(\text{stat.})_{-262}^{+208}(\text{sys.}) \mu\text{b}$

[1] C. Amsler et al. (Particle Data Group), PLB 667 (2008) 1.

[2] Fixed-Order Next-to-Leading Logarithm: M. Cacciari, PRL 95 (2005) 122001.

Charm cross section vs N_{bin}

YiFei Zhang, JPG 38, 124142 (2011)
 STAR arXiv:1204.4244.



All of the measurements are consistent.

Year 2003 d+Au : $D^0 + e$

Year 2009 p+p : $D^0 + D^*$

Year 2010 Au+Au: D^0

Assuming $N_{D^0} / N_{cc} = 0.56$ does not change.

Charm cross section in Au+Au 200 GeV:

Mid-rapidity:

186 ± 22 (stat.) ± 30 (sys.) ± 18 (norm.) μb

Total cross section:

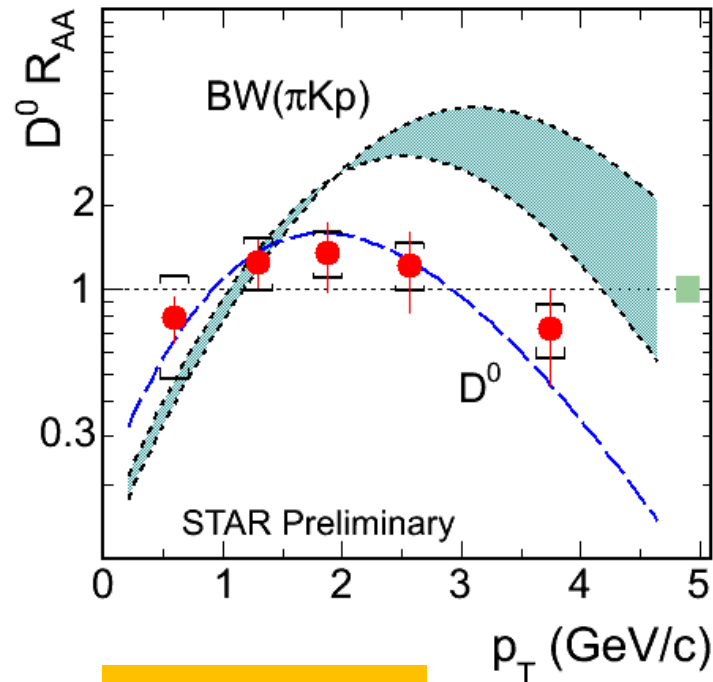
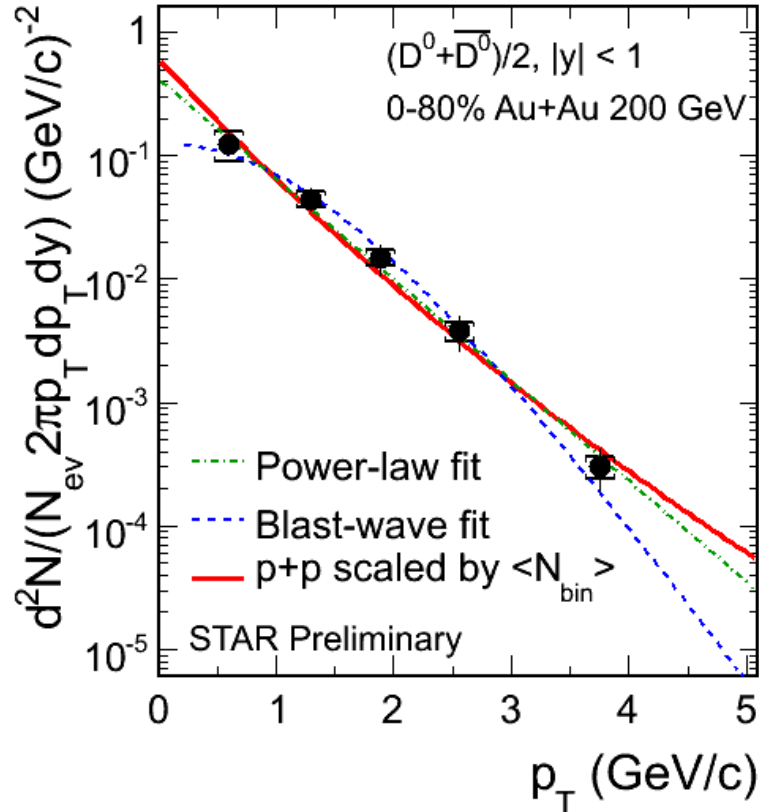
876 ± 103 (stat.) ± 211 (sys.) μb

- [1] STAR d+Au: J. Adams, et al., PRL 94 (2005) 62301
- [2] FONLL: M. Cacciari, PRL 95 (2005) 122001.
- [3] NLO: R. Vogt, Eur.Phys.J.ST 155 (2008) 213
- [4] PHENIX e: A. Adare, et al., PRL 97 (2006) 252002.

Charm cross section follows number of binary collisions scaling =>
Charm quarks are mostly produced via initial hard scatterings.

D⁰ nuclear modification factor Au+Au 200 GeV

BW ($\pi K\rho$): B. I. Abelev, et al., Phys. Rev. C 79 (2009) 34909.



STAR QM2011

Au+Au 200 GeV:

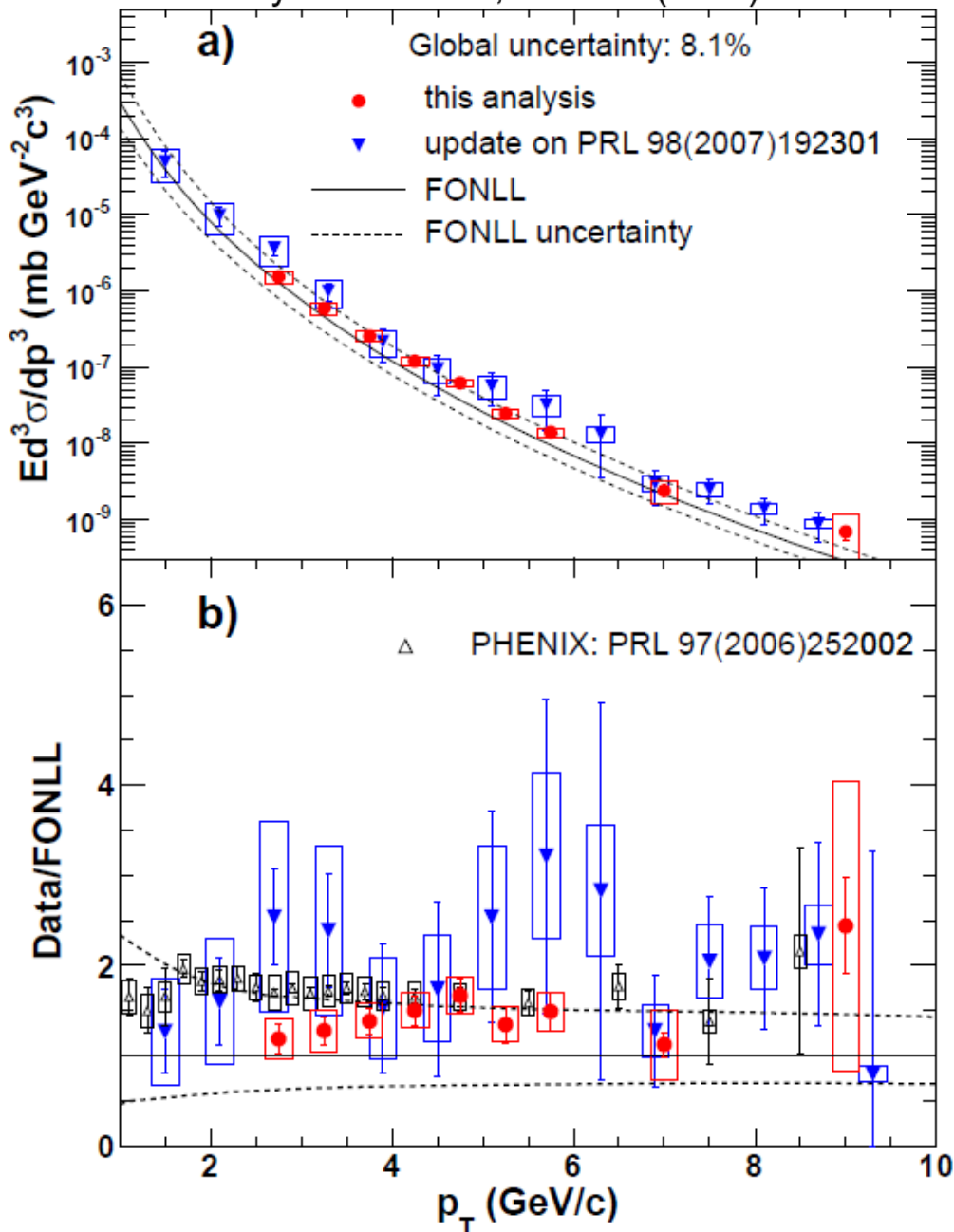
- No obvious suppression at $p_T < 3$ GeV/c.

$$R_{AA}(p_T) = \frac{\text{Yield}_{AA}(p_T)}{\langle N_{bin} \rangle_{AA} \text{Yield}_{pp}(p_T)}$$

Non-photonic electrons

Non-photonic electrons in p+p 200GeV

STAR Phys. Rev. D **83**, 052006 (2011)



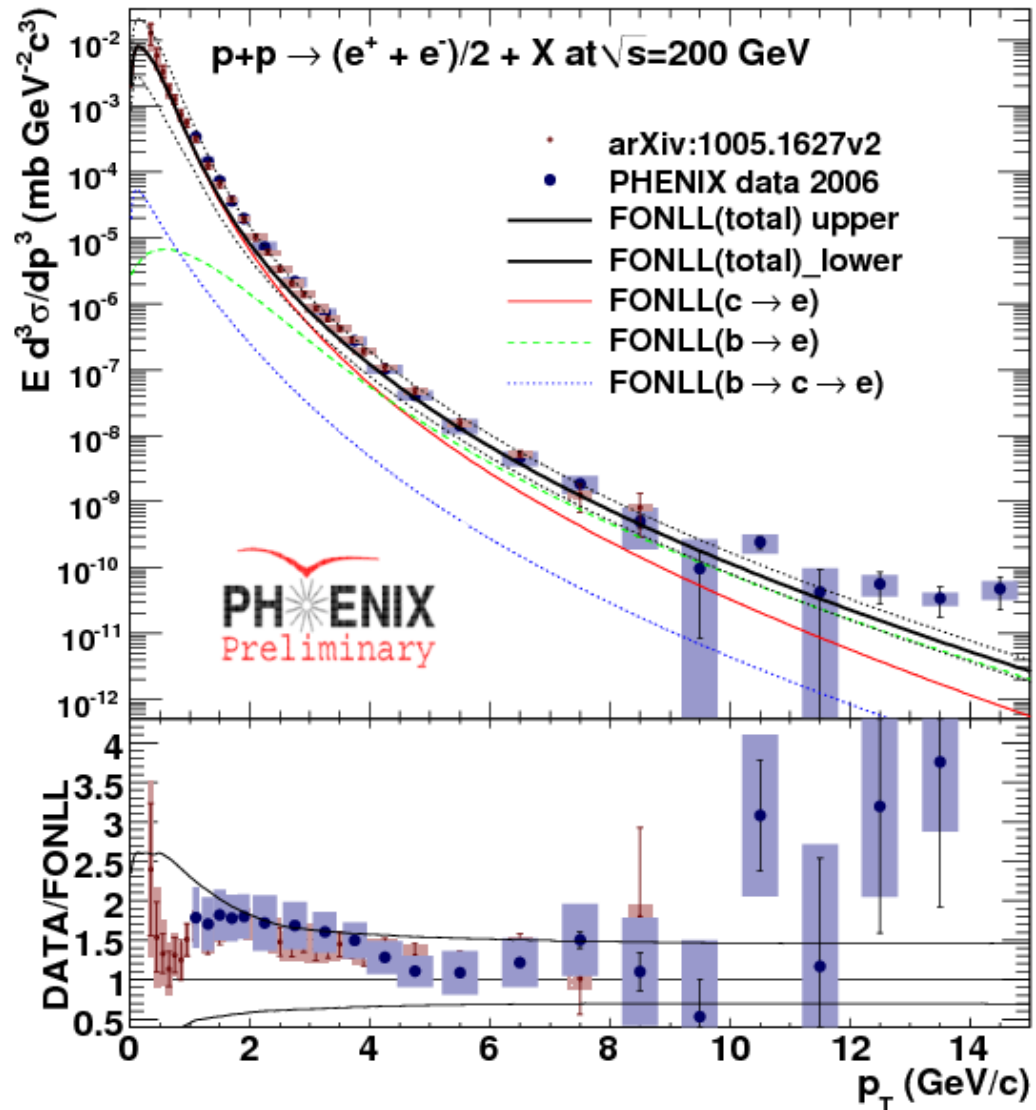
STAR and PHENIX NPE results in p+p 200GeV collisions

✓ Are consistent within errors at $p_T > 2.5$ GeV/c

NPE results are consistent with FONLL in p+p 200GeV collisions

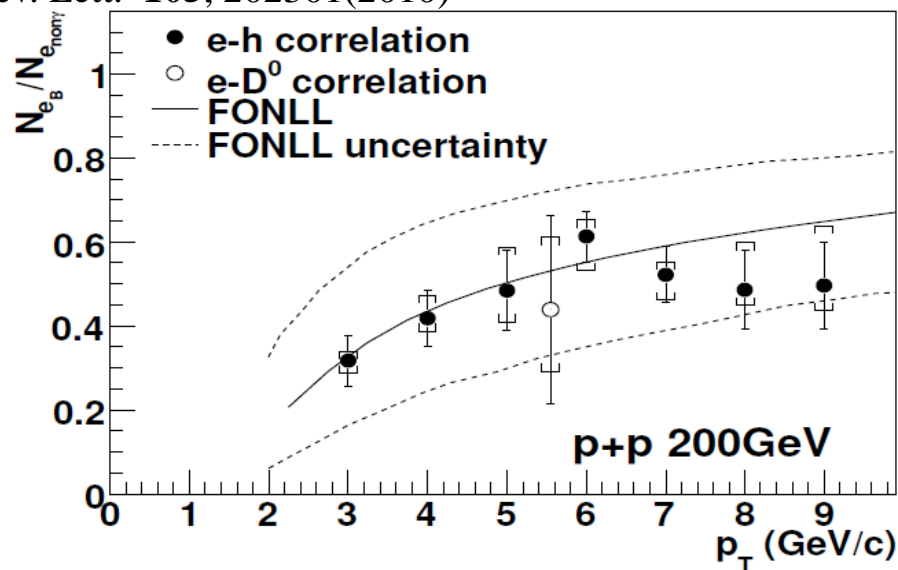
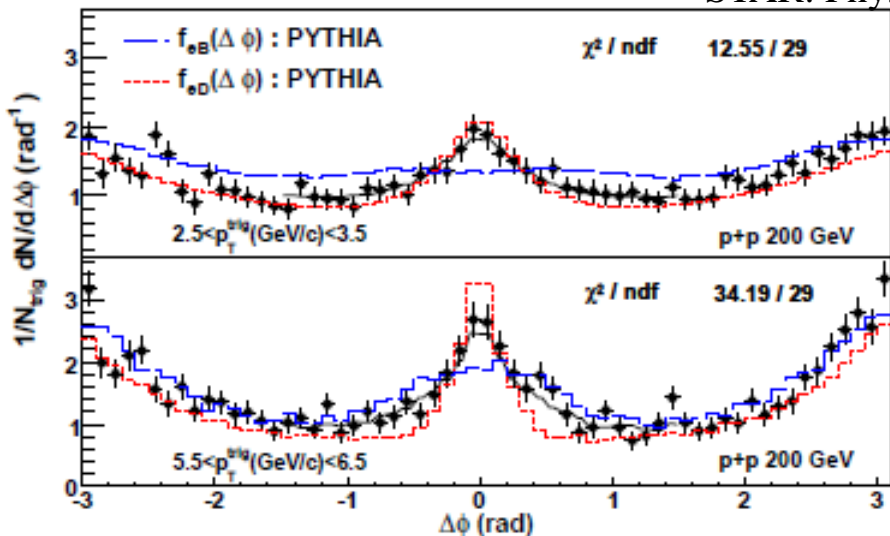
PHENIX update p+p 200 GeV

- Combined Run5 and Run6 p+p statistics
- Smaller uncertainties
 - Allows more precise R_{AA} comparisons
- Increased p_T range
- Consistent with previous results in overlap region

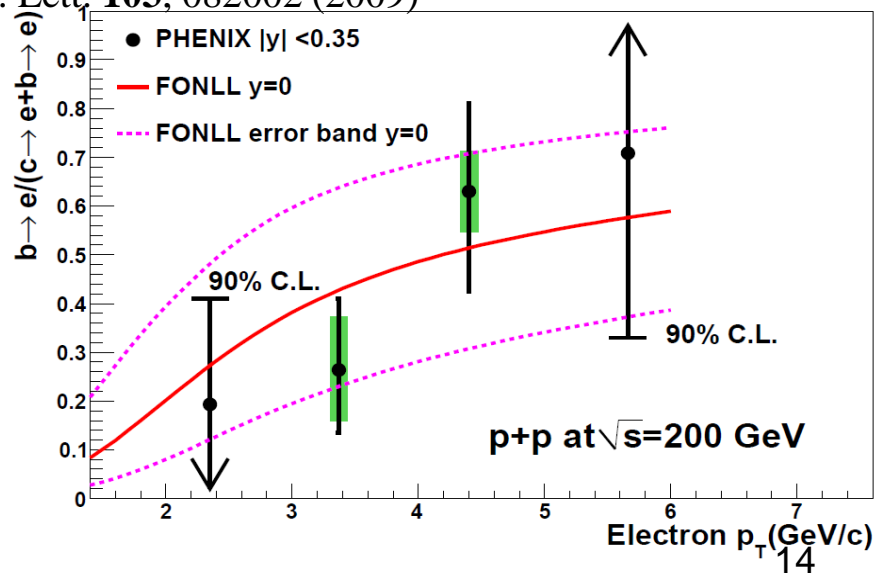
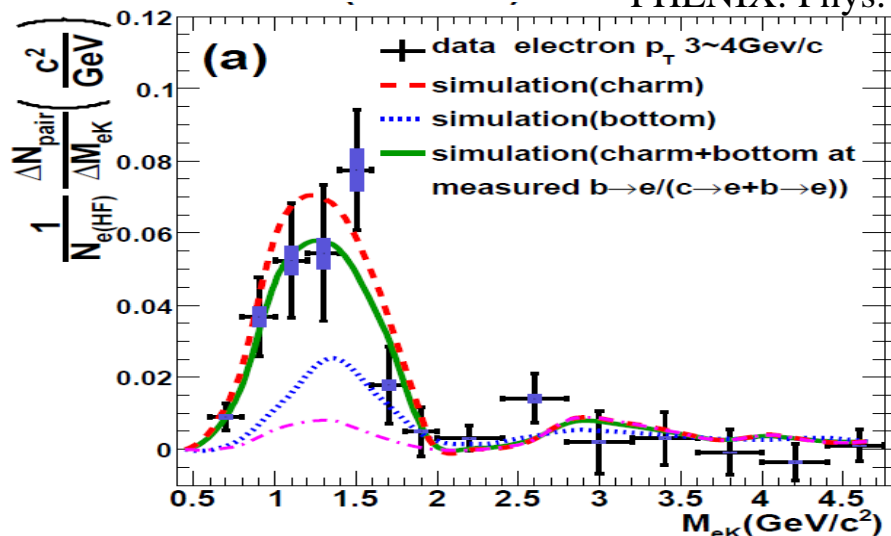


Separation of charm and bottom contribution

STAR: Phys. Rev. Lett. **105**, 202301(2010)



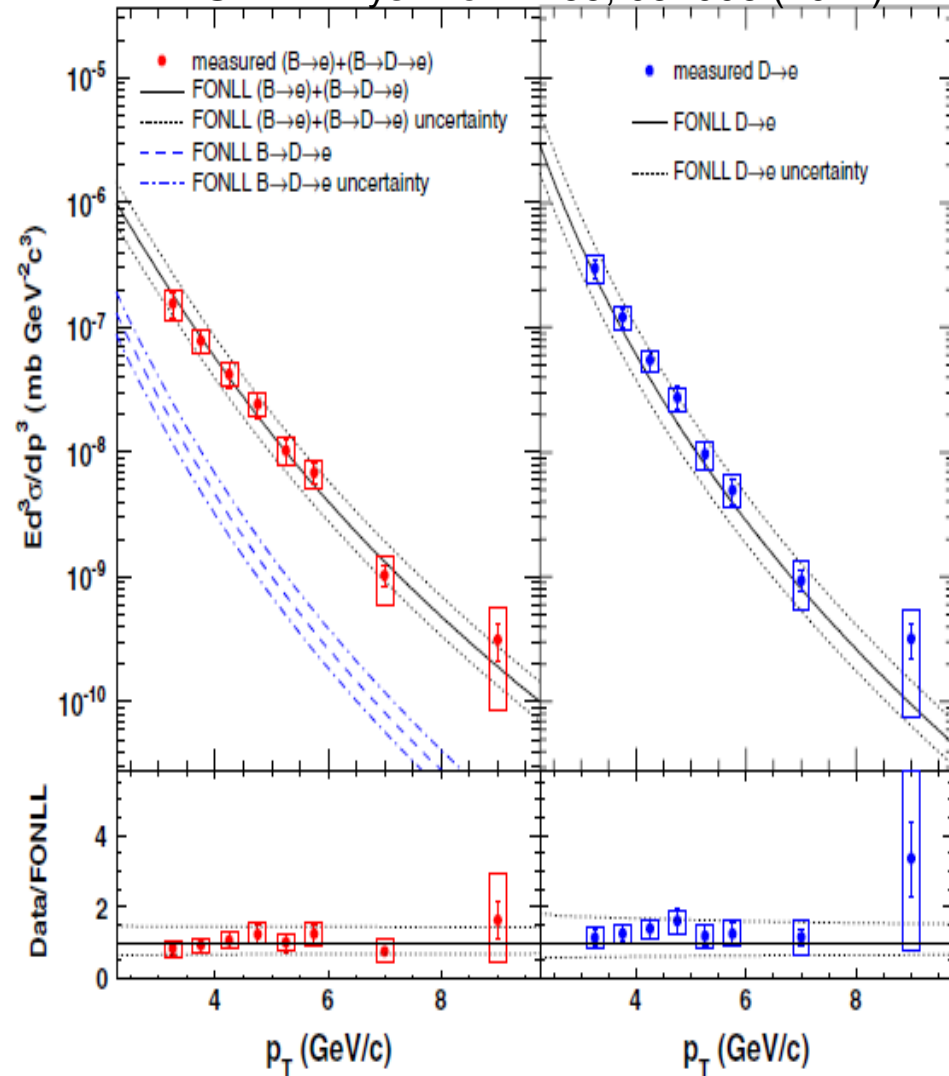
PHENIX: Phys. Rev. Lett. **103**, 082002 (2009)



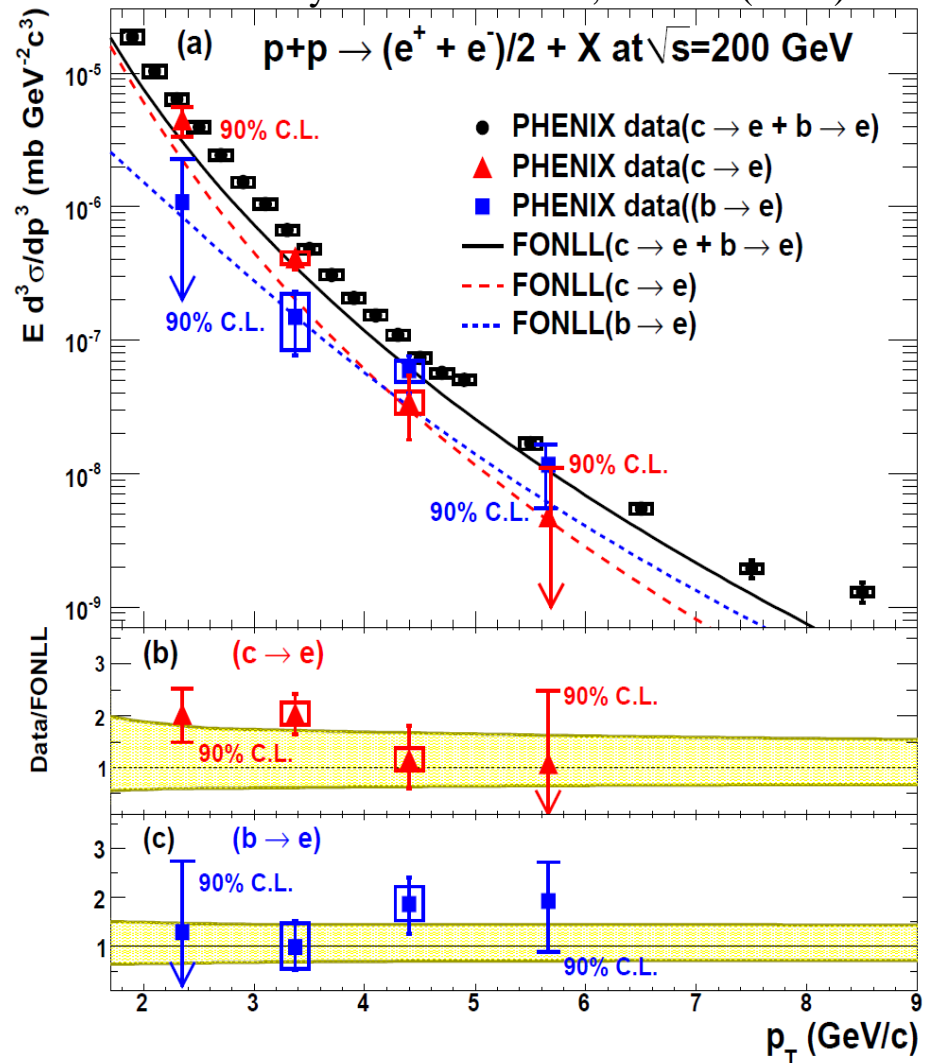
- ~30-60% of non-photonic electrons come from B meson in 200GeV p+p collisions.

Measurement of B and D meson spectra at RHIC

STAR Phys. Rev. D **83**, 052006 (2011)

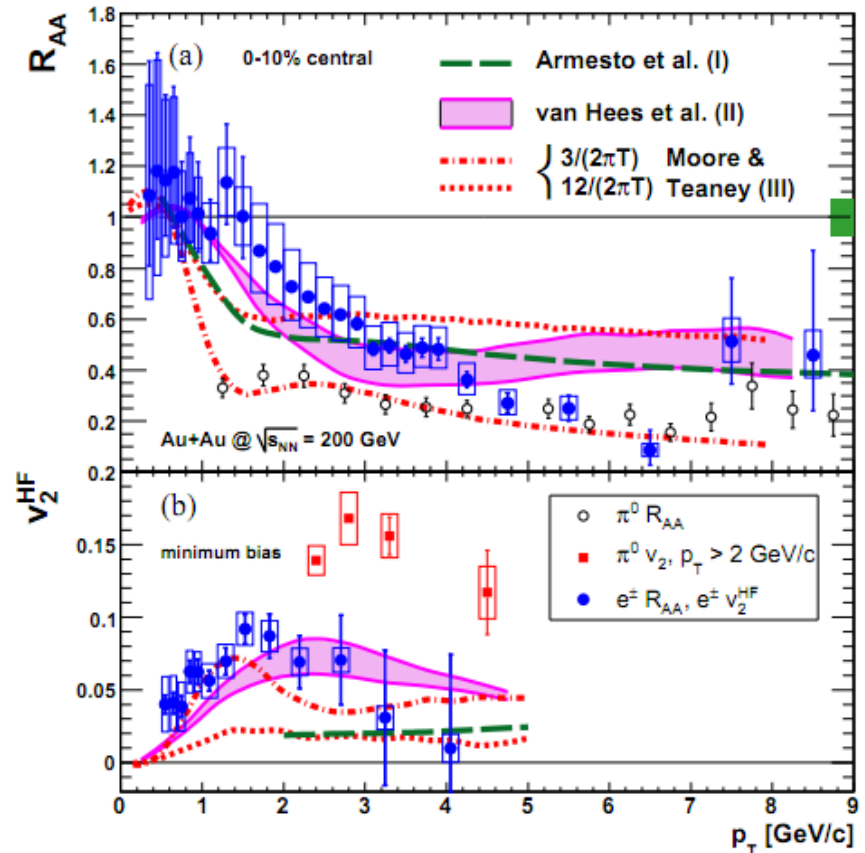
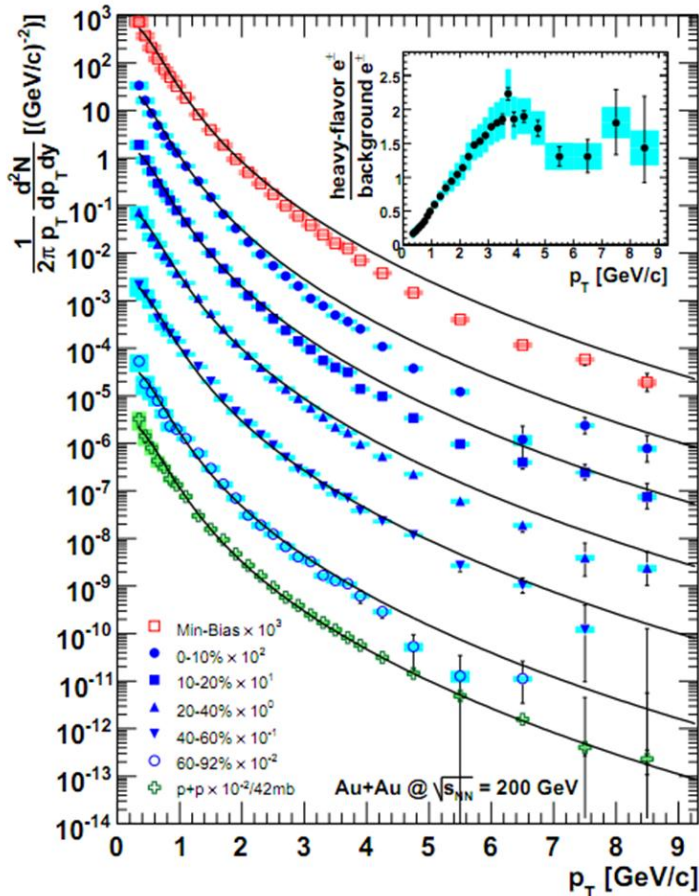


PHENIX Phys. Rev. Lett. **103**, 082002 (2009)



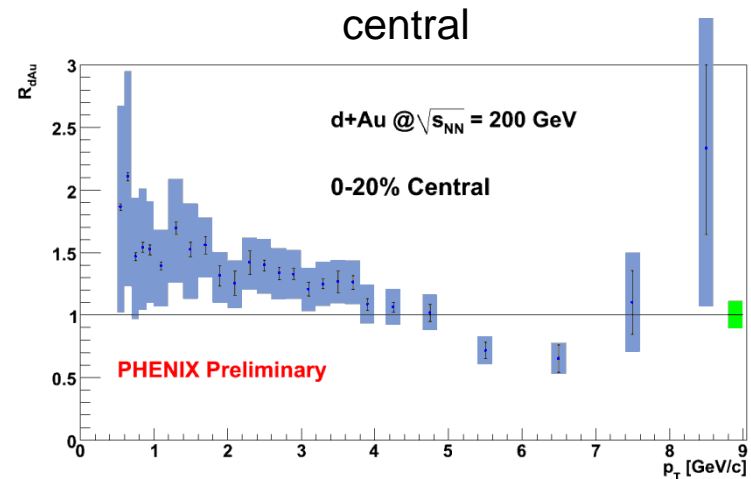
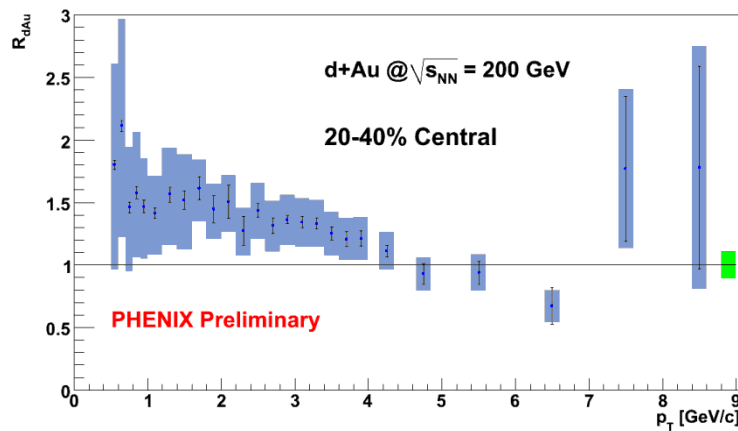
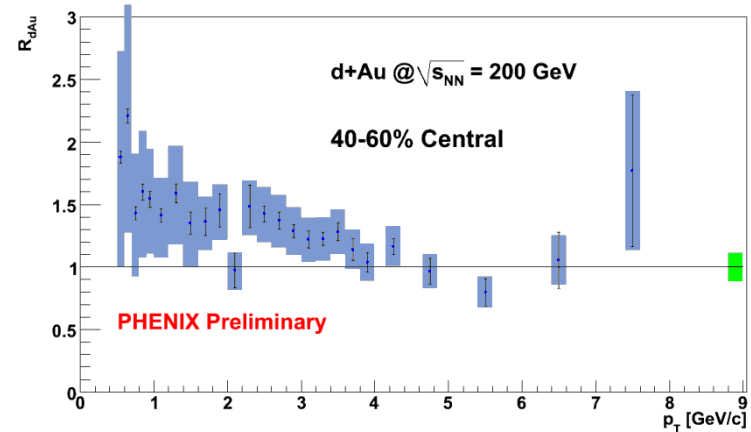
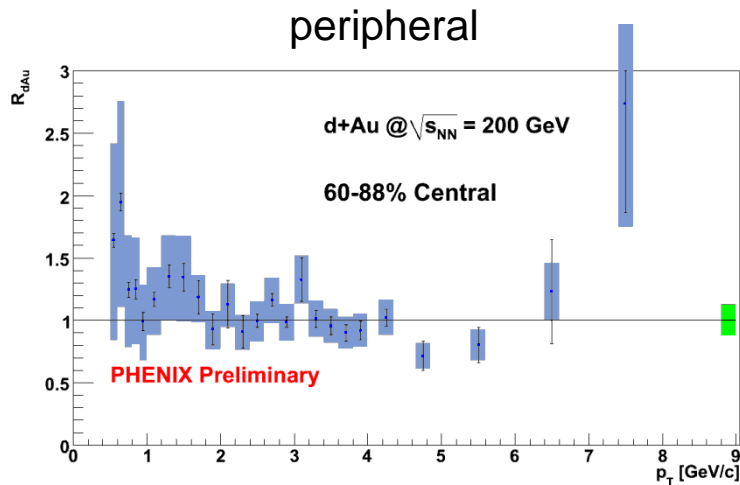
- p+p 200 GeV result
- ideal to get similar result from Au+Au but it's a lot harder.

Nuclear modification factor in Au+Au 200 GeV



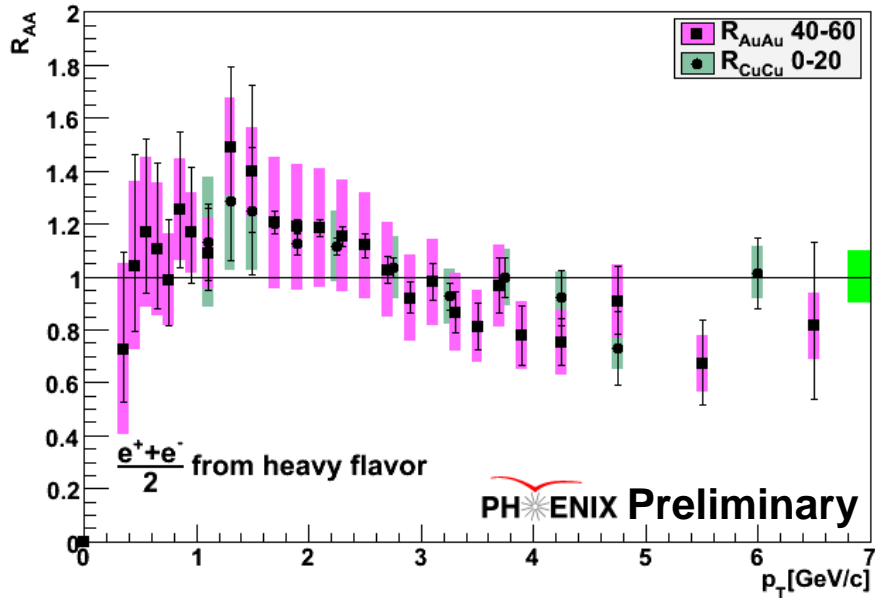
- Non-photonic electrons suppressed at high- p_T
- Flow of NPE was measured

Non-photonic electrons in d+Au 200 GeV



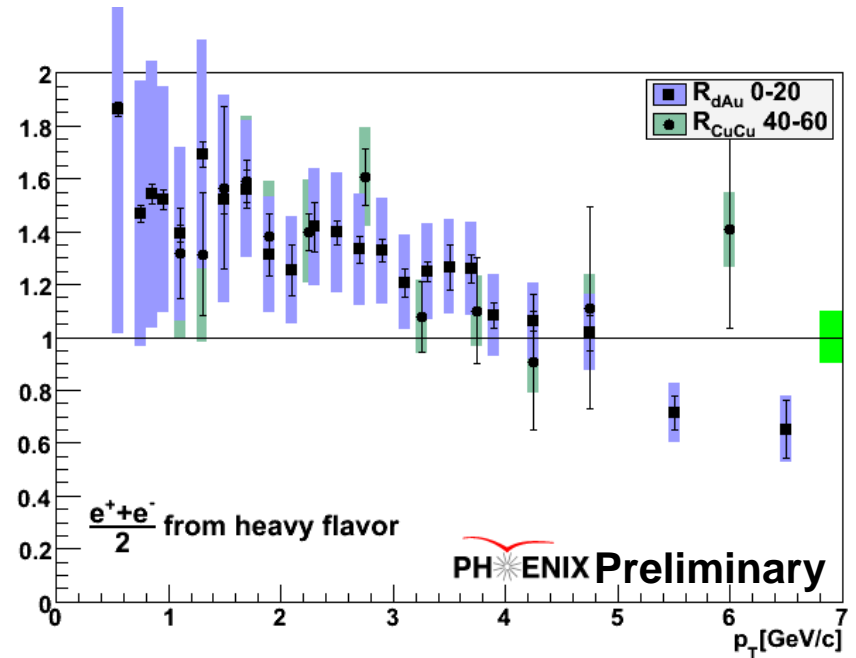
- Peripheral R_{dAu} consistent with 1.0
- Evidence of CNM effects on open HF yields at $1 < p_T < 4$ GeV/c for more central collisions

d+Au/Au+Au and Cu+Cu



$$\langle N_{coll} \rangle \text{ CuCu} = 150$$

$$\langle N_{coll} \rangle \text{ AuAu} = 91$$



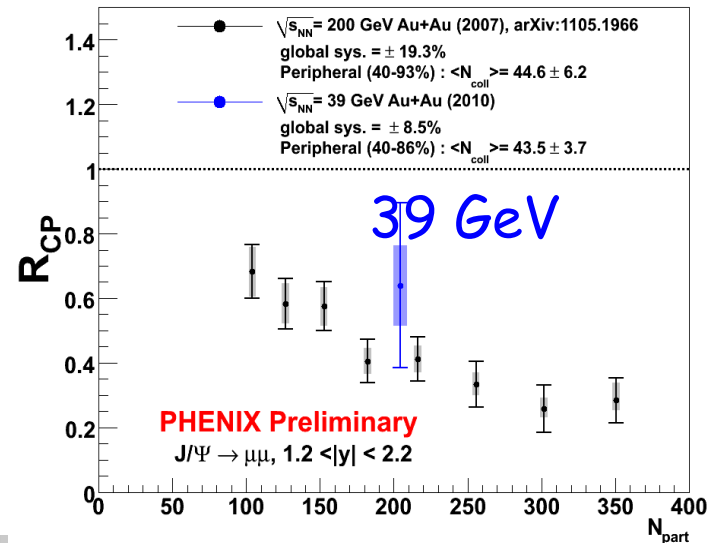
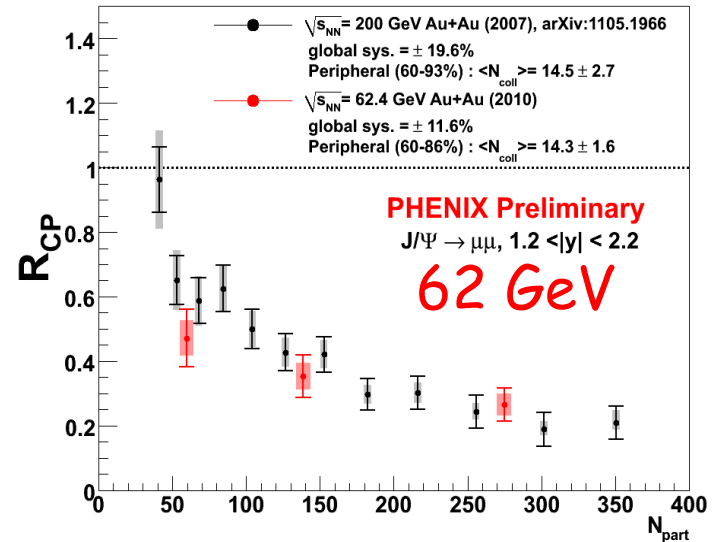
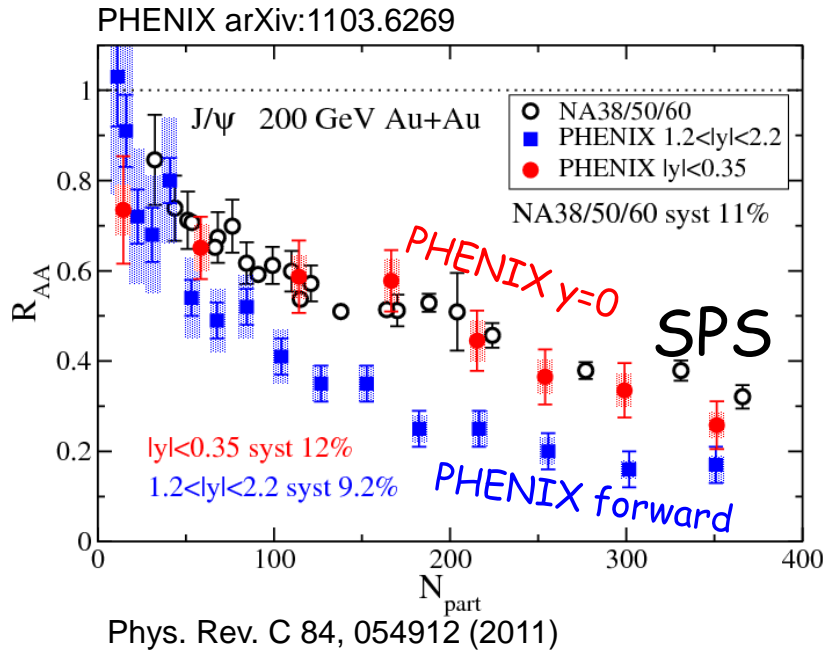
$$\langle N_{coll} \rangle \text{ dAu} = 15$$

$$\langle N_{coll} \rangle \text{ CuCu} = 22.3$$

QUARKONIA

J/ψ

Charmonium suppression

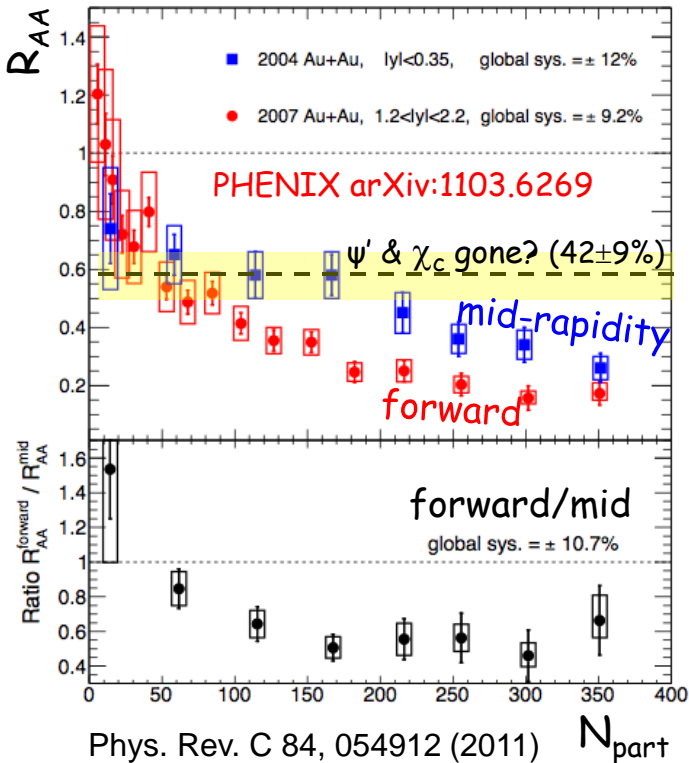


- Overall suppression of J/ψ is nearly identical between RHIC (200,62,39 GeV), SPS (17.2 GeV), (& LHC)
- Forward-rapidity is suppressed more than Mid-rapidity

Forward rapidity suppression

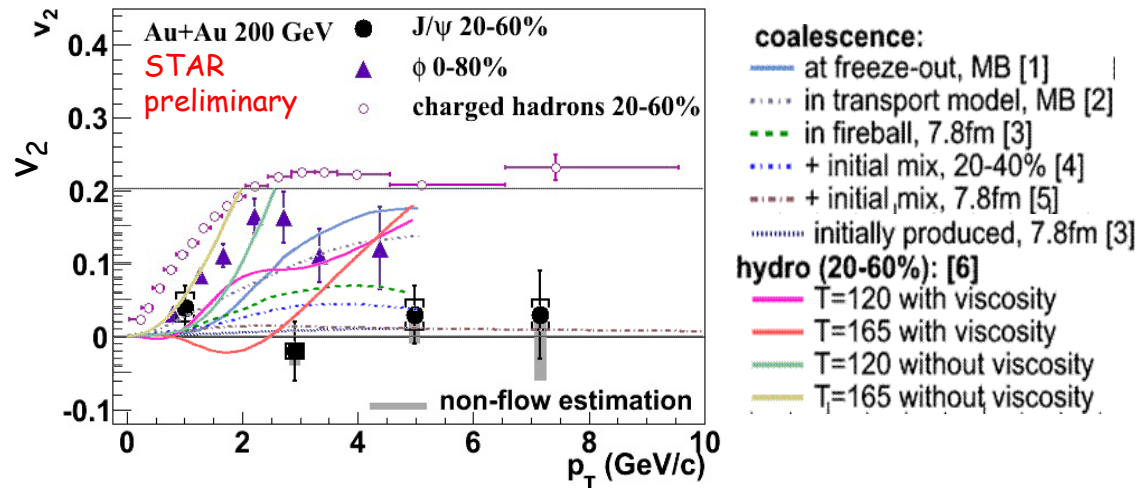
At RHIC - Forward-rapidity J/ψ's are suppressed more than Mid-rapidity – Why?

- 1) Stronger forward rapidity suppression due to CNM effects?
- 2) Regeneration at mid-rapidity reduces suppression relative to forward (and gives net suppression similar to SPS)?



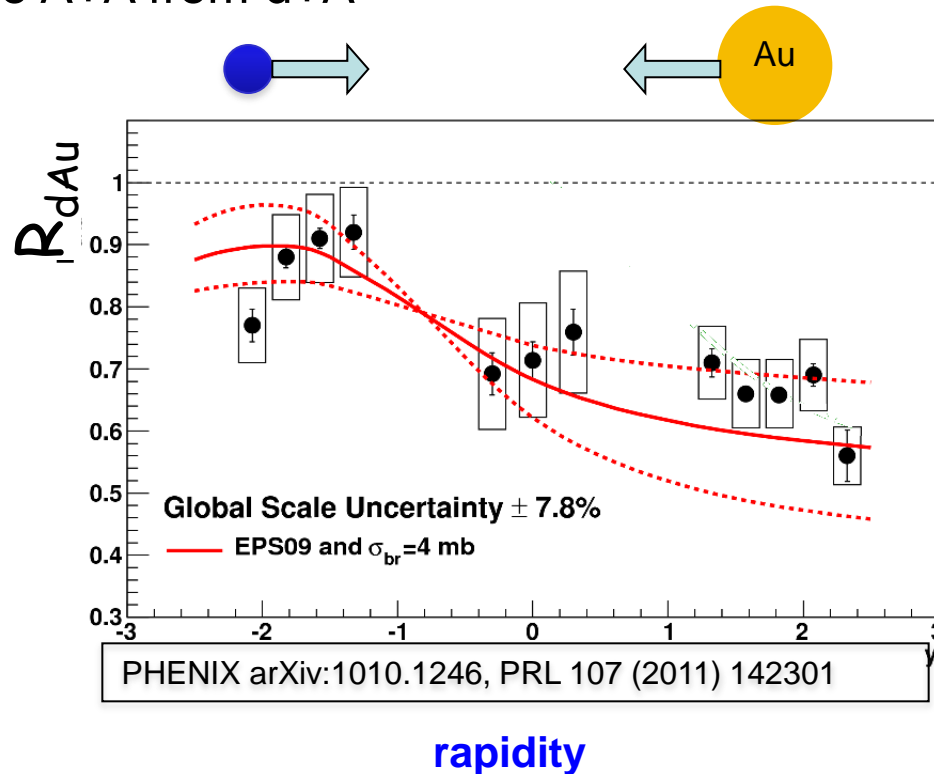
Small or no J/ψ flow at RHIC!

Many theoretical expectations & option #2 probably ruled out?



Cold nuclear matter effects – d+Au 200 GeV

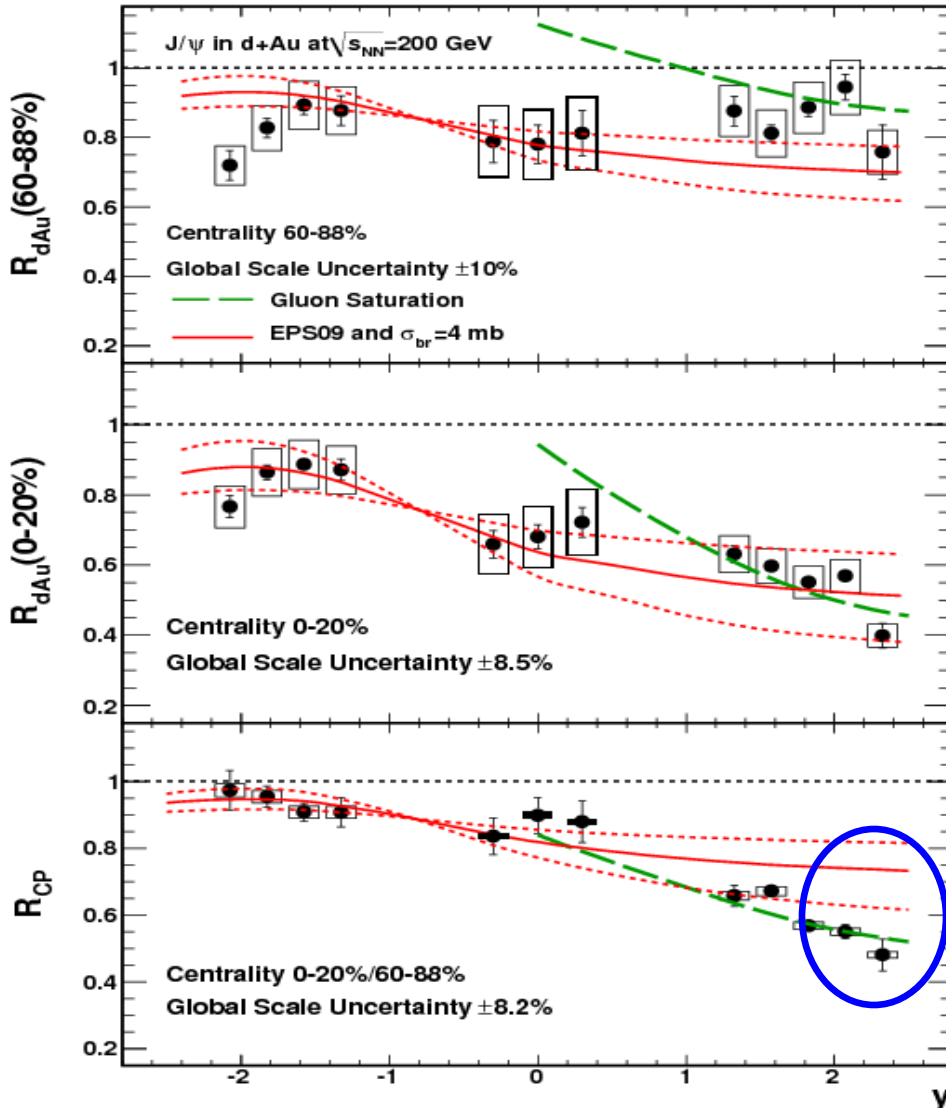
CNM effects appear to provide a large fraction of the observed suppression; so difficult to conclude much w/o a thorough understanding of CNM and its extrapolation to A+A from d+Au



- we have to understand CNM in a fundamental way in order to obtain reliable/quantitative extrapolations to A+A.

Cold nuclear matter effects – d+Au 200 GeV

centrality



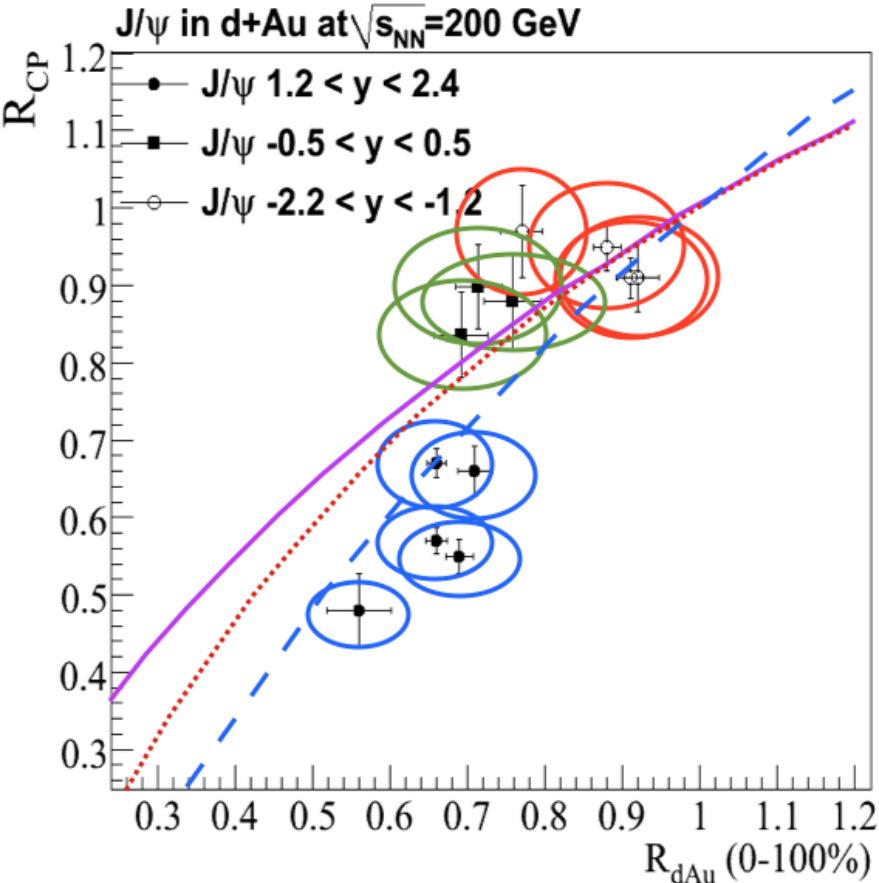
Reasonable agreement with **EPS09 nPDF + $\sigma_{br}=4$ mb** for central collisions but not peripheral

CGC calculations can't reproduce mid-rapidity (*Nucl. Phys. A 770(2006) 40*)

EPS09 with linear thickness dependence fails to describe centrality dependence of forward rapidity region.

Cold nuclear matter effects – d+Au 200 GeV

path-length dependence



PHENIX arXiv:1010.1246, PRL 107 (2011) 142301

- Assume modification is dependent on the nuclear thickness

$$\Lambda(r_T) = \frac{1}{\rho_0} \int dz \rho(z, r_T)$$

Woods-Saxon

Exponential: $M(r_T) = e^{-a\Lambda(r_T)}$

Linear: $M(r_T) = 1 - a\Lambda(r_T)$

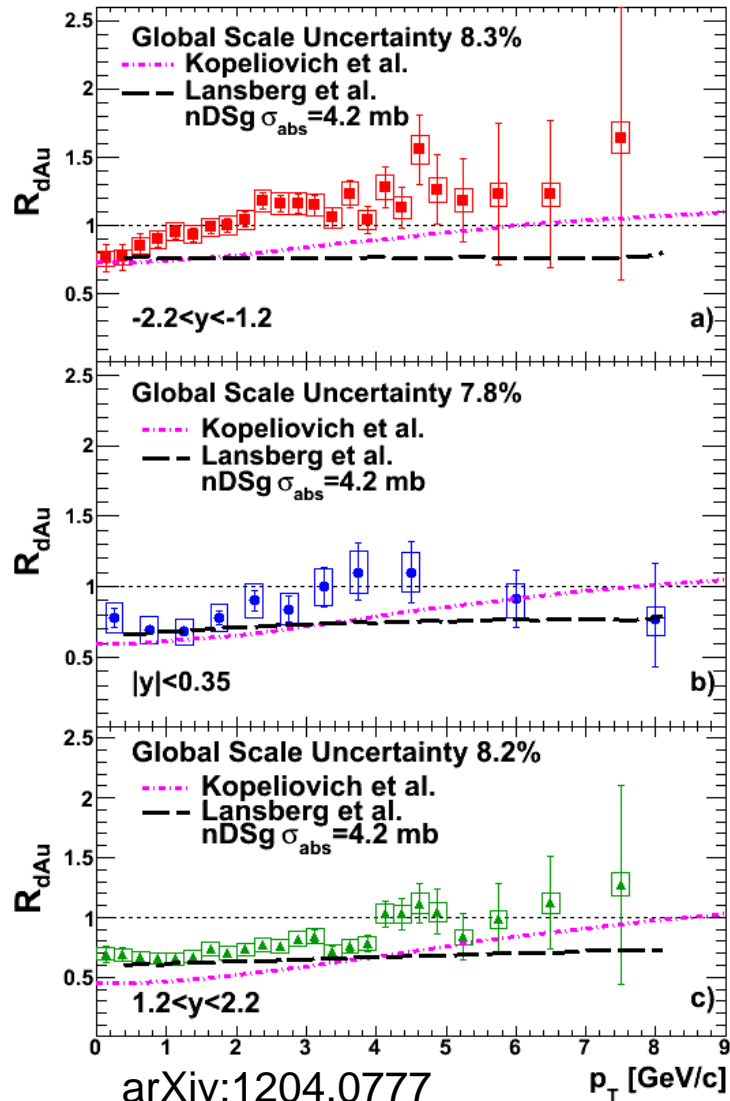
Quadratic: $M(r_T) = 1 - a\Lambda(r_T)^2$

- Break-up has exponential dependence
- EPS09 & initial-state dE/dx have unknown dependences

The forward rapidity points suggests a quadratic or higher geometrical dependence

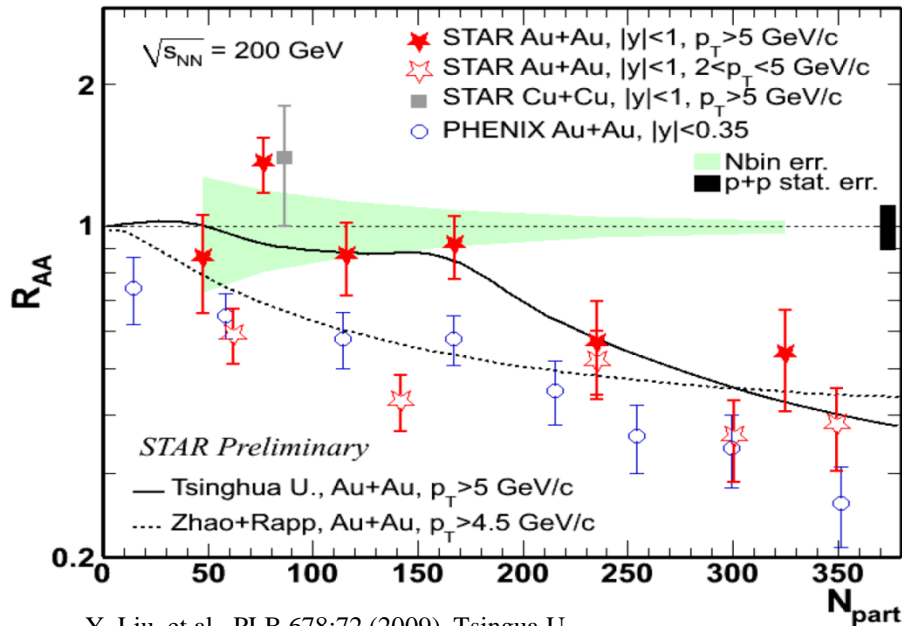
Cold nuclear matter effects – d+Au 200 GeV

Transverse momentum

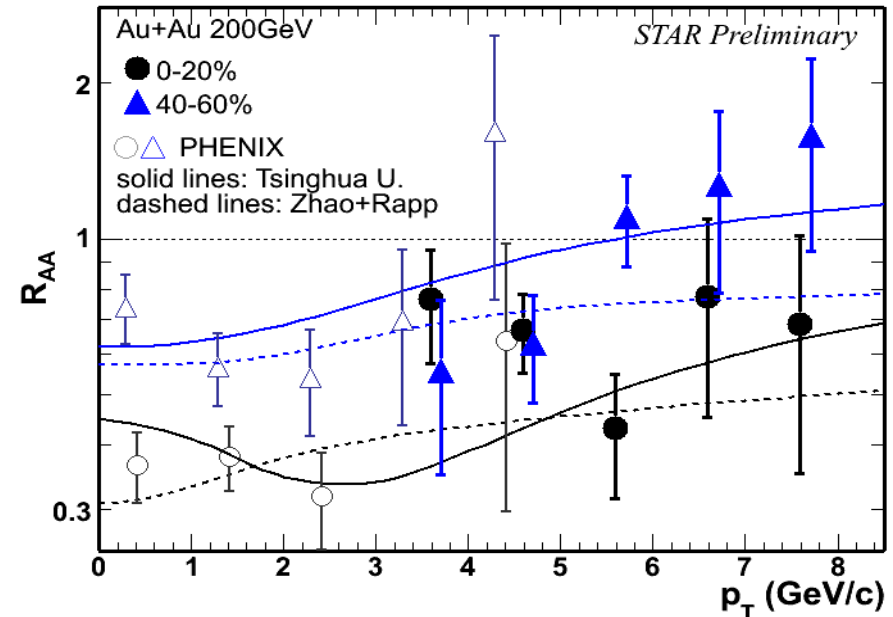


- Similar suppression at mid & forward rapidity
- Suppression for $p_T < 4$ GeV/c
- $R_{dAu} \approx 1$ for $p_T > 4$ GeV/c
- Backward rapidity: $R_{dAu} > 1$ for $p_T > 2$ GeV/c

High- p_T J/ψ measurement Au+Au 200 GeV



Y. Liu, et al., PLB 678:72 (2009), Tsinghua U.
 X. Zhao and R.Rapp, PRC 82, 064905(2010)



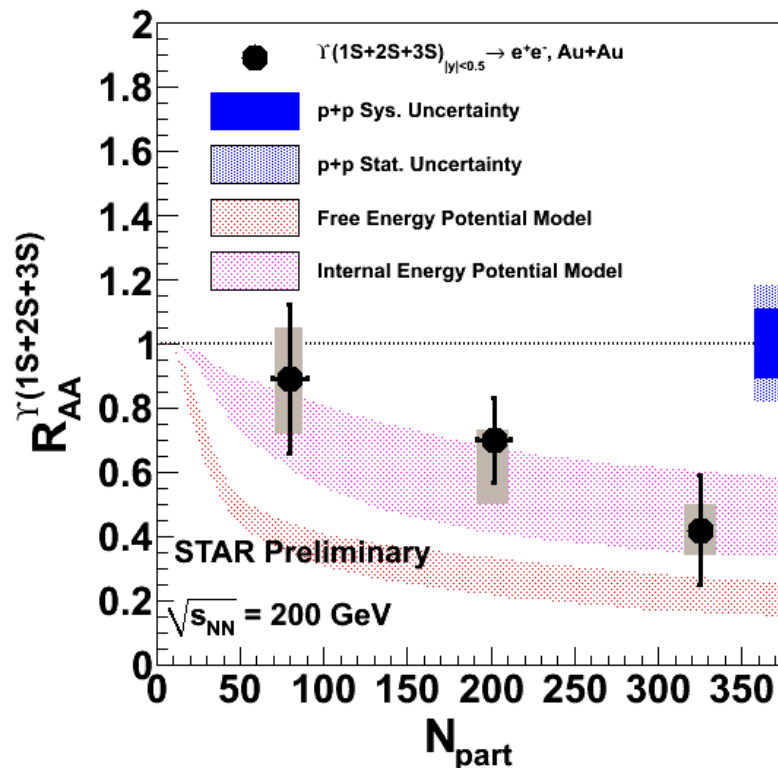
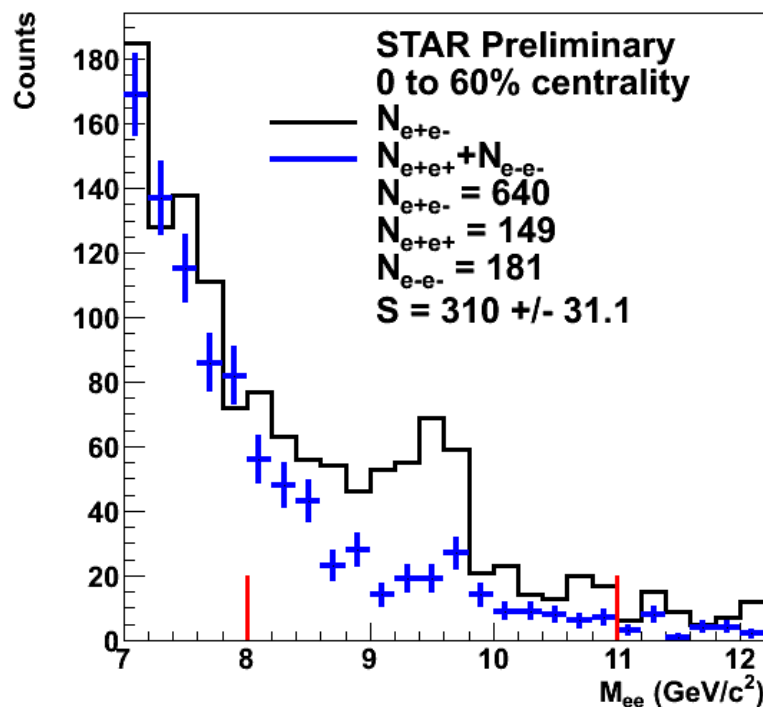
STAR CuCu: PRC80, 014922(R), PLB 678:72 (2009), PRC 82,064905(2010)
 PHENIX: PRL98, 232301

- Suppression of J/ψ in central and semi-central collisions is observed.
- R_{AA} increases with p_T and decreases with centrality.
- At high p_T suppression is present only in central collisions.

QUARKONIA

$$Y \rightarrow e^+e^-$$

Upsilon in Au+Au 200 GeV



- R_{AA} : Observation of Upsilon suppression.
(Including 2009 pp Preliminary $d\sigma/dy$)
 - Expect: Recombination: negligible, Hadronic co-mover absorption: negligible.
 - Suppression consistent with melting of excited states: **deconfinement effects**

Summary and outlook

- Heavy flavor is an important tool to understand medium properties.
- Results are interesting and challenging.

open charm measurement

- Charm hadrons; non-photonic electrons
- Charm production cross section.
- Separation of charm and bottom contribution.
- FONLL QCD describes the data rather well.

J/ψ

- SPS x RHIC mid x forward rapidity suppression.
- Systematic study of Cold nuclear effects.
- Less suppression at high- p_T in STAR.

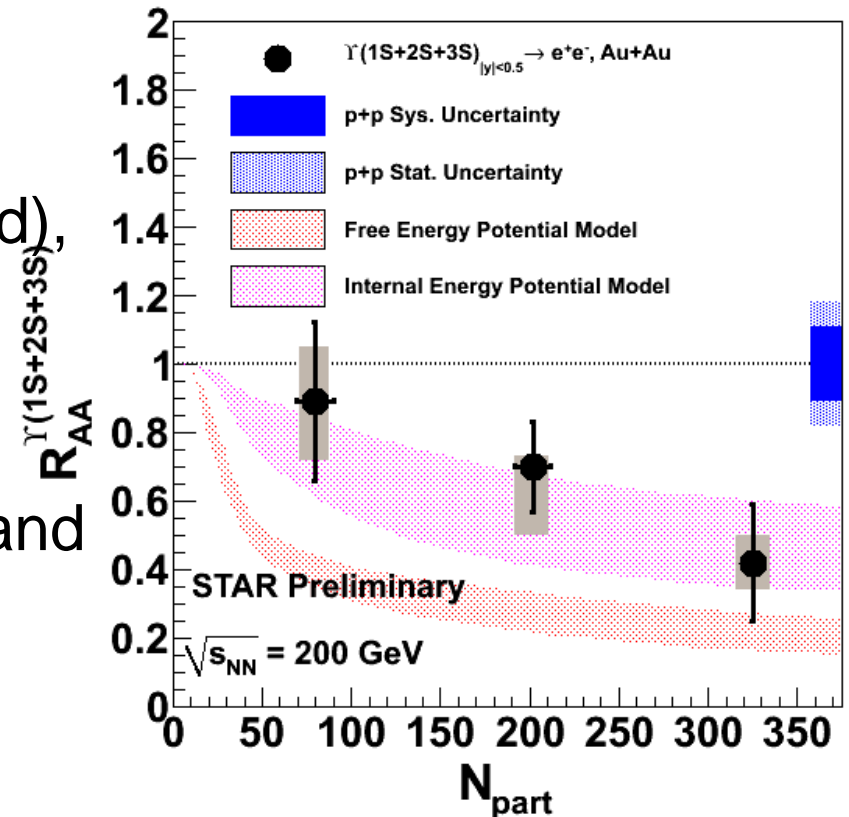
Y

- Suppression of $Y(1S+2S+3S)$ in central Au+Au observed.

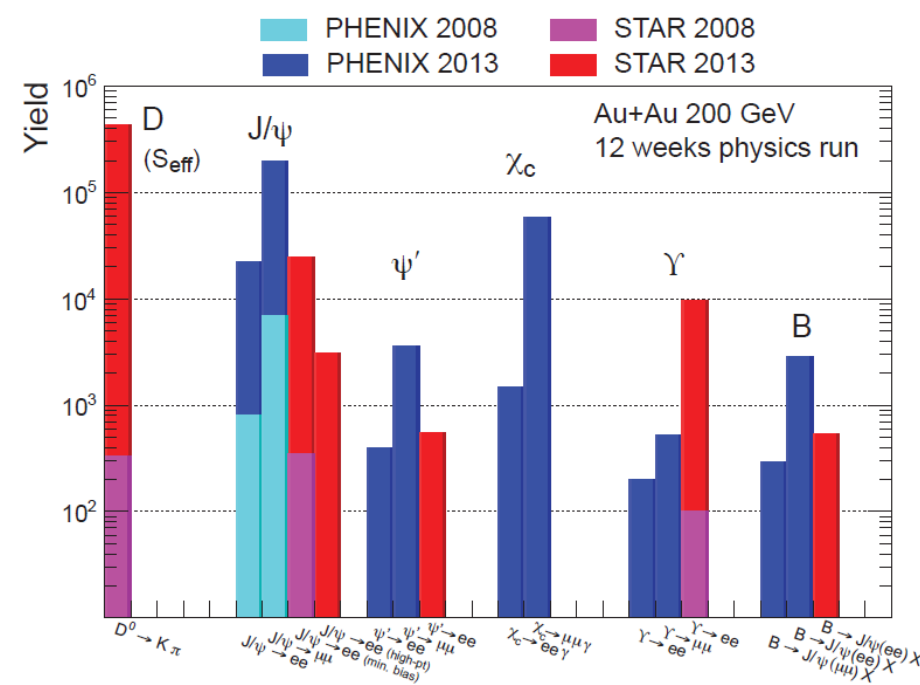
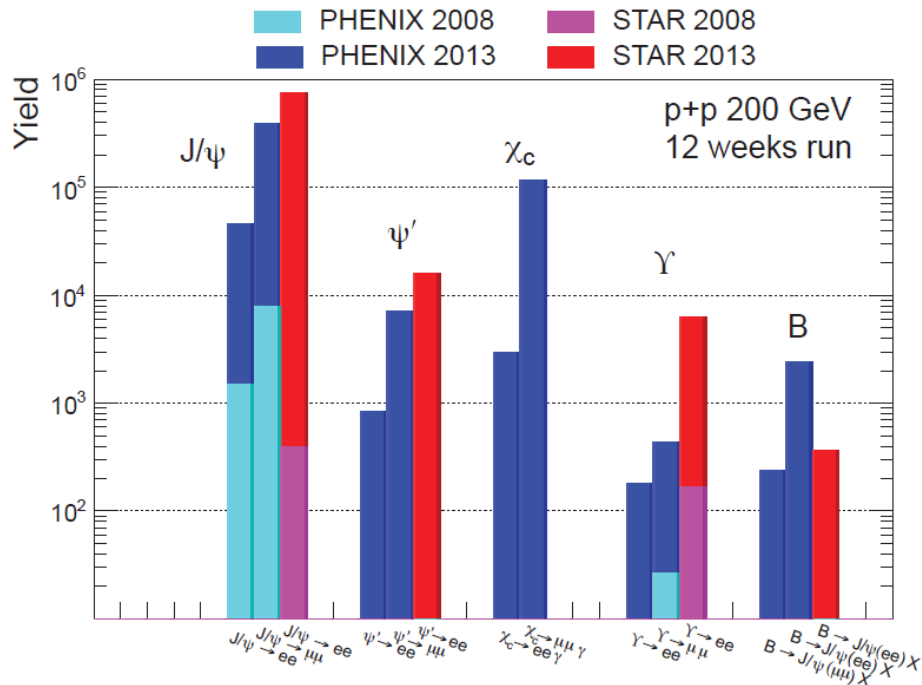
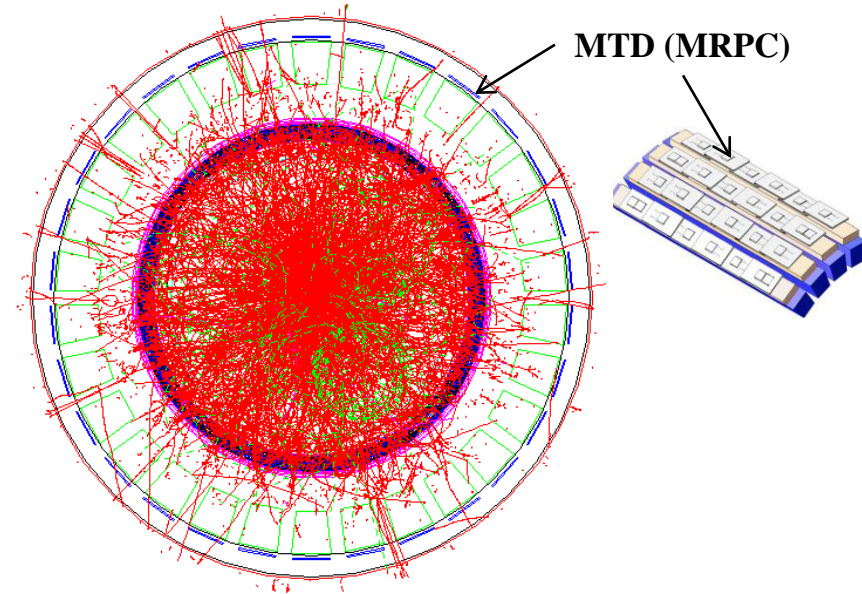
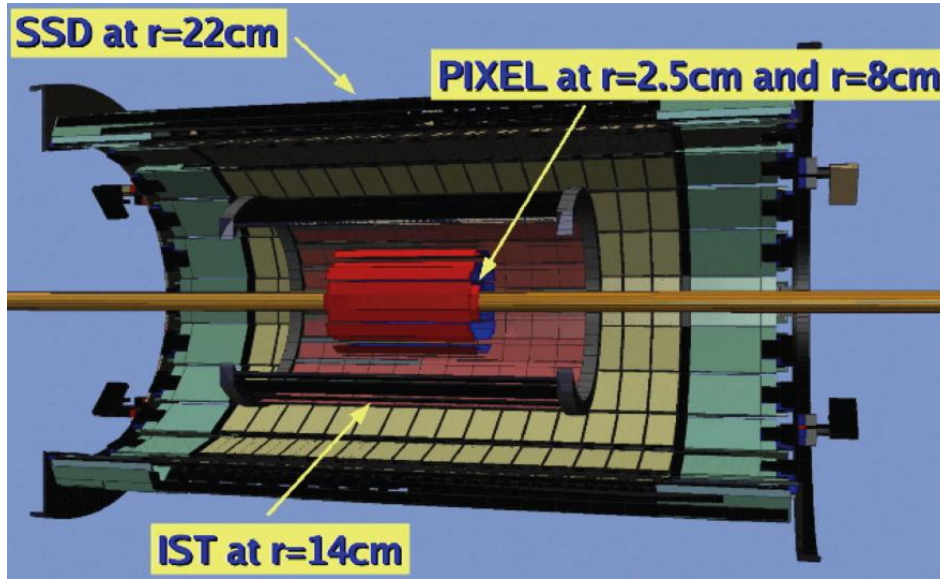
STAR/PHENIX upgrades: charm mesons flow; Y states separation
better charm/bottom separation

Υ model comparison

- Incorporating lattice-based potentials, including real and imaginary parts
 - A: Free energy (disfavored),
 - B: Internal energy (consistent with data vs. N_{part})
- Includes sequential melting and feed-down contributions
- Dynamical expansion, variations in initial conditions ($T, \eta/S$)
 - Data indicate: $428 < T_0 < 442 \text{ MeV}, 1 < 4\pi\eta/S < 3$

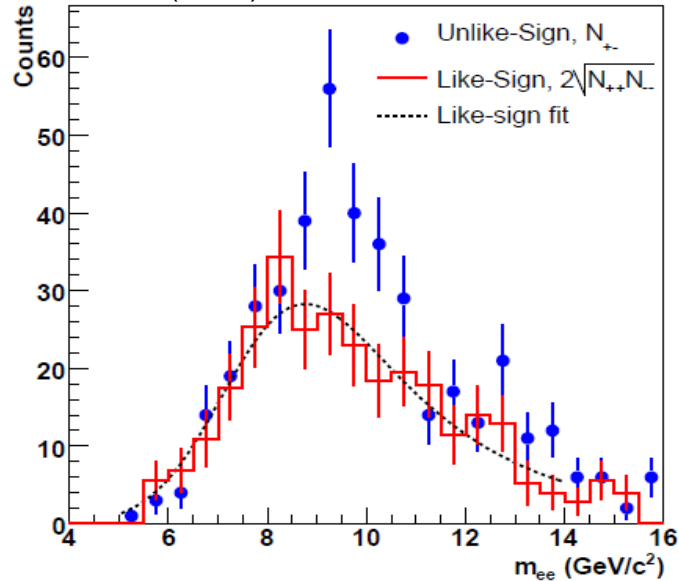


Future of Heavy Flavor Measurement at STAR

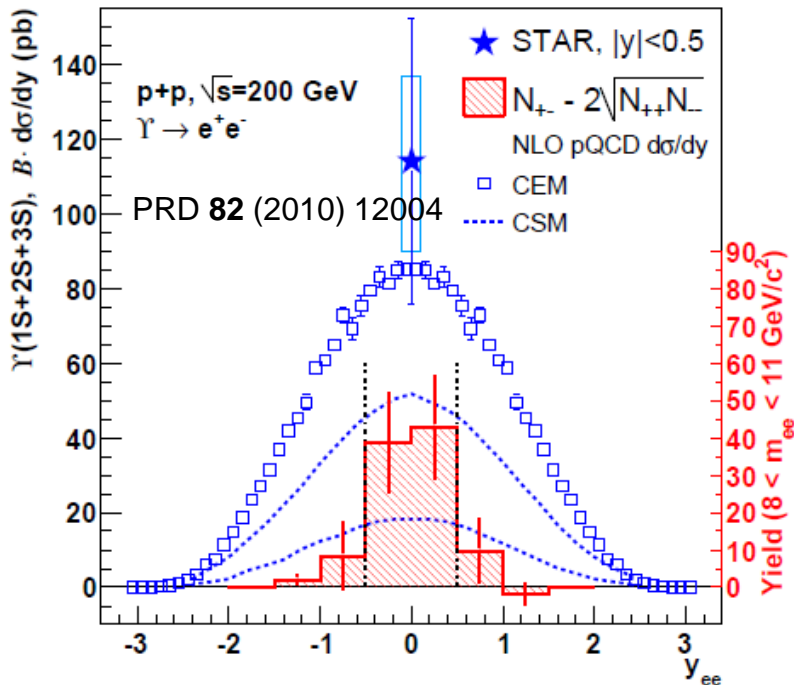
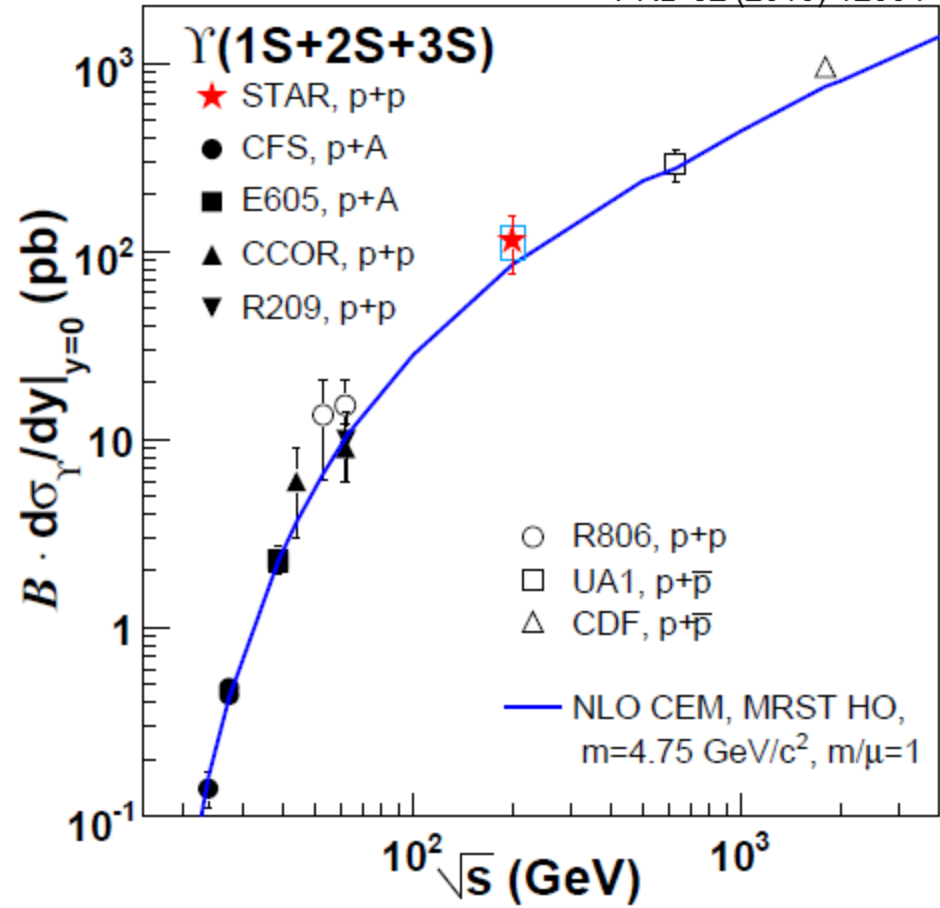


Upsilon in p+p 200GeV

PRD 82 (2010) 012004

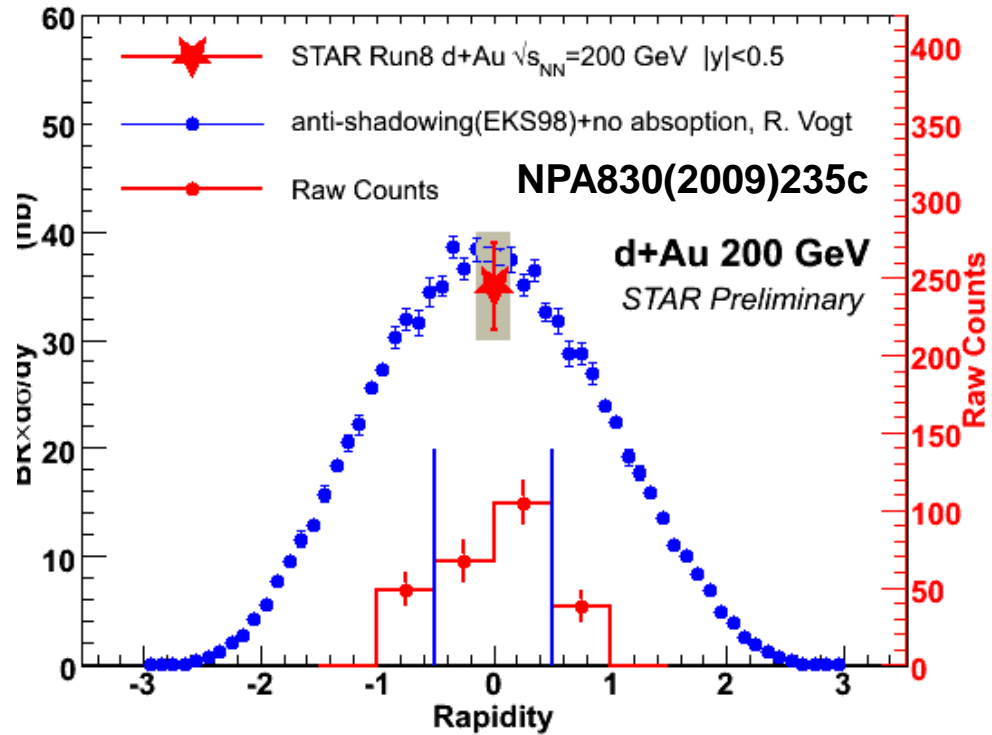
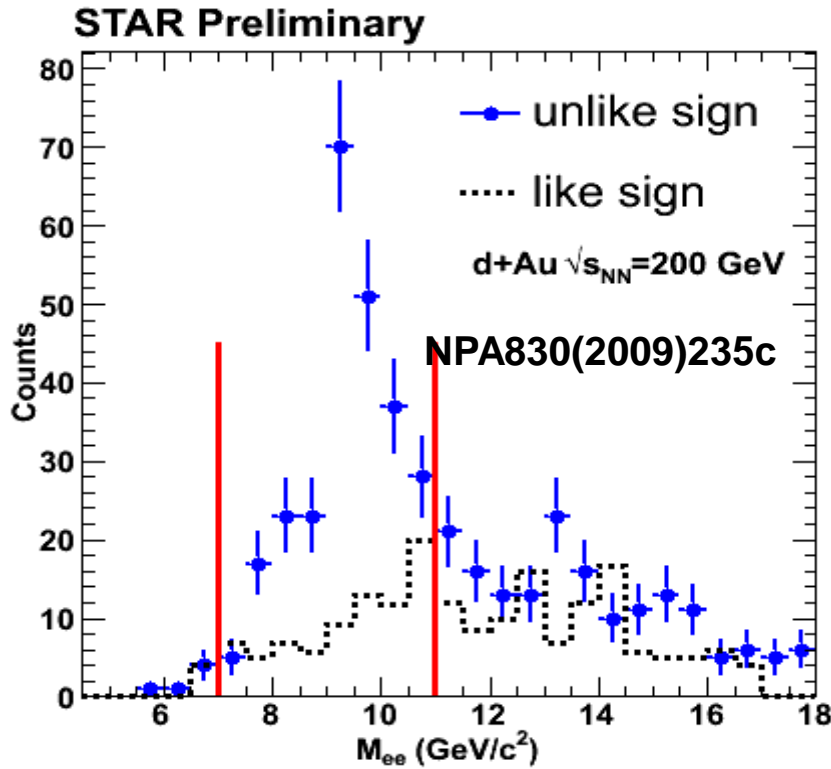


PRD 82 (2010) 12004



$$B_{ee} \left. \frac{d\sigma}{dy} \right|_{y=0} = 114 \pm 38(stat)_{-24}^{+23}(sys) \text{ pb}$$

Upsilon in d+Au 200GeV

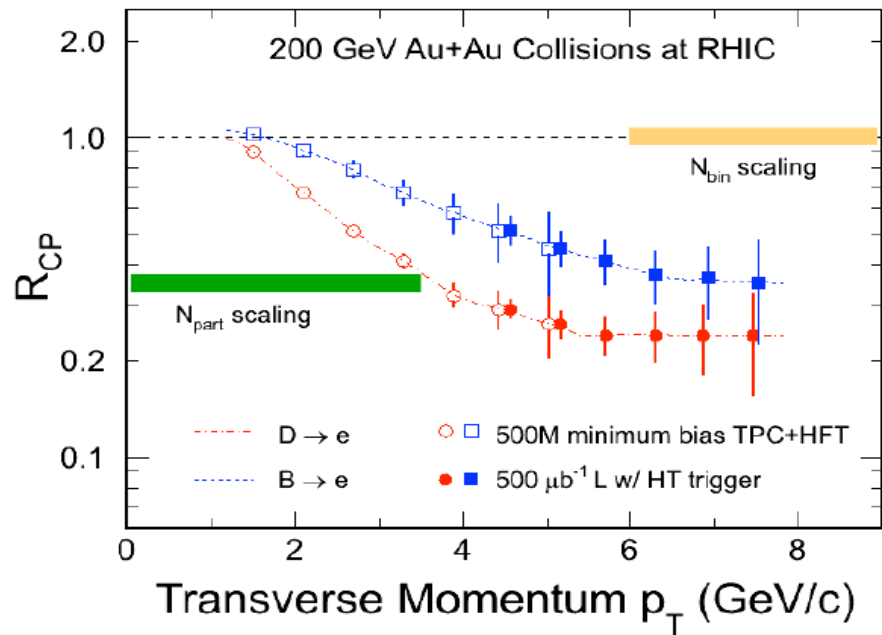
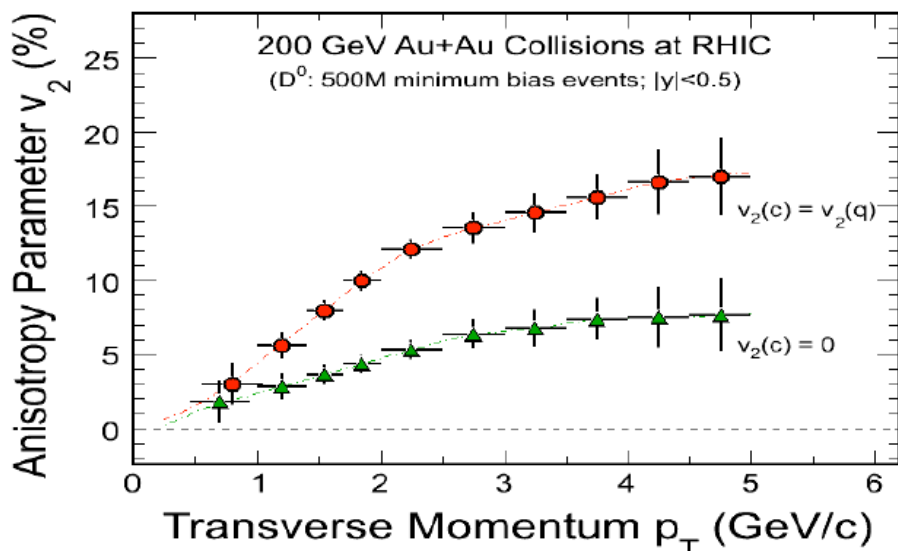
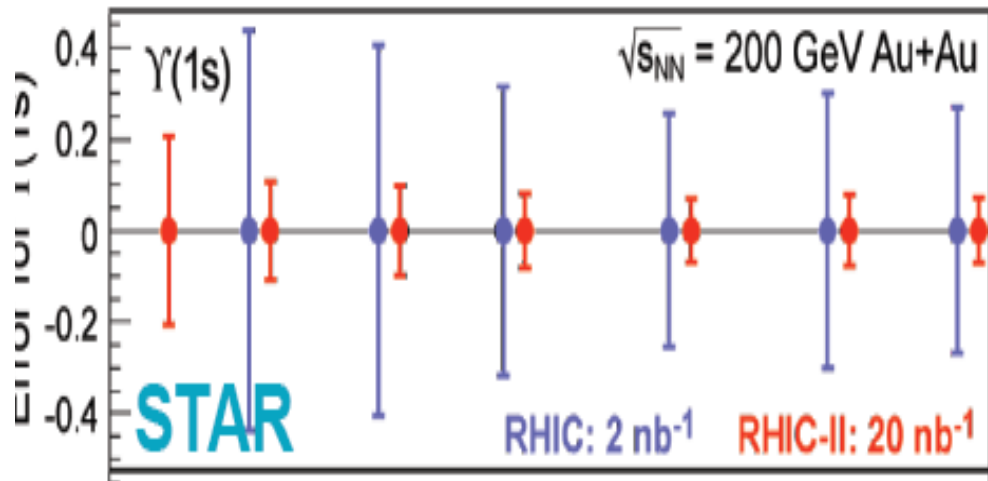
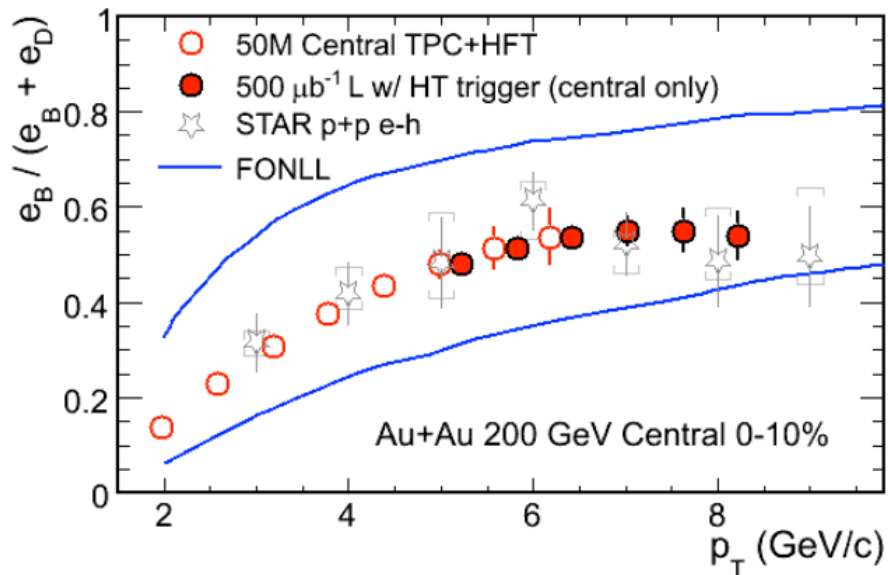


$$B_{ee} \left. \frac{d\sigma}{dy} \right|_{y=0} = 35 \pm 4(stat) \pm 5(sys) \text{ nb}$$

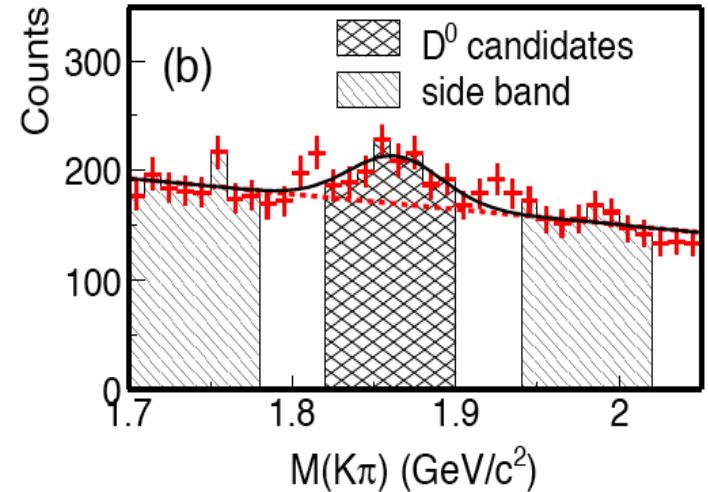
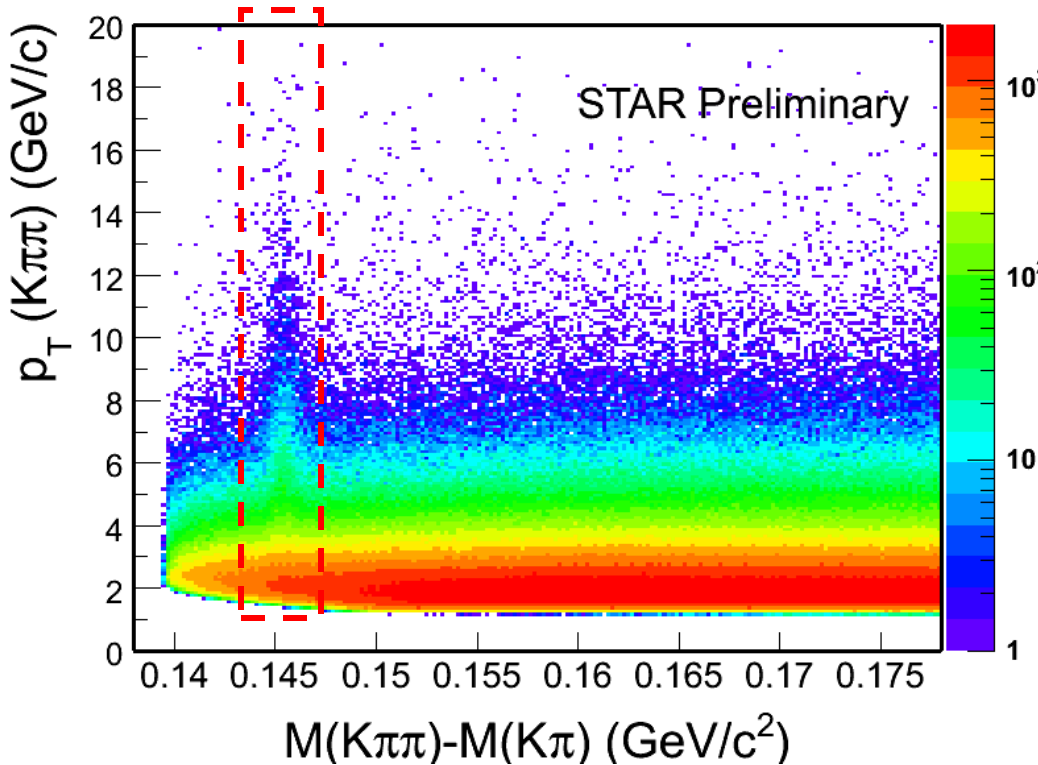
$$R_{dAu} = 0.8 \pm 0.3(stat) \pm 0.2(sys)$$

- Consistent with N_{bin} scaling of cross-section p+p - d+Au 200GeV

STAR with HFT



D* reconstruction



Background combinations:

Wrong sign:

D⁰ and π^- , D⁰bar and π^+

Side band:

$1.72 < M(K\pi) < 1.80$ or

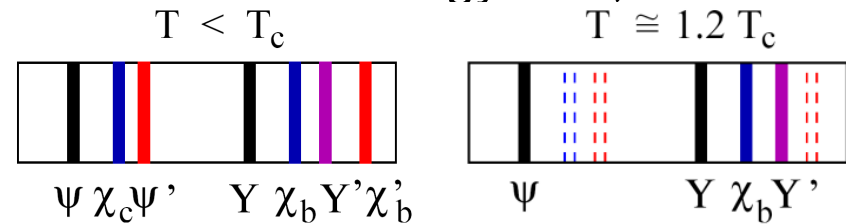
$1.92 < M(K\pi) < 2.0$ GeV/c²

All triggers included.

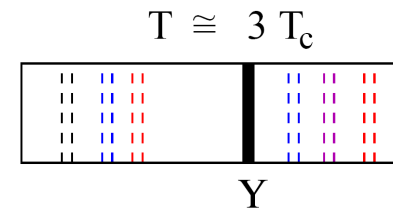
More than 4σ signal at low p_T and very significant at high p_T - mostly from EMC-based high neutral energy triggers.

Charmonia in nuclear matter

- Production mechanism is not clear
- Observed J/ψ is a mixture of direct production + feeddown
 - All $J/\psi \sim 0.6 J/\psi$ (Direct) + $\sim 0.3 \chi_c$ + $\sim 0.1 \psi'$
- Suppression and enhancement in the “cold” nuclear medium
 - Nuclear Absorption, Gluon shadowing, initial state energy loss, Cronin effect and gluon saturation

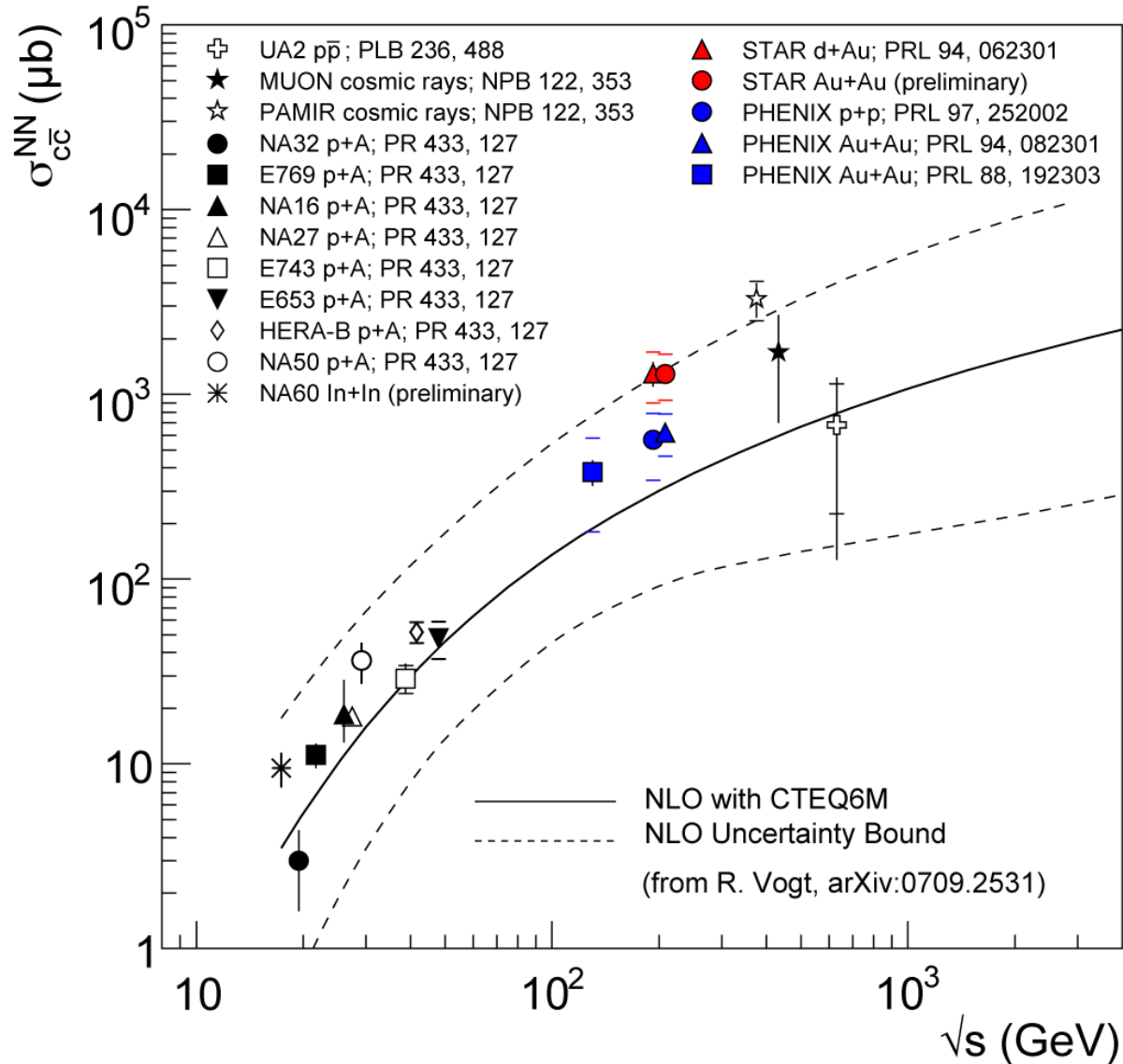


H. Satz, Nucl. Phys. A (783):249-260(2007)



- Hot/dense medium effect
 - J/ψ , Υ dissociation, i.e. suppression
 - Recombination from uncorrelated charm pairs

σ_{CC} : comparison with other measurements



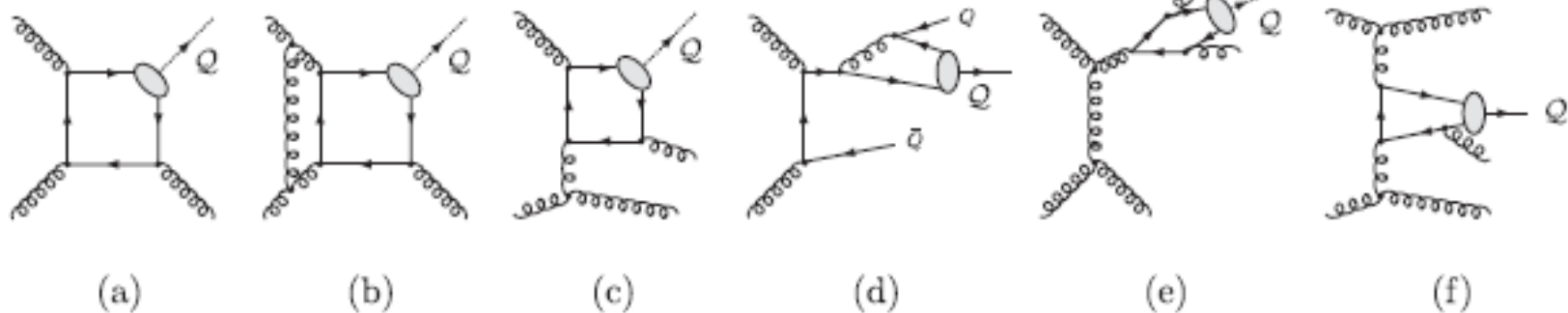
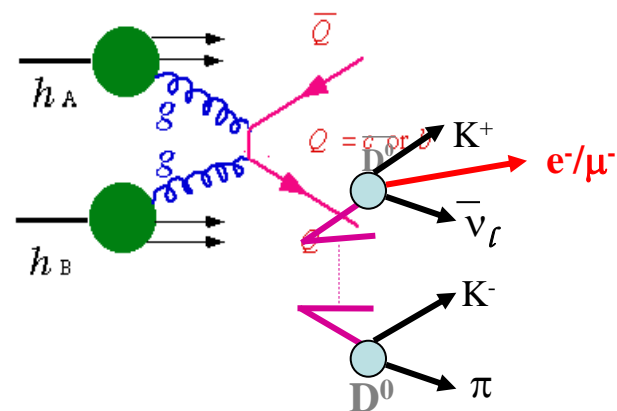
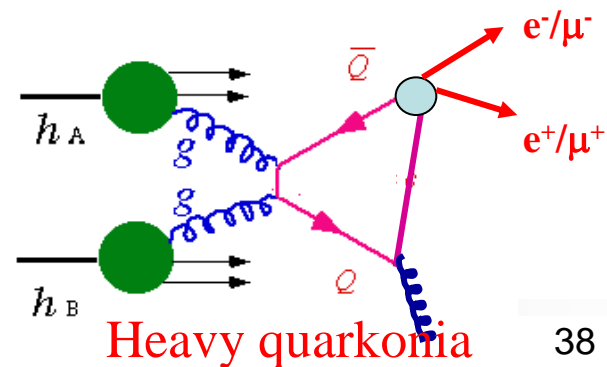
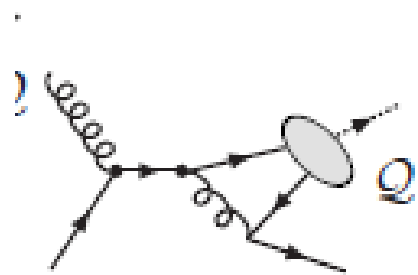


FIG. 1. Representative diagrams contributing to Y hadroproduction at orders α_S^3 (a), α_S^4 (b,c,d), α_S^5 (e,f). See discussions in the text.



Open heavy flavor



Heavy quarkonia

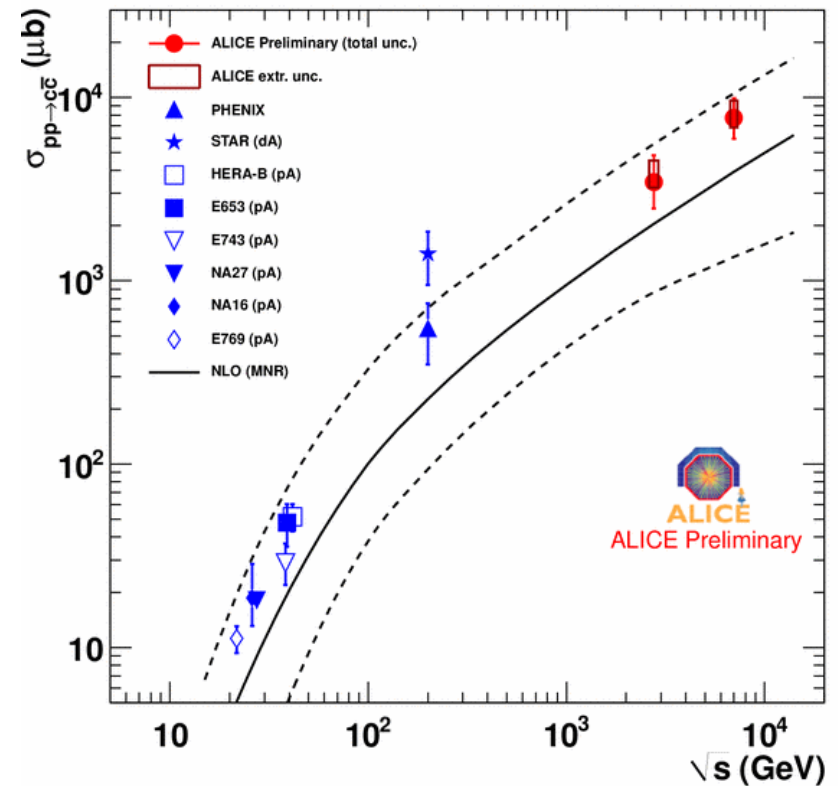
What can we learn at the LHC

- Higher c and b cross sections:
 - More abundant heavy flavour production
 - Better precision (reduced errors)

$$\sigma_{LHC}^{c\bar{c}} \approx 10 \cdot \sigma_{RHIC}^{c\bar{c}}$$

$$\sigma_{LHC}^{b\bar{b}} \approx 100 \cdot \sigma_{RHIC}^{b\bar{b}}$$

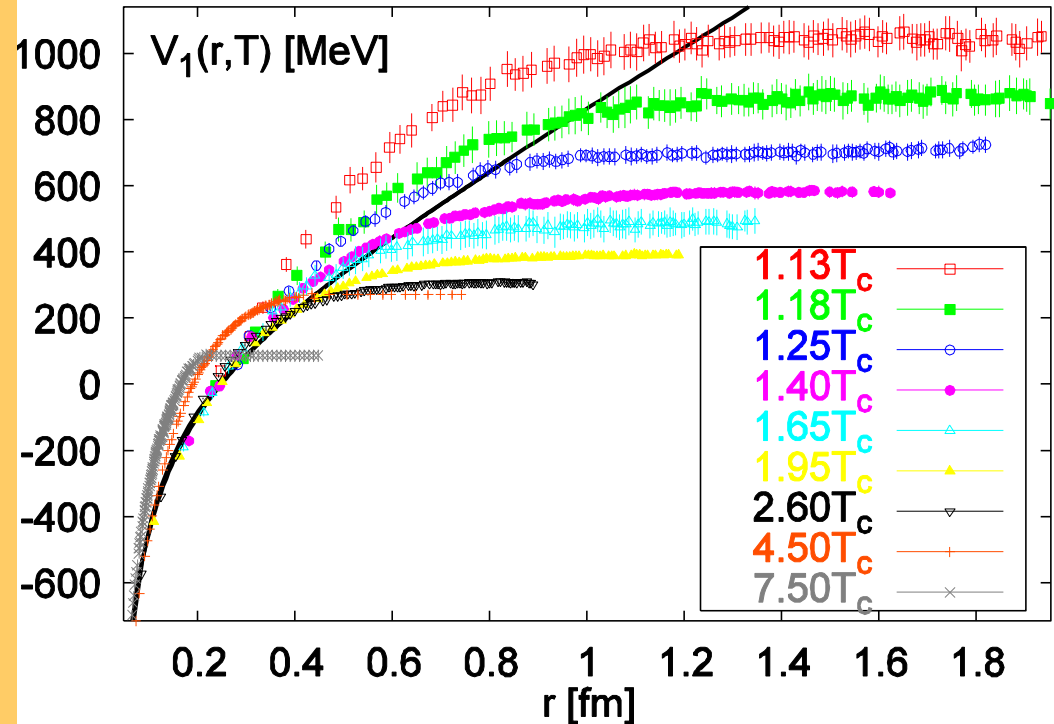
- High precision vertex detectors
 - Background removal
 - Separate c and b



ALI-PREL-8616

High T: the potential between the quarks is modified.

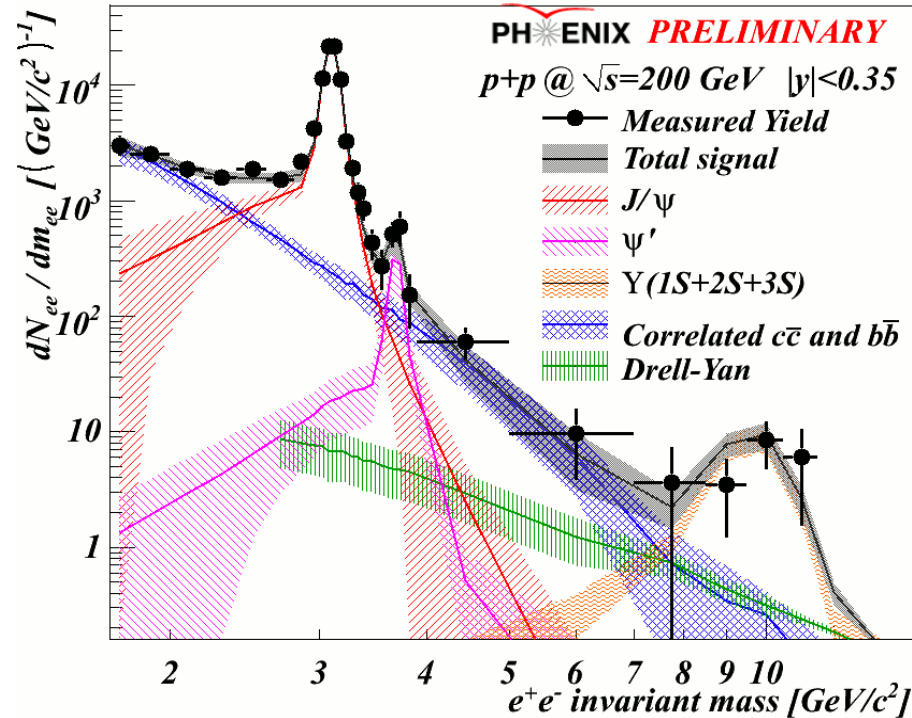
- Charmonium suppression: longstanding QGP signature
 - Original idea: High T leads to Debye screening
 - Screening prevents heavy quark bound states from forming!
 - **J/ψ suppression:**
 - Matsui and Satz, *Phys. Lett. B* 178 (1986) 416
 - lattice calculations confirm screening effects
 - Nucl.Phys.Proc.Suppl.129: 560-562,2004



O. Kaczmarek, et al.,
Nucl.Phys.Proc.Suppl.129:560-562,2004

Better Knowledge about the Baseline

$$\psi' \rightarrow e^+e^-$$

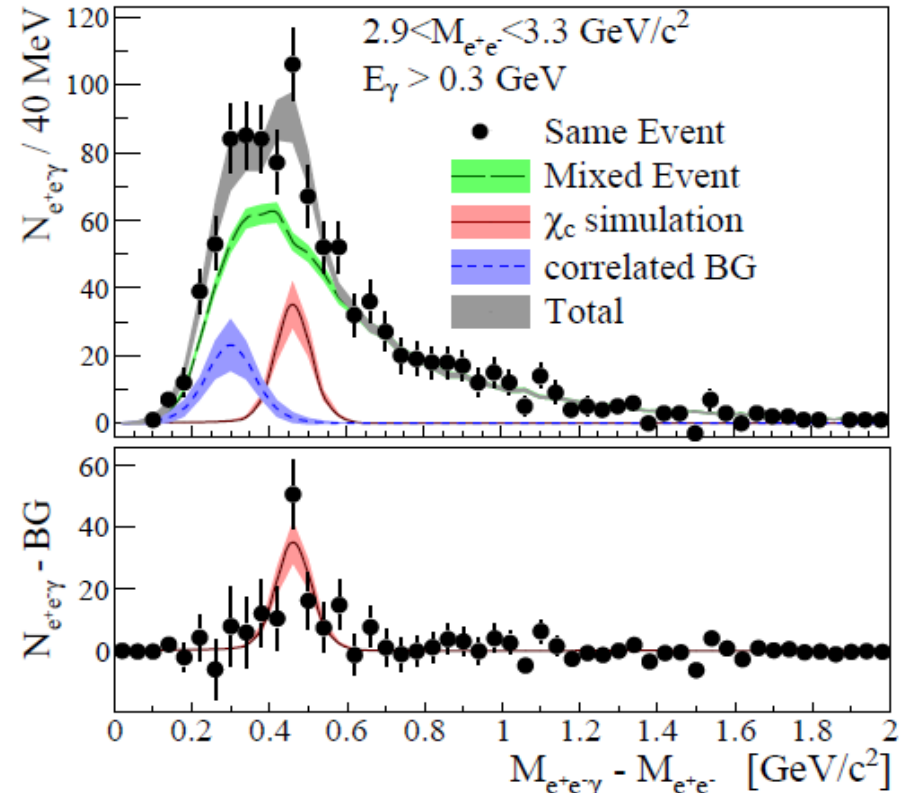


$$F_{\psi'}^{J/\psi} = \frac{B_{J/\psi}^{\psi'} \sigma_{\psi'}}{\sigma_{J/\psi}} = (9.6 \pm 2.4)\%$$

Consistent with world average!!!

9/19/2011

$$\chi_c \rightarrow J/\psi + \gamma$$



$$F_{\chi_c}^{J/\psi} = \frac{N_{\chi_c}}{N_{J/\psi}} \frac{1}{\langle \epsilon_{\chi_c} / \epsilon_{J/\psi} \rangle} = (32 \pm 9)\%$$

41

Non-photonic R_{AA} at RHIC

DGLV:

Djordjevic, PLB632, 81 (2006)

BDMPS:

Armesto, et al., PLB637, 362 (2006)

T-Matrix:

Van Hees et al., PRL100,192301(2008).

Coll. Dissoc.

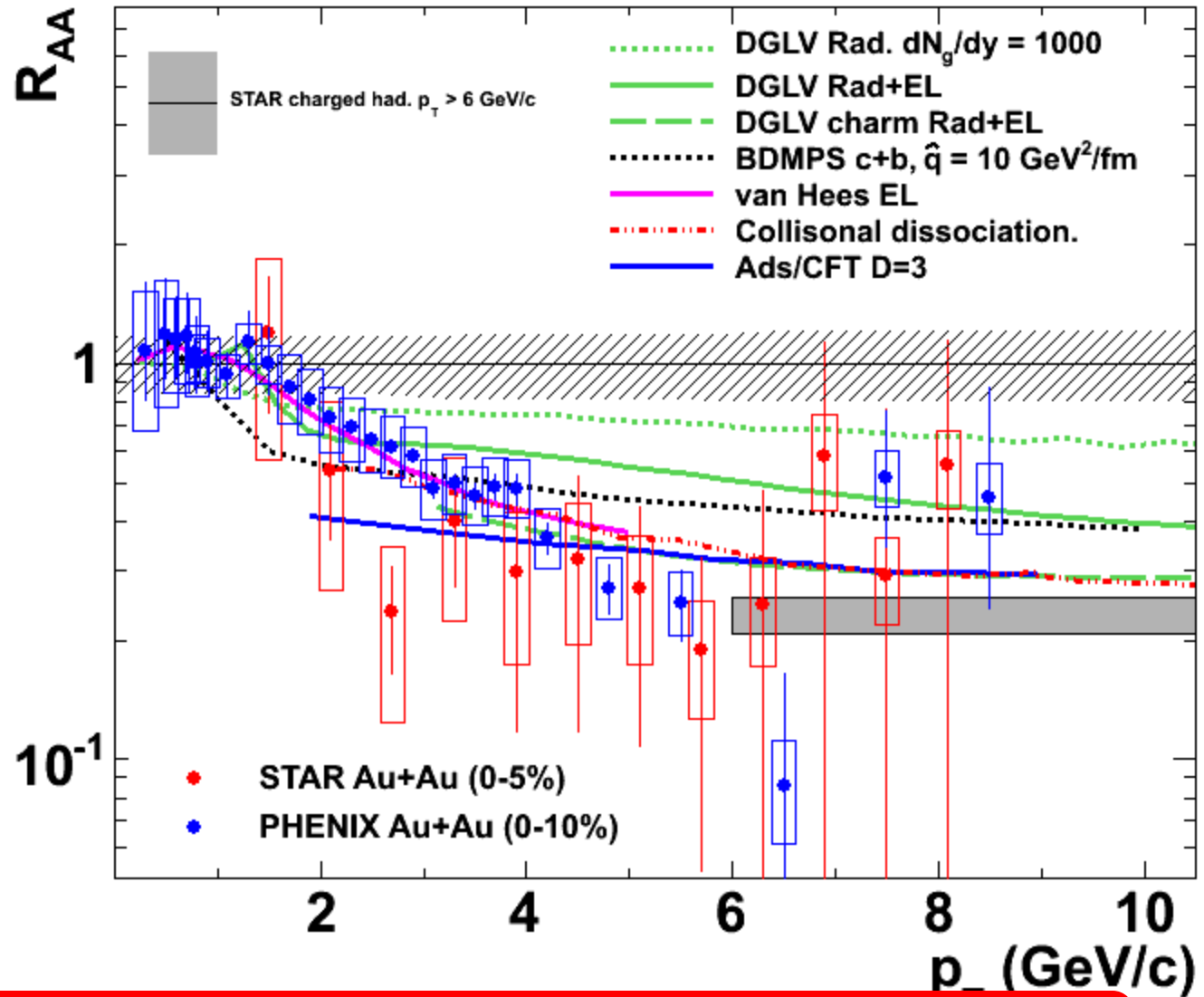
R. Sharma et al., PRC 80, 054902(2009).

Ads/CFT:

W. Horowitz Ph.D thesis.

RL.+ Coll.

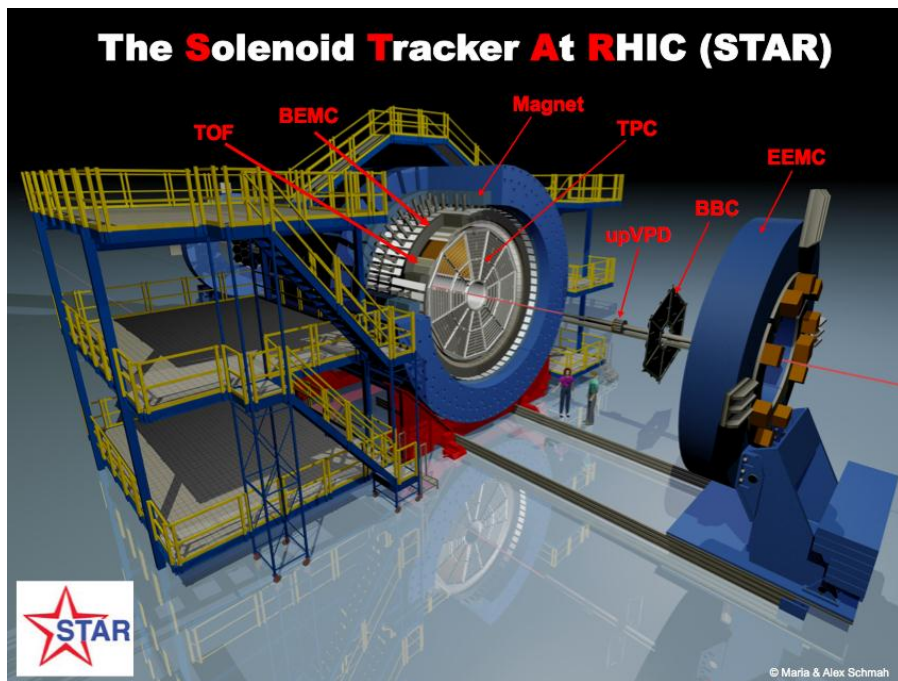
J. Aichelin et al., SQM11



Models with different or similar mechanisms can or can not describe the data

➤ Which one is right and what are missing?

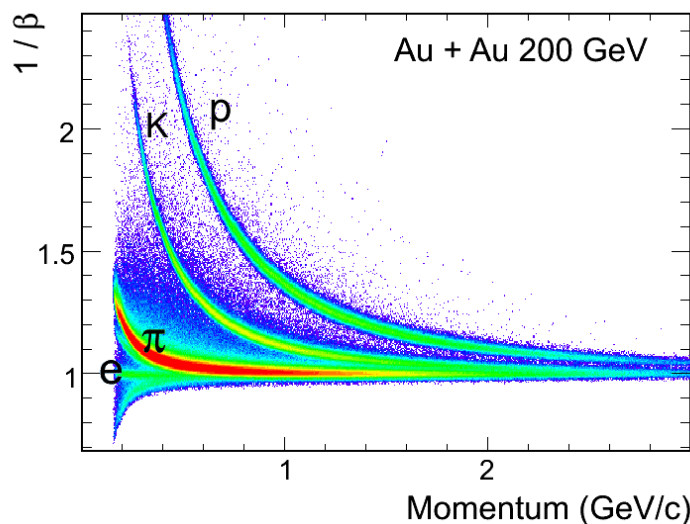
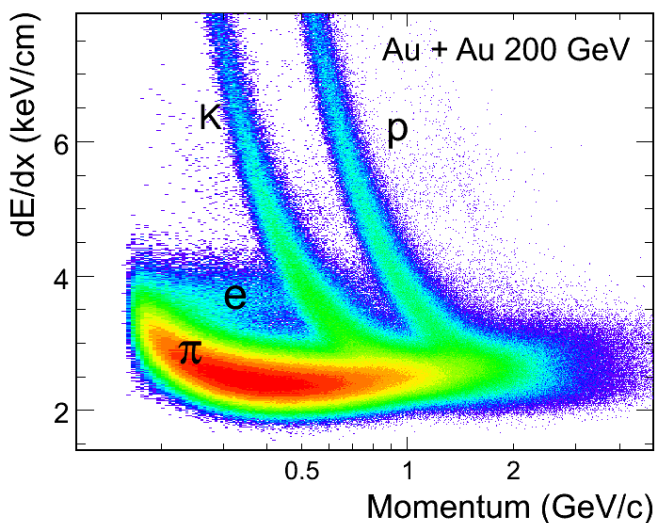
STAR detector and Particle ID



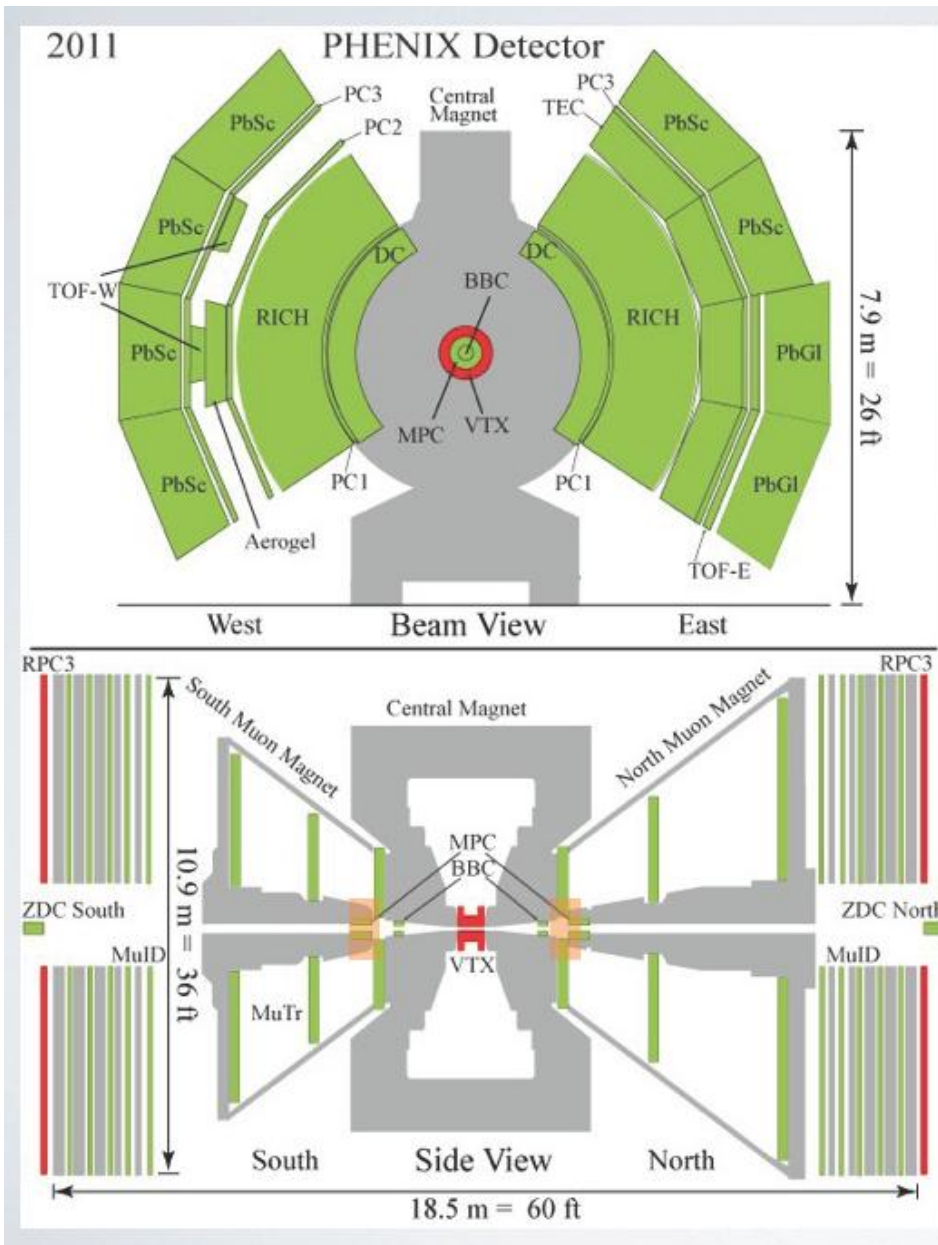
Large acceptance

$$|\eta| < 1, \quad 0 < \phi < 2\pi$$

- **T**ime **P**rojection **C**hamber
dE/dx, momentum
- **T**ime **O**f **F**light detector
particle velocity $1/\beta$
- **E**lectro**M**agnetical **C**alorimeter
E/p, single tower/topological Trigger



PHENIX detector and Particle ID

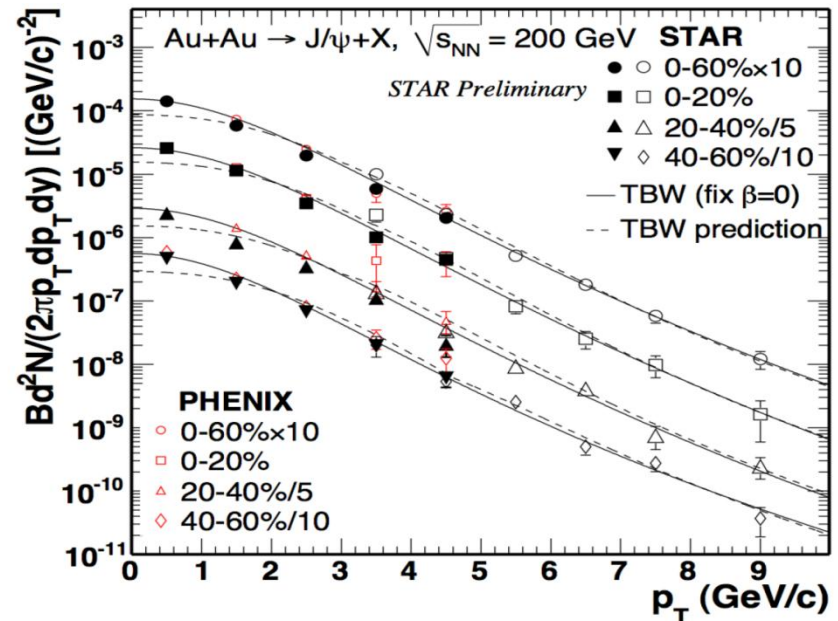
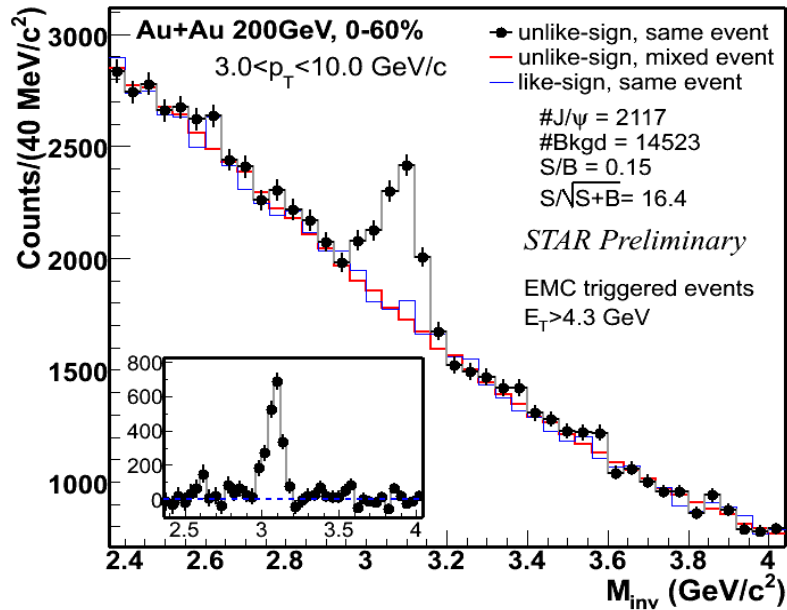


- Central Arms** $|\eta| < 0.35$
- Identified charged hadrons
 - e^\pm, π^0, η
 - **Direct Photon**
 - $J/\psi, \psi', \chi_c$
 - **Heavy Flavor**

- Muon Arms** $1.2 < |\eta| < 2.4$
- $J/\psi, \gamma$
 - Unidentified charged hadrons
 - **Heavy Flavor**

- MPC** $3.1 < |\eta| < 3.9$
- π^0, η

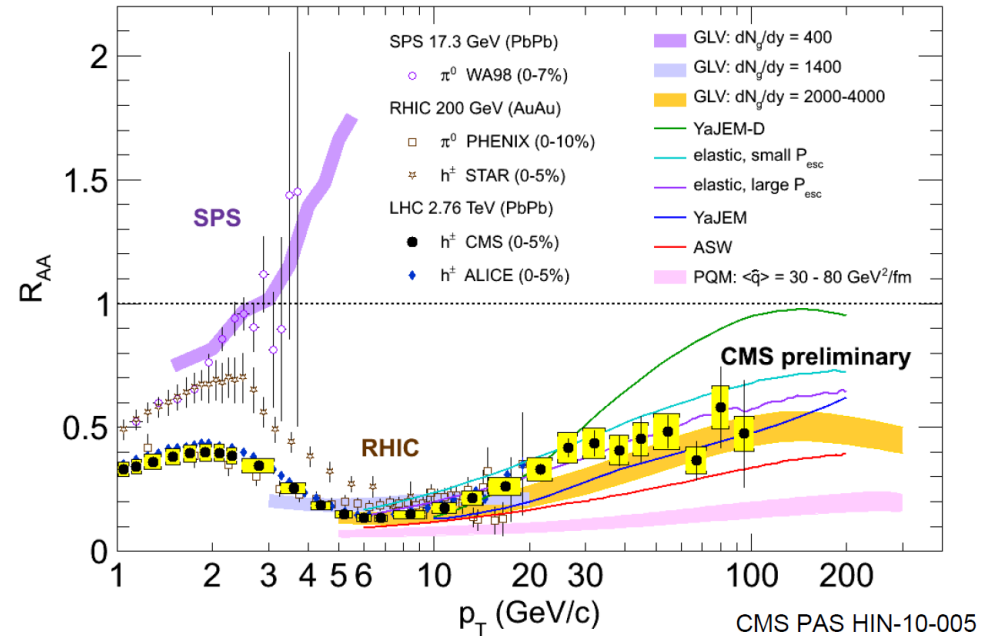
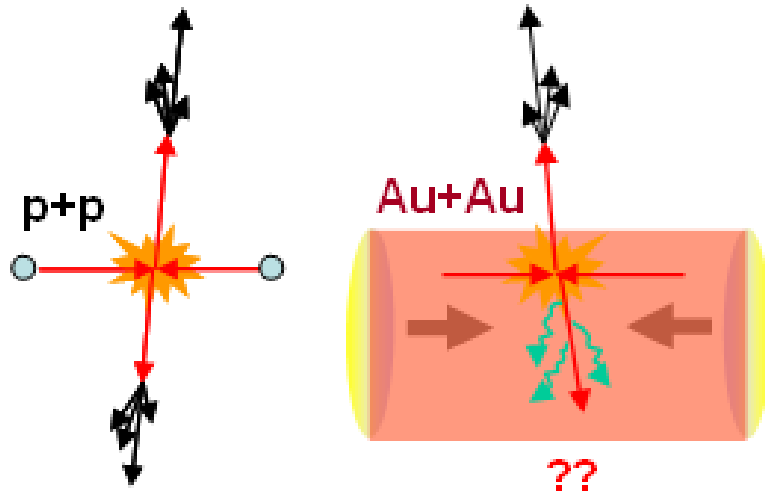
STAR J/ψ spectra in Au+Au 200GeV



Phys. Rev. Lett. 98, 232301 (2007)
 JPG 37, 085104 (2010)
 ArXiv:1101.1912 (2011)

Consistent with other RHIC measurements. Moreover we extend p_T region up to 10 GeV/c.

Nuclear modification factor



- **Hard probes** - produced in hard scatterings in initial phase of collision
- Nuclear matter influences the final particle production
e.g. production of particles at given p_T
supresion of particle production of particular type
- **Nuclear modification factor** - quantification of nuclear effects R_{AA}

$$R_{AA}(p_T) = \frac{\text{Yield}_{AA}(p_T)}{\langle N_{bin} \rangle_{AA} \text{Yield}_{pp}(p_T)}$$

Cold nuclear matter effects

- **Nuclear Shadowing** – Modification of PDF's for nucleons bound in nuclei.
Parametrizations of (mostly) DIS data (ex. EKS98, nDSg, EPS09).
- **Nuclear Break-up** – Break-up of cc pair through collisions with nucleons.
Usually parametrized using break-up cross section.
- **Cronin Effect** – Broadening of the p_T distribution through scattering of incoming partons.
- **Initial State Energy Loss** – decrease in parton momentum due to soft scatterings while propagating through colliding nucleus.

Cold nuclear matter effects

