

Charged-hadron production in the three- sources RDM at LHC energies

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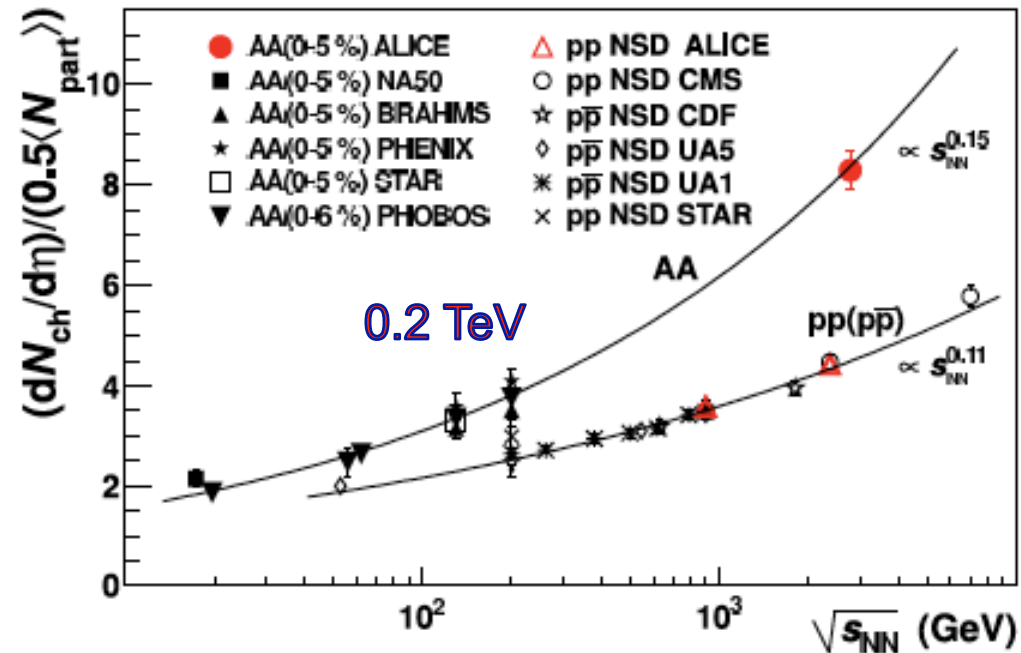


ICFP 2012

Topics

1. Introduction: PbPb @ LHC
2. Charged-hadron pseudorapidity distributions in a nonequilibrium-statistical approach (RDM)
3. Comparison with RHIC and LHC data
4. Conclusion

1. Intro: ALICE PbPb results @ LHC energy 2.76 TeV



$N_p \approx 382$

Talk by P.
Christagoklou
on June 12

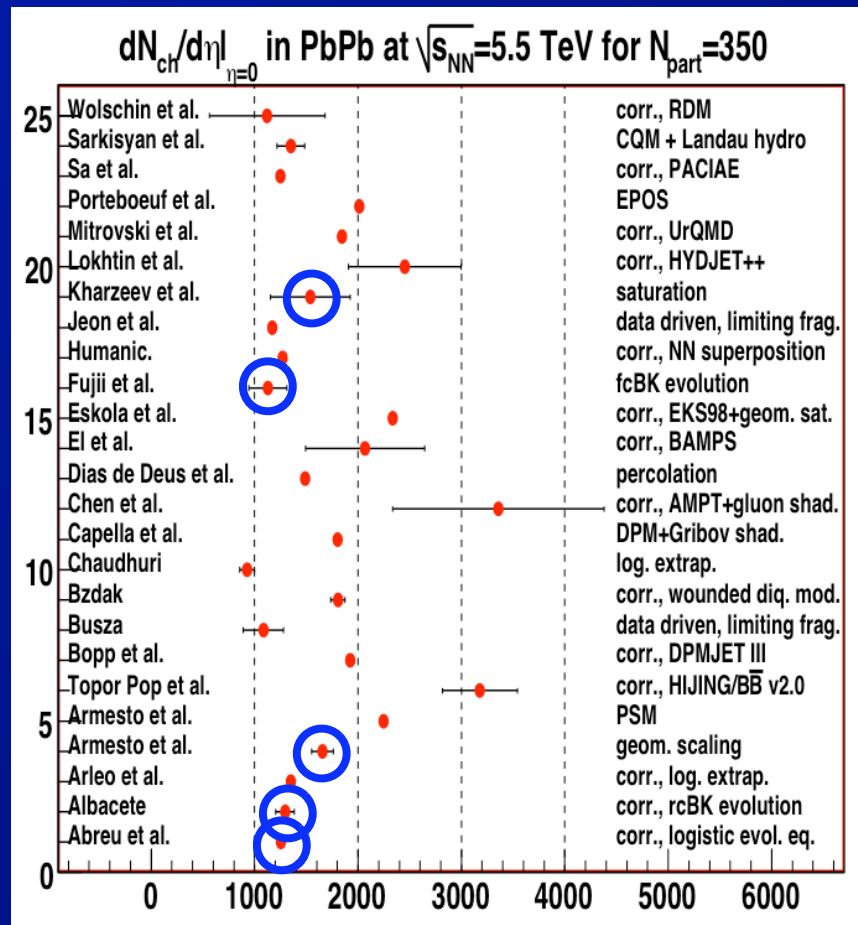
$$\eta = -\ln[\tan(\theta/2)]$$

FIG. 3. Charged particle pseudo-rapidity density per participant pair for central nucleus–nucleus [16–24] and non-single diffractive pp/p \bar{p} collisions [25–31], as a function of $\sqrt{s_{NN}}$. The energy dependence can be described by $s_{NN}^{0.15}$ for nucleus–nucleus, and $s_{NN}^{0.11}$ for pp/p \bar{p} collisions.

ALICE collab., PRL 105, 252301 (2010)(see also Eur. Phys. J. C 65, 111 (2010) for p + p)

$dN/d\eta(\eta=0) = 1584 \pm 4$ (stat.) ± 76 (sys.): 8.3 ± 0.4 per part. nucleon pair

LHC dN/dη Predictions



- Many different predictions for LHC Pb+Pb central dN/dη – @ 5.5 TeV
- Saturation (motivated) predictions at low end of range – 1200-1600

From: N. Armesto et al., J.Phys. G35, 05400 (2008); cf also B. Cole, ICHEP Paris 2010

Predictions vs. ALICE PbPb data of $dN/d\eta$ ($\eta \approx 0$) @ 2.76 TeV

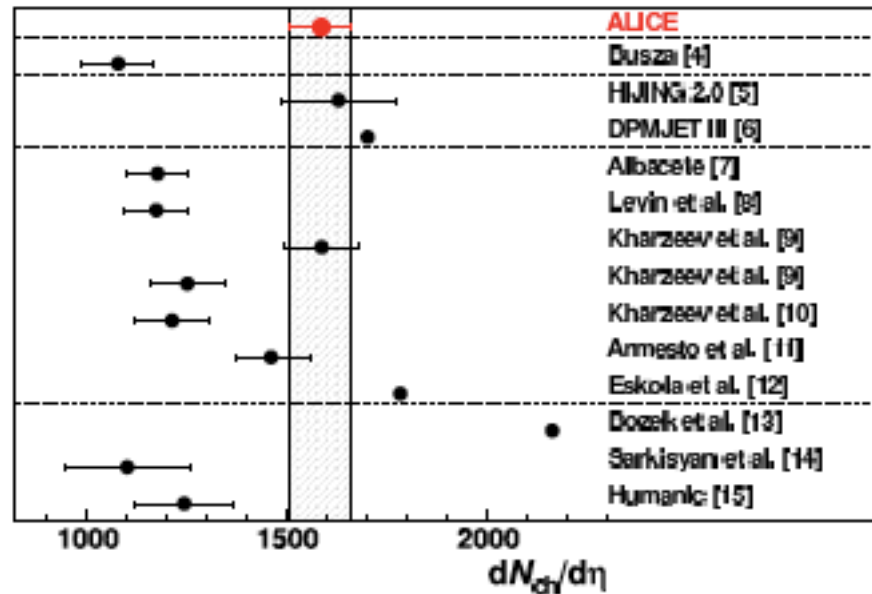


FIG. 4. Comparison of this measurement with model predictions. Dashed lines group similar theoretical approaches.

from: ALICE collab., PRL 105, 252301 (2010).

More interesting than the midrapidity value is the distribution function.

2. Relativistic Diffusion Model (RDM)

$$\frac{\partial}{\partial t} R(y, t) = -\frac{\partial}{\partial y} [J(y)R(y, t)] + D_y \frac{\partial^2}{\partial y^2} [R(y, t)]^{2-q}$$

R (y,t) Rapidity distribution function. The standard linear Fokker-Planck equation corresponds to $q = 1$, and a linear drift function. For the three components $k = 1,2,3$ of the rapidity distribution,

$$\frac{\partial}{\partial t} R_k(y, t) = -\frac{1}{\tau_y} \frac{\partial}{\partial y} [(y_{eq} - y) \cdot R_k(y, t)] + D_y^k \frac{\partial^2}{\partial y^2} R_k(y, t)$$

Linear drift term with relaxation time τ_y Diffusion term, $D_y = \text{const.}$

Relaxation time and diffusion coefficient are related through a **dissipation-fluctuation theorem**. The broadening is enhanced due to collective expansion.

$$\langle y_{1,2}(t) \rangle = y_{eq} [1 - \exp(-t/\tau_y)] \mp y_{max} \exp(-t/\tau_y) \quad \text{mean value}$$

$$\sigma_{1,2,eq}^2(t) = D_y^{1,2,eq} \tau_y [1 - \exp(-2t/\tau_y)] \quad \text{variance}$$

Linear Model: G. Wolschin, Eur. Phys. J. A5, 85 (1999); with 3 sources: Phys. Lett. B 569, 67 (2003); PLB 698, 411 (2011); M. Biyajima, M. Ide, M. Kaneyama, T. Mizoguchi, and N. Suzuki, Prog. Theor. Phys. Suppl. 153, 344 (2004)

Pseudorapidity distributions are then obtained through the Jacobian transformation

$$\frac{dN}{d\eta} = \frac{dN}{dy} \frac{dy}{d\eta} = \frac{p}{E} \frac{dN}{dy} \simeq J(\eta, \langle m \rangle / \langle p_T \rangle) \frac{dN}{dy}$$

$$J(\eta, \langle m \rangle / \langle p_T \rangle) = \cosh(\eta) \cdot$$

$$[1 + (\langle m \rangle / \langle p_T \rangle)^2 + \sinh^2(\eta)]^{-1/2}.$$

with the rapidity distribution

$$\frac{dN_{ch}(y, t = \tau_{int})}{dy} = N_{ch}^1 R_1(y, \tau_{int}) + N_{ch}^2 R_2(y, \tau_{int}) + N_{ch}^{eq} R_{eq}(y, \tau_{int}).$$

- Consider pions, kaons, protons
- Determine $J_{y=0}$ at $y = \eta = 0$ with the experimental values for $dN/dy(\pi, K, p)$ and $dN/d\eta$ (charged hadrons)

- Compute J using $\langle m \rangle = m_\pi$ and $\langle p_T \rangle = \langle p_{T,eff} \rangle$ with

$$\langle p_{T,eff} \rangle = m_\pi J_{y=0} / \sqrt{1 - J_{y=0}^2}$$

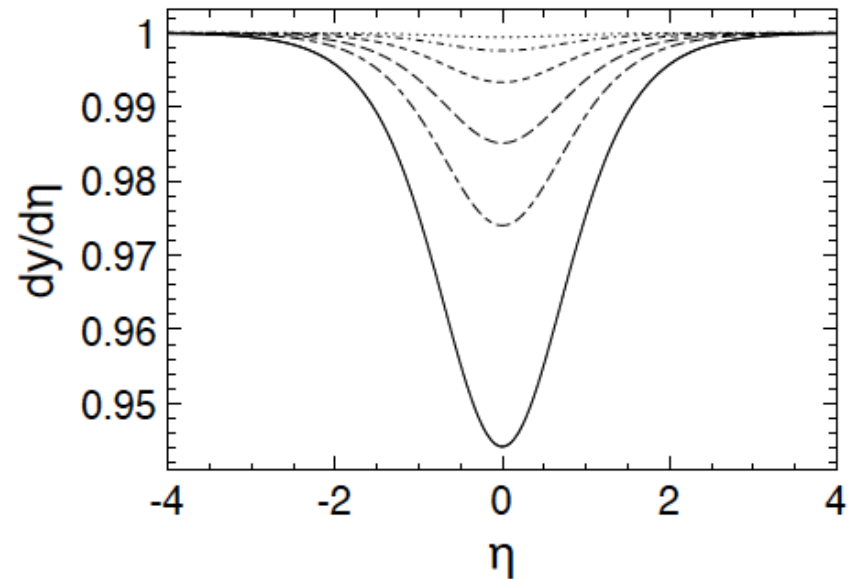
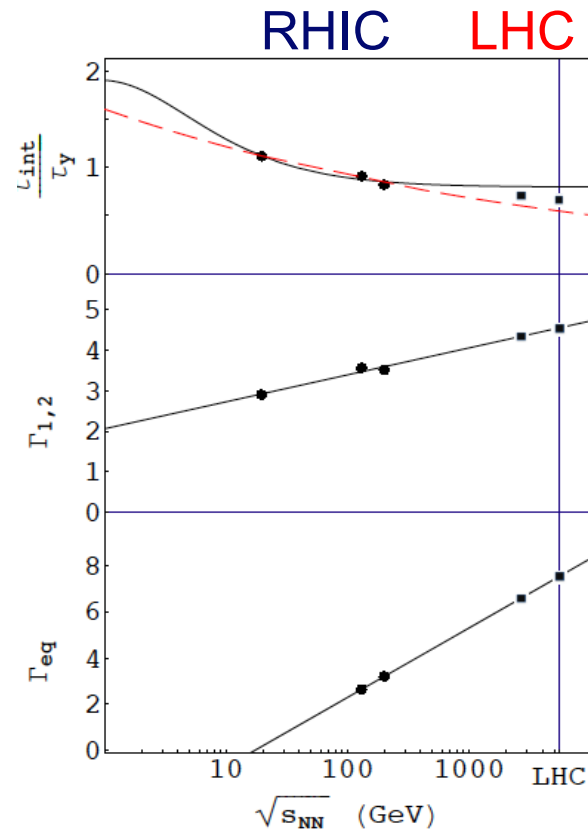


Figure 1: The Jacobian $dy/d\eta$ for $\langle m \rangle = m_\pi$ and average transverse momenta (bottom to top) $\langle p_T \rangle = 0.4, 0.6, 0.8, 1.2, 2$ and $4 \text{ GeV}/c$.
ICFP2C

Energy dependence of the RDM parameters



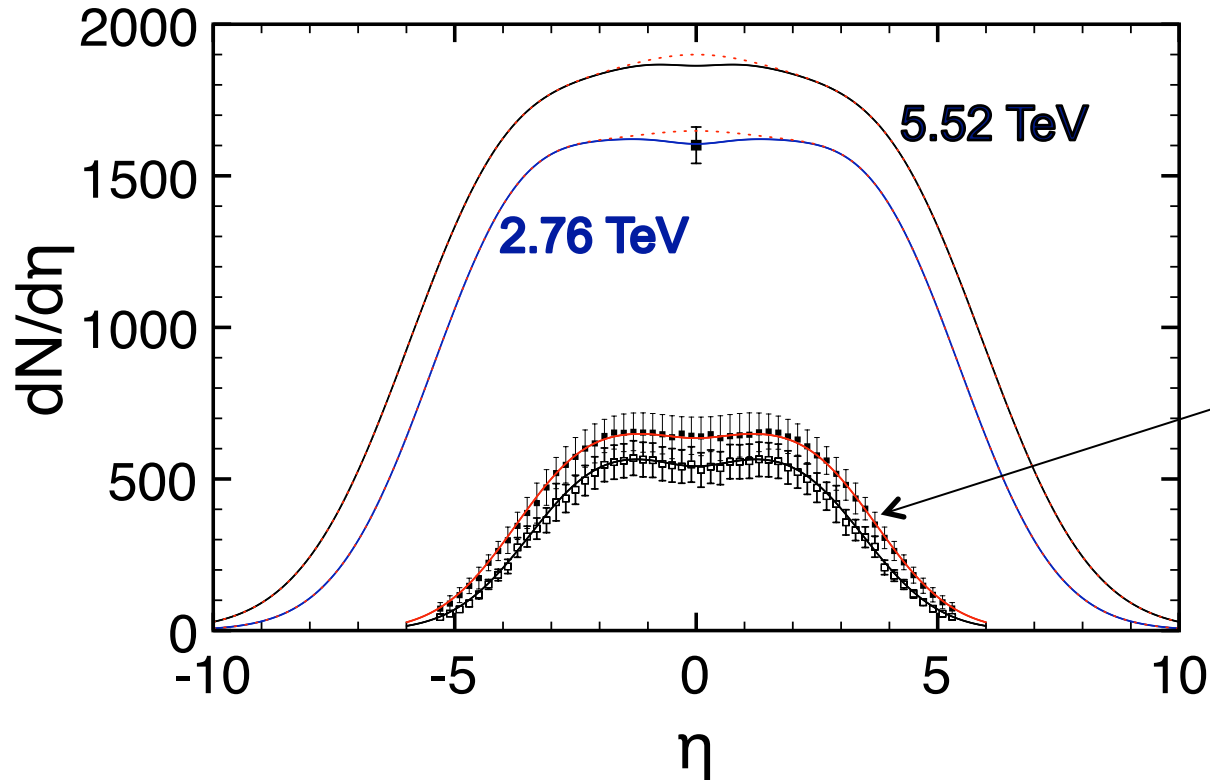
From: R. Kuiper and G.Wolschin,
Annalen Phys. 16, 67 (2006)

Dependence of the diffusion-model parameters for heavy systems at RHIC (dots) and LHC (squares) energies on $\ln \sqrt{s_{NN}}$:

- Quotient of interaction time and relaxation time for sinh- and exponential (dashed) extrapolation (upper frame);
- Effective widths of the peripheral sources, including collective expansion (middle frame);
- Effective width of the midrapidity source (lower frame).

The results are for charged-hadron pseudorapidity distributions.

3. Comparison with RHIC and LHC data



GW, Phys. Lett. B 698, 411 (2011)

$$dN/d\eta (\eta \approx 0) = 1584 \pm 4 \text{ (stat.)} \pm 76 \text{ (sys.)} \text{ [1]}$$

$$1601 \pm 60 \text{ [2]}$$

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LHC: PbPb 0-5%
@ 5.52 TeV and
2.76 TeV central
collisions, RDM-result
adjusted to ALICE [2]
 $dN/d\eta$ at $\eta=0$

RHIC: PHOBOS
AuAu data 0-6%
@ 0.13 and 0.2 TeV [3]
with RDM-results
(χ^2 -minimization)

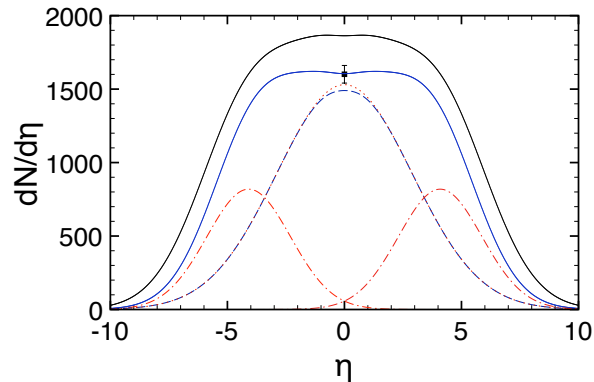
[1] ALICE collab., PRL 105, 252301 (2010)

[2] ALICE collab., PRL 106, 032301 (2011)

[3] B.B. Back et al., PHOBOS coll., PRL 87,
102303 (2001); PRL 91, 052303 (2003);
PRC 83, 024913 (2011)

LHC: Small fragmentation-source contributions at midrapidity

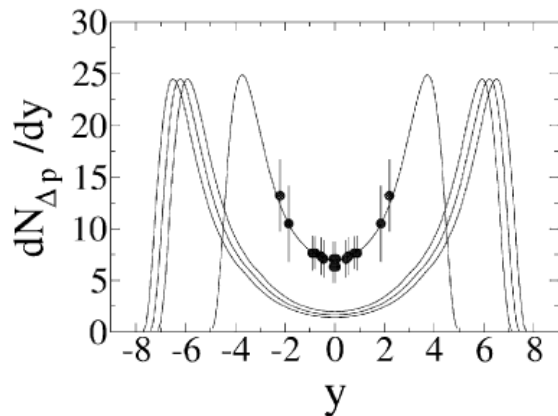
Charged hadrons



PbPb @ 2.76 TeV:

The smallness of the fragmentation sources at midrapidity is in qualitative agreement with results from our QCD-based microscopic model

Y. Mehtar-Tani and GW, Phys. Rev. Lett. 102,182301 (2009);
PRC C80, 054905 (2009)



for net-baryon distributions, which indicates a midrapidity net-baryon yield $dN/dy(y=0) \approx 4$, corresponding to 12 valence quarks, as cp. to 1248 valence quarks in the system (the net-baryon distribution has no gluon-gluon source)

Net protons

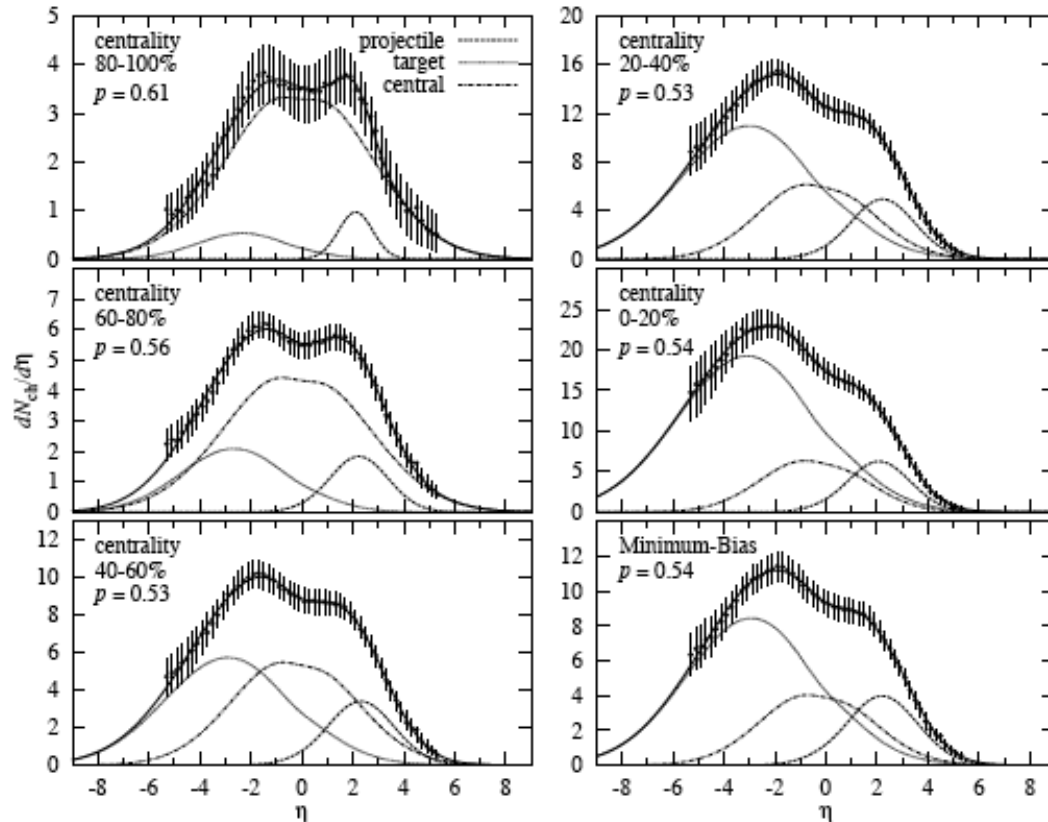
YMT&GW, Phys. Lett. B688, 174 (2010);
GW, Phys. Lett. B 698, 411 (2011)

Parameters of the 3-sources RDM at RHIC and LHC energies

Table 1: Three-sources RDM-parameters for 0–6% Au + Au at RHIC energies (upper two lines) and for 0–5% Pb + Pb at LHC energies (lower two lines). See Fig. 2 and text for the extrapolation of the time parameter τ_{int}/τ_y to LHC energies. Widths and particle numbers denoted by * are extrapolated linearly with $\log(\sqrt{s_{NN}})$. At RHIC energies the nonequilibrium sources from quark-gluon interactions with particle content $N_{ch}^{1,2}$ dominate. At LHC energies the local equilibrium source from gluon-gluon collisions with particle content N_{ch}^{eq} is the major origin of particle production at midrapidity. Experimental midrapidity values (last column) are from PHOBOS [22, 23] for $|\eta| < 1$ at RHIC energies and from ALICE [2] for $|\eta| < 0.5$ at 2.76 TeV.

$\sqrt{s_{NN}}$ (TeV)	y_{beam}	τ_{int}/τ_y	$\langle y_{1,2} \rangle$	$\Gamma_{1,2}$	Γ_{gg}	N_{ch}^{tot}	n_{gg}	$\frac{dN}{d\eta} _{\eta \simeq 0}$
0.13	∓ 4.93	0.89	∓ 2.02	3.43	2.46	4398	0.13	579 ± 23 [17]
0.20	∓ 5.36	0.82	∓ 2.40	3.48	3.28	5315	0.26	655 ± 49 [17]
2.76	∓ 7.99	0.87	∓ 3.34	4.99	6.24	17327	0.56	1601 ± 60 [2]
5.52	∓ 8.68	0.85*	∓ 3.70	5.16*	7.21*	21699*	0.67*	1940*

3-sources model (RDM): Centrality dependence of the asymmetric dAu system @ 0.2 TeV



200 GeV dAu

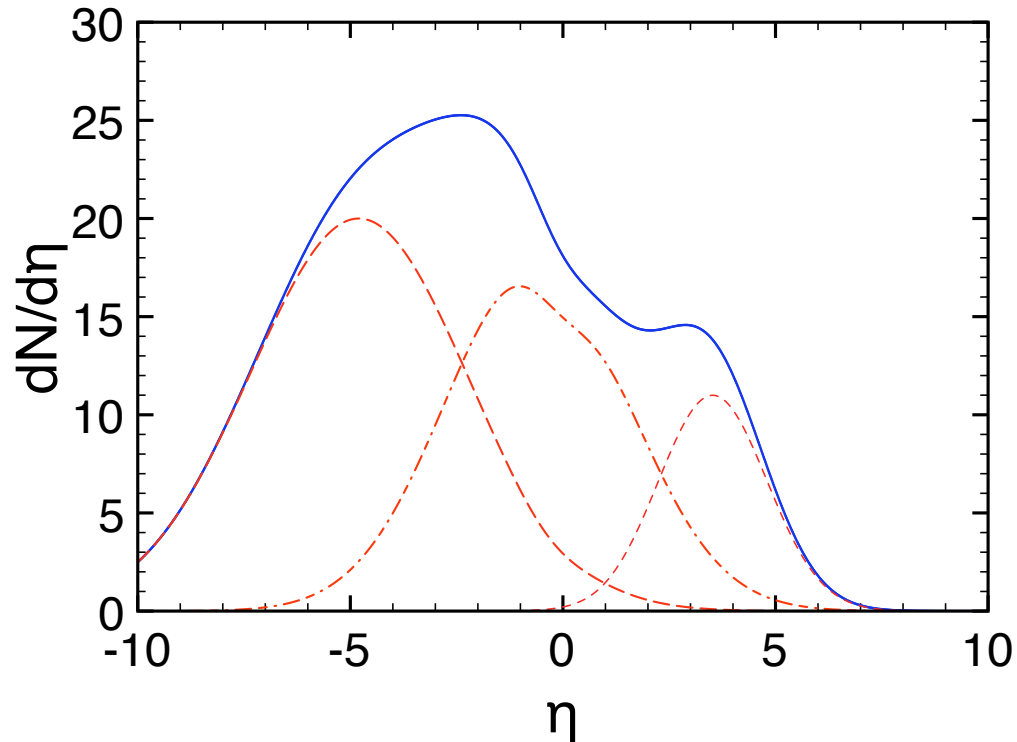
PHOBOS data Phys. Rev. C72, 031901 (2005)

G. Wolschin, M.Biyajima, T.Mizoguchi, N.Suzuki, Annalen Phys. 15, 369 (2006)

Asymmetric systems are more sensitive to details of the nonequilibrium-statistical evolution than symmetric systems

3-sources model (RDM): Preliminary calc. for central pPb @ 5.02 TeV

0-5% central 5.02 TeV pPb @ LHC



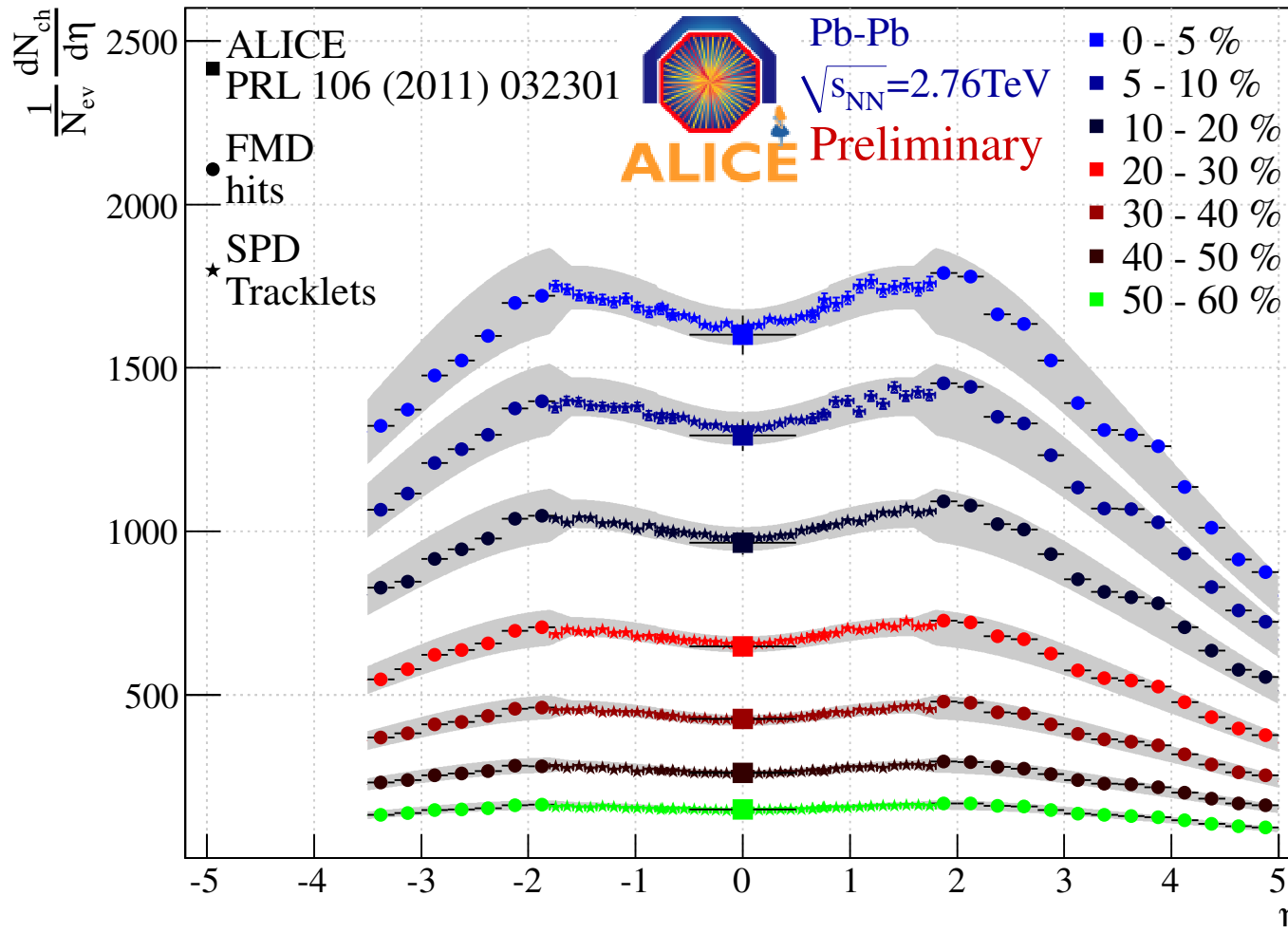
$$p_p = 4 \text{ TeV}/c$$

$$\sqrt{s_{NN}} = \sqrt{\frac{Z_1 * Z_2}{(A_1 * A_2)}} * 2p_p = 5.02 \text{ TeV}$$

$$y_{\text{beam}}^{cm} = \mp \ln(\sqrt{s_{NN}}/m_0) = \mp 8.586$$

(RDM parameters to be adjusted)

ALICE prel. results for pseudorapidity distributions of produced charged hadrons



H.H. Dalsgaard et al.,
ALICE Collab.;
cf talk **June 12**

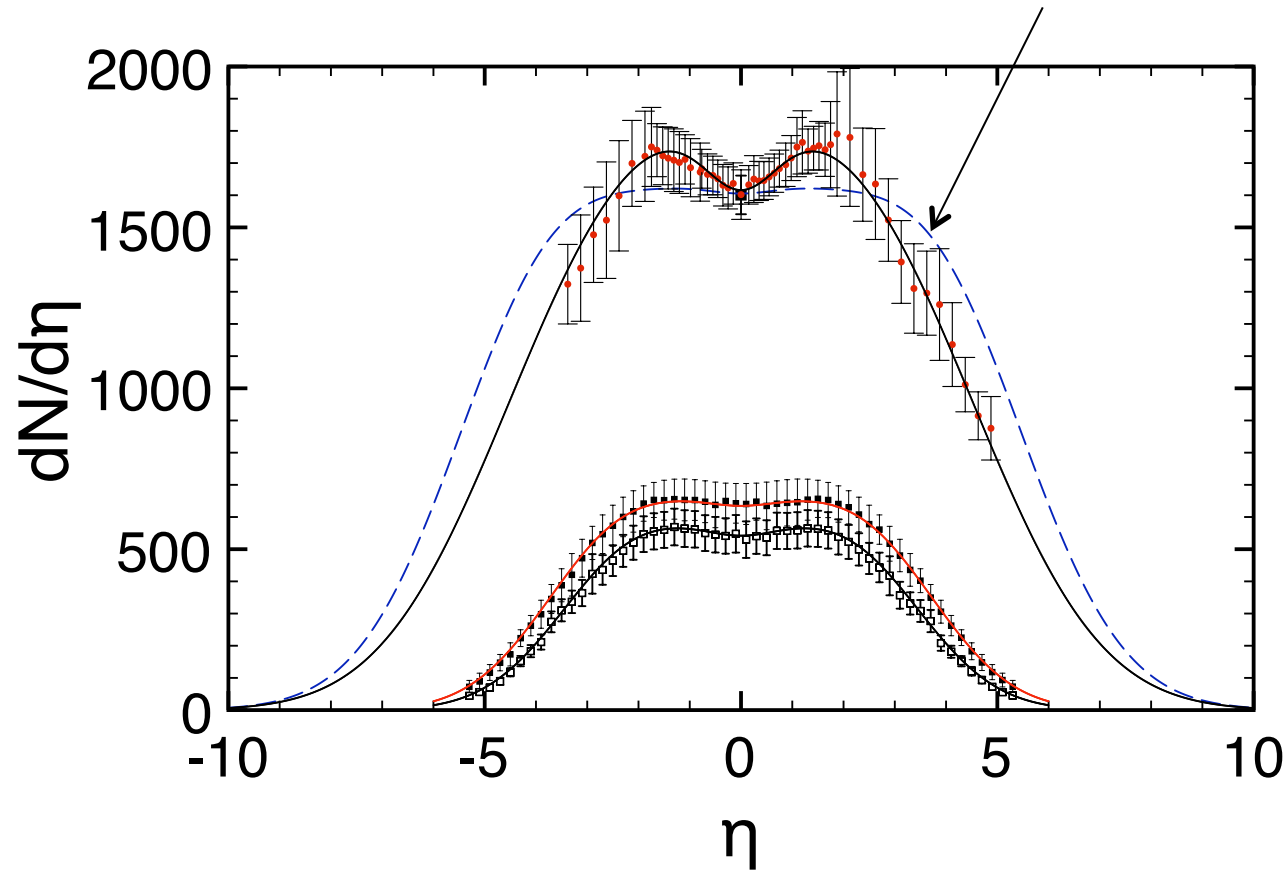
PRELIMINARY

- Central dip is more pronounced than in the prediction
- Peak at smaller eta values than expected

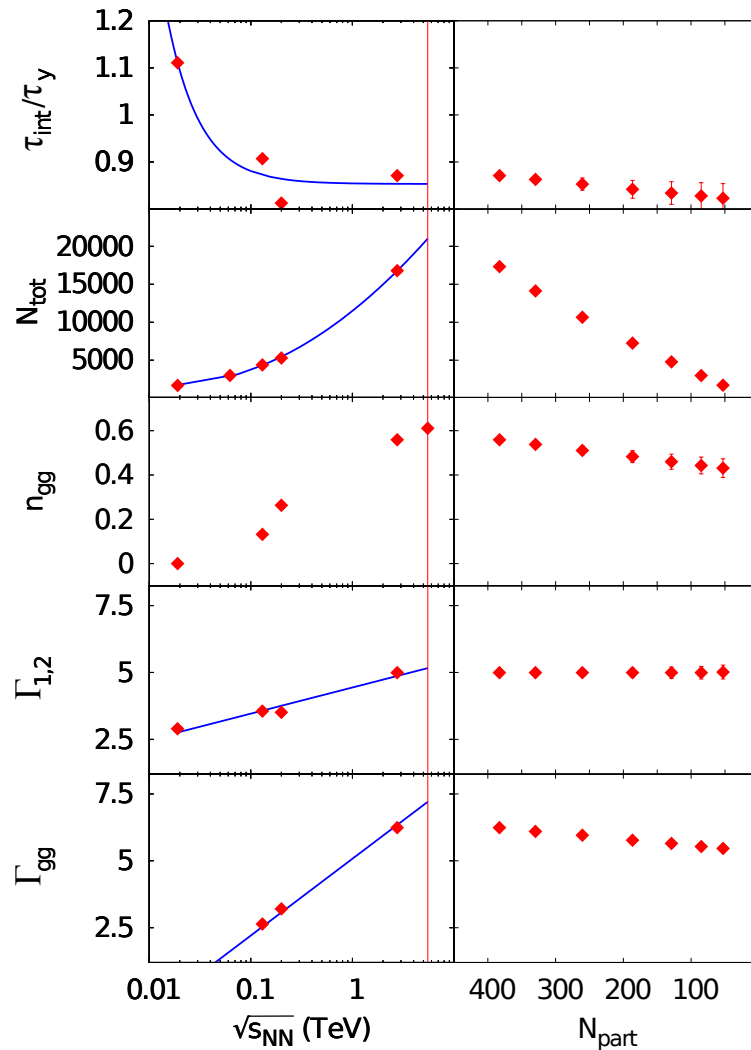
Comparison with the RDM prediction

Central PbPb @ 2.76 TeV

Prediction GW in PLB 698, 411 (2011)



Consequences for the RDM parameters

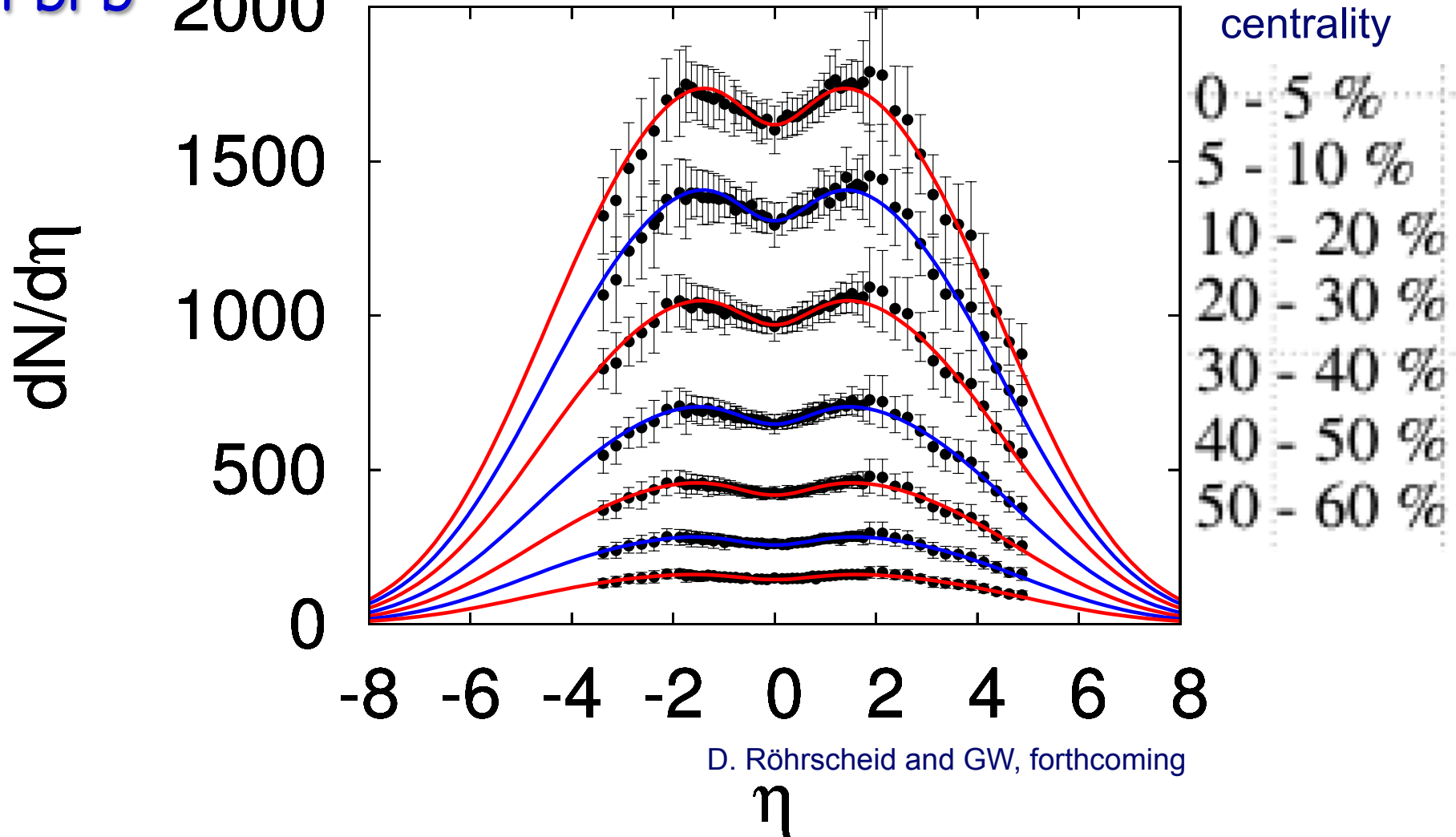


- The time parameter τ_{int}/τ_y is slightly larger than expected. It **saturates** at LHC energies because the rapidity relaxation time τ_y decreases (different from the original extrapolation!)
- The relative particle content in the gluon-gluon source is smaller than expected and hence, the midrapidity dip is more pronounced.

D. Röhrscheid and GW, preprint

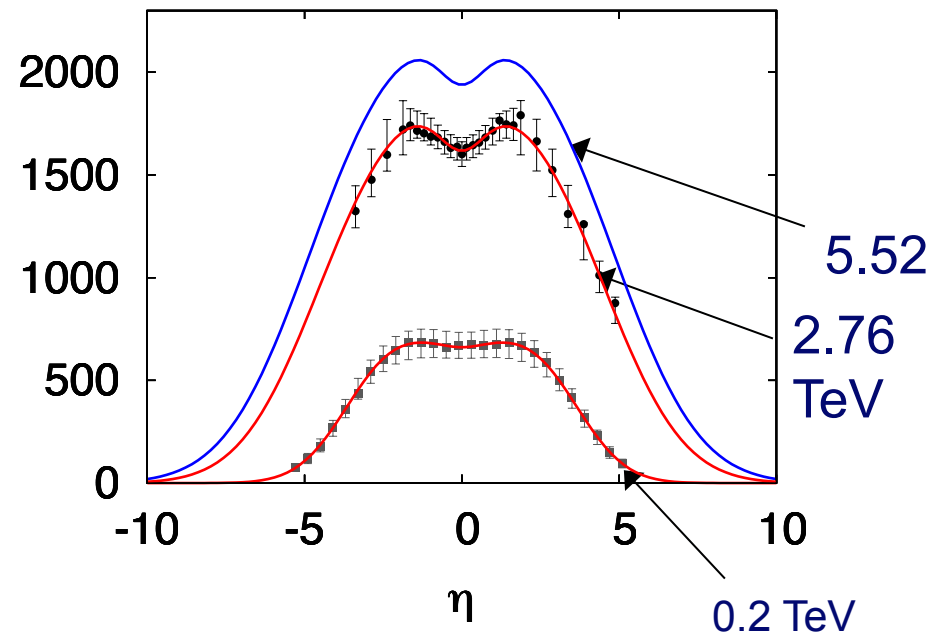
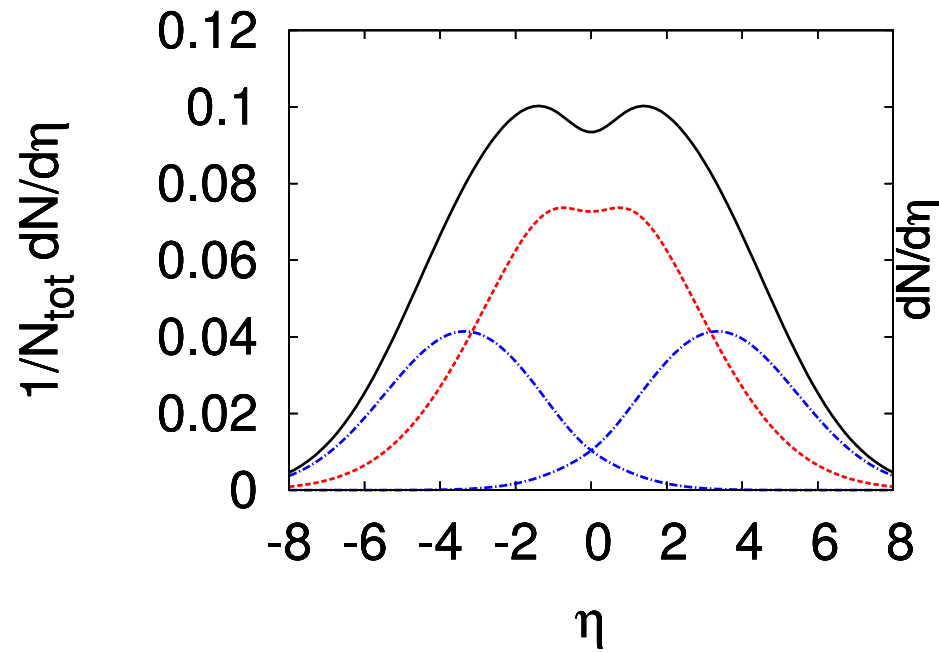
RDM χ^2 fits to LHC/ALICE results for 2.76 TeV

PbPb

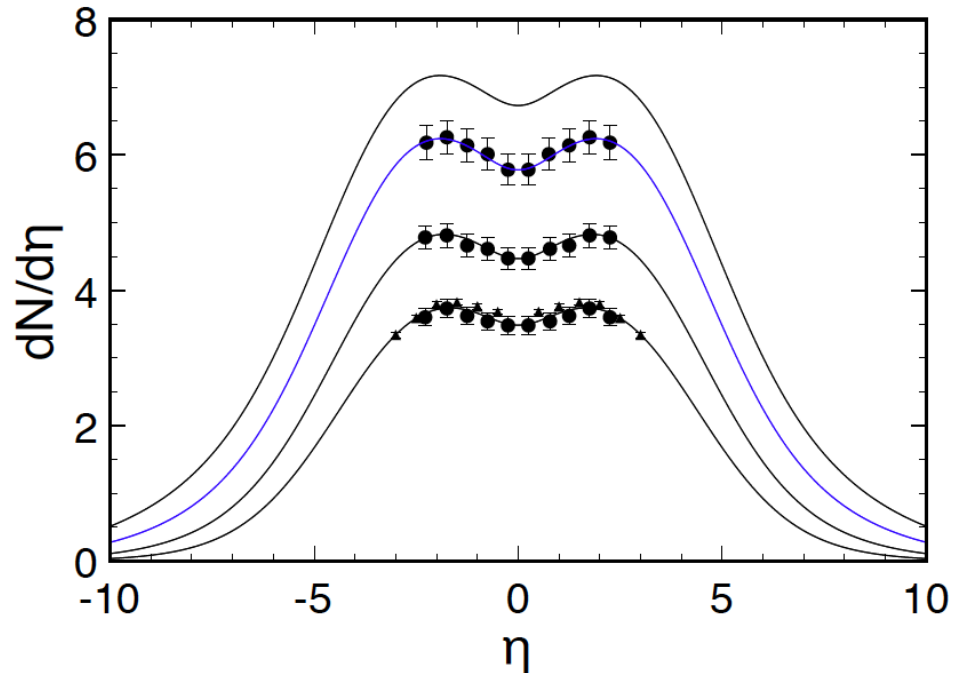


3 sources, and prediction for 5.52 TeV PbPb

Centrality 0-5%



Charged-hadron distributions in pp: 3-sources relativistic diffusion model (RDM)



14 TeV pp

7 TeV

2.36 TeV

0.9 TeV (LHC injection energy / UA5)

Data: CMS collab., V. Khachatryan et al.,
J. High Energy Phys. 02, 041 (2010);
Phys. Rev. Lett. 105, 022002 (2010);
UA5 collab., R. Ansorge et al., Z. Phys. C 83,
357 (1989)

pp charged-hadron
pseudorapidity
data beyond $\eta=2.5$
needed

3-sources RDM calculation: see GW, EPL 95, 61001 (2011) [Europhys. Lett.]

4. Conclusion

- ❖ Charged-hadron production at RHIC and LHC energies has been described in a Relativistic Diffusion Model (RDM).
- ❖ Predictions of pseudorapidity distributions $dN/d\eta$ of produced charged hadrons for various centralities in the 3-sources RDM at LHC energies rely on the extrapolation of the diffusion-model parameters with $\ln(\sqrt{s_{NN}})$
- ❖ In agreement with a QCD-based microscopic model, the contribution of the fragmentation sources from quark-gluon collisions at LHC energies is very small at midrapidity, but substantial at larger values of pseudorapidity η .
- ❖ The centrality dependence of the three sources has been investigated in direct comparison with the preliminary ALICE data.